Quark-model puzzles and Lattice QCD : From heavy-light mesons to the $\Lambda(1405)$

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TECHNISCHE UNIVERSITÄT DARMSTADT

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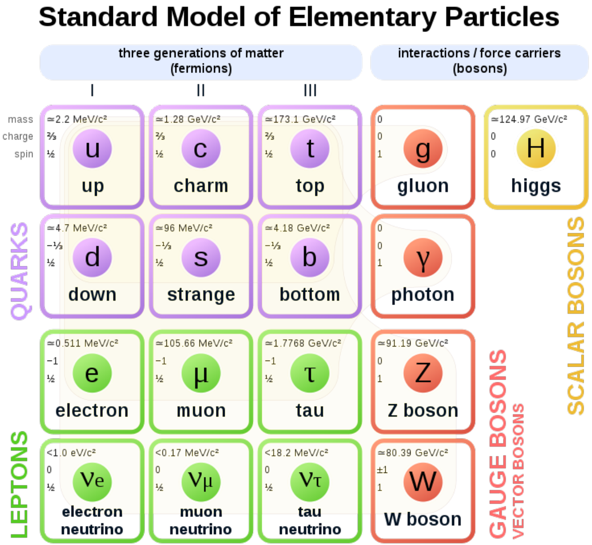
Outline

The Standard Model of Particle Physics - particle content



Quantum Chromodynamics (QCD) and the lattice

- QCD: Theory that describes the strong interaction between quarks and gluons within the Standard Model of Particle Physics

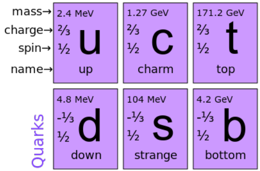
- Lattice QCD (Kenneth Wilson 1974) provides a method to calculate observables in the strong coupling regime of QCD

Important properties of QCD in different regimes:

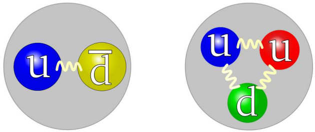
- *Asymptotic Freedom:*
  Theory perturbative at high energies/short distances
  $\rightarrow$ Perturbative calculations for high-energy physics
- *Confinement:* Color-charged particles do not exist individually but only confined into composite objects called hadrons.

Quarks, hadrons and confinement

- We do not observe free quarks and gluons but only color-neutral hadrons
- Textbook classification typically covers three-quark *baryons* and quark-antiquark *mesons*
- *Baryons:* qqq states in the quark model
  Examples: proton, neutron
- *Mesons:* $\bar{q}q$ states in the quark model
  Example: pion π^+
- More exotic objects like glueballs, hybrid mesons, four-quark states (tetraquarks or mesonic molecules), five-quark states (pentaquarks), . . .



$\longrightarrow$



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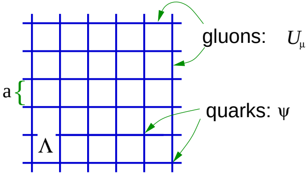
My method of choice: Lattice QCD

- Lattice QCD: Regularization of QCD by a 4-d Euclidean space-time lattice. Provides a calculational method.



Euclidean correlator of two Hilbert-space operators $\hat{O}_1$ and $\hat{O}_2$.

$$\left\langle \hat{O}_2(t)\hat{O}_1(0) \right\rangle = \sum_n e^{-t\Delta E_n} \langle 0|\hat{O}_2|n\rangle\langle n|\hat{O}_1|0\rangle$$
$$= \frac{1}{Z}\int \mathcal{D}[\psi,\bar{\psi},U] e^{-S_E} O_2[\psi,\bar{\psi},U] O_1[\psi,\bar{\psi},U]$$

- Path integral over the Euclidean action $S_{E,QCD}[\psi,\bar{\psi},U]$; (a sum over quantum fluctuations)
- Can be evaluated with *Markov Chain Monte Carlo* (using methods well established in statistical physics)

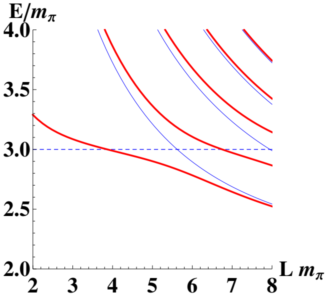
Progress from an old idea: Lüscher's finite-volume method

M. Lüscher Commun. Math. Phys. 105 (1986) 153;
Nucl. Phys. B 354 (1991) 531; Nucl. Phys. B 364 (1991) 237.

Basic observation: Finite-volume, multi-particle energies are shifted with regard to the free energy levels due to the interaction

$$E = E(p_1) + E(p_2) + \Delta_E$$

- Energy shifts encode scattering amplitude(s)
- Original method: Elastic scattering in the rest-frame in multiple spatial volumes L^3
- Coupled 2-hadron channels well understood
- $2 \leftrightarrow 1$ and $2 \leftrightarrow 2$ transitions well understood (example $\pi\pi \to \pi\gamma^*$)
- Significant progress for 3-particle scattering



Lattice QCD and quark-model puzzles

- Various kind of exotic/unconventional states (examples)
  - light scalar resonances (σ and κ)
  - $D^*_{s0}(2317)$, $D_{s1}(2460)$ and b-quark cousins
  - XYZ states
  - hybrid mesons
  - Roper resonance; $\Lambda(1405)$
  - Pentaquark states
  - Glueballs, ...

- Simple Lattice QCD calculations with $\bar{q}q$ and qqq interpolating fields struggle to make contact to experiment
  $\rightarrow$ Interpolating fields for multi-hadron states needed
- Need for all-to-all propagators (at least timeslice-to-timeslice) for meson-meson and meson-baryon scattering

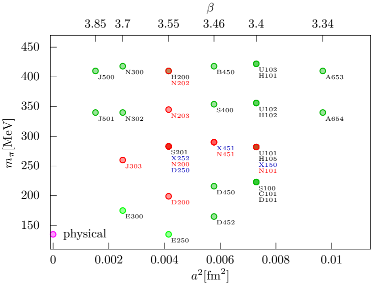
Challenges

- Hierarchy of difficulties
  - Meson systems are simpler than baryons
    (exponentially degrading signal to noise)
  - Cost of contractions/correlation functions much larger for systems with baryons
  - More complicated scattering amplitudes need many data points (volumes, frames)
    single two-hadron channel; coupled two-hadron channels; three-hadron scattering
- Hierarchy of projects:
  - Proof of principle
  - Explore quark mass dependence
  - Full spectroscopy calculation including continuum limit
  - Structure observables (transitions, form factors, . . .)
- Two examples:
  - Low-lying positive-parity heavy-light mesons
    Most systematics can be addressed
  - $\Lambda(1405)$ in coupled-channel $\pi\Sigma$–$\bar{K}N$-scattering
    Difficult but feasible with current methods

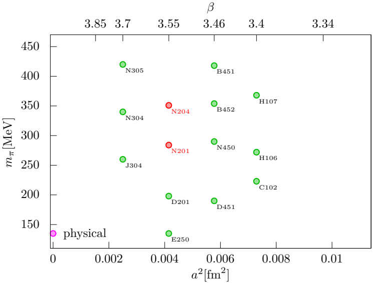
CLS gauge field ensembles

Bruno *et al.* JHEP 1502 043 (2015); Bali *et al.* PRD 94 074501 (2016)

$Tr(M) = \text{const.}$



$m_s = \text{const.}$



plot style by Jakob Simeth, RQCD

Important lattice systematics from

- Taking the *continuum limit*: $a(g, m) \to 0$
- Taking the *infinite volume limit*: $L \to \infty$
- Calculation at (or extrapolation to) physical quark masses

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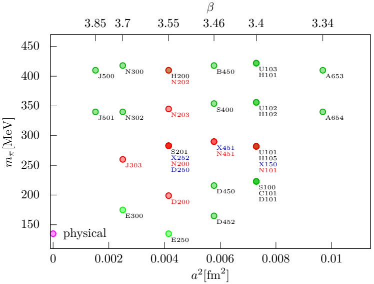

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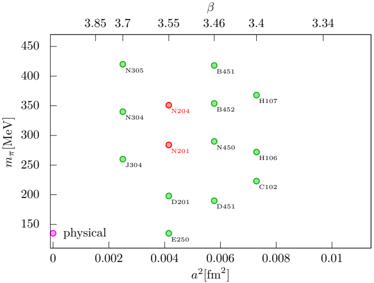


plot style by Jakob Simeth, RQCD

Important lattice systematics from

- Taking the *continuum limit*: $a(g, m) \to 0$
- Want to exploit (power law) finite volume effects (keeping exponential effects small)
- Calculation at (or extrapolation to) physical quark masses

Projects and People

- Heavy-quark exotics ($ud\bar{b}\bar{b}$ and $us\bar{b}\bar{b}$) and positive-parity heavy-light mesons

  `R.J. Hudspith, DM, PRD 107, 114510 (2023)`

  - TU Darmstadt/GSI: Jamie Hudspith, Daniel Mohler
- $\Lambda(1405)$ and meson-baryon scattering:
  - DESY Zeuthen $\rightarrow$ Bochum: John Bulava
  - BNL: Andrew Hanlon
  - Intel: Ben Hörz
  - North Carolina: Amy Nicholson, Joseph Moscoso
  - TU Darmstadt/GSI: Daniel Mohler, Barbara Cid Mora
  - CMU: Colin Morningstar, Sarah Skinner
  - MIT: Fernando Romero-López
  - LBNL: André Walker-Loud
- $\Lambda(1405)$ results are still preliminary, but close to publication

Outline

1. Introduction and Motivation
2. Positive-parity heavy-light hadrons
3. Coupled-channel scattering and the $\Lambda(1405)$
4. Conclusions and Outlook

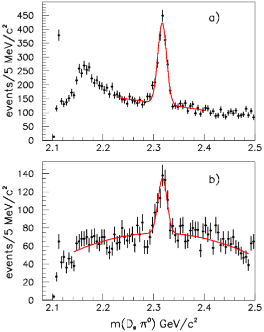
Exotic D_s and B_s candidates

Established s and p-wave hadrons:

D_s ($J^P = 0^-$) and D_s^* (1^-)
$D_{s0}^*(2317)$ (0^+), $D_{s1}(2460)$ (1^+),
$D_{s1}(2536)$ (1^+), $D_{s2}^*(2573)$ (2^+)

B_s ($J^P = 0^-$) and B_s^* (1^-)
?
$B_{s1}(5830)$ (1^+), $B_{s2}^*(5840)$ (2^+)

$D_{s0}^*(2317)$:
PRL 90 242001 (2003)



- Corresponding $D_0^*(2400)$ and $D_1(2430)$ are broad resonances
- Perceived peculiarity: $M_{c\bar{s}} \approx M_{c\bar{d}}$ (an old dispute; likely not the case)
- Additional exotic states are expected (in the sextet representation)

See for example Kolomeitsev, Lutz, PLB 582, 39 (2004)

- B_s cousins of the $D_{s0}^*(2317)$ and $D_{s1}(2460)$ not (yet) seen in experiment

$D_{s0}^*(2317)$: D-meson – Kaon s-wave scattering

M. Lüscher Commun. Math. Phys. 105 (1986) 153;
Nucl. Phys. B 354 (1991) 531; Nucl. Phys. B 364 (1991) 237.

Charm-light hadrons

$D_{s0}^*(2317)^\pm$ $I(J^P) = 0(0^+)$
J, P need confirmation.

J^P is natural, low mass consistent with 0^+.

Mass $m = 2317.7 \pm 0.6$ MeV
$m_{D_{s0}^*(2317)^\pm} - m_{D_s^\pm} = 349.4 \pm 0.6$ MeV
Full width $\Gamma < 3.8$ MeV, CL = 95%

$$p \cot \delta_0(p) = \frac{2}{\sqrt{\pi} L} Z_{00}\left(1; \left(\frac{L}{2\pi} p\right)^2\right)$$

$$\approx \frac{1}{a_0} + \frac{1}{2} r_0 p^2$$

$D_{s0}^{*}(2317)$: D-meson – Kaon s-wave scattering

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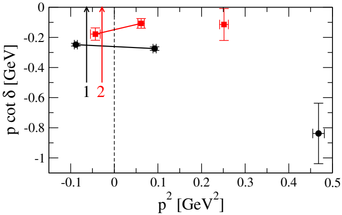
DM *et al.* PRL 111 222001 (2013)
Lang, DM *et al.* PRD 90 034510 (2014)

Results for ensembles (1) and (2)

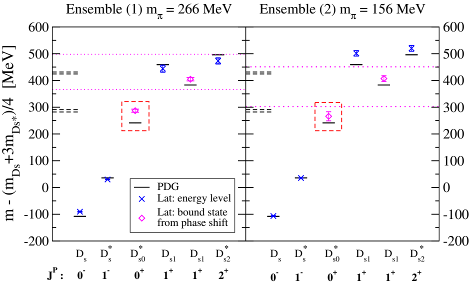


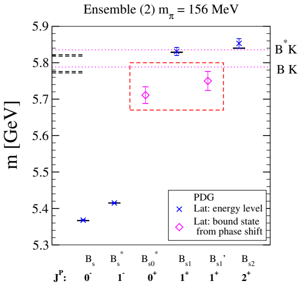
$$a_0 = -0.756 \pm 0.025\text{fm} \qquad (1)$$
$$r_0 = -0.056 \pm 0.031\text{fm}$$
$$a_0 = -1.33 \pm 0.20\text{fm} \qquad (2)$$
$$r_0 = 0.27 \pm 0.17\text{fm}$$

Positive-parity states in the D_s and B_s spectrum

DM *et al.* PRL 111 222001 (2013)
Lang, DM *et al.* PRD 90 034510 (2014)

Lang, DM, Prelovsek, Woloshyn PLB 750 17 (2015)



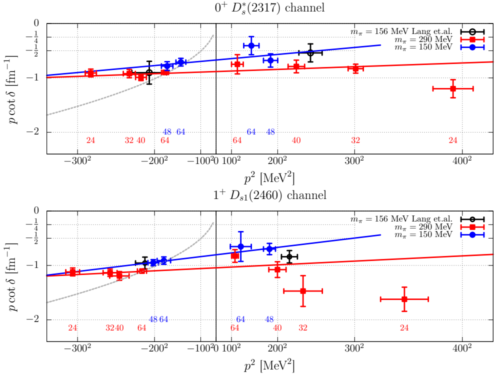


- Spectrum reliably extracted and agrees qualitatively with experiment
- Uncontrolled systematics sizable for the D_s states

- Full uncertainty estimate only for magenta B_s states
- Prediction of exotic states from Lattice QCD!

D_s results in multiple volumes from RQCD

Bali, Collins, Cox, Schäfer, PRD 96 074501 (2017)



- Study with different volumes at pion masses of 150, 290 MeV
- Results confirm basic behavior seen in a single volume
- Discretization effects remain unexplored

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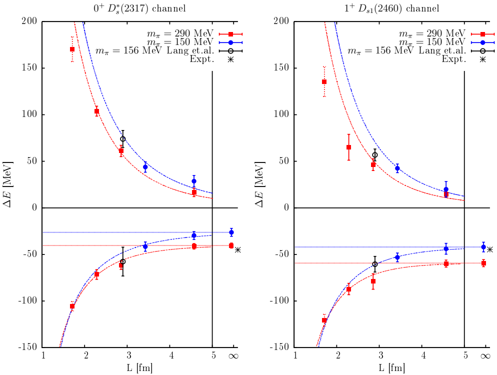

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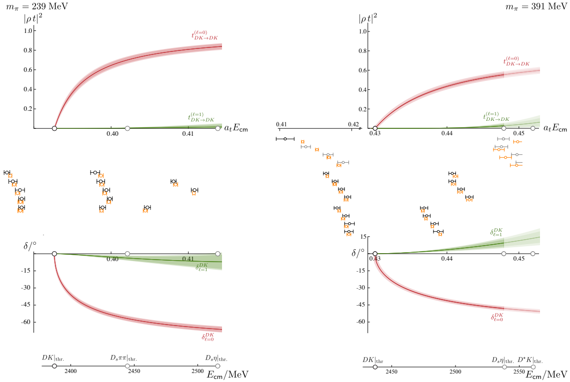
DK and $D\bar{K}$ scattering and the $D_{s0}^*(2317)$

Hadron spectrum collaboration, Cheung *et al.* JHEP 02 100 (2021)



- Study uses moving frames in addition
  results in large number of energy levels at $m_\pi = 238, 391$ MeV

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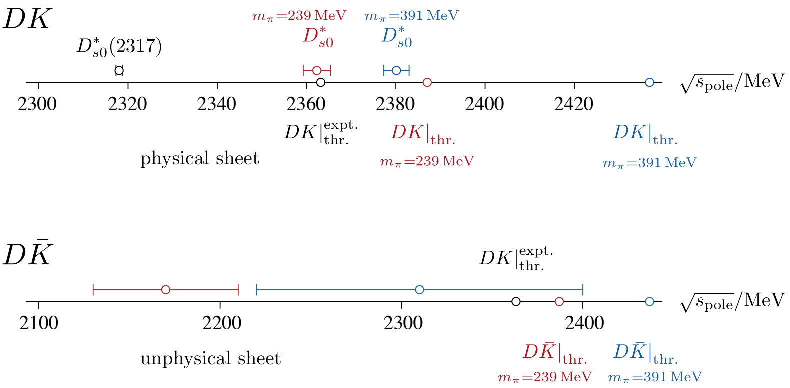


- Study uses moving frames in addition results in large number of energy levels at $m_\pi = 238, 391$ MeV

CLS ensembles used for heavy-light mesons

R.J. Hudspith, DM, PRD 107, 114510 (2023)

Ensemble	Mass trajectory	$L^3 \times L_T$	$N_{\text{Conf}} \times N_{\text{Prop}}$
U103	$\text{Tr}[M] = C$	$24^3 \times 128$	1000×23
H101	$\text{Tr}[M] = C$	$32^3 \times 96$	500×12
U102	$\text{Tr}[M] = C$	$24^3 \times 128$	732×18
H102	$\text{Tr}[M] = C$	$32^3 \times 96$	500×16
U101	$\text{Tr}[M] = C$	$24^3 \times 128$	600×18
H105	$\text{Tr}[M] = C$	$32^3 \times 96$	500×16
N101	$\text{Tr}[M] = C$	$48^3 \times 128$	537×18
C101	$\text{Tr}[M] = C$	$48^3 \times 96$	400×16
H107	$\widetilde{m_s} = \widetilde{m_s}^{\text{Phys.}}$	$32^3 \times 96$	500×16
H106	$\widetilde{m_s} = \widetilde{m_s}^{\text{Phys.}}$	$32^3 \times 96$	500×16
H200	$\text{Tr}[M] = C$	$32^3 \times 96$	500×28

NRQCD action

Typical tadpole-improved NRQCD action (here we will use n=4)

Lepage *et al.*, PRD 46, 4052-4067 (1992)

$$
\begin{aligned}
H_0 =& -\frac{1}{2aM_0}\Delta^2, \\
H_I =& \left(-c_1\frac{1}{8(aM_0)^2} - c_6\frac{1}{16n(aM_0)^2}\right)(\Delta^2)^2 + c_2\frac{i}{8(aM_0)^2}(\tilde{\Delta}\cdot\tilde{E} - \tilde{E}\cdot\tilde{\Delta}) + c_5\frac{\Delta^4}{24(aM_0)} \\
H_D =& -c_3\frac{1}{8(aM_0)^2}\sigma\cdot\left(\tilde{\Delta}\times\tilde{E} - \tilde{E}\times\tilde{\Delta}\right) - c_4\frac{1}{8(aM_0)}\sigma\cdot\tilde{B} \\
\delta H =& H_I + H_D.
\end{aligned}
$$

Propagators generated through symmetric evolution equation

$$
G(x,t+1) = \left(1-\frac{\delta H}{2}\right)\left(1-\frac{H_0}{2n}\right)^n \tilde{U}_t(x,t_0)^\dagger \left(1-\frac{H_0}{2n}\right)^n \left(1-\frac{\delta H}{2}\right) G(x,t).
$$

- We also tune a $\mathcal{O}(v^6)$ action with tree-level coefficients for the higher order terms

Neural net NRQCD tuning and setup

R.J. Hudspith, DM, PRD 106, 034508 (2022)
R.J. Hudspith, DM, PRD 107, 114510 (2023)

- Calculate runs with a random distribution for the action parameters
- Let the neural network make parameter predictions
- Due to additive mass we must only consider splittings $\rightarrow$ we subtract the η_B from all states
- Perform tuning at $SU(3)_f$-symmetric point
- Gauge-fixed wall sources
- Tuning precision is about 1%

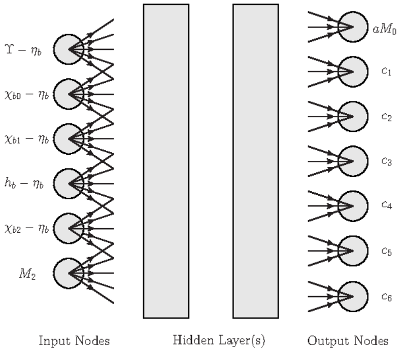


Figure: Schematic picture of our NRQCD setup

Input used for the tuning

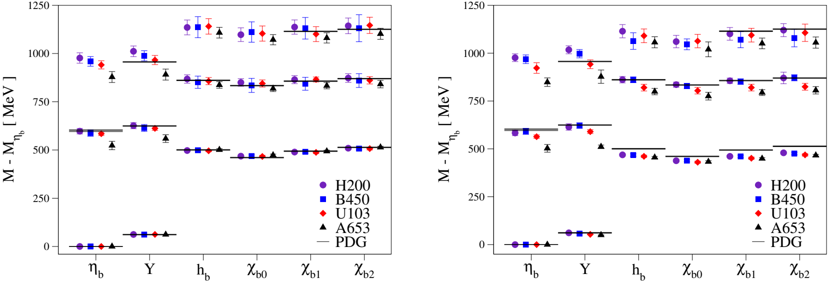
Consider only quark-line connected parts of simple meson operators

$$O(x) = (\bar{b}\Gamma(x)b)(x),$$

State	PDG mass [GeV]	$\Gamma(x)$
$\eta_b(1S)$	9.3987(20)	γ_5
$\Upsilon(1S)$	9.4603(3)	γ_i
$\chi_{b0}(1P)$	9.8594(5)	$\sigma \cdot \Delta$
$\chi_{b1}(1P)$	9.8928(4)	$\sigma_j \Delta_i - \sigma_i \Delta_j \; (i \neq j)$
$\chi_{b2}(1P)$	9.9122(4)	$\sigma_j \Delta_i + \sigma_i \Delta_j \; (i \neq j)$
$h_b(1P)$	9.8993(8)	Δ_i

Table: Table of lattice operators used and their continuum analogs.

NRQCD Neural Net Tuning:
Stable s- and p-wave bottomonia



- Higher S- and P-wave states serve as a check whether our tuning leads to reasonable results
- Main results from the lattice spacing of U103; H200 used to estimate systematics

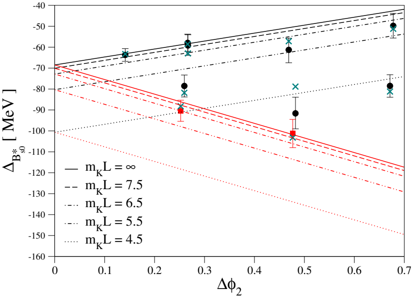
B_s: Chiral – infinite volume extrapolation

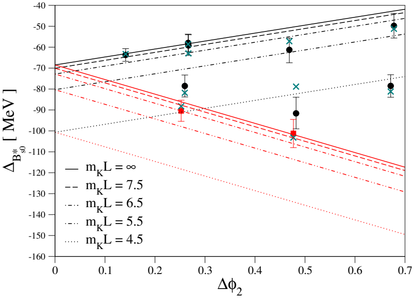
- We explore the previously predicted $J^P = 0^+$ and 1^+ bound states
- Mainly the CLS $\mathrm{Tr}M$ = const trajectory and $2\,m_S$ = const ensembles

Combined extrapolation:

$$\Delta_{B^*_{s0}/B_{s1}}(\Delta\phi_2, m_K L, a) = \Delta_{B^*_{s0}/B_{s1}}(0,\infty,a)\left(1 + A\Delta\phi_2 + Be^{-m_K L}\right)$$

$$\Delta\phi_2 = \phi_2^{\mathrm{Lat}} - \phi_2^{\mathrm{Phys}} \qquad ; \qquad \phi_2 = 8t_0 m_\pi^2$$





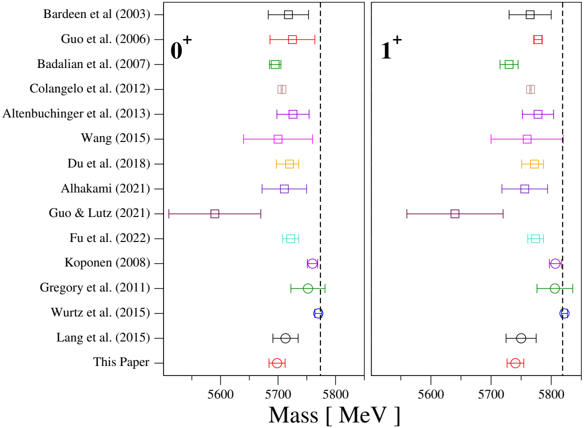
Systematic uncertainties and final result

Resulting binding energies:

$$\Delta_{B^*_{s0}}(0,\infty,0) = -75.4(3.0)_{\text{Stat.}}(13.7)_{\text{a}}\ [\text{MeV}],$$
$$\Delta_{B_{s1}}(0,\infty,0) = -78.7(3.7)_{\text{Stat.}}(13.4)_{\text{a}}\ [\text{MeV}].$$

- Small uncertainty from statistics + combined extrapolation
- Largest systematics from usage of NRQCD/discretization effects
- Central value shifted by applying half the mass difference between H200 and U103
- All other explored uncertainties (finite volume shapes, modified quark-mass dependence, etc.) small

Comparison to the literature



- Results agree well with models based on unitarized χPT
- Improved uncertainty estimate over older Lattice calculations

Outline

An old puzzle: $\Lambda(1405)$, $J^P = \frac{1}{2}^-$

- PDG (4 star resonance)

$$M_\Lambda = 1405^{+1.3}_{-1.0} MeV \qquad \Gamma_\Lambda = 50.5 \pm 2.0$$

(Some) quark models struggled to accommodate this state.

- However
  - Unitarized χPT + Model input yields 2 poles with $\Re \approx 1400$ MeV
    $\rightarrow$ Now new PDG state
  - CLAS observes different line shapes for $\Sigma^-\pi^+$, $\Sigma^+\pi^-$ and $\Sigma^0\pi^0$
    Interference between $I = 0$ and $I = 1$ amplitudes is the likely reason
  - Even the $\Sigma^0\pi^0$ is badly described by a single Breit-Wigner
  - CLAS data consistent with popular 2-pole picture
  - No satisfactory lattice results (although claims exist)
- Relevant channels: $\Sigma\pi$, $N\bar{K}$ (and maybe $\Lambda\eta$); simulation in isospin limit
- Goal: Explore coupled channel problem and extract scattering amplitudes from the low-lying energy spectrum

$\Lambda(1405)$ – Experimental developments

- Angular analysis of the process $\gamma + p \to K^+ + \Sigma + \pi$ by CLAS strongly favors the assignment of quantum numbers $J^P = \frac{1}{2}^-$

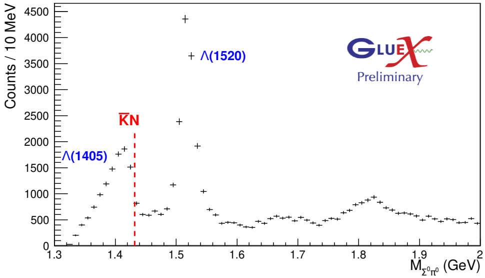
  Moriya *et al.*, PRC 87 035206 (2013)

- $K^- p$ scattering length determined by the SIDDHARTHA collaboration

  Bazzi *et al.*,PLB 704 (2011) 113

- A glimpse of the future: Preliminary analysis at GlueX

  Wickramaarachchi *et al.*, arXiv:2209.06230



Excited state energies and the variational method

Matrix of correlators projected to fixed momentum (will assume 0)

$$C(t)_{ij} = \sum_n \mathrm{e}^{-tE_n} \langle 0|O_i|n\rangle \left\langle n|O_j^\dagger|0\right\rangle$$

Solve the generalized eigenvalue problem:

$$C(t)\vec{\psi}^{(k)} = \lambda^{(k)}(t)C(t_0)\vec{\psi}^{(k)}$$
$$\lambda^{(k)}(t) \propto \mathrm{e}^{-tE_k}\left(1 + \mathcal{O}\left(\mathrm{e}^{-t\Delta E_k}\right)\right)$$

At large time separation: only a single state in each eigenvalue.
Eigenvectors can serve as a fingerprint.

`Michael Nucl. Phys. B259, 58 (1985)`

`Lüscher and Wolff Nucl. Phys. B339, 222 (1990)`

`Blossier et al. JHEP 04, 094 (2009)`

The "Distillation" method

Peardon et al. PRD 80, 054506 (2009)
Morningstar et al. PRD 83, 114505 (2011)

- Idea: Construct separable quark smearing operator using low modes of the 3D lattice Laplacian
  Spectral decomposition for an $N \times N$ matrix:

$$f(A) = \sum_{k=1}^{N} f(\lambda^{(k)})\, v^{(k)} v^{(k)\dagger}.$$

  With $f(\nabla^2) = \Theta(\sigma_s^2 + \nabla^2)$ (Laplacian-Heaviside (LapH) smearing):

$$q_s \equiv \sum_{k=1}^{N} \Theta(\sigma_s^2 + \lambda^{(k)}) v^{(k)} v^{(k)\dagger}\, q \;=\; \sum_{k=1}^{N_v} v^{(k)} v^{(k)\dagger}\, q\,.$$

- Advantages: momentum projection at source; large interpolator freedom, small storage
- Disadvantages: expensive; unfavorable volume scaling
- Stochastic approach (partly) eliminates bad volume scaling

Ensemble and group theory

Current data on CLS Ensemble D200

a [fm]	$T \times L^3$	m_π [MeV]	m_K [MeV]	$m_\pi L$	N_{cnfg}
0.0633(4)(6)	128×64^3	200	480	4.3	2000

Lattice irreducible representations for a given J^P

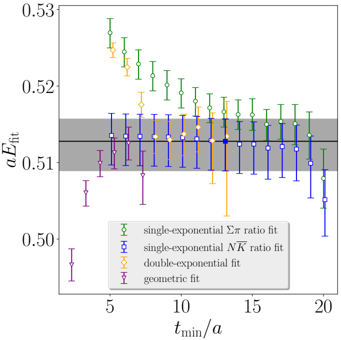
see Morningstar *et al.* arXiv:1303.6816

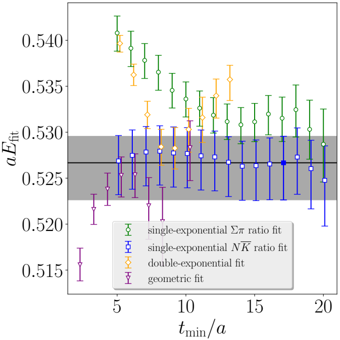
J^P	[000]	[00n]	[0nn]	[nnn]	
$\frac{1}{2}^+$	G_{1g}	G_1	G	G	Λ, $\Lambda(1600)$
$\frac{1}{2}^-$	G_{1u}	G_1	G	G	$\Lambda(1405)$, $\Lambda(1670)$
$\frac{3}{2}^+$	H_g	G_1, G_2	$2G$	F_1, F_2, G	$\Lambda(1690)$
$\frac{3}{2}^-$	H_u	G_1, G_2	$2G$	F_1, F_2, G	$\Lambda(1520)$, $\Lambda(1690)$

Specific setup on D200

- Combined basis of simple 3-quark structures and 2 hadron interpolators with the lowest few momentum combinations in each irrep
- Distillation setup:
  - $n_{ev} = 448$ eigenmodes of the Lattice Laplacian
  - Quark lines connecting source and sink:
    Noise dilution scheme with $(TF, SF, LI16)$ and 6 noises
  - Lines starting end ending on the same time slice:
    Noise dilution scheme with $(TI8, SF, LI16)$ and 2 noises
  - Four source time slices
  - Lattice Laplacian constructed on stout smeared links with $(\rho, n) = (0.1, 36)$

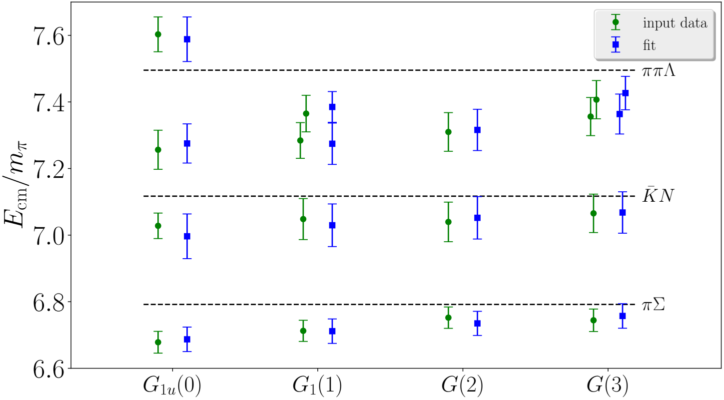
Extracting the spectrum (examples)





- We used various methods/cross checks
- Geometric series fit: $C(t) = \frac{A\mathrm{e}^{-E_0 t}}{1 - B\mathrm{e}^{-\Delta E t}}$
- Two students with two slightly different analysis methods

Finite-volume spectra



- Amplitude analysis uses ratios to extract energy differences with regard to non-interacting levels
- Blue squares indicate results from our preferred amplitude fit

Families of simple parameterizations

1. An ERE in the K matrix:

$$\frac{E_{\text{cm}}}{M_\pi}\tilde{K}_{ij} = A_{ij} + B_{ij}\Delta_{\pi\Sigma}, \qquad (3)$$

2. ERE in the K matrix without factors of E_{cm}

$$\tilde{K}_{ij} = \widehat{A}_{ij} + \widehat{B}_{ij}\Delta_{\pi\Sigma}, \qquad (4)$$

3. ERE in the inverse K matrix:

$$(\tilde{K}^{-1})_{ij} = \frac{E_{\text{cm}}}{M_\pi}\left(\widetilde{A}_{ij} + \widetilde{B}_{ij}\Delta_{\pi\Sigma}\right), \qquad (5)$$

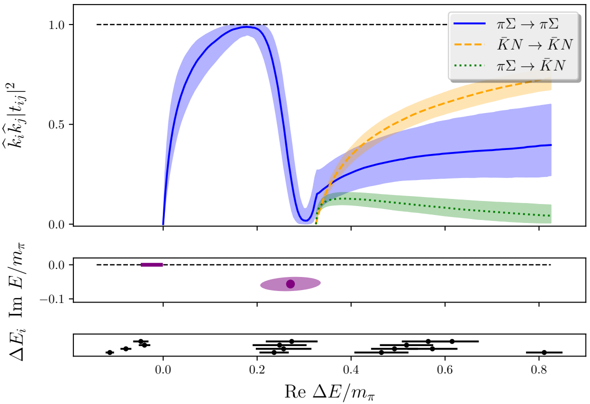
4. Blatt-Biederharn parametrization:

$$\tilde{K}_{ij} = \begin{pmatrix} \cos\epsilon & \sin\epsilon \\ -\sin\epsilon & \cos\epsilon \end{pmatrix} \begin{pmatrix} f_0(E_{\text{cm}}) & 0 \\ 0 & f_1(E_{\text{cm}}) \end{pmatrix} \begin{pmatrix} \cos\epsilon & -\sin\epsilon \\ \sin\epsilon & \cos\epsilon \end{pmatrix}$$

$$f_i(E_{\text{cm}}) = \frac{M_\pi}{E_{\text{cm}}}\frac{a_i + b_i\Delta_{\pi\Sigma}}{1 + c_i\Delta_{\pi\Sigma}}.$$

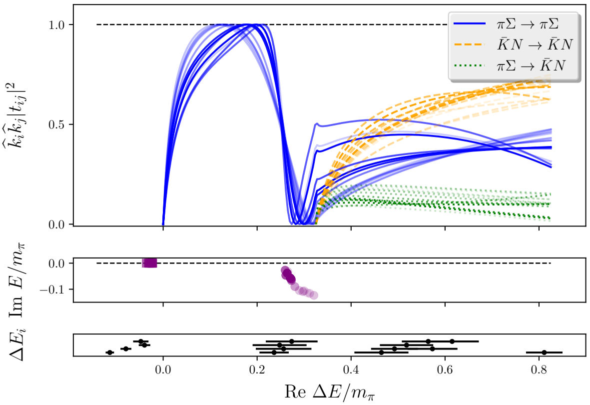
$\Delta_{\pi\Sigma}$ measures the distance from the $\pi\Sigma$ threshold

Our preferred amplitude and resulting poles



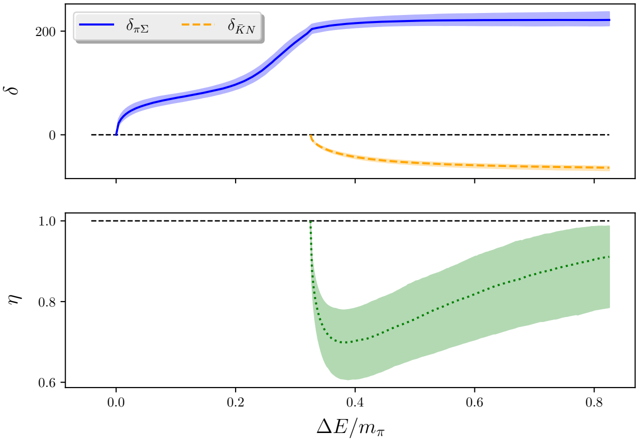
- Amplitudes evaluated with Akaike Information Criterion: $AIC = \chi^2 - 2\text{dof}$
- Sub-threshold levels pose strong constraints on the amplitude
- Limited data and therefore limited possibility to vary parameterizations

Some Variations of the used amplitude



- Results from varying parameterization/ omitting highest data point
- Amplitudes agnostic to the number of poles lead all yield 2 poles
- We also explored simple constraints for higher partial waves (negligible effect in range used)

Same thing different: Phases and inelasticity



- Alternative way of showing our results: 2 phases and inelasticity η

Pole positions and expectations from the literature

- Poles labeled as $(\pm, \pm)$
  depending on the signs of the imagianary part of $(k_{KN}, k_{\pi\Sigma})$
- Two poles are found on the $(-, +)$ sheet,
  the closest to physical scattering between the thresholds
- Our current result for the poles is

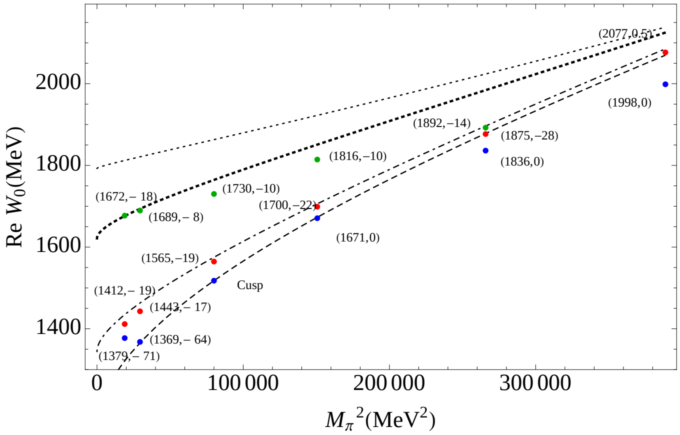
$$\begin{aligned} \text{Pole II} \quad & 1395(9)_{\text{stat}}(2)_{\text{model}}(16)_{\text{a}}\ \text{MeV} \\ \text{Pole I} \quad & 1456(14)_{\text{stat}}(2)_{\text{model}}(16)_{\text{a}}\ \text{MeV} \\ & - i \times 11.7(4.3)_{\text{stat}}(4)_{\text{model}}(0.1)_{\text{a}}\ \text{MeV} \end{aligned}$$

- Examples from the PDG review

approach	pole 1 [MeV]	pole 2 [MeV]
Refs. [14, 15], NLO	$1424^{+7}_{-23} - i\ 26^{+3}_{-14}$	$1381^{+18}_{-6} - i\ 81^{+19}_{-8}$
Ref. [17], Fit II	$1421^{+3}_{-2} - i\ 19^{+8}_{-5}$	$1388^{+9}_{-9} - i\ 114^{+24}_{-25}$
Ref. [18], solution #2	$1434^{+2}_{-2} - i\ 10^{+2}_{-1}$	$1330^{+4}_{-5} - i\ 56^{+17}_{-11}$
Ref. [18], solution #4	$1429^{+8}_{-7} - i\ 12^{+2}_{-3}$	$1325^{+15}_{-15} - i\ 90^{+12}_{-18}$

Expected quark-mass dependence

Molina, Döring, PRD 94 056010 (2016)



- Plots shows expected behavior for PACS-CS ensembles
- Qualitative agreement with regard to expected behavior

Outline

Conclusions and Outlook

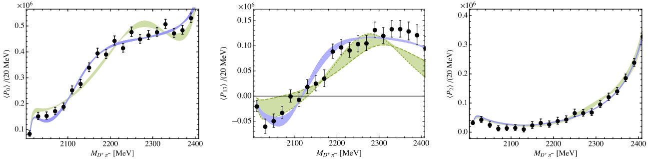
- Positive-parity heavy-light mesons
  - Presented NRQCD calculation with a full uncertainty estimate for B_0^* and B_{s1} mesons
    $\rightarrow$ refined predictions for LHCb, BelleII
  - Calculation could be further improved with RHQ action
  - Scattering amplitudes for the $D_{s0}^*(2317)$ and D_{s1} states using RHQ action planned
- Coupled-channel scattering and the $\Lambda(1405)$
  - First coupled-channel LQCD calculation in the baryon sector
  - Suitable K-matrix parameterizations suggest two poles at our m_π
  - Masses remarkably similar to physical situation in Unitarized χPT
    Consequence of $\mathrm{Tr}(M) = \mathrm{const.}$?
  - We would like to explore the quark-mass dependence, calculate more comprehensive spectra and use lattice data from additional volumes.
  - Other channels with strangeness?
  - Inconvenient things: discretization effects, chiral extrapolation, better parameterizations

Backup slides

The lightest $J^P = 0^+$ mesons

$D_0^*(2300)$ $I(J^P) = \frac{1}{2}(0^+)$ M.-L. Du *et al.*, PRL 126 192001 (2021)

- Unitarized ChiPT leads to a much lower mass than indicated by the PDG
- Authors compare data from LHCb to PDG (Breit Wigner) and Unitarized ChiPT scenarios



- Recent Lattice QCD results from HSC also obtain a much lighter state

HSC L. Gayer *et al.*, JHEP 07 (2021) 123