Understanding the tower weak gravity conjecture

Cesar Fierro Cota

Based on: 2212.09758 and 23XX.XXXXX

In collaboration with: Alessandro Mininno, Timo Weigand, and Max Wiesner.

Outline

- Part I: The weak gravity conjecture an overview
- Part II $\begin{cases} A : \text{The asymptotic weak gravity conjecture} \\ B : \text{The weak gravity conjecture with no infinite towers} \end{cases}$

Part I

-In this section statements are independent of string theory.

The weak gravity conjecture

Consider a U(1) gauge theory coupled to gravity.

The weak gravity conjecture:

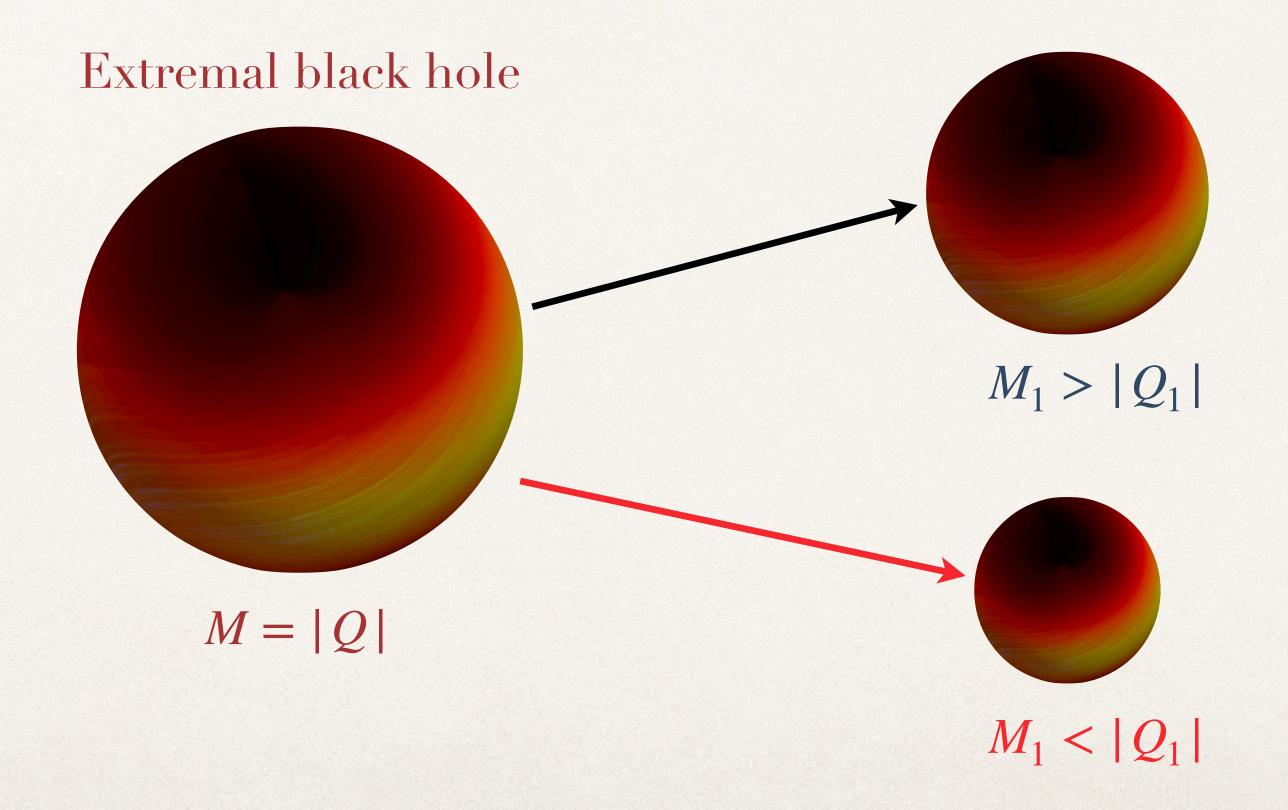
There must exist a super-extremal state of charge q and mass m, such that

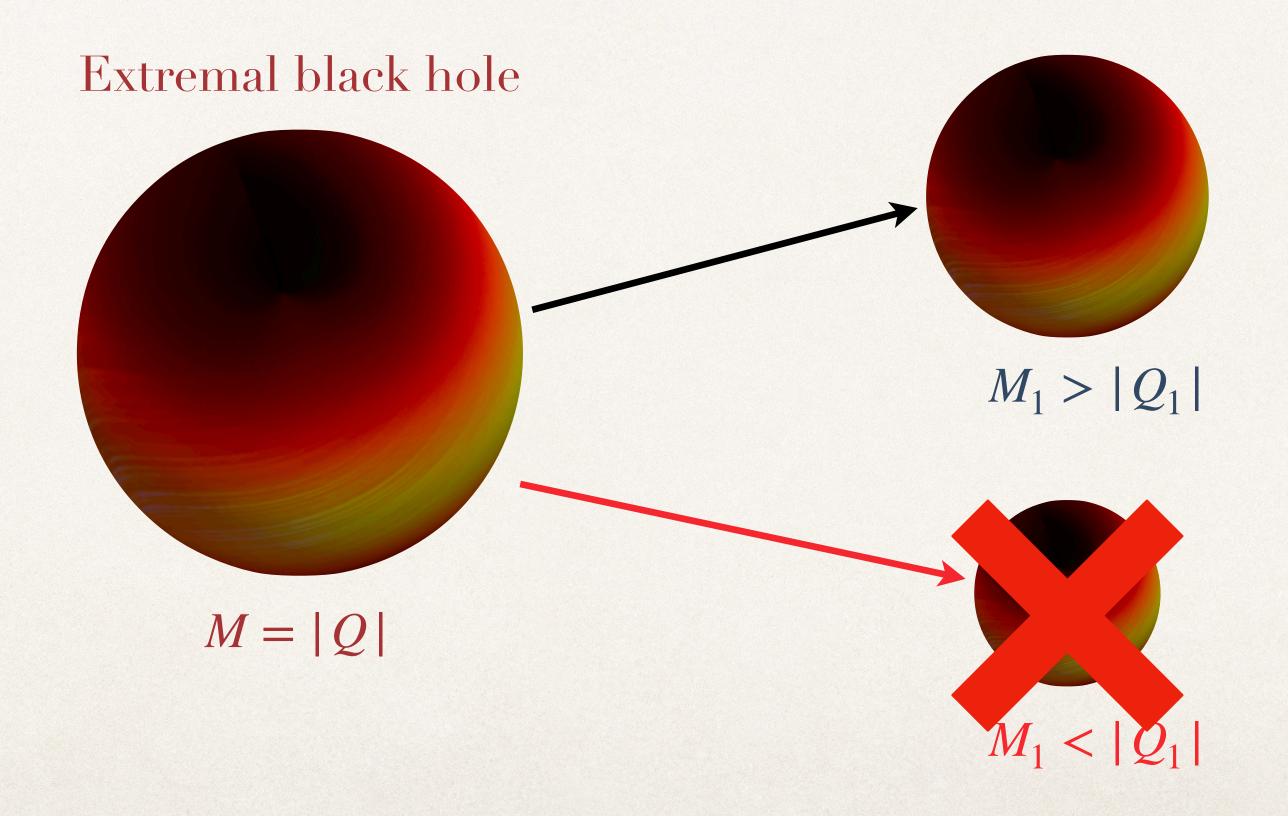
$$\frac{|q|}{m} \ge \frac{|Q|}{M} \Big|_{\text{Ext}} \,. \tag{1}$$

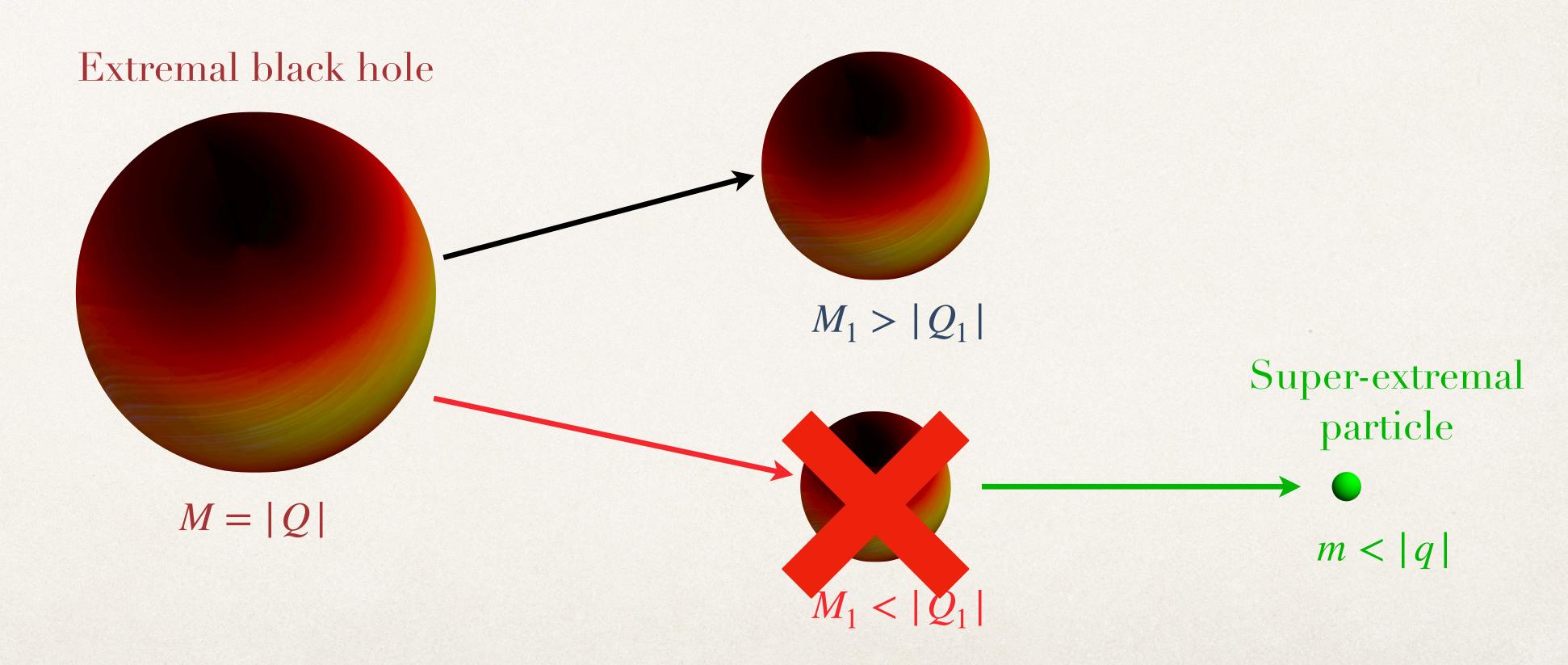
Q: Black hole charge

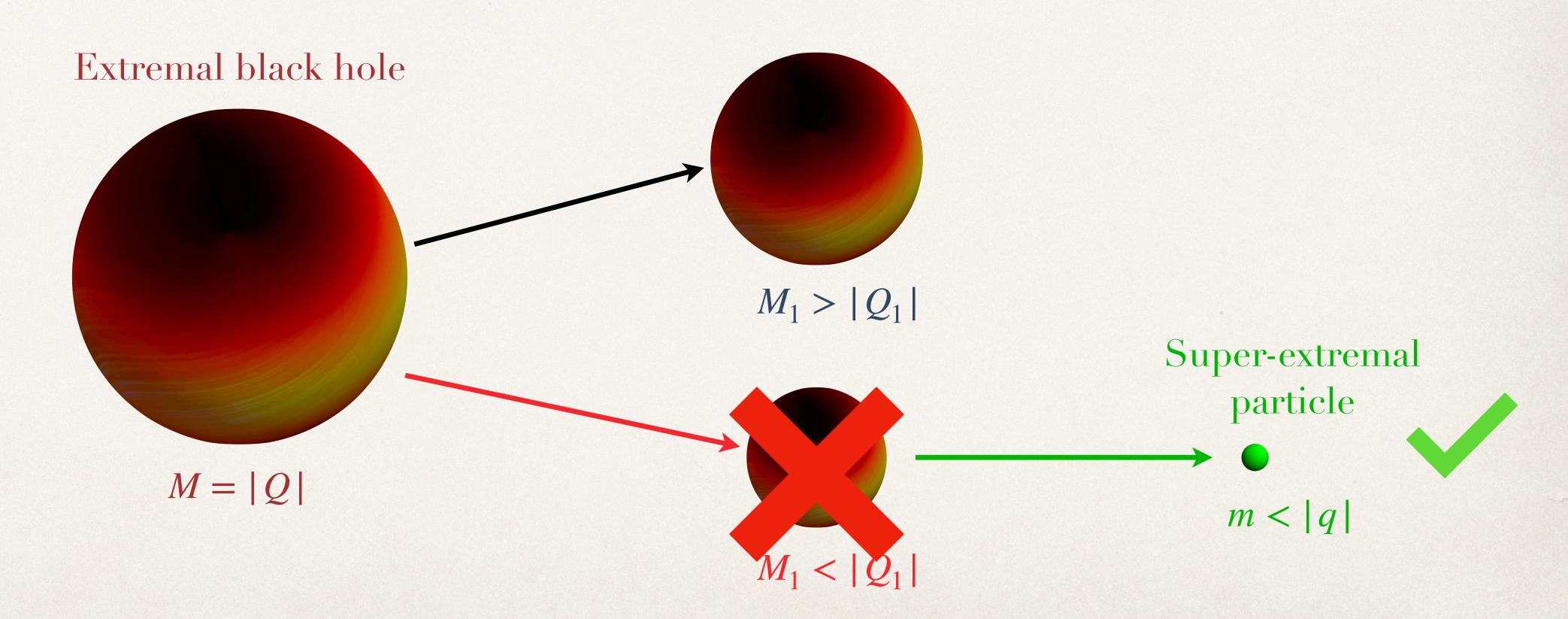
M: Black hole mass

[Arkani-Hamed, Motl, Nicolis, Vafa'06]









The weak gravity conjecture inequality can be generalized for a theory in D-dimensions with several abelian gauge factors, as well as several scalar fields.

However, one has to be careful about the multiple U(1) gauge group factors.

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e.g. consider a theory with $U(1)_A \times U(1)_B$ gauge group.

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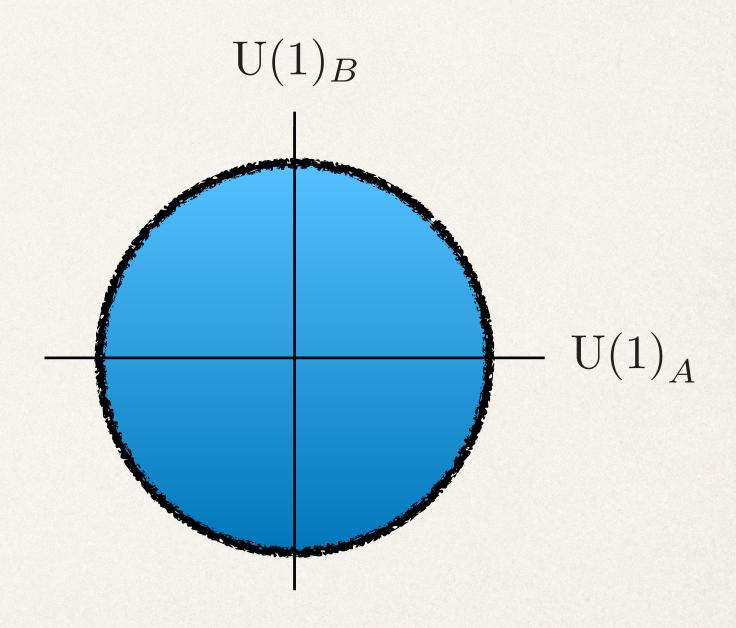
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$$\mathbf{z}_i := M_{ ext{Pl}}^{rac{d-2}{2}} rac{\mathbf{q}_i}{m_i}$$

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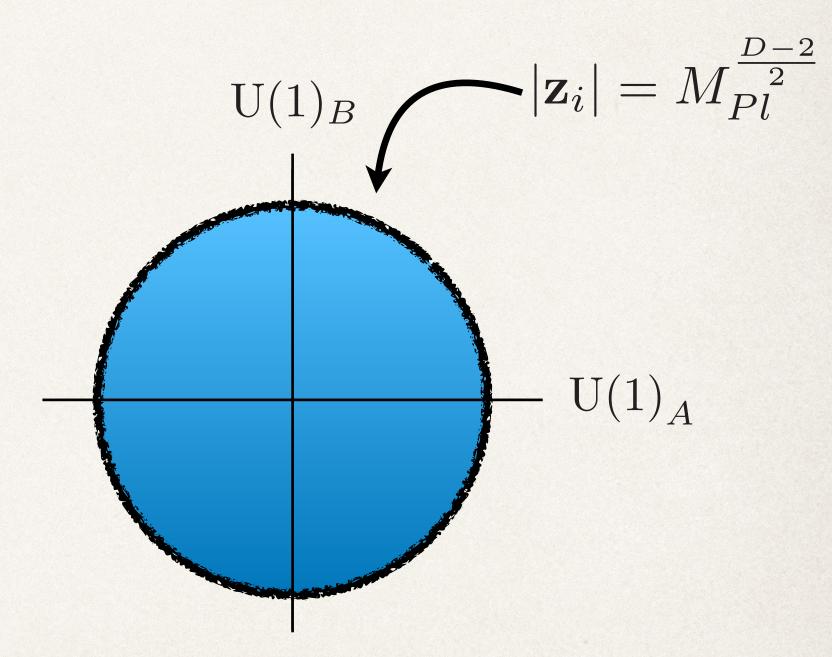
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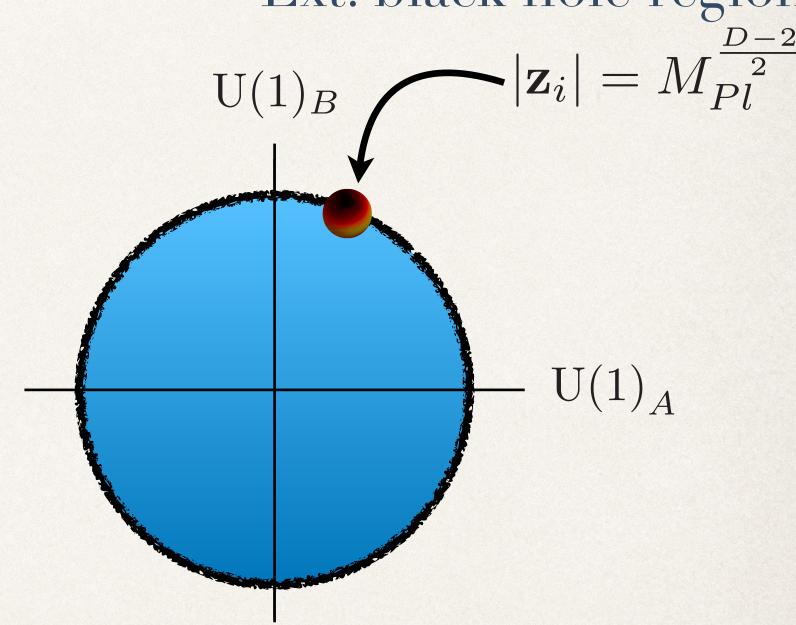


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Ext. black hole region:

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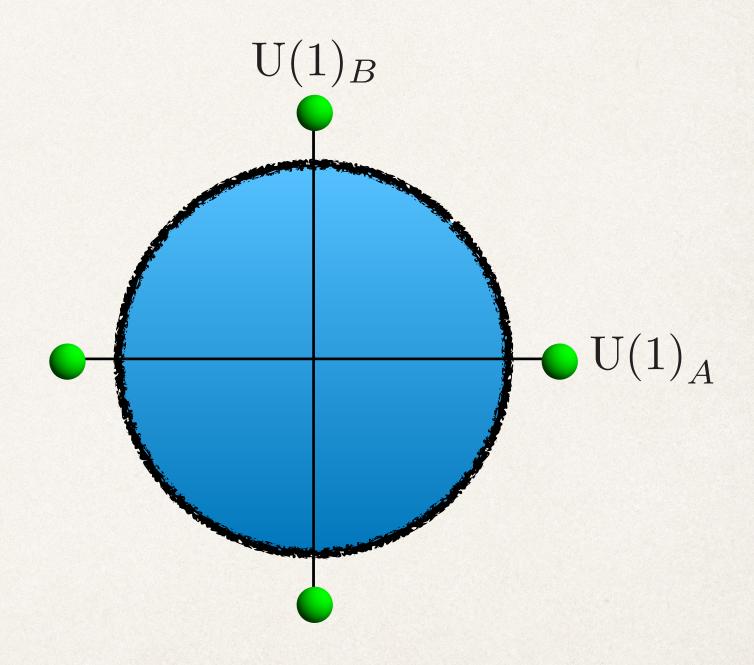
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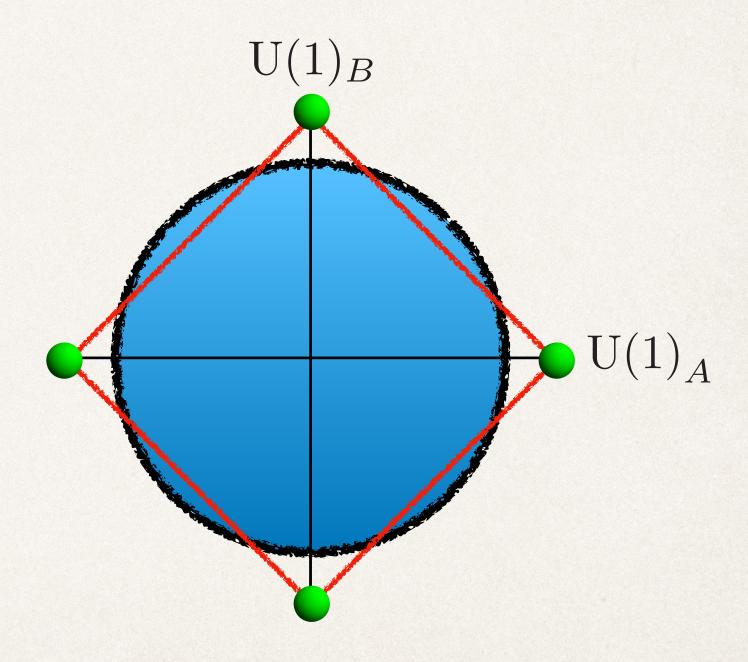
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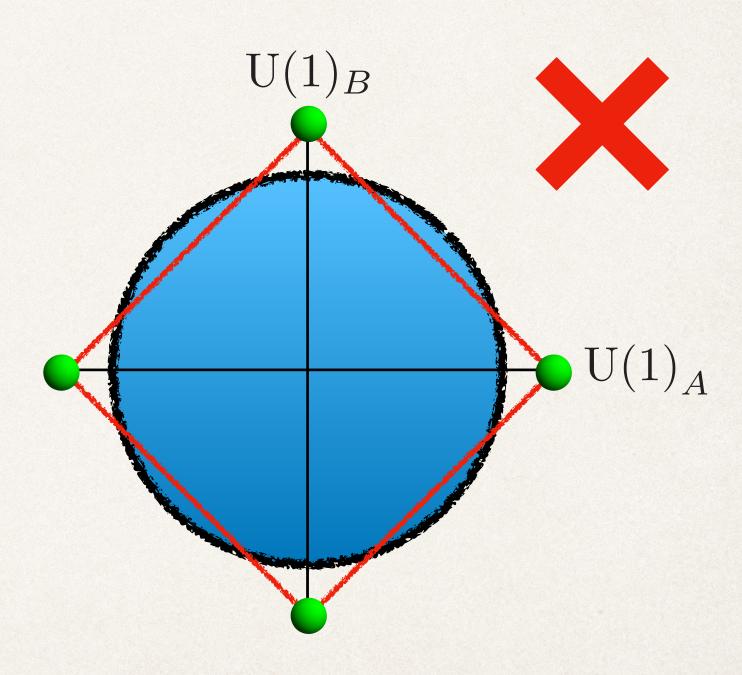
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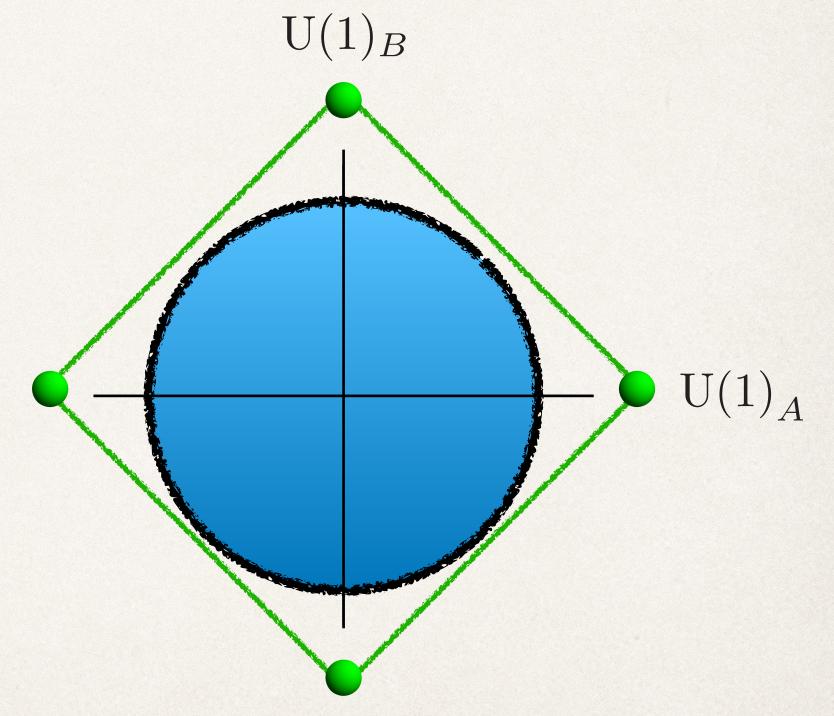
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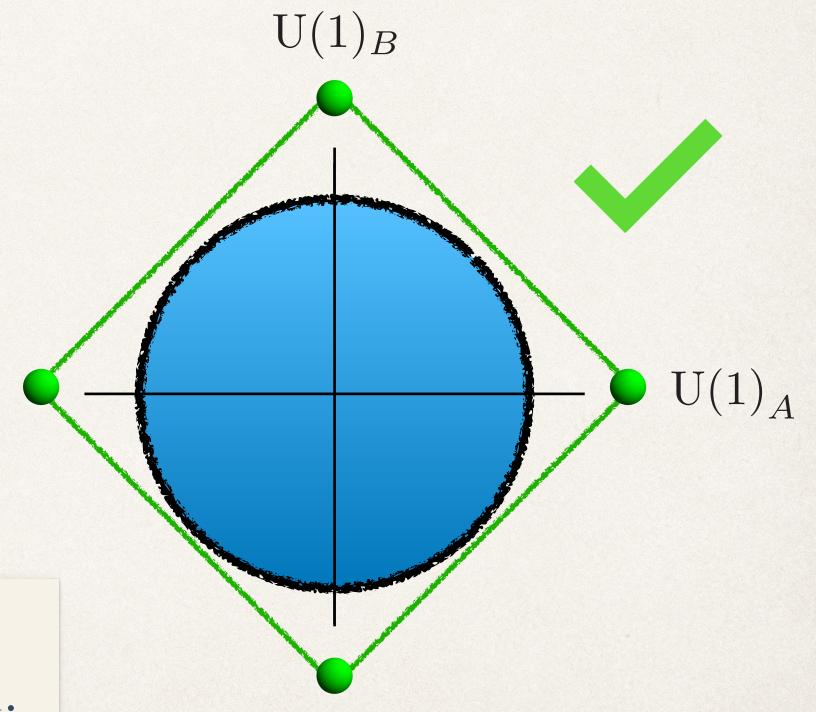
e.g. consider a theory with $U(1)_A \times U(1)_B$ gauge group.

Define the charge-to-mass ratio of an *i*-th particle:

$$\mathbf{z}_i := M_{\mathrm{Pl}}^{rac{d-2}{2}} rac{\mathbf{q}_i}{m_i}$$

Convex Hull Condition: [Cheung, Remmen'14]

There is a set $\{\mathbf{z}_i\}_{i\in\mathcal{I}}$ whose convex hull contains the unit ball.



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 S^1

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Convex hull condition satisfied whenever: [Heidenreich, Reece, Rudelius'15]

$$(m_D r_{S^1})^2 \ge \frac{1}{4z_D^2(z_D^2 - 1)}$$

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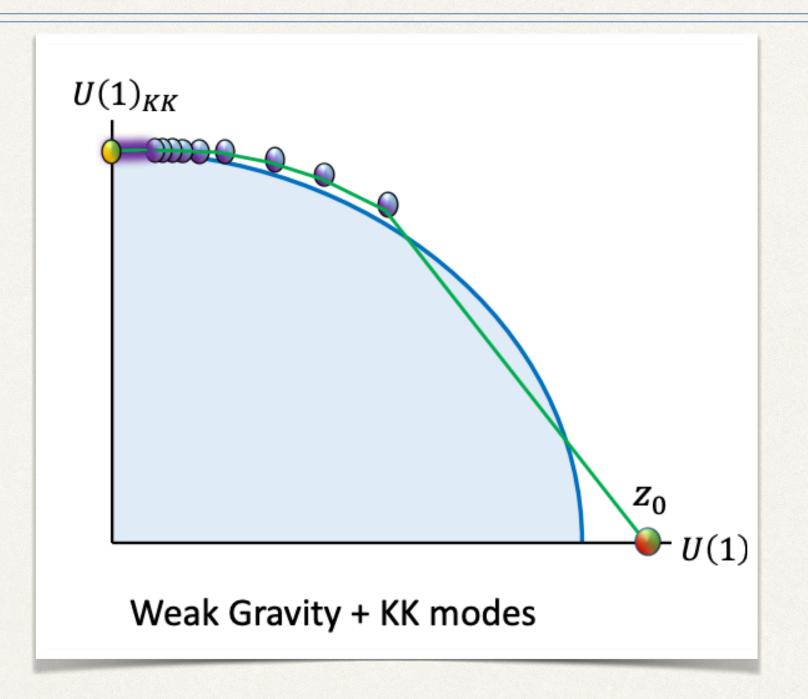
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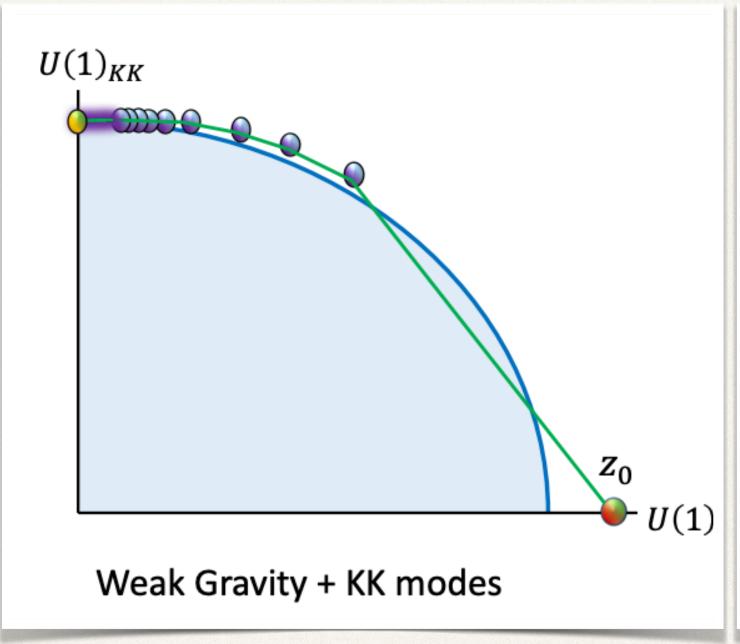
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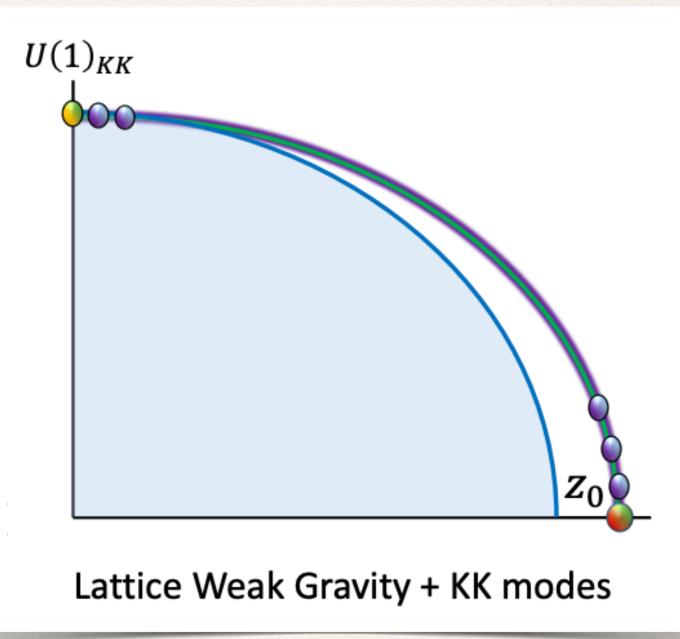


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[Pictures: Palti'19]

Resolution: [Heidenreich, Reece, Rudelius'15]

- 1. The effective field theory breaks down at $r_{S^1} = r_{\min} \sim \Lambda_D^{-1}$.
- 2. There exists an infinite tower of super-extremal states such that $z_D \ge 1$.

The tower weak gravity conjecture

Evidence for the tWGC in string theory:

Heidenreich, Reece, Rudelius'15'16'17

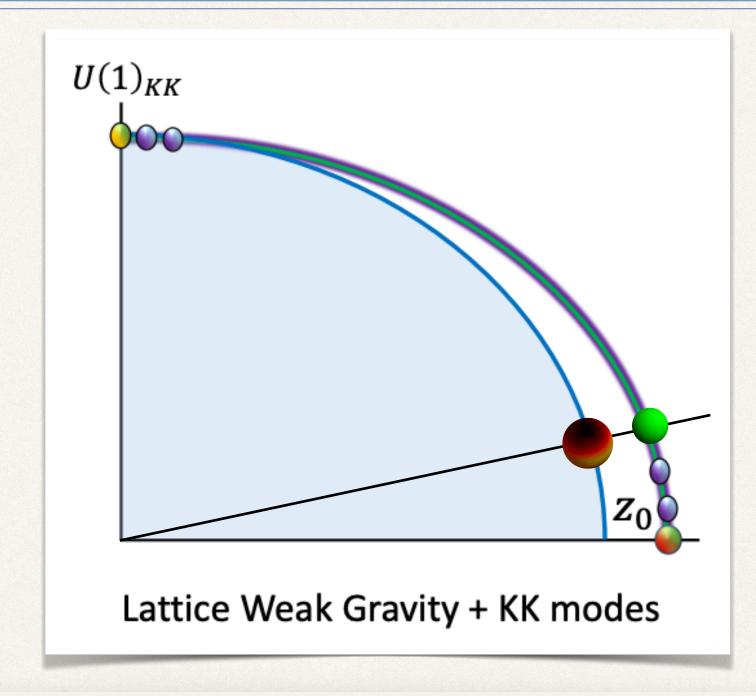
[Lee, Lerche, Weigand'18]

[Klaewer, Lee, Weigand, Wiesner'18]

[CFC, Klemm, Schimannek'20]

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Most examples are related to the heterotic string.



Also:

[Alim, Heidenreich, Rudelius'21]

[Gendler, Heidenreich, McAllister, Rudelius'22]

Super-extremal states here are black holes.

The tower weak gravity conjecture: Heidenreich, Reece, Rudelius'15'16 Andriolo, Junghans, Noumi, Shiu'16 For every rational direction \hat{Q} in charge space, a super-extremal state of charge and mass must exist, such that $\vec{q}/m \propto \hat{Q}$.

Part II

-From now on: F-theory/M-theory compactified on a Calabi-Yau threefold.

Part II — A: The asymptotic weak gravity conjecture

-Q: When can we relate super-extremal states with particle-like states in the EFT?

The weak gravity conjecture has an associated an energy scale: $\Lambda_{\rm WGC}=g_{U(1)}^2M_{\rm Pl}^{d-2}$

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Let Λ_{QG} be the energy scale where gravity becomes strongly coupled.

We call asymptotic weak gravity conjecture those limits in which the weak gravity conjecture can be described such that

$$rac{\Lambda_{
m WGC}}{\Lambda_{
m QG}}
ightarrow 0 \, .$$

In M-theory on a Calabi-Yau threefold X_3 , an asymptotic limit $\Lambda_{\text{WGC}}/\Lambda_{\text{QG}} \to 0$ is realized in the Kähler moduli space of a fibered X_3 by shrinking its fiber while blowing its base. [CFC, Mininno, Weigand, Wiesner'22]

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Gauge group: $U(1) = c_{\alpha}U(1)^{\alpha} \longleftrightarrow C = c_{\alpha}C^{\alpha}$; $\{C^{\alpha}\}$ is basis of curves in X_3 and $c_{\alpha} \in \mathbb{Z}_{\geq 0}$.

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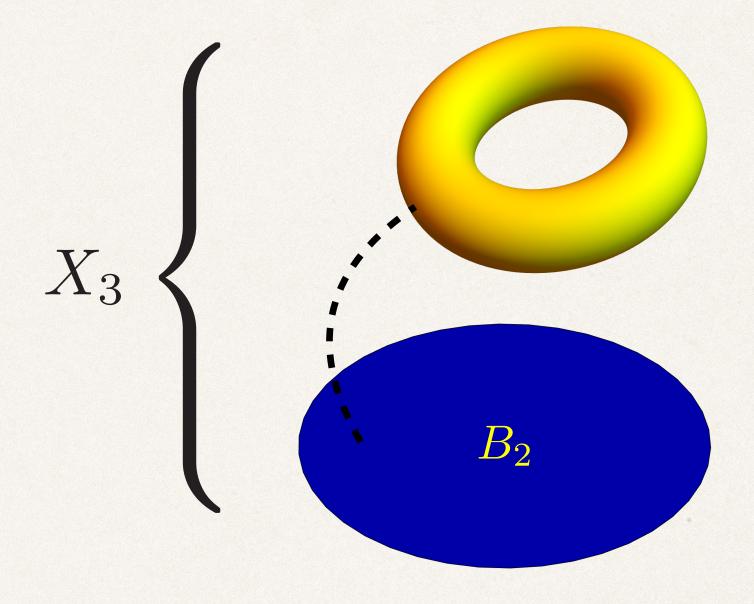
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Tower weak gravity conjecture in the asymptotic limit holds for [CFC, Minimo, Weigand, Wiesner'22]

- 1. A T^2 -fibered Calabi-Yau threefold $\pi: X_3 \to B_2$.
- 2. A K3-fibered Calabi-Yau threefold $\rho: X_3 \to \mathbb{P}^1$.

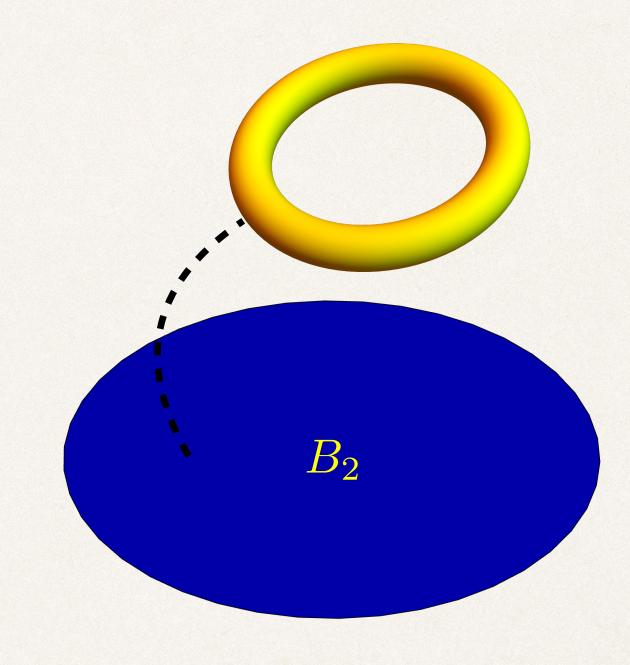
1. T^2 -fibered Calabi-Yau threefold $\pi: X_3 \to B_2$. Here U(1) group determined by T^2 curve.

Super-extremal states counted by Gopakumar-Vafa invariants $n_{nT^2}^0 = -\chi(X_3), \ n \in \mathbb{N}.$



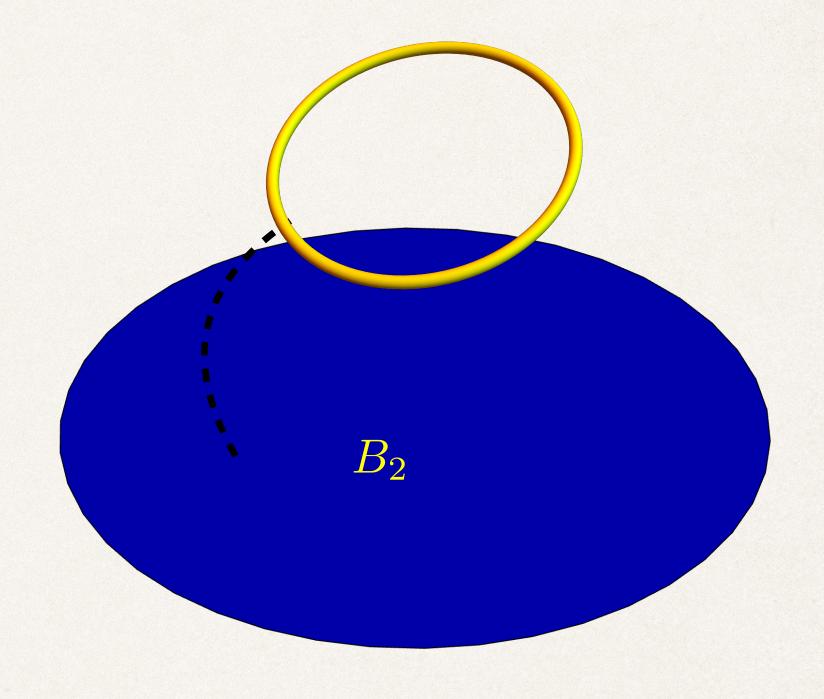
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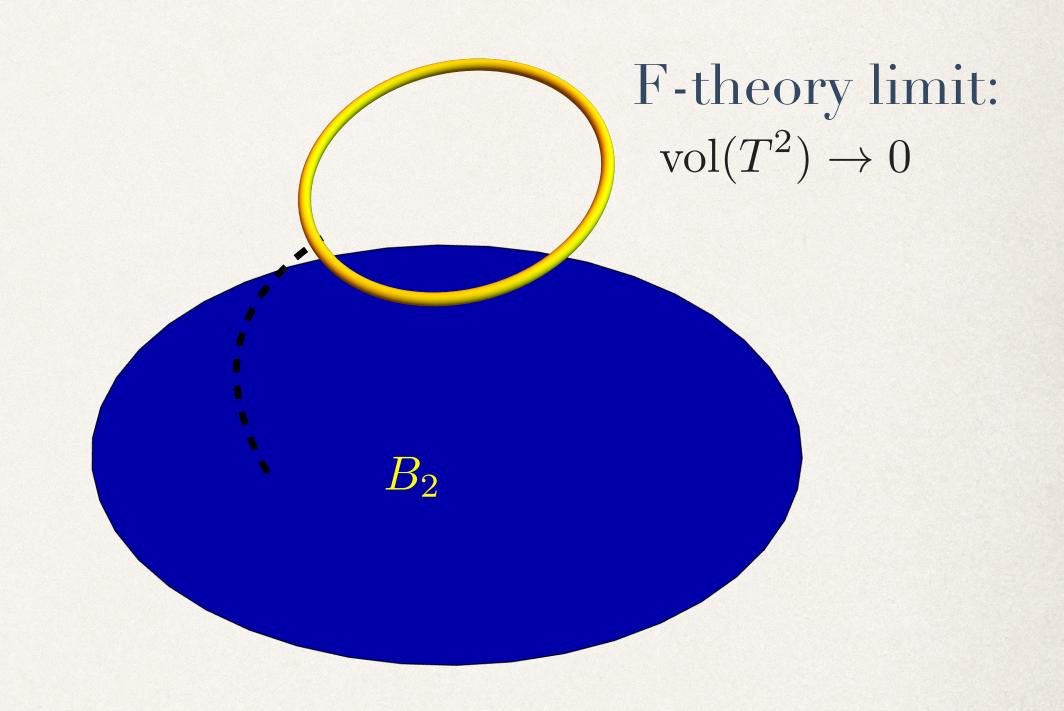
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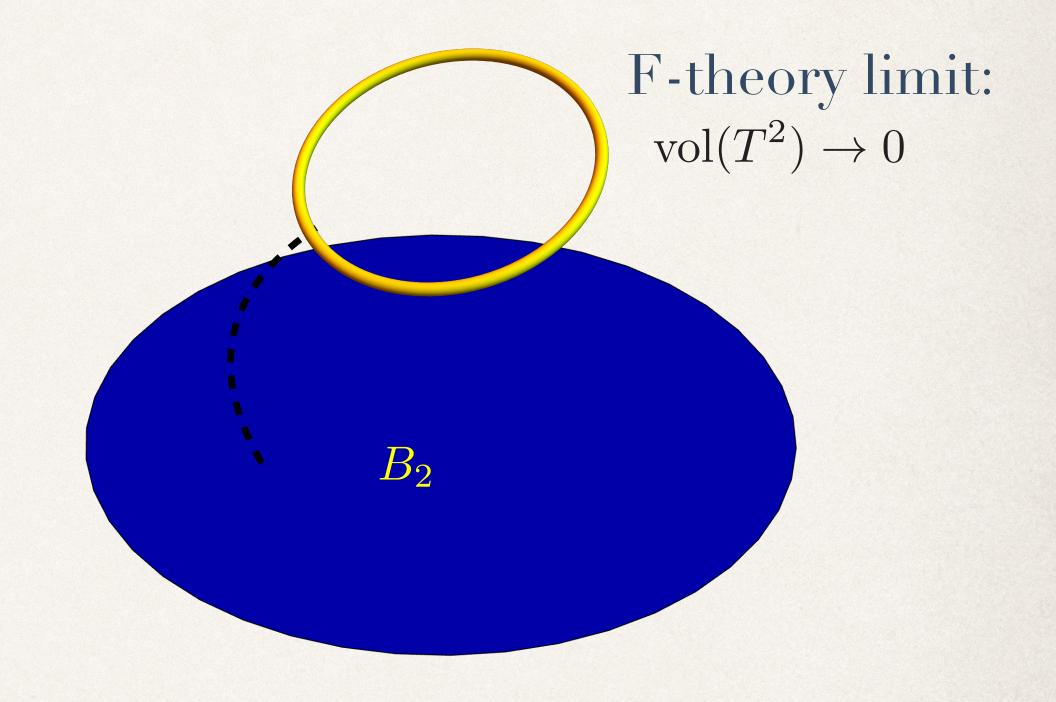
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[Oehlmann, Schimannek'19]



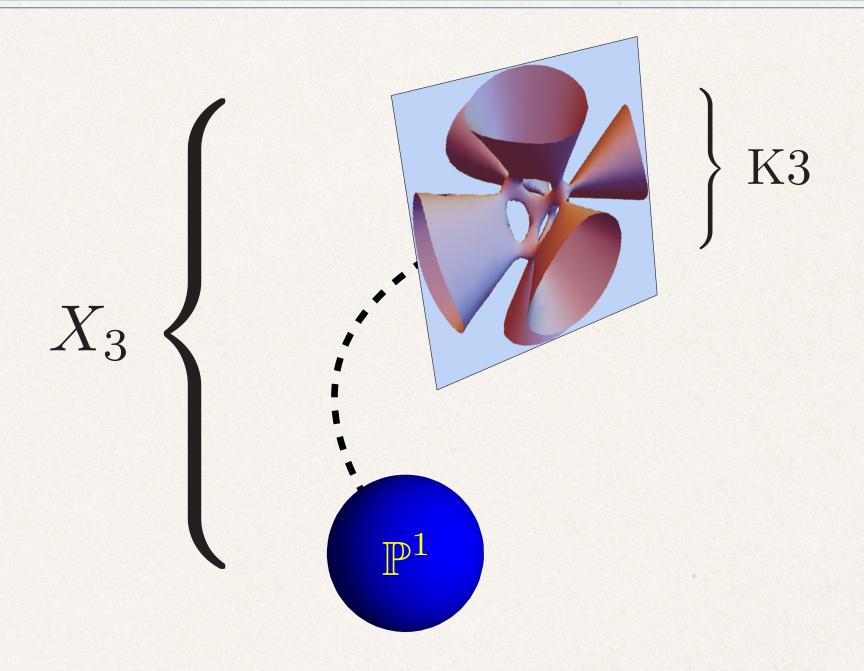
Super-extremal states are Kaluza-Klein modes of a six-dimensional theory realized by F-theory.

[Lee, Lerche, Weigand'19]

2. K3-fibered Calabi-Yau threefold $\rho: X_3 \to \mathbb{P}^1$. Here the U(1)_C gauge group is given by a curve C inside the generic K3-fiber.

Super-extremal states counted by Donaldson-Thomas invariants related to K3-fiber.

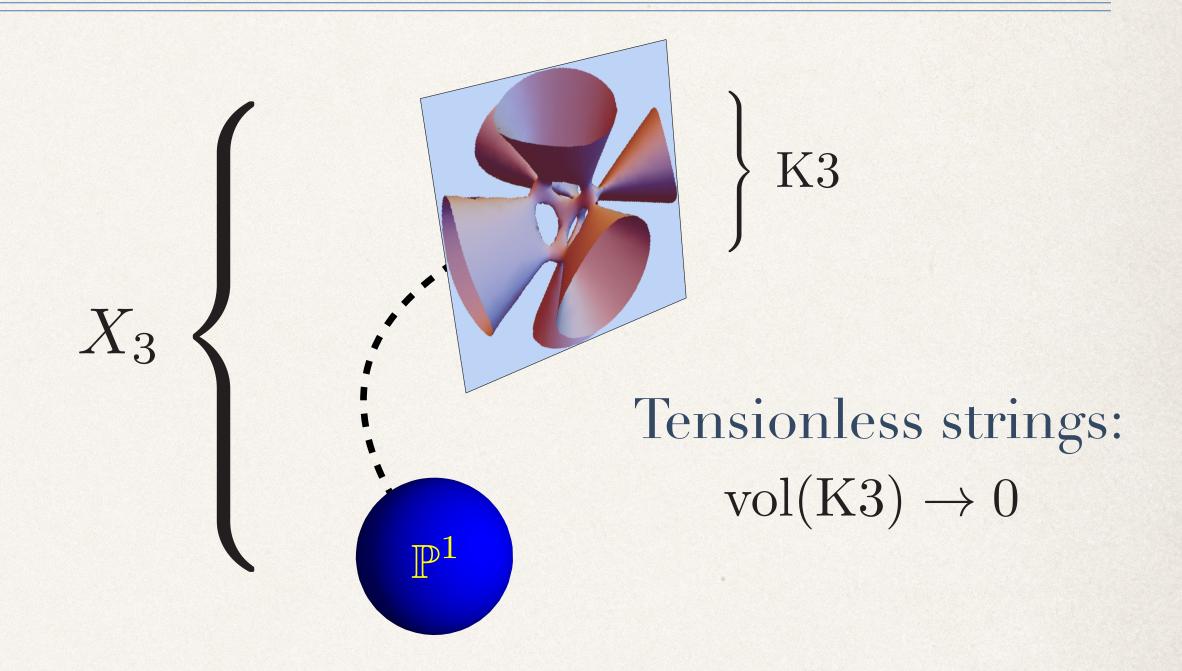
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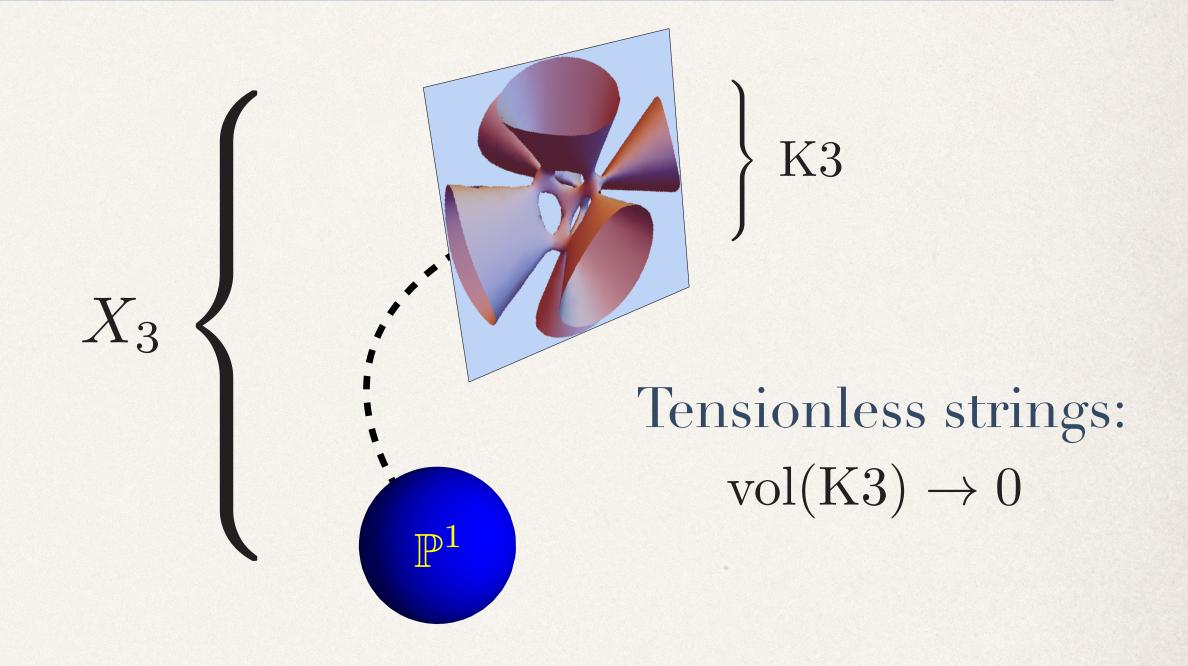
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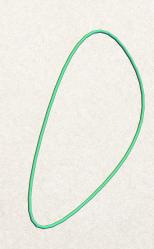
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Super-extremal states are excitations of heterotic strings in five dimensions.

[Lee, Lerche, Weigand'19]



Part II — B: The weak gravity conjecture with no infinte tower

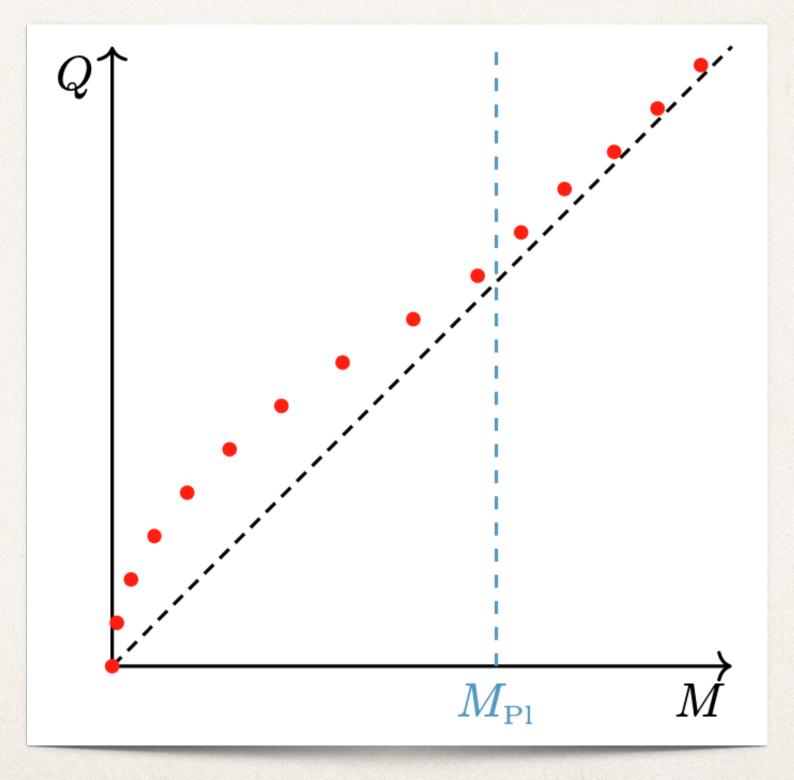
-Q: What about U(1) strongly coupled super-extremal states?

Philosophy: Not every U(1) gauge theory requires a tower of super-extremal states in the spectrum of the theory.

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Super-extremal states can be either particle-like states or black holes.

Extremal black boles: $g_{\mathrm{U}(1)}^2 q^2 M_{\mathrm{Pl}}^{D-2} = \gamma m^2$



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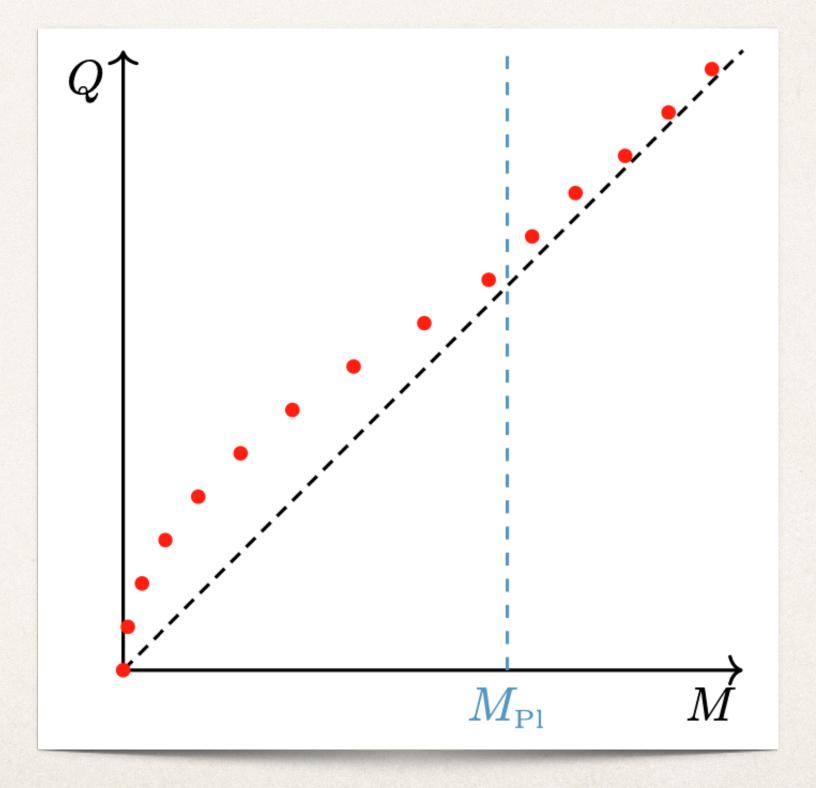
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Black holes can transition into particles when:

1.
$$g_{\mathrm{U}(1)}^2 M_{\mathrm{Pl}}^{D-2} \to 0$$
.

 $2. \gamma \rightarrow \infty$.



Conjecture: [CFC, Mininno, Weigand, Wiesner' to appear]

In a *D*-dimensional U(1) gauge theory coupled to gravity, there **can** exist a tower of superextremal states **if**:

- i) An infinite distance limit exists in the field moduli space, such that $g_{U(1)}^2 M_{\rm Pl}^{D-4} \to 0$.
- Emergent string conjecture limits.
- $|ii\rangle$ A finite distance limit exists in the field moduli space, such that $g_{U(1)}^2M_{\rm Pl}^{D-4}\to\infty$.
- Strongly coupled dynamic limits.

Examples when the tower of super-extremal states is NOT required.

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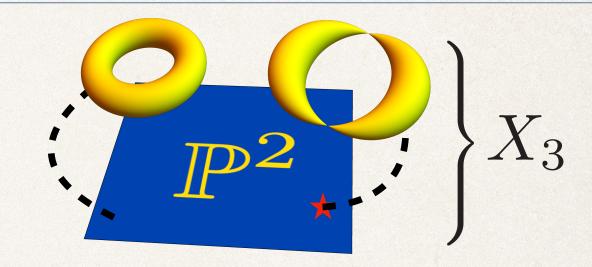
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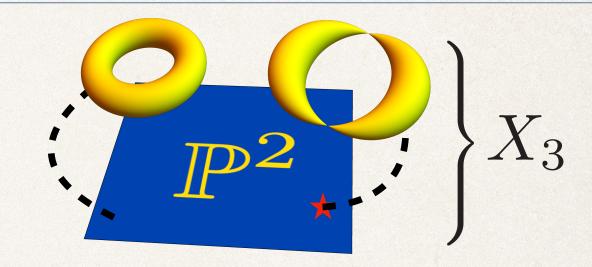
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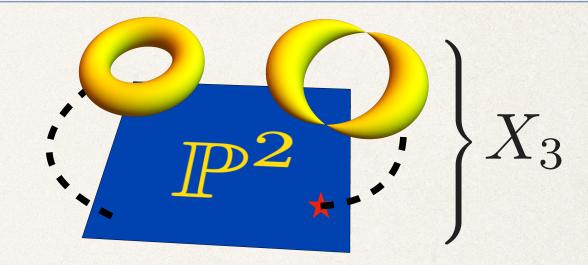
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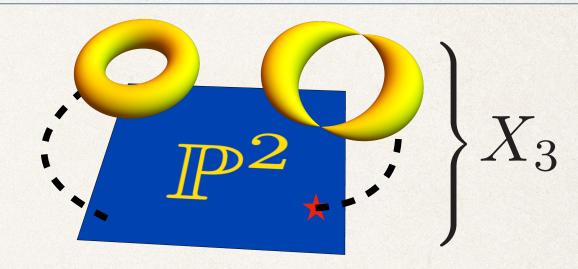
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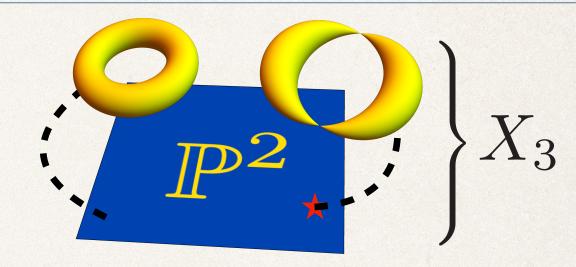
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 \Rightarrow no weak or strong coupling limit.

$$\oint \operatorname{vol}(T^2) = \frac{1}{r_{S^1}}$$
M-theory on $\pi: X_3 \to \mathbb{P}^2$.

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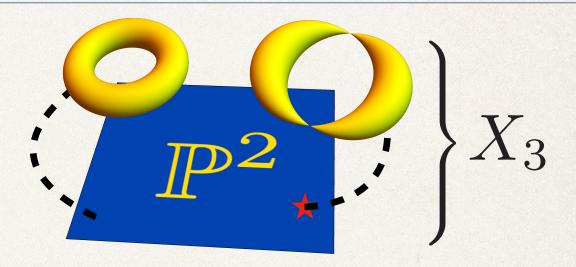
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 \Rightarrow no weak or strong coupling limit.

Upshot: we can find $\operatorname{vol}(T^2)|_{\max} \ge \operatorname{vol}(T^2)$, or $r_{\min} \le r_{S^1}$, for which EFT description breaks.

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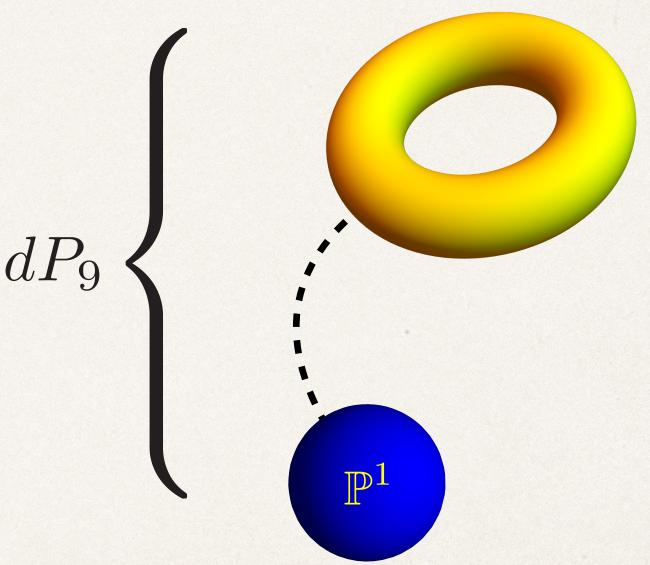
Upshot: we can find $\operatorname{vol}(T^2)|_{\max} \ge \operatorname{vol}(T^2)$, or $r_{\min} \le r_{S^1}$, for which EFT description breaks. The convex hull condition holds for all $r_{S^1} \ge r_{\min}$ via non-trivial charged massless matter.

Examples when the tower of super-extremal states is NOT required.

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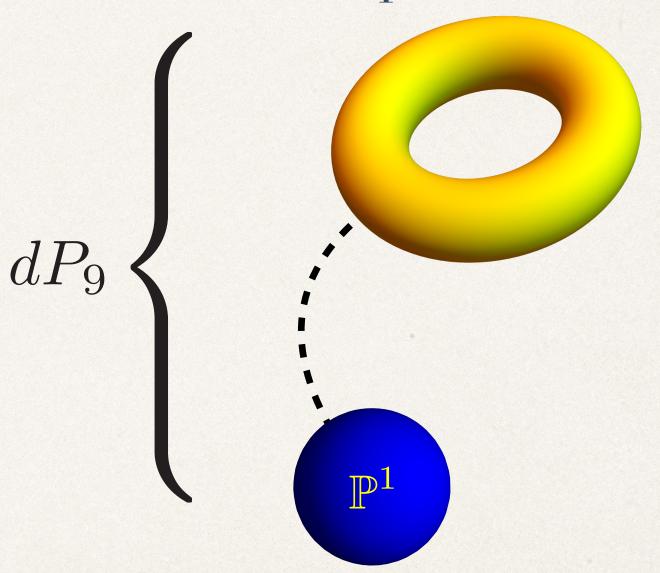
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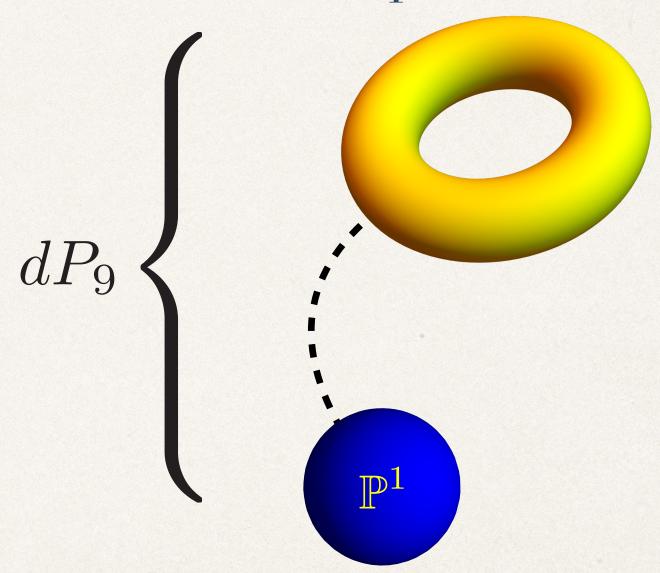
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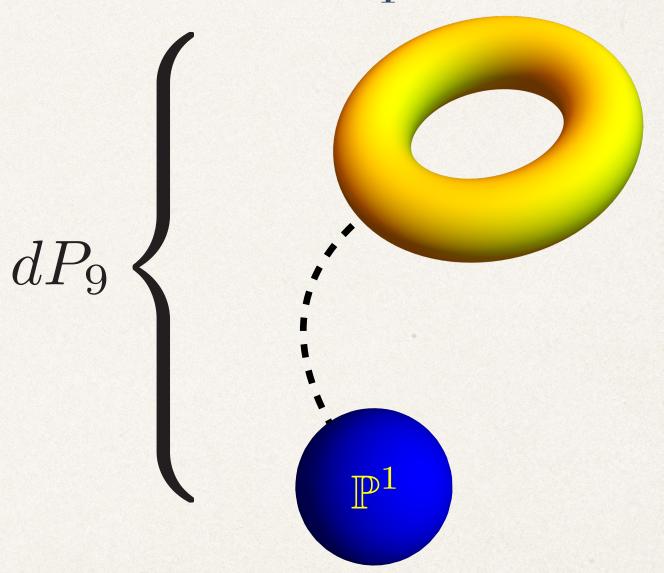
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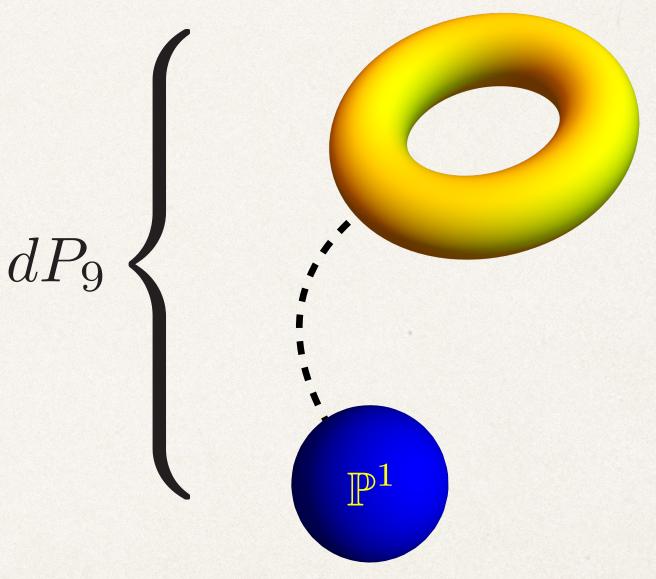
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Rational elliptic surface:



Upshot: Only a finite set of massive super-extremal states from [p,q]-string or E-strings.

Outlook

- Part I: We introduced the tower weak gravity conjecture as a constraint for the consistency of theories under Kaluza-Klein reduction.
- Part II A: In five-dimensional theories, we ask ourselves when is the tower of super-extremal states relevant for the EFT. For similar results for four-dimensional $\mathcal{N}=1$ theories, see [CFC, Mininno, Weigand, Wiesner '22x2].
- Part II B: Based on black hole arguments, we formulated a criterion to find when is the tower of super-extremal states necessary.

[CFC, Mininno, Weigand, Wiesner' to appear]