

The CS Kernel for dynamical z_{max}

April 6, 2023

The CS Kernel

The CSS factorisation formula:

$$\frac{d\sigma}{dQ^2 dy dq_T^2} = \frac{2\pi}{9} \frac{\alpha_{\text{em}}^2(Q)}{sQ^2} |C_V(Q, \mu_Q)|^2 \int_0^\infty db b J_0(bq_T) \\ \times \sum_q e_q^2 f_{q,h_1}(x_1, b; \mu_Q, Q^2) f_{\bar{q},h_2}(x_2, b; \mu_Q, Q^2), \quad (1)$$

The CS Kernel is defined as

$$\frac{df_{q,h}(x, b; \mu, \zeta)}{d \ln \zeta} = -\mathcal{D}(b, \mu) f_{q,h}(x, b; \mu, \zeta)$$

The CS Kernel from ratios

$$\mathcal{D}(b, \mu_0) = \frac{\ln\left(\frac{\Sigma_1}{\Sigma_2}\right) - \ln Z(Q_1, Q_2) - 2\Delta_R(Q_1, Q_2; \mu_0)}{4 \ln(Q_2/Q_1)} - 1$$

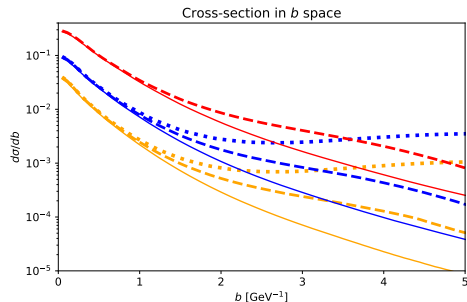
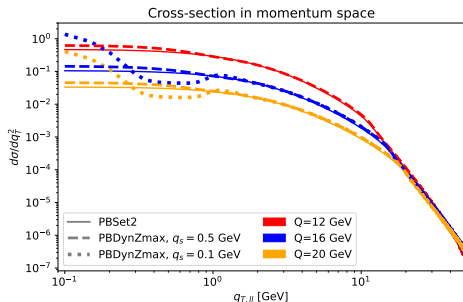
where Σ is the Hankel transformation of the cross-section:

$$\Sigma(s, y, Q, b) = \int_0^\infty dq_T q_T J_0(q_T b) \frac{d\sigma}{dQ^2 dy dq_T^2}.$$

and

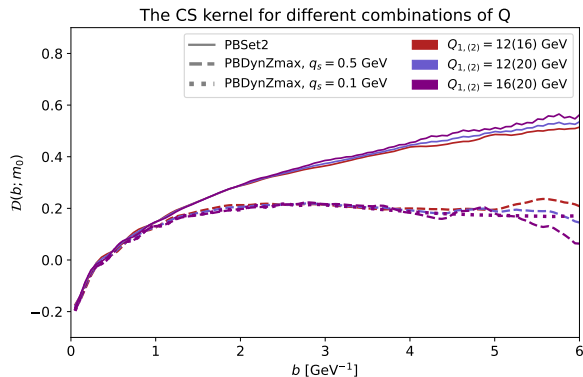
$$\Delta_R(Q_1, Q_2; \mu_0) = \int_{\mu_{Q_2}}^{\mu_{Q_1}} \frac{d\mu}{\mu} \gamma_F(\mu, Q_1) - 2 \ln\left(\frac{Q_1}{Q_2}\right) \int_{\mu_0}^{\mu_{Q_2}} \frac{d\mu}{\mu} \Gamma_{\text{cusp}}(\mu) \quad Z(Q_1, Q_2) = \frac{\alpha_{\text{em}}^2(Q_1) |C_V(Q_1, \mu_{Q_1})|^2}{\alpha_{\text{em}}^2(Q_2) |C_V(Q_2, \mu_{Q_2})|^2}$$

The DY cross section: PBSet2 vs Dynamical Z_{max}



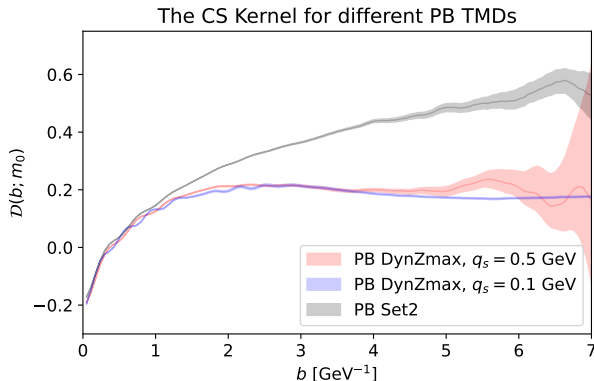
- The three calculations agree at large q_T (small b). The effect of the different TMDs comes at low q_T (large b)
- Dynamical Z_{max} TMDs do not allow emissions with $q < q_0 = 1$ GeV \rightarrow It creates a valley between q_0 and q_s in momentum space (a peak in b space for $b > 1$ GeV^{-1})

The CS kernel in CASCADE: PBSet2 vs Dynamical z_{max}



- All three TMDs and the different combinations agree at small b (pert. calculation)

The CS kernel in CASCADE: Flat behaviour at large b



- The dependence on b of the CS Kernel comes from emissions $q_0 < 1 \text{ GeV}$
- We can observe that there is no dependence on the intrinsic k_T
- The CS Kernel extracted from PB DynZmax are constant at a value of $D \sim 0.2$

Conclusions

- We have extracted the CS Kernel for dynamical z max TMDs with $q_0 = 1$ GeV, which does not allow emissions below q_0
- The perturbative calculations ($b < 1$ GeV $^{-1}$) agree between the different TMDs
- The CS Kernel for dynamical z max TMDs saturate at $D \sim 0.2$, no dependence on b for $b > 1$ GeV $^{-1}$
- The flat behaviour arises from the fact that at low q_T the distribution is the same for different values of Q
- Could we derive an uncertainty using a variation of q_0 (q_{cut}) for dynamical z max (PBSet2) TMDs?

The End