

Evaluating Parametric Integrals in the Minkowski Regime without Contour Deformation

Thomas Stone

In collaboration with Stephen Jones (IPPP) & Anton Olsson (KIT)

(and the rest of the pySecDec collaboration:
G. Heinrich, M. Kerner, V. Magerya, J. Schlenk)

16th April 2024



Outline

- 1 Introduction & Motivation
- 2 Massless Integrals
 - 1-Loop Off-Shell Box
 - 2-Loop Non-Planar Box
 - 3-Loop Non-Planar Box
- 3 Massive Integrals
 - Massive Bubble
 - Massive Triangle
- 4 Outlook

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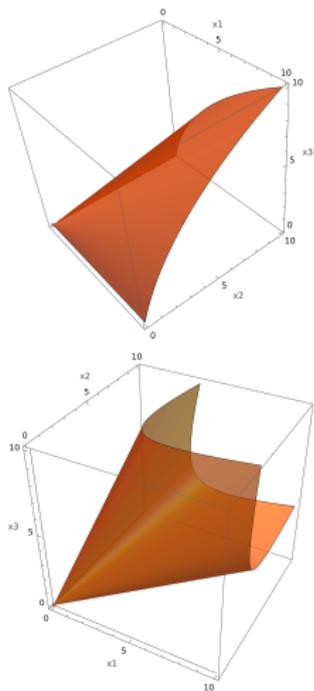
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- Many loop integrals appearing in state-of-the-art amplitude calculations are analytically intractable
- Numerical methods developed to tackle these integrals (Monte Carlo techniques, differential equation methods etc.)
- Exploring singularity structure of Feynman integrals can help us understand how to integrate in the Minkowski regime



Minkowski Regime

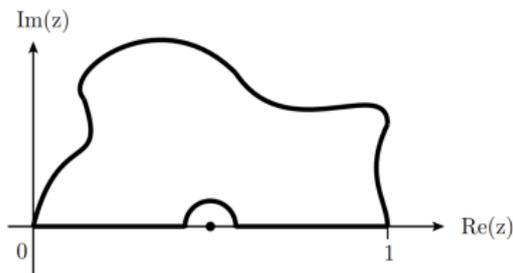
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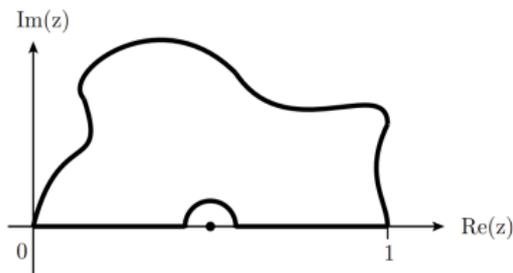
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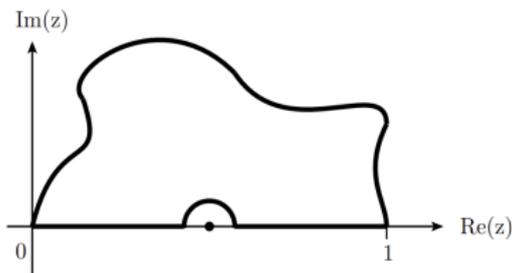
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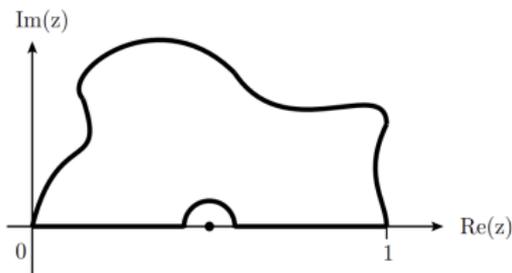


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- Methods being explored to remove need for contour deformation in momentum space
[Buchta, Chachamis, Draggiotis, Rodrigo; Anastasiou, Haindl, Sterman, Yang, Zeng; Aguilera-Verdugo, Hernandez-Pinto, Sborlini, Torres Bobadilla; Capatti, Hirschi, Kermanschah, Pelloni, Ruijl; ...]

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- Can we do the same in Feynman parameter space?

Recap: Feynman Parameterisation

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Momentum Space Integral

$$I = \int_{-\infty}^{+\infty} \left(\prod_{l=1}^L \frac{d^D k_l}{i\pi^{\frac{D}{2}}} \right) \prod_{j=1}^N \frac{1}{P_j^{\nu_j}(\{k\}, \{p\}, m_j^2)}$$

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Feynman-Parameterised Integral

$$I = \frac{(-1)^\nu \Gamma(\nu - LD/2)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_{\mathbb{R}_{\geq 0}^N} \left(\prod_{j=1}^N dx_j x_j^{\nu_j - 1} \right) \frac{\mathcal{U}(\mathbf{x})^{\nu - (L+1)D/2}}{(\mathcal{F}(\mathbf{x}, \mathbf{s}) - i\delta)^{\nu - LD/2}} \delta\left(1 - \sum_{j=1}^N x_j\right)$$

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\mathcal{U} & \mathcal{F} constructable directly from Feynman diagrams with $\mathcal{U} \geq 0$
and \mathcal{F} depending on both parameters **and** kinematics

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Implementation in Parameter Space

$$\mathcal{F}(\vec{z}) = \mathcal{F}(\vec{x}) - i \sum_j \tau_j \frac{\partial \mathcal{F}(\vec{x})}{\partial x_j} \quad \tau_j = \lambda_j x_j (1 - x_j) \frac{\partial \mathcal{F}(\vec{x})}{\partial x_j}$$

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- There are even cases where this procedure fails entirely!

[see Stephen Jones' talk]

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Maybe...?

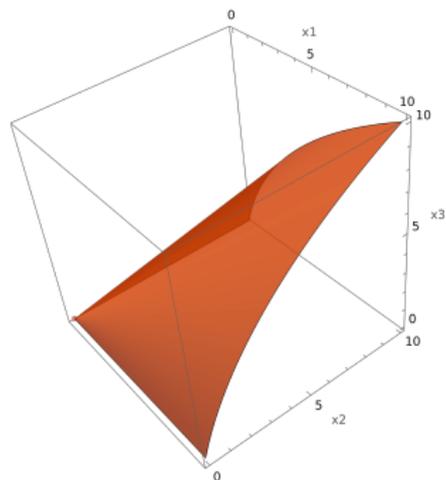
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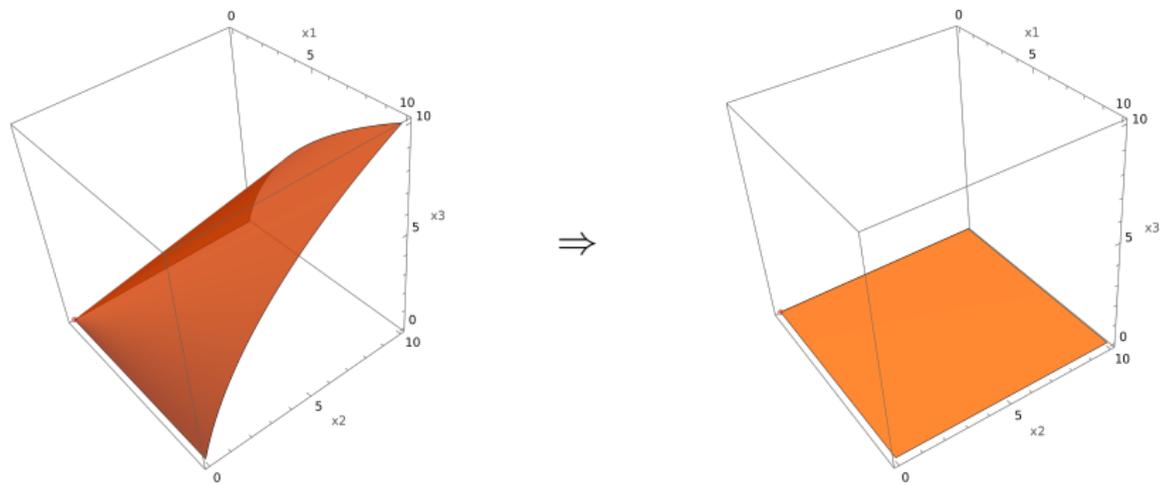
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$$\forall S \subseteq \{1, \dots, N\} \wedge S \neq \emptyset : \\ \delta \left(1 - \sum_{j=1}^N x_j \right) \rightarrow \delta \left(1 - \sum_{j \in S} x_j \right) \text{ leaves } I \text{ invariant}$$

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Heaviside Identity

Under the integral sign: $\theta(x_a - x_b) + \theta(x_b - x_a) = 1$

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Overall Construction

$$I = \sum_{n_+=1}^{N_+} I_{n_+}^+ + (-1 - i\delta)^{-(\nu-LD/2)} \sum_{n_-=1}^{N_-} I_{n_-}^-$$

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Overall Construction

$$I = \sum_{n_+=1}^{N_+} I_{n_+}^+ + (-1 - i\delta)^{-(\nu-LD/2)} \sum_{n_-=1}^{N_-} I_{n_-}^-$$

- Can be **much** faster numerically to only calculate the manifestly non-negative integrals $\{I_{n_+}^+, I_{n_-}^-\}$!

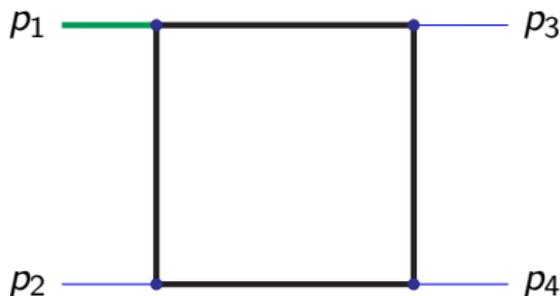
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Lets consider the one-loop massless box with an offshell leg to make this concrete:

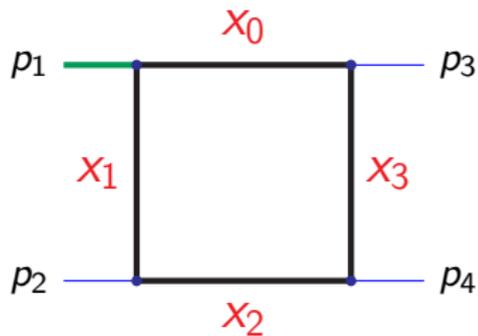
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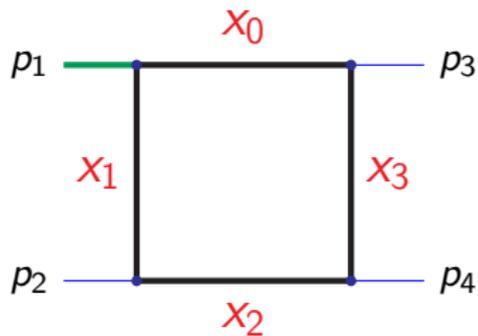


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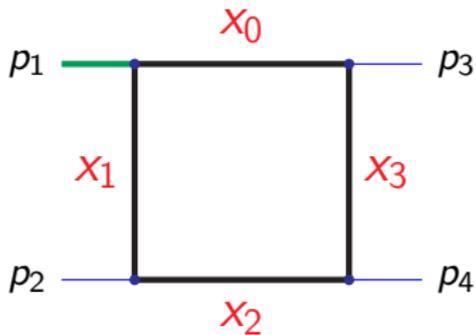


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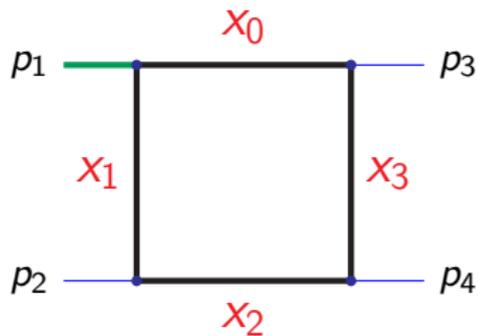
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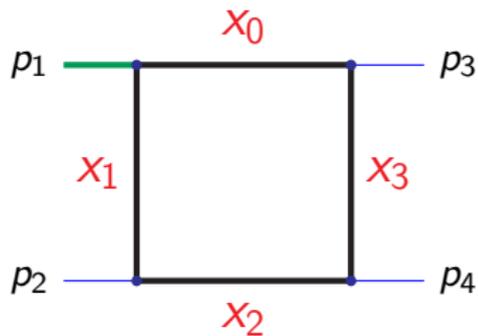
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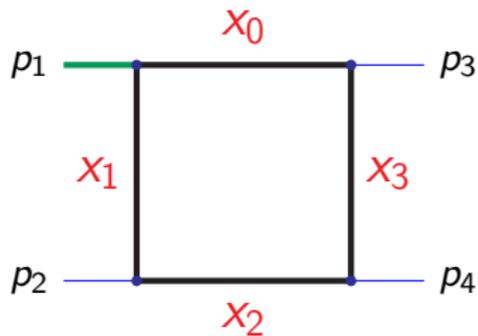
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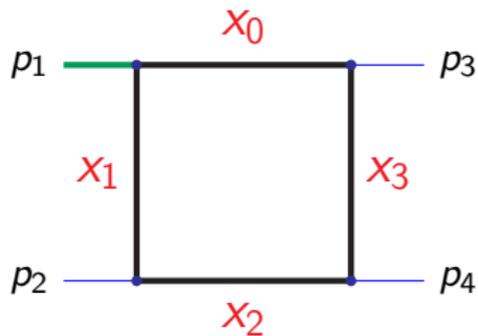
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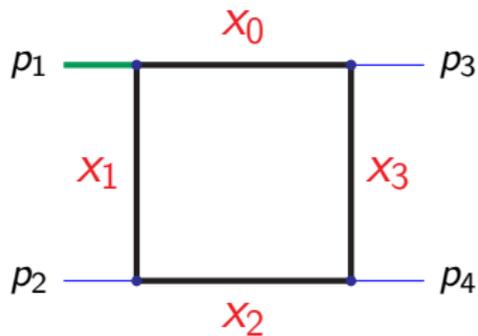
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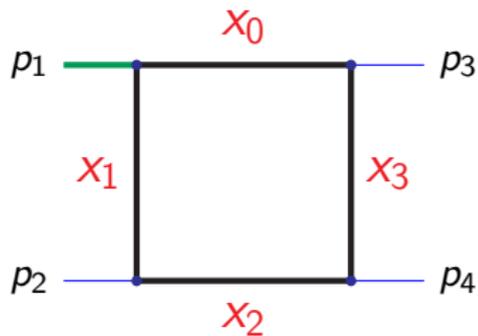
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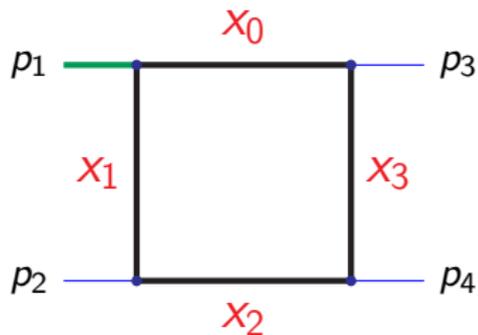
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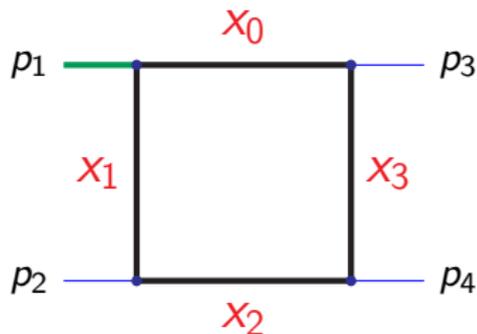
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Let's consider the regime: $s > 0$, $p_1^2 > 0$ & $t < 0 \Rightarrow$ zeroes of \mathcal{F} *within* the integration volume for $\{x_0, x_1, x_2, x_3\} \in \mathbb{R}_{\geq 0}^4$

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- $\mathcal{F} \rightarrow x_1 \left(x_2 (x_3 - x_0) - \frac{p_1^2}{s} x_0 x_1 \right)$
- Introduce the hierarchy $x_0 > x_3$ by *shifting* x_0 : $x_0 \rightarrow x_0 + x_3$
- $\mathcal{F} \rightarrow -\frac{1}{s} \left(x_1 (s x_0 x_2 + p_1^2 x_1 (x_0 + x_3)) \right) =: -\mathcal{F}_1^-$
- \mathcal{F}_1^- non-negative as required!
- To cover the whole original space, the Heaviside identity tells us we need to consider the converse case ($x_3 > x_0$) so *shift* x_3 :
 $x_3 \rightarrow x_0 + x_3$
- $\mathcal{F} \rightarrow x_1 \left(-\frac{p_1^2}{s} x_0 x_1 + x_2 x_3 \right)$
- This is **not** of uniform sign \Rightarrow needs further work!

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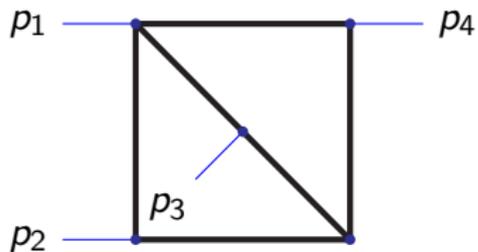
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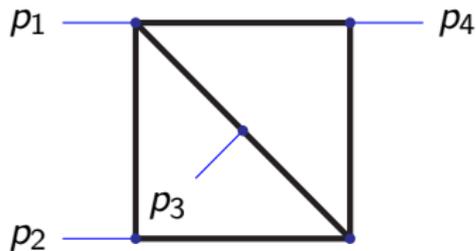
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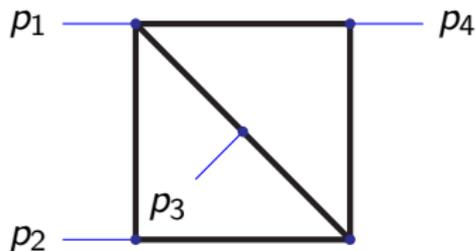


2-Loop Non-Planar Box



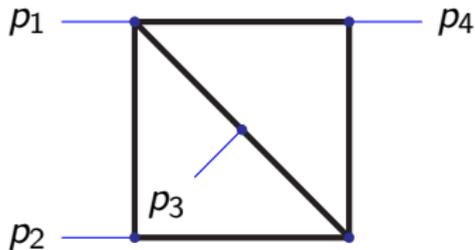
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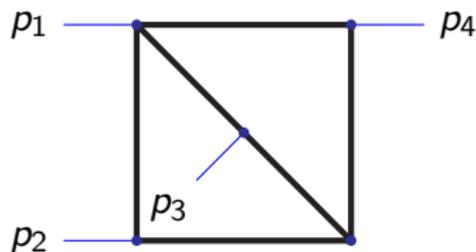
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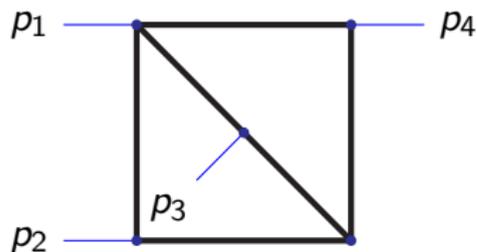


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Hence, \mathcal{F} can be 0 *within* $\{x_i\} \in \mathbb{R}_{\geq 0}^6$ even with $s > 0, t > 0$

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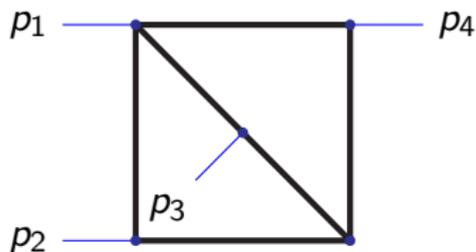
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Nevertheless, the method works

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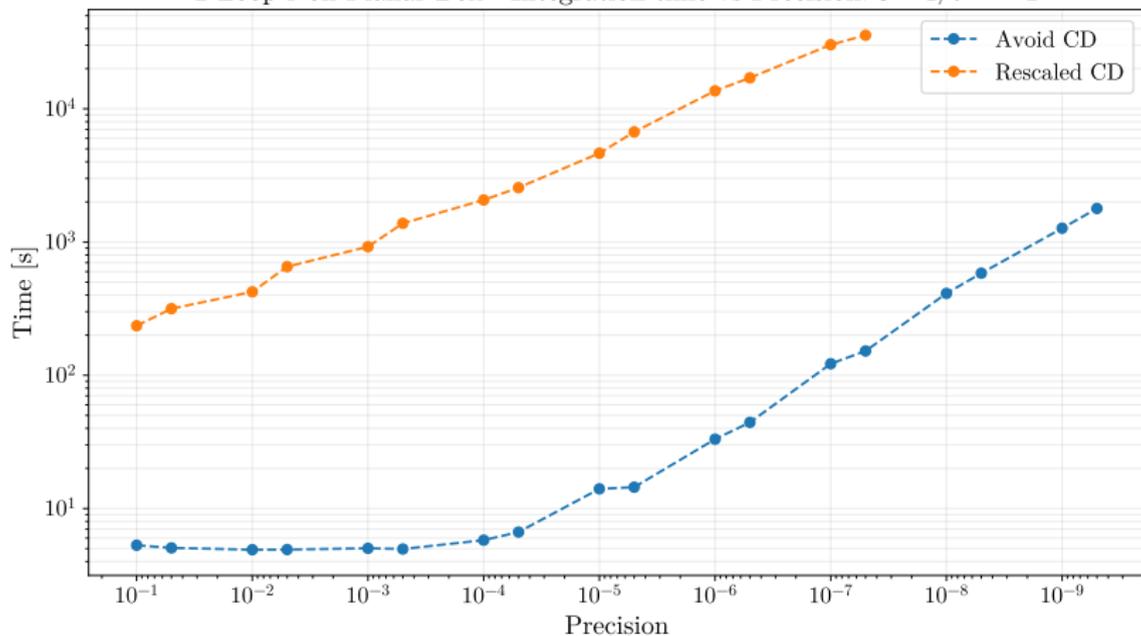
δ -Function Issue

Some rescalings can modify our δ -function in an impractical way for our numerical integration setup - let's avoid that!

2-Loop Non-Planar Box: pySecDec Timing Comparison

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2-Loop Non-Planar Box - Integration time vs Precision: $s = 4, t = -1$



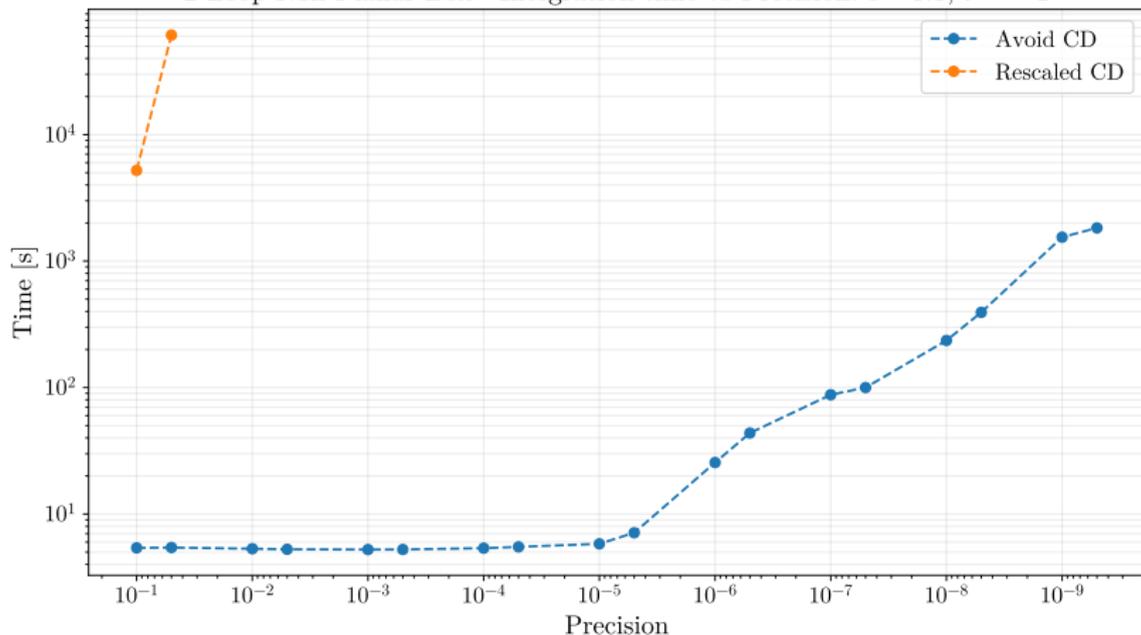
Evaluating up-to-and-including finite order with pySecDec

[Heinrich, Jones, Kerner, Magerya, Olsson, Schlenk]

2-Loop Non-Planar Box: pySecDec Timing Comparison

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Evaluating up-to-and-including finite order with pySecDec

3-Loop Non-Planar Box

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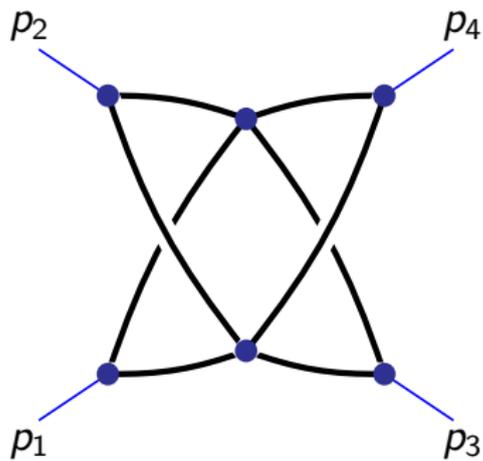


Diagram by Yao Ma

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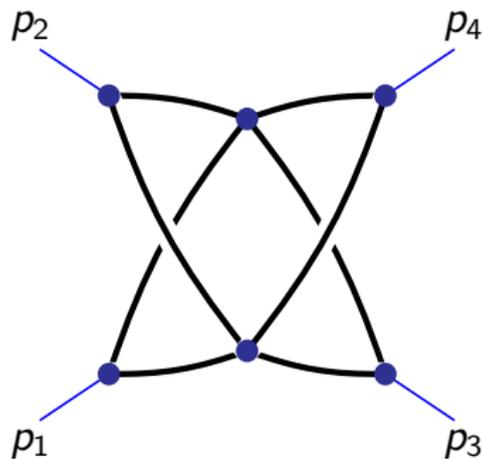


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$$\mathcal{F} = -s(x_1x_4 - x_0x_5)(x_3x_6 - x_2x_7) \\ - t(x_1x_2 - x_0x_3)(x_5x_6 - x_4x_7)$$

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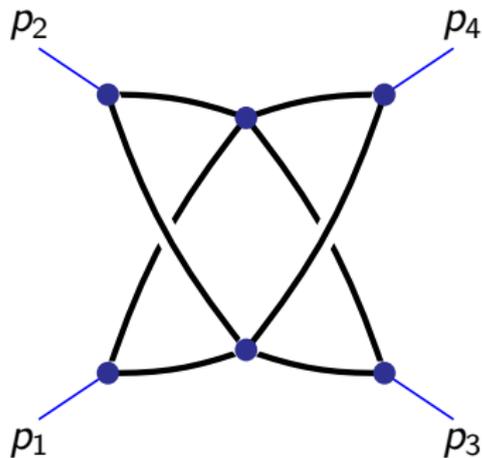


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[see Stephen Jones' talk]



$$\mathcal{F} \rightarrow -sx_1x_3x_5x_7(x_4 - x_0)(x_6 - x_2) \\ - tx_1x_3x_5x_7(x_2 - x_6)(x_6 - x_4)$$

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- This gives us 12 integrals to compute, none of which require contour deformation!

3-Loop Non-Planar Box: Putting the Pieces Together

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$$I = \sum_{n_+=1}^8 I_{n_+}^+ + (-1 - i\delta)^{-2-3\epsilon} \sum_{n_-=1}^4 I_{n_-}^-$$

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- Avoiding contour deformation allowed us to go to higher orders (tried up to $\frac{1}{\epsilon^2}$)

3-Loop Non-Planar Box: pySecDec Precision Comparison

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$$I_a^{\text{NOCD}} = \varepsilon^{-4} \left[(18.51948920208488 - 15.70796326794897i) \pm (4.032 * 10^{-11} + 4.592 * 10^{-11}i) \right] + \dots$$

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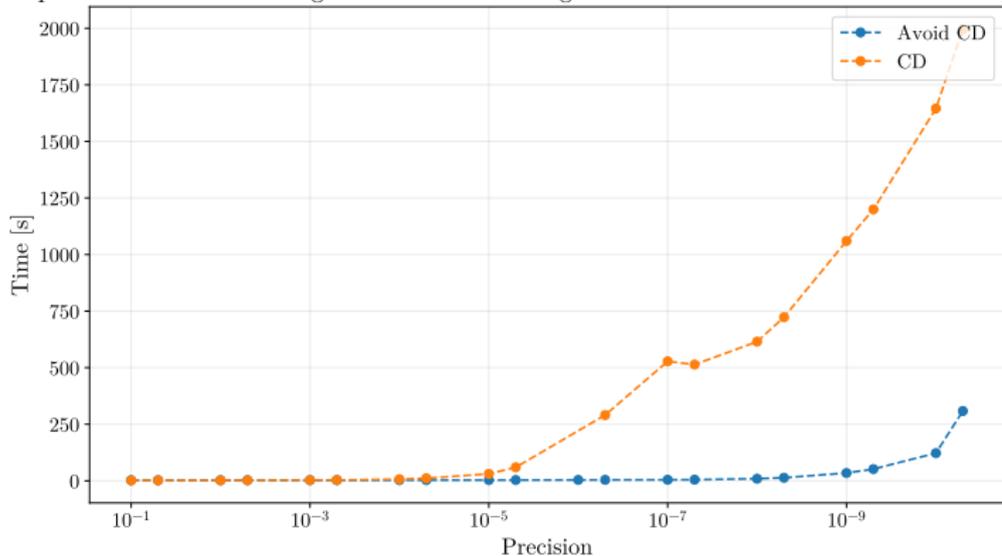
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3-Loop Non-Planar Box: pySecDec Timing Comparison

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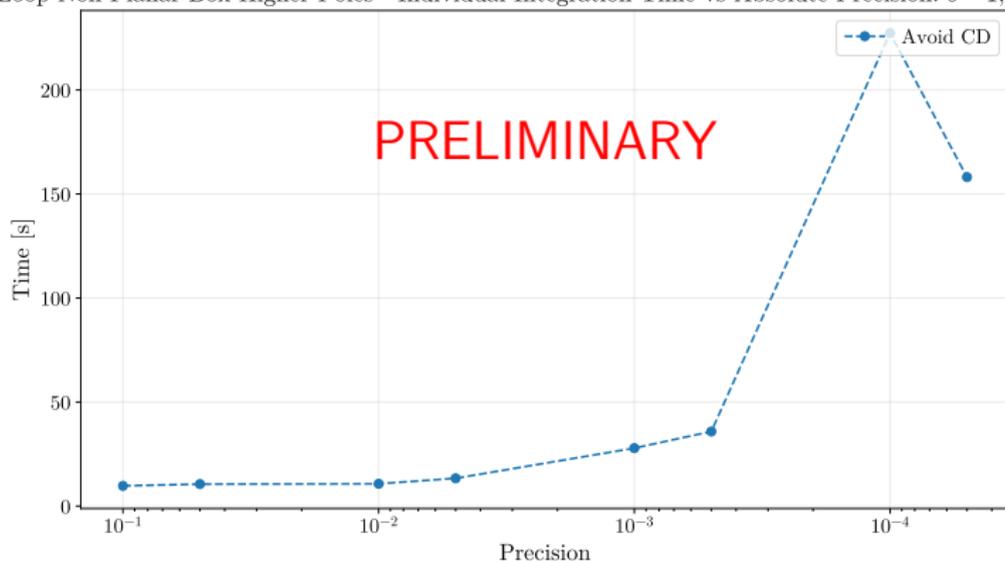
3-Loop Non-Planar Box Leading Pole - Individual Integration Time vs Absolute Precision: $s=1$, $t=-1/5$



Evaluating leading pole ($\frac{1}{\epsilon^4}$) with pySecDec

3-Loop Non-Planar Box: pySecDec Timing

3-Loop Non-Planar Box Higher Poles - Individual Integration Time vs Absolute Precision: $s = 1$, $t = -1/5$



Evaluating up-to-and-including $\frac{1}{\epsilon^2}$ with pySecDec

Table of Contents

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Massive Integrals

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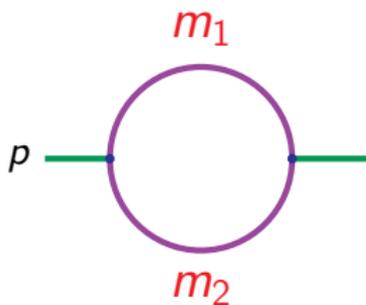
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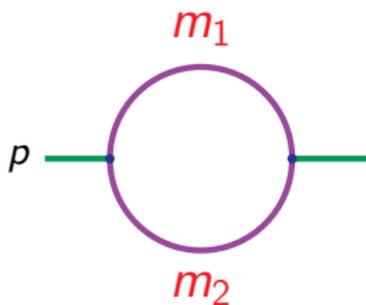
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- Viable transformations difficult even for trivial integrals
- Can we use geometry to guide us in the right direction?

Massive Bubble

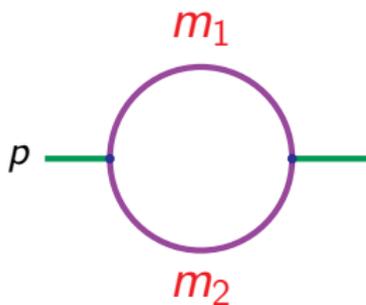


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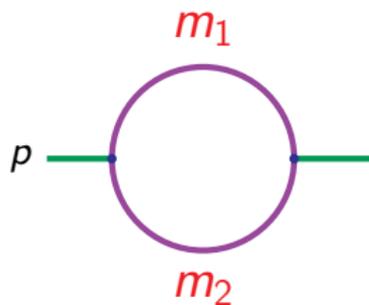
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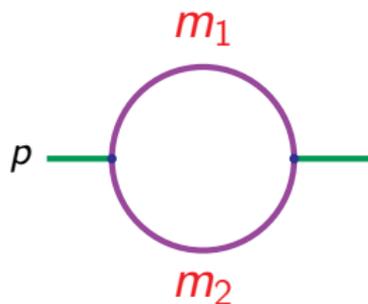
- $\mathcal{F} = -p^2 x_1 x_2 + (x_1 + x_2) (m_1^2 x_1 + m_2^2 x_2)$

Massive Bubble



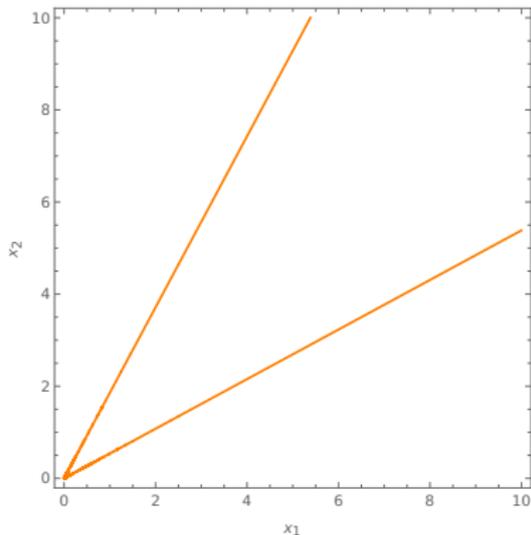
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Massive Bubble



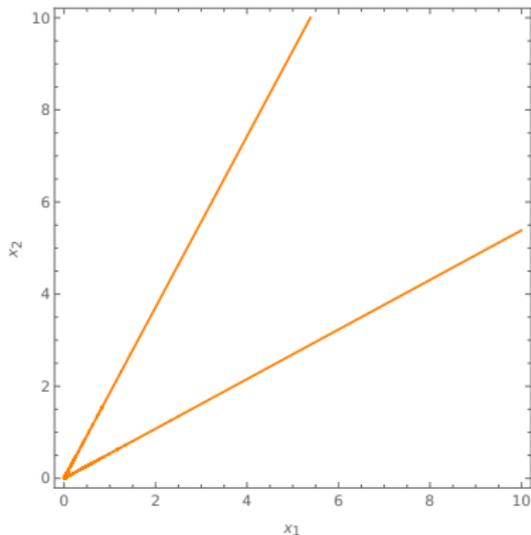
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- Scale out dimension of \mathcal{F} via $x_i \rightarrow \frac{x_i}{m_i}$

Massive Bubble



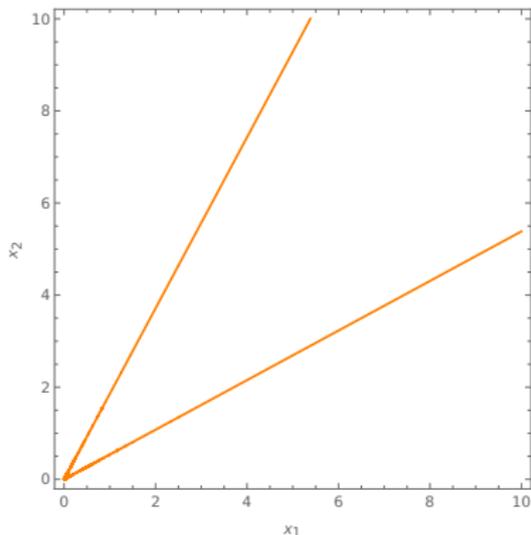
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Massive Bubble



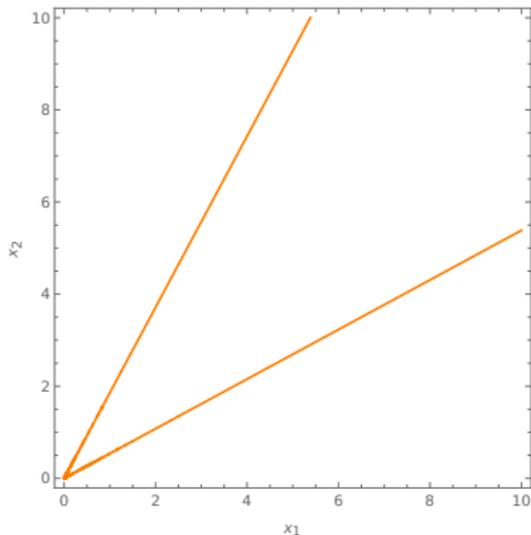
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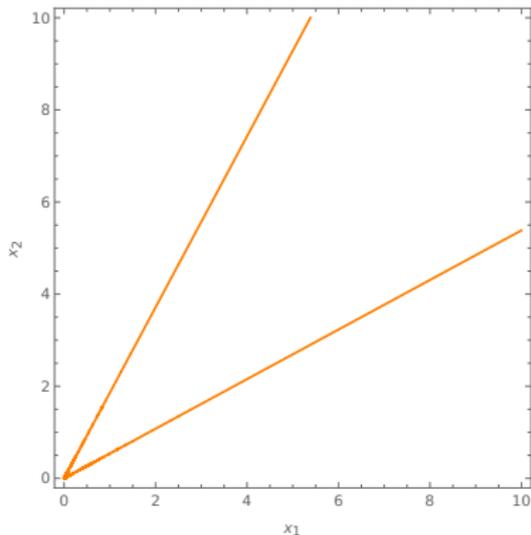


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Massive Bubble

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- Construct transformations which directly send the variety to the integration boundary

Massive Bubble

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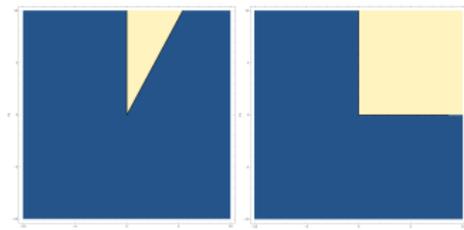
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Massive Bubble

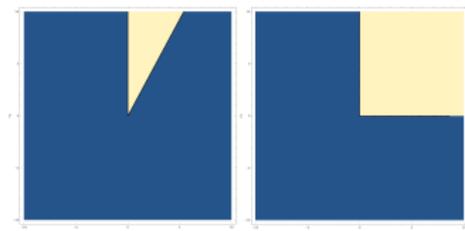
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$$y_1 \stackrel{!}{=} x_1, y_2 \stackrel{!}{=} x_2 - \frac{1+\beta}{1-\beta}x_1 \Rightarrow x_1 \rightarrow y_1, x_2 \rightarrow y_2 + \frac{1+\beta}{1-\beta}y_1$$

Massive Bubble

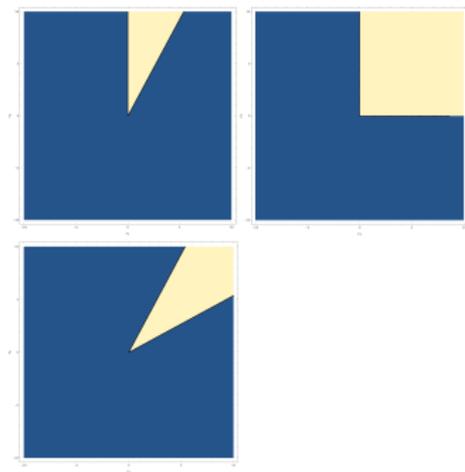


Massive Bubble



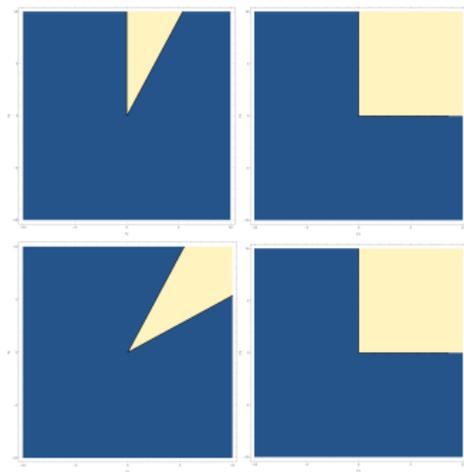
$$\tilde{\mathcal{F}}_1^+ = y_2 \left(y_2 + \frac{4\beta}{1-\beta^2} y_1 \right)$$

Massive Bubble



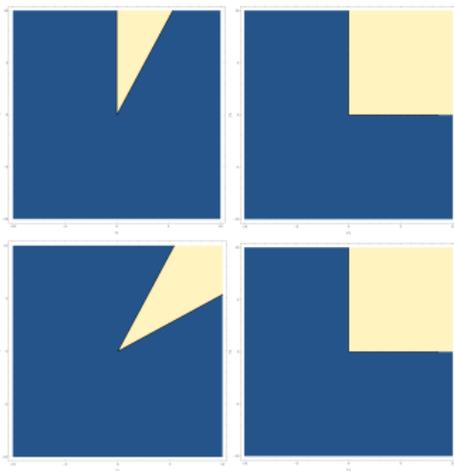
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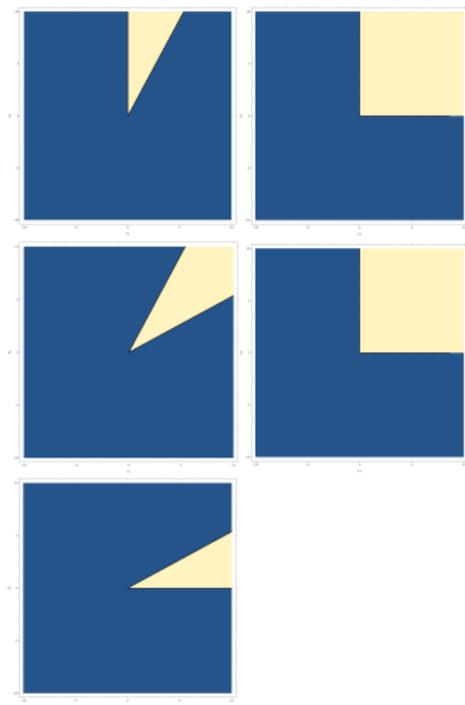
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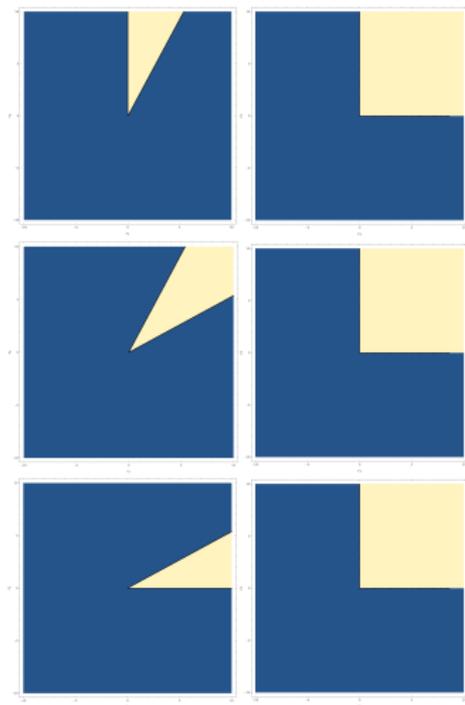
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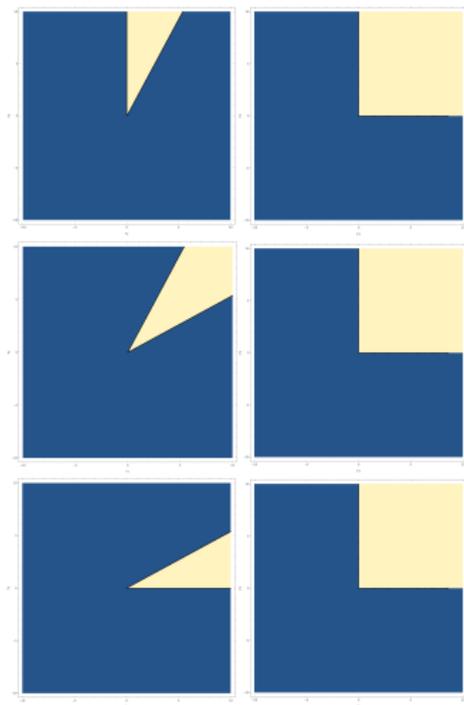
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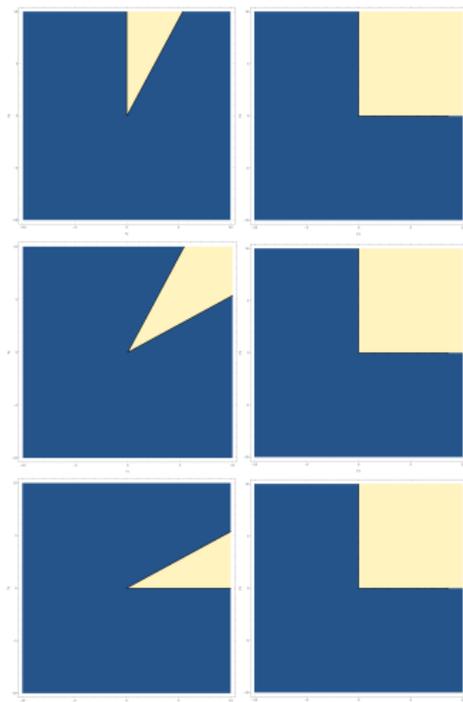


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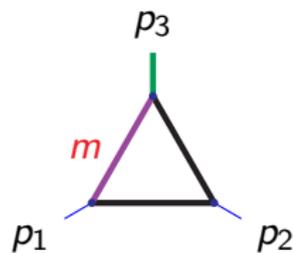
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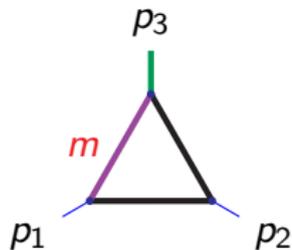
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Verified result numerically & analytically ✓

Massive Triangle

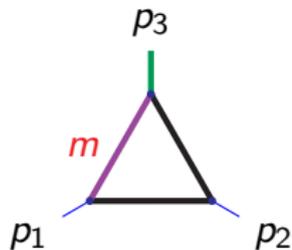


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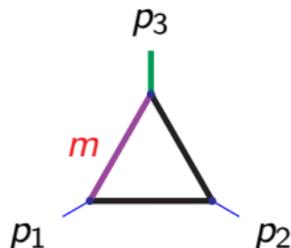
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Massive Triangle



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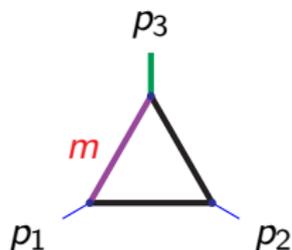
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Can solve this using **rescalings** and **shifts** as before (verified ✓)

Massive Triangle



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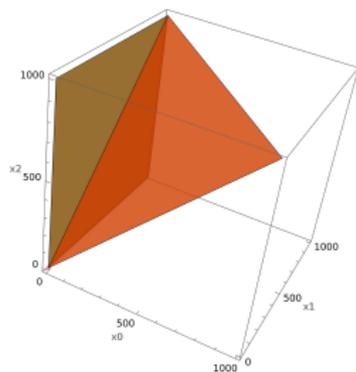
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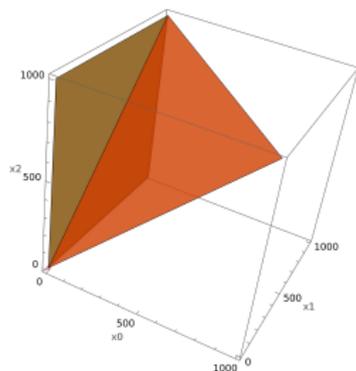
$$I = I_1^+ + I_2^+ + (-1 - i\delta)^{-1-\varepsilon} I_1^-$$

Massive Triangle

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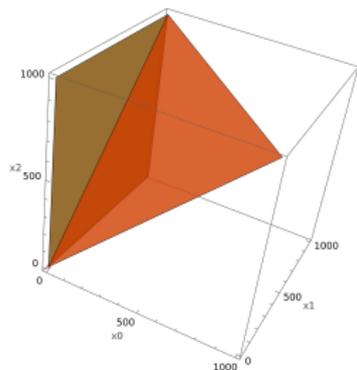


Massive Triangle



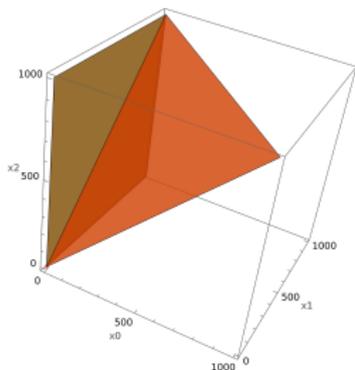
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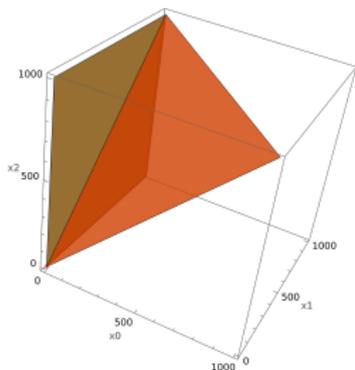
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Massive Triangle



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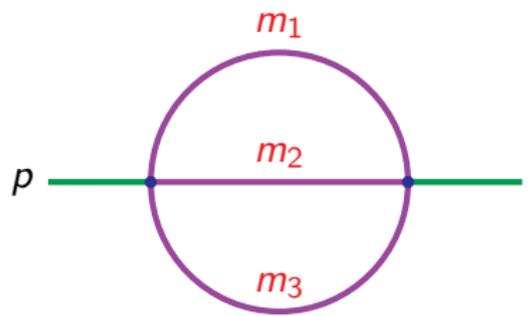


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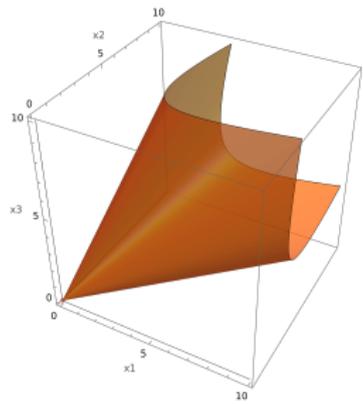
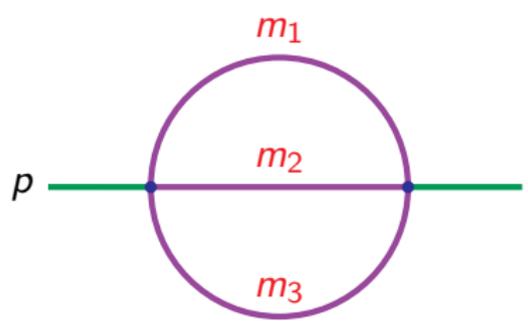
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- 4 Outlook

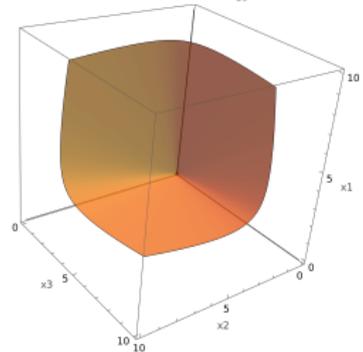
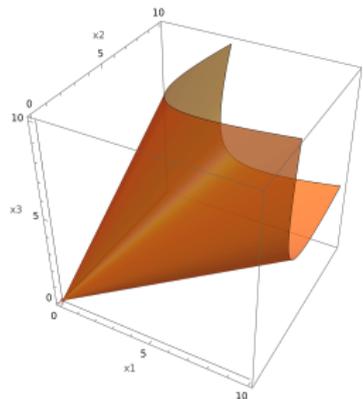
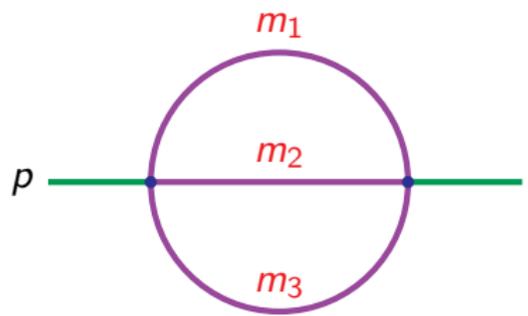
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- Thank you for listening!**