

Indirect detection as a probe of the spectrum of primordial perturbations on small scales

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Based on Bringmann, PS, Akrami, in prep.

PS & Sivertsson, Phys Rev Lett 103:211301 (arXiv:0908.4082)

Slides available from <http://www.physics.mcgill.ca/~patscott>



What is an *ultracompact* minihalo?

- Small-scale, large amplitude density perturbations in the early Universe create ultracompact minihalos (UCMHs)
Ricotti & Gould (2009), Berezhinsky et al. (multiple earlier papers)
- ‘Failed primordial black holes’
- Very dense dark matter minihalos
- Enhanced perturbations produced by e.g. inflaton potential or phase transitions
- Non-baryonic, diffuse MACHOs
- Also excellent indirect detection targets



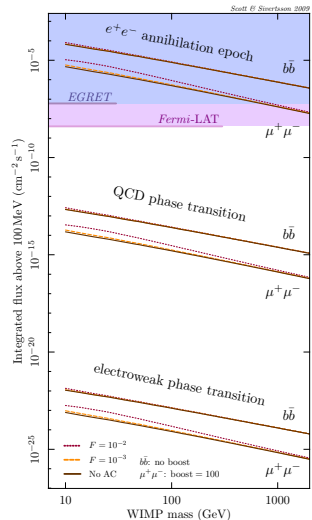
UCMH formation

- Formed (or at least ‘seeded’) well before matter-radiation equality
- Full gravitational collapse before $z = 1000$
- Requires $\delta \gtrsim \mathcal{O}(10^{-3})$
(compare with normal inflationary perturbations $\delta \sim 10^{-5}$)
- $\longrightarrow$ much more likely than PBH formation ($\delta \gtrsim 0.3$)
- UCMH mass is set by horizon scale at time of transition
- $\implies$ specific UCMH mass $\equiv$ specific cosmological scale
- $\implies$ limit on abundance of specific mass halo $\equiv$ limit on power on specific scale k
(Recent limits from Josan & Green (2010), improved analysis here)



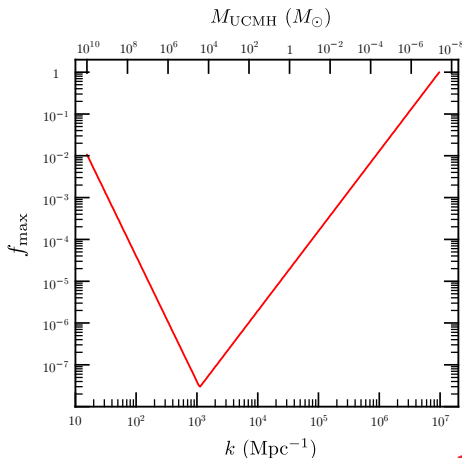
Gamma-ray fluxes from UCMHs

- Take e.g. $d = 100$ pc
- Minihalos produced at e.g. time of the e^+e^- annihilation epoch should be visible today by *Fermi* or HESS, or even in *EGRET* data
- Strongly constrains spectrum of perturbations formed at time of e^+e^- annihilation phase
- Gravitational contraction of UCMHs makes little difference



Limits on present-day UCMH abundance with *Fermi*

- 1-year, 5σ upper limits
- Based on public *Fermi* point source sensitivity
- Proper statistical treatment of observable limit
- Rather conservative assumptions:
 - 100% $b\bar{b}$
 - $\langle\sigma v\rangle = 3 \cdot 10^{-26} \text{ cm}^3 \text{ s}^{-1}$
 - $m_\chi = 5 \text{ TeV}$
 - no DM minihalo detected
- Limits essentially identical for e.g. $\mu^+\mu^-$ with $B = 100$

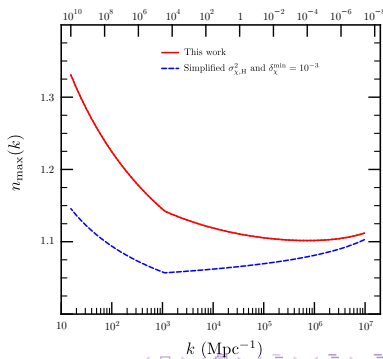


Scale-free (power law) spectrum

$$\sigma_{\chi,H}^2(R) \approx \delta_H^2(t_{k_0}) \left(\frac{k}{k_0} \right)^{n-1} \underbrace{\int_0^\infty x^{n+2} T_\chi^2(x) W_{\text{TH}}^2(x) / T_\chi^2(1) dx}_{\alpha_\chi^2} \quad (1)$$

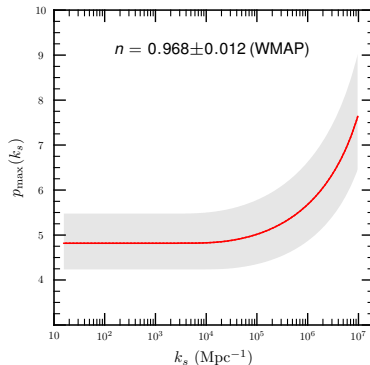
$M_{\text{UCMH}} (M_\odot)$

- Improved $\sigma_{\chi,H}^2$ using top hat window function
- Explicit calculation of $\delta_{\min}$
 (Solution to linear growth eqs;
 $\delta = 1.686$ during matter
 domination in linear approx
 $\Rightarrow \delta \rightarrow \infty$ in non-linear
 regime)



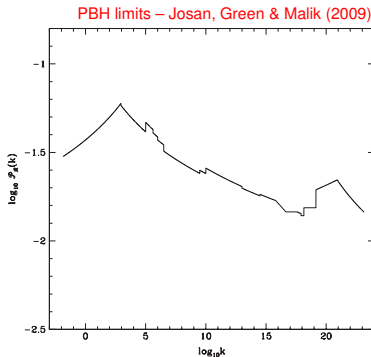
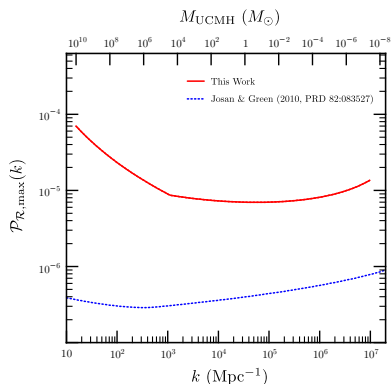
Spectrum with a step

$$\alpha_{\chi}^2 \rightarrow \int_0^{\infty} x^{n+2} T_{\chi}^2(x) W_{\text{TH}}^2(x) / T_{\chi}^2(1) \left\{ \theta\left(\frac{k_s}{k} - x\right) + p^2 \theta\left(x - \frac{k_s}{k}\right) \right\} dx \quad (2)$$



General spectrum and curvature perturbation limits

$$\sigma_{\chi, H}^2(R) = \frac{1}{9} \mathcal{P}_{\mathcal{R}}(k_R) \int_0^\infty x^{n_R+2} W_{\text{TH}}^2(x) W_{\text{TH}}^2(x/\sqrt{3}) dx \quad (3)$$



Summary

- Ultracompact minihalos are promising indirect detection targets
- Could be visible by *Fermi*/VERITAS/HESS/CTA
- Assuming DM annihilates, non-observation places limits on primordial perturbations at small scales
- Derived limits are **much** tighter than existing ones from primordial black holes
- ... but weakened slightly by careful treatment of mass variance, radial profile and minimum density contrast required for formation



UCMH density profiles

- Initially single UCMH per horizon, minimal initial angular momentum
- $\implies$ UCMHs form via radial infall
- $\implies$ Very steep radial density profile

$$\rho_{\chi}(r) = \frac{3f_{\chi}M_{\text{UCMH}}}{16\pi R_{\text{UCMH}}^{\frac{3}{4}}r^{\frac{9}{4}}}, \quad (4)$$

- Truncated at self-annihilation radius r_{SA} , and radius r_{AM} where gas angular momentum violates radial infall approx.

$$\rho(r_{\text{SA}}) = \frac{m_{\chi}}{\langle\sigma v\rangle(t-t_i)}, \quad r_{\text{AM}} = \frac{\sigma_{\text{DM}}^{\frac{8}{7}}R_{\text{UCMH}}^{\frac{11}{7}}}{G^{\frac{4}{7}}M_{\text{UCMH}}^{\frac{4}{7}}} \quad (5)$$



UCMH relic density calculation

- For some distribution of perturbations $\text{pdf}(\delta)$

$$f_{\text{UCMH}} = \left(\frac{1 + z_{\text{eq}}}{1 + z_{\text{stop}}} \right) \int_{\delta_{\text{min}}}^{\delta_{\text{PBH}}} \text{pdf}(\delta) d\delta \quad (6)$$

- For Gaussian perturbations,

$$\text{pdf}(\delta) = \frac{1}{\sqrt{2\pi}\sigma_{\chi,H}^2(z_X, R)} \exp\left(-\frac{\delta^2}{2\sigma_{\chi,H}^2(z_X, R)^2}\right) \quad (7)$$

- Improved $\sigma_{\chi,H}^2$ using top hat window function
- Explicit calculation of δ_{min}



Best scale for limits on step spectrum

$$\alpha_{\chi}^2 \rightarrow \int_0^{\infty} x^{n+2} T_{\chi}^2(x) W_{\text{TH}}^2(x) / T_{\chi}^2(1) \left\{ \theta\left(\frac{k_s}{k} - x\right) + p^2 \theta\left(x - \frac{k_s}{k}\right) \right\} dx$$

