Aspects of Flat-Space Holography and Higher Spins

DESY Theory Workshop 2023

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1 Motivation: Farewell to AdS

- 2 Peculiarities of Flat-Space Holography
- 3 Higher Spins in Flat Space: Massless and Massive
- 4 Discussion: What's next?



Motivation: Farewell to AdS







Objective I: Quantum Gravity

- emergence of gravity
- unitarity of time evolution
- string-theoretic embedding



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- promising: d = 2 + 1
- dualities involving higher spin





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Towards Quantum Gravity:







Towards understanding Holography: (in 3d flat space)



absence of natural length scale

missing entries in holographic dictionary couple to massive fields





Lessons from AdS/CFT

Following results: three steps at once

go to flat space

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- allow higher spins
- introduce matter coupling

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Necessary ingredients:

- · definition of asymptotics
- · higher-spin symmetry algebra
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First-principle considerations, no $\Lambda \rightarrow 0$ limit!

couple to

massive fields



Peculiarities of Flat-Space Holography

Carrollian View on Flat Spacetimes





Carrollian View on Flat Spacetimes



Special Features:

- lightlike conformal boundary
- different boundary regions
- particular boundary conditions, asymptotic symmetry: BMS [Bondi/van der Burg/Metzner, 1962] [Sachs, 1962]
- dual field theory: Carrollian CFT [Lévy-Leblond, 1965]
 [Duval/Gibbons/Horvathy, 1403.4213]
- in 3d: Flat-Space Cosmologies [Prohazka/Salzer/Schöller, <u>1701.06573</u>]



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(alternatively: **Celestial Holography**) See also Charlotte's talk!



M Pannier

Asymptotically Flat Spacetimes

[Cornalba/Costa, hep-th/0203031; Barnich/Troessaert, 1001.1541]

Most General Asymptotically Flat Metric

 $ds^{2} = M(\phi)du^{2} - 2dudr + 2N(u,\phi)dud\phi + r^{2}d\phi^{2}$

Eddington-Finkelstein coordinates (u, r, ϕ)

- *u*: outgoing retarded time (light-like)
- *r*: radial coordinate ($r \rightarrow \infty$ is asymptotic direction)
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Chern-Simons Formulation

$$\begin{split} \omega &= \left(J_1 - \frac{\mathsf{M}(\phi)}{4}J_{-1}\right)\mathsf{d}\phi\,,\\ \mathsf{e} &= \left(\mathsf{P}_1 - \frac{\mathsf{M}(\phi)}{4}\mathsf{P}_{-1}\right)\mathsf{d}u + \frac{1}{2}\mathsf{P}_{-1}\mathsf{d}r + \left(\mathsf{r}\mathsf{P}_0 - \frac{\mathsf{N}(\mathsf{u},\phi)}{2}\mathsf{P}_{-1}\right)\mathsf{d}\phi \end{split}$$

[Afshar/Bagchi/Fareghbal/Grumiller/Rosseel, 1307.4768]

Isometries: Poincaré algebra iso $(2,1) \simeq isl(2,\mathbb{R}) \simeq sl(2,\mathbb{R}) \in \mathbb{R}^3$ $[J_m, J_n] = (m-n)J_{m+n}$ $[J_m, P_n] = (m-n)P_{m+n}$ $[P_m, P_n] = 0$

Higher Spins in Flat Space: Massless and Massive



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second-order Casimir elements:

$$\begin{split} \mathcal{M}^2 &= (\mathcal{P}_0)^2 - \mathcal{P}_1 \mathcal{P}_{-1} \,, \\ \mathcal{S} &= J_0 \mathcal{P}_0 - \frac{1}{2} \left(J_1 \mathcal{P}_{-1} + J_{-1} \mathcal{P}_1 \right) \end{split}$$



[Ammon/MP/Riegler, 2009.14210]

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(l = s - 1, s - 2)

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[Ammon/MP, 2211.12530]

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$$\mathsf{dC} + \left[\omega, \mathcal{C}\right]_{\star} + \mathbf{e} \star \mathcal{C} = 0$$

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Fierz-Pauli Equations

massive, spin- σ field:

 $(\Box - \mathcal{M}^2) \phi_{\mu_1 \dots \mu_\sigma} = 0$ $\nabla^{\mu} \phi_{\mu \mu_2 \dots \mu_\sigma} = 0$ $g^{\mu \nu} \phi_{\mu \nu \mu_3 \dots \mu_\sigma} = 0$





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Discussion: What's next?

Technical aspects:



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• include higher-spin interactions



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- connection to different unfolding approaches

[Ponomarev/Vasiliev, <u>1001.0062</u>] [Boulanger/Ponomarev/Sezgin/Sundell, <u>1412.8209</u>]



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- representation theory of $\mathfrak{iso}(2,1)$ and \mathfrak{bms}_3 (and higher-d)

[Ammon/Gray/Moran/MP/Wölfl, 2012.09173] [MP, 2023]



Thanks for your attention!

Boundary Conditions for Flat-Space Higher-Spin

$$\begin{split} \omega &= \left(J_1 - \frac{1}{4}\sum_{s=2}^{\infty} Z^{(s)}(\phi) J_{-1}(P_{-1})^{s-2}\right) d\phi \,, \\ \mathbf{e} &= \left(P_1 - \frac{1}{4}\sum_{s=2}^{\infty} Z^{(s)}(\phi) (P_{-1})^{s-1}\right) d\mathbf{u} + \frac{1}{2} P_{-1} d\mathbf{r} + \left(\mathbf{r} P_0 - \frac{1}{2}\sum_{s=2}^{\infty} W^{(s)}(\mathbf{u}, \phi) (P_{-1})^{s-1}\right) d\phi \end{split}$$

 $M(\phi) \equiv Z^{(2)}(\phi) ,$ $N(u, \phi) \equiv W^{(2)}(u, \phi)$

 $\partial_{\phi} Z^{(s)}(\phi) = 2 \partial_{u} W^{(s)}(u,\phi)$

