

# On the Boundary Conformal Field Theory Approach to Symmetry-Resolved Entanglement 2212.09767

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### Outline

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- 3 Symmetry-Resolved Entanglement Entropy
- 4 The BCFT Approach to Symmetry Resolution
- **5** Conclusion and Outlook

- Motivation

- 4 The BCFT Approach to Symmetry Resolution

# Symmetry Resolution-Main Idea

- Symmetries in quantum theory induce superselection rules
- The theory decomposes into sectors labeled by a fixed charge
- Conservation law constrains entanglement between two subregion
- Only sum of subregion charges is conserved
- Symmetry-resolved entanglement is entanglement in sector with fixed subregion charge

#### Motivation

- Contains more information than entanglement entropy
- Describes many body localization [Lukin et al.: 1805.09819]
- Understanding entanglement in AdS/CFT
- Applications to AQFT: Symmetry resolution of modular flow and modular operator [Di Giulio, Erdmenger: 2305.02343]

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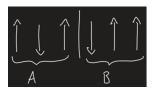
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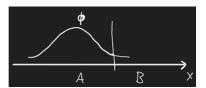
### Regularization

Definition of entanglement requires bipartition of Hilbert space

$$\mathscr{H}_{\text{tot}} = \mathscr{H}_{A} \otimes \mathscr{H}_{B} \tag{1}$$

- Evident for lattice theories
- Problematic in quantum field theory [Cardi, Tonni: 1608.01283]
- Regularization in field theory spirit requires boundary conditions



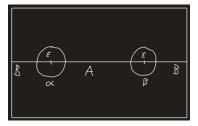


# The Cutting Operation

- Cutting out small circles around the entangling points provides regularization
- Boundary conditions  $\alpha, \beta$  specify the cutting operation [Ohmori, Tachikawa: 1406.4167]

$$\mathscr{H} \to \mathscr{H}_{\alpha,A,\beta} \otimes \mathscr{H}_{\beta,B,\alpha}$$
 (2)

 The factorized Hilbert space is different from other regularizations-only the lowest lying states coincide



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#### **Definitions**

- Consider CFT with additional Abelian symmetry group A
- Choose conformal boundary conditions that preserve at least parts of the symmetry
- Reduced density matrix decomposes into charge blocks

$$\rho_A = \bigoplus_{Q_A} P(Q_A) \rho_A(Q_A) \tag{3}$$

Symmetry-resolved entanglement entropy (SREE)

$$S_1(Q_A) = -\operatorname{tr}\left(\rho_A(Q_A)\log\rho_A(Q_A)\right) \tag{4}$$

#### Definition of SRRE

Define charged partition functions [Goldstein, Sela: 1711.09418]

$$Z_n(Q_A) = \operatorname{tr}\left(\rho_A^n \Pi_{Q_A}\right) \tag{5}$$

- Replica partition function of all states with fixed subregion charge
- Symmetry-resolved Rényi entropy (SRRE)

$$S_n(Q_A) = \frac{1}{1-n} \log \left\lfloor \frac{Z_n(Q_A)}{Z_1(Q_A)^n} \right\rfloor \tag{6}$$

- SREE can be recovered by taking  $n \to 1$  limit
- The key objects to compute are the charged partition functions

# Representation Theory

- The subregion Hilbert space will decompose into irreducible representations  $\mathcal{H}_{\mathcal{O}}$  of the symmetry group  $\mathcal{A}$
- The charged partition functions can be rewritten as

$$Z_n(Q) = \operatorname{tr}_{\mathcal{H}_Q}\left(q^{n\left(L_0 - \frac{c}{24}\right)}\right) \equiv \chi_Q(q^n) \tag{7}$$

- $\chi_O(q)$  are known in the literature as **characters**
- Boundary conditions determine which charges enter the spectrum and their multiplicities

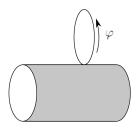
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### Example: Compact Free Boson

Bosonic field φ, action

$$S = \frac{g}{2} \int_{\Sigma} d^2 x \partial_{\mu} \varphi \partial^{\mu} \varphi, \qquad \varphi = \varphi + 2\pi R$$
 (8)

- Translation  $U(1) \otimes U(1)$  symmetry  $\varphi \rightarrow \varphi + a$
- ullet Conservation of momentum and winding quantum numbers m and w
- $\mathbb{Z}_2$  symmetry  $\varphi \to -\varphi$



# **Boundary Conditions**

- Separate boundary conditions need to be chosen at the two entangling points for the cutting operation
- Two possible boundary conditions that preserve conformal symmetry: Dirichlet (D) and Neumann (N)
- DD and NN boundary conditions preserve one conservation law of either wor m
- ND and DN break both conservation laws, only  $\mathbb{Z}_2$  symmetry remaining

#### Results

### DD and NN boundary conditions

- A single U(1) of the original two copies remains
- Symmetry-resolved Rényi entropy

$$S_n(Q) = \frac{1}{1-n} \log \left( \frac{\eta(q)^n}{\eta(q^n)} \right)$$
 (9)

- **Exact** equipartition to all orders in the cutoff
- Equipartition due to the exact form of the characters
- To leading order in the cutoff, familiar terms appear

$$S_1(Q) = \frac{1}{3} \log \frac{L}{\epsilon} - \frac{1}{2} \log \left( \frac{2}{\pi} \log \frac{L}{\epsilon} \right) - \frac{1}{2} + \mathcal{O}(\epsilon)$$
 (10)

#### Results

#### DN and ND boundary conditions:

- Both copies of U(1) are broken, resolve w.r.t. remaining  $\mathbb{Z}_2$
- SRRE agrees to results in the literature to lowest order in the cutoff

$$S_n(\pm) = \frac{1}{6} \frac{1+n}{n} \log \frac{L}{\epsilon} - \log 2 + \dots$$
 (11)

Equipartition broken at higher orders

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#### Conclusions

- Symmetry-resolved entanglement entropy calculates entanglement in a sector with fixed subregion charge
- SREE can be calculated directly by calculating characters
- Gives correct leading term, also correction to arbitrary order
- Equipartition exact to all orders for U(1) symmetry
- Mixed boundary conditions break U(1) symmetry, resolution w.r.t.  $\mathbb{Z}_2$ still possible

### Outlook

 Formalism resolves spacetime symmetries [Northe: 2303.07724]

### Open questions:

- non-Abelian symmetries
- Holographic interpretation of higher order terms
- Holographic implementation of boundary conditions
- Connection to Wilson line prescription