

Baryogenesis via relativistic bubble expansion

Simone Blasi
Vrije Universiteit Brussel (VUB)

Based on:

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

"How fast does the bubble grow?" 15-17 May 2023, DESY Hamburg

Introduction

Baryon asymmetry of the Universe:

$$Y_B = \frac{n_b - n_{\bar{b}}}{s} \sim 10^{-10}$$

Abundance set by $B - \bar{B}$ annihilation:

$$\frac{n_b}{s} \left(= \frac{n_{\bar{b}}}{s} \right) \sim 10^{-20}$$

Solution: some **initial quark-anti quark asymmetry** prevents strong annihilation, and explains absence of antimatter in our Universe.

Dynamical generation of the asymmetry: **baryogenesis**, a challenge for particle physics!

Introduction

Sakharov conditions

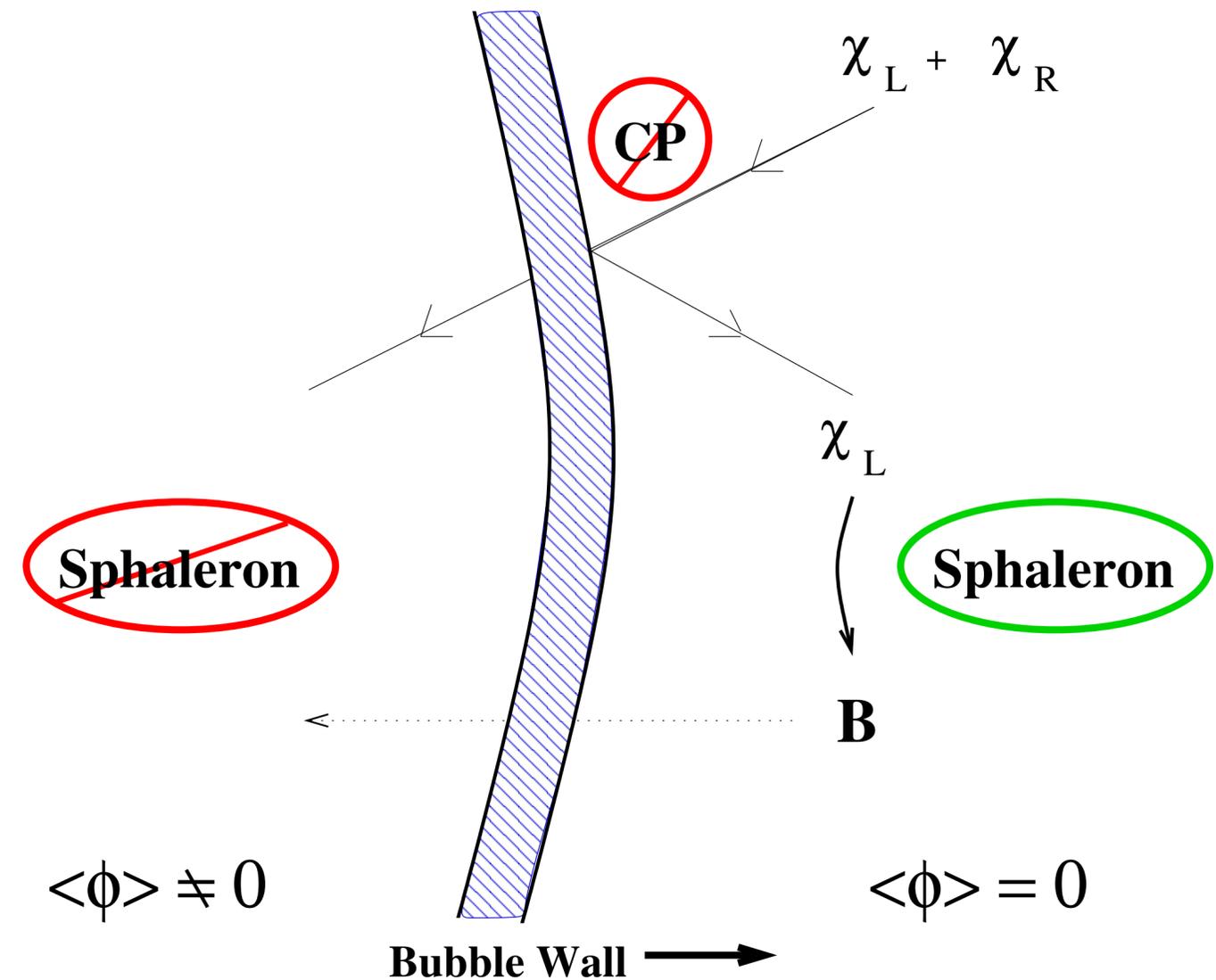
- Baryon number violation
- C and CP violation
- Departure from thermal equilibrium

Introduction

Electroweak baryogenesis

- SM sphalerons violate B+L
- Beyond the SM violation of CP
- First order electroweak phase transition

Fig. from Morrissey, Ramsey-Musolf
[1206.2942]

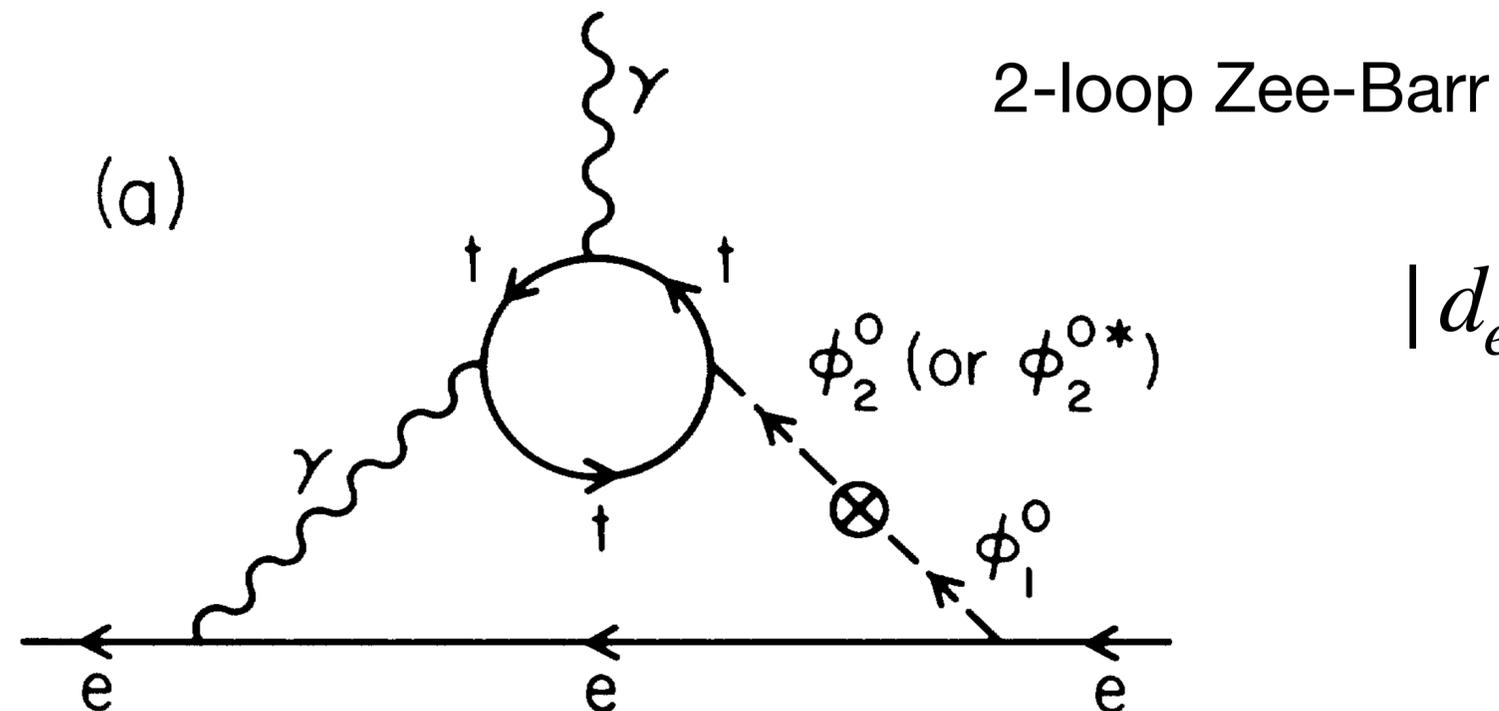


Introduction

Successful electroweak baryogenesis

- In tension with electric dipole moment measurement

See also: [Postma, van de Vis, White \[2206.01120\] JHEP](#)



$$|d_e| \sim 10^{-29} \theta_{CP} \left(\frac{2.5 \text{ TeV}}{\Lambda} \right)^2 e \cdot \text{cm}$$

[Glioti, Rattazzi, Vecchi \[1811.11740\] JHEP](#)

A way out: CP only *temporarily* broken during the phase transition

Introduction

Successful electroweak baryogenesis

- BAU suppressed for $v_w \rightarrow 0$ and $v_w \rightarrow 1$ (but possible for $v_w > c_s$)

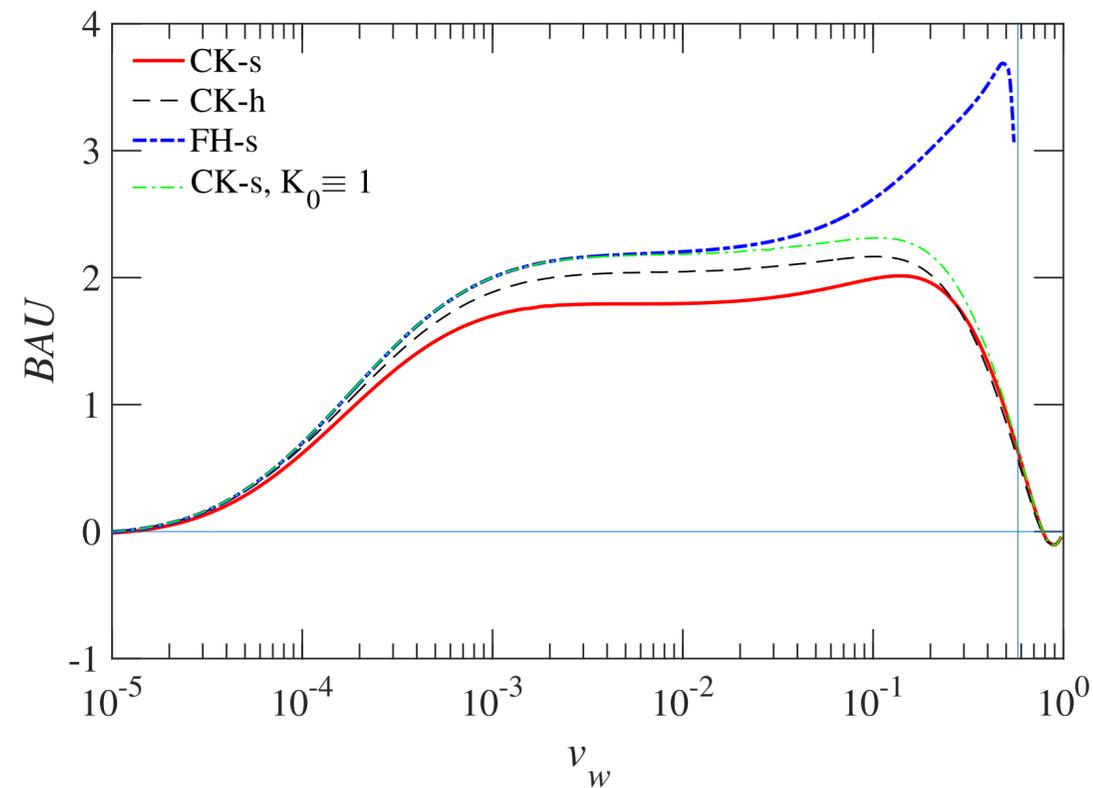


Fig. from Cline, Kainulainen [2001.00568] PRD
See also Dorsch, Huber, Konstandin [2106.06547]

- BAU and observable gravitational waves?

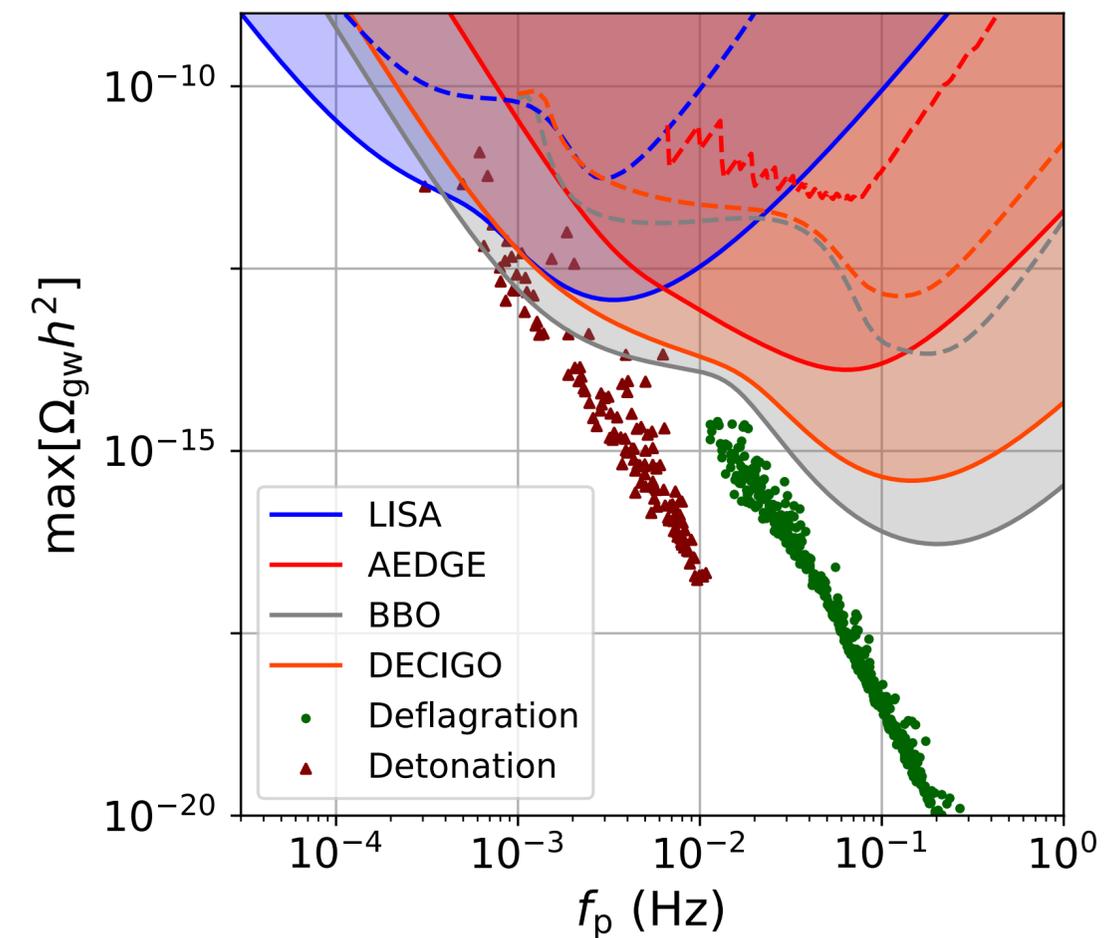


Fig. from Laurent, Cline et al. [2102.12490] PRD

- ➡ Can we have baryogenesis with (ultra) relativistic bubble walls?
- ➡ Can this mechanism lead to a large GW signal?

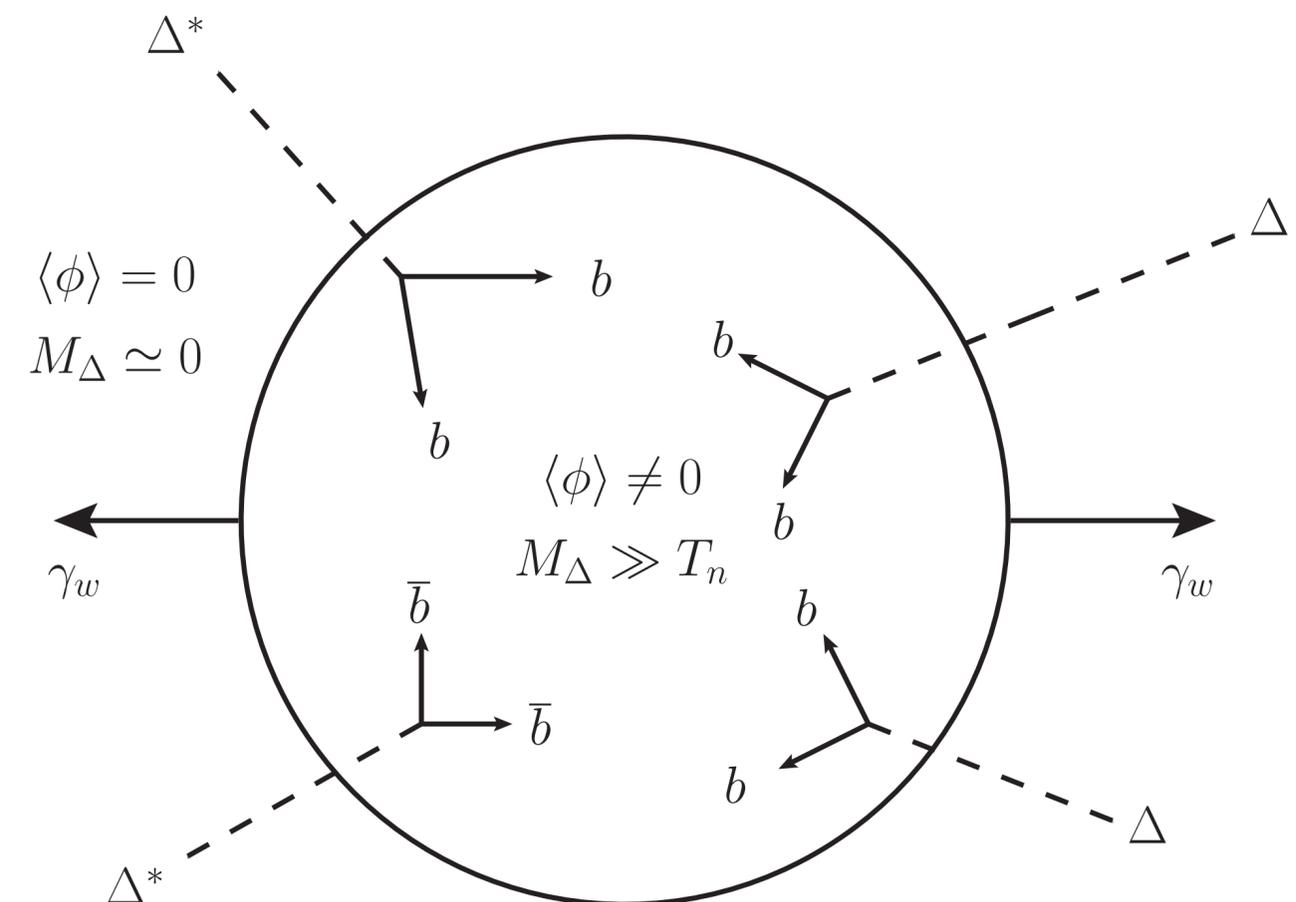
Baryogenesis with relativistic walls

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

Mass gain mechanism

- A particle Δ in thermal equilibrium gains a large mass M_Δ across the wall

$$\mathcal{L} \supset -\frac{\lambda}{2}\phi^2|\Delta|^2$$



Δ particles crossing the ϕ wall and gaining a mass undergo asymmetric decays.

Baryogenesis with relativistic walls

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

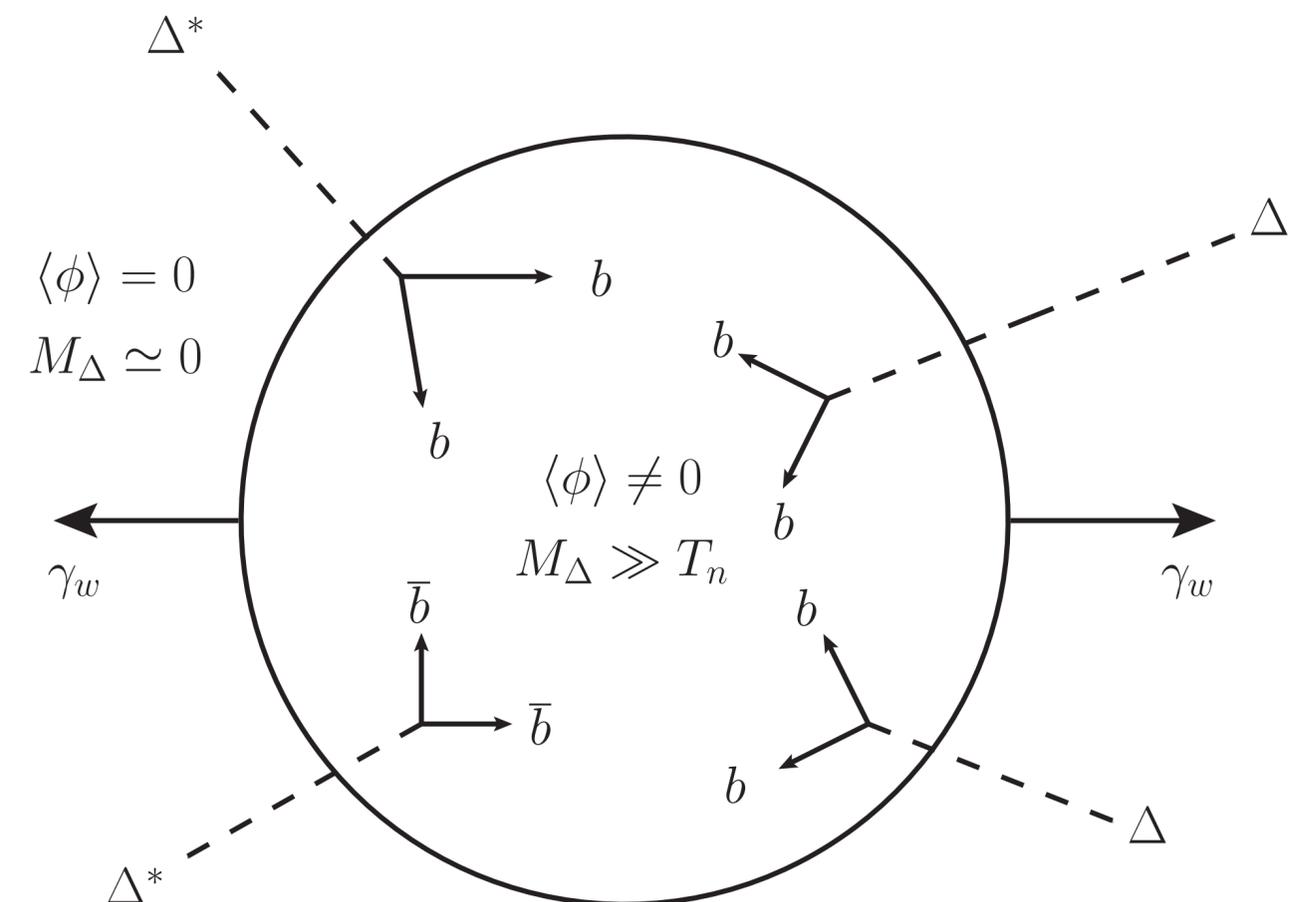
Mass gain mechanism

- A particle Δ in thermal equilibrium gains a large mass M_Δ across the wall

$$\mathcal{L} \supset -\frac{\lambda}{2}\phi^2|\Delta|^2$$

- Kinematics requires fast enough walls

$$\gamma_w > \frac{M_\Delta}{T_n}$$



Δ particles crossing the ϕ wall and gaining a mass undergo asymmetric decays.

Baryogenesis with relativistic walls

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

Mass gain mechanism

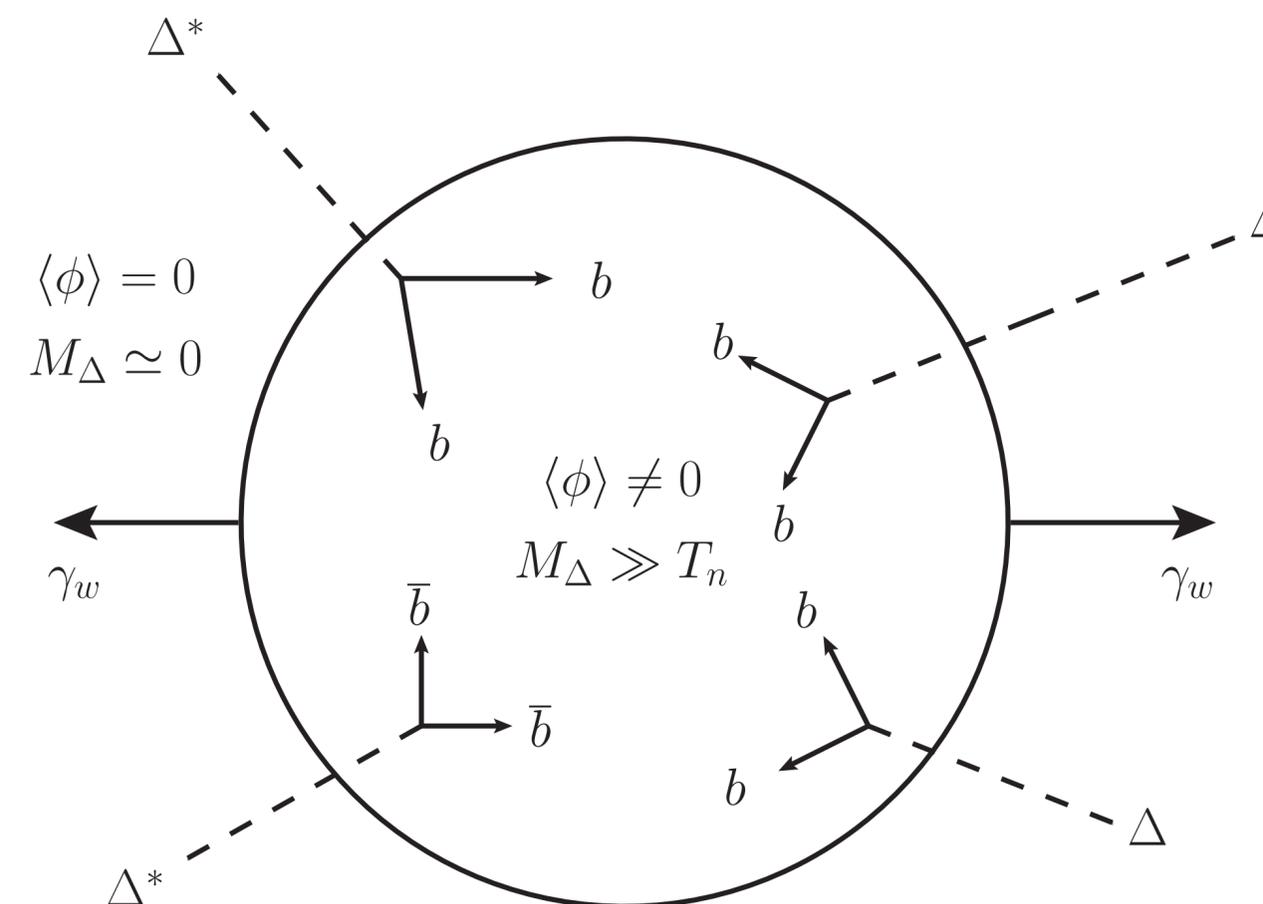
- A particle Δ in thermal equilibrium gains a large mass M_Δ across the wall

$$\mathcal{L} \supset -\frac{\lambda}{2}\phi^2|\Delta|^2$$

- Kinematics requires fast enough walls

$$\gamma_w > \frac{M_\Delta}{T_n}$$

- Δ inside the bubble is out of equilibrium, and undergoes CP and B violating decays



Δ particles crossing the ϕ wall and gaining a mass undergo asymmetric decays.

Baryogenesis with relativistic walls

Mass gain mechanism

$$Y_B = \epsilon_{\Delta} \cdot \kappa_{Sph.} \left(\frac{n_{\Delta}(M_{\Delta} \simeq 0)}{s_e} \right) \Big|_{eq} \left(\frac{T_n}{T_{RH}} \right)^3$$

Fraction of asymmetric decays (1-loop effect) $\rightarrow \epsilon_{\Delta}$
 Sphaleron reprocessing factor = 28/79 $\rightarrow \kappa_{Sph.}$
 Large equilibrium abundance $\rightarrow \left(\frac{n_{\Delta}(M_{\Delta} \simeq 0)}{s_e} \right)$
 Dilution from reheating (if strong supercooling) $\rightarrow \left(\frac{T_n}{T_{RH}} \right)^3$

Baryogenesis with relativistic walls

Mass gain mechanism

$$Y_B \sim \left(\frac{100}{g^*} \right) \left(\frac{\epsilon_\Delta}{1/16\pi} \right) \left(\frac{T_n}{T_{RH}} \right)^3 10^5 Y_{B \text{ obs.}}$$

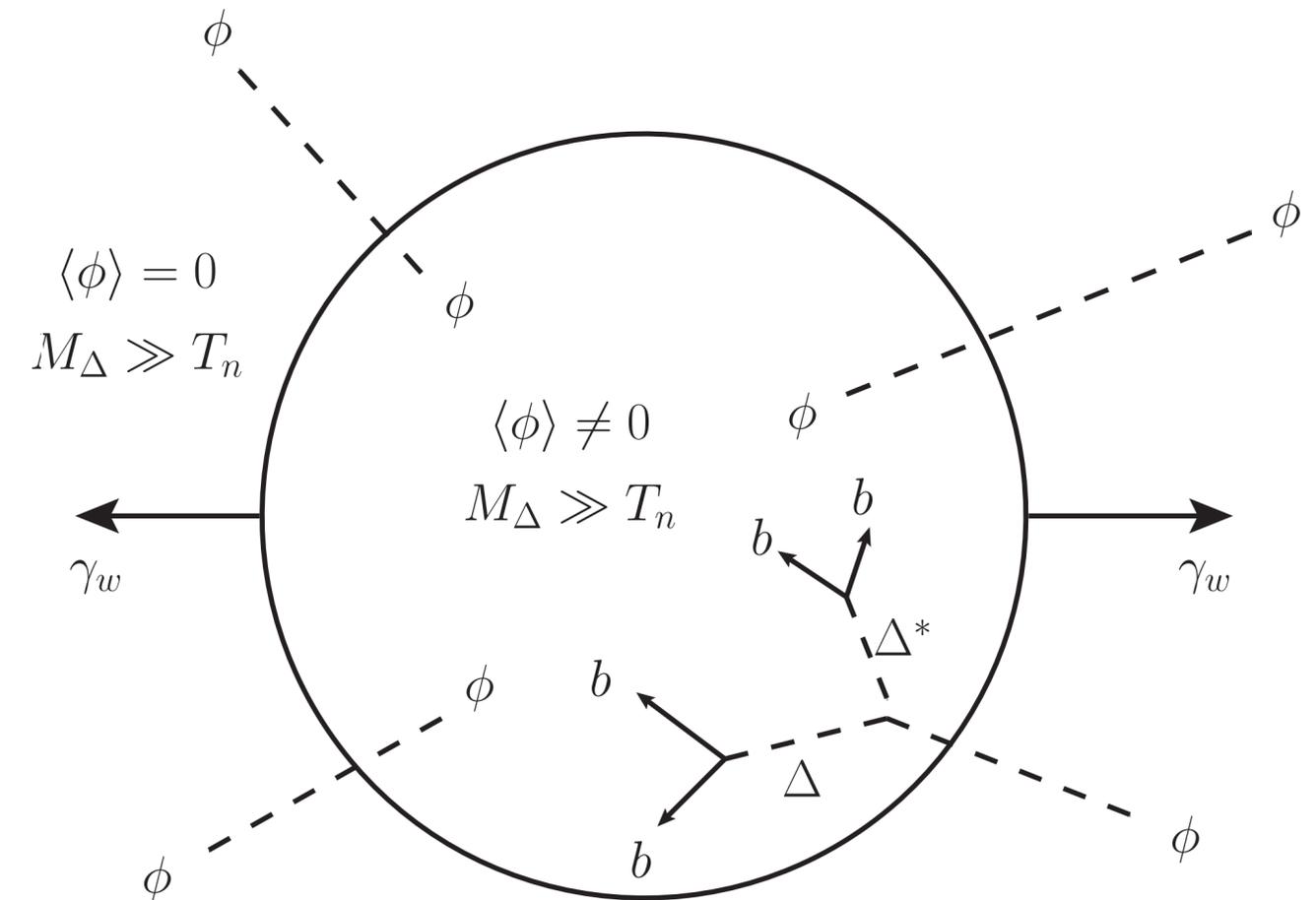
Baryogenesis with relativistic walls

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

Azatov-Vanvlasselaer mechanism

- Heavy Δ particles are possibly produced from ϕ interacting with the wall

$$\mathcal{L} \supset -\frac{\lambda}{2}\phi^2|\Delta|^2 - M_\Delta^2|\Delta|^2$$



Δ particles produced at the ϕ wall undergo asymmetric decays.

Baryogenesis with relativistic walls

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

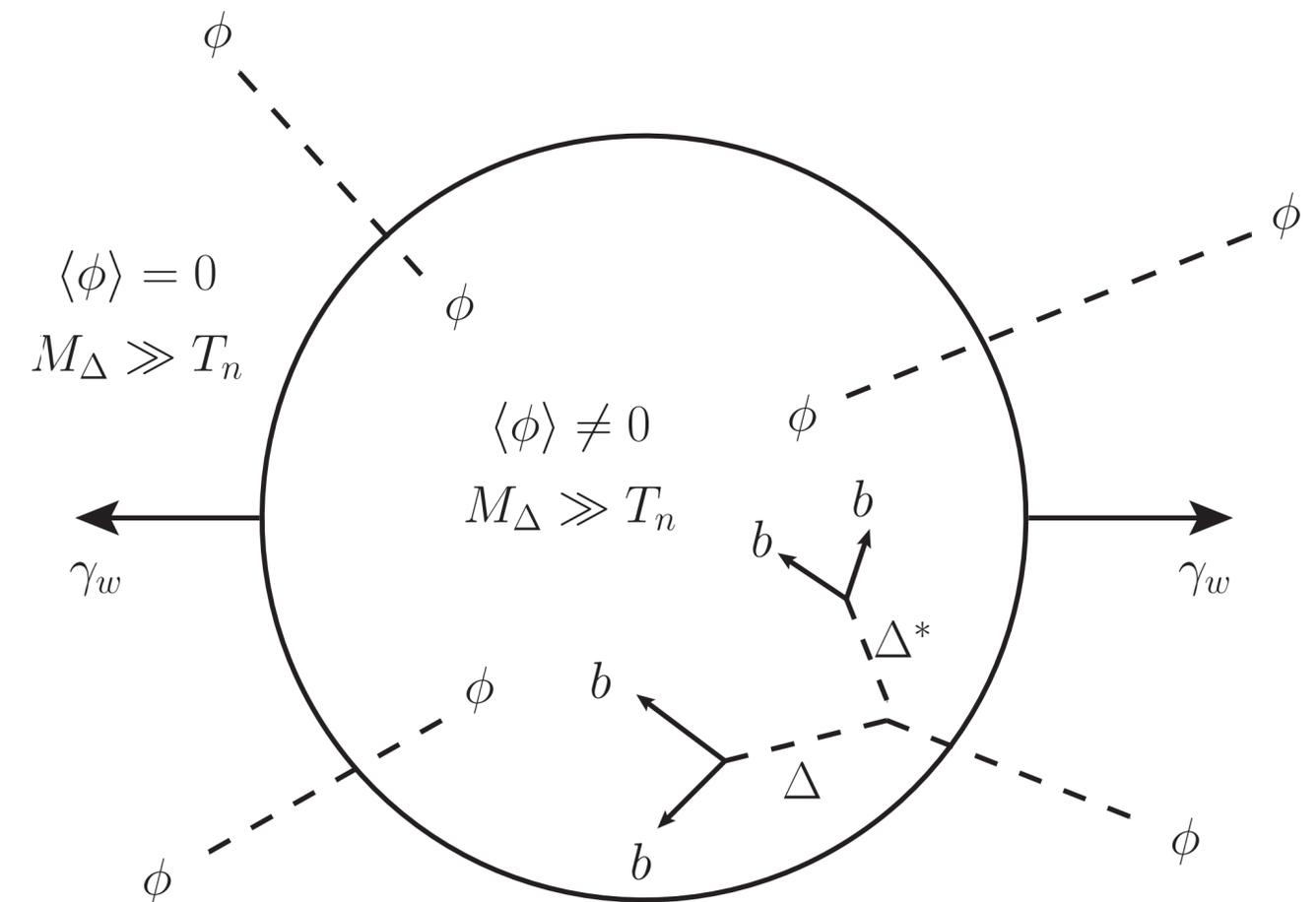
Azatov-Vanvlasselaer mechanism

- Heavy Δ particles are possibly produced from ϕ interacting with the wall

$$\mathcal{L} \supset -\frac{\lambda}{2}\phi^2|\Delta|^2 - M_\Delta^2|\Delta|^2$$

- Fast enough walls for the anti-adiabatic regime

$$\gamma_w > (M_\Delta \delta_w) \cdot \frac{M_\Delta}{T_n}$$



Δ particles produced at the ϕ wall undergo asymmetric decays.

Baryogenesis with relativistic walls

Baldes, **SB**, Mariotti, Sevrin, Turbang
[2106.15602] PRD

Azatov-Vanvlasselaer mechanism

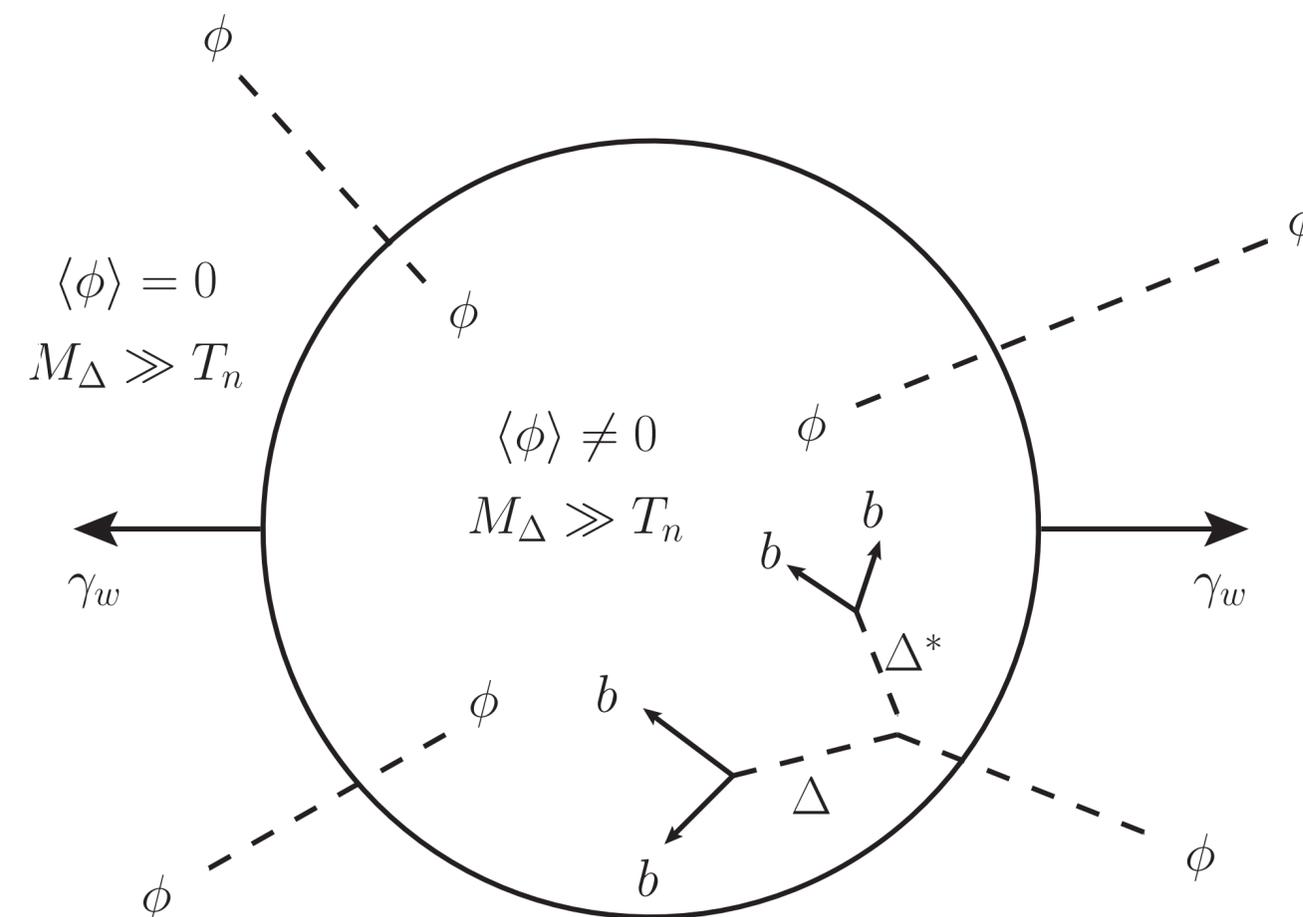
- Heavy Δ particles are possibly produced from ϕ interacting with the wall

$$\mathcal{L} \supset -\frac{\lambda}{2}\phi^2|\Delta|^2 - M_\Delta^2|\Delta|^2$$

- Fast enough walls for the anti-adiabatic regime

$$\gamma_w > (M_\Delta \delta_w) \cdot \frac{M_\Delta}{T_n}$$

- Δ decays as the “mass gain” mechanism



Δ particles produced at the ϕ wall undergo asymmetric decays.

Baryogenesis with relativistic walls

Azatov-Vanvlasselaer mechanism

$$Y_B = \epsilon_\Delta \cdot \kappa_{Sph.} \cdot \mathcal{P}(\phi \rightarrow \Delta\Delta^*) \left(\frac{n_\phi}{s_e} \right) \Big|_{eq} \left(\frac{T_n}{T_{RH}} \right)^3$$

Azatov, Vanvlasselaer
[2010.02590] JCAP

Azatov, Vanvlasselaer, Yin
[2101.05721] JHEP

Prob. of 1 to 2 transition:

$$\mathcal{P} \sim \frac{\lambda^2 v_\phi^2}{96\pi^2 M_\Delta^2}$$

Large equilibrium
abundance

Baryogenesis with relativistic walls

Azatov-Vanvlasselaer mechanism

$$Y_B \sim \left(\frac{100}{g^*} \right) \left(\frac{\epsilon_\Delta}{1/16\pi} \right) \left(\frac{T_n}{T_{RH}} \right)^3 \lambda^2 \left(\frac{v_\phi}{M_\Delta} \right)^2 10^2 Y_{B \text{ obs.}}$$

Related work

- [Azatov, Vanvlasselaer, Yin \[2106.14913\] JHEP](#) (see also [\[2207.02230\]](#)): AV mechanism with CP violation at the wall

Related work

- [Azatov, Vanvlasselaer, Yin \[2106.14913\] JHEP](#) (see also [\[2207.02230\]](#)): AV mechanism with CP violation at the wall
- [Huang, Xie \[2206.04691\] JHEP](#); [Dasgupta, Dev, Ghoshal, Mazumdar \[2206.07032\] PRD](#): high and low scale leptogenesis-style realization

Related work

- [Azatov, Vanvlasselaer, Yin \[2106.14913\] JHEP](#) (see also [\[2207.02230\]](#)): AV mechanism with CP violation at the wall
- [Huang, Xie \[2206.04691\] JHEP](#); [Dasgupta, Dev, Ghoshal, Mazumdar \[2206.07032\] PRD](#): high and low scale leptogenesis-style realization
- [Shuve, Tamarit \[1704.01979\] JHEP](#): N_R gains a mass during the PT and decays, but option of ultra relativistic walls not considered

Related work

- [Azatov, Vanvlasselaer, Yin \[2106.14913\] JHEP](#) (see also [\[2207.02230\]](#)): AV mechanism with CP violation at the wall
- [Huang, Xie \[2206.04691\] JHEP](#); [Dasgupta, Dev, Ghoshal, Mazumdar \[2206.07032\] PRD](#): high and low scale leptogenesis-style realization
- [Shuve, Tamarit \[1704.01979\] JHEP](#): N_R gains a mass during the PT and decays, but option of ultra relativistic walls not considered
- [Falkowski, No \[1211.5615\] JHEP](#); [Katz, Riotto \[1608.00583\] JCAP](#): production of heavy particles as a result of elastic wall collisions, leading to DM, BAU and GWs

Bubble wall dynamics

- Vacuum pressure:

$$\mathcal{P}_{dr} \sim \Delta V = c_{vac} \cdot v_\phi^4$$

- Runaway bubbles: $\Delta V > \mathcal{P}_{LO}$, $\mathcal{P}_{NLO} = 0$

$$\gamma_* \sim \frac{R_*}{R_c} \sim \frac{1}{\beta} \frac{T_n}{H_*} \gg 1$$

- Pressure from the plasma:

$$\mathcal{P}_{fr} = \mathcal{P}_{LO} + \gamma \mathcal{P}_{NLO}$$

1 → 1 transmission of particles
changing mass across the wall

1 → 2 scatterings
involving gauge bosons

$$\mathcal{P}_{LO} \sim g_a \frac{v_\phi^2 T^2}{24} + \mathcal{O}(\gamma^{-1})$$

Bodeker, Moore [0903.4099] JCAP Bodeker, Moore [1703.08215] JCAP

See also Hoche, Kozaczuk, Long, Turner, Wang [2007.10343] JCAP

Summary of kinematic conditions

- **Runaway bubbles:**

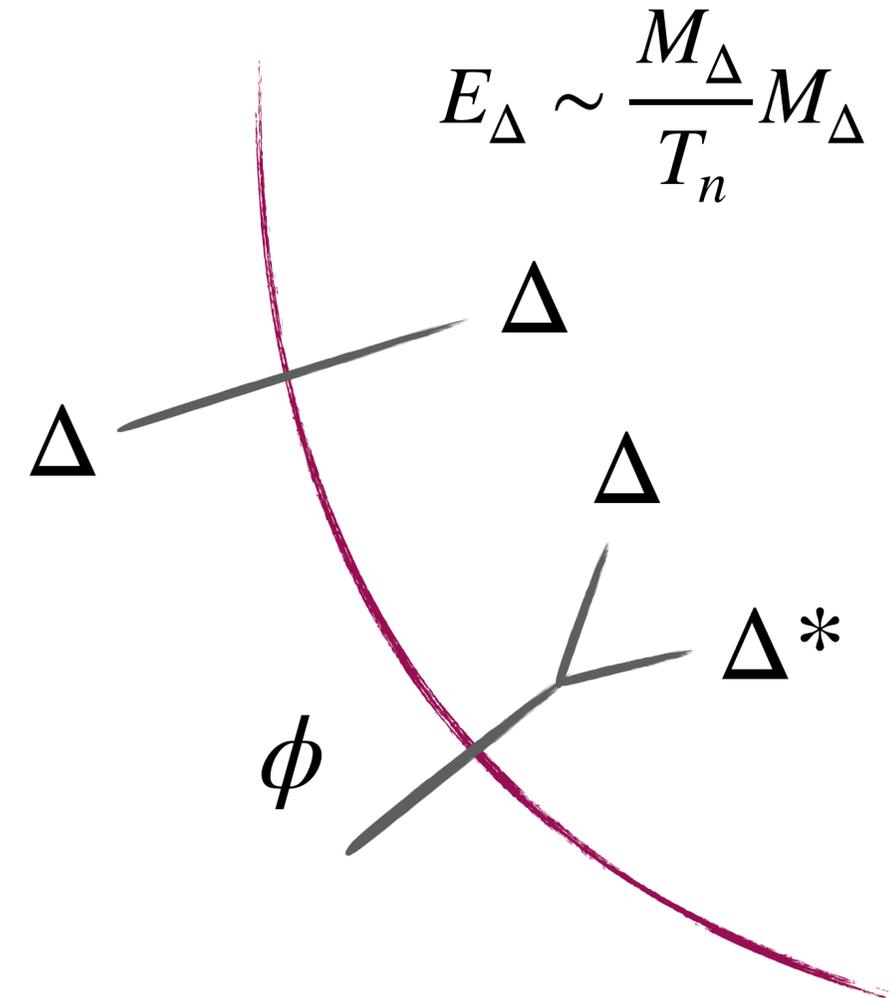
Moderate supercooling

$$\frac{T_n}{v_\phi} < 0.5 \left(\frac{c_{vac}}{0.1} \cdot \frac{10}{g_a} \right)^{1/2}$$

- **Efficient Δ crossing or production:**

$$\frac{M_\Delta}{T_n} < \left(\frac{1}{\beta} \frac{T_n}{H_*} \right)^a, \quad a = \begin{cases} 1 & MG \\ \frac{1}{2} & AV \end{cases}$$

Anti-adiabatic condition
(taking $v_\phi \sim T_n$)



A model for relativistic baryogenesis

- Two “flavors” $\Delta_{1,2}$ with QCD and hypercharge interactions
- B violation due to **simultaneous diquark** and **leptoquark** interactions

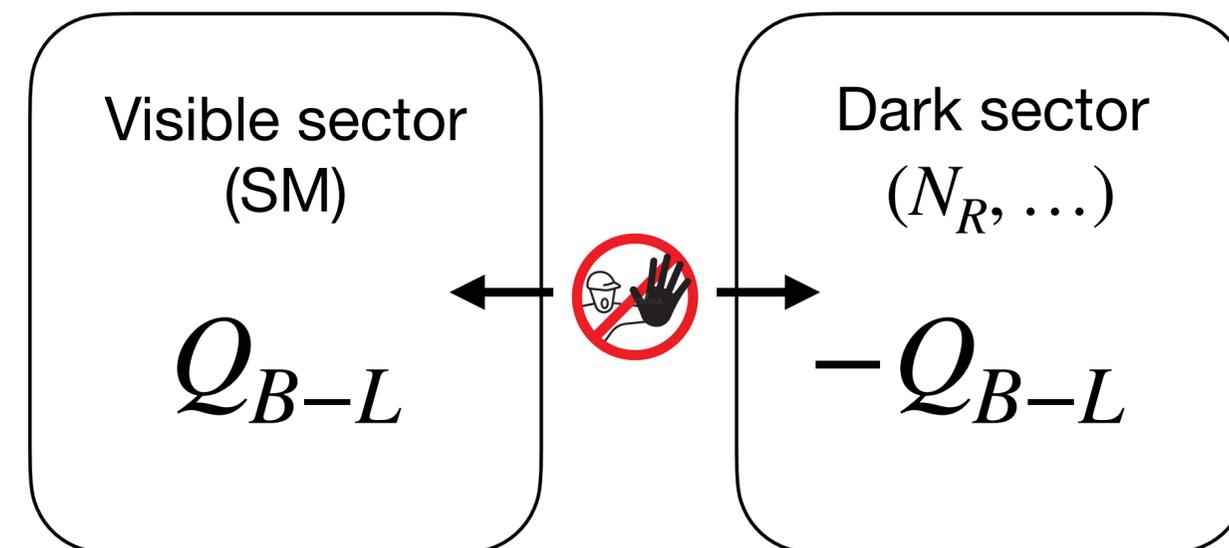
$$\mathcal{L}_{B \text{ viol.}} =$$

$$+ \Delta_i \begin{array}{c} \nearrow \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \searrow \end{array} \begin{array}{c} \bar{d}'_R \\ \\ d_R^c \end{array} + \Delta_i^* \begin{array}{c} \nearrow \\ \text{---} \text{---} \text{---} \\ \text{---} \text{---} \text{---} \\ \searrow \end{array} \begin{array}{c} u_R^c \\ \\ \bar{N}_R \end{array}$$

A model for relativistic baryogenesis

- Two “flavors” $\Delta_{1,2}$ with QCD and hypercharge interactions
 - B violation due to **simultaneous diquark and leptoquark** interactions
 - **B-L cons. Δ decays**: charge carried by N_R needs to be “hidden” until T_{EW}
- (Explicit B-L interactions may be included as well)

$$\mathcal{L}_{B \text{ viol.}} = \Delta_i \begin{array}{c} \nearrow \\ \bullet \text{ } y_i^d \\ \searrow \\ d_R^c \end{array} \bar{d}'_R + \Delta_i^* \begin{array}{c} \nearrow \\ \bullet \text{ } y_i^u \\ \searrow \\ \bar{N}_R \end{array} u_R^c$$



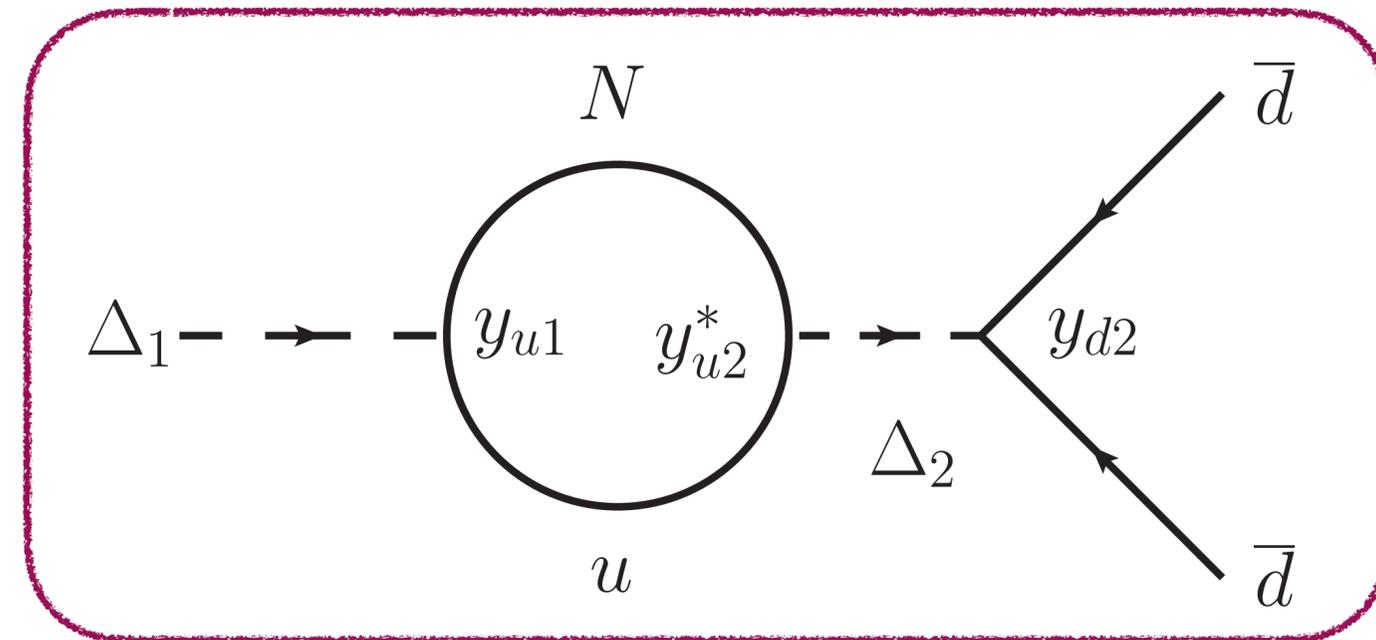
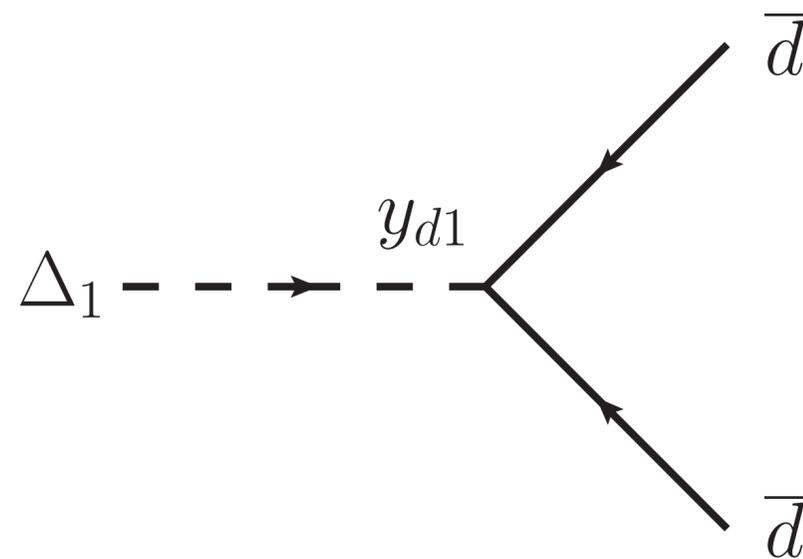
Generating the asymmetry

Right after wall crossing:

- CP violating Δ **decays** generate a B-L asymmetry in the visible sector

$$\epsilon_{\Delta} \sim \frac{\text{Im}[y^2]}{6\pi} \left(\frac{M_{\Delta_1}}{M_{\Delta_2}} \right)^2$$

$$\Gamma_{\Delta} \sim \frac{3y^2}{16\pi} M_{\Delta}$$



Generating the asymmetry

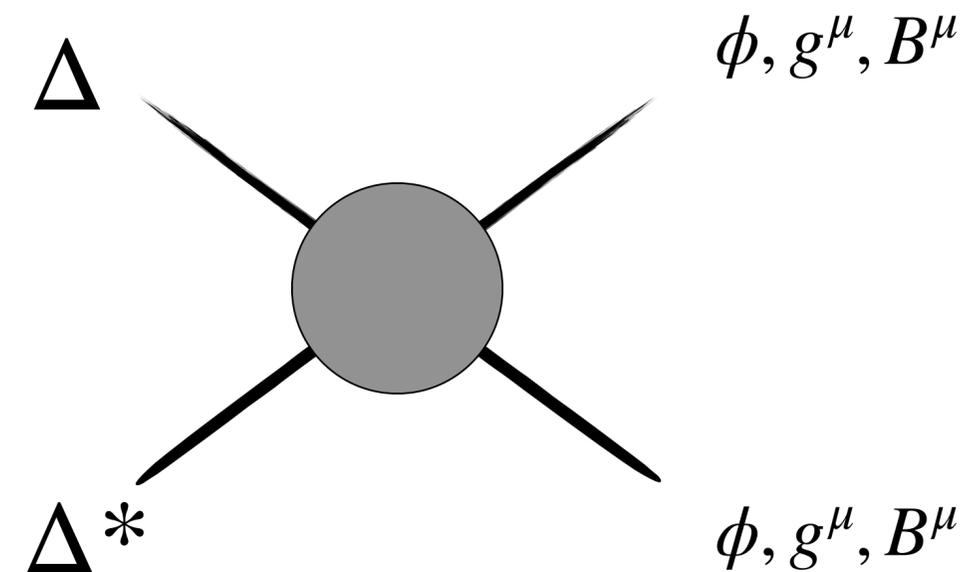
Right after wall crossing:

- CP violating Δ **decays** generate a B-L asymmetry in the visible sector

- Δ **should not annihilate** before they undergo asymmetric decays: $\Gamma_{ann} < \Gamma_{\Delta}$

- ▶ In their own gas frame, Δ particles are squeezed by a factor M_{Δ}/T_n

$$\Gamma_{ann} \sim n_{\Delta} \langle \sigma \cdot v_{rel} \rangle$$



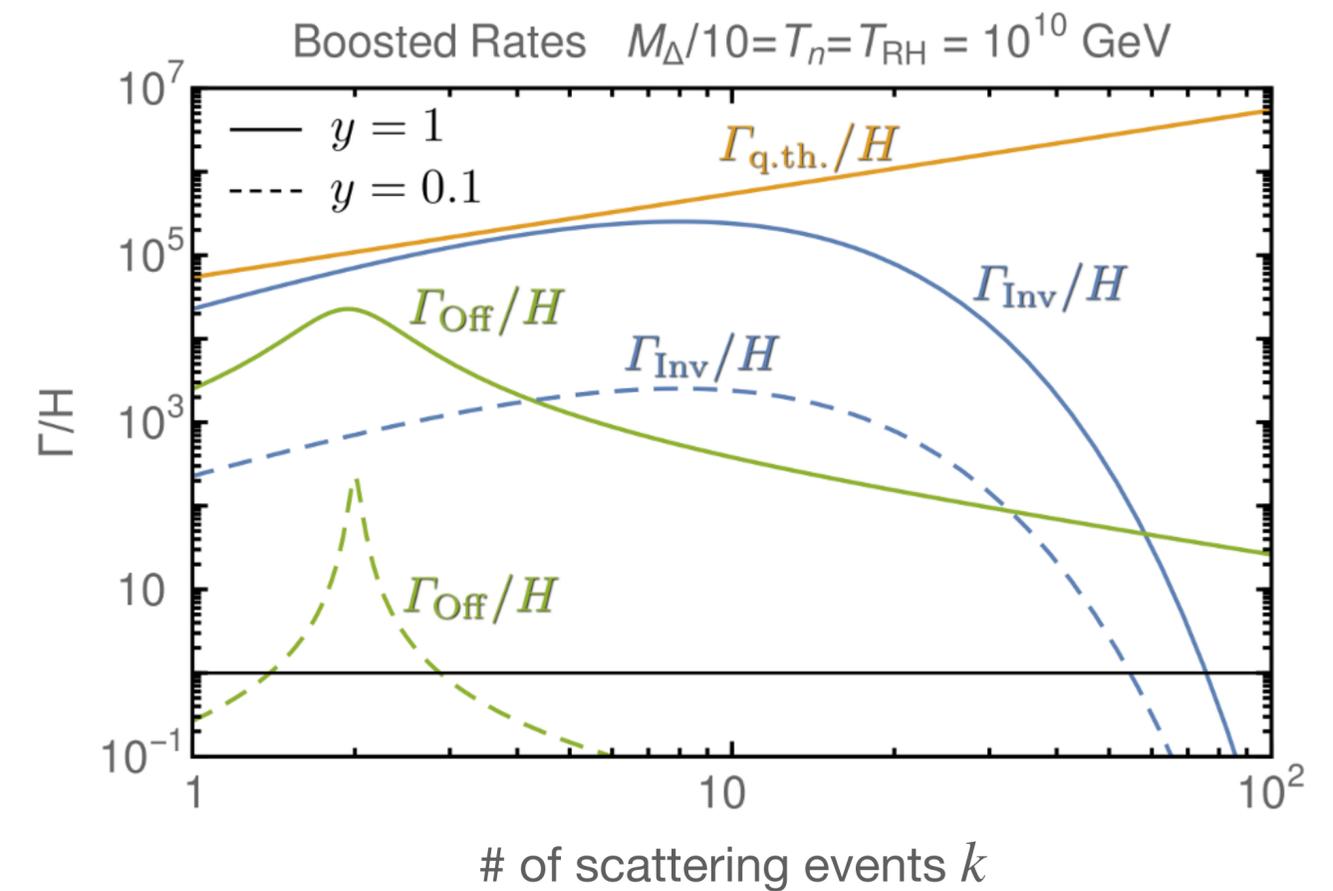
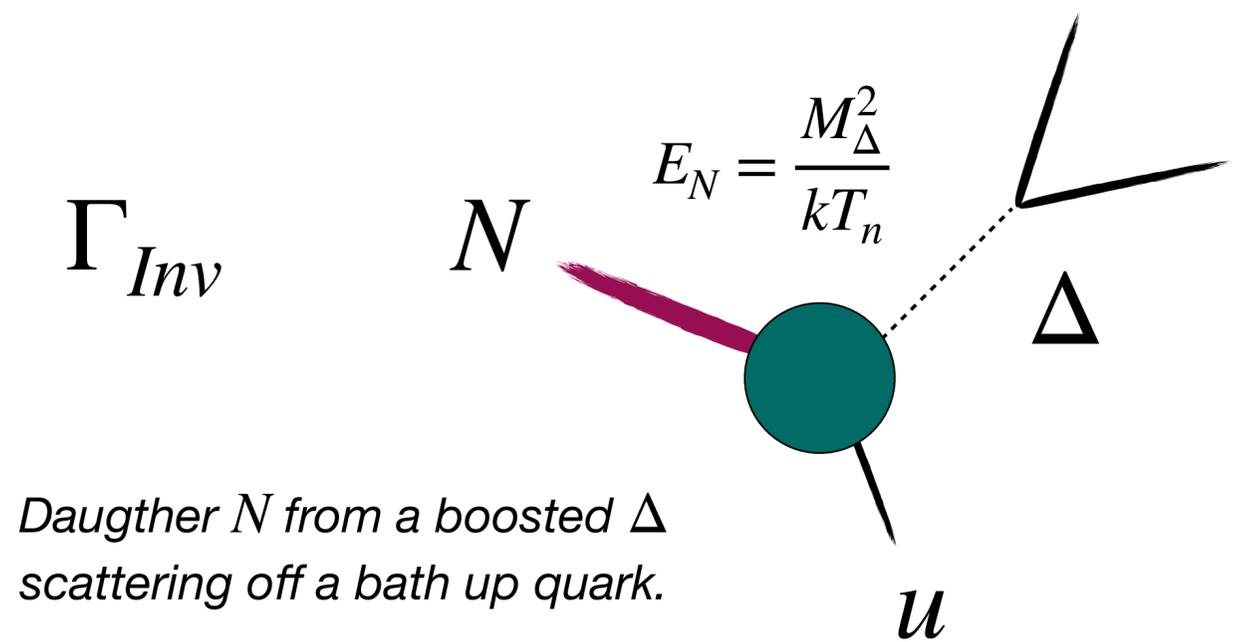
- ▶ The MG relative velocity is $v_{\Delta} \sim T_n/M_{\Delta} \ll 1$, Sommerfeld enhancement may be relevant

Preserving the asymmetry

Right after decay:

- Δ decay products should not erase the asymmetry on the way to kinetic equilibrium

▶ Δ decays with boost $\sim M_\Delta/T_n$ intact



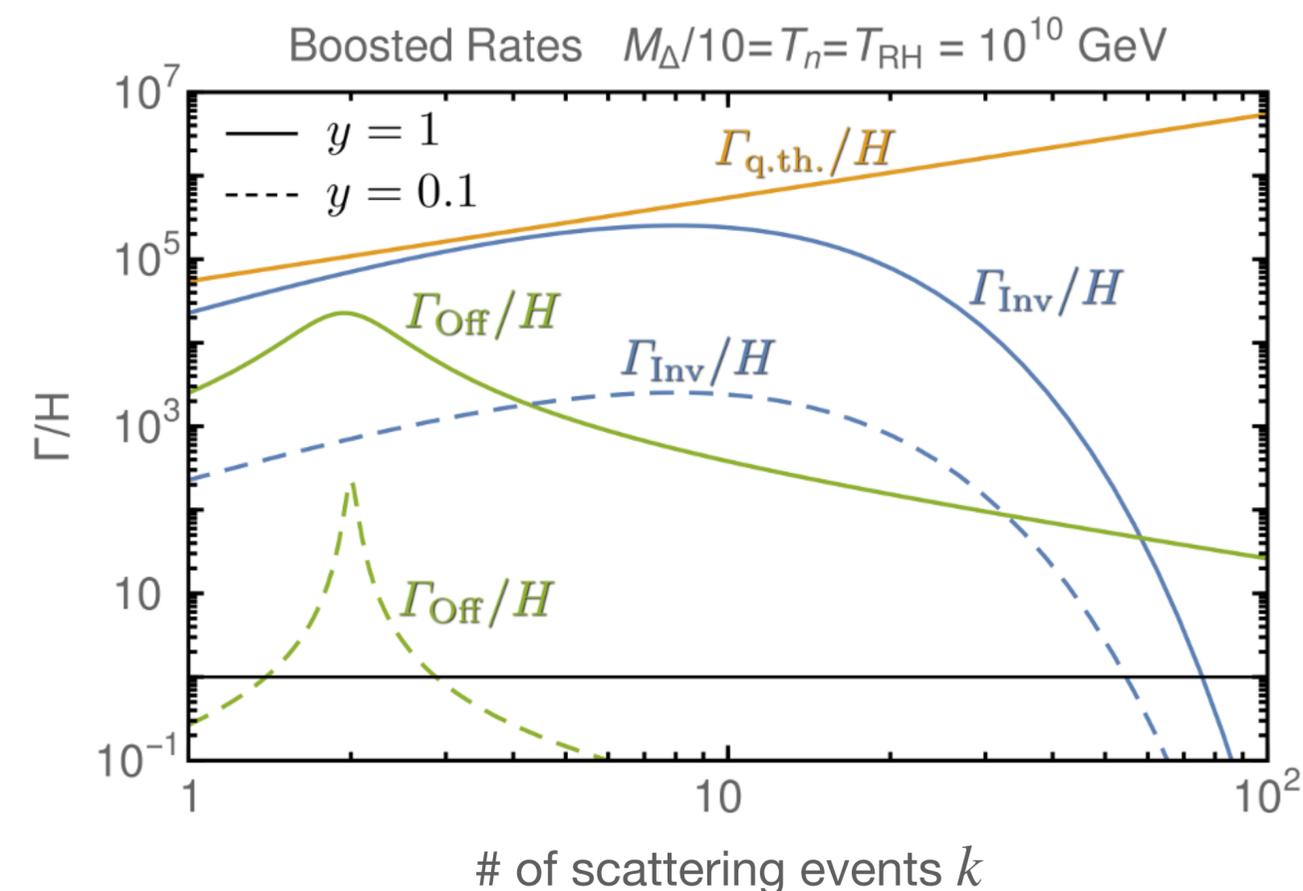
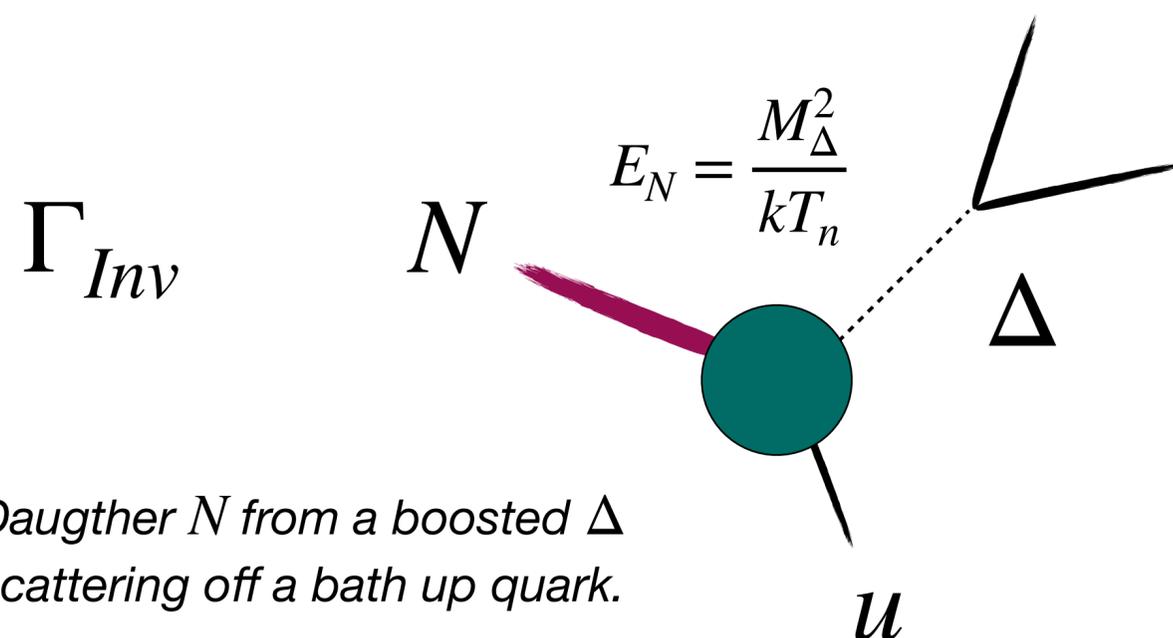
- ◆ All the inverse decays are fast, but quarks thermalize faster
- ◆ N_R needs to thermalize fast as well to not erase the asymmetry: additional interactions

Preserving the asymmetry

Right after decay:

- Δ decay products should not erase the asymmetry on the way to kinetic equilibrium

- ▶ Δ decays with boost $\sim M_\Delta/T_n$ intact
- ▶ Δ decays at rest

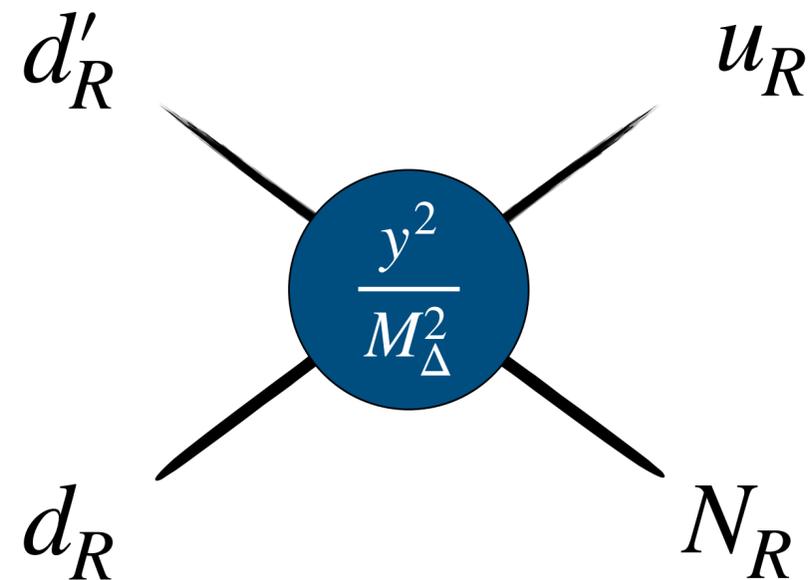


- ◆ All the inverse decays are fast, but quarks thermalize faster
- ◆ N_R needs to thermalize fast as well to not erase the asymmetry: additional interactions

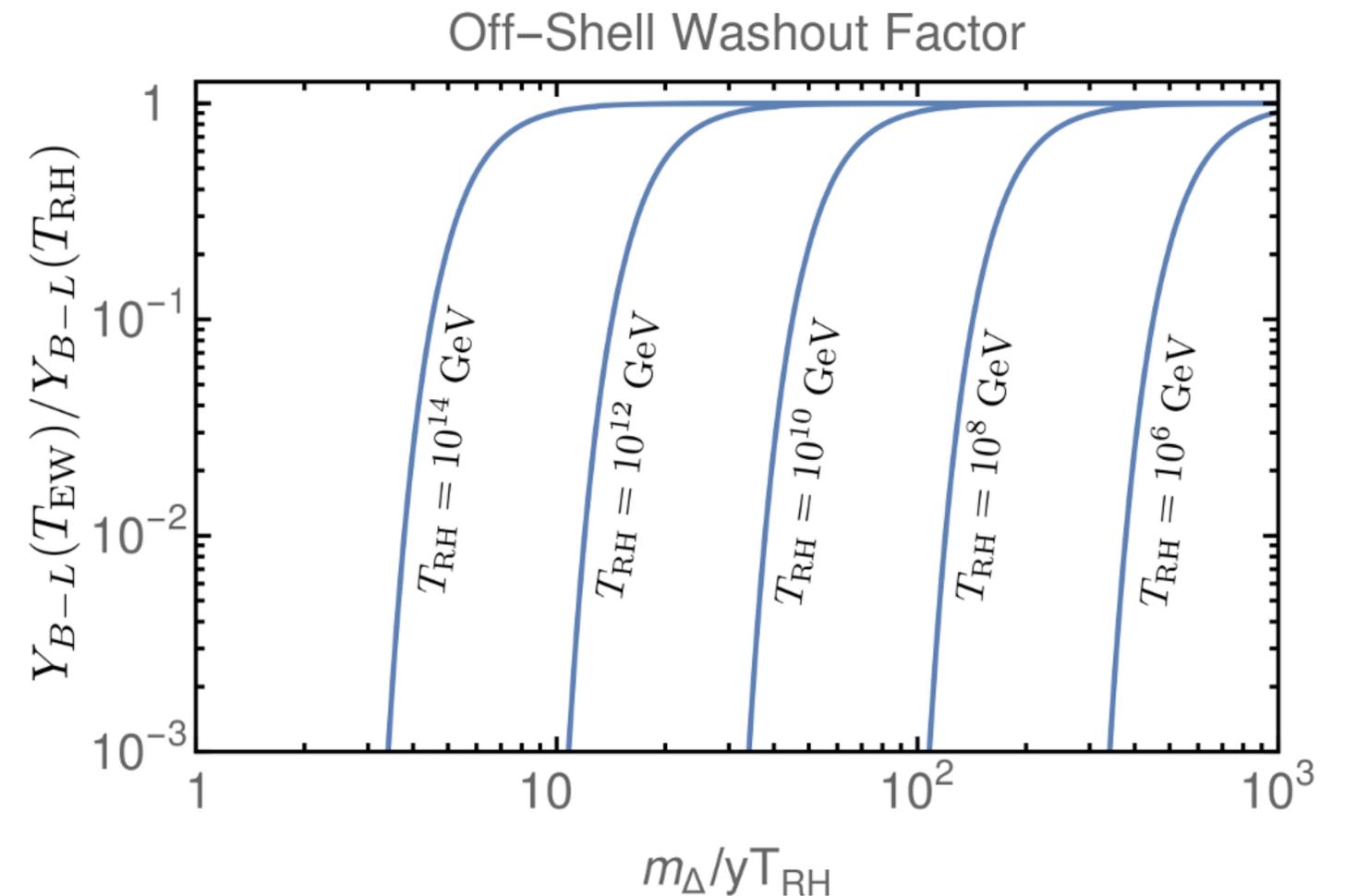
Preserving the asymmetry

After equilibrium is reached at $T = T_{RH}$

- Interactions mediated by Δ can re-equilibrate the B-L charge between dark and visible sector



Effective vertex contributing to the washout from integrating out Δ at tree level.



Not shown: wash out from Boltzmann-suppressed inverse decay rates.

Gravitational waves

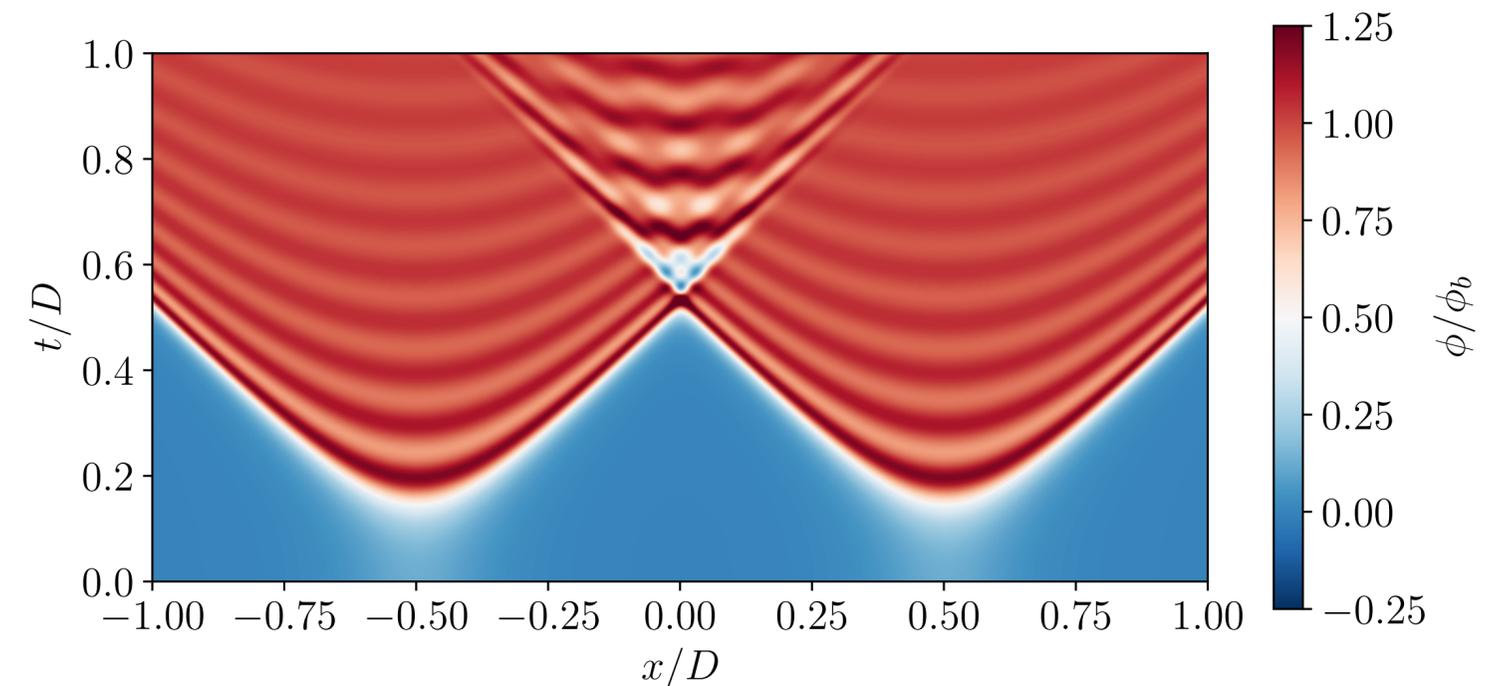
Runaway bubbles

- Latent heat goes into kinetic/gradient energy of expanding scalar field bubbles
- Peak amplitude:

$$h^2 \Omega_{GW} \simeq 4 \cdot 10^{-7} \left(\frac{1}{\beta_H} \right)^2 \left(\frac{\alpha}{1 + \alpha} \right)^2$$

- Peak frequency:

$$f_{peak} \simeq 260 \text{ Hz} \left(\frac{T_{RH}}{10^{10} \text{ GeV}} \right)$$



Simulation of thick-walled bubble collisions with $\gamma = 4$.
 Fig. from Cutting, Escartin, Hindmarsh, Weir [2005.13537] PRD

Gravitational waves

Runaway bubbles

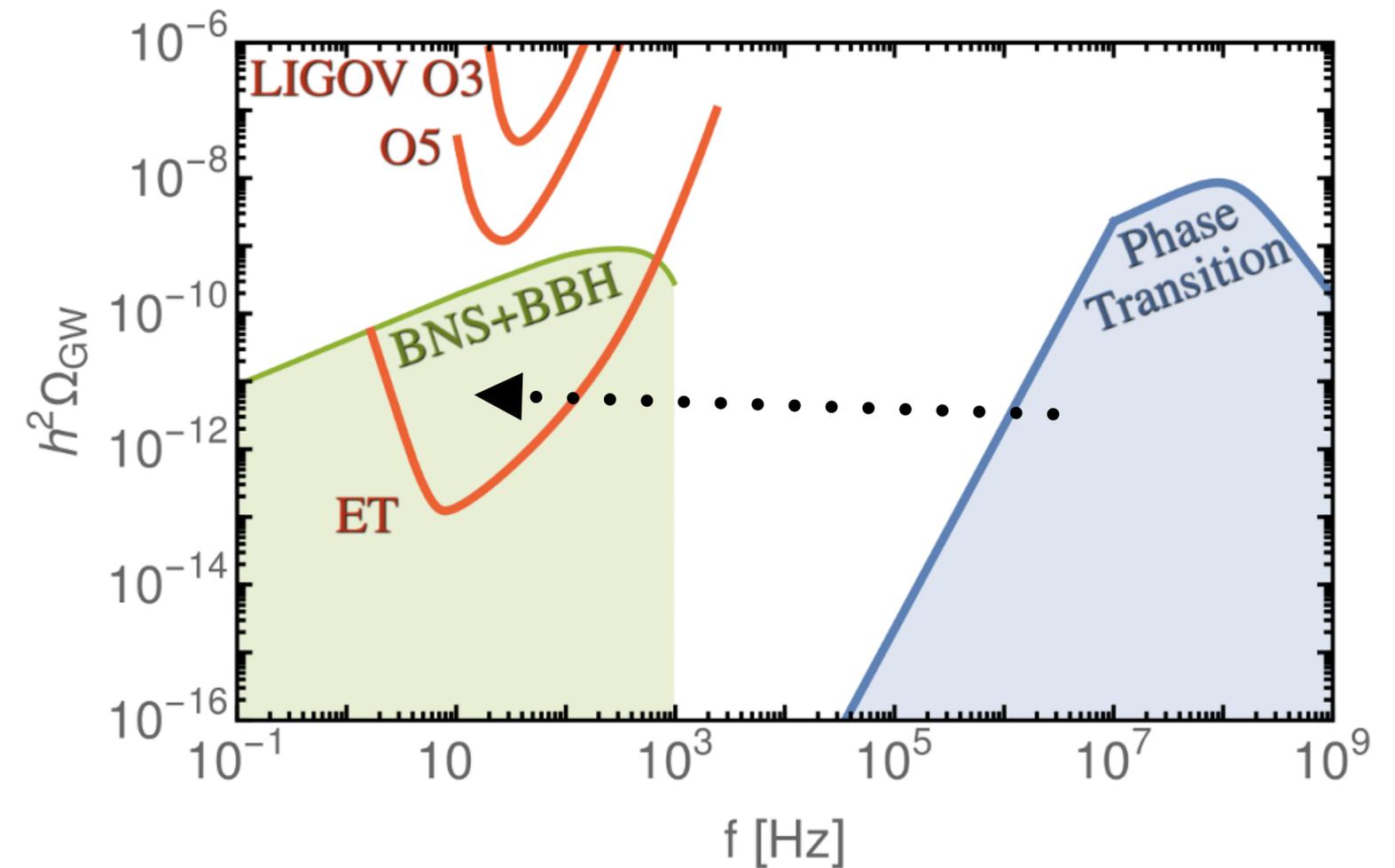
- Latent heat goes into kinetic/gradient energy of expanding scalar field bubbles

- Peak amplitude:

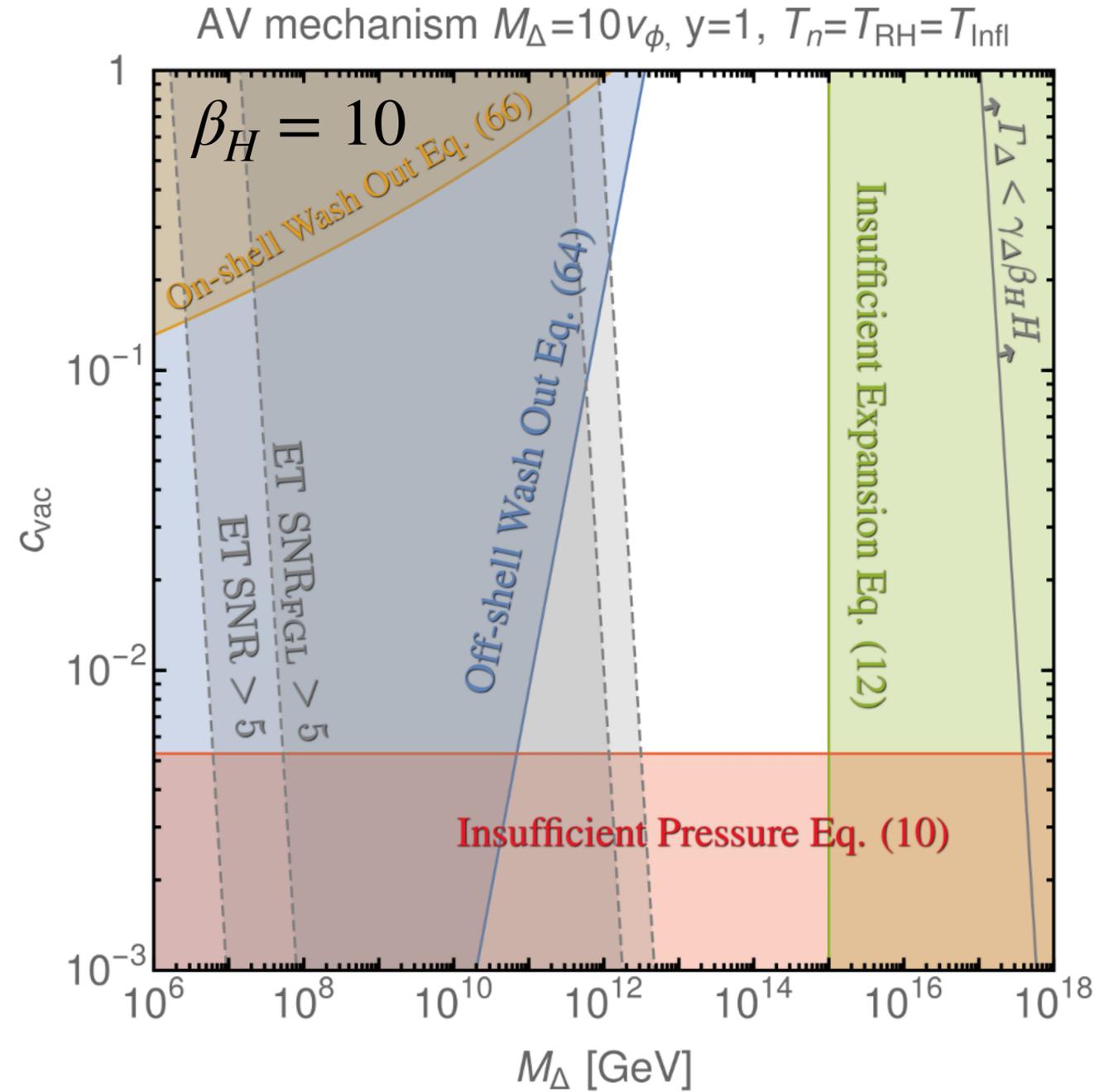
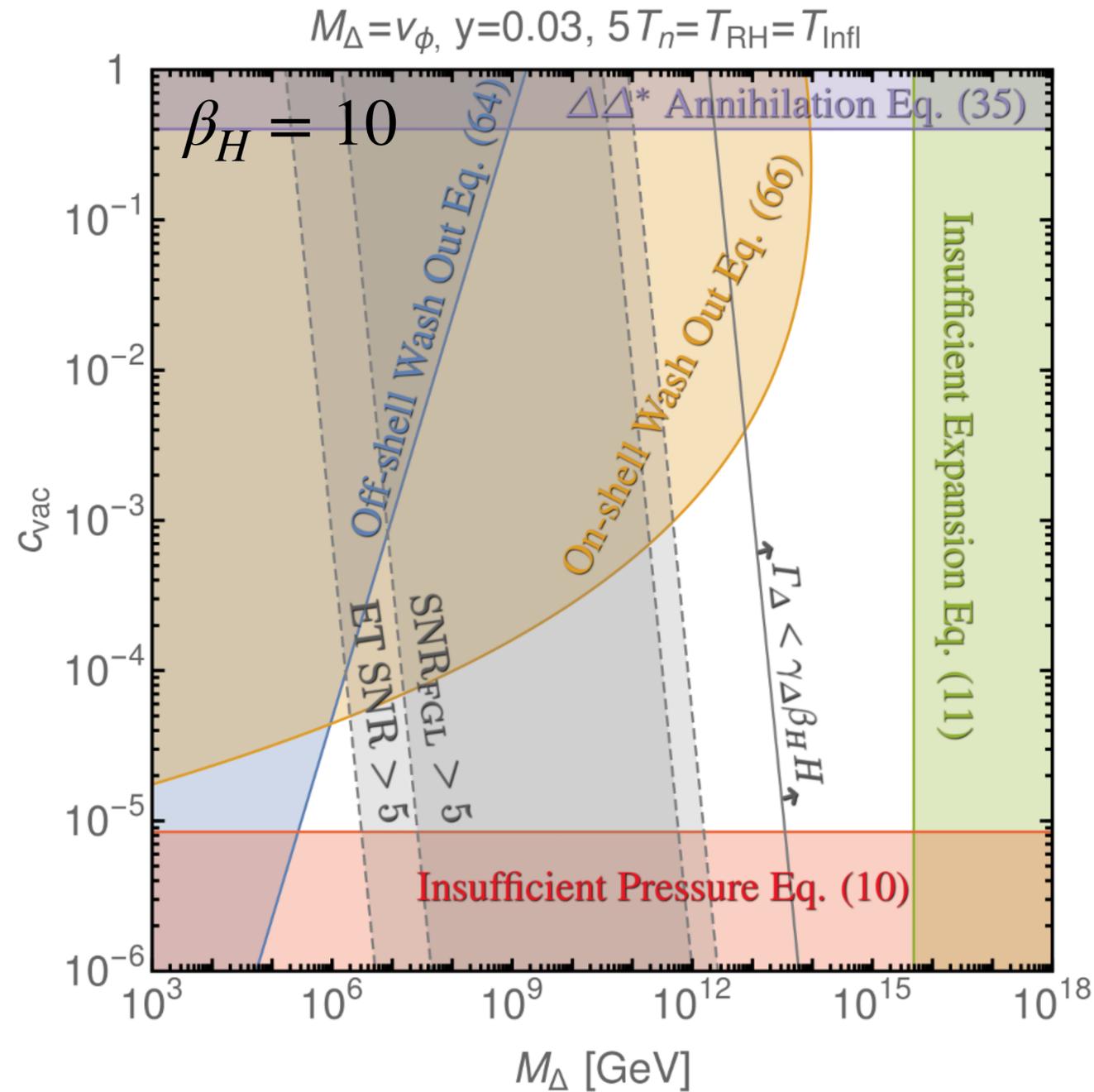
$$h^2 \Omega_{GW} \simeq 4 \cdot 10^{-7} \left(\frac{1}{\beta_H} \right)^2 \left(\frac{\alpha}{1 + \alpha} \right)^2$$

- Peak frequency:

$$f_{peak} \simeq 260 \text{ Hz} \left(\frac{T_{RH}}{10^{10} \text{ GeV}} \right)$$



Summary of the results



Conclusion

- We have presented two mechanisms for baryogenesis with (ultra) relativistic bubble walls during FOPTs: mass gain (MG) and Azatov-Vanvlasselaer heavy particle production (AV)
- We have derived requirements on the phase transition properties, mass spectrum, and CP violation for reproducing the observed baryon asymmetry
- Some of the features/challenges of these mechanisms apply beyond our specific choice of the model
- Interesting interplay with gravitational wave production, even though our setups work best at high scales to avoid washout

Conclusion

- We have presented two mechanisms for baryogenesis with (ultra) relativistic bubble walls during FOPTs: mass gain (MG) and Azatov-Vanvlasselaer heavy particle production (AV)
- We have derived requirements on the phase transition properties, mass spectrum, and CP violation for reproducing the observed baryon asymmetry
- Some of the features/challenges of these mechanisms apply beyond our specific choice of the model
- Interesting interplay with gravitational wave production, even though our setups work best at high scales to avoid washout

Thank you!