#### Multigrid Multilevel Monte Carlo

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Seminar @ DESY/Berlin



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#### Some Essentials

Multigrid and DD- $\alpha$ AMG

Multigrid Multilevel Monte Carlo 3

- MGMLMC for a lattice with a displacement



#### Some Essentials

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The problem at hand: compute  $tr(\Gamma D^{-1})$ . We focus on  $\Gamma = I$  here.

- D the Dirac operator
- do it stochastically

We focus here on the Wilson discretization:

$$D_{W} = \underbrace{m_{0}I}_{\text{mass shift}} + \frac{1}{2} \sum_{\mu=1}^{4} \left( \gamma_{\mu} \otimes \underbrace{(\partial_{\mu} + \partial^{\mu})}_{\text{centralized fd}} - \underbrace{al_{4} \otimes \partial_{\mu}\partial^{\mu}}_{\text{stabilization term}} \right)$$

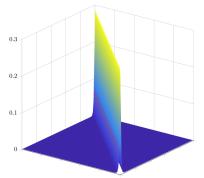
(note:  $(\Gamma_5 D_W)^H = \Gamma_5 D_W$ ).



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UNIVERSITÄT WUPPERTAL Another problem at hand:

compute  $tr(\tilde{P}^H D^{-1})$ , where D is the Dirac matrix and  $\tilde{P}$  is a shift in the rows (which comes from a displacement on the lattice in one of the four spacetime directions).



- to illustrate how this can be difficult, note here a plot of the entries of A<sup>-1</sup>, with A the 2D Laplace problem
- in LQCD, we see a rapid decay in the entries of  $D^{-1}$



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(taken from Claudia Schimmel's PhD thesis)

How do we compute the trace stochastically?

• use the Hutchinson estimator:

$$\operatorname{tr}(A) \approx \frac{1}{N} \sum_{i=1}^{N} (x^{(i)})^{H} A x^{(i)},$$

• in such a Monte Carlo computation, the variance determines the work, and in this particular case the variance takes the form (if we choose Rademacher vectors):

$$\mathbb{V}[x^{H}Ax] = \frac{1}{2} \|\text{offdiag}(A + A^{T})\|_{F}^{2}$$
$$= \frac{1}{2} \left( \|A + A^{T}\|_{F}^{2} - \sum_{i=1}^{n} (A_{ii} + (A^{T})_{ii}) \right) \xrightarrow{\text{Bergische Universität}}_{\text{WUPPERTAL}}$$

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How do we compute the trace stochastically?

- the SVD of a matrix  $A \in \mathbb{C}^{n \times n}$  has the form:  $A = U \Sigma V^H$
- there is a relation between the Frobenius norm and the singular values of a matrix:

$$\|A\|_F^2 = \sum_{i=1}^n \sigma_i^2$$

• so, to reduce the variance of the Hutchinson estimator we need to remove the largest singular values of A. In the case of  $tr(D^{-1})$ , this means removing the smallest singular values of D



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- set the orthogonal projector:  $\Pi = U_k U_k^H$
- if the columns of  $U_k$  are left singular vectors of D to very high accuracy:

$$\begin{split} \operatorname{tr}(D^{-1}) &= \operatorname{tr}(D^{-1}(I - \Pi)) + \operatorname{tr}(D^{-1}\Pi) \\ &= \operatorname{tr}(D^{-1}(I - \Pi)) + \operatorname{tr}(U_k^H V_k \Sigma_k^{-1}) \end{split}$$



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• note (remember:  $D = U\Sigma V^H$ ):

$$D^{-1}(I-\Pi) = (V\Sigma^{-1}U^H)(I-U_kU_k^H)$$

• therefore, when computing  $tr(D^{-1}(I - \Pi))$ , this deflated form will have a smaller variance compared to  $tr(D^{-1})$ 



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- BUT: computing singular vectors of *D* to high accuracy can be very expensive
- if we compute the vectors in  $U_k$  to low accuracy, then:

$$\begin{aligned} \operatorname{tr}(D^{-1}) &= \operatorname{tr}(D^{-1}(I - \Pi)) + \operatorname{tr}(D^{-1}\Pi) \\ &= \operatorname{tr}(D^{-1}(I - \Pi)) + \operatorname{tr}(U_k^H D^{-1} U_k) \end{aligned}$$

which requires extra k inversions!



• in this case of low accuracy for the vectors in  $U_k$ , we might want to switch from an orthogonal to an oblique projector:

$$\operatorname{tr}(D^{-1}) = \operatorname{tr}(D^{-1} - U_k(V_k^H D U_k)^{-1} V_k^H) + \operatorname{tr}((V_k^H D U_k)^{-1} V_k^H U_k)$$



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UNIVERSITÄT WUPPERTAL Multigrid and DD- $\alpha$ AMG: DD- $\alpha$ AMG = Domain Decomposition Aggregation-Based  $\alpha$ daptive Algebraic Multigrid

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 $DD-\alpha AMG = Domain Decomposition Aggregation-Based \alpha daptive Algebraic Multigrid.$ 

The main problem with iterative solvers when dealing with Ax = b

 condition number of a matrix A, κ(A), gives a measure on how difficult it is to solve a system of equations of the form Ax = b

• 
$$\kappa_p(A) = ||A||_p ||A^{-1}||_p$$

• 
$$\kappa_2(A) = \frac{\sigma_1}{\sigma_n}$$

iter. count	$\kappa(A)$
37	100
592	400
9,472	1,600
151,552	6,400

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## Multigrid and DD- $\alpha$ AMG

Iterative solvers as smoothers: Gauss-Seidel applied to the Laplace 2D problem

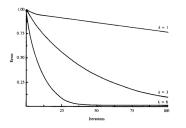
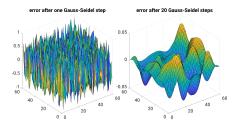


Image taken from A Multigrid Tutorial by Briggs

et.al. Three different right hand sides.

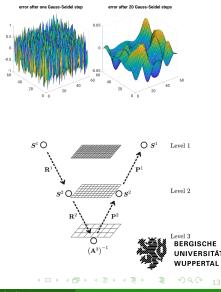


Starting with a random initial guess.



# Multigrid and DD- $\alpha$ AMG

- the error can be split as:
  - $e = e_{low} + e_{high}$ .
- a good smoother eliminates *e<sub>high</sub>* effectively.
- multigrid solvers can lead to convergence independent of κ(A).



Algebraic multigrid: non-geometric.

Two-level AMG:

- Petrov-Galerkin projection:  $D_c = RDP$
- coarse-grid correction:  $\psi \leftarrow \psi + PD_c^{-1}Rr$ , with  $r = \eta - D\psi$
- error propagation:  $E = I - PD_c^{-1}RD$
- ⇒ choose range(P) to contain the near kernel

#### Algorithm 1 Two-level V-cycle

Input :  $\psi$ ,  $\eta$ ,  $\nu$ Output :  $\psi$ 

$$r \leftarrow \eta - D\psi \\ \psi \leftarrow \psi + PD_c^{-1}Rr \\ r \leftarrow \eta - D\psi \\ \psi \leftarrow \psi + M_{SAP}^{(\nu)}r$$

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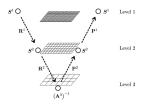


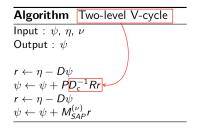
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#### Multigrid and DD- $\alpha$ AMG

Multilevel AMG:

this extension is possible by decomposing  $e_{low} = e_{c,low} + e_{c,high}$  into the second level.





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• 
$$\mathcal{A}_i := \mathcal{L}_{j(i)} \times \mathcal{W}_i$$

• standard aggregation:  $\{\mathcal{A}_{j,\tau} : j = 1, ..., n_{\mathcal{L}_c}, \tau = 0, 1\}$ , such that:

$$\mathcal{A}_{j,0} := \mathcal{L}_j imes \{0,1\} imes \mathcal{C} \ \mathcal{A}_{j,1} := \mathcal{L}_j imes \{2,3\} imes \mathcal{C}$$

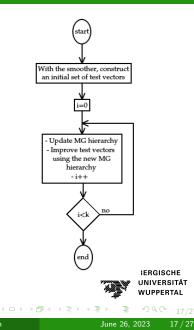
- when using the standard aggregation: Γ<sub>5</sub>P = PΓ<sub>5</sub><sup>c</sup>
- also, it makes sense to choose: D<sub>c</sub> = P<sup>H</sup>DP

A section of P:

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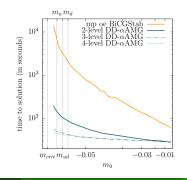


- test vectors: the approximate eigenvectors used in the construction of *P*, i.e. {v<sub>i</sub> : i = 1,..., N}
- the setup phase in DD-αAMG builds these test vectors
- the method updates the multigrid hierarchy as the test vectors improve



# Multigrid and DD- $\alpha$ AMG

Comparison of computational cost for solving linear systems with a configuration from a BMW collaboration configuration using DD- $\alpha$ AMG and a Krylov subspace method. The left plot reports on timings for the solve only, whereas the right plot includes the multigrid setup time. Both plots were generated on the JUROPA high performance computer from the Jülich Supercomputing Centre (taken from Matthias Rottmann's PhD thesis).





#### MGMLMC: Multigrid Multilevel Monte Carlo

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Assumming a multigrid construction such that  $R_{\ell} = P_{\ell}^{H}$  and  $P_{\ell}^{H}P_{\ell} = I$ , then (Frommer, Nasr, R-H, 2022):

• two levels:

$$\operatorname{tr}(A_0^{-1}) = \operatorname{tr}(A_0^{-1}) - \operatorname{tr}(P_0 A_1^{-1} R_0) + \operatorname{tr}(P_0 A_1^{-1} R_0)$$
(1)  
= 
$$\operatorname{tr}(A_0^{-1} - P_0 A_1^{-1} P_0^H) + \operatorname{tr}(A_1^{-1})$$
(2)



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(1)  
= 
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(2)

 we can connect this to inexact/approximate deflation, which is discussed in *Multigrid deflation for Lattice QCD* (Romero, Stathopoulos and Orginos), and already in Lüscher's local coherence paper



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We move from a two-level to a multilevel method by means of Multilevel Monte Carlo (Giles, 2008), and then write (Frommer, Nasr, R-H, 2022):

multilevel:

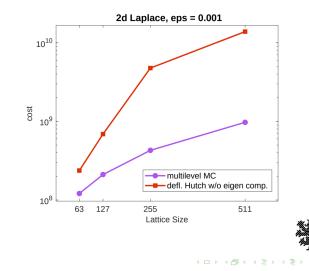
$$\operatorname{tr}(A_0^{-1}) = \sum_{\ell=0}^{L-2} \operatorname{tr}\left(A_\ell^{-1} - P_\ell A_{\ell+1}^{-1} R_\ell\right) + \operatorname{tr}(A_{L-1}^{-1})$$
(3)
$$\left(\operatorname{cost}(D_1^{-1} b_1) \sim 0.1 \operatorname{cost}(D_0^{-1} b_0)\right)$$



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MGMLMC vs exactly deflated Hutchinson (Frommer, Nasr, R-H, 2022):



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A short note on deflation in LQCD:

if we want to deflate as  $tr((I - \Pi)D^{-1})$ , we can use the following to compute approximate deflation vectors:

• the Dirac operator fulfills  $\Gamma_5$ -hermiticity:  $(\Gamma_5 D)^H = \Gamma_5 D$ 



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- this gives us access to the SVD:  $\Gamma_5 D = \frac{V \Sigma V^H}{\Gamma_5 V \operatorname{sign}(\Sigma)} \operatorname{abs}(\Sigma) \frac{V^H}{\Gamma_5 V \operatorname{sign}(\Sigma)} + \frac{V \operatorname{abs}(\Sigma^{-1})}{\Gamma_5 V \operatorname{sign}(\Sigma)} + \frac{V \operatorname{abs}(\Sigma^{$



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- the Dirac operator fulfills  $\Gamma_5$ -hermiticity:  $(\Gamma_5 D)^H = \Gamma_5 D$
- this gives us access to the SVD:  $\Gamma_5 D = \frac{V \Sigma V^H}{P} \Rightarrow D = (\Gamma_5 V \operatorname{sign}(\Sigma)) \operatorname{abs}(\Sigma) V^H$   $\Rightarrow D^{-1} = \frac{V}{\operatorname{abs}(\Sigma^{-1})} (\Gamma_5 V \operatorname{sign}(\Sigma))^H$
- $\bullet$   $\Rightarrow$  block pow. it. on  $(\Gamma_5 D)^{-1}$  gives columns of V , ready for deflation



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A short note on deflation in LQCD:

if we want to deflate in MGMLMC as e.g.  $tr((I - \Pi)(D^{-1} - P_1D_2^{-1}P_1^H))$ , do we have something similar to the previous slide to compute the deflation vectors?

• in our multigrid construction:

$$\Gamma_{5}^{\ell} P_{\ell} = P_{\ell} \Gamma_{5}^{\ell+1}, \ (\Gamma_{5}^{\ell} D_{\ell})^{H} = \Gamma_{5}^{\ell} D_{\ell}$$
(4)  
which implies:  $((D_{\ell}^{-1} - P_{\ell} D_{\ell+1}^{-1} P_{\ell}^{H}) \Gamma_{5}^{\ell})^{H} = (D_{\ell}^{-1} - P_{\ell} D_{\ell+1}^{-1} P_{\ell}^{H}) \Gamma_{5}^{\ell}$ 



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# MGMLMC for a lattice with a displacement (i.e. $tr(\tilde{P}^H D_0^{-1})$ )

Two-level decomposition:

$$\operatorname{tr}(\widetilde{P}^{H}D^{-1}) = \operatorname{tr}\left(\widetilde{P}^{H}(D_{0}^{-1} - P_{0}D_{1}^{-1}P_{0}^{H})\right) + \operatorname{tr}(\widetilde{P}^{H}P_{0}D_{1}^{-1}P_{0}^{H})$$
$$= \operatorname{tr}\left(\widetilde{P}^{H}(D_{0}^{-1} - P_{0}D_{1}^{-1}P_{0}^{H})\right) + \operatorname{tr}(P_{0}^{H}\widetilde{P}^{H}P_{0}D_{1}^{-1})$$

Let us introduce:  $B_1 := P_0^H \widetilde{P}^H P_0$ .



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I have recently implemented a minimal version of DD- $\alpha \rm AMG$  in MATLAB, which is available here:

https://github.com/Gustavroot/MGMLMC\_for\_mat\_with\_displ

In that MATLAB code, also orthogonal deflation with low-accuracy vectors has been implemented for both Hutchinson and MGMLMC. This code allows one to compute variances and compare both methods.



Some results (4-D problem, LQCD,  $16^4$  lattice  $\rightarrow$  786, 432 matrix dim.,  $5 \times 12$  matrix displacement  $\widetilde{P} - \beta = 3.55$ ,  $\kappa = 0.137$ ,  $m_{\pi} \approx 420$  MeV,  $N_f = 3$ , p.b.c.):



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• sampled variances (at  $\ell = 0$ ):

• 
$$\operatorname{var}(\Pi D_0^{-1} \widetilde{P}^H)$$
:  
 $\boxed{n_{defl} \quad \operatorname{var} (\times 10^3)}$   
 $\boxed{0 \quad 145.5}$   
 $32 \quad 131.8$   
 $128 \quad 120.5$   
 $512 \quad 115.3$   
•  $\operatorname{var}\left(\Pi(D_0^{-1} - P_0 D_1^{-1} P_0^H) \widetilde{P}^H\right)$   
 $\boxed{n_{defl} \quad \operatorname{var} (\times 10^3)}$   
 $\boxed{0 \quad 110.7}$   
 $32 \quad 109.8$   
 $128 \quad 105.5$ 

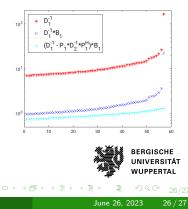
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- sampled variances (at  $\ell = 0$ ):
  - $\operatorname{var}(\Pi D_0^{-1} \widetilde{P}^H)$ : var  $(\times 10^3)$ n<sub>defl</sub> 0 145.5 32 131.8 128 120.5 512 115.3 • var  $\left( \prod (D_0^{-1} - P_0 D_1^{-1} P_0^H) \widetilde{P}^H \right)$ : var ( $\times 10^3$ ) n<sub>defl</sub> 0 110.7 32 109.8 128 105.5
- singular values (at  $\ell = 1$ , ascending order):



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- more details of these results for a lattice with a displacement will be made available in a proceedings at the end of this week
- large-scale tests (C, MPI, OpenMP, large lattices) are under preparation and will be released soon in a paper, as part of the work of Jose Jimenez-Merchan (PhD student in our group)



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Thank you!

Questions?

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