The Tameness of QFT and CFT

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- What is Tameness?
- Tameness of QFT
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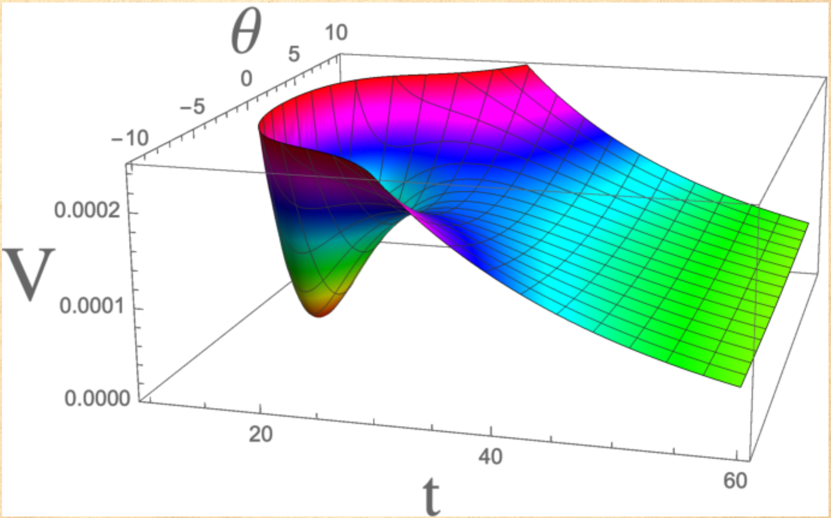
  - Tameness of the observables in CFTs
  - Tameness of the observables in EFTs

What is tameness?

- Tameness is a generalized finiteness principle
- Forbids *discrete* infinities
- Idea: Functions appearing in physics should only behave in finitely many different ways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o-minimal structures

- Allow only functions which are definable in an o-minimal structure S

Definition of a Structure

Collections $S = (S_n)_{\geq 1}$ *of sets in* $\mathbb{R}^n$ *closed under* $\cup, \cap, \times, /$ *and linear projections containing at least all algebraic sets (= zero sets of polynomials).*

Definition of o-minimality

A structure is o-minimal if the definable subsets of $\mathbb{R}$ *are finite unions of intervals and points*

Definable subsets of $\mathbb{R}$

- Only finitely many points and intervals.
- But the intervals can be infinitely long.
- Higher dimensional sets have to project down to these.

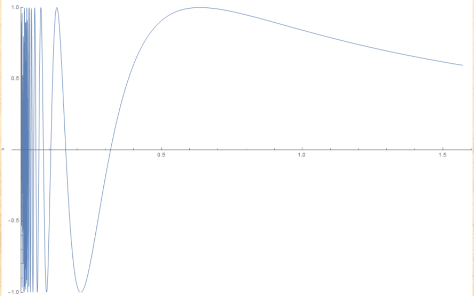
The language

- sets in o-minimal structure: tame sets
- functions whose graph is a tame set: tame functions

$\rightarrow$ tame manifolds, tame bundles, tame geometry

What does this mean in practice?

- o-minimal structures forbid anything infinite discrete
  - no integers $\mathbb{Z}$
  - no periodic functions
  - no sin(x) and cos(x) for $x \in \mathbb{R}$
  - no error or gamma functions on $\mathbb{R}$
  - also restricts functions on finite intervals



Examples of o-minimal structures

- $\mathbb{R}_{\mathrm{alg}}$: semi-algebraic sets ($P(x) \geq 0$ instead of $P(x) = 0$)
- $\mathbb{R}_{\mathrm{an}}$: restricted analytic functions
- $\mathbb{R}_{\mathrm{exp}}$: real exponential function
- $\mathbb{R}_{\mathrm{an,exp}}$: combination of the two above
- $\mathbb{R}_{\mathrm{Pfaff}}$: structure of all Pfaffian functions
- $\ldots$

Applications of tamness

- Used in proofs of many deep mathematical conjectures
  - Ax-Schanuel for Hodge Structures [Bakker, Tsimerman'17]
  - Griffiths'conjecture [Bakker, Brunebarbe, Tsimerman'18]
  - Andre-Oort conjecture [Pila, Shankar, Tsimerman'21]
  - Geometric André-Grothendieck Period Conjecture [Bakker, Tsimerman'22]
- Finiteness of of vacua [Bakker, Grimm, Schnell, Tsimerman'21]

Tameness of QFT

What is the right question?

- Basic questions
  - Which objects are tame?
  - To be able to talk about tameness we need structures, what are the right structures?
  - Do different objects live in different structures or does there exist a common structure?
  - Is there one overarching structure like $\mathbb{R}_{\mathrm{an,exp}}$ or does every QFT define its own structure?
  - Is every QFT tame?

Structures from QFTs

- Interesting objects of a QFT include
  - The Lagrangian/action
  - The partition function Z
  - The correlators
  - Amplitudes/observables
- To be able to talk about a QFT we need a language to formulate everything in
  - A set of theories $\mathcal{T}$, e.g parameter space of specified Lagrangians
  - Set $\mathcal{S}$ of Euclidean spacetimes with metric (Σ, g)
  - Both are definable in some structure $\mathbb{R}^{\text{def}}_{\mathcal{T},\mathcal{S}}$
- Simplest example: Polynomial Lagrangians in $\mathbb{R}^d \to \mathbb{R}^{\text{def}}_{\mathcal{T},\mathcal{S}} = \mathbb{R}_{\text{alg}}$

Structures from QFTs

- The theory is defined using $\mathbb{R}^{\text{def}}_{\mathcal{T},\mathcal{S}}$.
- Add the partition function and correlators to the original structure $\to \mathbb{R}_{\mathcal{T},\mathcal{S}}$
- We often take $\mathbb{R}^d, T^d, S^d$ as our spacetime and drop the explicit dependence on $\mathcal{S}$, e.g $\mathbb{R}_{\text{QFT}}$, $\mathbb{R}_{\text{CFT}}$, $\mathbb{R}_{\text{EFT}}$.

Questions about $\mathbb{R}_{\mathcal{T},\mathcal{S}}$

- If $\mathbb{R}^{\mathrm{def}}_{\mathcal{T},\mathcal{S}}$ is o-minimal, when is $\mathbb{R}_{\mathcal{T},\mathcal{S}}$ o-minimal?.
  - Are observables tame?
- Under which conditions is $\mathbb{R}^{\mathrm{def}}_{\mathcal{T},\mathcal{S}}$ o-minimal?
  - Tameness of the space of theories?

Tameness of perturbative QFT

Theorem: For any renormalizable QFT with finitely many fields and interactions all finite-loop amplitudes are tame functions of the masses, external momenta and coupling constants definable in $\mathbb{R}_{\mathrm{an,exp}}$. [Douglas,Grimm,LS - Part I]

- Feynman integrals are periods
- periods are definable in $\mathbb{R}_{\mathrm{an,exp}}$ [Bakker,Klingler,Tsimerman][Bakker,Mullane '22]

  $\rightarrow$ Feynman integrals are definable

  $\rightarrow$ Amplitudes are definable
- If the Lagrangian is tame the *perturbative* corrections will not destroy this tameness!

$\rightarrow$ perturbative QFTs are tame if the Lagrangian is tame

Feynman integrals

- Can write a l-loop Feynman integral in a d-dimensional theory in Lee-Pomeransky representation [Lee,Pomeransky 13']

$$I = \frac{\Gamma(\frac{d}{2})}{\Gamma\left(\frac{(\ell+1)d}{2} - \nu\right)\prod_{j=1}^{n}\Gamma(\nu_j)} \int_{x_j \geq 0} \prod_{j=1}^{n} \mathrm{d}x_j x_j^{\nu_j - 1} G^{-\frac{d}{2}} .$$

- G is the Lee-Pomeranski polynomial depending on the masses and external momenta.
- We can interpret this as the defining polynomial of variety in projective space.

$$\omega_i = \int_{\gamma_i} \Omega \equiv \int_{\gamma_i} \frac{\mathrm{d}x_1 \wedge \mathrm{d}x_2 \wedge \ldots \wedge \mathrm{d}x_n}{P(a_j, x_i)} ,$$

- Lots of technical details involved in the identification! (integration contours, open vs closed chains, divergences ...). In a detailed analysis these work out.

What about non-perturbative effects?

- Instantons appear to produce cos potentials$\rightarrow$ appear to be dangerous.
- The Feynman diagram argument does not help due to the non-perturbative nature.
- But: Tameness is not conserved under power series expansion!

$$x^2 = \frac{\pi^2}{3} - 4cos(x) + cos(2x) + \ldots$$

- Look at some examples of exactly solvable theories.

Gauged linear sigma models

- 2d theory with $\mathcal{N} = 2$ supersymmetry.
- Exactly solvable by supersymmetric localization.
- The sphere partition function is given in terms of the Kähler potential of the described geometry [Jockers et al. 12']

$$Z_{S_2} = e^{-K} = \overline{\Pi}\Sigma\Pi$$

- As the partition function is given in terms of periods it is definable in $\mathbb{R}_{\mathrm{an,exp}}$!

Solvable 0d QFTs

- On points the path integral reduces to usual integrals
- Many 0d QFTs are solvable like the Sine-Gordon model or the ϕ^4 theory

$$Z(m,\lambda) = \int_{-\infty}^{\infty} \mathrm{d}\phi \; e^{-\frac{m^2}{2}\phi^2 - \frac{\lambda}{4!}\phi^4} = \sqrt{\frac{3}{\lambda}} e^{\frac{3m^4}{4\lambda}} \, m \, K_{1/4}\left(\frac{3m^4}{4\lambda}\right) ,$$

$$Z(g) = \int_{-\pi}^{\pi} \mathrm{d}\phi \; e^{-g \; sin(\phi)^2} = 2e^{-g/2}\pi I_0(g/2) ,$$

- I_0 is a tame function, period of an explicit geometry
- $K_{1/4}$ is an exponential period, tameness of these is an open question!
- The explicit functions in the model form a Pfaffian chain, definable in $\mathbb{R}_{\mathrm{Pfaff}}$ [Van den Dries, private communication]

Other tame examples

- 1d harmonic oscillator

$$Z(\beta, m) = \frac{1}{sinh(\frac{\beta}{2m})}$$

- 2d string theories, 3d non-critical M-theory

$$F_{3d}(\omega, \mu) = -\frac{1}{6\omega^2}\mu^3 + \frac{\Lambda}{4\omega}\mu^2 - \frac{1}{2\pi\omega_0}\mu^2 \log(1 - e^{-2\pi\mu/\omega}) + \frac{1}{2\pi^2}\mu Li_2(1 - e^{-2\pi\mu/\omega}) + \frac{\omega}{4\pi^3} Li_3(1 - e^{-2\pi\mu/\omega})$$

- 2d Yang-Mills theory

$$Z_{SU(2)} = e^{A\lambda/16}(\theta_3(e^{-A\lambda/16}) - 1)$$

- Klein-Gordon field in *d*-dimensional AdS

$$O^{(d)}(y_1, y_2) = (2\pi)^{-d/2} \left(\frac{(y_2 - y_1)^2}{\sqrt{m}} \right)^{\frac{d-2}{2}} K_{\frac{d-2}{2}}(\sqrt{m}(y_2 - y_1)^2) \,.$$

Is every QFT tame?

No! Can construct explicit counterexamples:

- infinite discrete symmetries $Z(g \cdot \lambda) = Z(\lambda)$
  - Need to be gauged or broken
  - Fits with no global symmetries conjecture [Banks,Dixon 88'][Banks,Seiberg 10']
- non-tame Lagrangian
  - Simple example : $V(\theta) = A\cos(\theta) + B\cos(\alpha\theta) \qquad \alpha$ irrational
  - Allows for infinite spirals $\rightarrow$ tension with distance conjecture [Grimm,Lanza,Li 22]
- Observables also need not be tame
  - Neutrino oscillations

Tameness of EFTs

Conjecture: All effective theories valid below a fixed finite energy cut-off scale Λ that can be coupled to QG are labelled by a tame parameter space and have scalar field spaces and Lagrangians that are tame in an o-minimal structure [Grimm 21']

Conjecture: $\mathbb{R}_{\mathrm{EFTd}}[\Lambda]$ are o-minimal structures, i.e. observables are also tame [Douglas,Grimm,LS - Part II]

What about string theory

- Perturbative string theory has infinitely many fields$\rightarrow$ not a tame theory
- The partition functions are expressible via θ functions $\rightarrow$ tame functions
- Any effective theory with finite cutoff Λ is tame.
- In 2d string theory one can understand what happens to the infinite discrete modes
- Goldstone modes of broken area-preserving diffeomorphisms of 3d theory
- Only a toy model!

Tameness of conformal field theories

Tameness of CFTs

Conjecture 1: All observables of a tame set $\mathcal{T}_{\mathrm{CFT}}$ are tame functions.
[Douglas,Grimm,LS - Part II]

Conjecture 2(a): The theory space $\mathcal{T}_{\mathrm{CFT}}$ in d=2 is tame if
- the central charge is bounded
- lowest operator dimension is bounded from below.

[Douglas,Grimm,LS - Part II]

Conjecture 2(b): The theory space $\mathcal{T}_{\mathrm{CFT}}$ in $d > 2$ is tame if
- an appropriate measure for the degrees of freedom is bounded
- theories differing by discrete gaugings are identified.

[Douglas,Grimm,LS - Part II]

Evidence for the conjectures - Observables

- Conformal symmetry fixes the form of the 2- and 3-point correlators to

$$\langle \mathcal{O}_i(x_1)\mathcal{O}_j(x_2)\rangle = \frac{\delta_{ij}}{(x_1 - x_2)^{\Delta_i+\Delta_j}} \, ,$$

$$\langle \mathcal{O}_1(x_1)\mathcal{O}_2(x_2)\mathcal{O}_3(x_3)\rangle = \frac{C_{1,2,3}}{x_{12}^{\Delta_1+\Delta_2-\Delta_3} x_{23}^{-\Delta_1+\Delta_2+\Delta_3} x_{13}^{\Delta_1-\Delta_2+\Delta_3}} \, ,$$

- Trivially tame in the positions and operator dimensions
- First non-trivial case is the 4-point correlator

Evidence for the conjectures - Observables

- Conformal symmetry fixes the dependence on the positions and weights in terms of conformal partial waves (blocks) $W_{\mathcal{O}}$:

$$\langle \mathcal{O}_1(x_1)\mathcal{O}_2(x_2)\mathcal{O}_3(x_3)\mathcal{O}_4(x_4)\rangle = \sum_{\mathcal{O}\in\mathcal{O}_1\times\mathcal{O}_2} C_{1,2,\mathcal{O}} C_{3,4,\mathcal{O}} W_{\mathcal{O}} \,,$$

$$W_{\mathcal{O}} = \frac{1}{C_{1,2,\mathcal{O}} C_{3,4,\mathcal{O}}} \sum_{\alpha\in\text{descendants}} \langle 0| \mathcal{O}_1\mathcal{O}_2 |\alpha\rangle \langle\alpha| \mathcal{O}_3\mathcal{O}_4 |0\rangle$$

- Conformal partial waves are Lauricella type hypergeometric functions
- In 2d Gauss hypergeometric functions

$$W^{2d}_{\mathcal{O}}(\Delta) = \left(\frac{x_{14}^2}{x_{13}^2}\right)^{\frac{\Delta_{34}}{2}} \left(\frac{x_{24}^2}{x_{14}^2}\right)^{\frac{\Delta_{12}}{2}} \frac{u^{\Delta/2} v^{\Delta/2}}{x_{12}^{\Delta_1+\Delta_2} x_{34}^{\Delta_3+\Delta_4}} .$$

$${}_2F_1\left(\frac{\Delta-\Delta_{12}}{2}, \frac{\Delta+\Delta_{34}}{2}; \Delta; u\right) {}_2F_1\left(\frac{\Delta-\Delta_{12}}{2}, \frac{\Delta+\Delta_{34}}{2}; \Delta; v\right) ,$$

Evidence for the conjectures - Observables

$$W^{2d}_{\mathcal{O}}(\Delta) = \left(\frac{x^2_{14}}{x^2_{13}}\right)^{\frac{\Delta_{34}}{2}} \left(\frac{x^2_{24}}{x^2_{14}}\right)^{\frac{\Delta_{12}}{2}} \frac{u^{\Delta/2} v^{\Delta/2}}{x_{12}^{\Delta_1+\Delta_2} x_{34}^{\Delta_3+\Delta_4}}.$$

$$ {}_2F_1\left(\frac{\Delta-\Delta_{12}}{2}, \frac{\Delta+\Delta_{34}}{2}; \Delta; u\right) {}_2F_1\left(\frac{\Delta-\Delta_{12}}{2}, \frac{\Delta+\Delta_{34}}{2}; \Delta; v\right),$$

- Tame in the positions
- Tameness in Δ more complicated
- Analysis shows that the differences in operator dimensions $\Delta_{i,j}$ need to be bounded!
- Fits nicely with the no parametric separation of scales conjecture [Lüst, Palti, Vafa 19']

Evidence for the conjectures - Space of theories - 2d

- The space of 2d CFTs is clearly not o-minimal
- Many infinite discrete sets exist, e.g. unitary minimal models

$$c = 1 - \frac{6}{(p+1)(p+2)} \xrightarrow{p\to\infty} 1 \qquad \Delta_1 = \frac{3}{4(p+1)(p+2)} \xrightarrow{p\to\infty} 0$$

- WZW models

$$c = 3 - \frac{6}{(p+2)} \xrightarrow{p\to\infty} 1 \qquad \Delta_1 = \frac{3}{4(p+2)} \xrightarrow{p\to\infty} 0$$

- For fixed lower bound on Δ_1 only finitely many theories

Evidence for the conjectures - Space of theories - 3d

- Again many families of theories parameterized by discrete choices of parameters
- 3d Chern-Simons theory: gauge group N and level k
  $\rightarrow$ naively a lattice $\mathbb{Z}^2$ of theories
- Dualities identify different choices, e.g level-rank duality
$$F(N,k) = F(k,N) = \frac{N}{2} log(k+N) + \dots$$
- For fixed upper bound of F only finitely many theories!

Outlook

- Study consequences of tameness
- Connection to complexity
- Relations between amplitudes(Ax-Schanuel)

Summary

- Perturbative QFTs are tame if the Lagrangian is tame
- Non-perturbative tameness requires restrictions on the theories
  - CFT with bounded degrees of freedom and finite gap
    $\rightarrow$ leads to bounds on operator dimensions
  - EFT originating in QG
    $\rightarrow$Tameness as a swampland conjecture
- Dualities play an important role in the tameness of the theory space

Beyond Large Complex Structure: Quantized Periods and Boundary Data for One-Modulus Singularities

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Abstract

We study periods near boundaries in one-dimensional complex structure moduli spaces of Calabi-Yau threefolds. Near the large complex structure point these asymptotic periods are well understood in terms of the topological data of a mirror Calabi-Yau manifold. The aim of this work is to characterize the period data near other boundaries in moduli space such as generalized conifold points and so called K-points. We provide general models for these asymptotic periods in a quantized three-form basis. Moreover, we elucidate the geometrical meaning of the model-dependent coefficients that appear in these expressions: we find that these can be identified with certain topological and arithmetic numbers associated to the singular Calabi-Yau geometry at the moduli space boundary. To be more precise, for conifold points we encounter L-function values of the modular form associated to the conifold itself; for K-points we find a correspondence between the log-monodromy matrix and intersection forms of rigid K3 surfaces.