

HIGHER ORDER CORRECTIONS FOR JET OBSERVABLES IN POLARIZED DIS

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Tübingen University

Loops and Legs in Quantum Field Theory
Wittenberg - April 18th 2024

EBERHARD KARLS
UNIVERSITÄT
TÜBINGEN



HIGHER ORDER CORRECTIONS FOR JET OBSERVABLES IN POLARIZED DIS

Introduction

Why polarized processes?

Jet production in polarized DIS at NNLO

In collaboration with D. de Florian and I. Pedron

NLO+PS implementation for polarized DIS

In collaboration with B. Jäger

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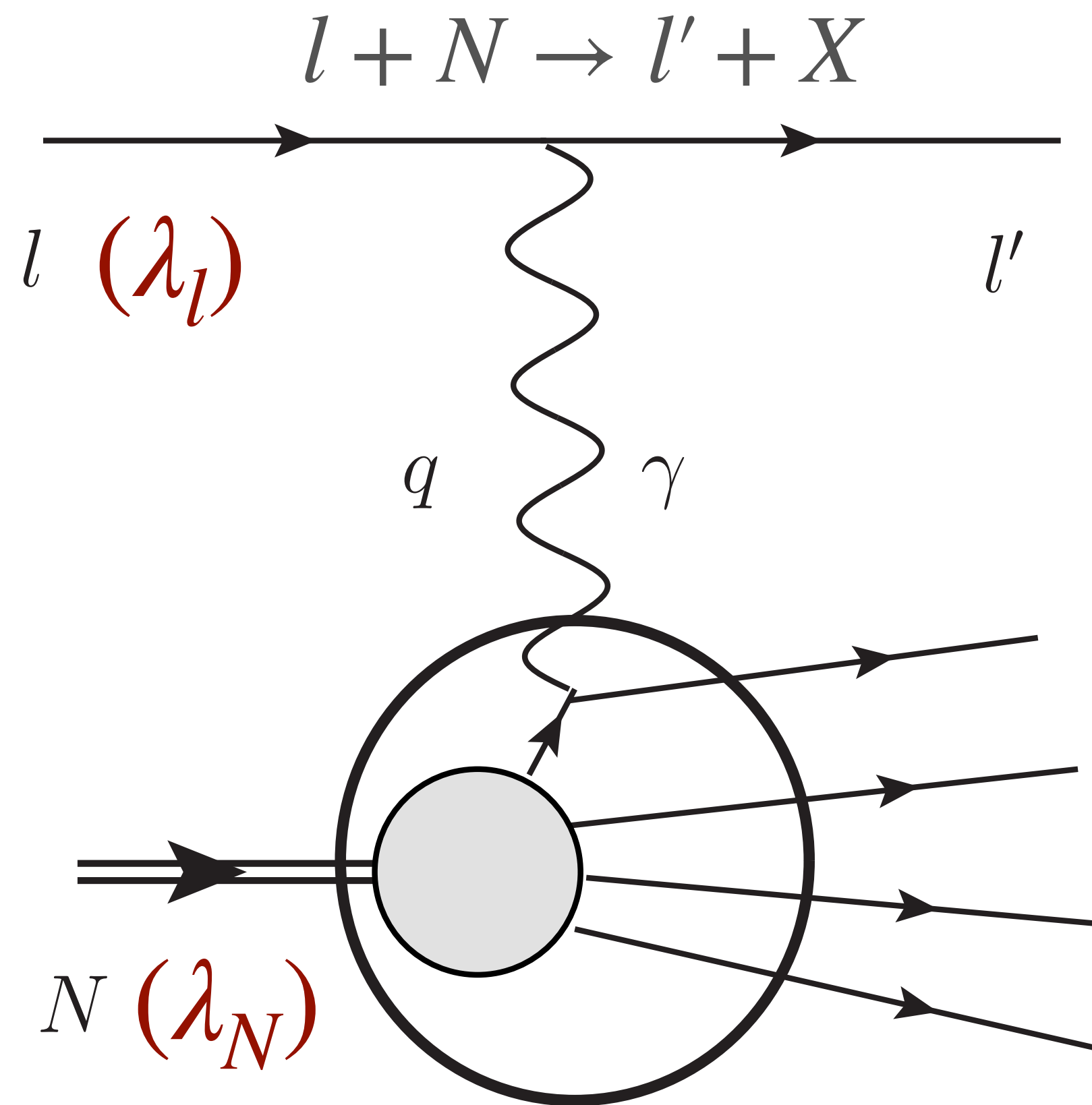
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MOTIVATION - THE PROTON'S SPIN STRUCTURE



$$\Delta\sigma \equiv \sigma^{\uparrow\uparrow} - \sigma^{\uparrow\downarrow}$$

$$\Delta\sigma = (\sigma^{++} + \sigma^{--}) - (\sigma^{+-} + \sigma^{-+})$$

$$\Delta\sigma = \sum_a \int dz \Delta f_a(z, \mu_F^2) \Delta \hat{\sigma}_i(\alpha_S(\mu_R), \mu_F^2, \mu_R^2)$$

Helicity PDFs

$$\Delta f_a \equiv f_a^{\uparrow} - f_a^{\downarrow}$$

Polarized Partonic Cross-Section

$$\Delta \hat{\sigma} \equiv \frac{1}{2} [\hat{\sigma}^{\uparrow\uparrow} - \hat{\sigma}^{\uparrow\downarrow}]$$

$$\Delta f_a(\mu_F^2) = \int_0^1 \Delta f_a(x, \mu_F^2) dx$$

Contribution of parton a to the proton's spin

$$\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta g + \mathcal{L}$$

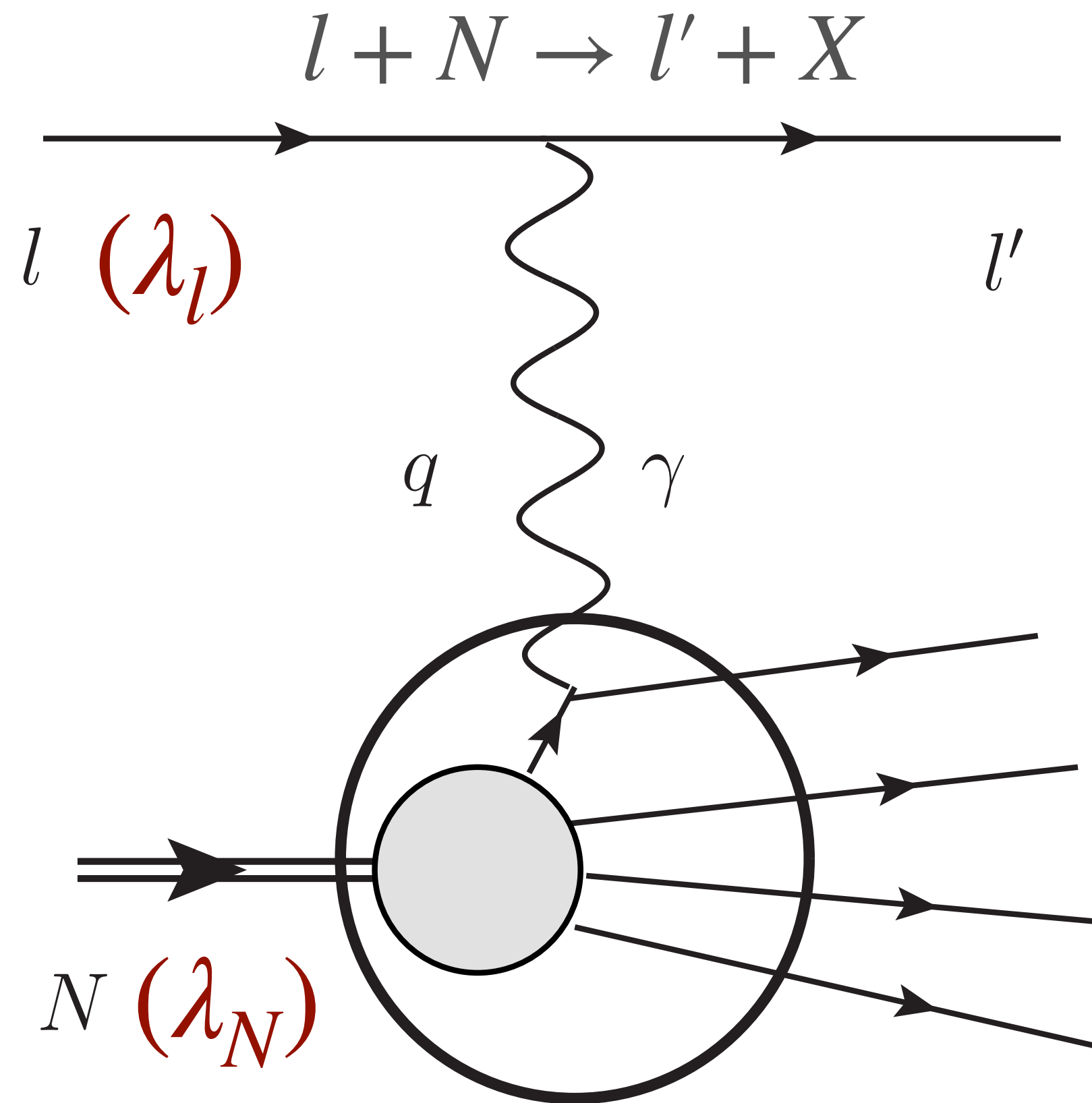
$$\Delta\Sigma \sim 0.25$$

European Muon Collaboration (1989) - "Proton spin crisis"

Evidence for positive contribution from Δg

De Florian, Sassot, Stratmann, Vogelsang (2014)

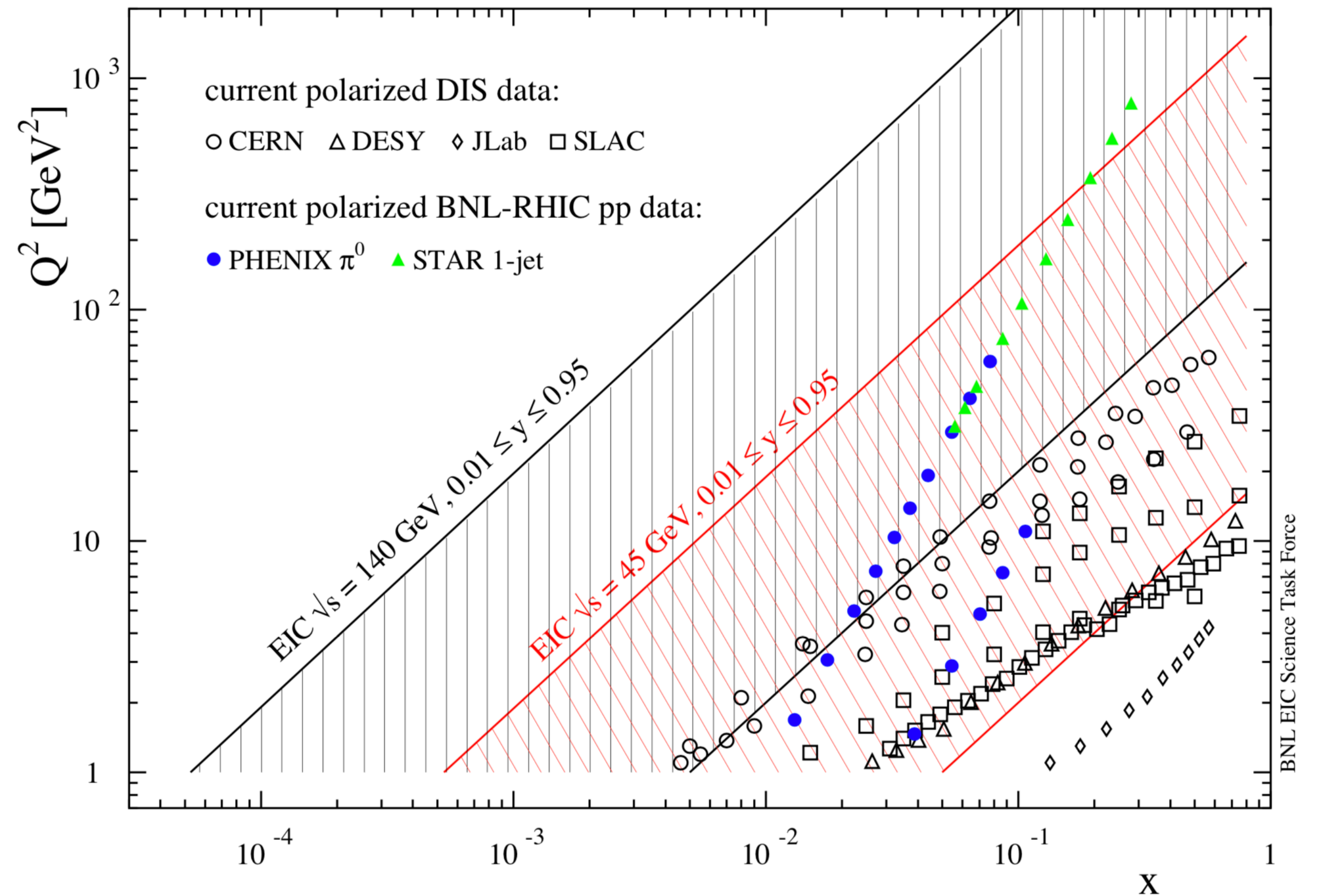
MOTIVATION - THE PROTON'S SPIN STRUCTURE



$$\Delta\sigma \equiv \sigma^{\uparrow\uparrow} - \sigma^{\uparrow\downarrow}$$

$$\Delta\sigma = (\sigma^{++} + \sigma^{--}) - (\sigma^{+-} + \sigma^{-+})$$

New constraints coming from the future Electron-Ion Collider (EIC)



STATUS - HIGHER ORDER CALCULATIONS FOR **POLARIZED** DIS

- NNLO structure functions g_1 (photon exchange)
van Neerven, Zijlstra (1994)
- NLO Photoproduction single-jet (small-cone approximation)
Jäger (2008)
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Hinderer, Schlegel, Vogelsang (2017)
- NLO Single-Jet production (polarized N-jetiness)
Boughezal, Petriello, Xing (2018)
- NNLO Single-Jet production (polarized Catani-Seymour Dipoles+Projection to Born)
Photon- IB, de Florian, Pedron (2020)
- NNLO NC & CC structure functions g_1, g_4, g_5
IB, de Florian, Pedron (2022)
- NNLO Single-Jet production (polarized Catani-Seymour Dipoles+Projection to Born)
NC and CC- IB, de Florian, Pedron (2023)
- (approx. N3LO) Semi-Inclusive DIS
Abele, de Florian, Vogelsang (2022)
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Blümlein, Marquard, Schneider, Schönwald (2022)
[See J. Blümlein talk for HF corrections]

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Bonino, Gehrmann, Löchner, Schönwald, Stagnitto (2024)
Goyal, Moch, Pathak, Rana, Ravindran (2024)

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IMPLEMENTATION - NLO: DIPOLE SUBTRACTION

Based on the construction of counter-terms $d\sigma^A$, that satisfy:

- **Exactly** reproduce the divergent behavior of the real emission part
- Simple enough to be **integrated analytically** over the divergent part of the phase space

$$\sigma^{NLO} = \int_{m+1} \left[(d\sigma^R)_{\epsilon=0} - (d\sigma^A)_{\epsilon=0} \right] + \int_m \left[\int_1 d\sigma^A + d\sigma^V \right]_{\epsilon=0}$$

Polarized Dipole subtraction

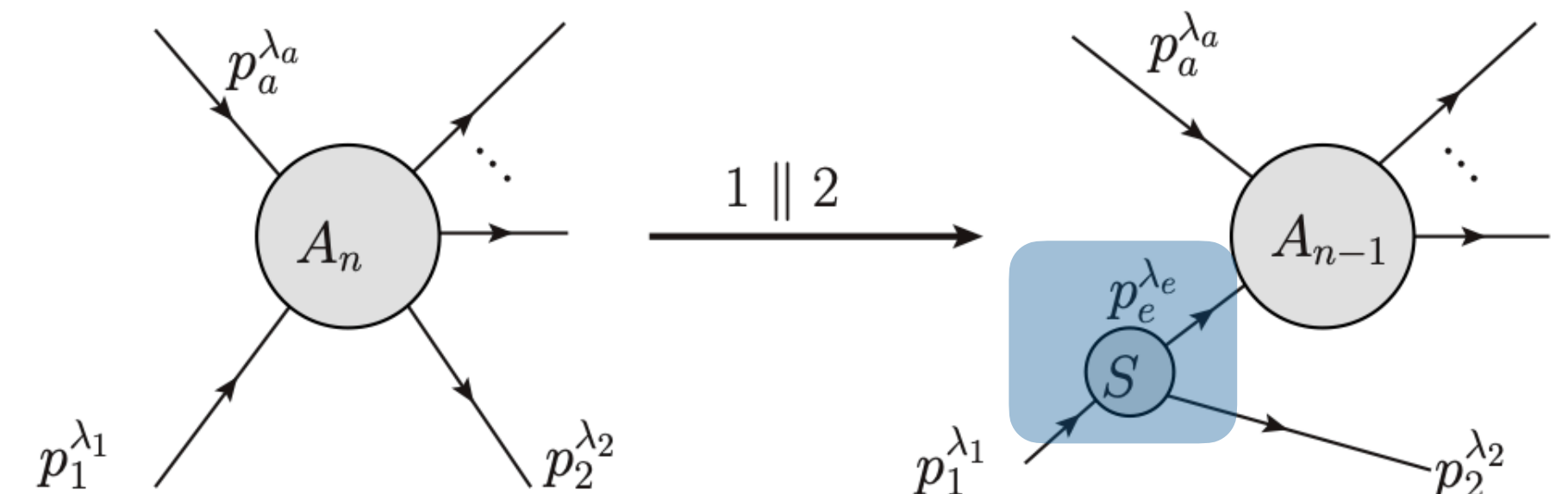
Based on the **polarized** dipole factorization formula

$$d\Delta\hat{\sigma}_{(m+1)}^{NLO} = d\Delta\hat{\sigma}_{(m)}^{\text{Born}} \otimes d\Delta V_{\text{dipole}}$$

↓
Universal dipole factor ; Interpolates between collinear AP splittings and soft eikonal factor

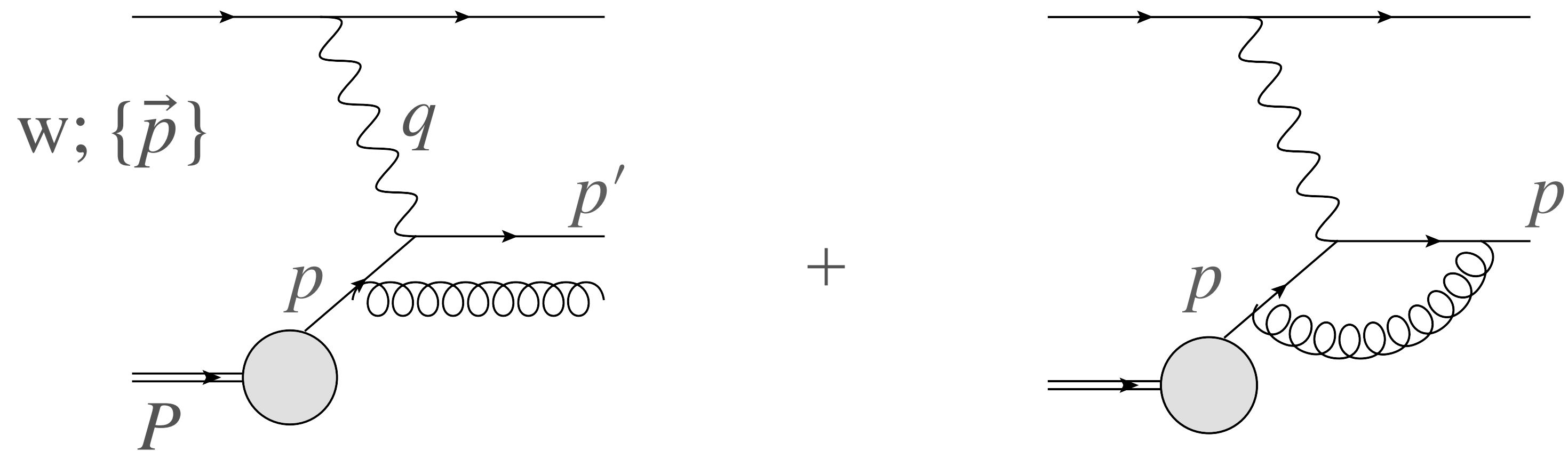
$$dA = \sum d\Delta\hat{\sigma}^{\text{Born}} \otimes d\Delta V_{\text{dipole}} \rightarrow \text{General NLO subtraction scheme for (initially) polarized processes}$$

IB, de Florian, Pedron. *Phys.Rev.Lett.* 125 (2020)
IB, de Florian, Pedron. *Phys.Rev.D* 103 (2021)



IMPLEMENTATION - NNLO: PROJECTION TO BORN (P2B)

Allows to obtain fully differential cross section for \mathcal{O} if the **inclusive cross section for \mathcal{O} at that order** and for **\mathcal{O} +jet at the previous order** are known

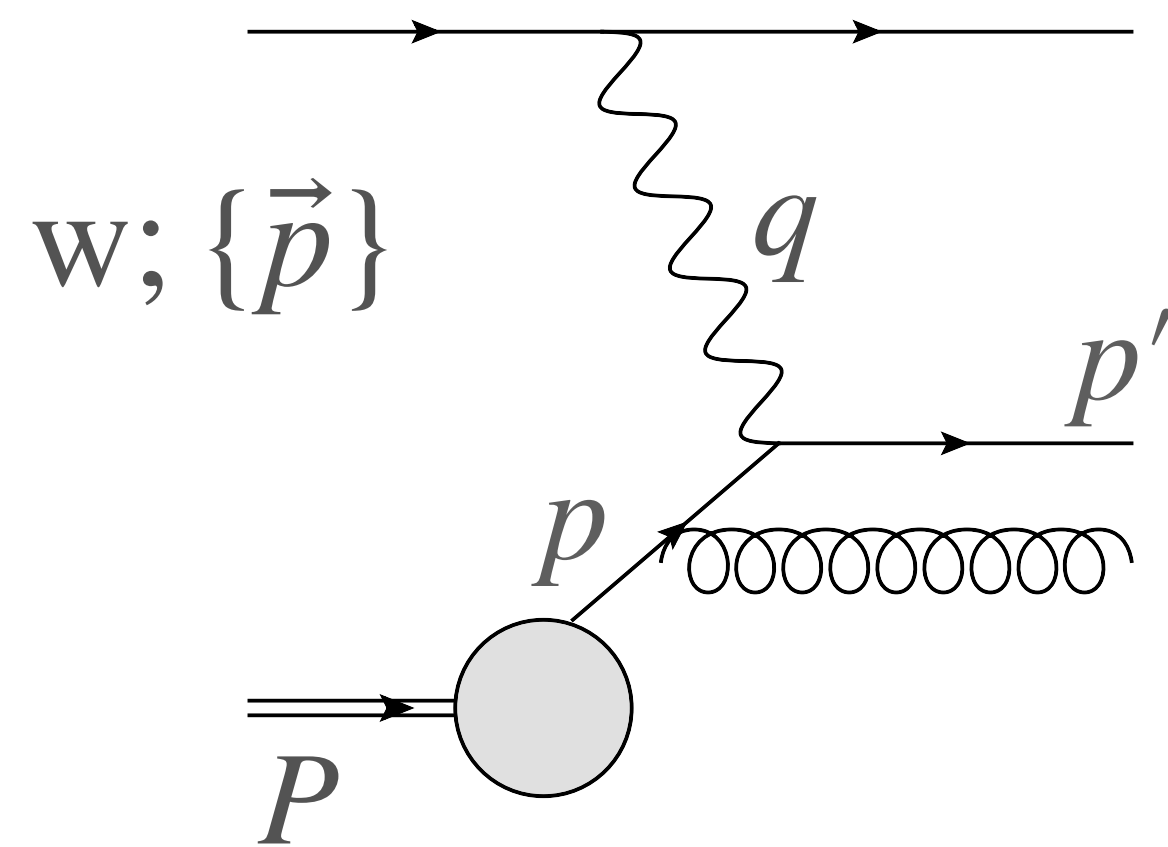


Born kinematic can be inferred from non-QCD particles

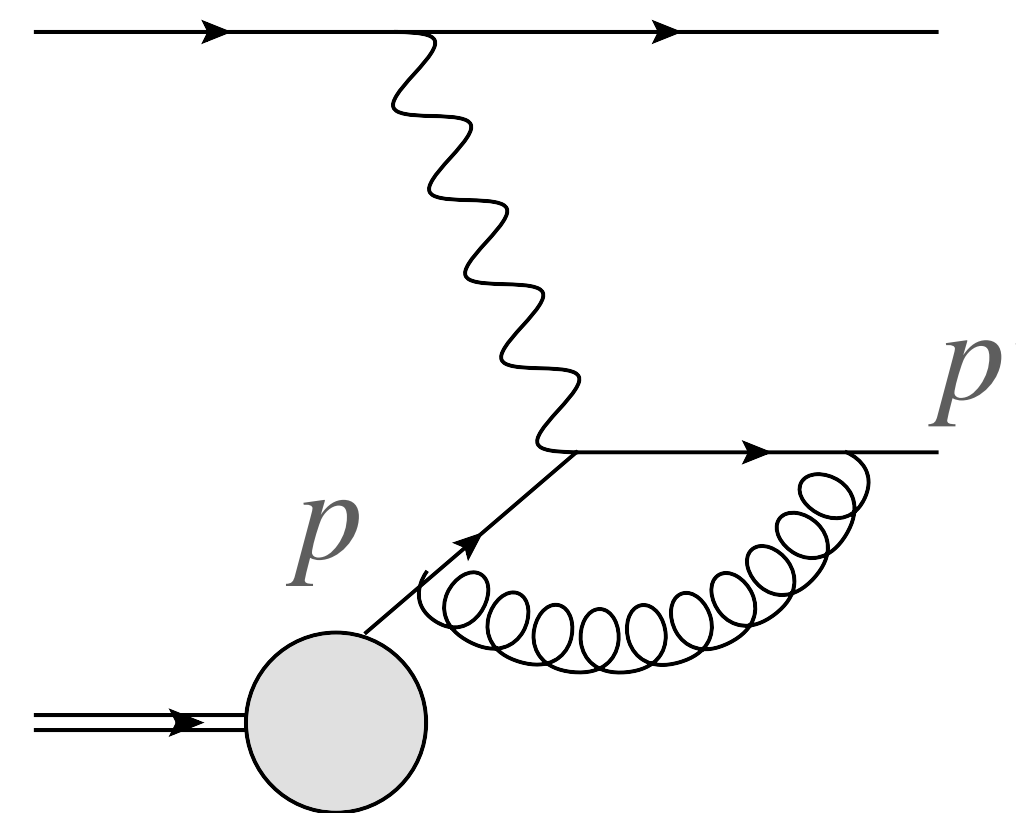
$$p_B = xP$$
$$p'_B = xP + q$$

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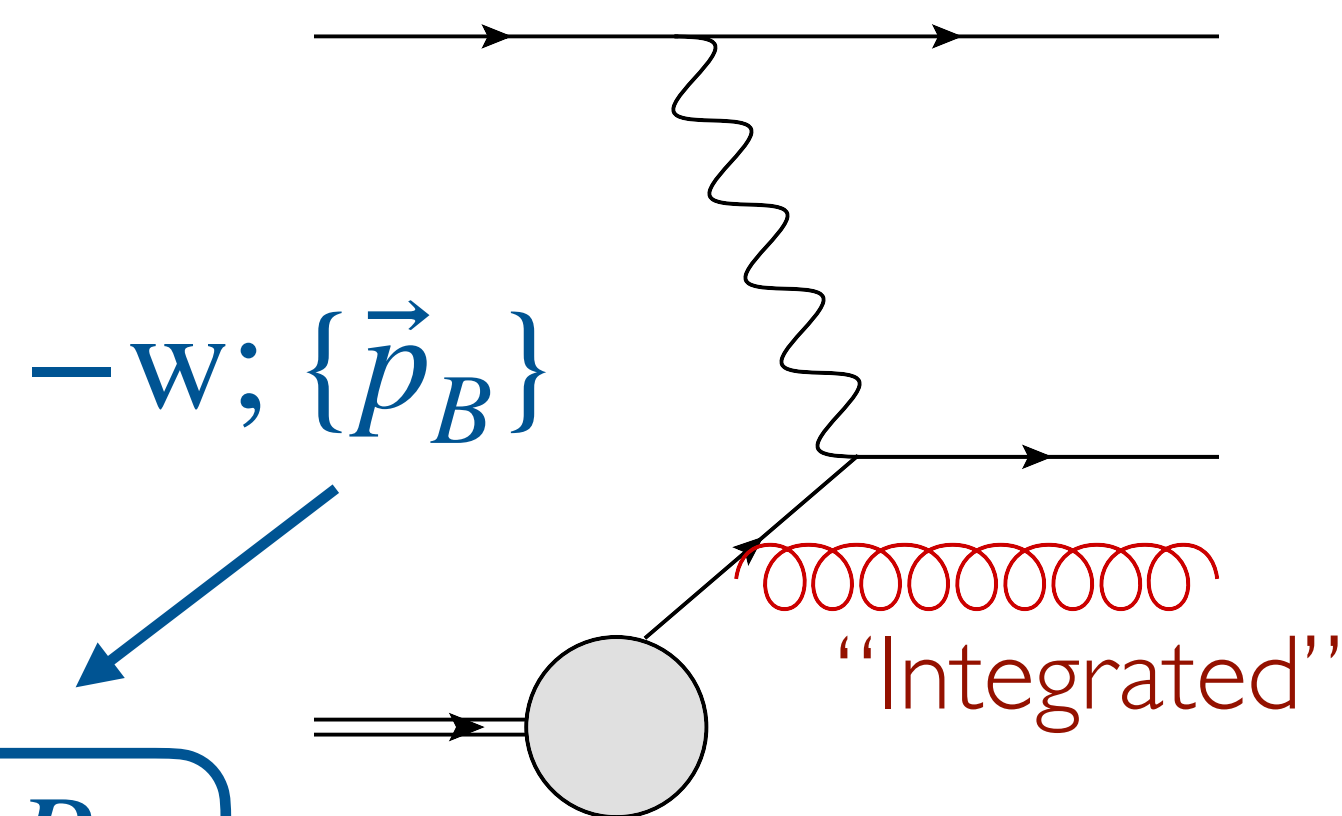
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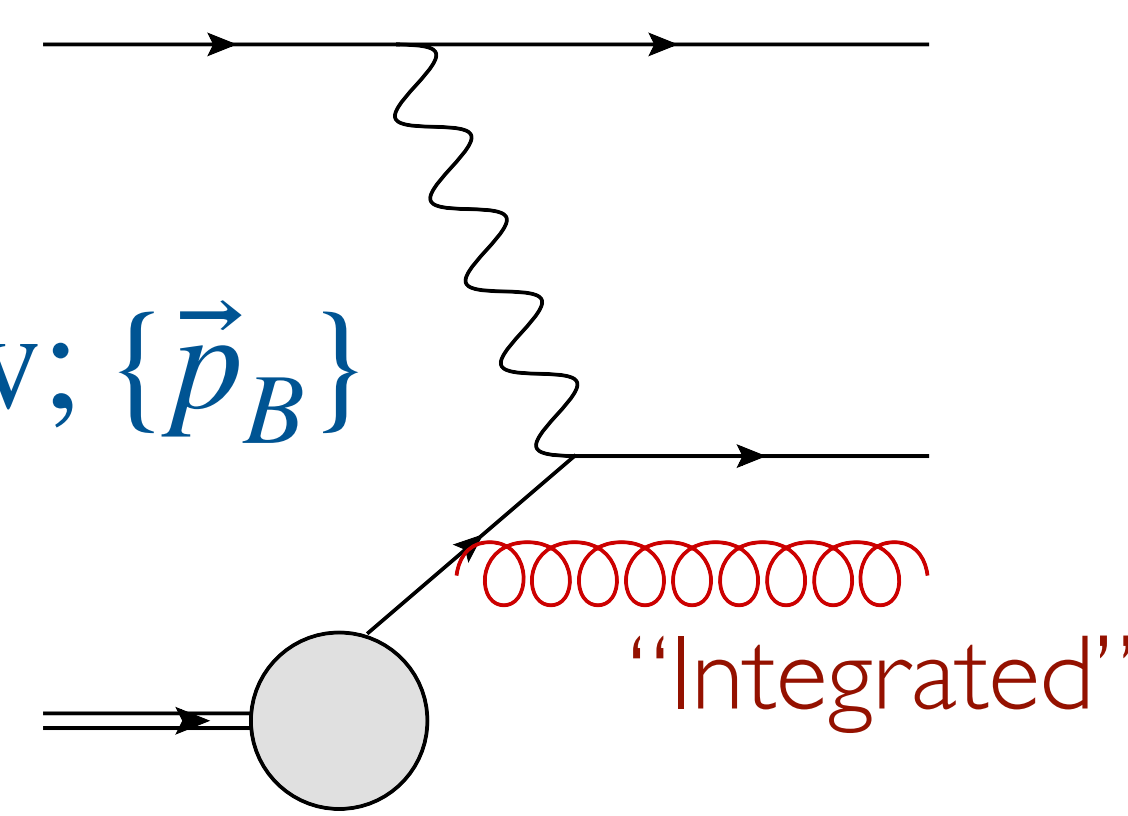
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$+w; \{\vec{p}_B\}$



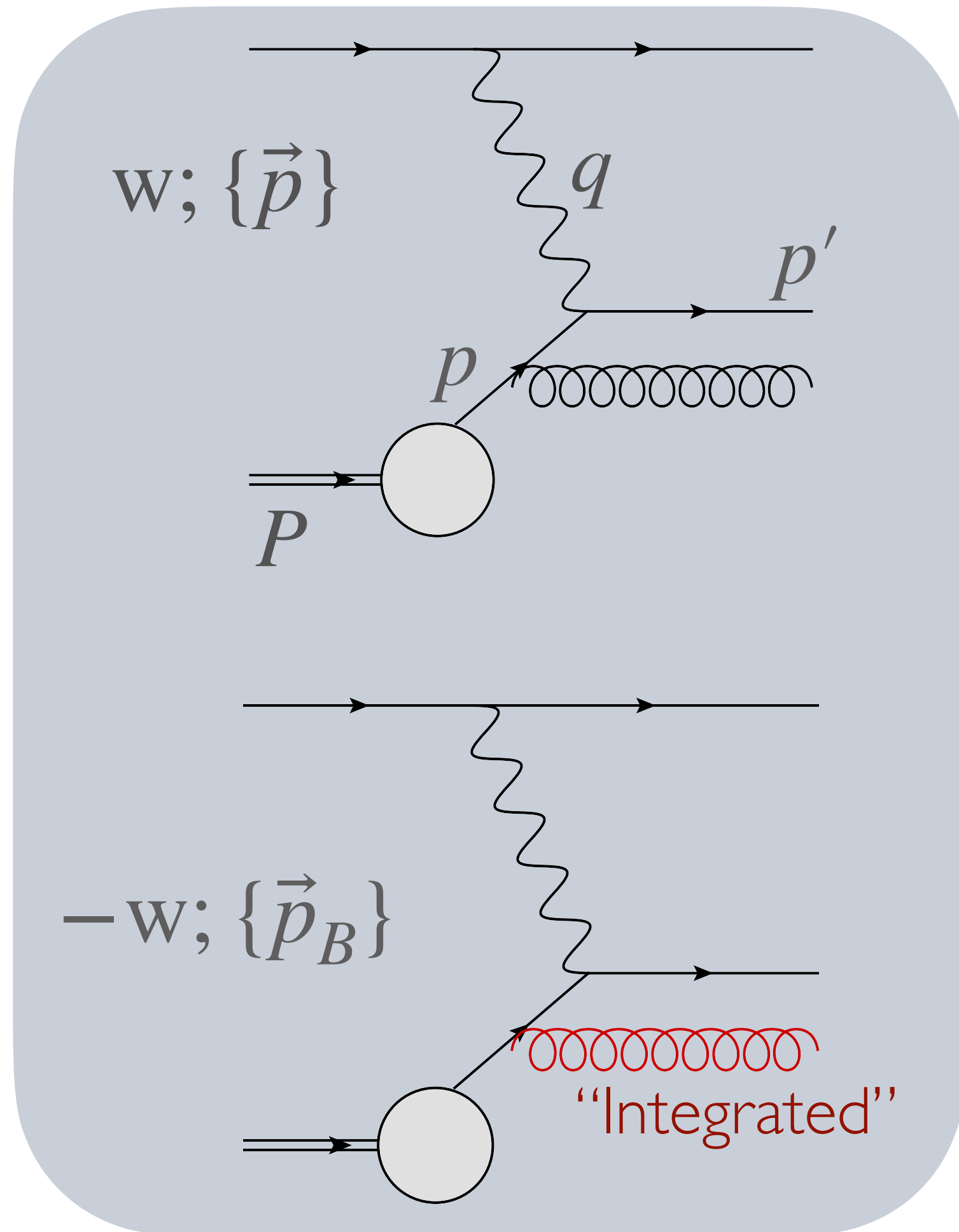
Same weights as the real emission process, but **binned with Born kinematics**

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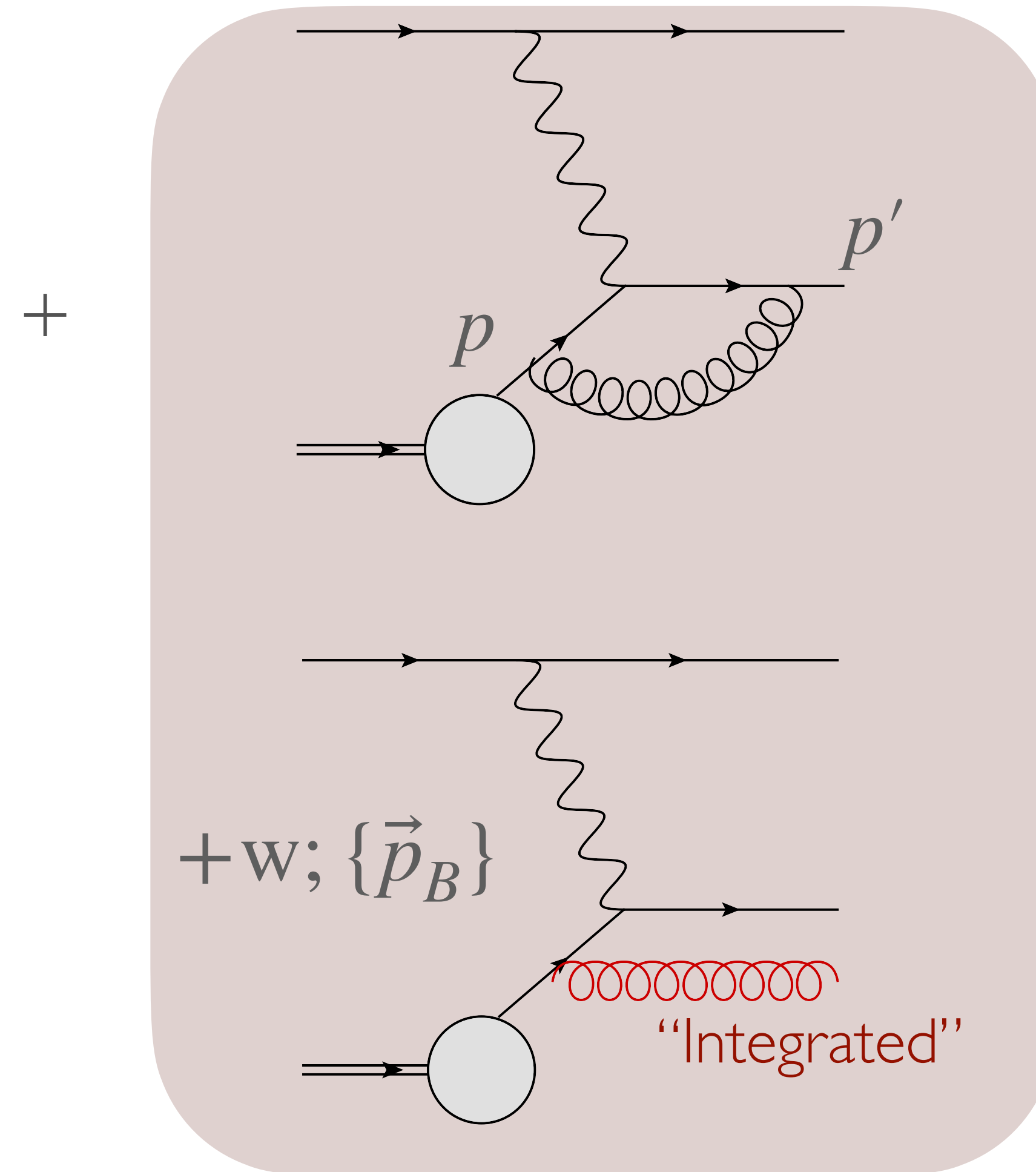
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Finite & Integrable in 4D



Contribution to the inclusive cross section



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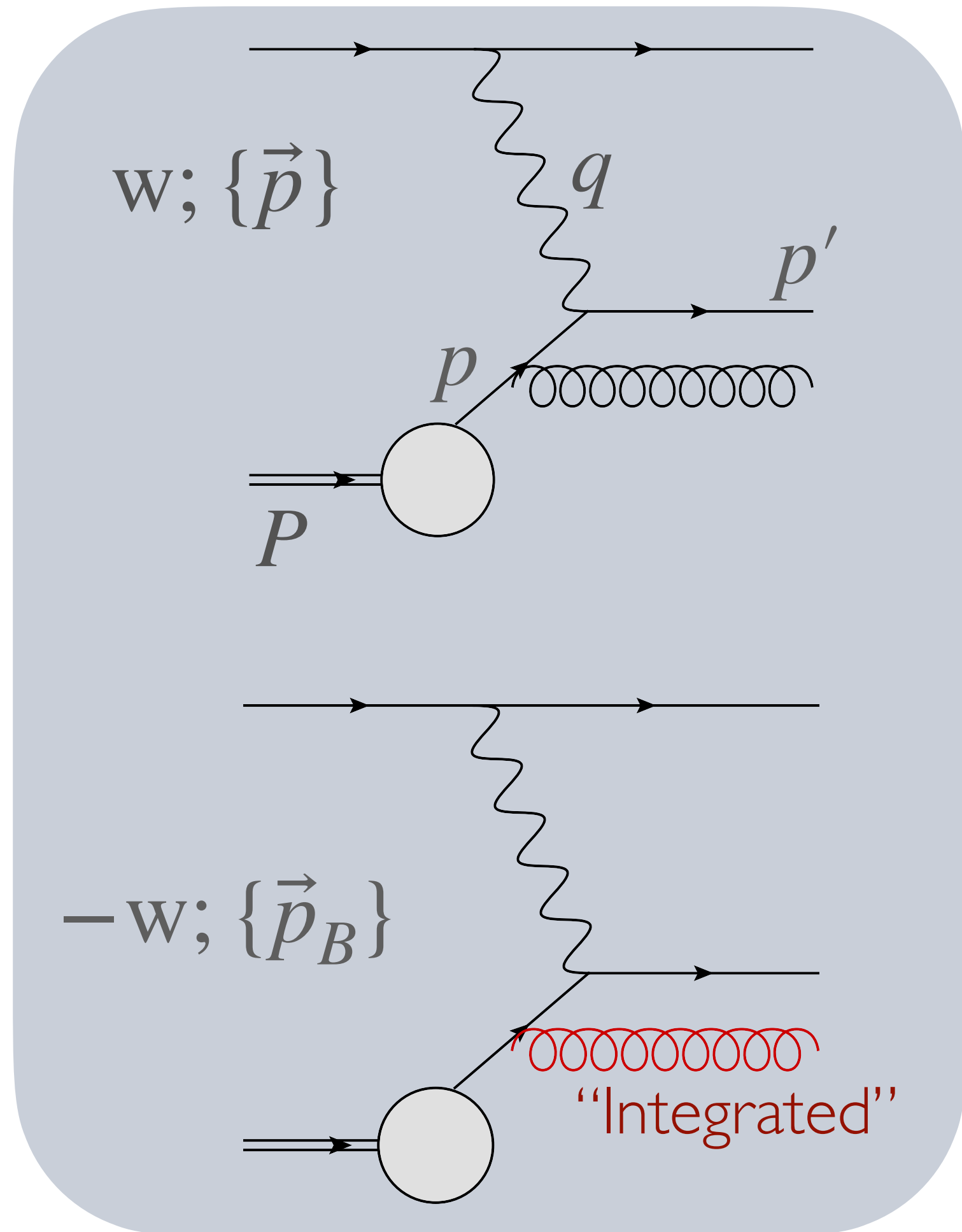
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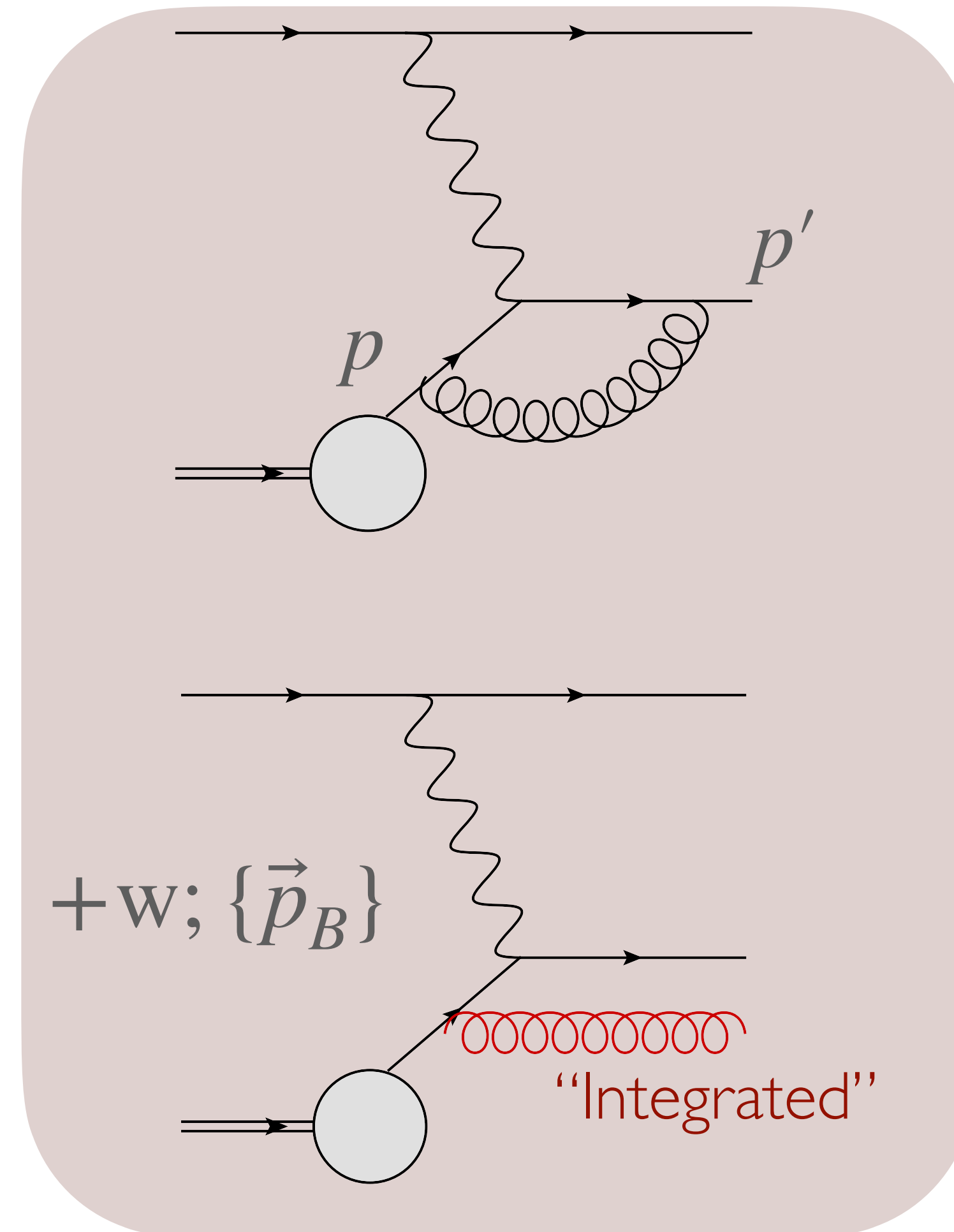
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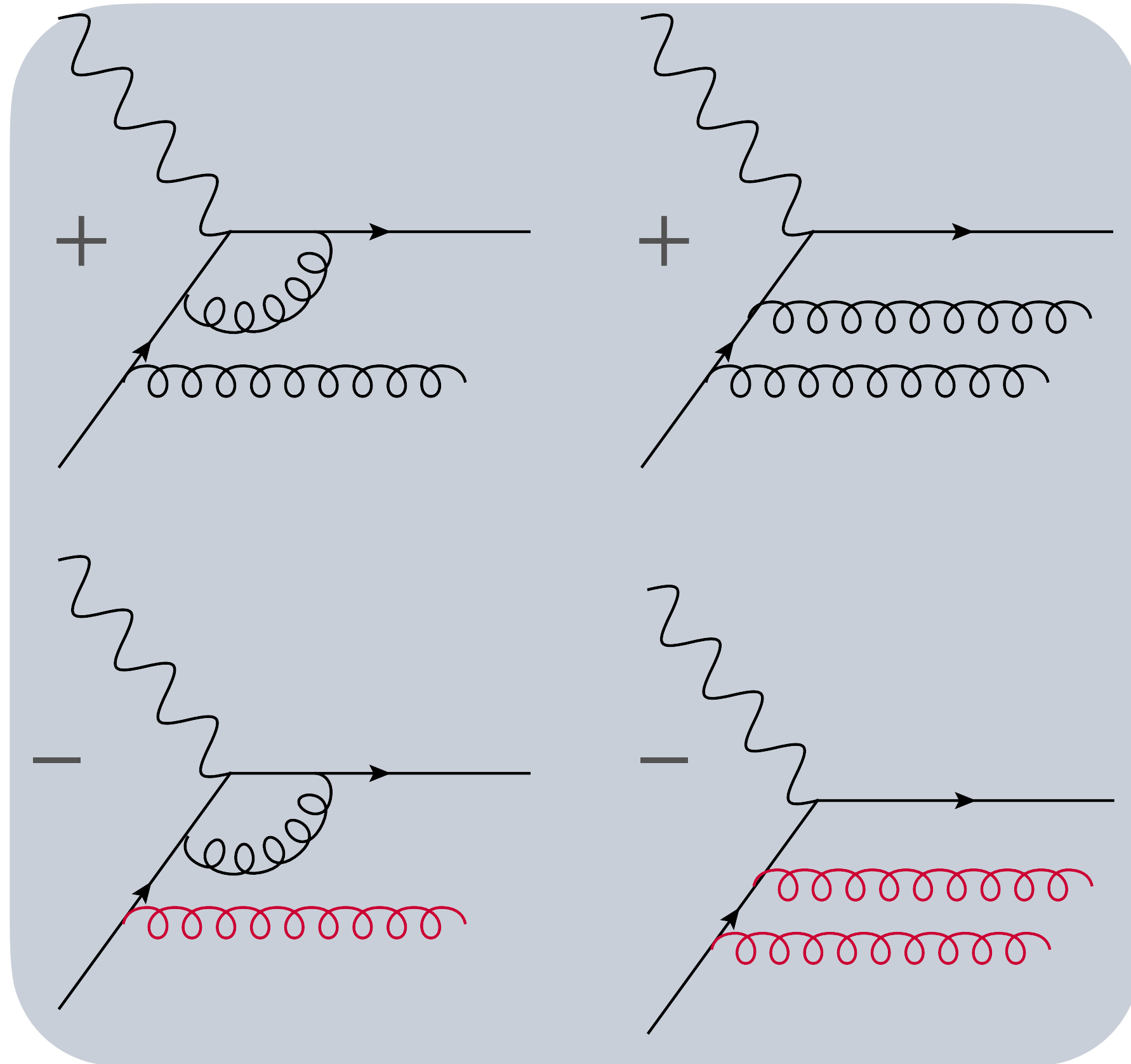
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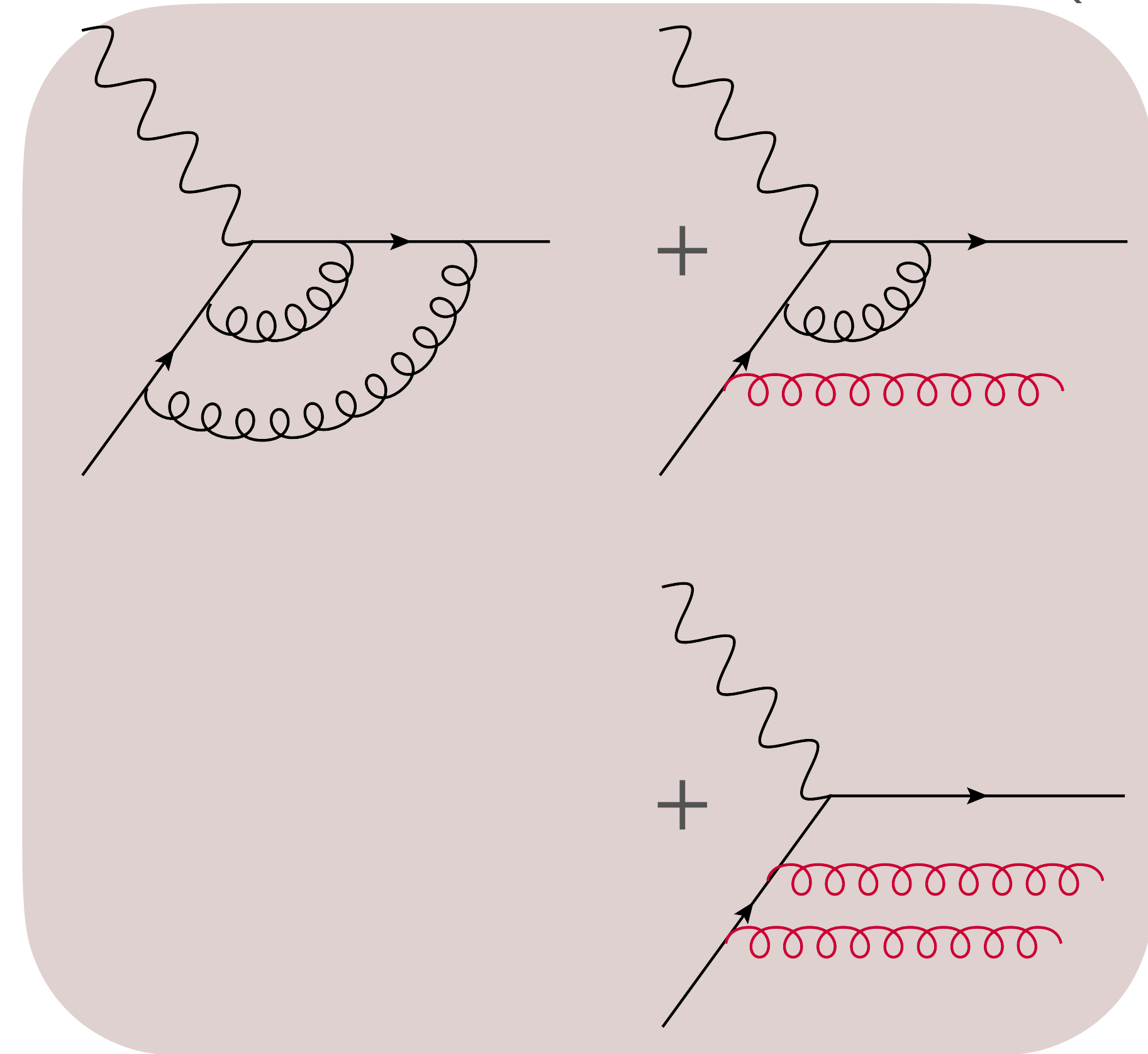
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$$d\sigma_{1\text{jet}}^{\text{NLO}} = d\sigma_{2\text{jets}}^{\text{LO}} - d\sigma_{2\text{jets P2B}}^{\text{LO}} + d\sigma_{1\text{jet P2B}}^{\text{NLO,incl}}$$

IMPLEMENTATION - NNLO: PROJECTION TO BORN (P2B)



Finite & Integrable in 4D

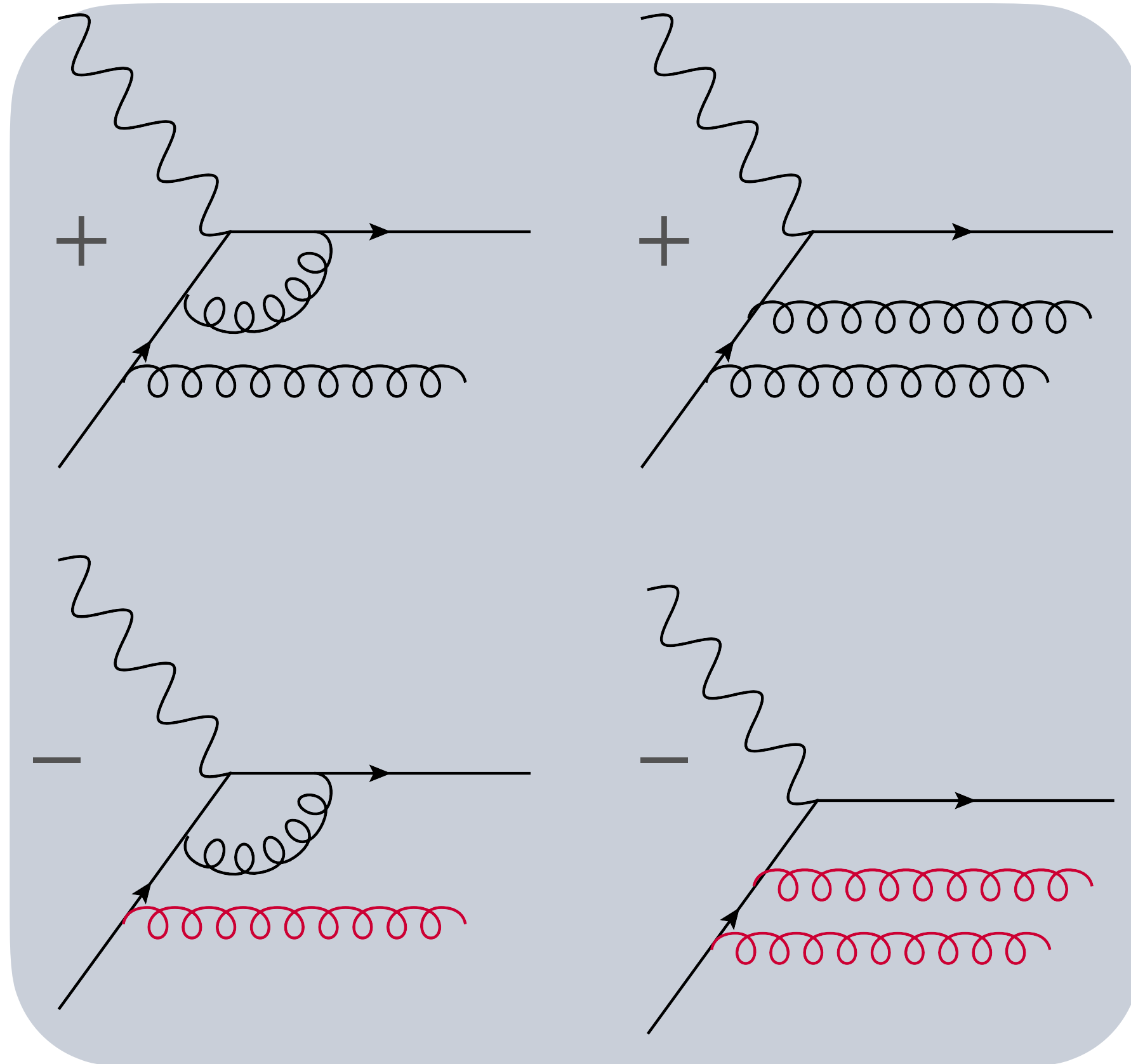


Contribution to the inclusive cross section

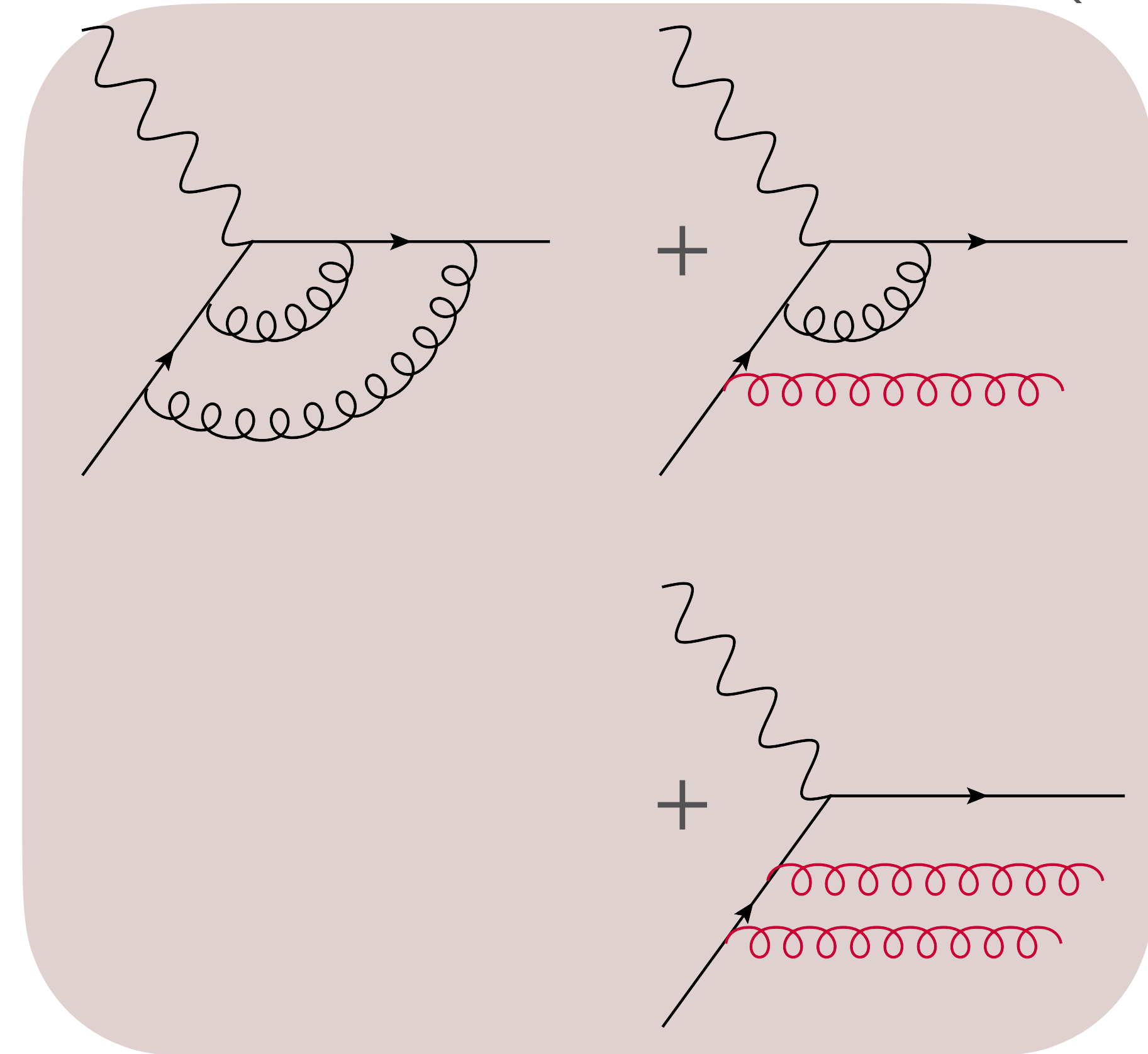
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Can compute the **NNLO cross section for 1-jet**, given that the **NLO calculation for 2-jets** and the **NNLO inclusive cross sections** are available

IMPLEMENTATION - NNLO: PROJECTION TO BORN (P2B)



Finite & Integrable in 4D



Contribution to the inclusive cross section

Calculated using polarized dipoles

$$d\sigma_{1\text{jet}}^{\text{NNLO}} = d\sigma_{2\text{jets}}^{\text{NLO}} - d\sigma_{2\text{jets P2B}}^{\text{NLO}} + d\sigma_{1\text{jet P2B}}^{\text{NNLO,incl}}$$

van Neerven, Zijlstra (1994)

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IMPLEMENTATION - POLDIS CODE

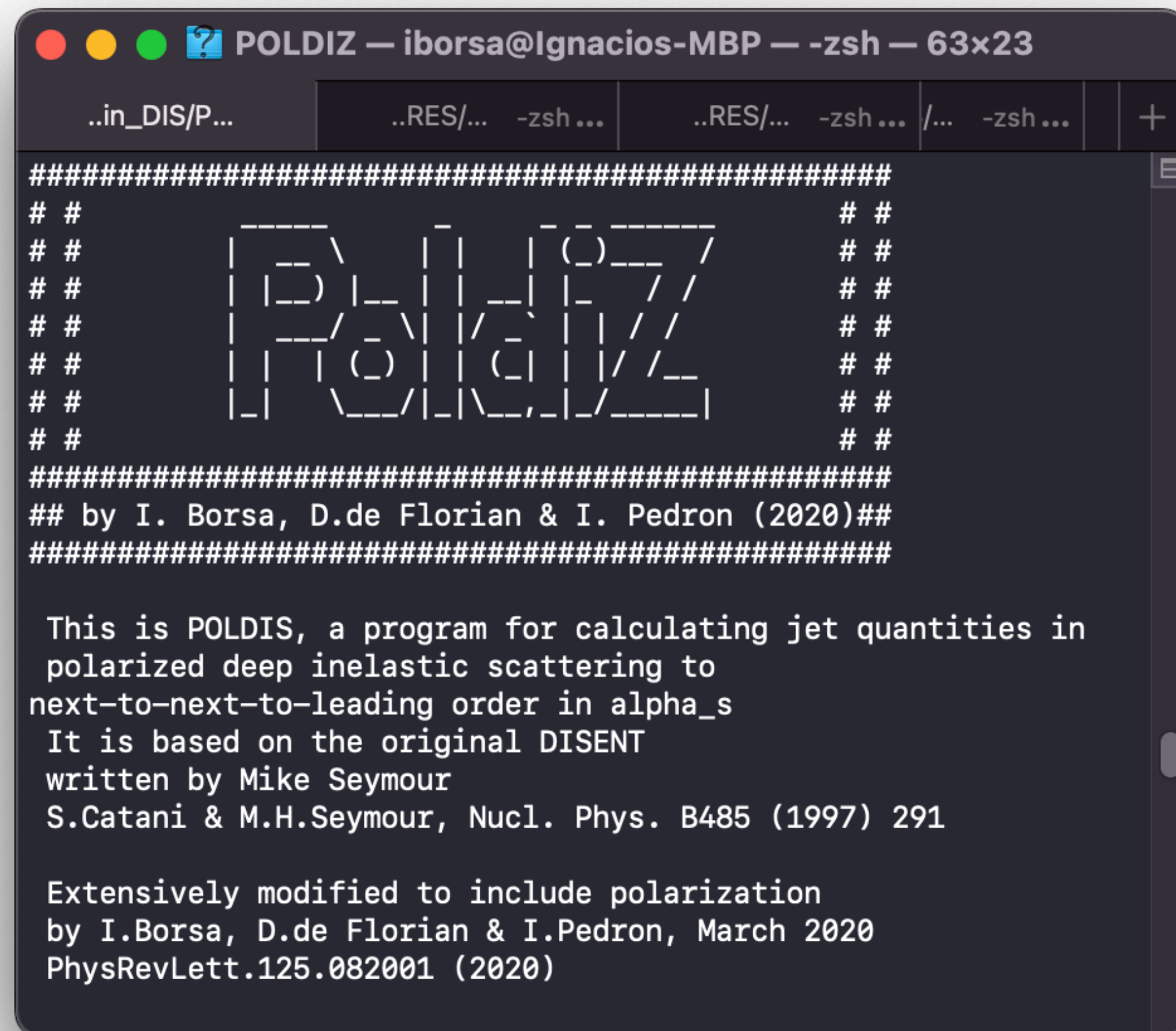
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Based on DISENT → Fixed bug in the gluon channel

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POLDIZ — iborsa@Ignacios-MBP — -zsh — 63x23
..in_DIS/P... ..RES/... -zsh... ..RES/... -zsh... /... -zsh... +
#####
# #                               # #
# #                               # #
# #                               # #
# #                               # #
# #                               # #
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#####
## by I. Borsa, D.de Florian & I. Pedron (2020)##
#####

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written by Mike Seymour
S.Catani & M.H.Seymour, Nucl. Phys. B485 (1997) 291

Extensively modified to include polarization
by I.Borsa, D.de Florian & I.Pedron, March 2020
PhysRevLett.125.082001 (2020)
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POLDIS: Photon exchange

- Completely differential 1-jet observables up to NNLO
- Completely differential 2-jet observables up to NLO

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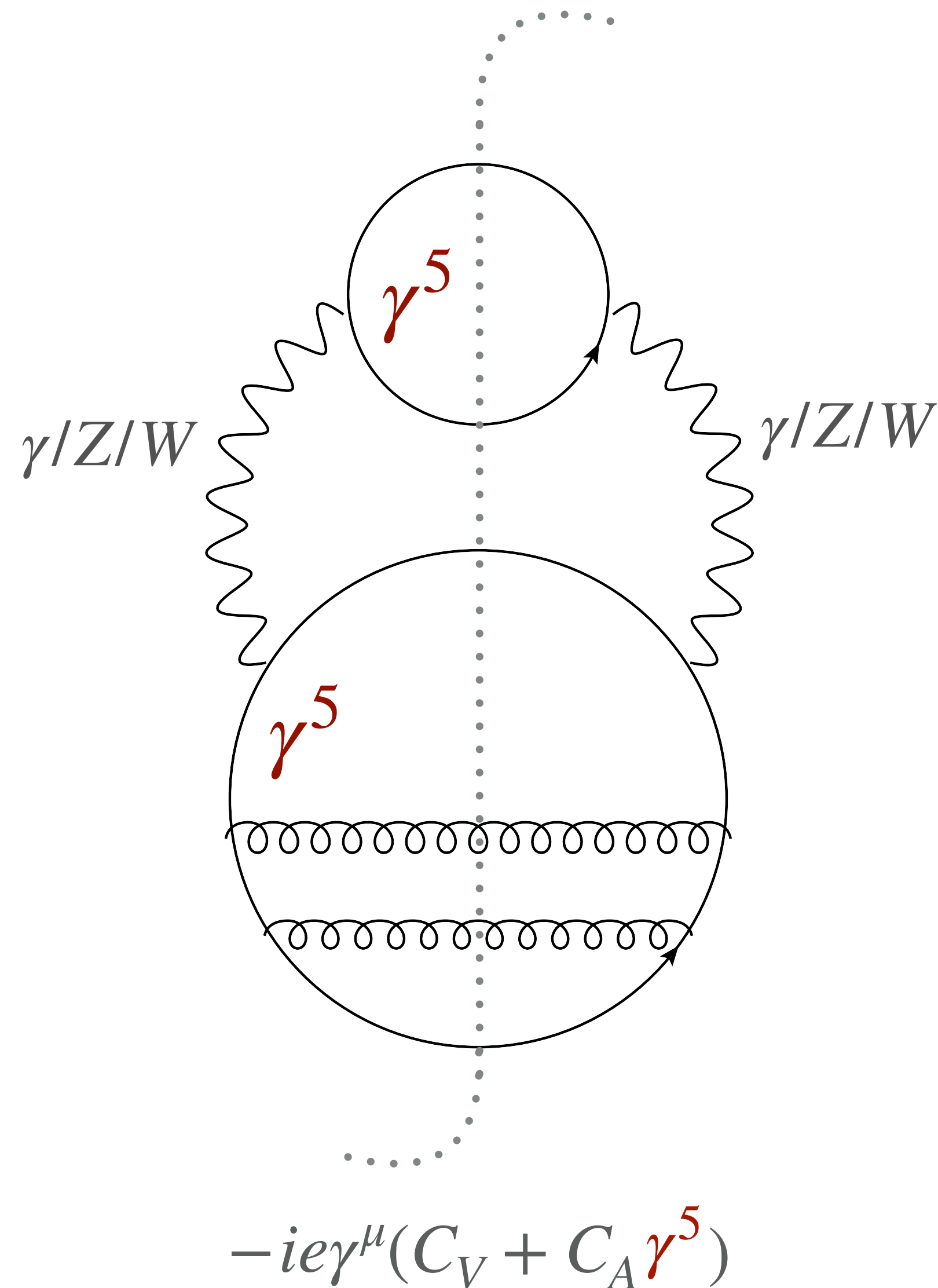
- + Full Neutral current
- + Charged current

IB, de Florian, Pedron. *Phys.Rev.D* 105 (2022)

IB, de Florian, Pedron. *Phys.Rev.D* 107 (2023)

IMPLEMENTATION - INCLUSION OF EW CURRENTS

$$\hat{\sigma}_q = \hat{\sigma}_q^{\text{PV}} + \hat{\sigma}_q^{\text{NPV}}$$



HVBM scheme

- **Quark channel:** Relations between polarized and unpolarized contributions

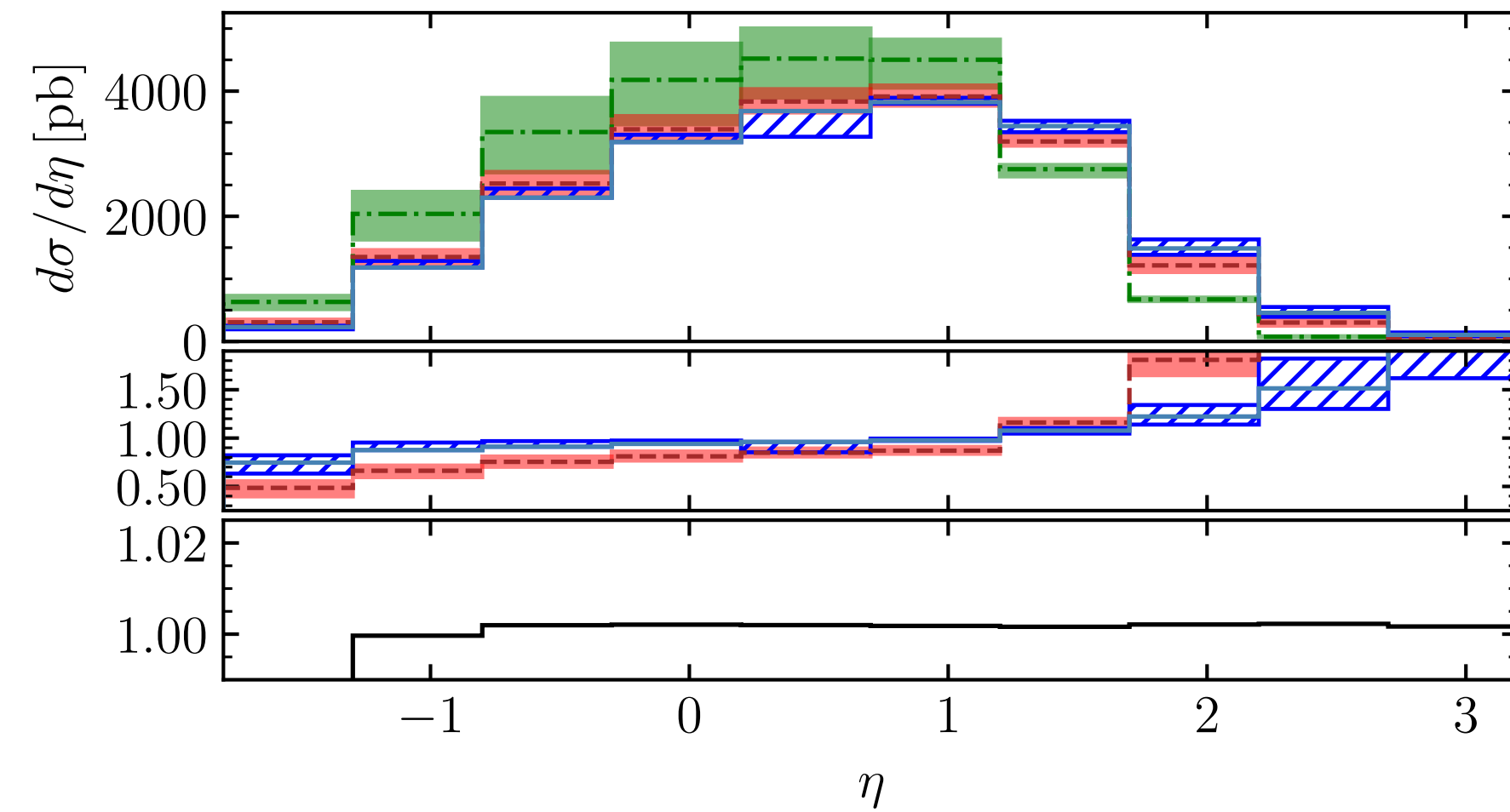
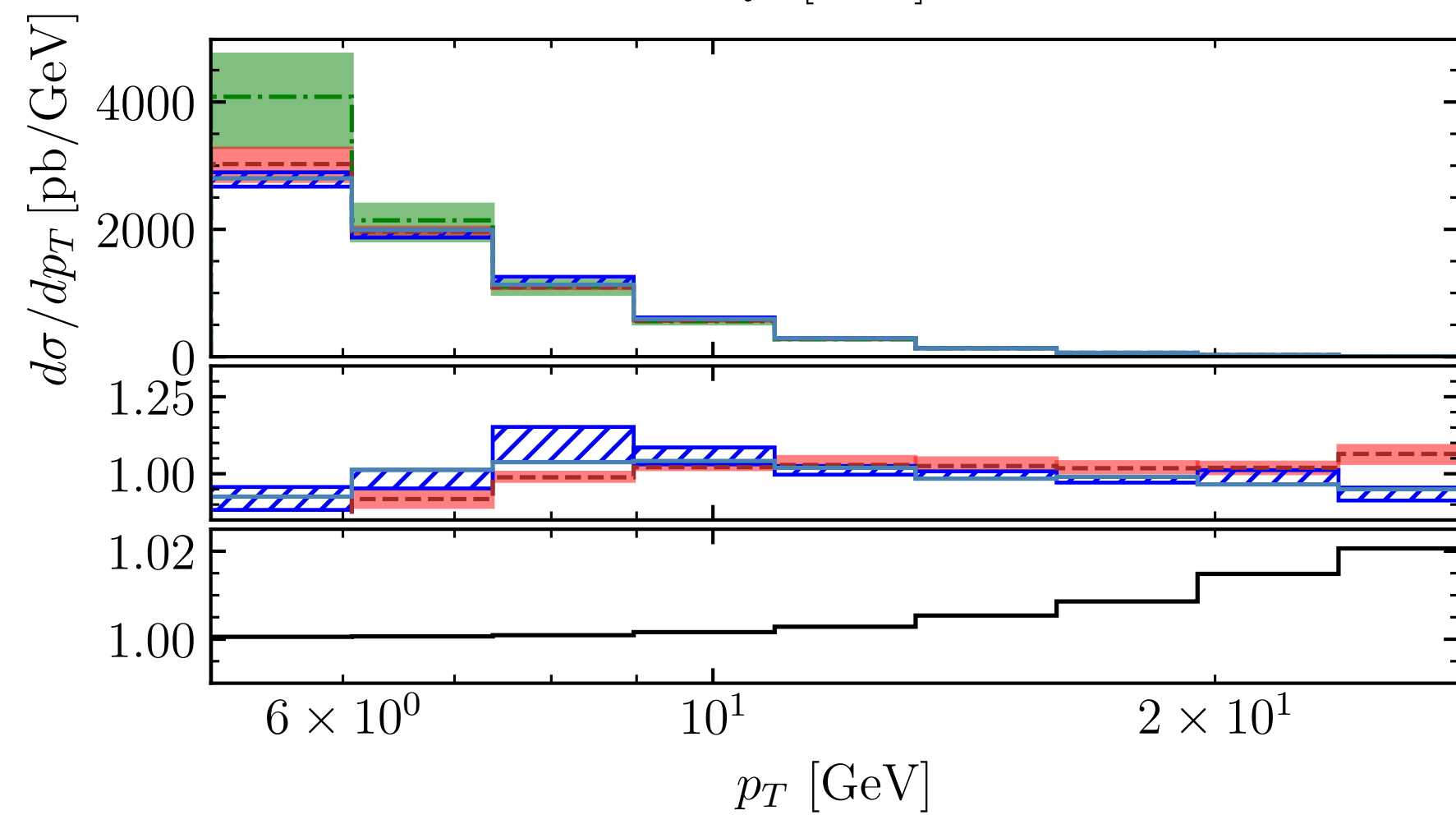
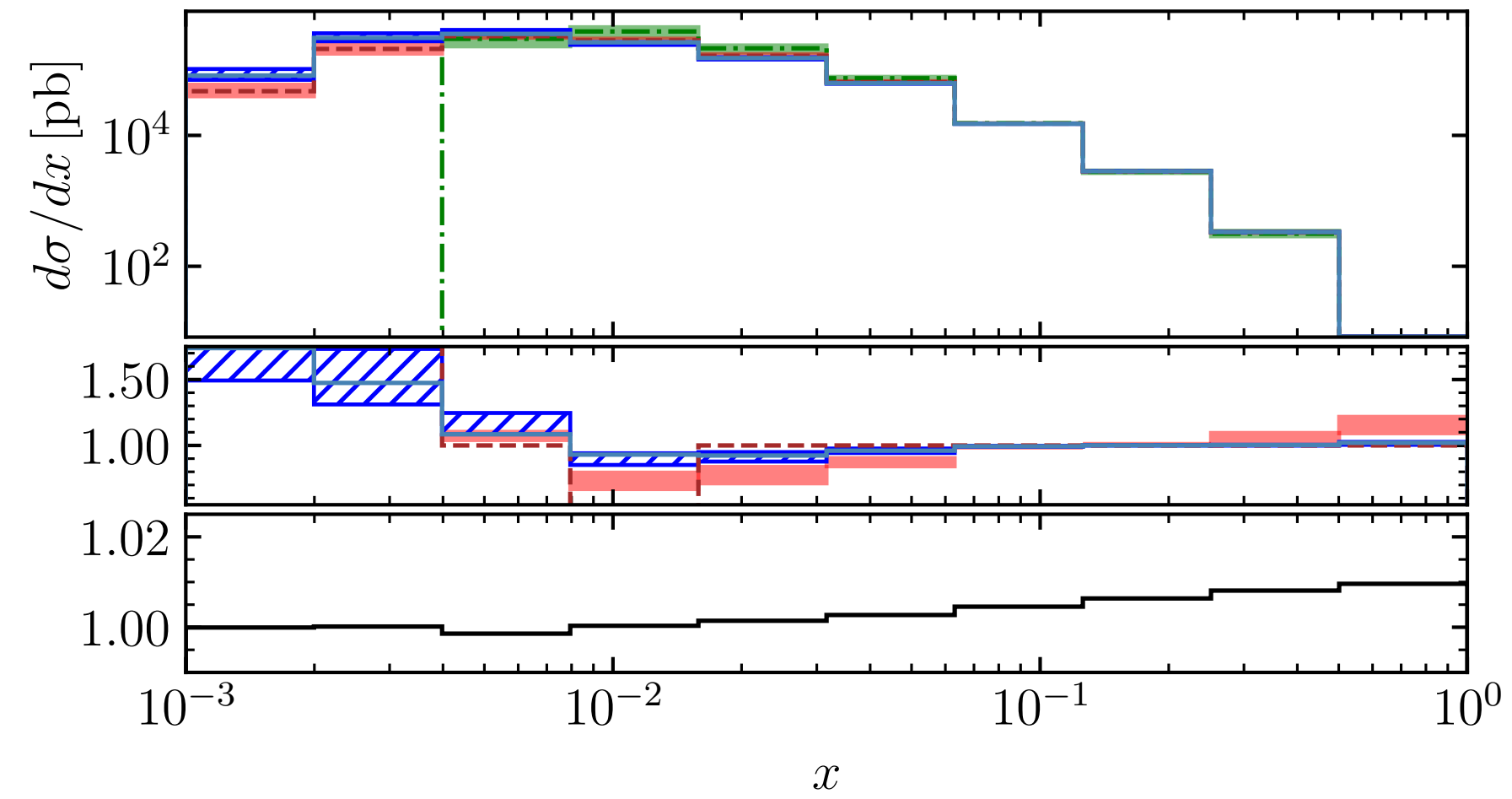
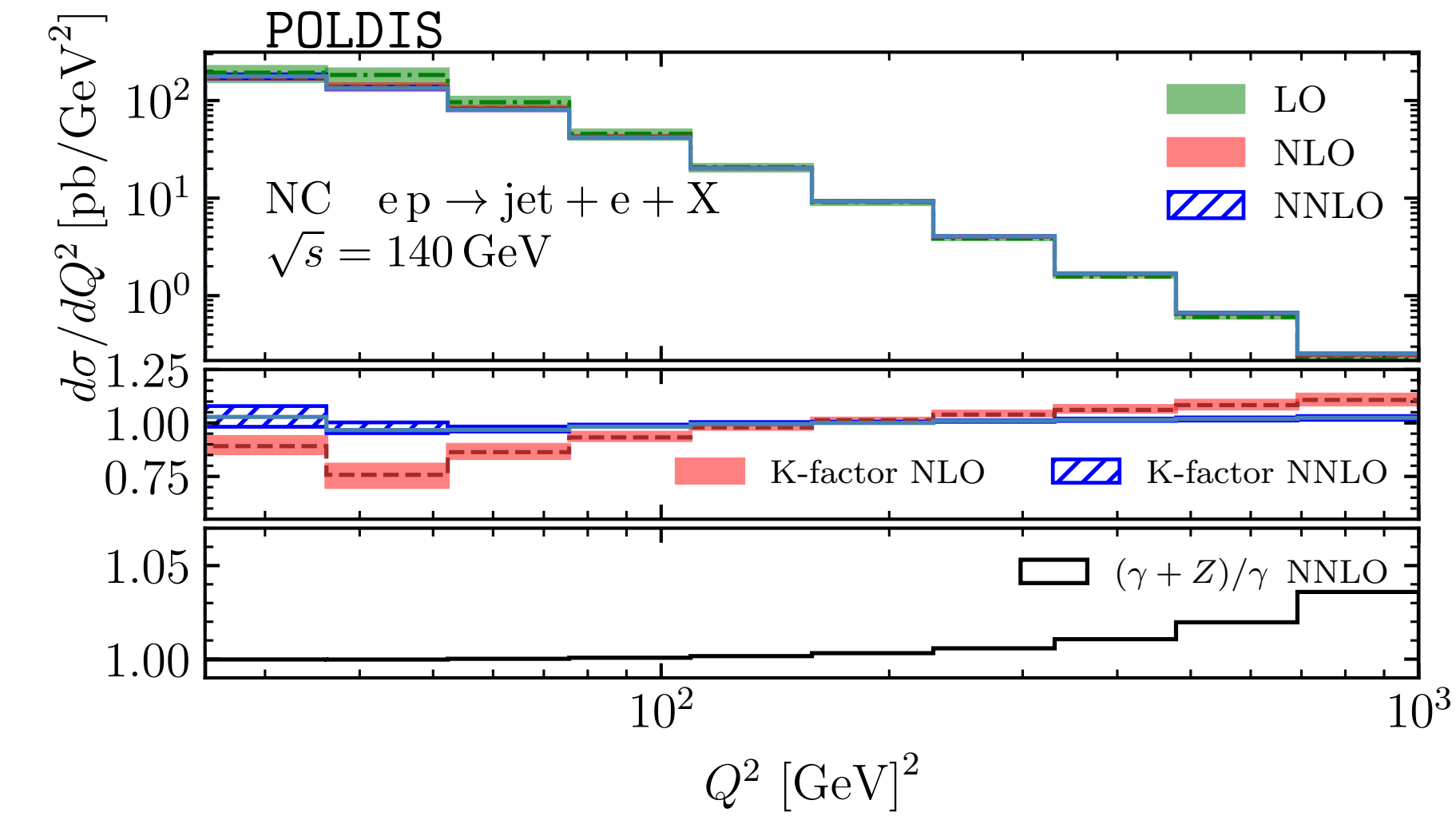
$$\begin{aligned} \text{ODD } \gamma^5 : & \quad \hat{\sigma}_q^{\text{PV}} \Leftrightarrow \Delta\hat{\sigma}_q^{\text{NPV}} \\ \text{EVEN } \gamma^5 : & \quad \Delta\hat{\sigma}_q^{\text{PV}} \Leftrightarrow \hat{\sigma}_q^{\text{NPV}} \end{aligned}$$

- Easy to see for real emission contributions (4 dimensional)
- Virtual contributions require additional finite renormalization from HVBM
- **Gluon channel:** Contribution cancels if the charged of the jet is not identified

Parity-violating structure functions g_4, g_5
[IB, de Florian, Pedron \(2022\)](#)

PHENOMENOLOGY - NNLO SINGLE-JET PRODUCTION IN DIS

Neutral Current DIS

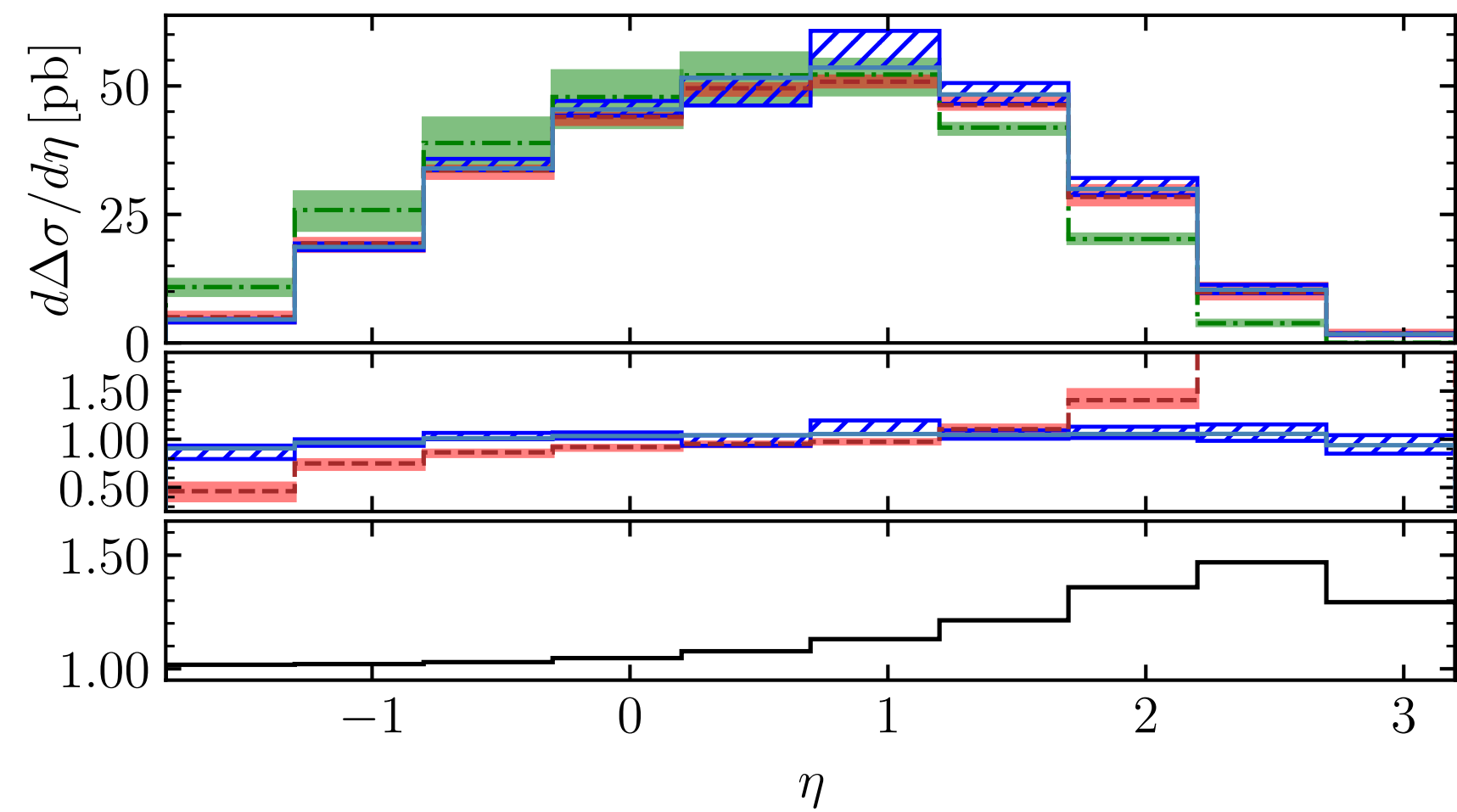
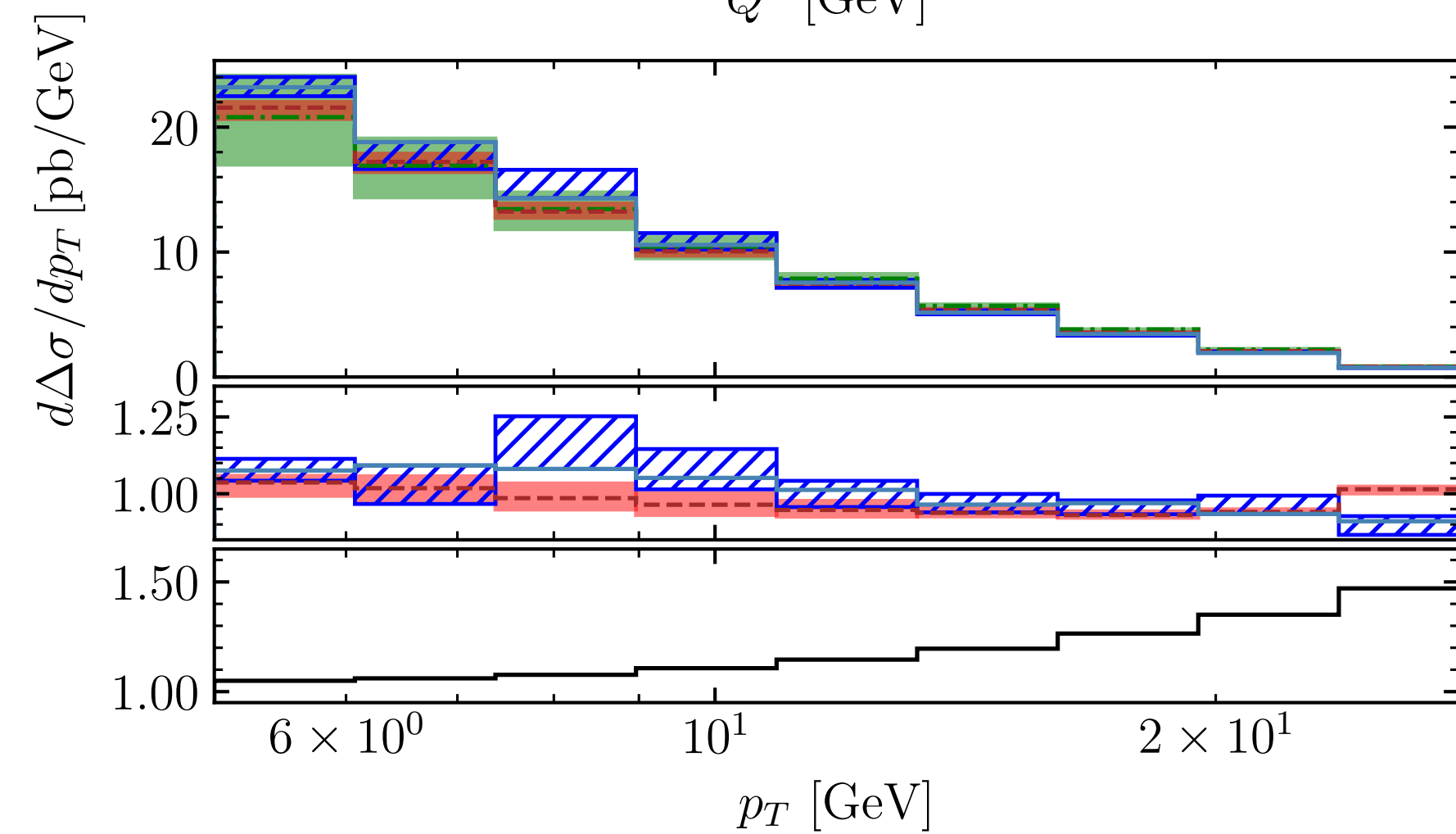
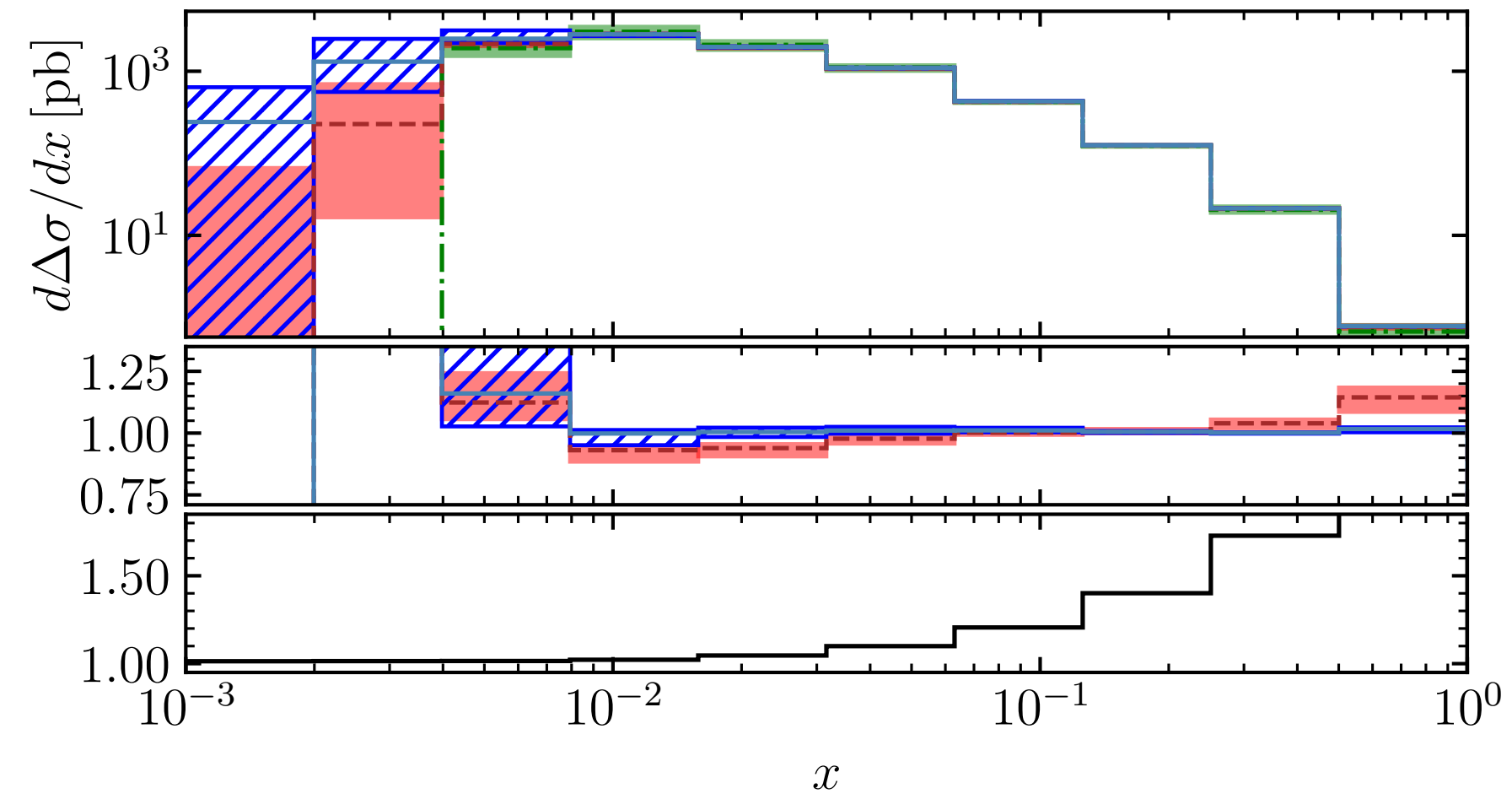
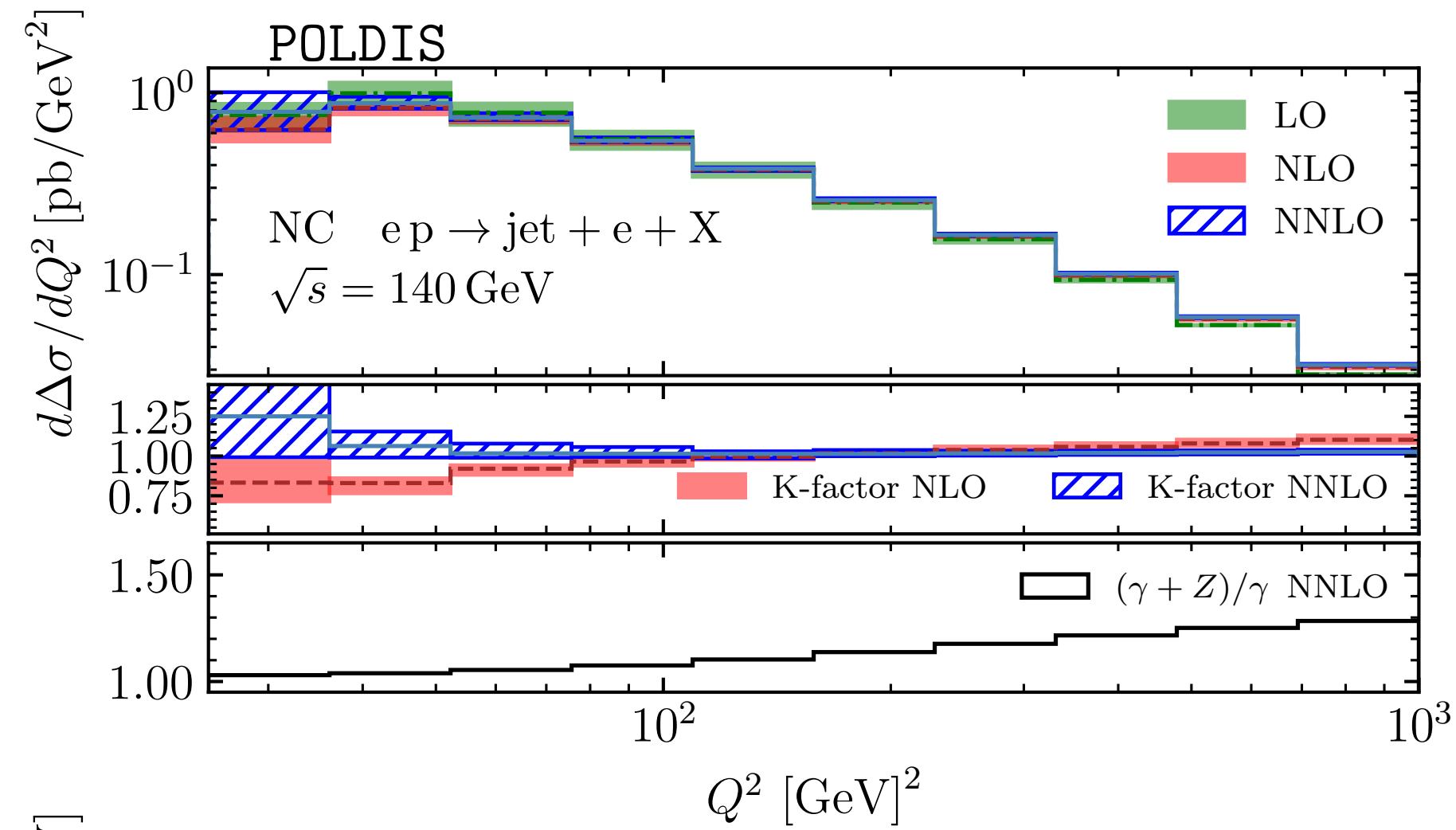


- Improved perturbative convergence
- Small contribution from Z-boson exchange

IB, de Florian, Pedron. Phys.Rev.D 107 (2023)

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Neutral Current DIS

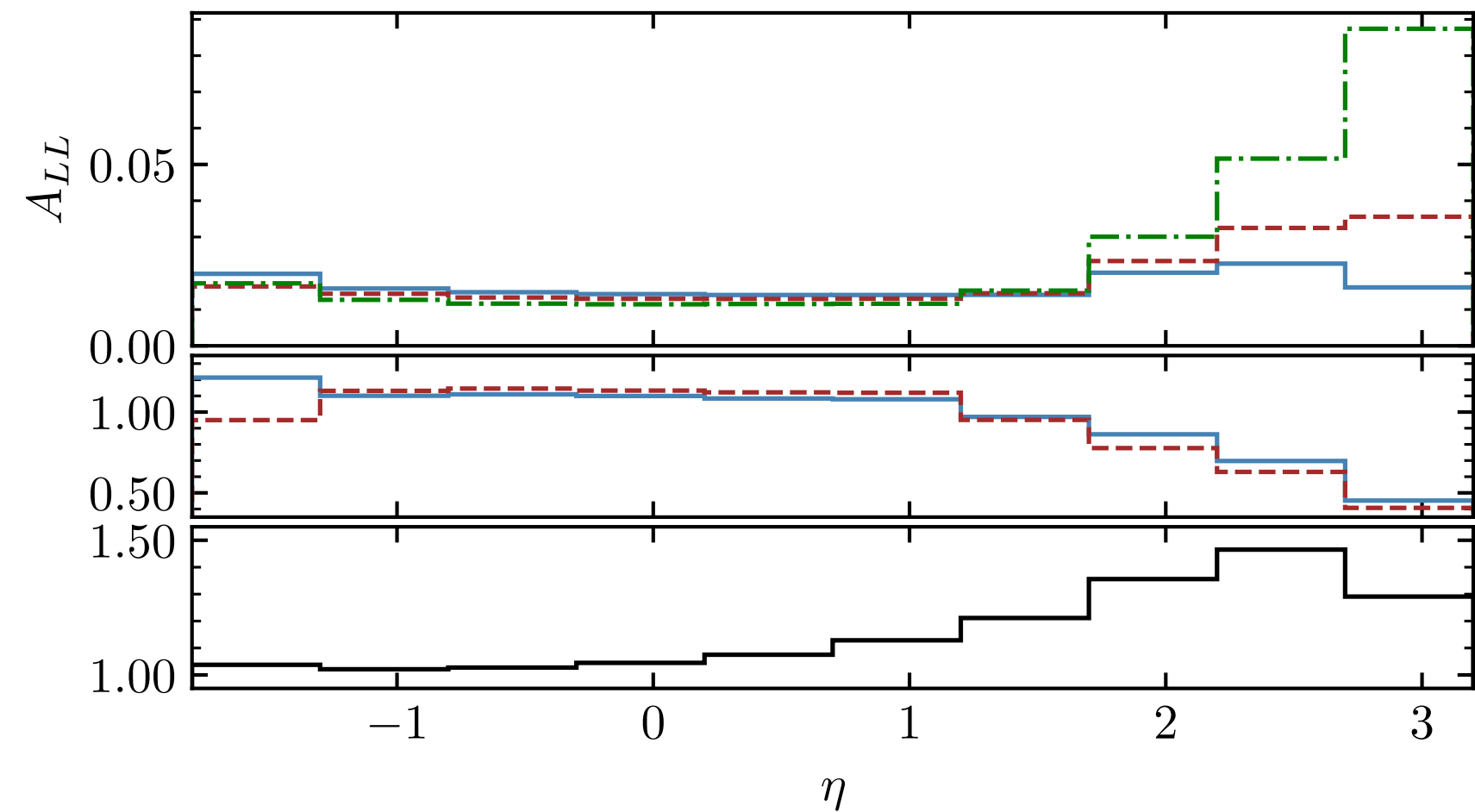
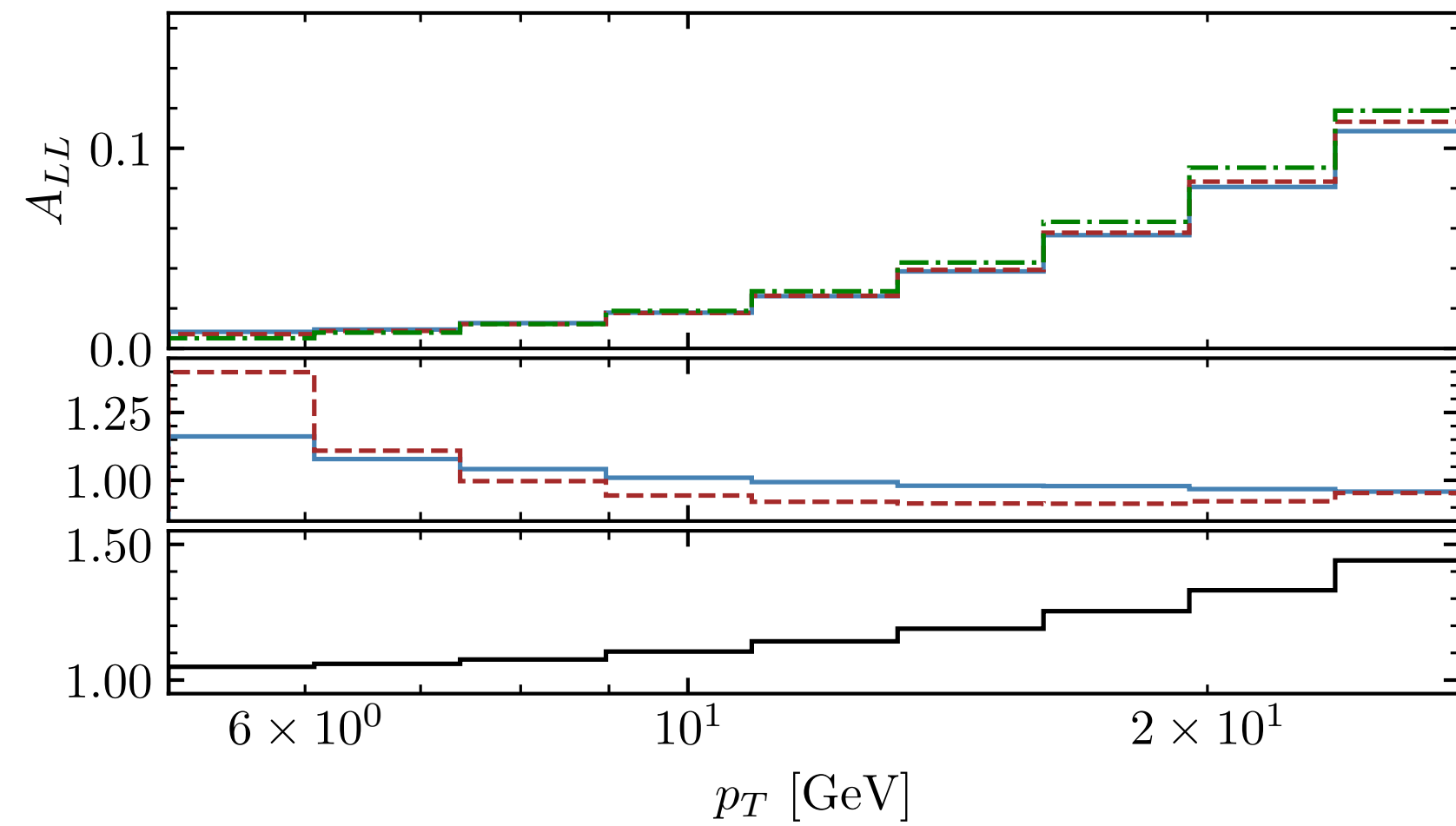
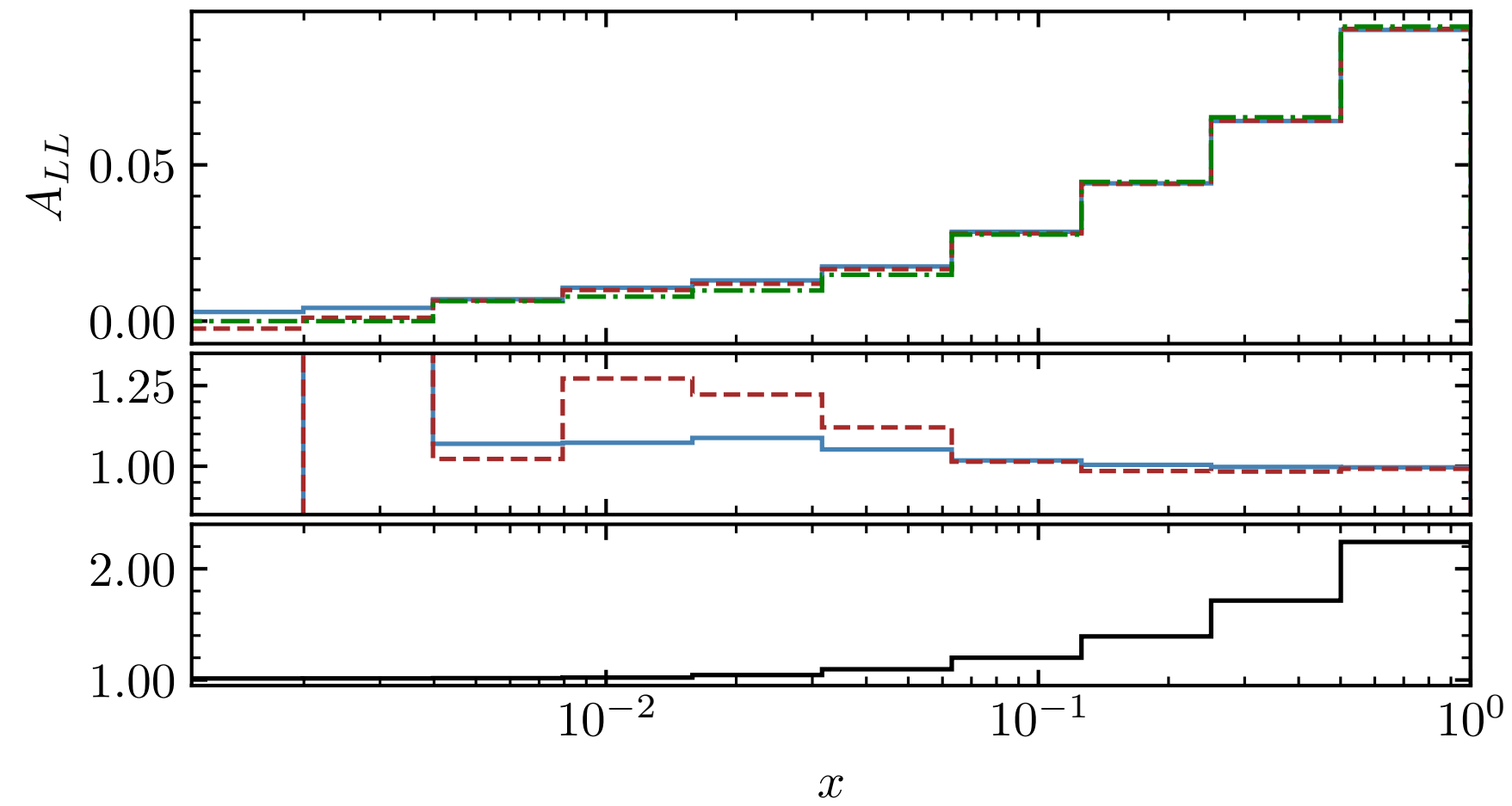
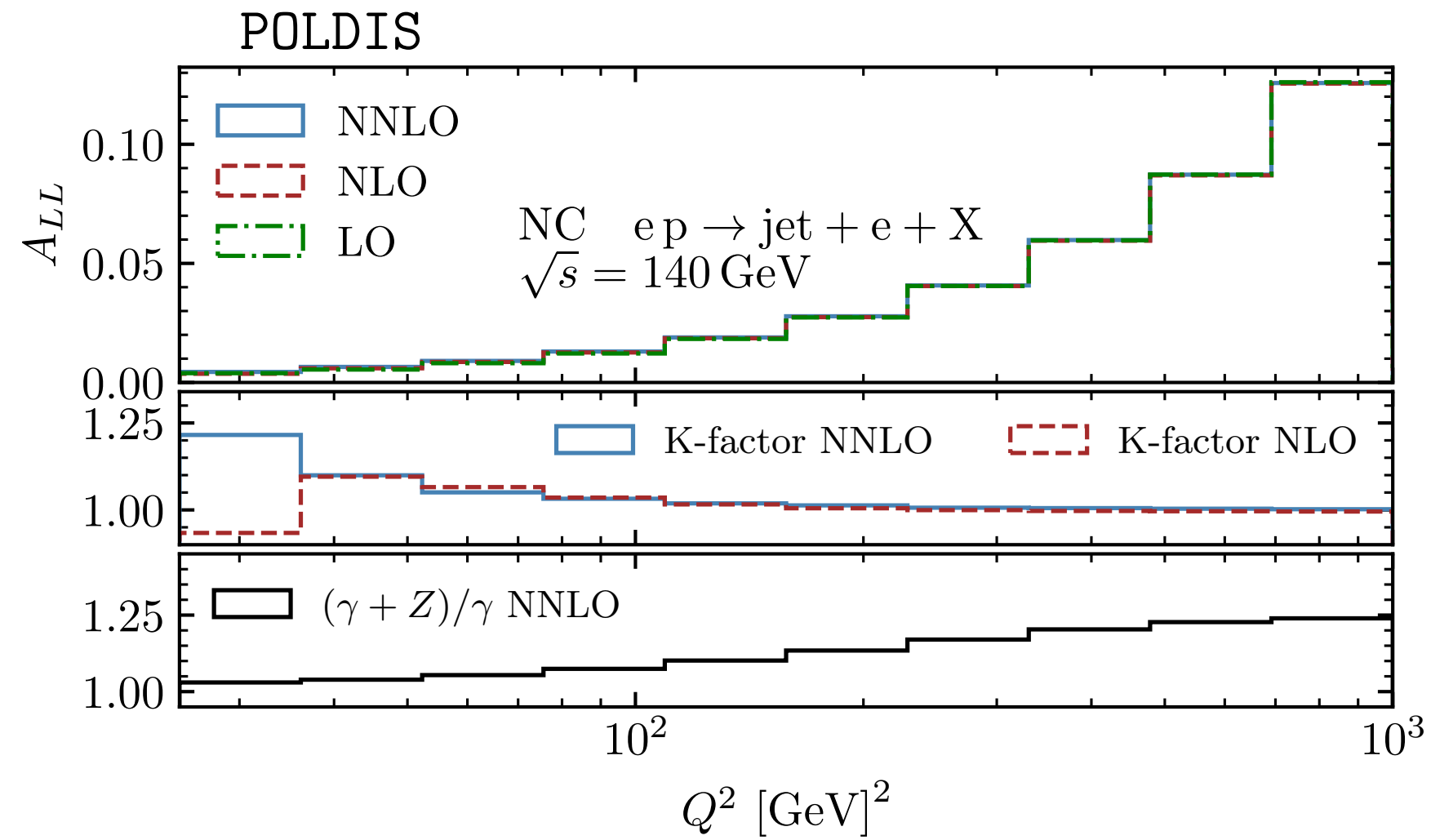


- Enhanced contributions at high Q^2 , x , and p_T compared to the unpolarized case (unsuppressed parity-violating terms for Z exchange)
- Enhancement is translated to increased spin asymmetries

IB, de Florian, Pedron. Phys.Rev.D 107 (2023)

PHENOMENOLOGY - NNLO SINGLE-JET PRODUCTION IN DIS

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IB, de Florian, Pedron. Phys.Rev.D 107 (2023)

HIGHER ORDER CORRECTIONS FOR JET OBSERVABLES IN POLARIZED DIS

Introduction

What can we learn from polarized processes?
Why do we need higher order corrections?

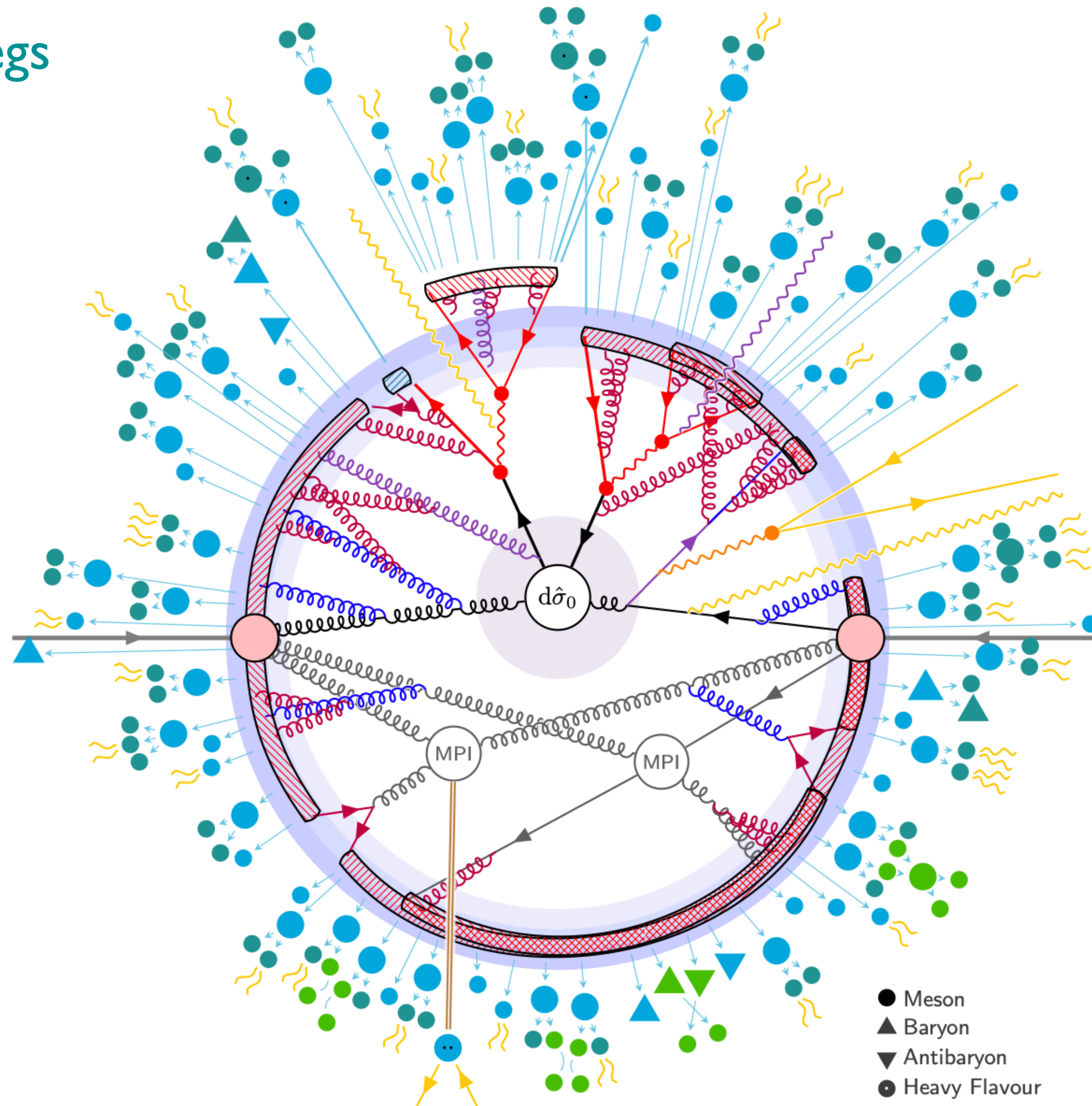
Fixed-order approach to jet observables

Jet production in polarized DIS at NNLO

Beyond fixed order calculations

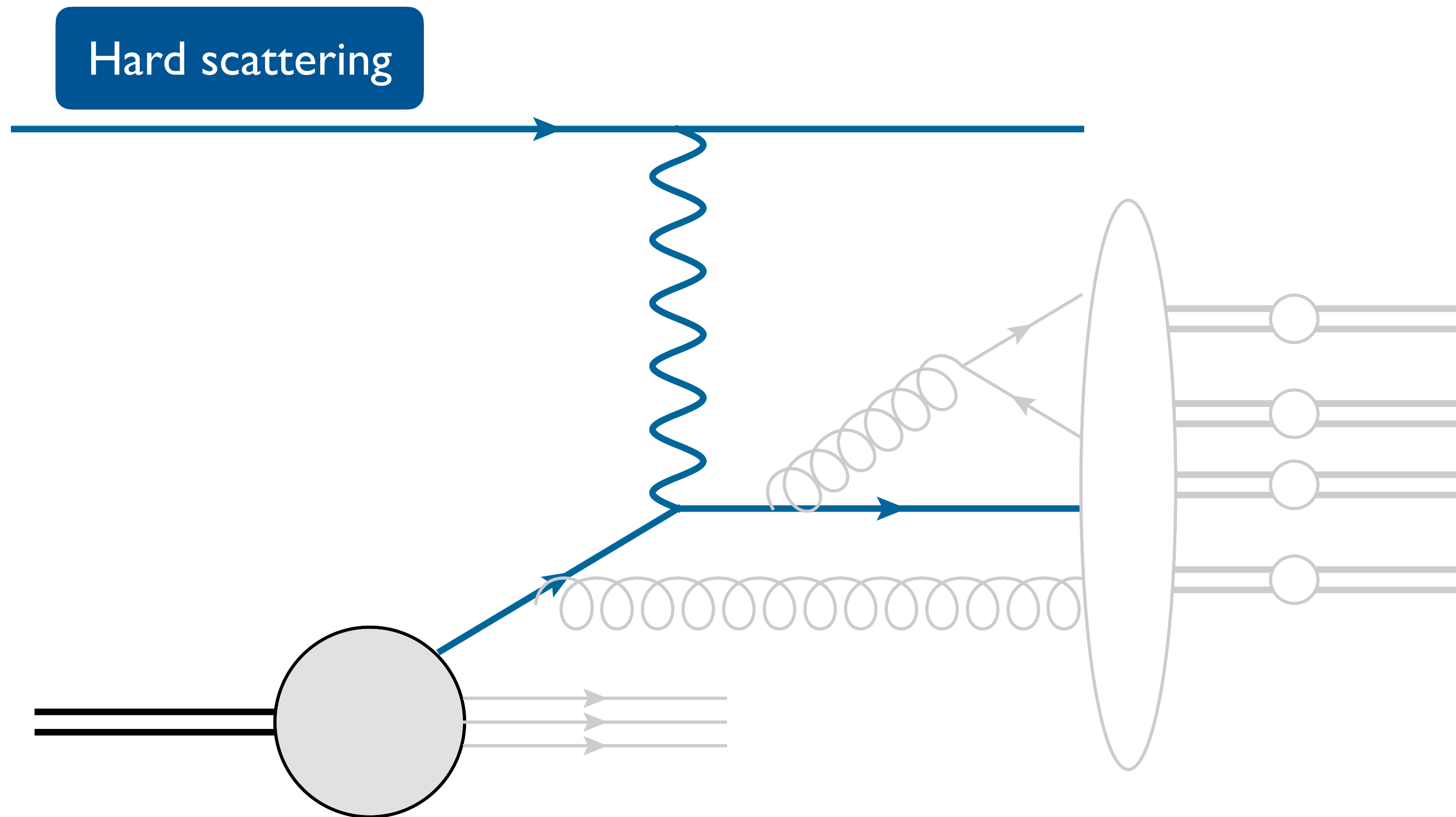
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(In collaboration with B. Jäger)

Loops \ll Legs

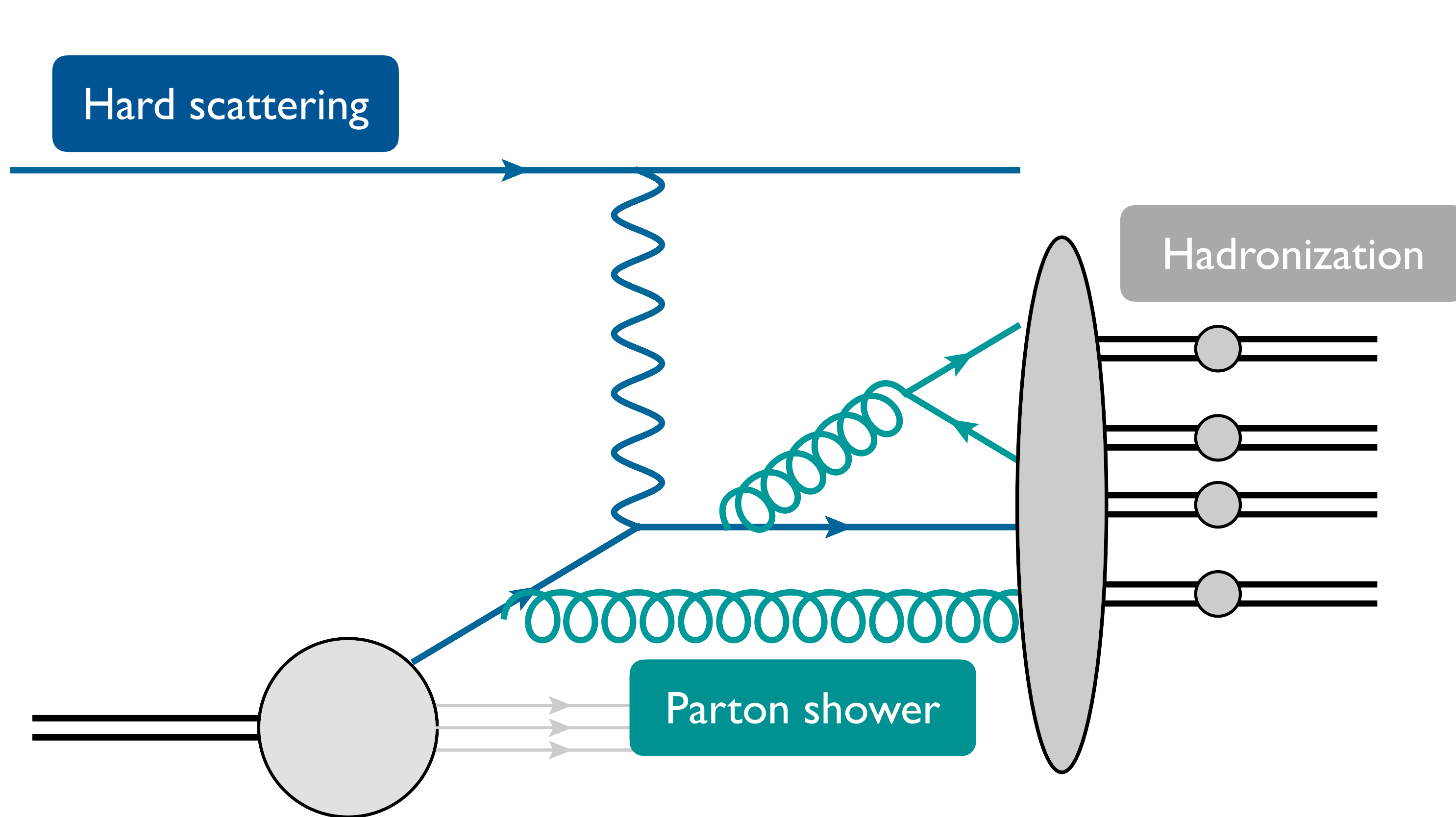


MONTE CARLO EVENT GENERATORS

[See Florian Herren's talk]



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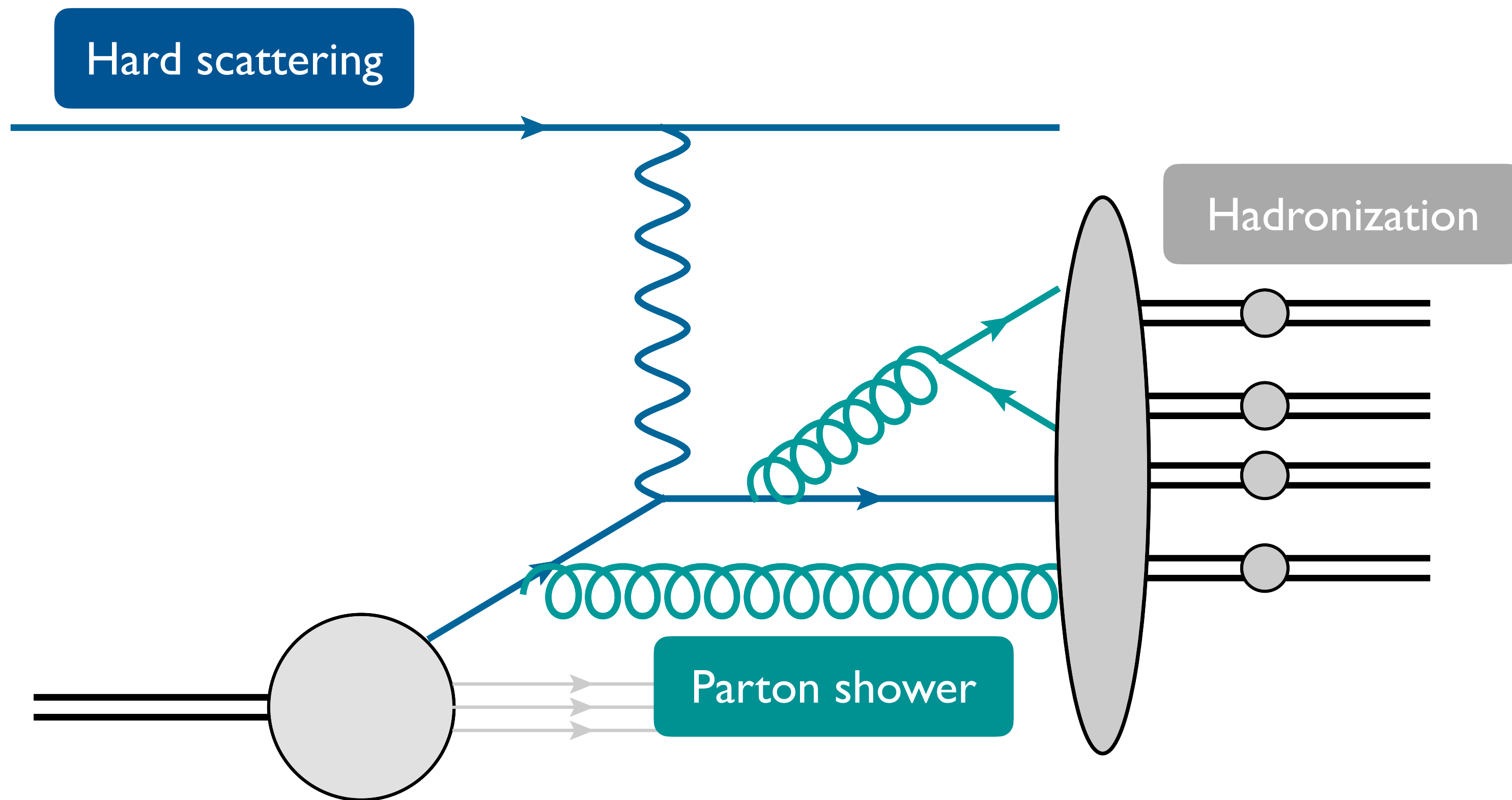
Leading Log accurate

- Hadronization
- Hadron decays
- EW radiation

...
PYTHIA - Bengtsson, Sjöstrand (1987)
HERWIG - Marchesini, Webber (1988)
ARIADNE - Lönnblad (1992)
SHERPA - Gleisberg, Höhe, Krauss, Schlicke,
Schumann, Winter (2004)
...

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[See Florian Herren's talk]



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Polarized event generators (unpolarized parton shower)

PEPSI Mankiewicz, Schafer, Veltri (1992)

DJANGO Charchula, Schuler, Spiesberger (1994)

FIXED-ORDER VS SMC

NLO

- Accurate distributions at high p_T
- Normalization accurate at NLO
- Wrong distributions at small p_T
- Description only at the parton level

SMC (LO + Parton Shower)

- Incorrect distributions at high p_T
- Normalization accurate at LO
- Correct Sudakov suppression at small p_T
- Possible to simulate events at the hadron level

Try to merge the two approaches, trying to keep the desirable features of both
potential problems with double counting of real emission!

MC@NLO - Frixione, Webber (2001)

POWHEG - Nason (2004)

MiNNLOps - Monni, Nason, Re, Wiesemann, Zanderighi (2020)

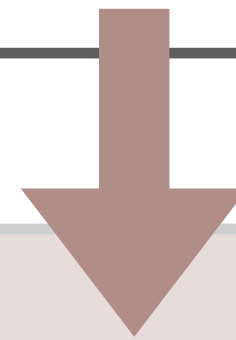
[See talk by Vasily Sotnikov]

NLO + PARTON SHOWER - POWHEG SCHEME

$$d\sigma_{\text{SMC}} = B(\Phi_n) d\Phi_n \left\{ \Delta_{t_0} + \frac{\alpha_S}{2\pi} P(z) \frac{1}{t} \Delta_t d\Phi_r \right\}$$
$$d\sigma_{\text{NLO}} = d\Phi_n \left\{ B(\Phi_n) + \left[V(\Phi_n) + \int d\Phi_r C(\Phi_n, \Phi_r) \right] + \left[R(\Phi_n, \Phi_r) - C(\Phi_n, \Phi_r) \right] d\Phi_r \right\}$$

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$$d\sigma_{\text{SMC}} = B(\Phi_n) d\Phi_n \left\{ \Delta_{t_0} + \frac{\alpha_S}{2\pi} P(z) \frac{1}{t} \Delta_t d\Phi_r \right\}$$
$$d\sigma_{\text{NLO}} = d\Phi_n \left\{ B(\Phi_n) + \left[V(\Phi_n) + \int d\Phi_r C(\Phi_n, \Phi_r) \right] + [R(\Phi_n, \Phi_r) - C(\Phi_n, \Phi_r)] d\Phi_r \right\}$$



$$d\sigma_{\text{POWHEG}} = \bar{B}(\Phi_n) d\Phi_n \left\{ \Delta(\Phi_n, p_{Tmin}) + \frac{R(\Phi_n, \Phi_r)}{B(\Phi_n)} \Delta(\Phi_n, p_T) d\Phi_r \right\}$$

NLO + PARTON SHOWER - POWHEG SCHEME

$$d\sigma_{\text{SMC}} = B(\Phi_n) d\Phi_n \left\{ \Delta_{t_0} + \frac{\alpha_S}{2\pi} P(z) \frac{1}{t} \Delta_t d\Phi_r \right\}$$

$$d\sigma_{\text{NLO}} = d\Phi_n \left\{ B(\Phi_n) + \left[V(\Phi_n) + \int d\Phi_r C(\Phi_n, \Phi_r) \right] + \left[R(\Phi_n, \Phi_r) - C(\Phi_n, \Phi_r) \right] d\Phi_r \right\}$$

$$d\sigma_{\text{POWHEG}} = \bar{B}(\Phi_n) d\Phi_n \left\{ \Delta(\Phi_n, p_{Tmin}) + \frac{R(\Phi_n, \Phi_r)}{B(\Phi_n)} \Delta(\Phi_n, p_T) d\Phi_r \right\}$$

POWHEG
SUDAKOV

$$\bar{B}(\Phi_n) = B(\Phi_n) + \left[V(\Phi_n) + \int d\Phi_r C(\Phi_n, \Phi_r) \right] + \int d\Phi_r \left[R(\Phi_n, \Phi_r) - C(\Phi_n, \Phi_r) \right] d\Phi_r$$

$$\Delta(\Phi_n, p_T) \sim \exp \left[- \int d\Phi'_r \frac{R(\Phi_n, \Phi'_r)}{B(\Phi_n)} \right]$$

NLO + PARTON SHOWER - POWHEG SCHEME

$$d\sigma_{\text{SMC}} = B(\Phi_n) d\Phi_n \left\{ \Delta_{t_0} + \frac{\alpha_s}{2\pi} P(z) \frac{1}{t} \Delta_t d\Phi_r \right\}$$

$$d\sigma_{\text{NLO}} = d\Phi_n \left\{ B(\Phi_n) + \left[V(\Phi_n) + \int d\Phi_r C(\Phi_n, \Phi_r) \right] + \left[R(\Phi_n, \Phi_r) - C(\Phi_n, \Phi_r) \right] d\Phi_r \right\}$$

$$d\sigma_{\text{POWHEG}} = \bar{B}(\Phi_n) d\Phi_n \left\{ \Delta(\Phi_n, p_{T\text{min}}) + \frac{R(\Phi_n, \Phi_r)}{B(\Phi_n)} \Delta(\Phi_n, p_T) d\Phi_r \right\}$$

POWHEG
SUDAKOV

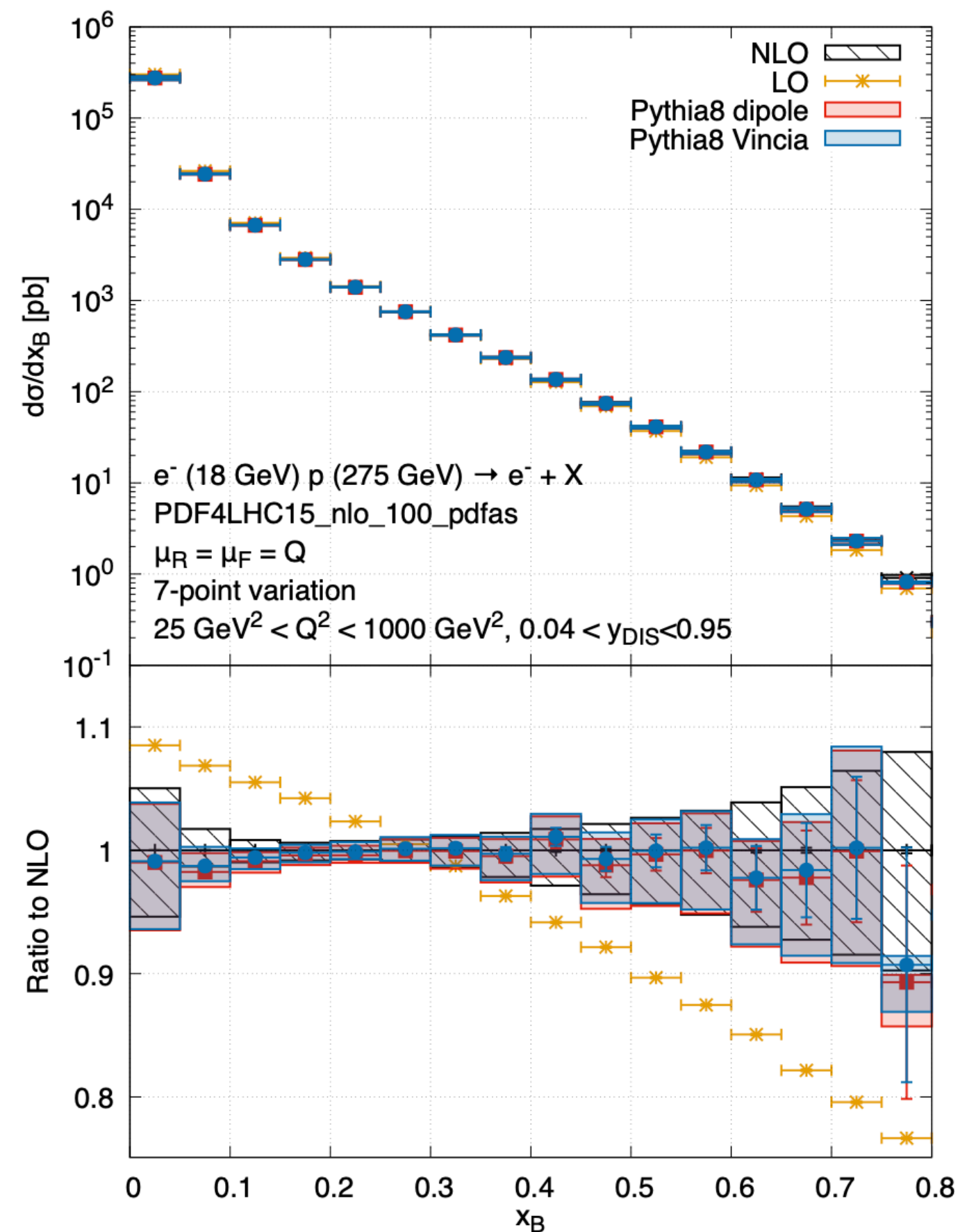
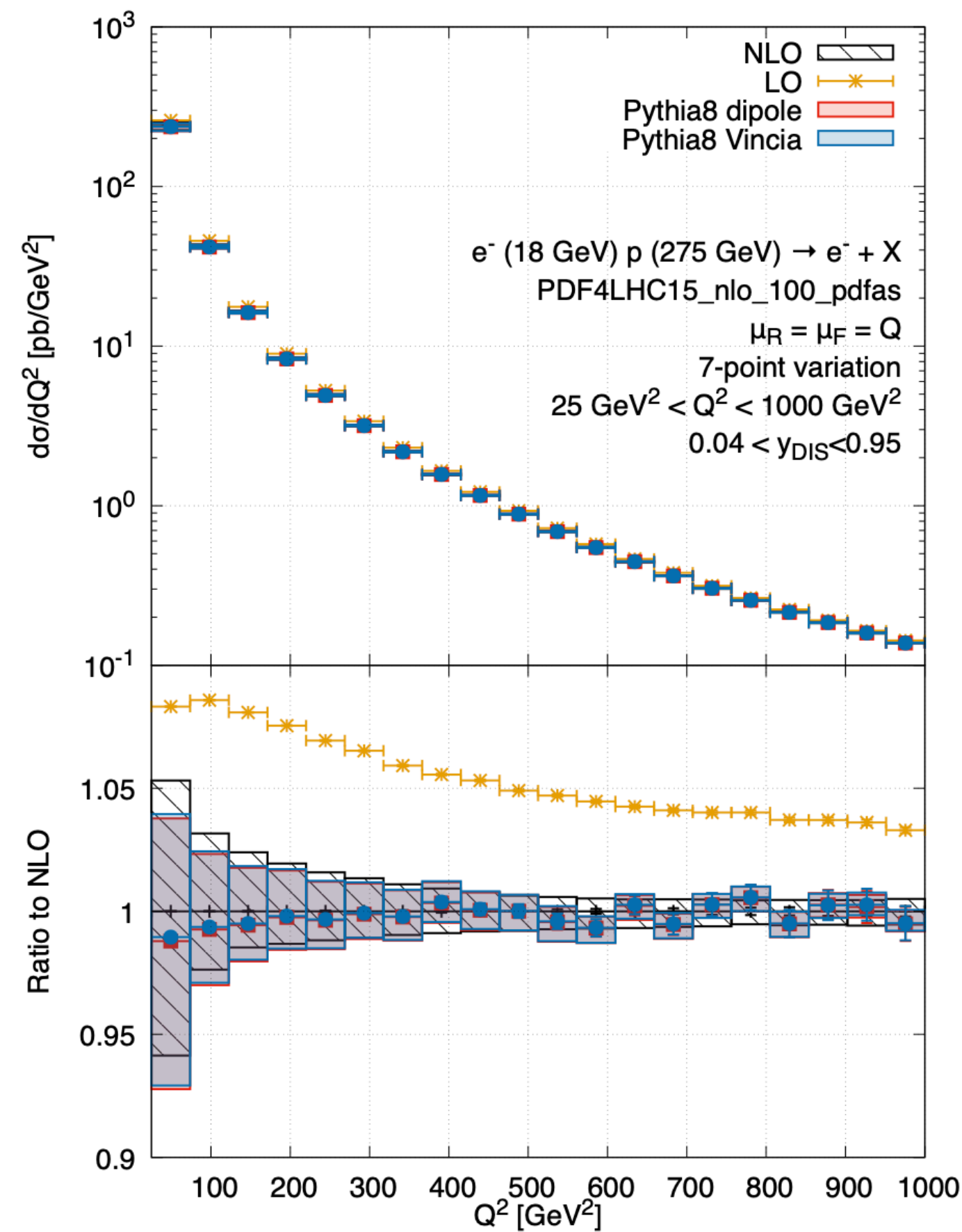
$$\bar{B}(\Phi_n) = B(\Phi_n) + \left[V(\Phi_n) + \int d\Phi_r C(\Phi_n, \Phi_r) \right] + \int d\Phi_r \left[R(\Phi_n, \Phi_r) - C(\Phi_n, \Phi_r) \right] d\Phi_r$$

$$\Delta(\Phi_n, p_T) \sim \exp \left[- \int d\Phi'_r \frac{R(\Phi_n, \Phi'_r)}{B(\Phi_n)} \right]$$

- Hardest emission generated according to the POWHEG Sudakov and $\bar{B}(\Phi_n)$
 - Subsequent (less hard) radiation generated using parton-shower programs (p_T veto)
- NLO accuracy on integrated quantities
 Leading log accurate

NLO + PARTON SHOWER - POWHEG SCHEME

$$d\sigma_{\text{POWHEG}} = \bar{B}(\Phi_n) d\Phi_n \left\{ \Delta(\Phi_n, p_{Tmin}) + \frac{R(\Phi_n, \Phi_r)}{B(\Phi_n)} \Delta(\Phi_n, p_T) d\Phi_r \right\}$$



Unpolarized DIS (NC and CC) in POWHEG
 Banfi, Ferraro Ravasio, Jäger, Karlberg, Reichenbach, Zanderighi (2023)

NLO + PARTON SHOWER - **POLARIZED** POWHEG SCHEME

$$d\sigma_{\text{POWHEG}} = \Delta\bar{B}(\Phi_n) d\Phi_n \left\{ \Delta^{\text{pol}}(\Phi_n, p_{Tmin}) + \frac{\Delta R(\Phi_n, \Phi_r)}{\Delta B(\Phi_n)} \Delta^{\text{pol}}(\Phi_n, p_T) d\Phi_r \right\}$$

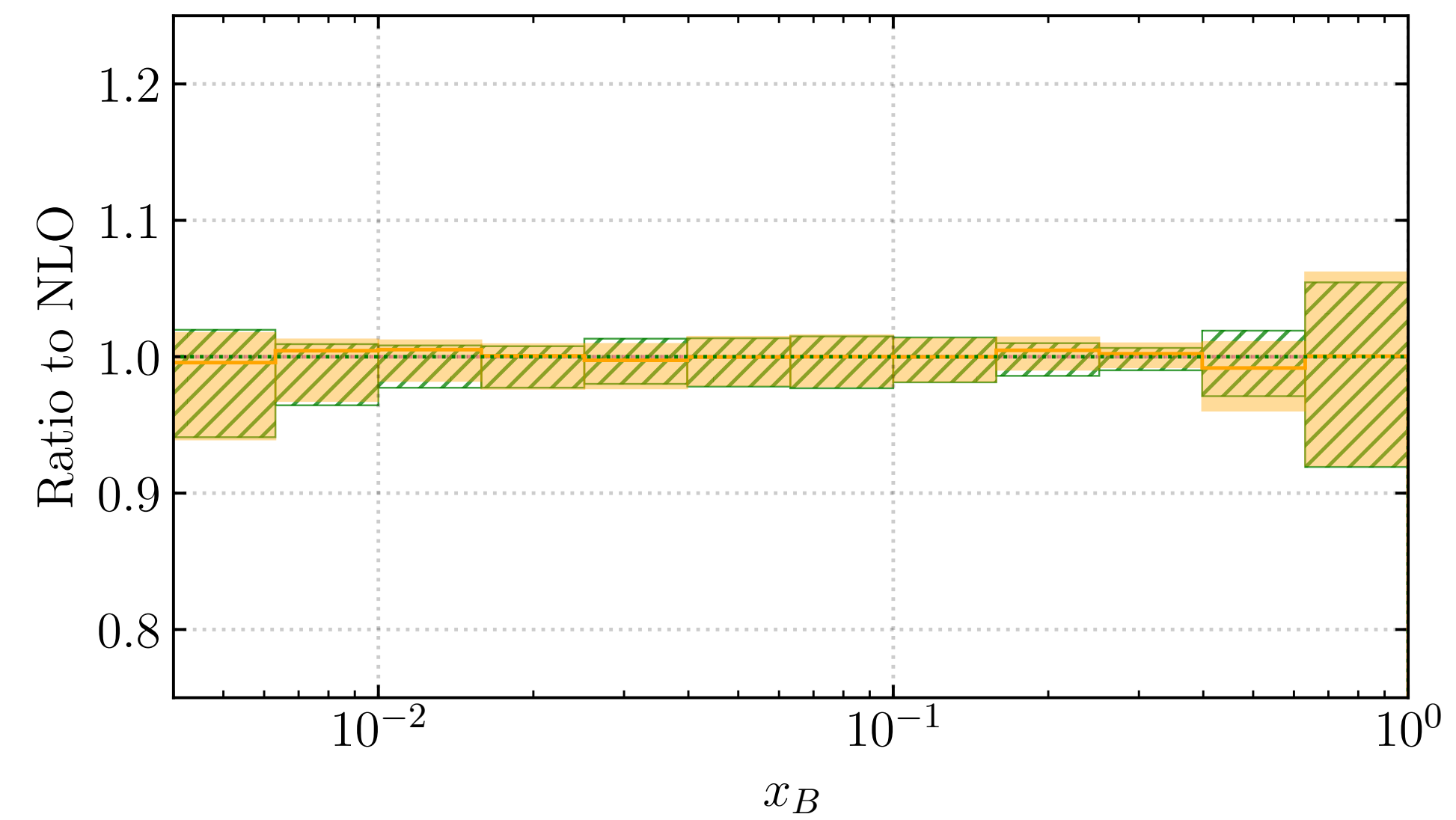
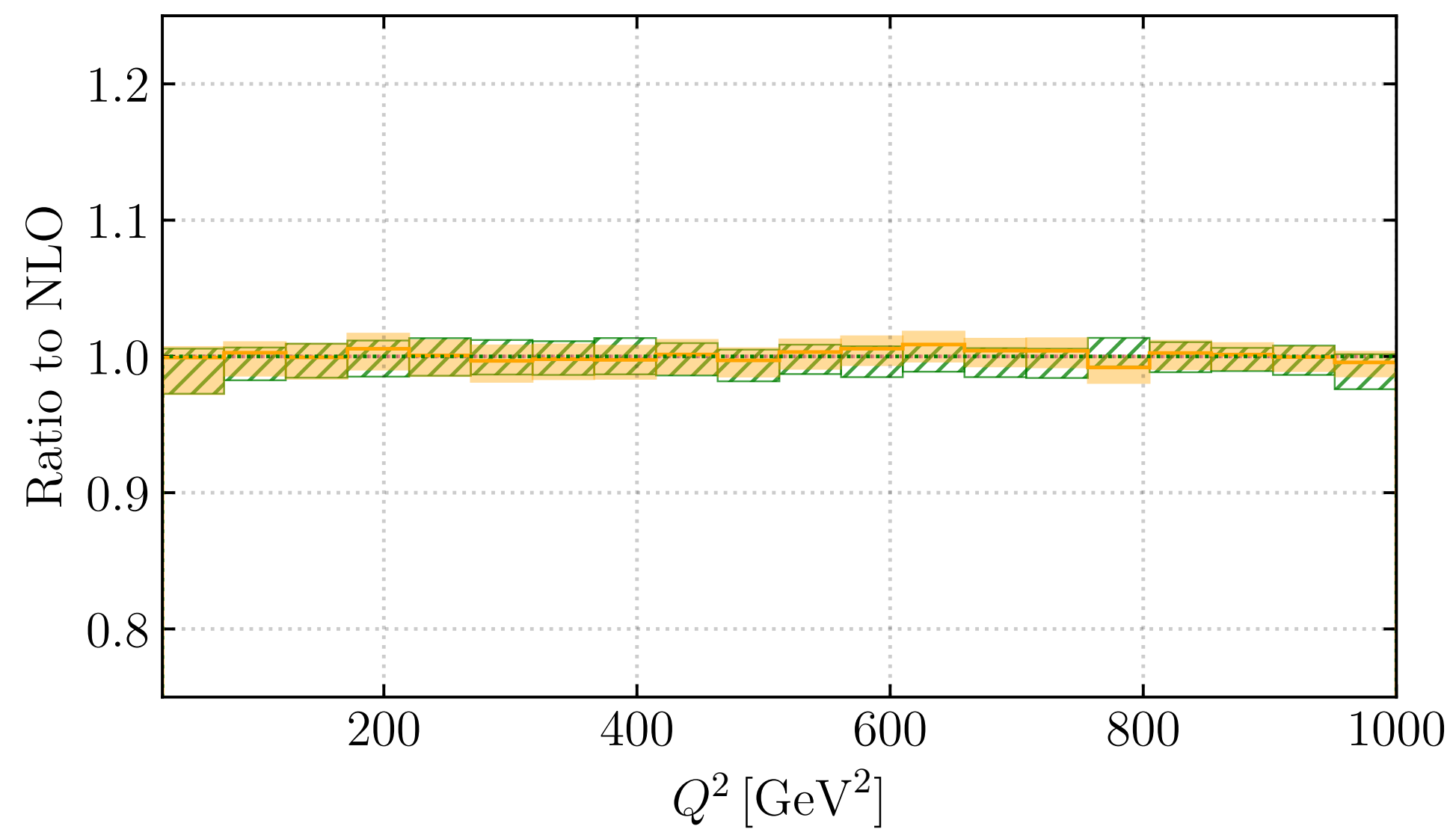
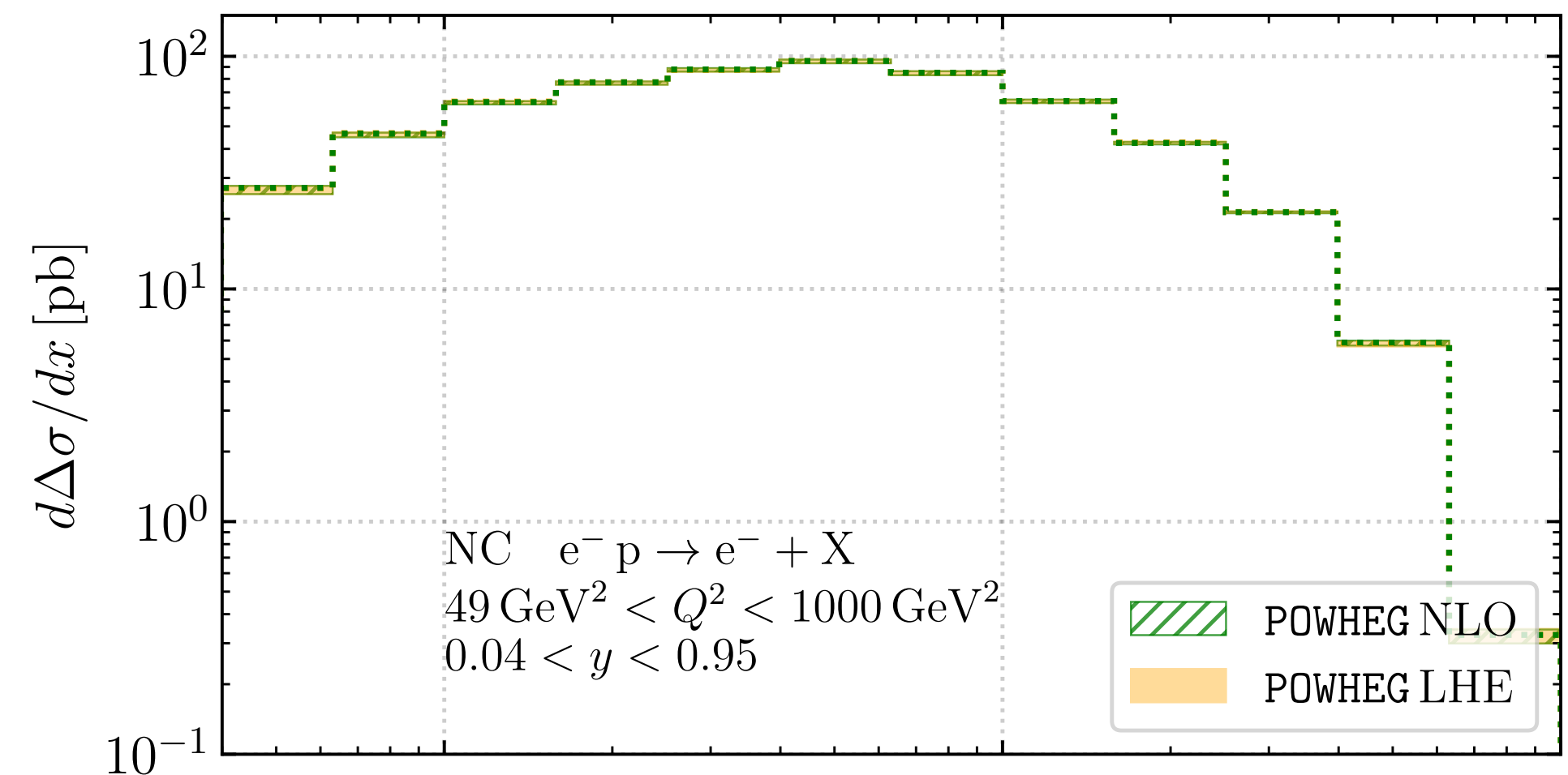
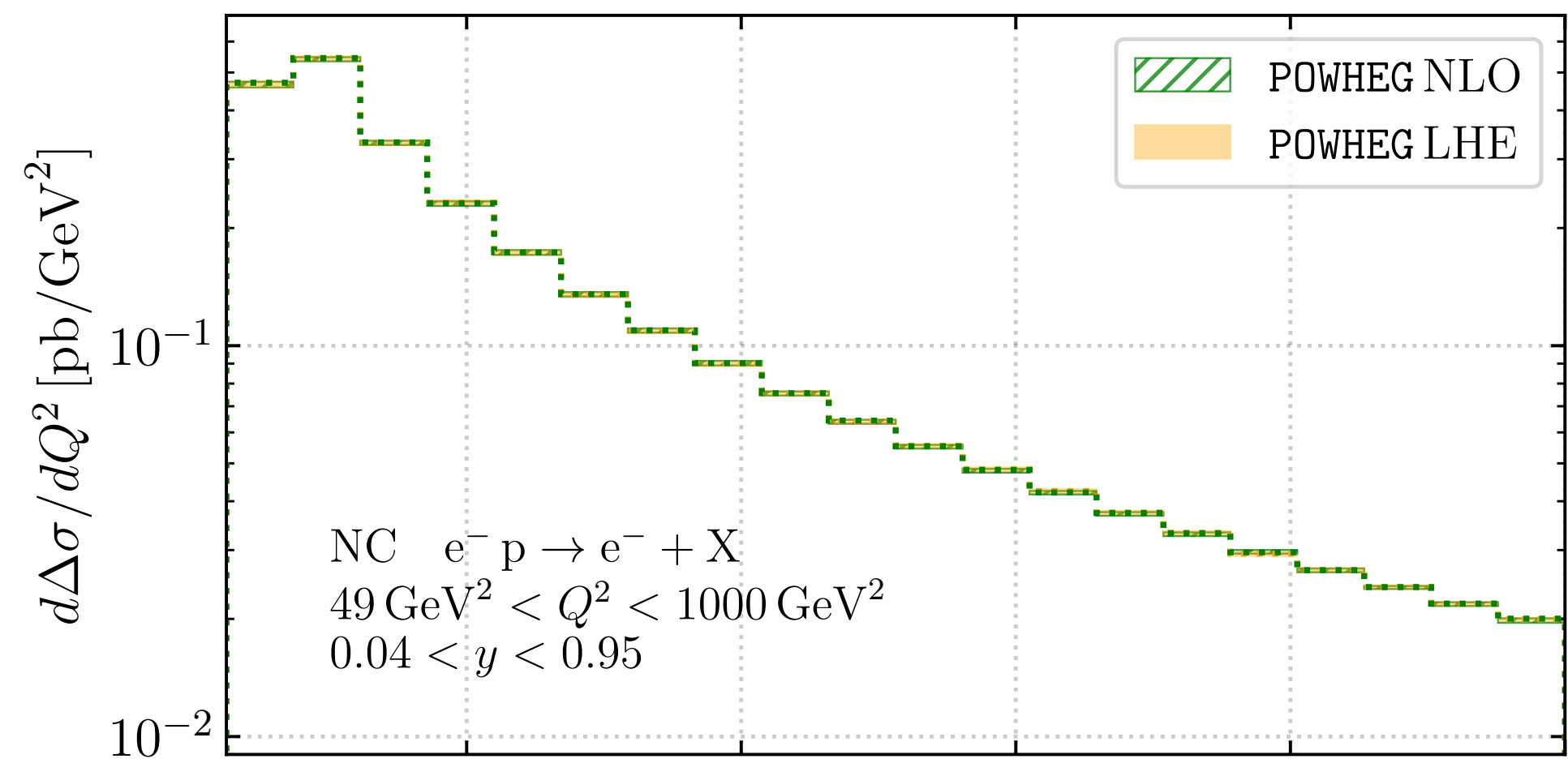
$$\Delta\bar{B}(\Phi_n) = \Delta B(\Phi_n) + \left[\Delta V(\Phi_n) + \int d\Phi_r \Delta C(\Phi_n, \Phi_r) \right] + \int d\Phi_r [\Delta R(\Phi_n, \Phi_r) - \Delta C(\Phi_n, \Phi_r)] d\Phi_r$$

$$\Delta(\Phi_n, p_T) \sim \exp \left[- \int d\Phi_r' \frac{\Delta R(\Phi_n, \Phi_r')}{\Delta B(\Phi_n)} \right]$$

Modifications to handle polarized DIS

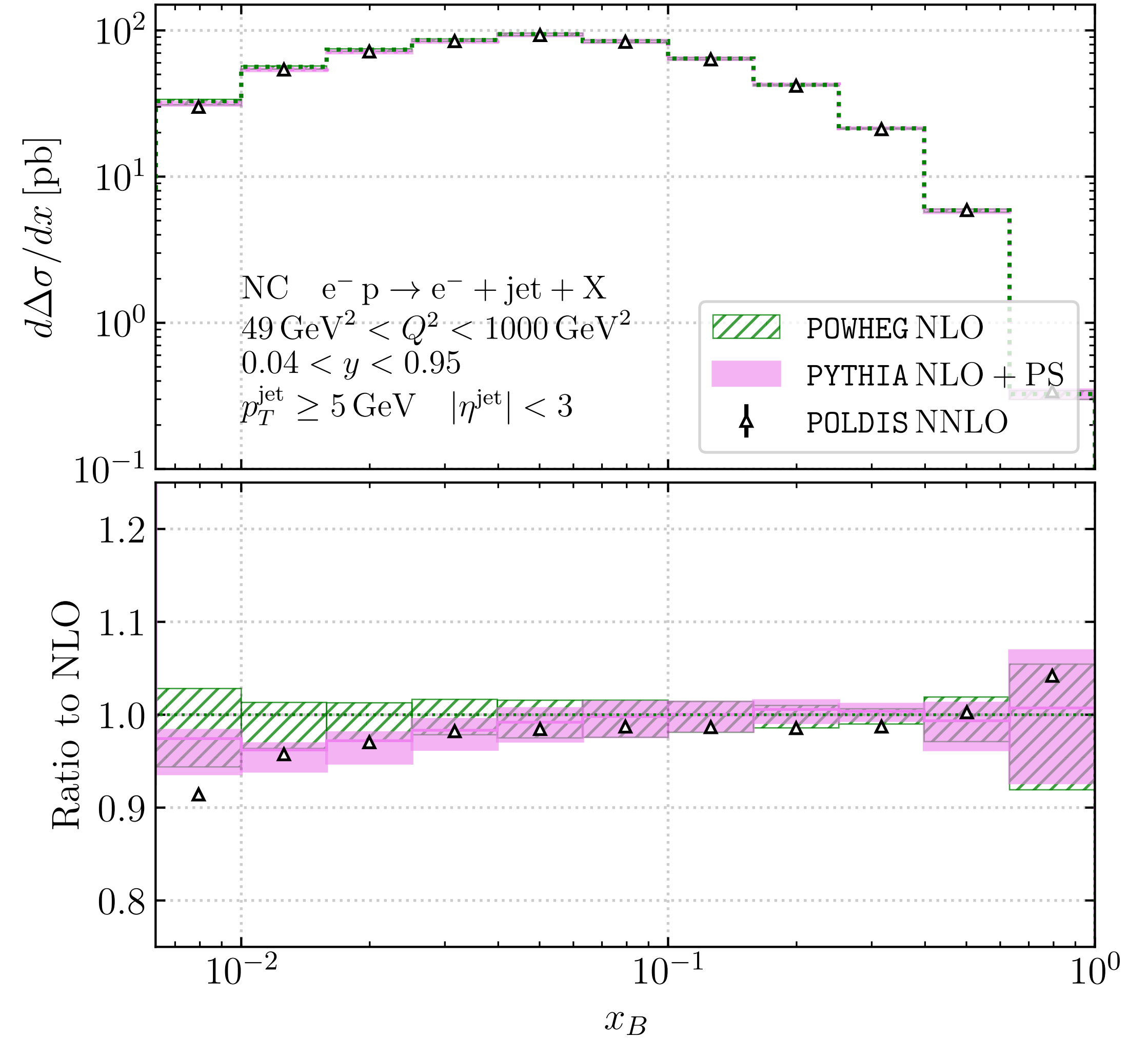
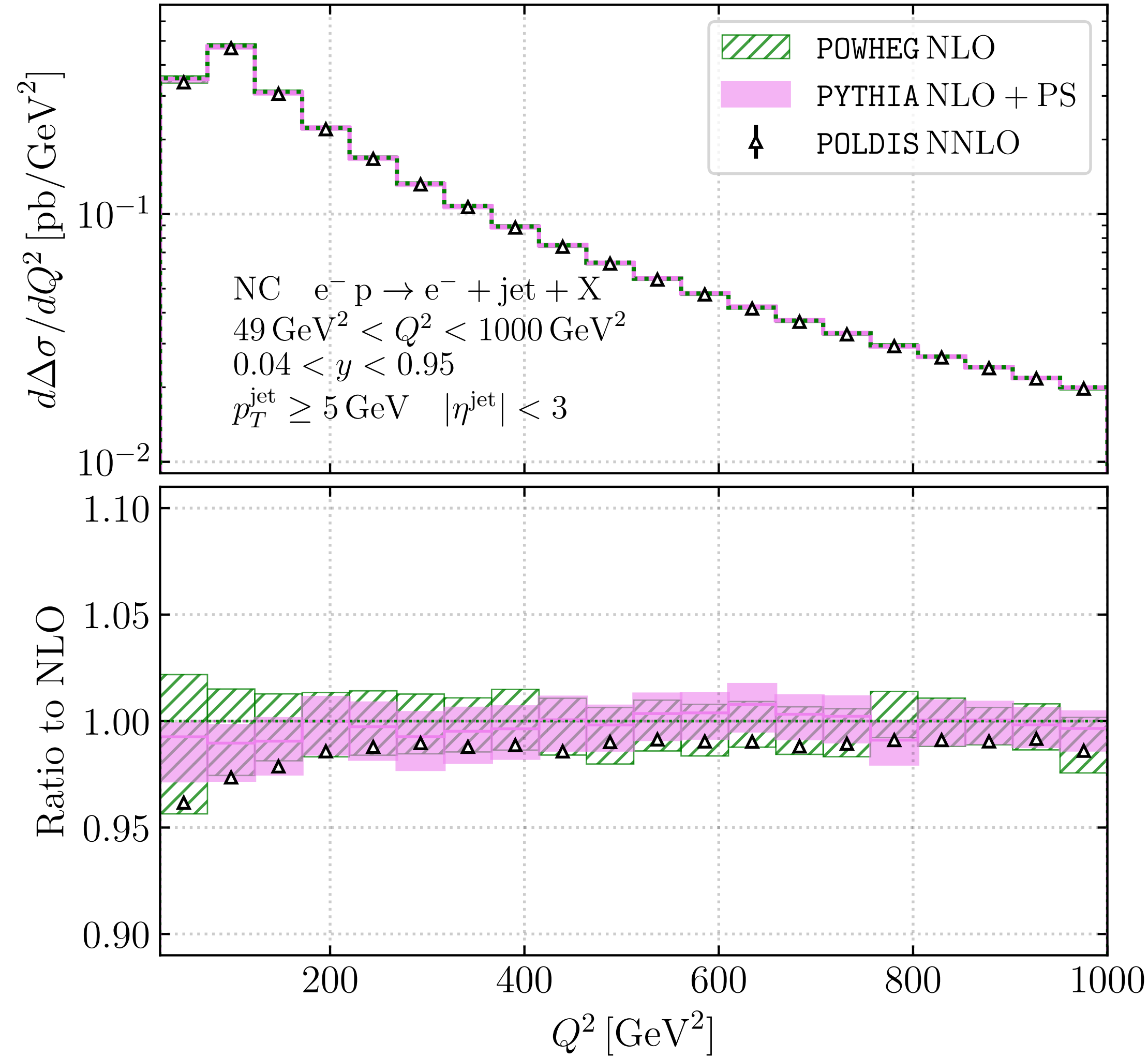
- Polarized Matrix elements & PDFs
- NLO Subtraction scheme → implementation of polarized FKS subtraction
Frixione, Kunszt, Signer(1996)
de Florian, Frixione, Signer, Vogelsang (1999)
- Modifications to handle negative weights
- **Polarized shower**

PHENOMENOLOGY



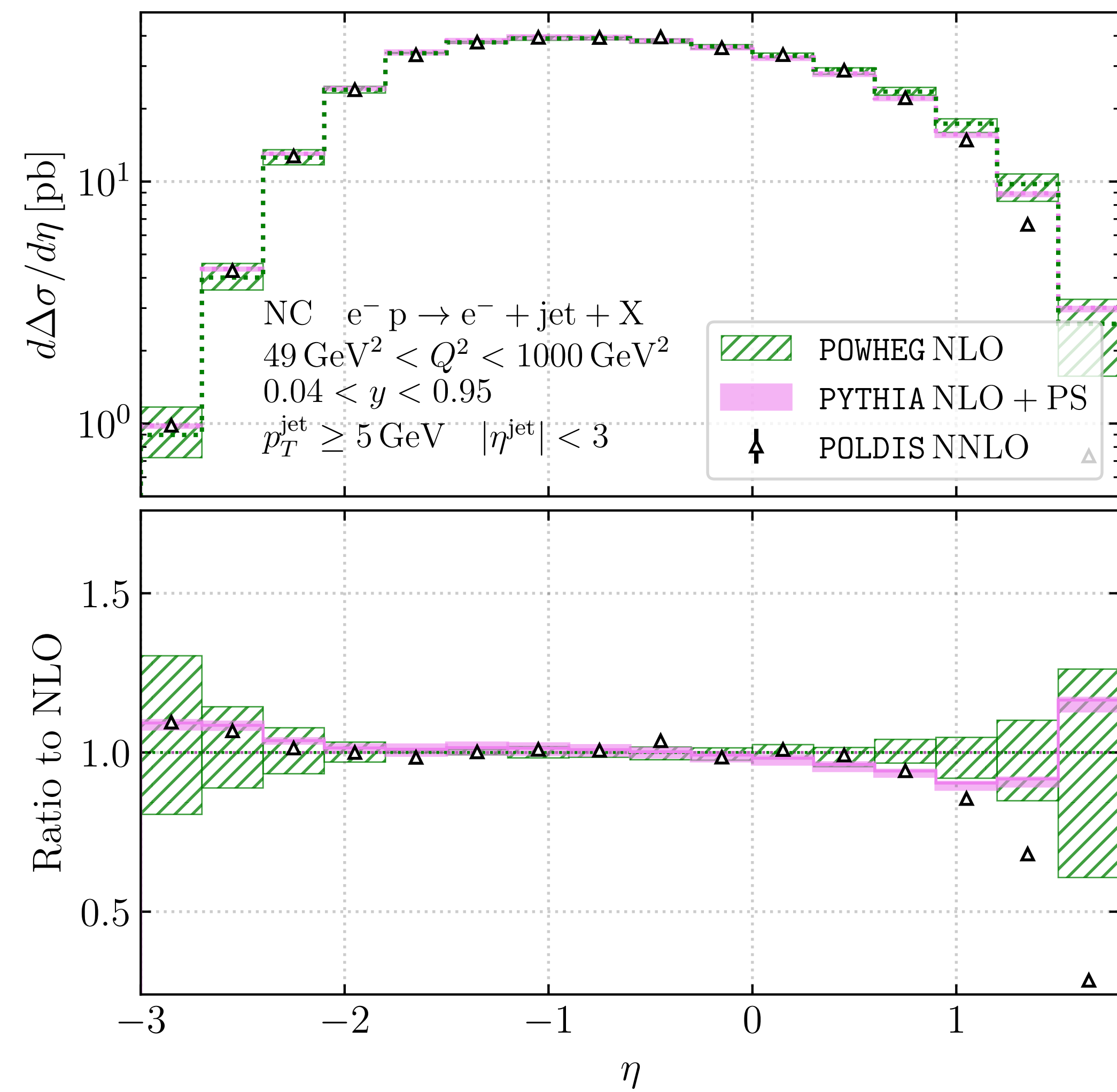
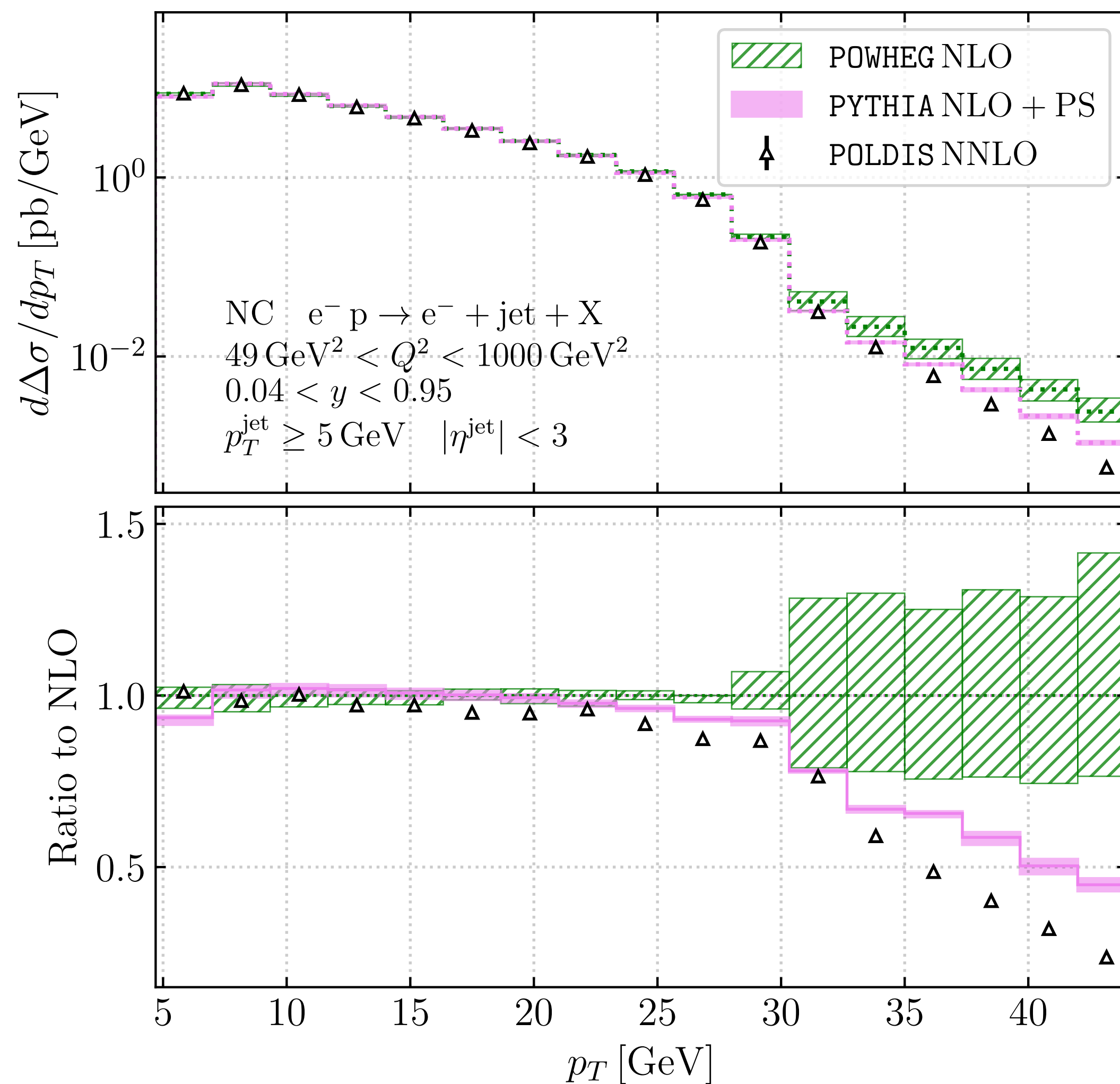
- Small corrections from the PS to more inclusive observables
- Slight improvement in the agreement with NNLO results

PHENOMENOLOGY



- Small corrections from the PS to more inclusive observables
- Slight improvement in the agreement with NNLO results

PHENOMENOLOGY



- Sizable corrections in the kinematical suppressed regions
- Slight improvement in the agreement with NNLO results

SUMMARY

- Higher order corrections instrumental for a precise description of polarized jet observables, and will play a central role in the improvement of our picture of the proton's spin.

NNLO calculation for jet production in Polarized DIS - full NC & CC

- Numerical implementation in code POLDIS
(also allows NLO calculation for di-jet production in pol DIS).
- For EIC kinematics, Increased perturbative stability is observed, but corrections are still sizable.
- Significant corrections in double spin asymmetries.

NLO+PS implementation of polarized DIS in POWHEG

- Extended POWHEG scheme to account for the helicities of the initial-state particles.
- Sizable PS effects in selected regions of phase space for EIC kinematics.
- Important step towards the development of polarized parton-shower generators for EIC, with polarization included in all stages of simulation.

THANK YOU

BACK-UP SLIDES

NNLO SINGLE-JET PRODUCTION IN DIS

EIC kinematics: $E_p = 275 \text{ GeV}$ $E_e = 18 \text{ GeV}$

Laboratory Frame (Needed for the P2B mapping)

Kinematical cuts: $0.04 < y < 0.95,$ $5 \text{ GeV} < p_T^L < 36 \text{ GeV},$
 $25 \text{ GeV}^2 < Q^2 < 1000 \text{ GeV}^2$ $|\eta^L| < 3,$

NLO PDFS:

DSSV14 MC (polarized case)/PDF4LHC15 (unpolarized case)

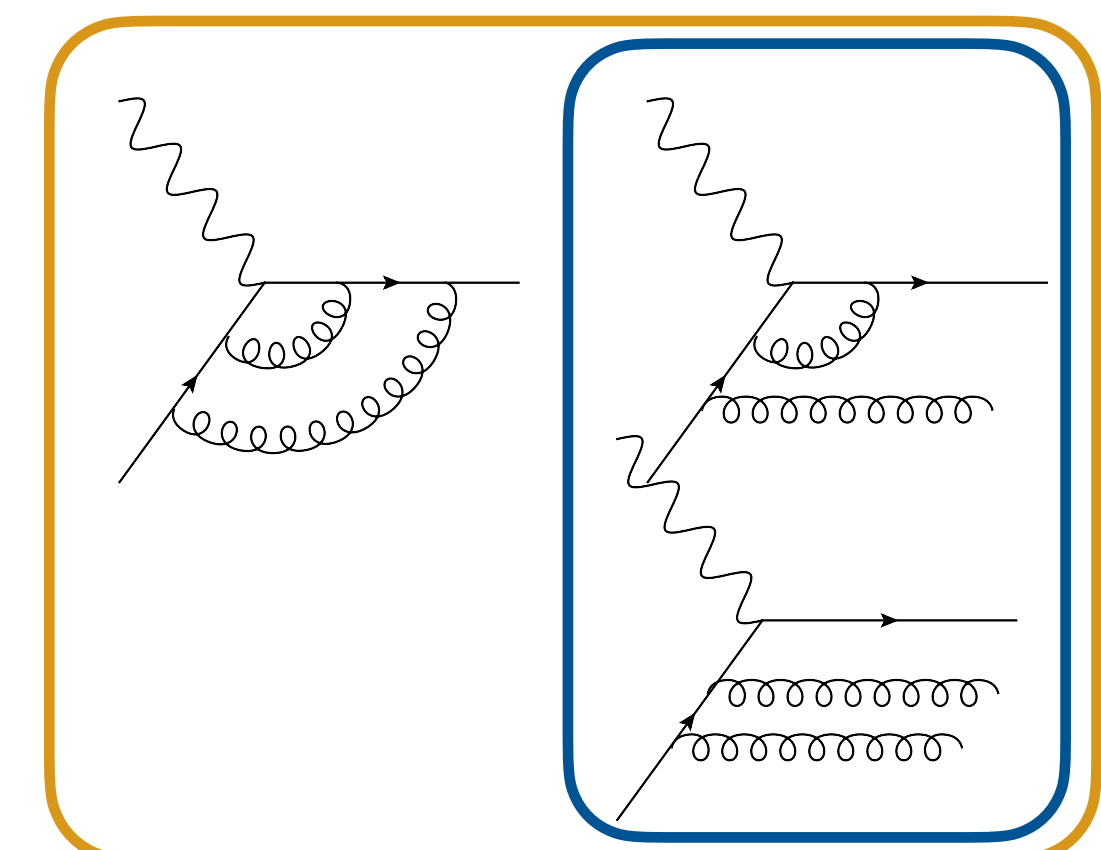
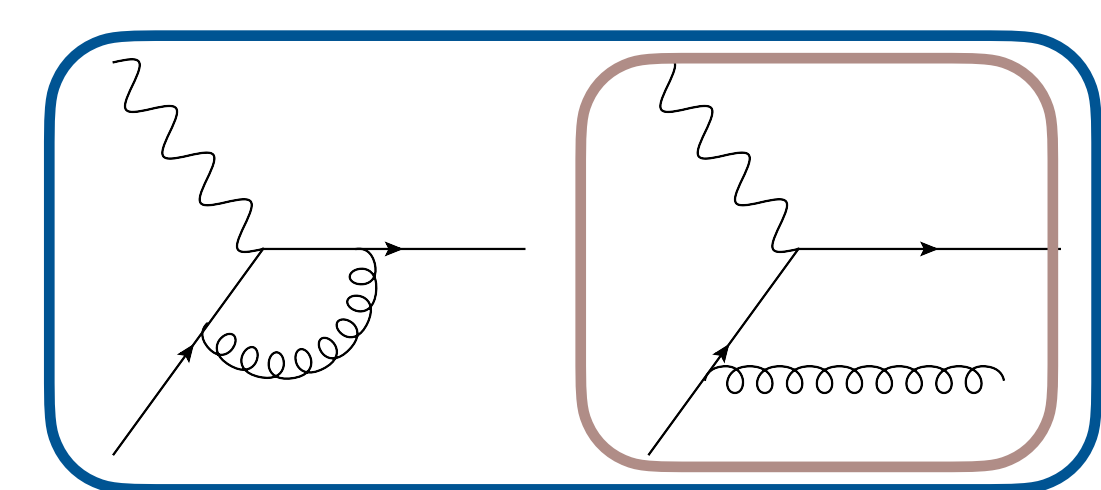
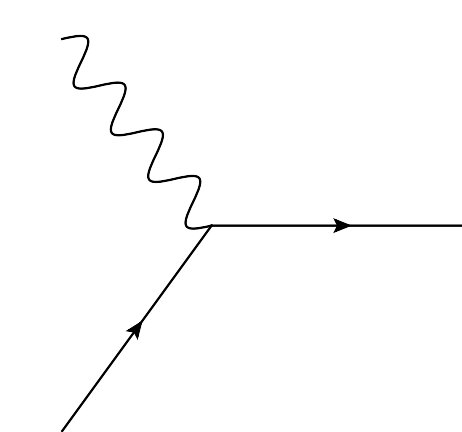
Central Scale: $\mu_F^2 = \mu_R^2 = Q^2 \equiv \mu_0$

k_T jet algorithm (R=0.8)

Seven-Point independent variation to estimate theoretical uncertainty

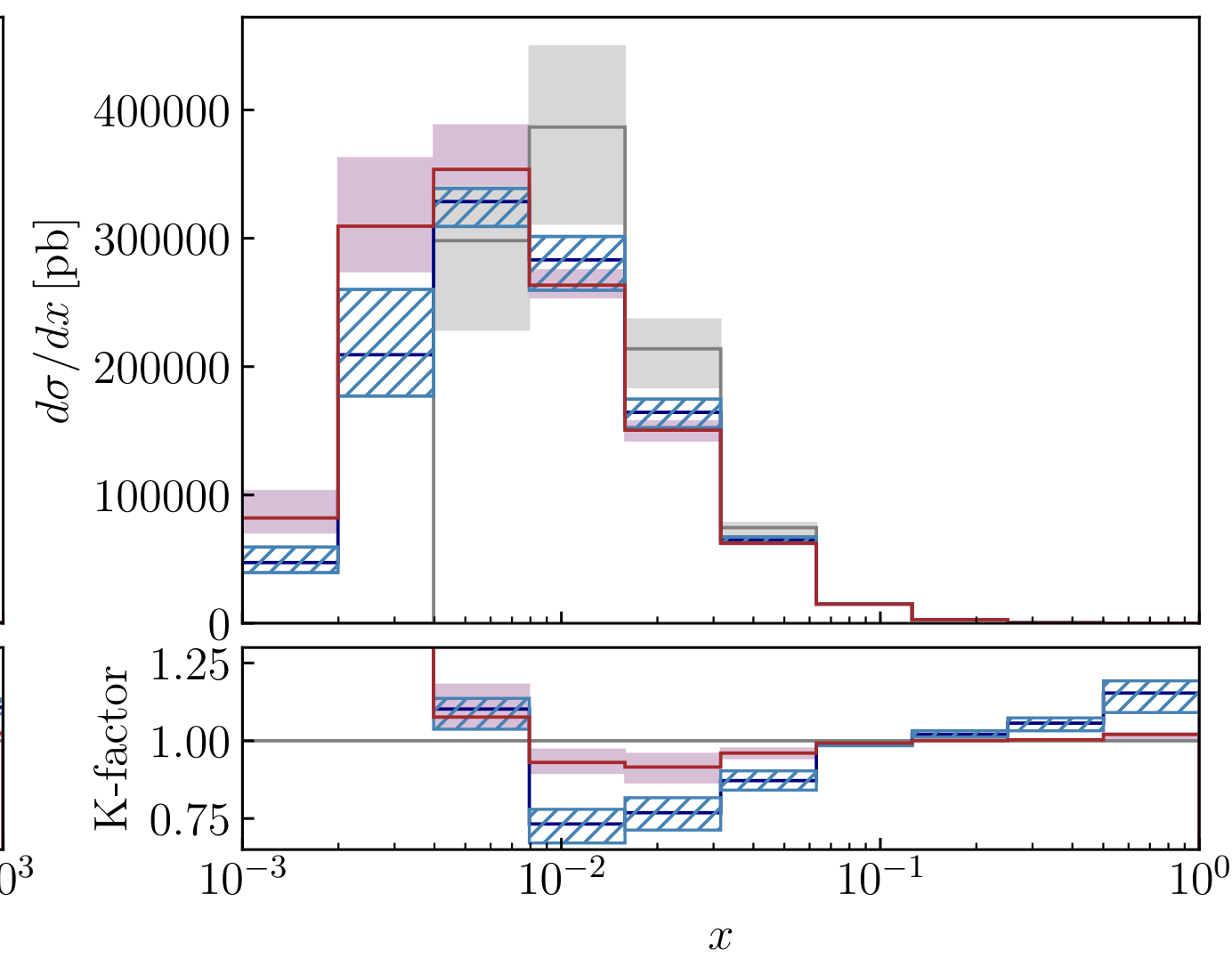
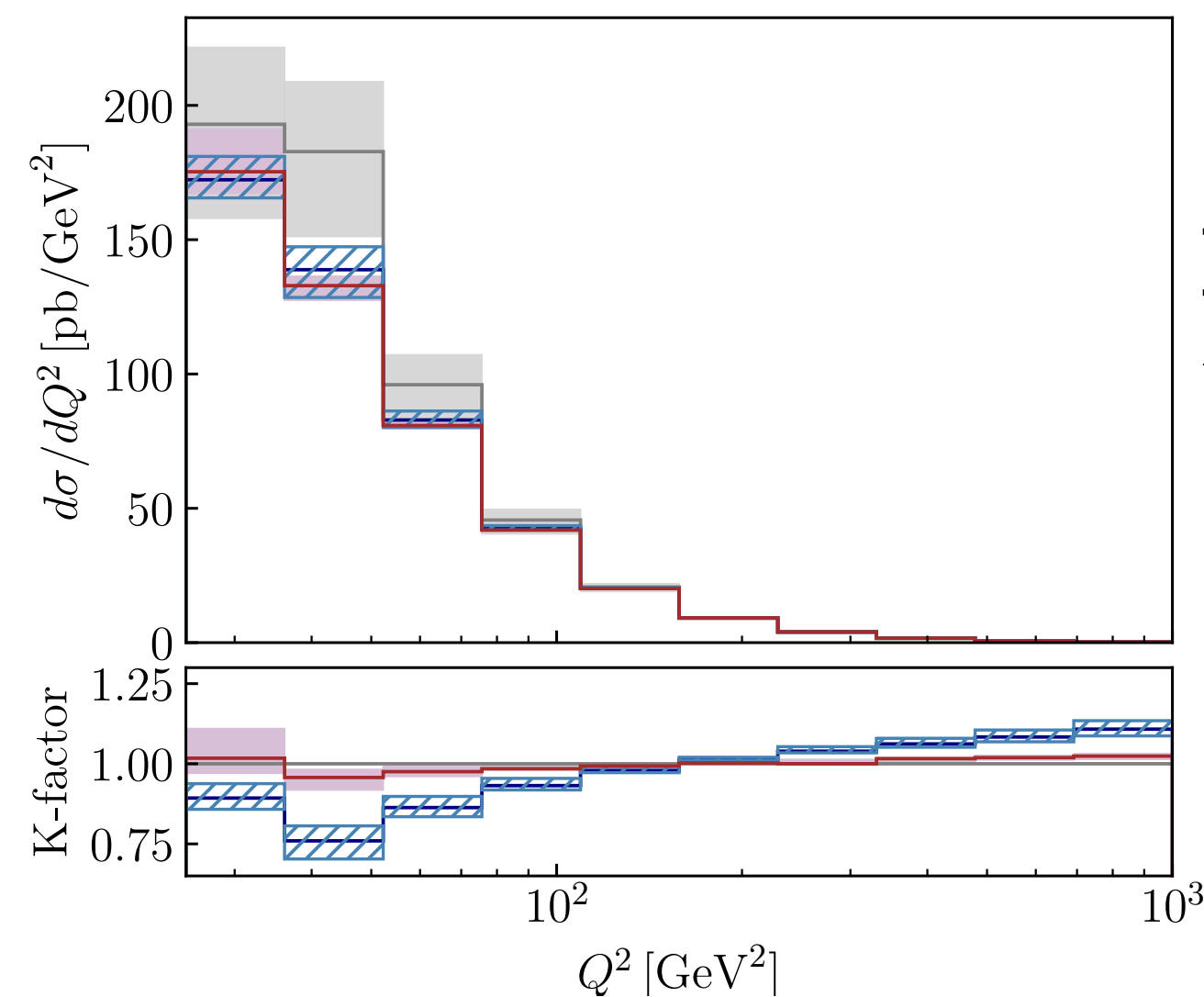
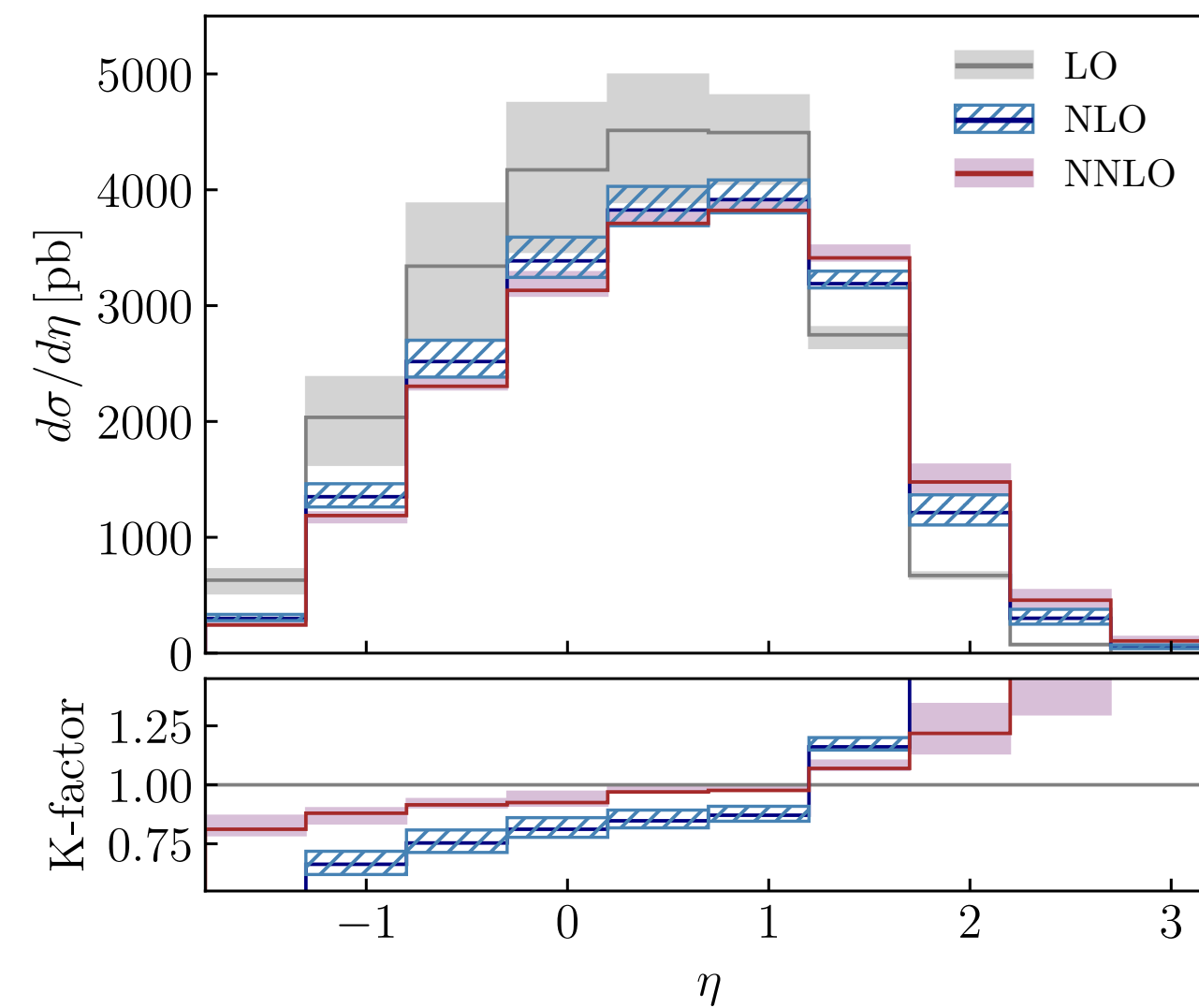
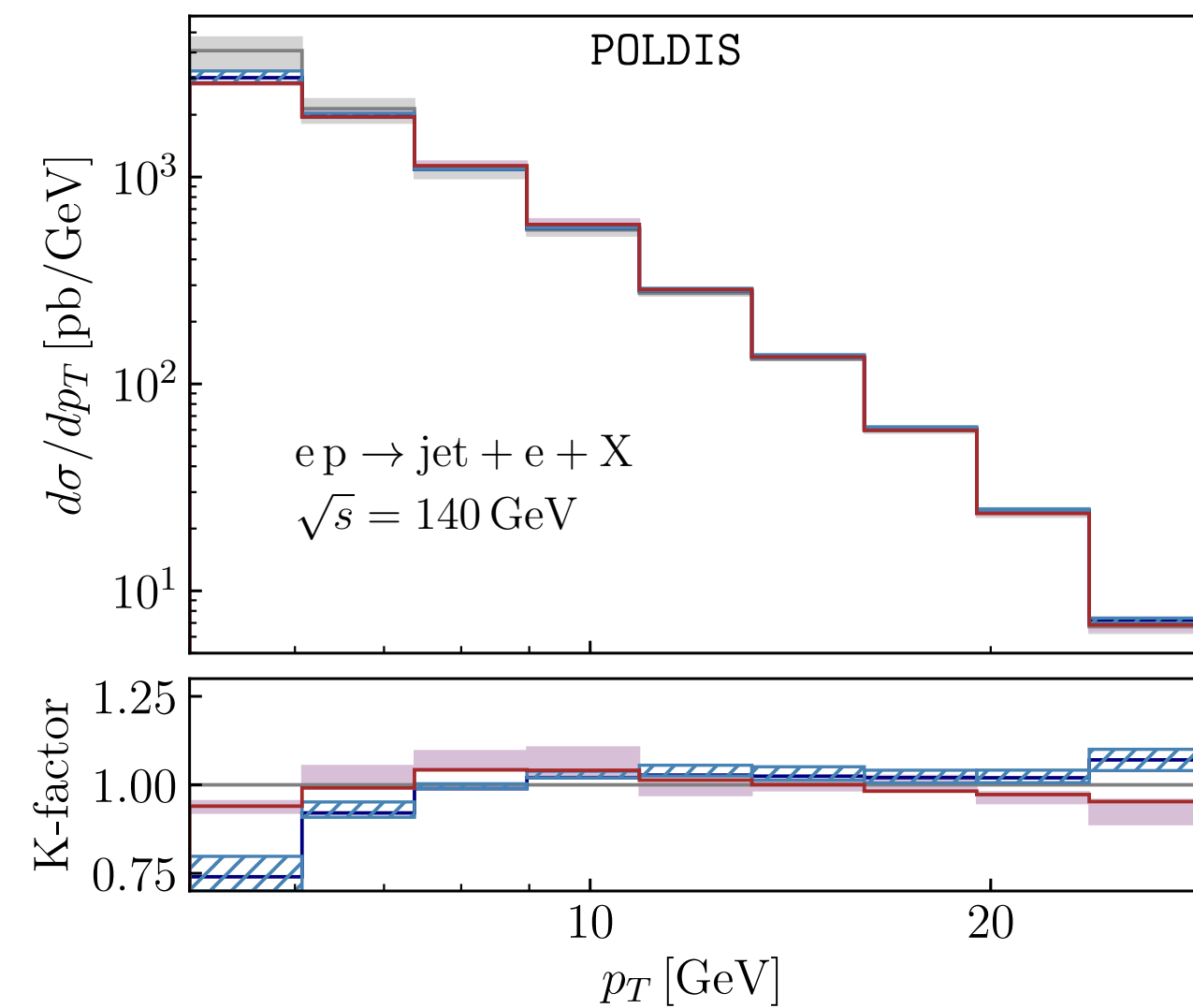
IMPLEMENTATION - PROJECTION TO BORN (P2B)

	Inclusive cross section	Single-Jet	Di-jet
α_S^0	$q\gamma^* \rightarrow q$ LO	$q\gamma^* \rightarrow q$	
α_S^1	$q\gamma^* \rightarrow q$ 1 loop $q\gamma^* \rightarrow qg$ $g\gamma^* \rightarrow q\bar{q}$ NLO	$q\gamma^* \rightarrow q$ 1 loop $q\gamma^* \rightarrow qg$ $g\gamma^* \rightarrow q\bar{q}$	$q\gamma^* \rightarrow qg$ $g\gamma^* \rightarrow q\bar{q}$ LO
α_S^2	$q\gamma^* \rightarrow q$ 2 loops $q\gamma^* \rightarrow qg$ 1 loop $g\gamma^* \rightarrow q\bar{q}$ 1 loop $q\gamma^* \rightarrow qgg$ $q\gamma^* \rightarrow qq\bar{q}$ $g\gamma^* \rightarrow q\bar{q}g$ NNLO	$q\gamma^* \rightarrow q$ 2 loops $q\gamma^* \rightarrow qg$ 1 loop $g\gamma^* \rightarrow q\bar{q}$ 1 loop $q\gamma^* \rightarrow qgg$ $q\gamma^* \rightarrow qq\bar{q}$ $g\gamma^* \rightarrow q\bar{q}g$	$q\gamma^* \rightarrow qg$ 1 loop $g\gamma^* \rightarrow q\bar{q}$ 1 loop $q\gamma^* \rightarrow qgg$ $q\gamma^* \rightarrow qq\bar{q}$ $g\gamma^* \rightarrow q\bar{q}g$ NLO



NNLO SINGLE-JET PRODUCTION IN DIS

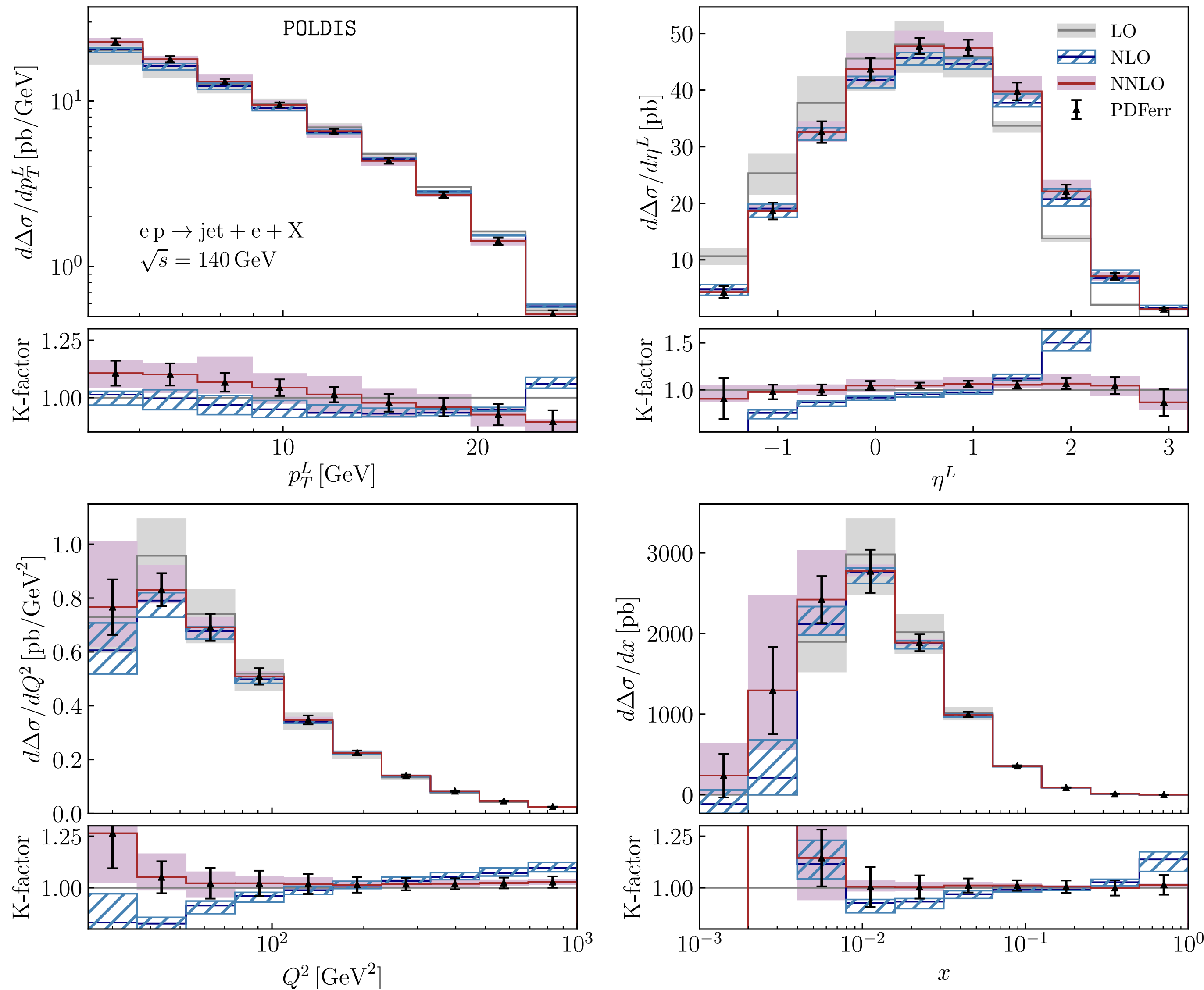
Unpolarized distributions



- Improved convergence at NNLO: reduced NNLO/NLO K-factor
- Shift towards larger rapidities and lower p_T , as the emission of extra partons populates those regions
- Larger NLO and NNLO corrections for the forward and backward regions
- Milder corrections for central rapidity bins

NNLO SINGLE-JET PRODUCTION IN DIS

Polarized distributions

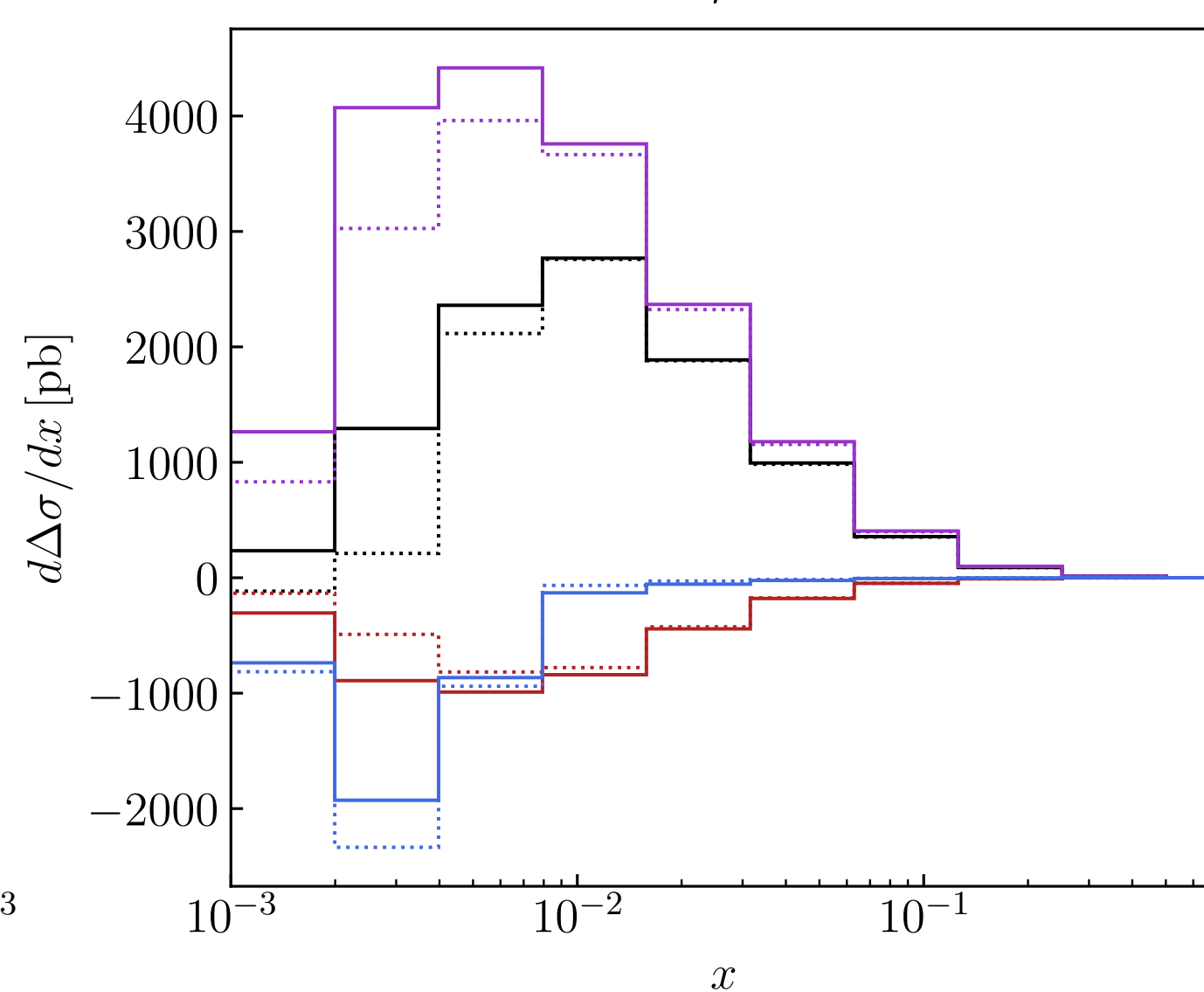
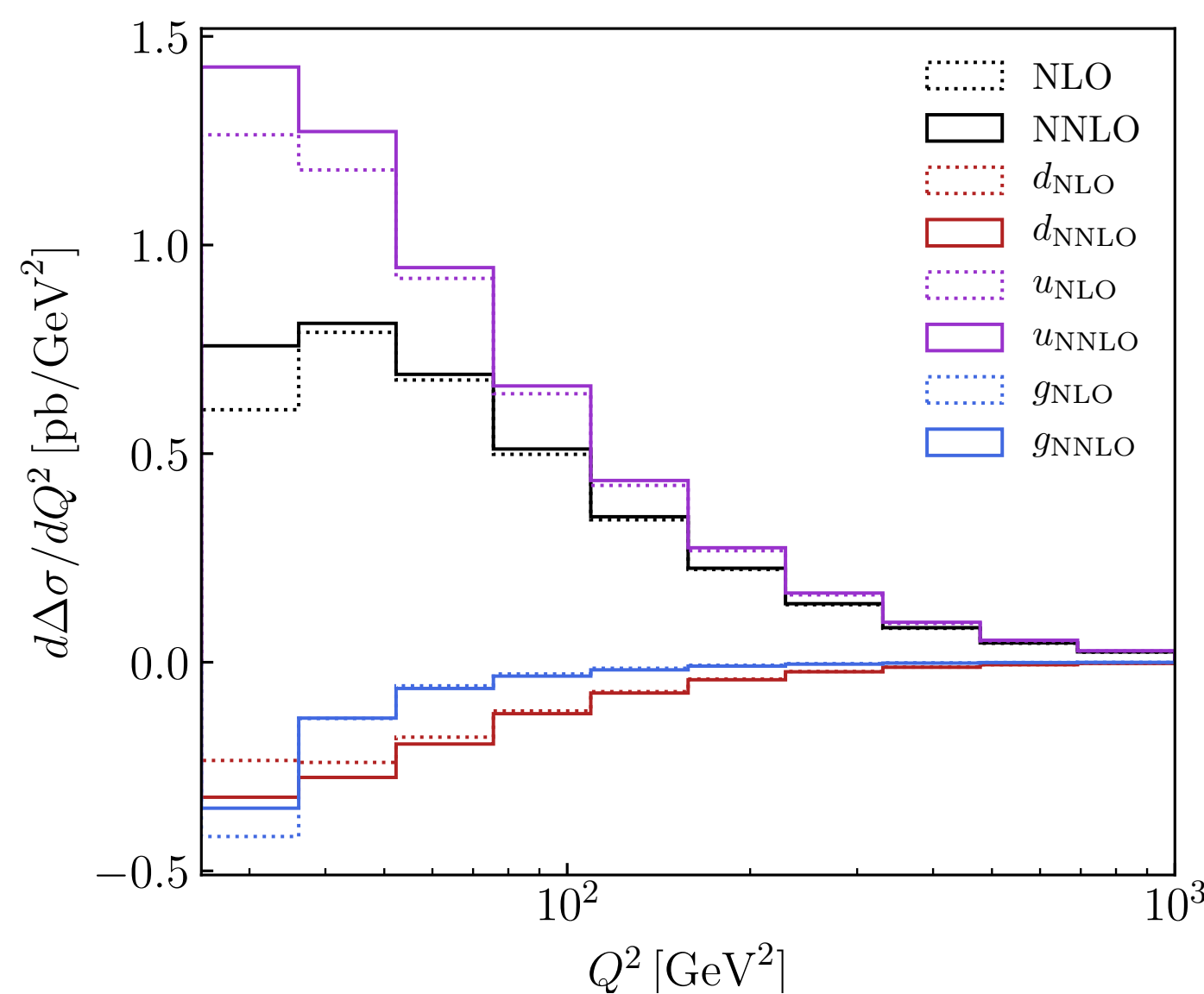
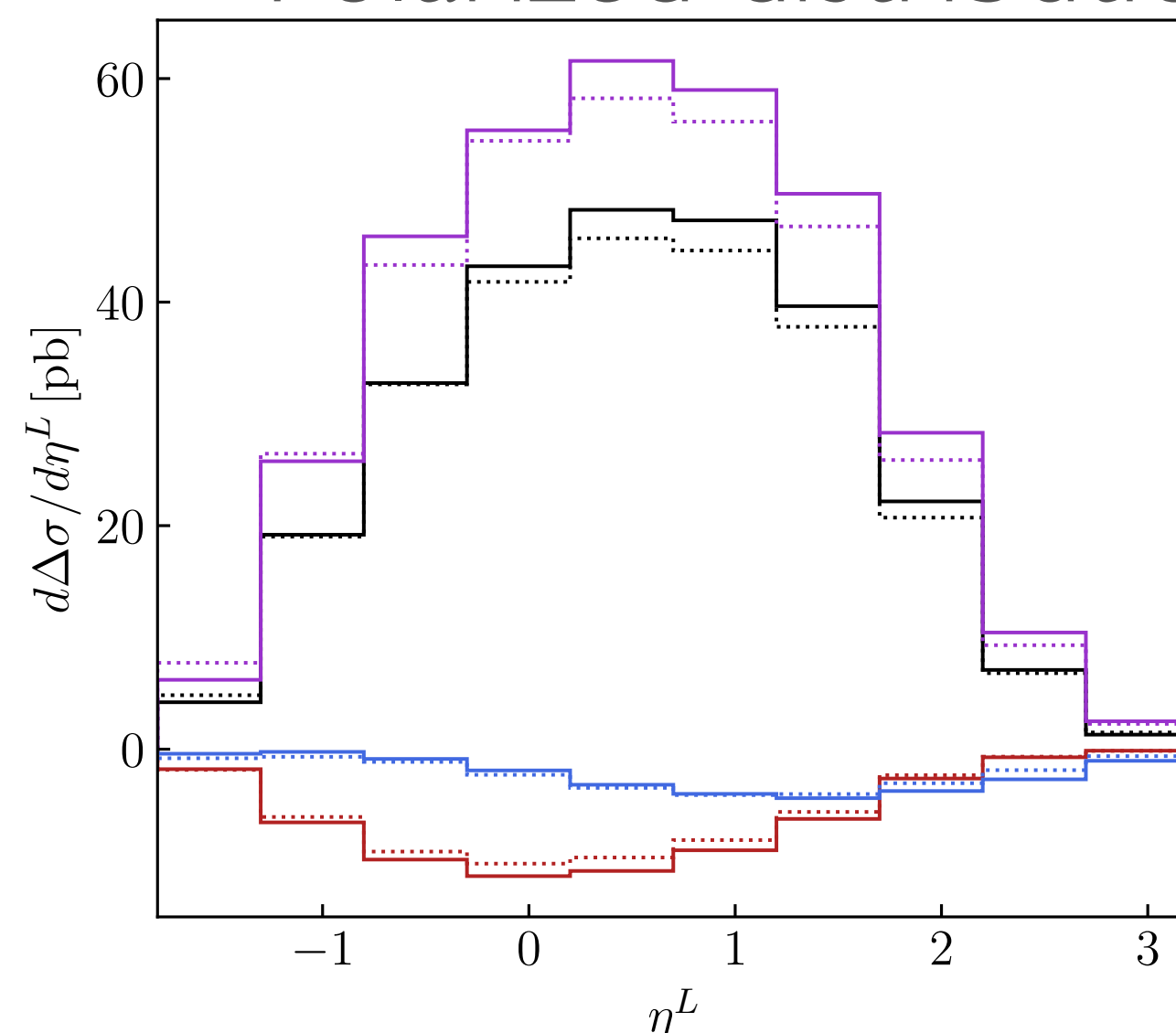
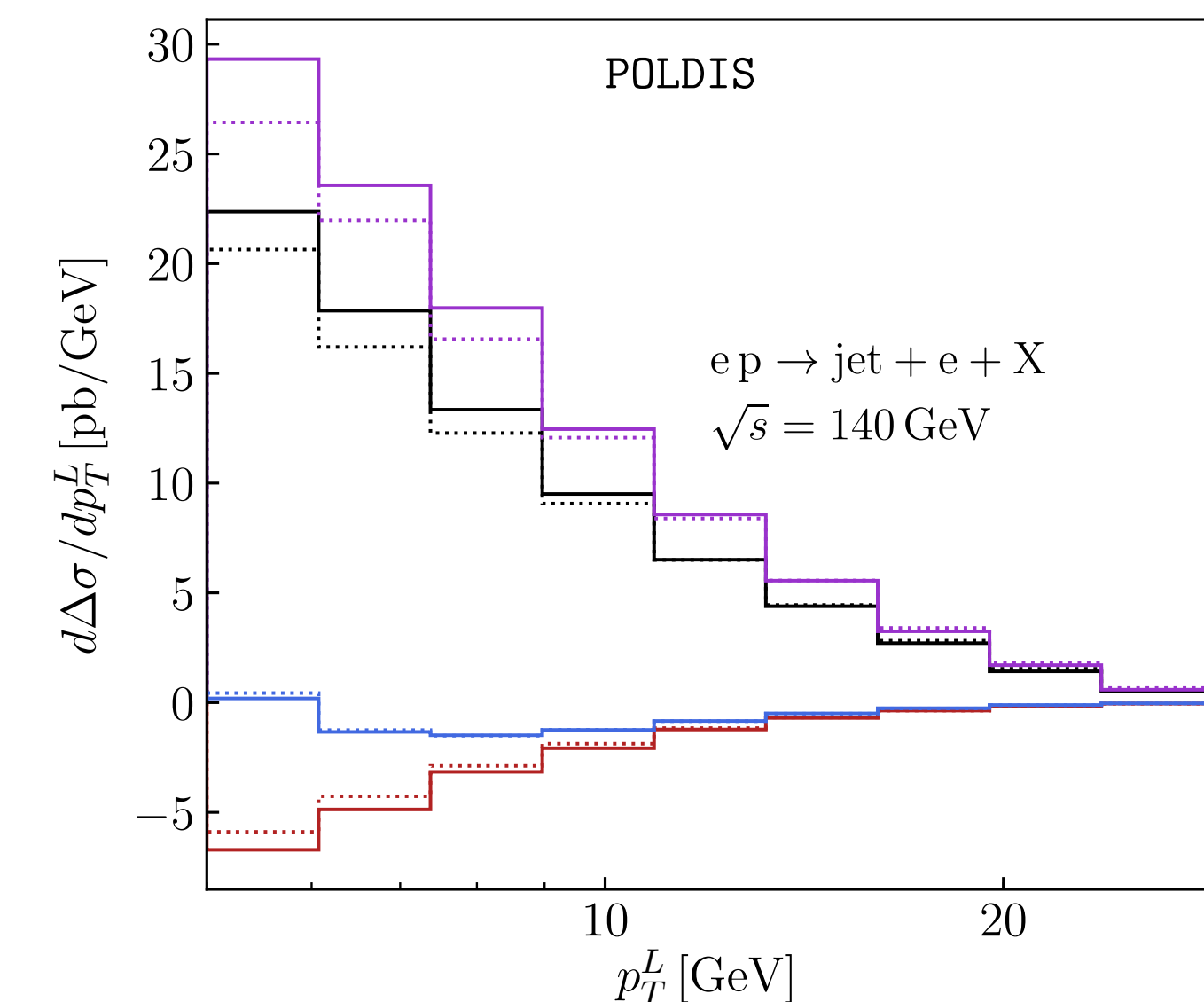


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Stronger scale dependence because of cancellations between partonic channels

NNLO SINGLE-JET PRODUCTION IN DIS

Polarized distributions



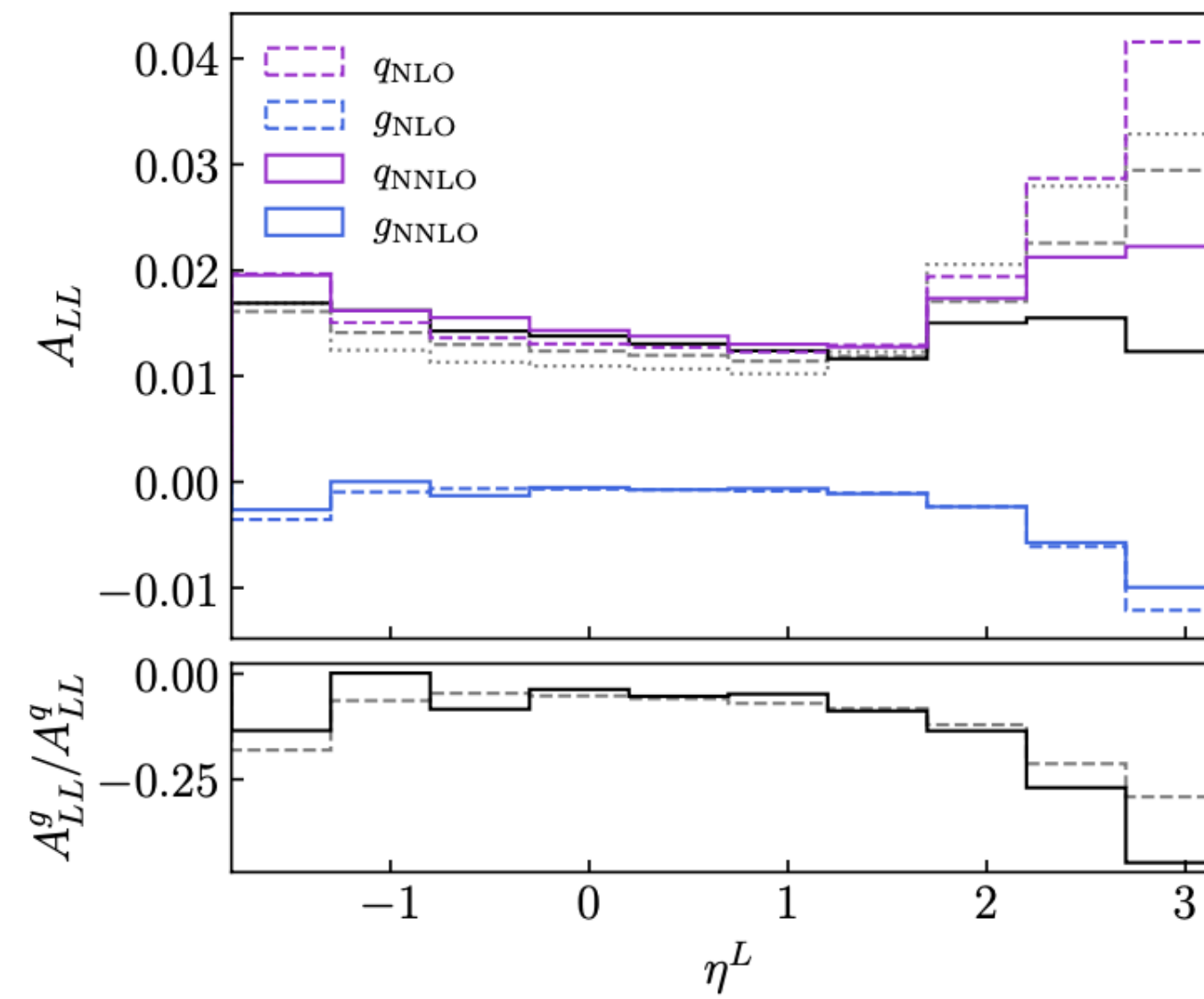
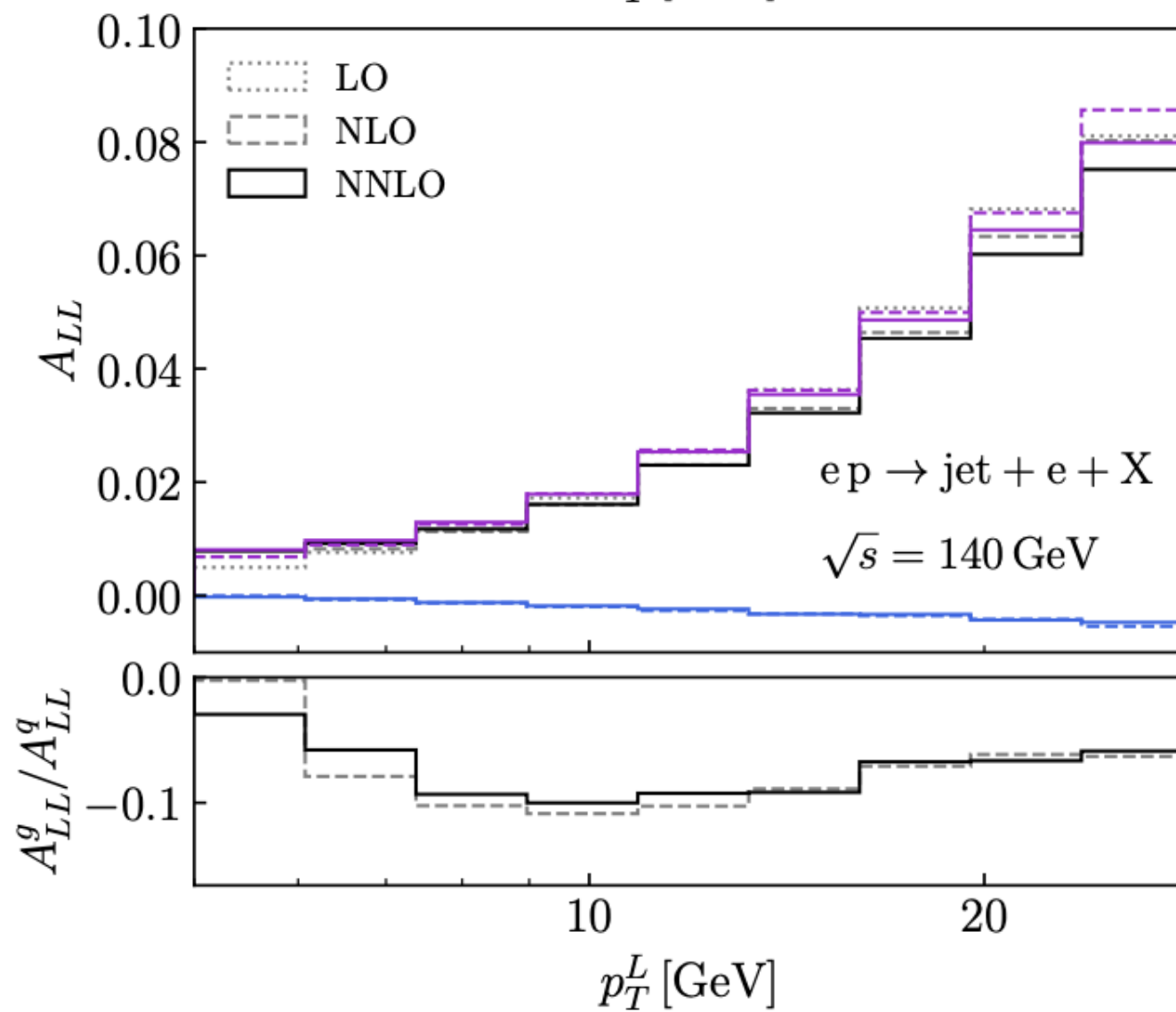
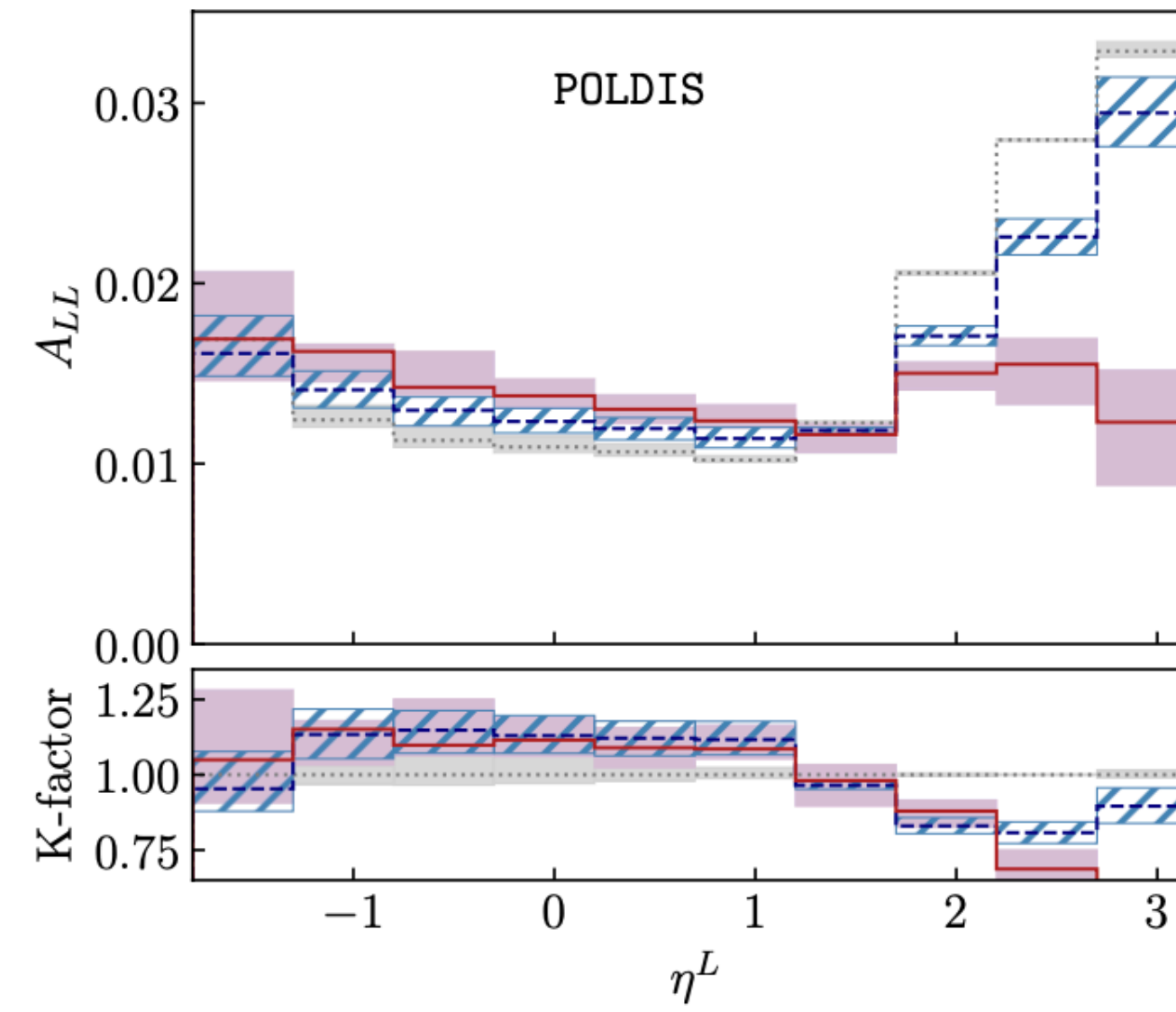
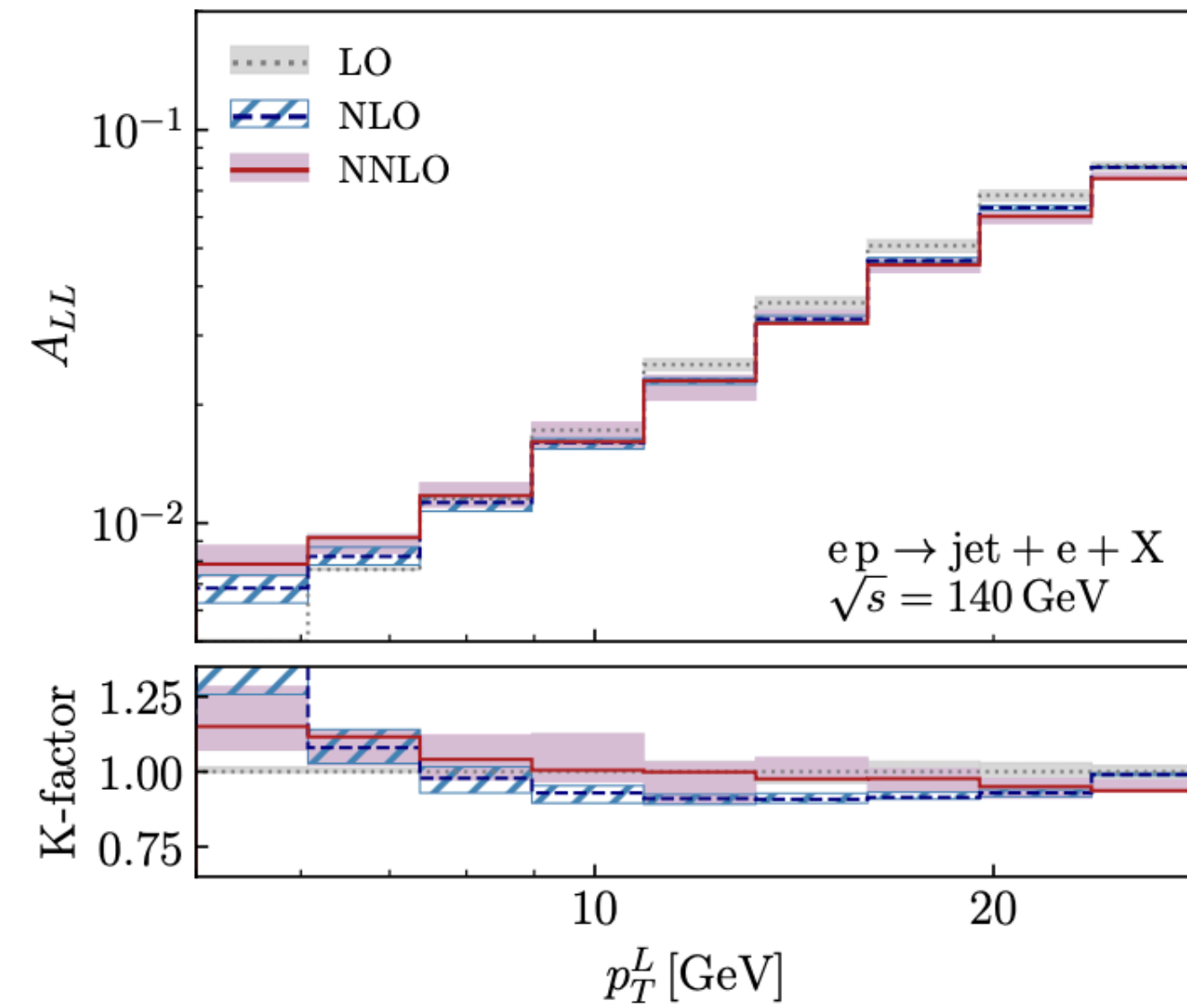
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Stronger scale dependence because of cancellations between partonic channels

NNLO SINGLE-JET PRODUCTION IN DIS

Double spin asymmetries

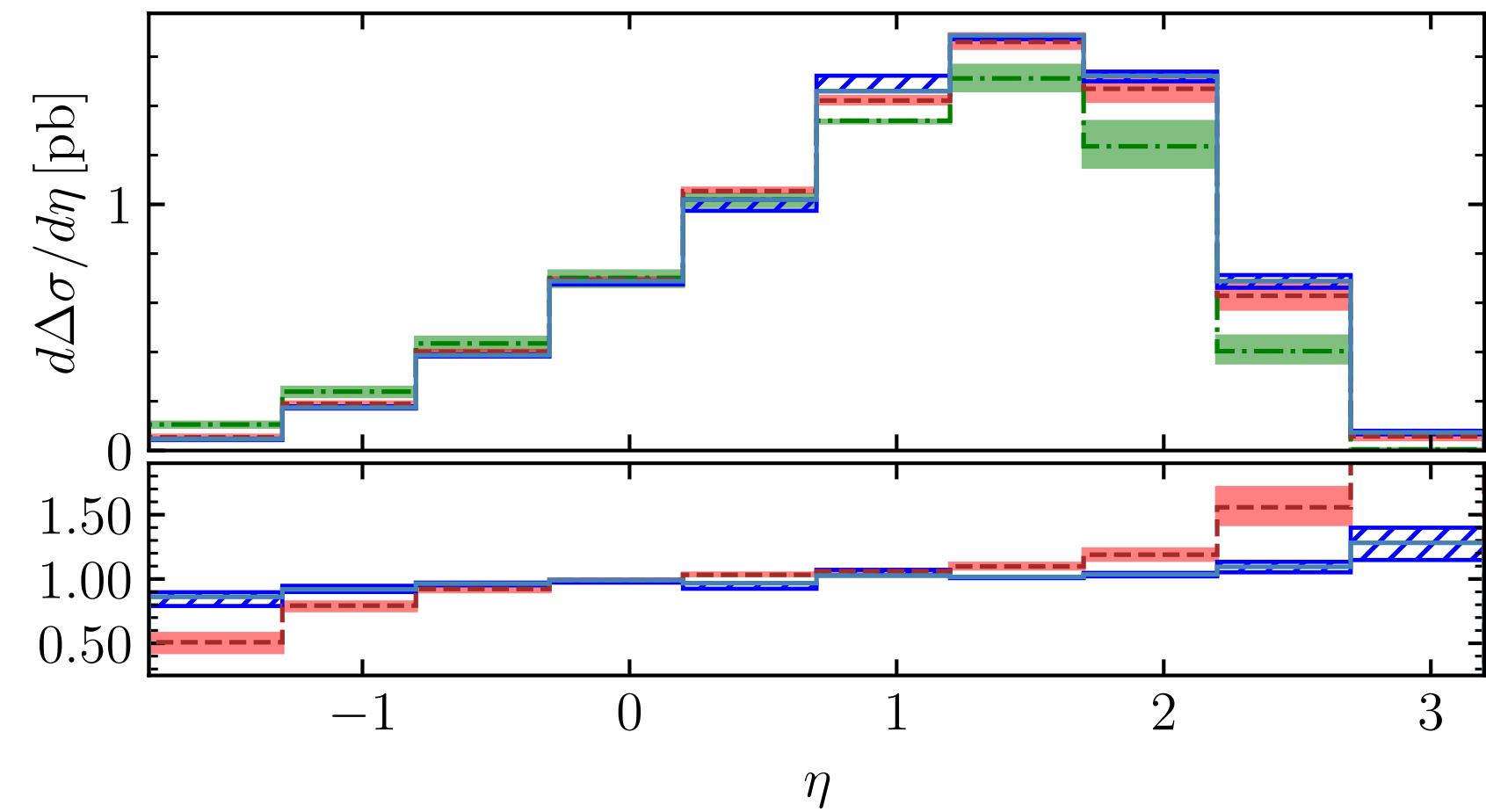
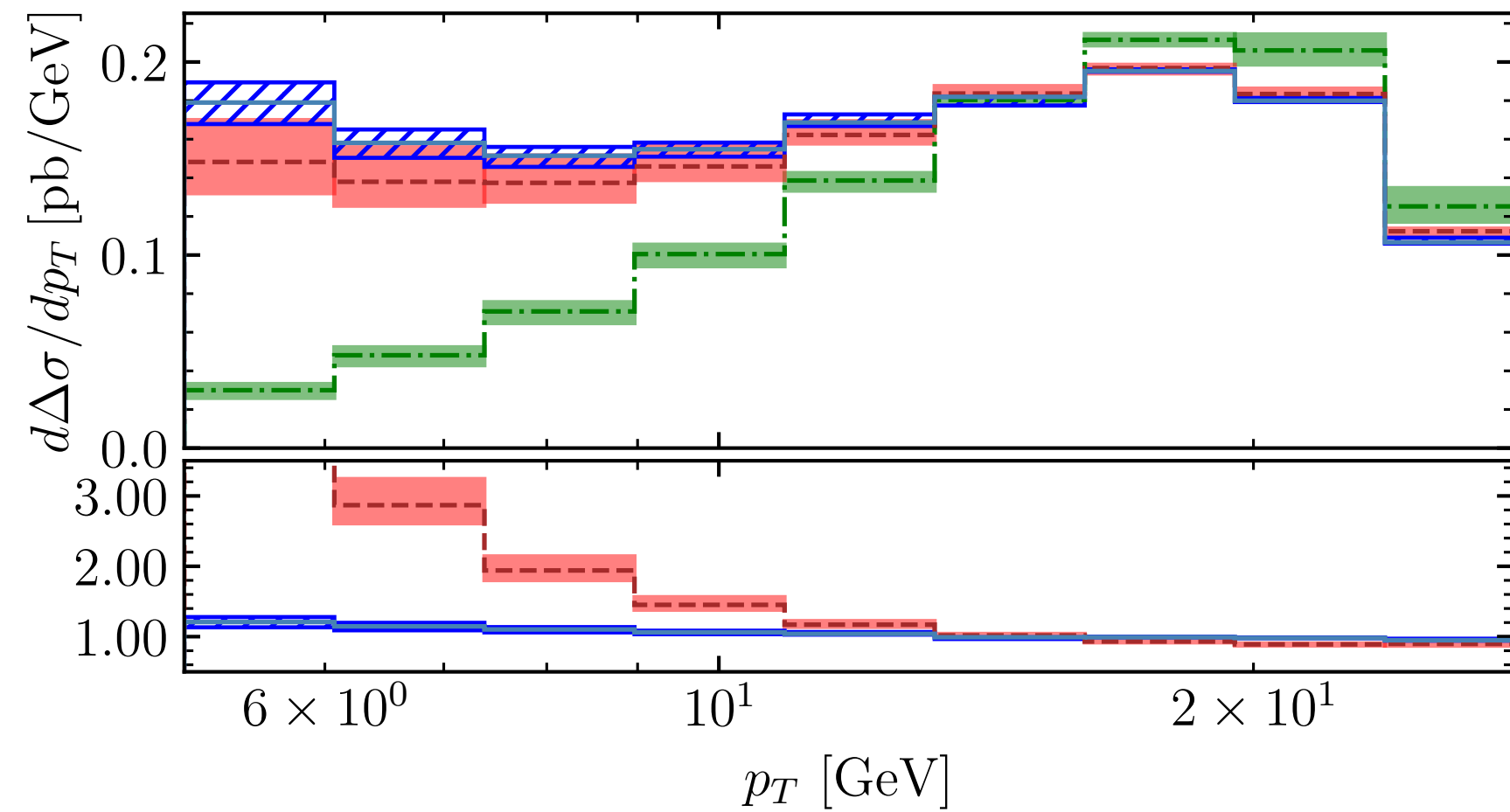
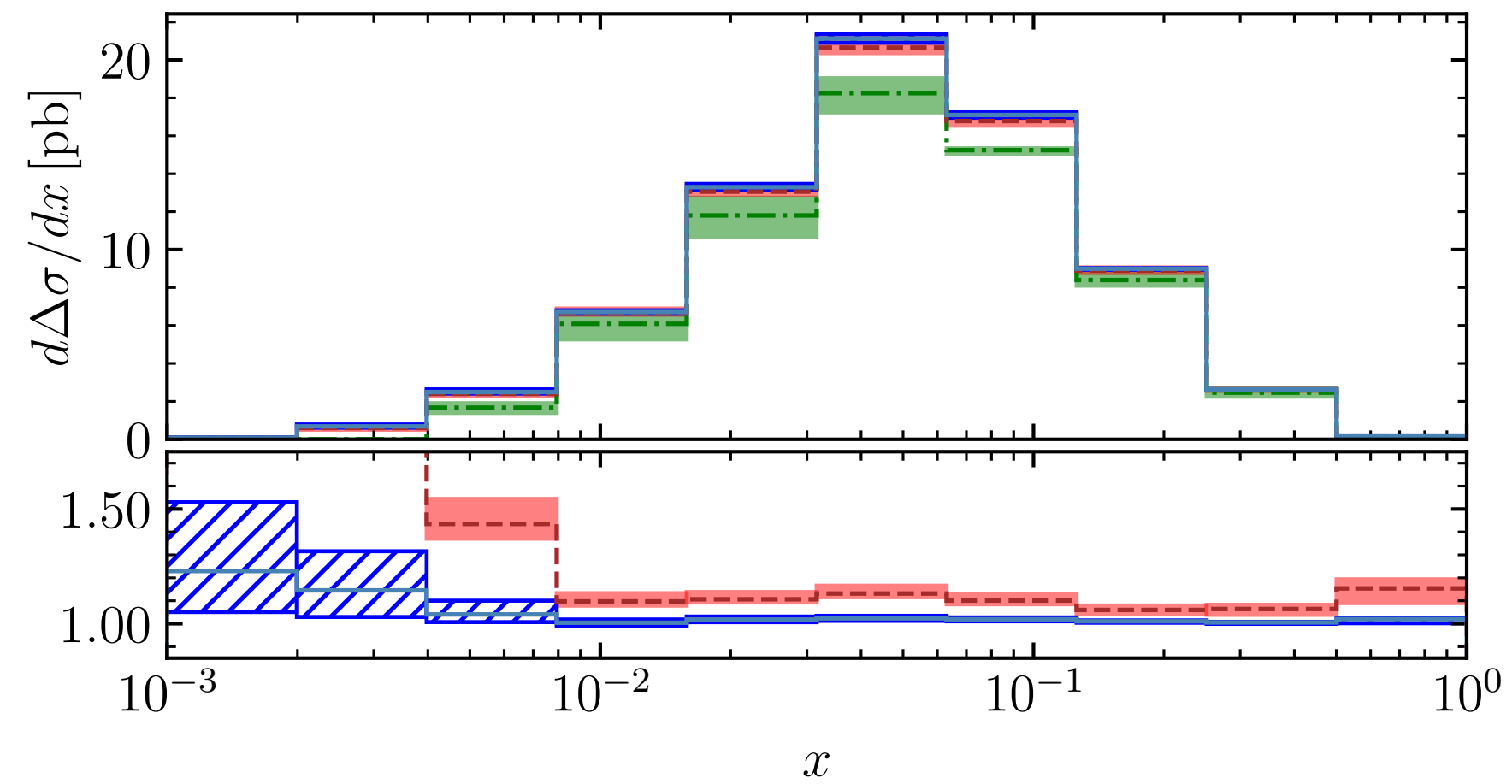
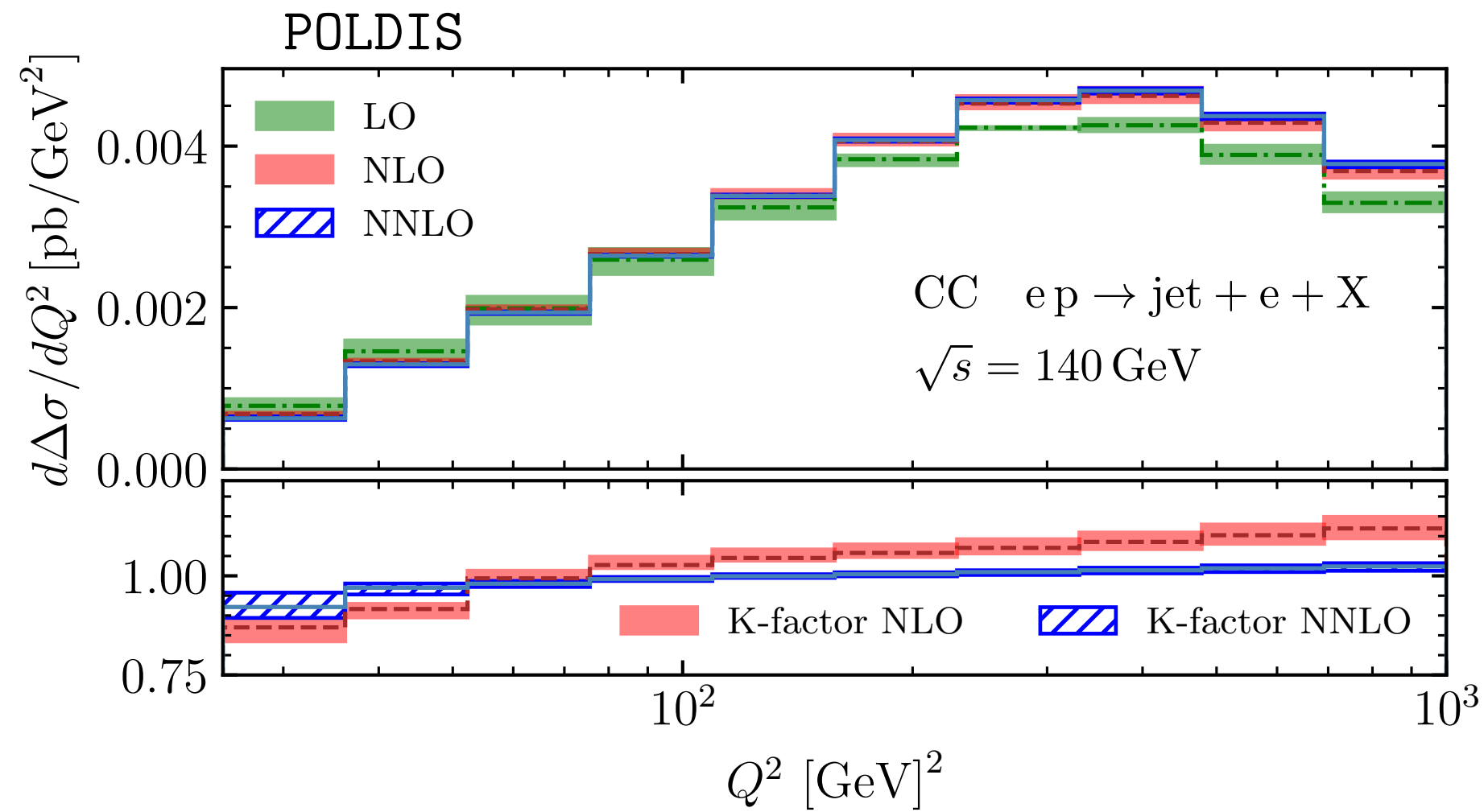
$$\rightarrow A_{LL} = \frac{\Delta\sigma}{\sigma}$$



- Cancellations in the polarized cross section lead to small asymmetries $\mathcal{O} \sim 2\%$
- Different behavior between polarized and unpolarized leads to significant corrections in the asymmetries, even at NNLO

PHENOMENOLOGY - NNLO SINGLE-JET PRODUCTION IN DIS

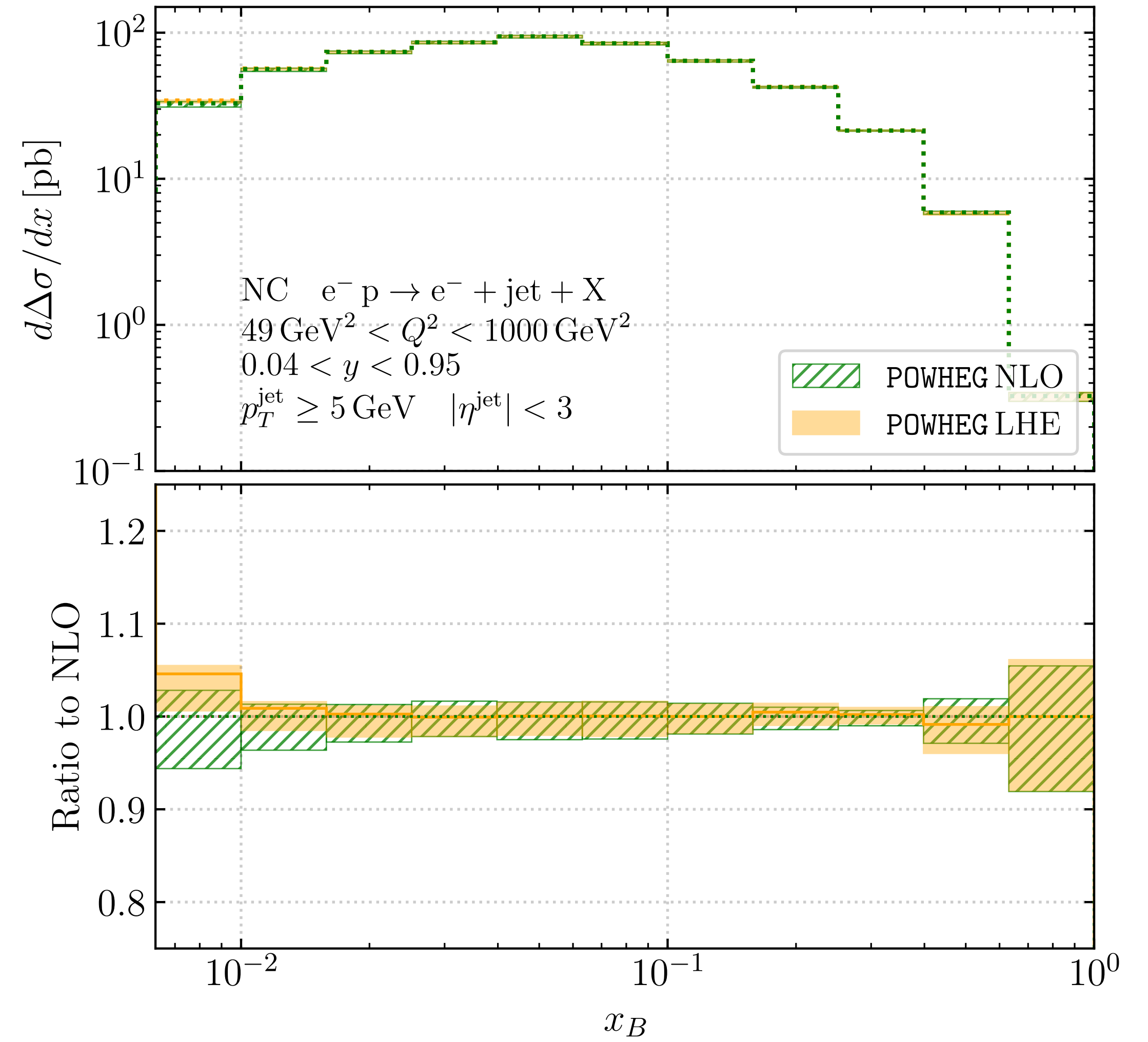
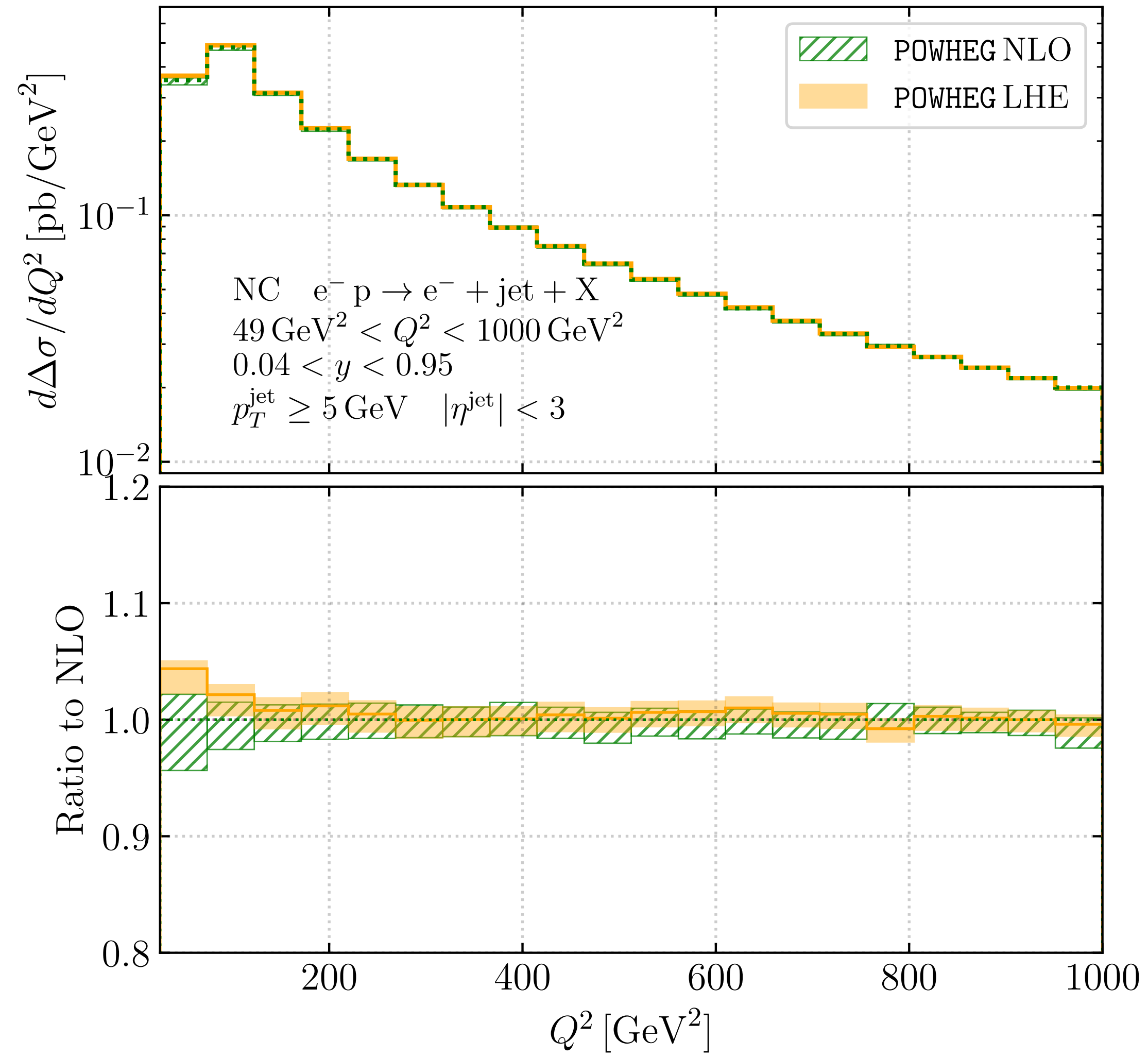
Charged Current DIS



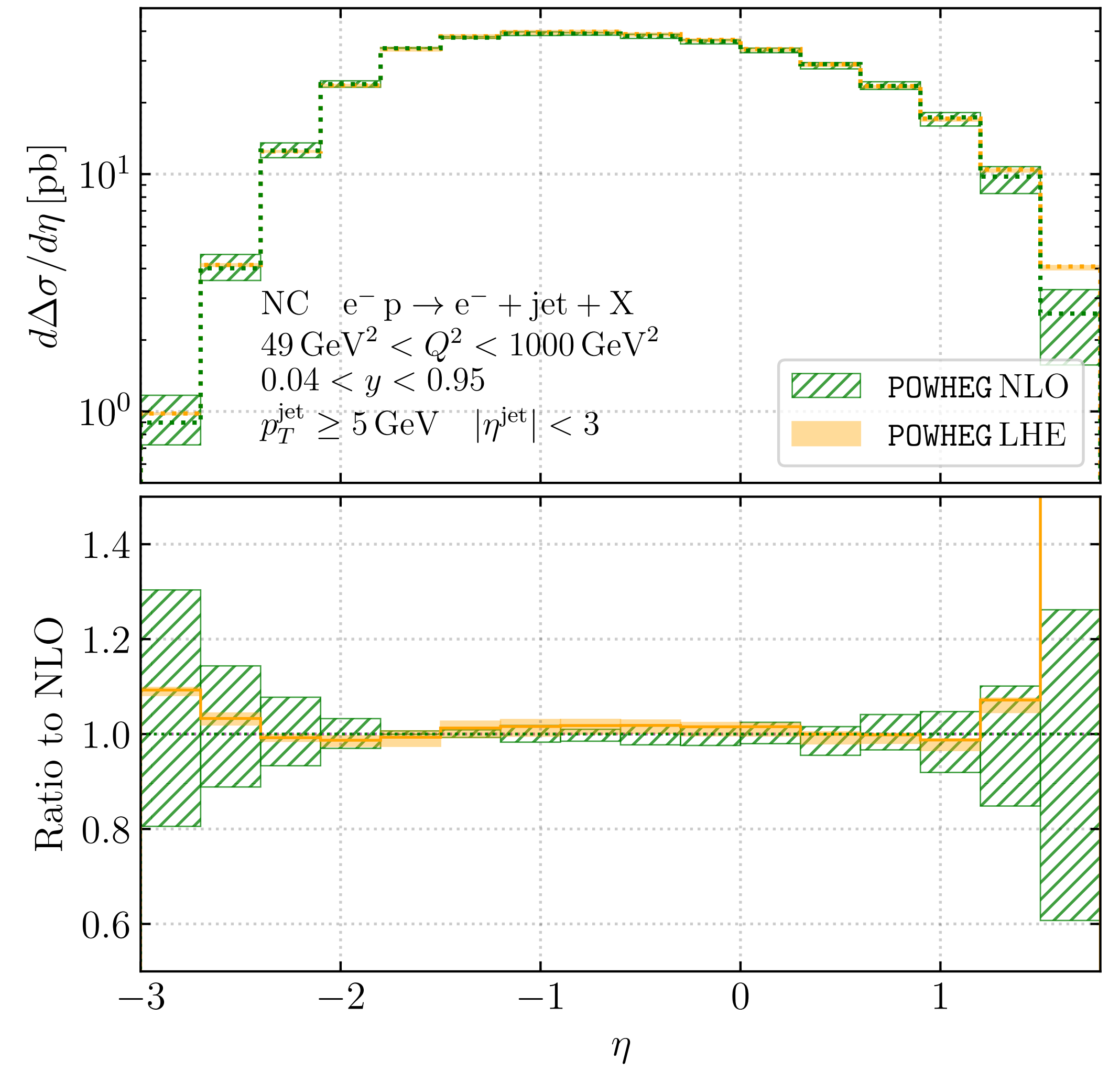
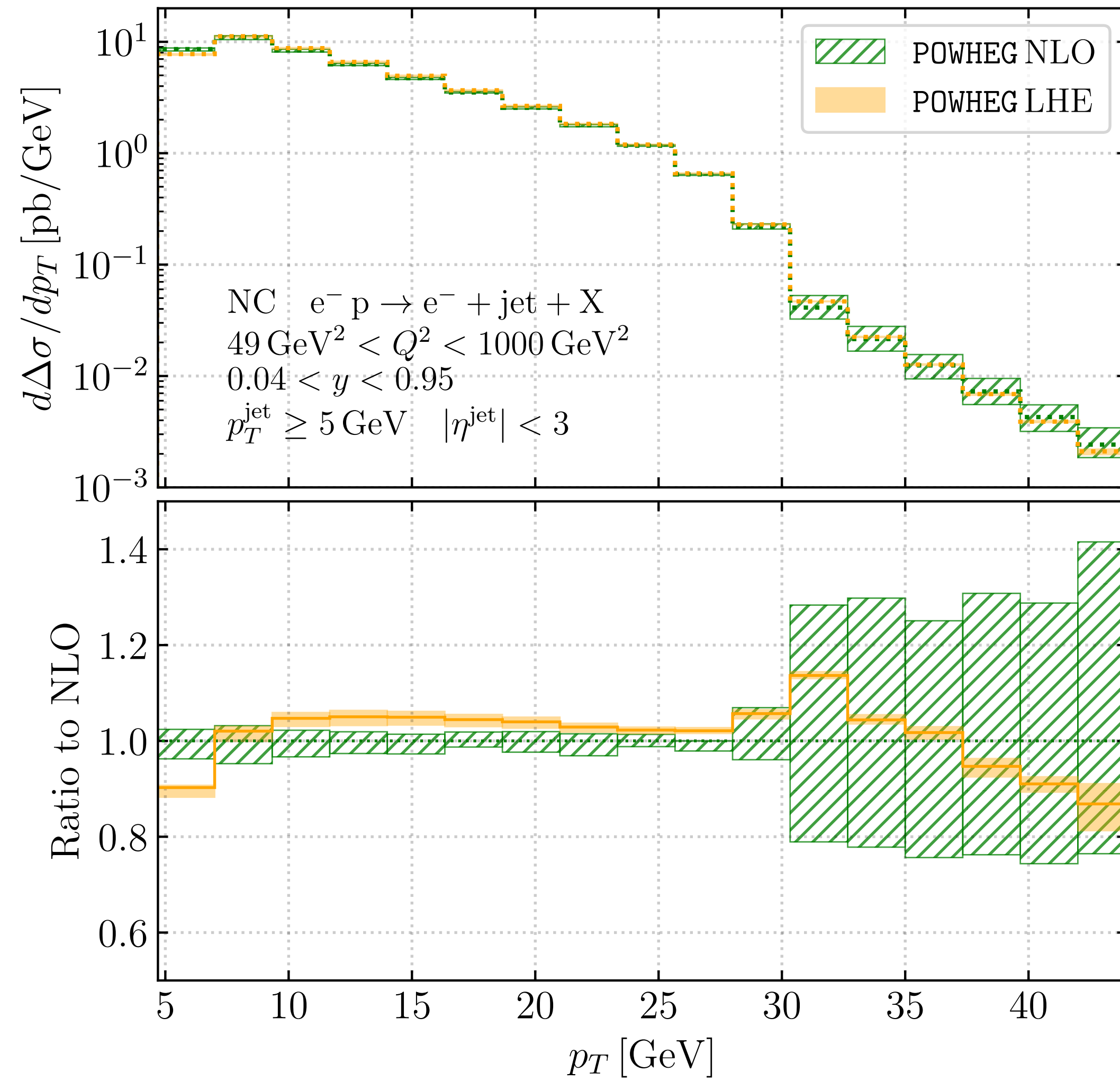
- Suppression at low Q^2 due to the massive propagator.
- LO distribution suppressed at low p_T , since at that order p_T is proportional to Q^2 . Sizable corrections at higher orders.

IB, de Florian, Pedron. Phys.Rev.D 107 (2023)

NLO + PARTON SHOWER - POWHEG



NLO + PARTON SHOWER - POWHEG



NLO + PARTON SHOWER - POWHEG

