



מכון
וַיַּצְמַח
לִמְדָע

WEIZMANN
INSTITUTE
OF SCIENCE

Primordial Gravitational Waves and Pulsar Timing Array Data

E.M., Morgante, Puchades Ibáñez, Ramberg, Ratzinger, Schenk, and Schwaller

JHEP **10** (2023), 171 [arXiv:2306.14856 [hep-ph]]

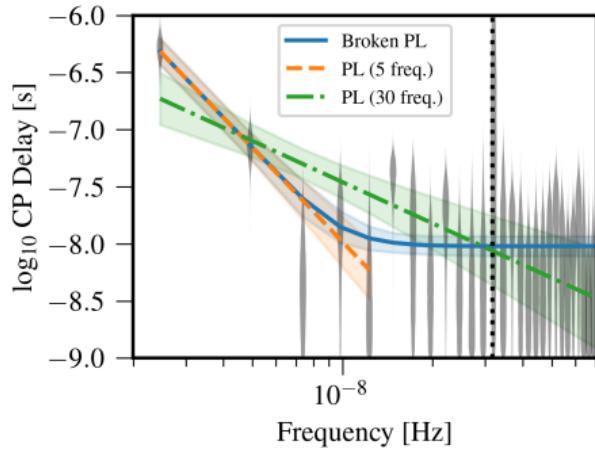
Eric Madge

DESY – November 1, 2023

Motivation

NANOGrav, 12.5 year dataset

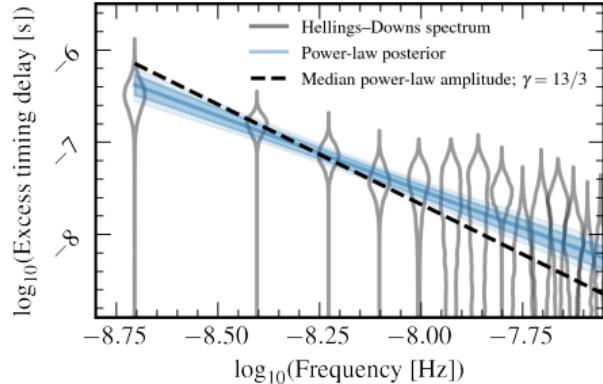
[NANOGrav (APJL 2020)]



astrophysical?

NANOGrav, 15 year dataset

[NANOGrav (APJL 2023)]

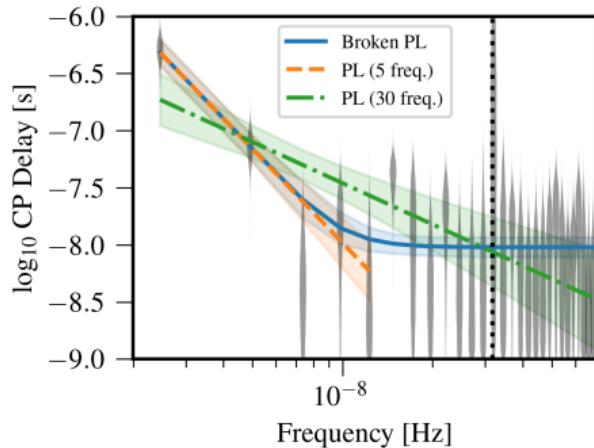


cosmological?

Motivation

NANOGrav, 12.5 year dataset

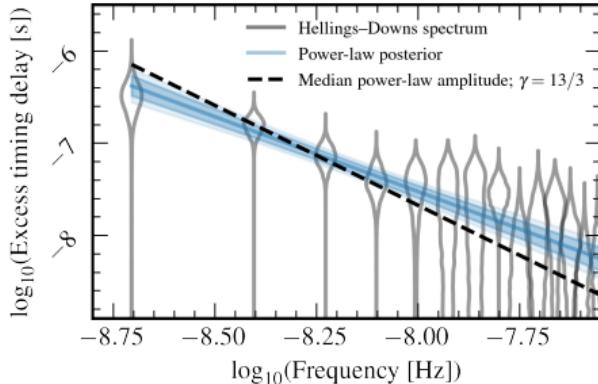
[NANOGrav (APJL 2020)]



astrophysical?

NANOGrav, 15 year dataset

[NANOGrav (APJL 2023)]



SGWB

cosmological?

- here: benchmark models
 cosmological constraints

Outlook

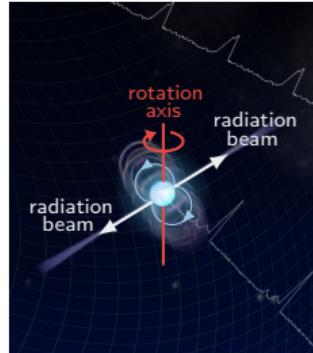
1. Pulsar Timing Arrays
2. Cosmological Phase Transitions
3. Domain Walls
4. Bosonic Instabilities
5. Conclusions

Pulsar Timing Arrays

1. Pulsar Timing Arrays
2. Cosmological Phase Transitions
3. Domain Walls
4. Bosonic Instabilities
5. Conclusions

Pulsar Timing Arrays

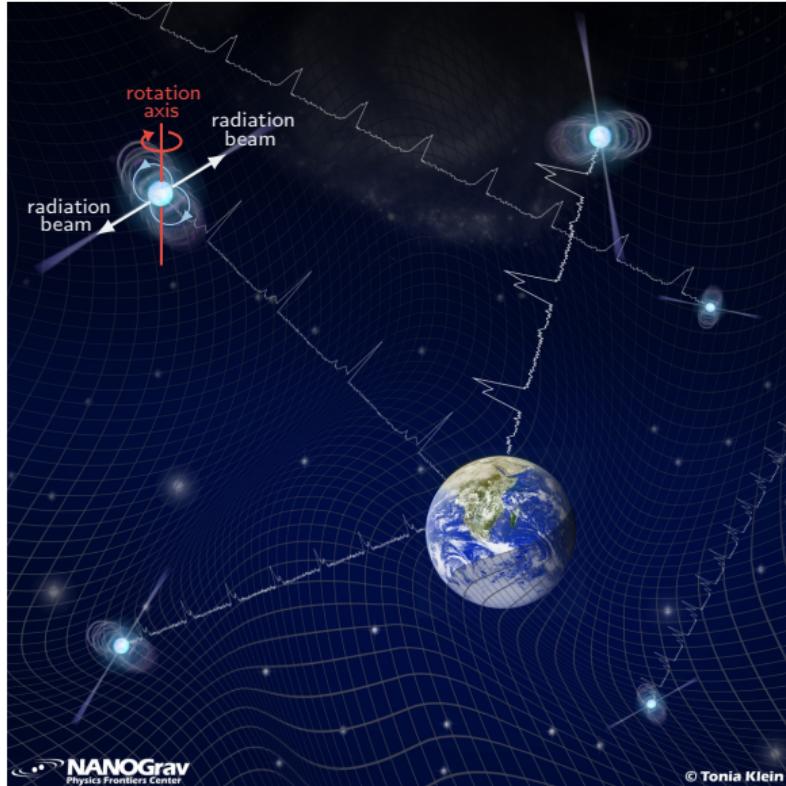
- pulsar = rapidly rotating neutron star
➡ cosmic lighthouse



[Credit: T. Klein, NANOGrav; modified]

Pulsar Timing Arrays

- pulsar = rapidly rotating neutron star
 ➡ cosmic lighthouse
- pulsar timing arrays (PTAs)
GW detection via correlations in timing residuals ➡ $f \sim \text{nHz}$



[Credit: T. Klein, NANOGrav; modified]

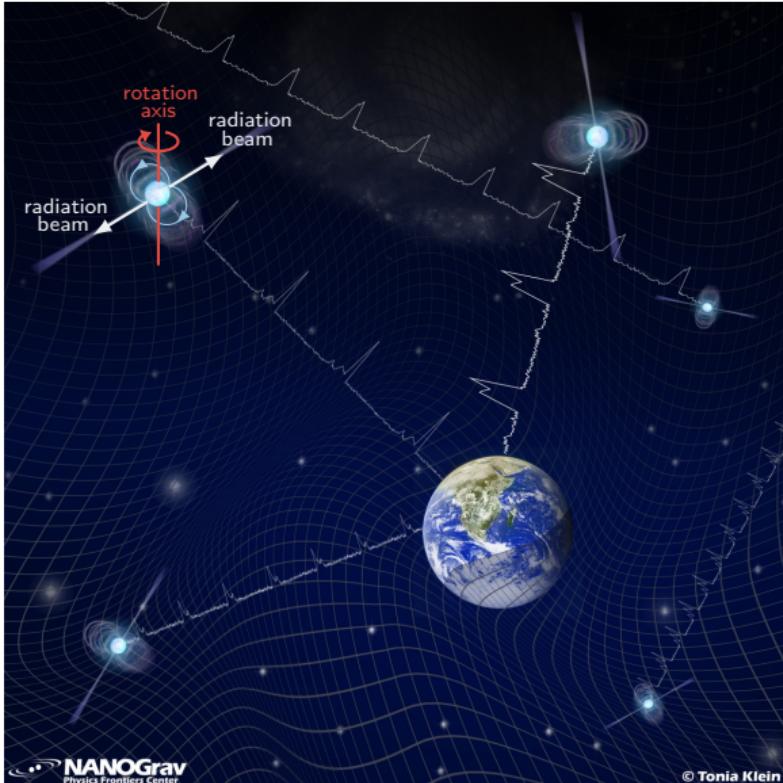
Pulsar Timing Arrays

- pulsar = rapidly rotating neutron star
 ➡ cosmic lighthouse

- pulsar timing arrays (PTAs)
GW detection via correlations in timing residuals ➡ $f \sim \text{nHz}$

- operating PTAs:

collaboration	N_P	T_{obs}	ref.
EPTA	25	25 yrs	[A&A 2023]
NANOGrav	67	15 yrs	[APJL 2023]
PPTA	30	18 yrs	[PASA 2023]
InPTA	14	3.5 yrs	[PASA 2022]
CPTA	57	3.5 yrs	[in prep.]
IPTA	65	20 yrs	[MNRAS 2019]



[Credit: T. Klein, NANOGrav; modified]

PTA Signal

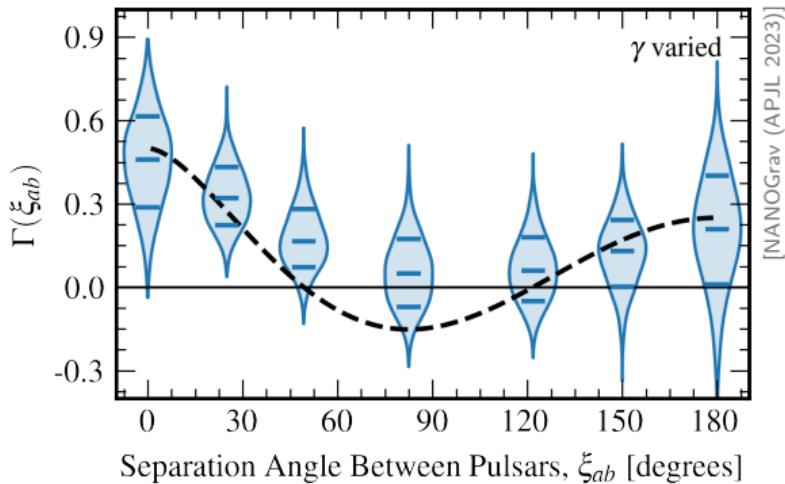
- since Sept. 2020:
strong evidence for common red-noise process

[NANOGrav (APJL 2020); PPTA (APJL 2021); EPTA (MNRAS 2021); IPTA (MNRAS 2022)]

- since June 28, 2023:
strong evidence for Hellings-Downs correlations

[NANOGrav (APJL 2023); EPTA (A&A 2023); PPTA (APJL 2023); CPTA (RAA 2023)]

⇒ GW background



PTA Signal

- since Sept. 2020:
strong evidence for common red-noise process

[NANOGrav (APJL 2020); PPTA (APJL 2021); EPTA (MNRAS 2021); IPTA (MNRAS 2022)]

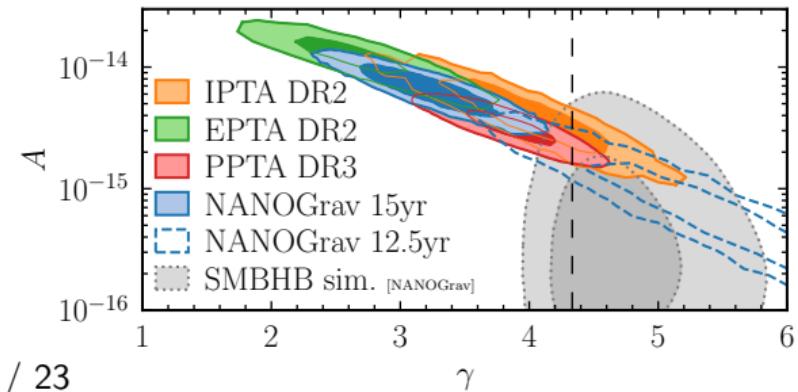
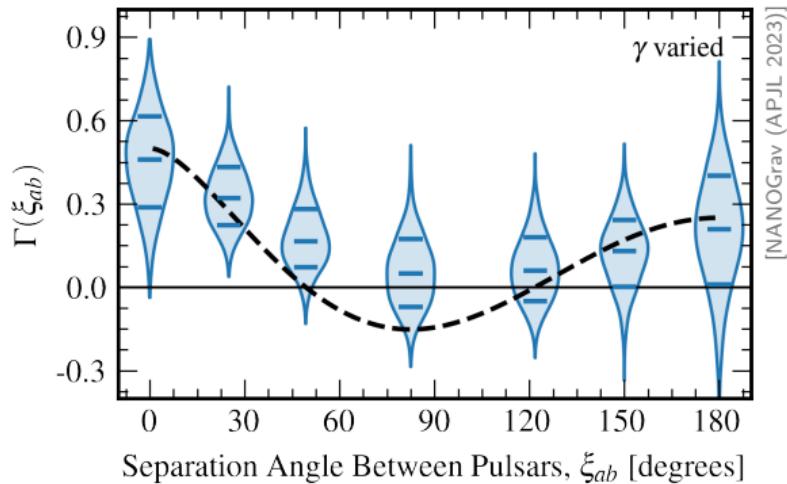
- since June 28, 2023:
strong evidence for Hellings-Downs correlations

[NANOGrav (APJL 2023); EPTA (A&A 2023); PPTA (APJL 2023); CPTA (RAA 2023)]

⇒ GW background

- consistent with expected SGWB from SMBHBs:

$$h_c = A \left(\frac{f}{1 \text{ yr}^{-1}} \right)^{\frac{3-\gamma}{2}} \quad \text{with } \gamma = \frac{13}{3}$$



Primordial Sources of SGWBs

phase transitions, cosmic strings, domain walls, bosonic instabilities, inflation, . . .

Primordial Sources of SGWBs

phase transitions, cosmic strings, domain walls, bosonic instabilities, inflation, . . .

○ frequency:

$$f_0 \sim \frac{a_0}{a_*} H_* \underbrace{\frac{f_*}{H_*}}_{\lesssim 1} \sim 1 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{f_*}{H_*}$$

Primordial Sources of SGWBs

phase transitions, cosmic strings, domain walls, bosonic instabilities, inflation, . . .

○ frequency: $f_0 \sim \frac{a_0}{a_*} H_* \underbrace{\frac{f_*}{H_*}}_{\lesssim 1} \sim 1 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{f_*}{H_*}$

○ amplitude: $\rho_{\text{GW}} \sim M_{\text{pl}}^2 \langle \dot{h}_{ij} \dot{h}_{ij} \rangle, \quad \square \bar{h}_{\mu\nu} \sim \frac{T_{\mu\nu}}{M_{\text{pl}}^2}$

$$\implies \Omega_{\text{GW}} \sim \left(\frac{a_*}{a_0}\right)^4 \frac{1}{\rho_c} \frac{\rho_{\text{source}}^2}{f_*^2 M_{\text{pl}}^2} \sim \left(\frac{a_*}{a_0}\right)^4 \left(\frac{H_0}{H_*}\right)^2 \left(\frac{H_*}{f_*}\right)^2 \Omega_{\text{source}}^2$$

Primordial Sources of SGWBs

phase transitions, cosmic strings, domain walls, bosonic instabilities, inflation, . . .

○ frequency: $f_0 \sim \frac{a_0}{a_*} H_* \underbrace{\frac{f_*}{H_*}}_{\lesssim 1} \sim 1 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{f_*}{H_*}$

○ amplitude: $\rho_{\text{GW}} \sim M_{\text{pl}}^2 \langle \dot{h}_{ij} \dot{h}_{ij} \rangle, \quad \square \bar{h}_{\mu\nu} \sim \frac{T_{\mu\nu}}{M_{\text{pl}}^2}$

$$\implies \Omega_{\text{GW}} \sim \left(\frac{a_*}{a_0}\right)^4 \frac{1}{\rho_c} \frac{\rho_{\text{source}}^2}{f_*^2 M_{\text{pl}}^2} \sim \left(\frac{a_*}{a_0}\right)^4 \left(\frac{H_0}{H_*}\right)^2 \left(\frac{H_*}{f_*}\right)^2 \Omega_{\text{source}}^2$$

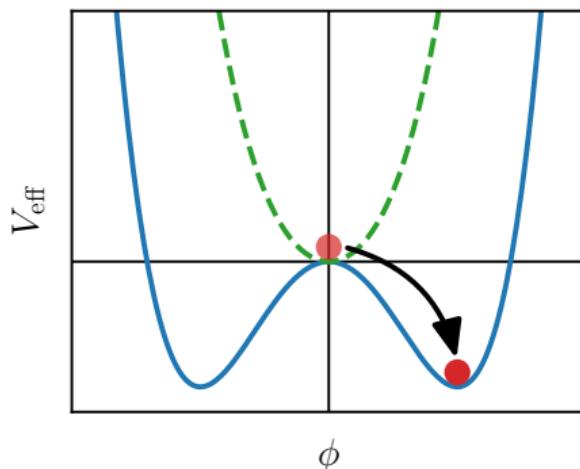
$$\implies \text{for PTAs: } T_* \sim \text{few MeV}, \quad \Omega_{\text{source}} \gtrsim 0.1$$

Cosmological Phase Transitions

1. Pulsar Timing Arrays
2. Cosmological Phase Transitions
3. Domain Walls
4. Bosonic Instabilities
5. Conclusions

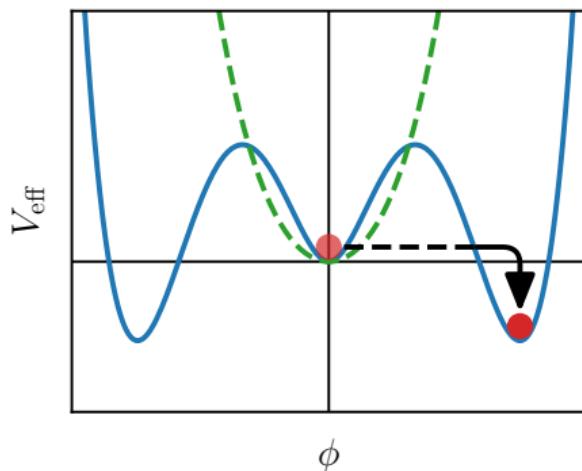
Cosmological Phase Transitions

- thermal corrections typically restore spontaneously broken symmetries at high temperatures
➡ symmetry breaking phase transition
- can be **crossover** or first-order



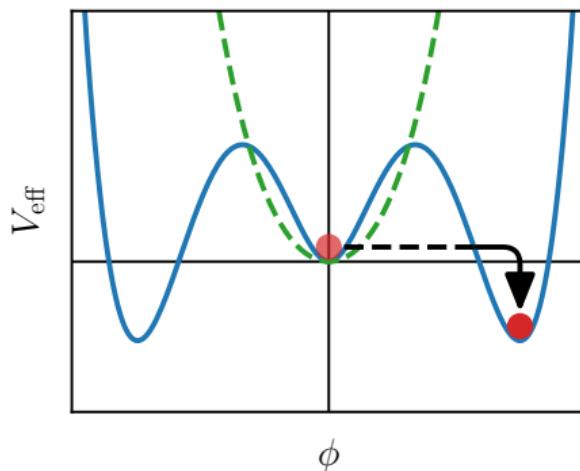
Cosmological Phase Transitions

- thermal corrections typically restore spontaneously broken symmetries at high temperatures
➡ symmetry breaking phase transition
- can be crossover or first-order



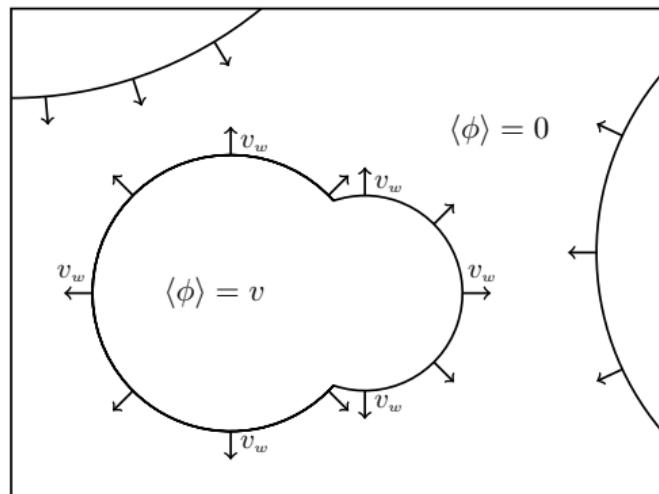
Cosmological Phase Transitions

- thermal corrections typically restore spontaneously broken symmetries at high temperatures
➡ symmetry breaking phase transition
- can be crossover or first-order



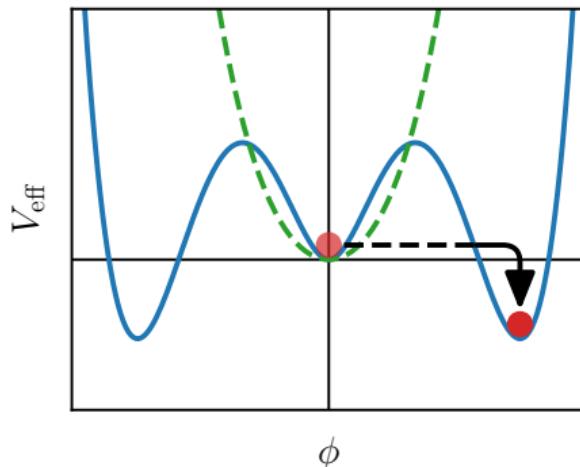
GW production:

1. vacuum bubble collisions



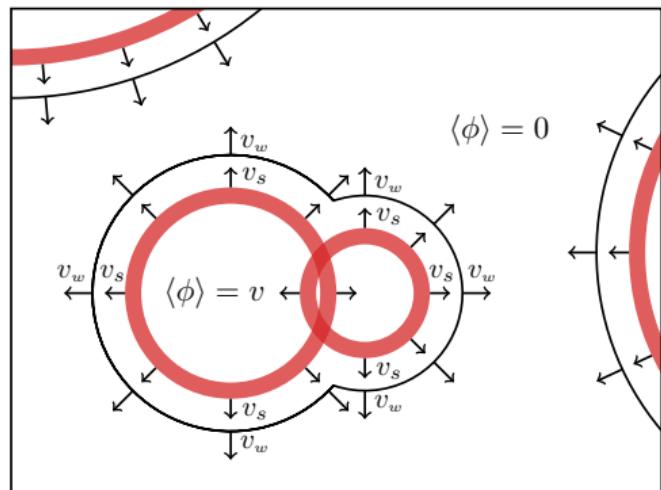
Cosmological Phase Transitions

- thermal corrections typically restore spontaneously broken symmetries at high temperatures
➡ symmetry breaking phase transition
- can be crossover or first-order



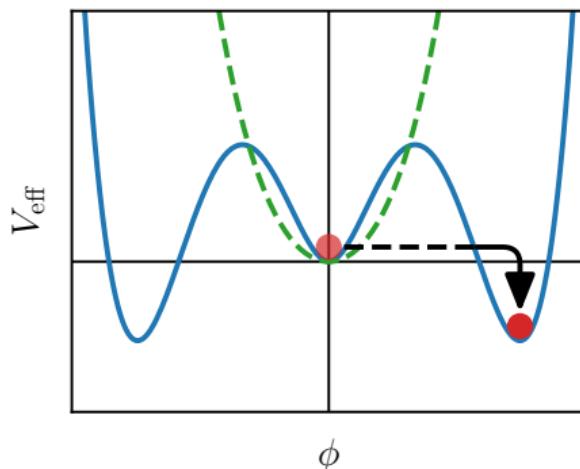
GW production:

1. vacuum bubble collisions
2. sound waves collisions



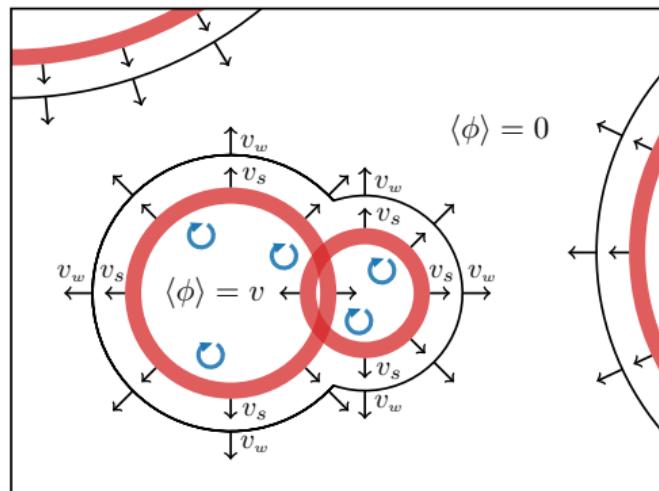
Cosmological Phase Transitions

- thermal corrections typically restore spontaneously broken symmetries at high temperatures
➡ symmetry breaking phase transition
- can be crossover or first-order



GW production:

1. vacuum bubble collisions
2. sound waves collisions
3. turbulence and vortical motion



Gravitational Wave Spectrum

- peak frequency:

$$f_p \propto T_* \frac{1}{R_* H_*} \sim 10 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{\beta/H_*}{10}$$

$$\alpha \approx \frac{\Delta V}{\rho_R}, \quad \beta = \frac{\dot{\Gamma}}{\Gamma} \approx \frac{3}{R_*}$$

Gravitational Wave Spectrum

- peak frequency:

$$f_p \propto T_* \frac{1}{R_* H_*} \sim 10 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{\beta/H_*}{10}$$

- peak amplitude: (for collisions)

$$h^2 \Omega_p \propto \left(\frac{\Delta V}{\Delta V + \rho_R} \right)^2 (R_* H_*)^2 \sim 10^{-8} \left(\frac{10}{\beta/H_*} \right)^2 \left(\frac{\alpha}{1+\alpha} \right)^2$$

$$\alpha \approx \frac{\Delta V}{\rho_R}, \quad \beta = \frac{\dot{\Gamma}}{\Gamma} \approx \frac{3}{R_*}$$

Gravitational Wave Spectrum

- peak frequency:

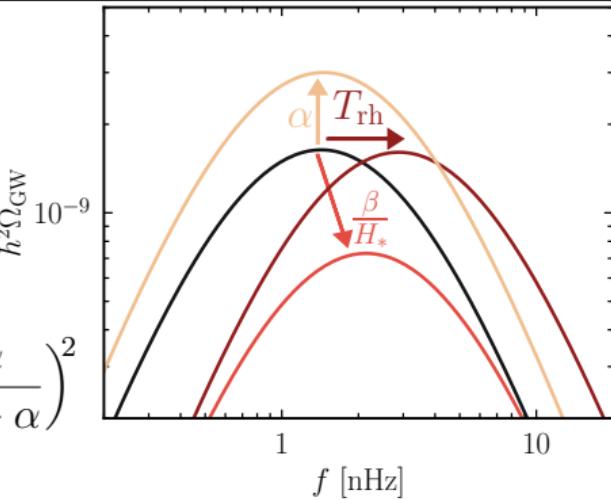
$$f_p \propto T_* \frac{1}{R_* H_*} \sim 10 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{\beta/H_*}{10}$$

- peak amplitude: (for collisions)

$$h^2 \Omega_p \propto \left(\frac{\Delta V}{\Delta V + \rho_R} \right)^2 (R_* H_*)^2 \sim 10^{-8} \left(\frac{10}{\beta/H_*} \right)^2 \left(\frac{\alpha}{1+\alpha} \right)^2$$

- spectral shape: [from Lewicki, Vaskonen (EPJC 2023)]

$$\Omega_{\text{GW}} = \Omega_p \frac{(a+b)^c}{\left[b(f/f_p)^{-a/c} + a(f/f_p)^{b/c} \right]^c}, \quad a = 2.41, \quad b = 2.42, \quad c = 4.08$$



$$\alpha \approx \frac{\Delta V}{\rho_R}, \quad \beta = \frac{\dot{\Gamma}}{\Gamma} \approx \frac{3}{R_*}$$

Gravitational Wave Spectrum

- peak frequency:

$$f_p \propto T_* \frac{1}{R_* H_*} \sim 10 \text{ nHz} \frac{T_*}{10 \text{ MeV}} \frac{\beta/H_*}{10}$$

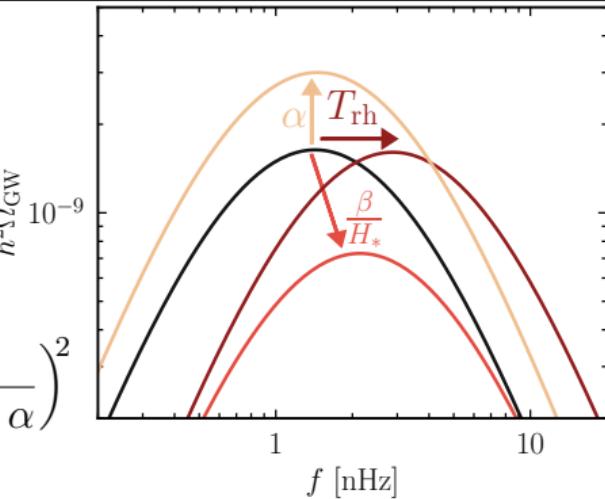
- peak amplitude: (for collisions)

$$h^2 \Omega_p \propto \left(\frac{\Delta V}{\Delta V + \rho_R} \right)^2 (R_* H_*)^2 \sim 10^{-8} \left(\frac{10}{\beta/H_*} \right)^2 \left(\frac{\alpha}{1+\alpha} \right)^2$$

- spectral shape: [from Lewicki, Vaskonen (EPJC 2023)]

$$\Omega_{\text{GW}} = \Omega_p \frac{(a+b)^c}{\left[b(f/f_p)^{-a/c} + a(f/f_p)^{b/c} \right]^c}, \quad a = 2.41, \quad b = 2.42, \quad c = 4.08$$

- for PTAs: $h^2 \Omega \sim 10^{-9} \implies \alpha \gtrsim 1$



$$\alpha \approx \frac{\Delta V}{\rho_R}, \quad \beta = \frac{\dot{\Gamma}}{\Gamma} \approx \frac{3}{R_*}$$

Gravitational Wave Spectrum

- peak frequency:

$$f_p \propto T_{\text{rh}} \frac{1}{R_* H_*} \sim 10 \text{ nHz} \frac{T_{\text{rh}}}{10 \text{ MeV}} \frac{\beta/H_*}{10}$$

- peak amplitude: (for collisions)

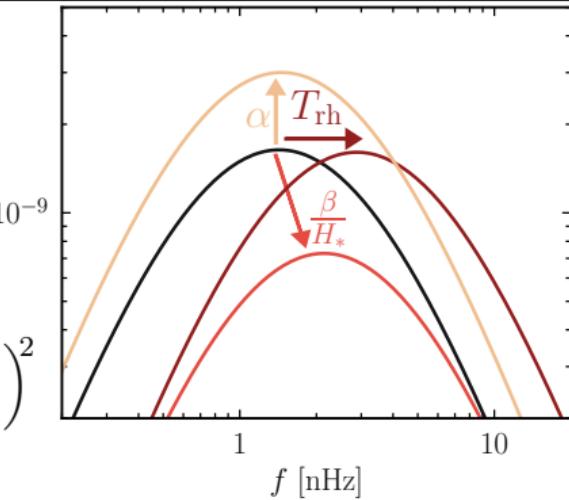
$$h^2 \Omega_p \propto \left(\frac{\Delta V}{\Delta V + \rho_R} \right)^2 (R_* H_*)^2 \sim 10^{-8} \left(\frac{10}{\beta/H_*} \right)^2 \left(\frac{\alpha}{1+\alpha} \right)^2$$

- spectral shape: [from Lewicki, Vaskonen (EPJC 2023)]

$$\Omega_{\text{GW}} = \Omega_p \frac{(a+b)^c}{\left[b(f/f_p)^{-a/c} + a(f/f_p)^{b/c} \right]^c}, \quad a = 2.41, \quad b = 2.42, \quad c = 4.08$$

- for PTAs: $h^2 \Omega \sim 10^{-9} \implies \alpha \gtrsim 1$

\implies supercooled PTs \implies nearly conformal models



$$\alpha \approx \frac{\Delta V}{\rho_R}, \quad \beta = \frac{\dot{\Gamma}}{\Gamma} \approx \frac{3}{R_*}$$

Coleman-Weinberg Model

- classically scale-invariant gauged U(1) model

[Coleman, Weinberg (PRD '73)]

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + D_\mu\phi^\dagger D^\mu\phi - \frac{\lambda}{4}\left(\phi^\dagger\phi\right)^2$$

- 3 parameters: g, λ, M

$$M = \frac{\mu_R}{4\pi} e^{\gamma_E - \frac{1}{3}}$$


Coleman-Weinberg Model

- classically scale-invariant gauged U(1) model

[Coleman, Weinberg (PRD '73)]

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + D_\mu\phi^\dagger D^\mu\phi - \frac{\lambda}{4}\left(\phi^\dagger\phi\right)^2$$

- 3 parameters: g , ~~M~~

$$M = \frac{\mu_R}{4\pi} e^{\gamma_E - \frac{1}{3}}$$


- for strong PhT: $\lambda \ll g^4$
(loop-induced SSB \rightarrow need flat direction)

Coleman-Weinberg Model

- classically scale-invariant gauged U(1) model

[Coleman, Weinberg (PRD '73)]

$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + D_\mu\phi^\dagger D^\mu\phi - \frac{\lambda}{4}\left(\phi^\dagger\phi\right)^2$$

- 3 parameters: g , ~~M~~

$$M = \frac{\mu_R}{4\pi} e^{\gamma_E - \frac{1}{3}}$$


- for strong PhT: $\lambda \ll g^4$
(loop-induced SSB \rightarrow need flat direction)

- $m_\phi \propto g^2 \langle\phi\rangle \propto gM \sim (1 - 100) \text{ MeV}$
 $m_A \propto g\langle\phi\rangle \propto M$

Coleman-Weinberg Model

- classically scale-invariant gauged U(1) model

[Coleman, Weinberg (PRD '73)]

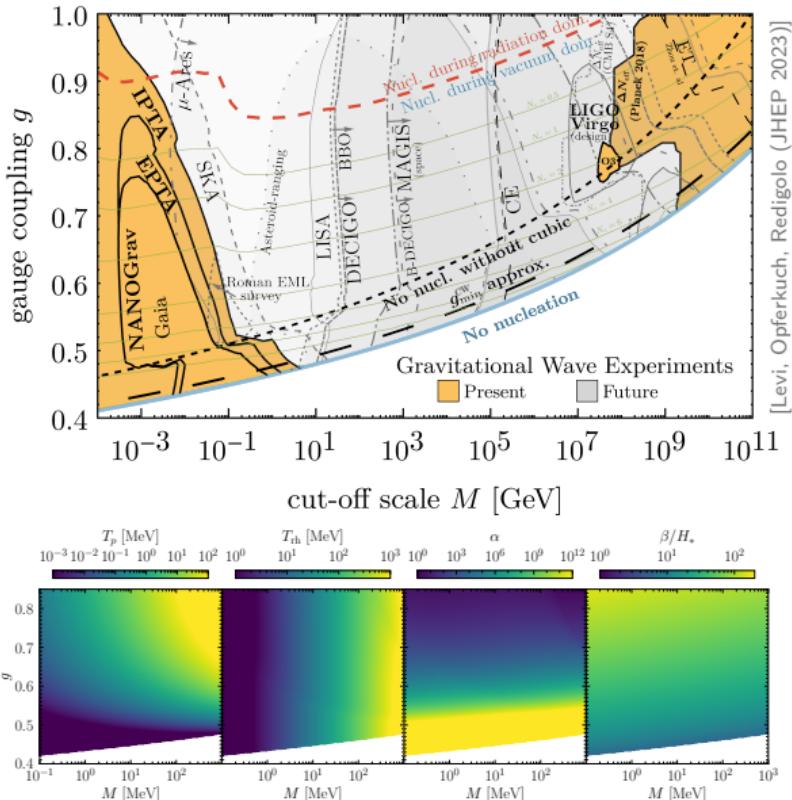
$$\mathcal{L} \supset -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + D_\mu\phi^\dagger D^\mu\phi - \frac{\lambda}{4}(\phi^\dagger\phi)^2$$

- 3 parameters: g , \cancel{M}

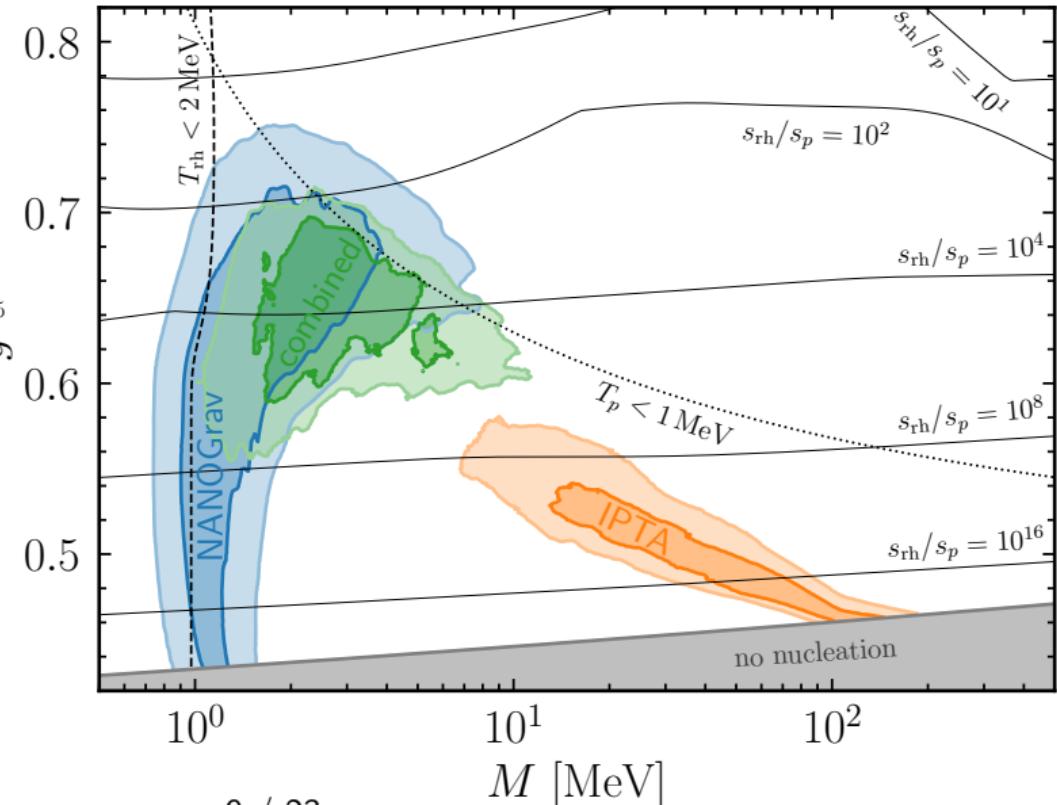
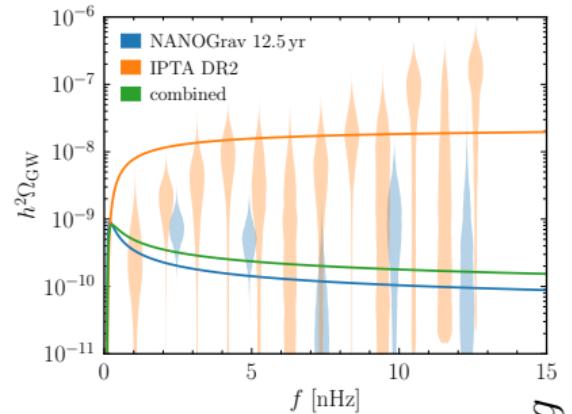
$$M = \frac{\mu_R}{4\pi} e^{\gamma_E - \frac{1}{3}}$$

- for strong PhT: $\lambda \ll g^4$
(loop-induced SSB \rightarrow need flat direction)

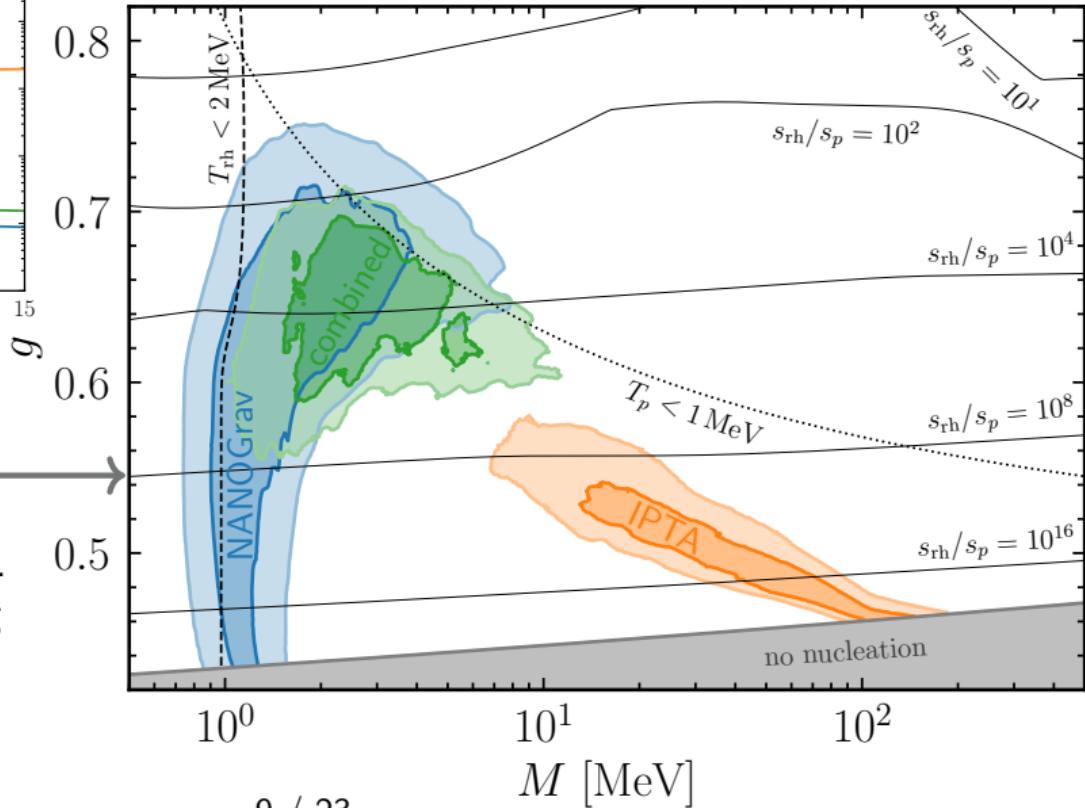
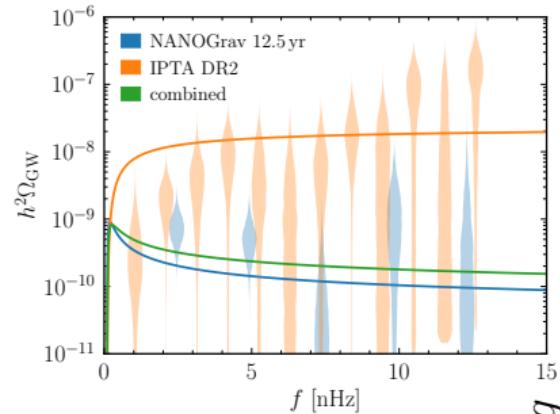
- $m_\phi \propto g^2 \langle\phi\rangle \propto gM \sim (1 - 100) \text{ MeV}$
 $m_A \propto g\langle\phi\rangle \propto M$



PTA fit - PhT in CW Model

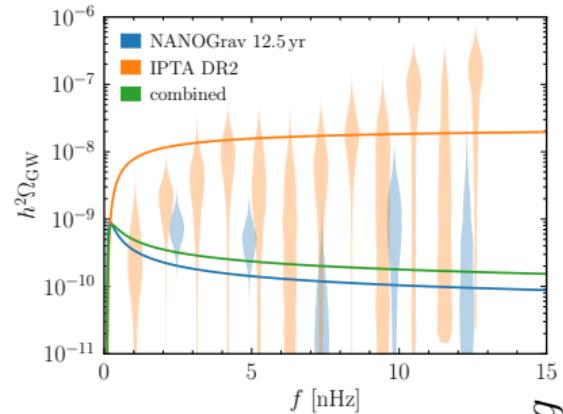


PTA fit - PhT in CW Model

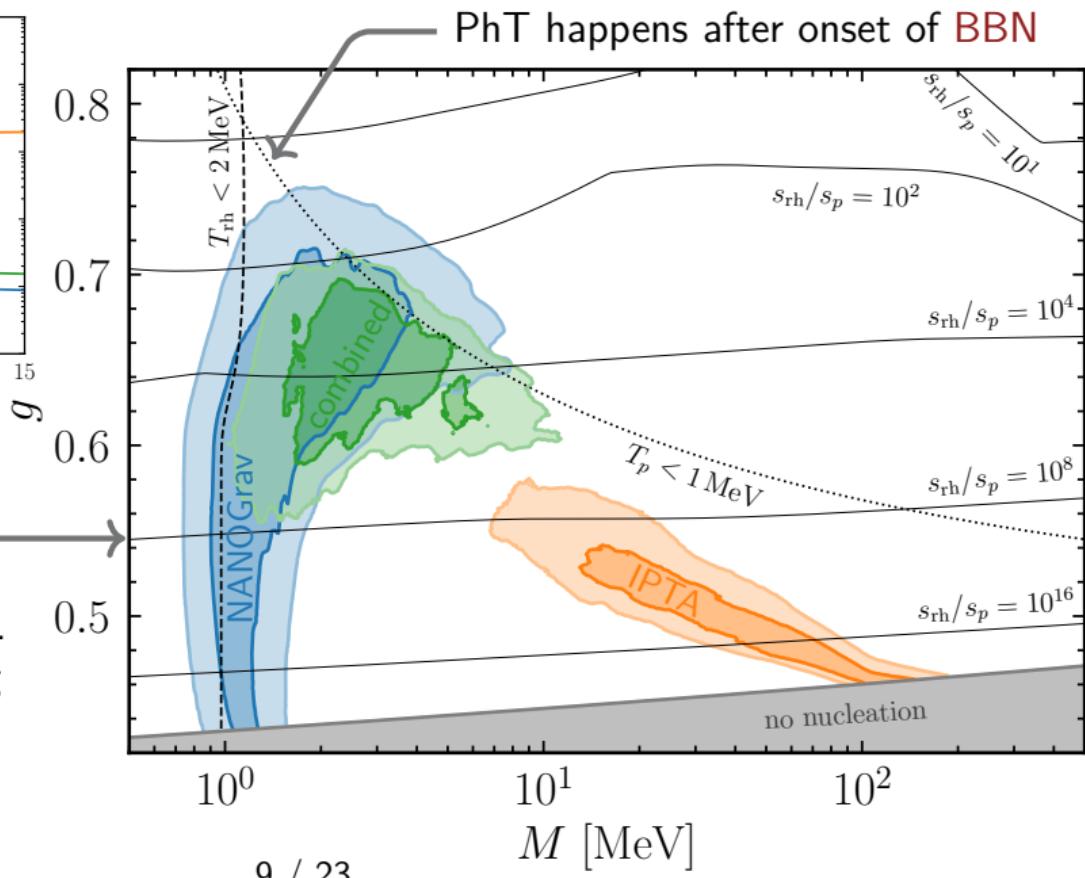


BAU: $\mathcal{O}(1)$ baryon asym.
required before reheating
(if $\frac{s_{\text{rh}}}{s_p} \sim 10^{10}$)

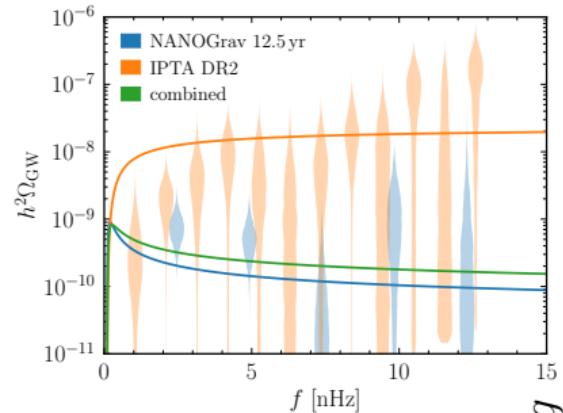
PTA fit - PhT in CW Model



BAU: $\mathcal{O}(1)$ baryon asym.
required before reheating
(if $\frac{s_{\text{rh}}}{s_p} \sim 10^{10}$)

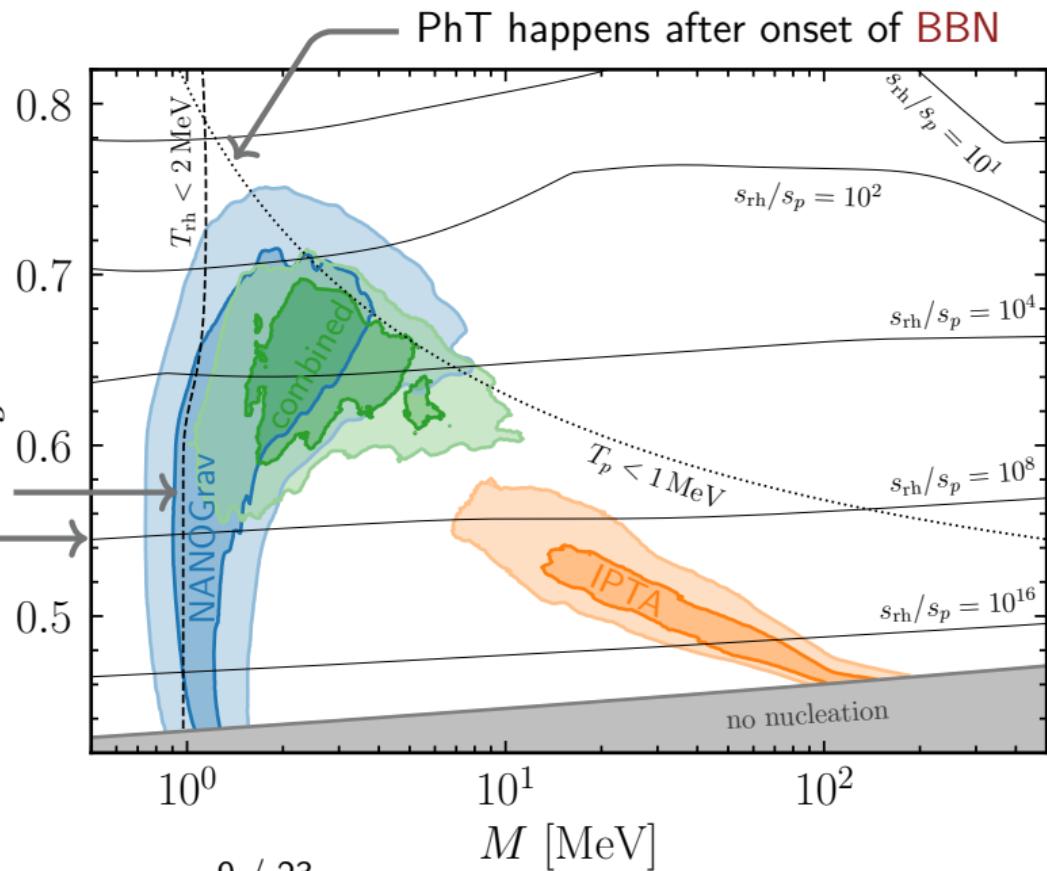


PTA fit - PhT in CW Model



ν -rethermalization

BAU: $\mathcal{O}(1)$ baryon asym.
required before reheating
(if $\frac{s_{\text{rh}}}{s_p} \sim 10^{10}$)



Reheating

- reheating: $\Delta V \longrightarrow \rho_R$ \Longrightarrow requires $\Gamma_\phi \gtrsim H_{\text{rh}} \sim \frac{T_{\text{rh}}^2}{M_{\text{pl}}} \gtrsim 4 \times 10^{-20} \text{ MeV}$

Reheating

- reheating: $\Delta V \longrightarrow \rho_R$ \implies requires $\Gamma_\phi \gtrsim H_{\text{rh}} \sim \frac{T_{\text{rh}}^2}{M_{\text{pl}}} \gtrsim 4 \times 10^{-20} \text{ MeV}$
- **Higgs portal:** $\mathcal{L} \supset \lambda_p |\phi|^2 |H|^2 \longrightarrow$ mixing
 - $\lambda_p \sim 10^{-4} \implies \Delta m_\phi^2 \sim \lambda_p v_H^2 \sim \text{GeV}^2$

Reheating

- reheating: $\Delta V \rightarrow \rho_R \implies$ requires $\Gamma_\phi \gtrsim H_{\text{rh}} \sim \frac{T_{\text{rh}}^2}{M_{\text{pl}}} \gtrsim 4 \times 10^{-20} \text{ MeV}$
- Higgs portal: $\mathcal{L} \supset \cancel{\lambda_p |\phi|^2 |H|^2} \rightarrow$ mixing
 - $\lambda_p \sim 10^{-4} \implies \Delta m_\phi^2 \sim \lambda_p v_H^2 \sim \text{GeV}^2$
 - violates scale invariance

Reheating

- reheating: $\Delta V \rightarrow \rho_R \implies$ requires $\Gamma_\phi \gtrsim H_{\text{rh}} \sim \frac{T_{\text{rh}}^2}{M_{\text{Pl}}} \gtrsim 4 \times 10^{-20} \text{ MeV}$
- Higgs portal: $\mathcal{L} \supset \lambda_p |\phi|^2 |H|^2 \rightarrow$ mixing
 - $\lambda_p \sim 10^{-4} \implies \Delta m_\phi^2 \sim \lambda_p v_H^2 \sim \text{GeV}^2$
 - violates scale invariance
- direct coupling to e, γ : $\mathcal{L} \supset c_e \frac{|\phi|^2}{\Lambda^2} L H \bar{e} + c_\gamma \frac{|\phi|^2}{\Lambda^2} F_{\mu\nu} F^{\mu\nu}$ ✓
 - does not violate scale invariance
 - $\phi \rightarrow e^+ e^-$, $\gamma\gamma$ searches e.g. at MESA

Domain Walls

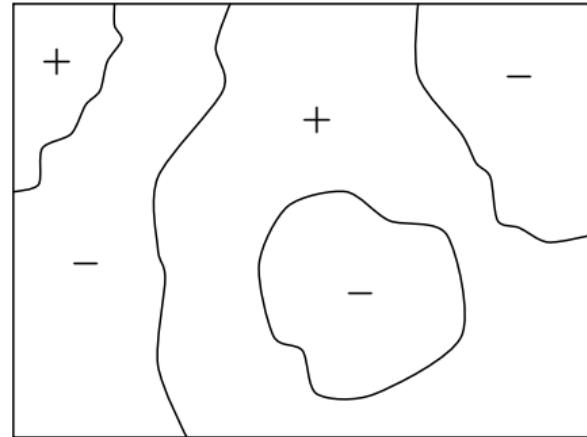
1. Pulsar Timing Arrays
2. Cosmological Phase Transitions
3. Domain Walls
4. Bosonic Instabilities
5. Conclusions

Domain Walls

- from spontaneously broken discrete symmetry (e.g. Z_2) after inflation
- DWs stretched to horizon size if friction negligible
⇒ scaling regime: ~ 1 DW/Hubble

$$\rho_{\text{DW}} \propto \frac{\sigma R^2}{R^3} \sim \sigma H$$

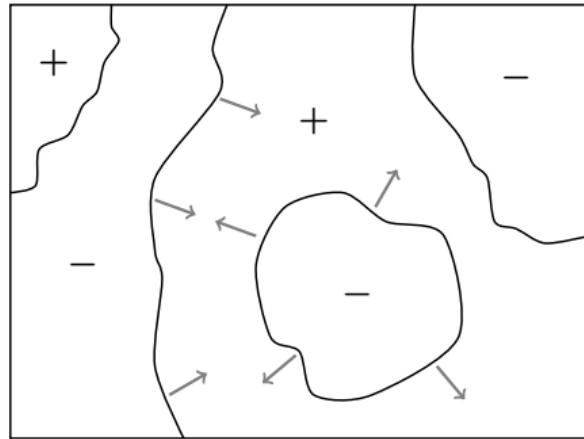
- DW domination: $\Omega_{\text{DW}} = \frac{\rho_{\text{DW}}}{3M_{\text{pl}}^2 H^2} \sim \frac{\sigma}{M_{\text{pl}} T^2}$ ⇒ $T_{\text{dom}} \sim \sqrt{\frac{\sigma}{M_{\text{pl}}}}$



Domain Walls

- from spontaneously broken discrete symmetry (e.g. Z_2) after inflation
- DWs stretched to horizon size if friction negligible
⇒ scaling regime: ~ 1 DW/Hubble

$$\rho_{\text{DW}} \propto \frac{\sigma R^2}{R^3} \sim \sigma H$$



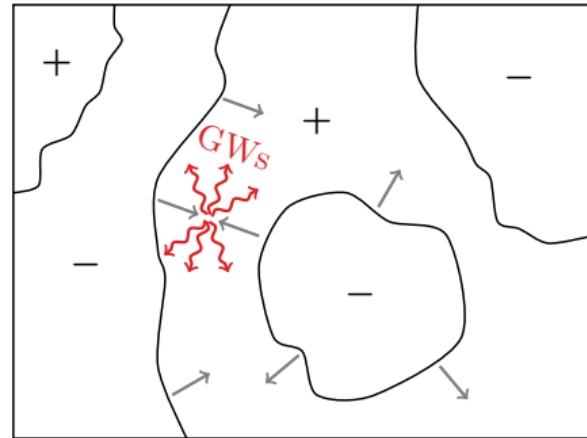
- DW domination: $\Omega_{\text{DW}} = \frac{\rho_{\text{DW}}}{3M_{\text{pl}}^2 H^2} \sim \frac{\sigma}{M_{\text{pl}} T^2}$ ⇒ $T_{\text{dom}} \sim \sqrt{\frac{\sigma}{M_{\text{pl}}}}$
- need to decay ⇒ add bias V_b
⇒ DWs annihilate when $V_b \sim \frac{\sigma}{R}$ ⇒ $T_{\text{ann}} \sim 20 \text{ MeV} \sqrt{\frac{\text{TeV}^3}{\sigma} \frac{V_b}{\text{MeV}^4}}$

Domain Walls

- from spontaneously broken discrete symmetry (e.g. Z_2) after inflation
- DWs stretched to horizon size if friction negligible
⇒ scaling regime: ~ 1 DW/Hubble

$$\rho_{\text{DW}} \propto \frac{\sigma R^2}{R^3} \sim \sigma H$$

- DW domination: $\Omega_{\text{DW}} = \frac{\rho_{\text{DW}}}{3M_{\text{pl}}^2 H^2} \sim \frac{\sigma}{M_{\text{pl}} T^2}$ ⇒ $T_{\text{dom}} \sim \sqrt{\frac{\sigma}{M_{\text{pl}}}}$
- need to decay ⇒ add bias V_b
⇒ DWs annihilate when $V_b \sim \frac{\sigma}{R}$ ⇒ $T_{\text{ann}} \sim 20 \text{ MeV} \sqrt{\frac{\text{TeV}^3}{\sigma} \frac{V_b}{\text{MeV}^4}}$
- GWs dominantly produced in DW collisions around T_{ann}



Gravitational Wave Spectrum

GW production at T_{ann}

- peak frequency:

$$f_p \sim \frac{a}{a_0} \frac{1}{R} \sim 1 \text{ nHz} \frac{T_{\text{ann}}}{10 \text{ MeV}}$$

Gravitational Wave Spectrum

GW production at T_{ann}

- peak frequency:

$$f_p \sim \frac{a}{a_0} \frac{1}{R} \sim 1 \text{ nHz} \frac{T_{\text{ann}}}{10 \text{ MeV}}$$

- peak amplitude:

$$h^2 \Omega_p \sim \left(\frac{a}{a_0} \right)^4 \frac{(\rho_{\text{DW}} R)^2}{\rho_c M_{\text{pl}}^2} \sim 10^{-18} \left(\frac{\sigma}{\text{TeV}^3} \right)^2 \left(\frac{T_{\text{ann}}}{10 \text{ MeV}} \right)^{-4}$$

Gravitational Wave Spectrum

GW production at T_{ann}

- peak frequency:

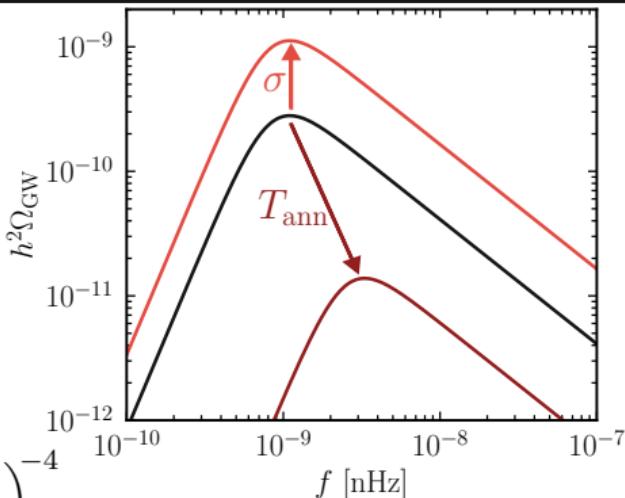
$$f_p \sim \frac{a}{a_0} \frac{1}{R} \sim 1 \text{ nHz} \frac{T_{\text{ann}}}{10 \text{ MeV}}$$

- peak amplitude:

$$h^2 \Omega_p \sim \left(\frac{a}{a_0} \right)^4 \frac{(\rho_{\text{DW}} R)^2}{\rho_c M_{\text{pl}}^2} \sim 10^{-18} \left(\frac{\sigma}{\text{TeV}^3} \right)^2 \left(\frac{T_{\text{ann}}}{10 \text{ MeV}} \right)^{-4}$$

- spectral shape:

$$\Omega_{\text{GW}} = \Omega_p \frac{4}{(f/f_p)^{-3} + 3 (f/f_p)}$$



[Ferreira et al. (JCAP 2023)]
[Hitamatsu et al. (JCAP 2014)]

Gravitational Wave Spectrum

GW production at T_{ann}

- peak frequency:

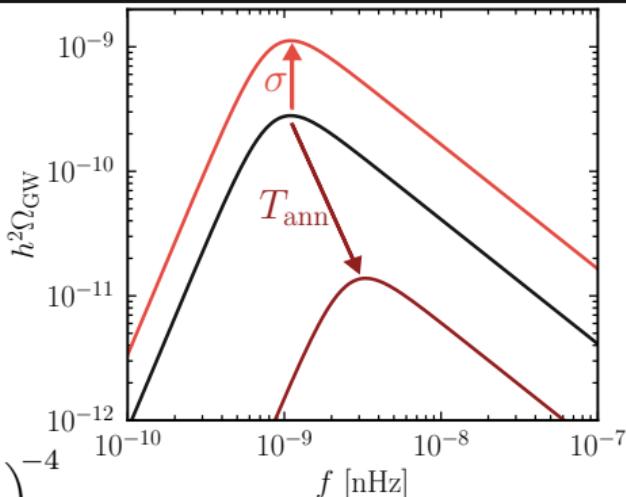
$$f_p \sim \frac{a}{a_0} \frac{1}{R} \sim 1 \text{ nHz} \frac{T_{\text{ann}}}{10 \text{ MeV}}$$

- peak amplitude:

$$h^2 \Omega_p \sim \left(\frac{a}{a_0} \right)^4 \frac{(\rho_{\text{DW}} R)^2}{\rho_c M_{\text{pl}}^2} \sim 10^{-18} \left(\frac{\sigma}{\text{TeV}^3} \right)^2 \left(\frac{T_{\text{ann}}}{10 \text{ MeV}} \right)^{-4}$$

- spectral shape: $\Omega_{\text{GW}} = \Omega_p \frac{4}{(f/f_p)^{-3} + 3 (f/f_p)}$

- for PTAs: $T_{\text{ann}} \sim \text{few MeV}, h^2 \Omega_{\text{DW}} \sim 0.1$



[Ferreira et al. (JCAP 2023)
Hitamatsu et al. (JCAP 2014)]

ALP Domain Walls

$$\mathcal{L} \supset \partial_\mu \phi^\dagger \partial^\mu \phi - \underbrace{\lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2}_{\begin{array}{l} \text{breaks U(1) spontaneously} \\ \text{w/ } \phi = \frac{v+\rho}{\sqrt{2}} e^{i \frac{a}{v}} \end{array}} - \underbrace{m_a^2 f_a^2 \left[1 - \cos \frac{a}{f_a} \right]}_{\begin{array}{l} \text{explicit breaking} \\ \text{U(1)} \rightarrow Z_N \\ \text{w/ } f_a = v/N \\ \Rightarrow \text{DW stable if } N \geq 2 \end{array}}$$

ALP Domain Walls

$$\mathcal{L} \supset \partial_\mu \phi^\dagger \partial^\mu \phi - \underbrace{\lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2}_{\begin{array}{l} \text{breaks U(1) spontaneously} \\ \text{w/ } \phi = \frac{v+\rho}{\sqrt{2}} e^{i \frac{a}{v}} \end{array}} - \underbrace{m_a^2 f_a^2 \left[1 - \cos \frac{a}{f_a} \right]}_{\begin{array}{l} \text{explicit breaking} \\ \text{U(1)} \rightarrow Z_N \\ \text{w/ } f_a = v/N \\ \Rightarrow \text{DW stable if } N \geq 2 \end{array}}$$

- **wall tension:** $\sigma = 8 \frac{m_a^2 f_a^2}{m_a}$ } size of potential barriers
} wall thickness $\sim \lambda_{\text{Coulomb}} \sim m_a^{-1}$

ALP Domain Walls

$$\mathcal{L} \supset \partial_\mu \phi^\dagger \partial^\mu \phi - \underbrace{\lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2}_{\begin{array}{l} \text{breaks U(1) spontaneously} \\ \text{w/ } \phi = \frac{v+\rho}{\sqrt{2}} e^{i \frac{a}{v}} \end{array}} - \underbrace{m_a^2 f_a^2 \left[1 - \cos \frac{a}{f_a} \right]}_{\begin{array}{l} \text{explicit breaking} \\ \text{U(1)} \rightarrow Z_N \\ \text{w/ } f_a = v/N \\ \Rightarrow \text{DW stable if } N \geq 2 \end{array}}$$

- **wall tension:** $\sigma = 8 \frac{m_a^2 f_a^2}{m_a}$ } size of potential barriers
} wall thickness $\sim \lambda_{\text{Coulomb}} \sim m_a^{-1}$
- additional breaking of $Z_N \implies V_b$

ALP Domain Walls

$$\mathcal{L} \supset \partial_\mu \phi^\dagger \partial^\mu \phi - \underbrace{\lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2}_{\begin{array}{l} \text{breaks U(1) spontaneously} \\ \text{w/ } \phi = \frac{v+\rho}{\sqrt{2}} e^{i \frac{a}{v}} \end{array}} - \underbrace{m_a^2 f_a^2 \left[1 - \cos \frac{a}{f_a} \right]}_{\begin{array}{l} \text{explicit breaking} \\ \text{U(1)} \rightarrow Z_N \\ \text{w/ } f_a = v/N \\ \Rightarrow \text{DW stable if } N \geq 2 \end{array}}$$

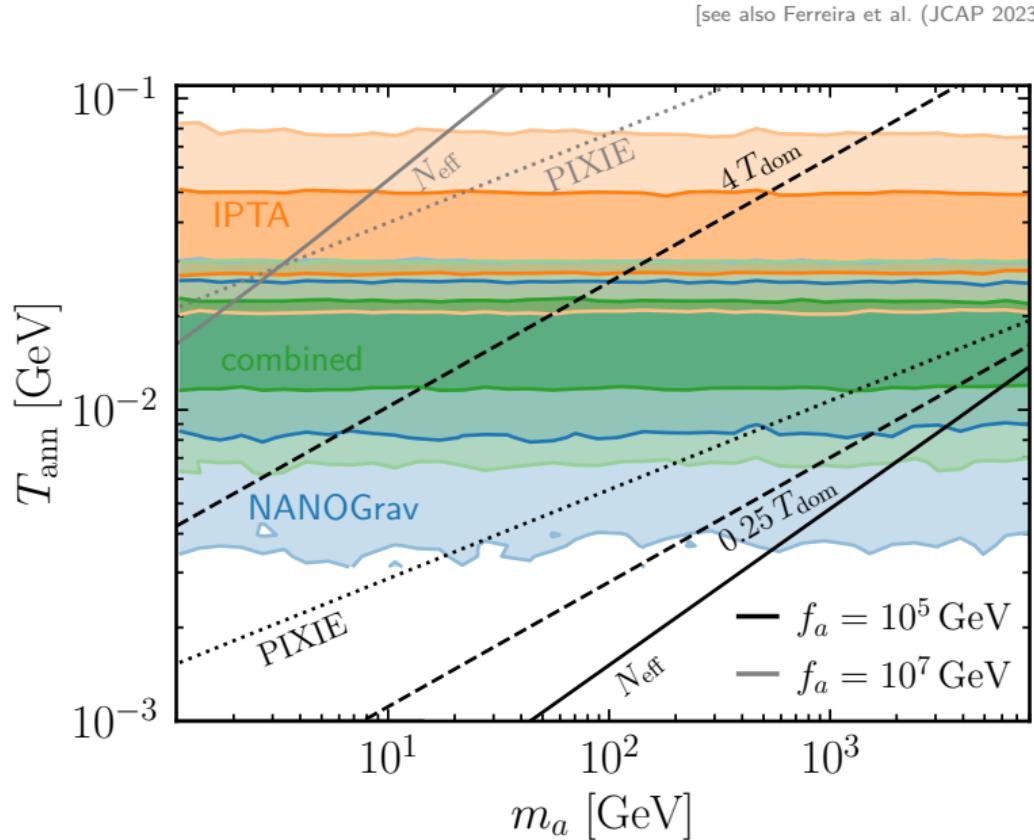
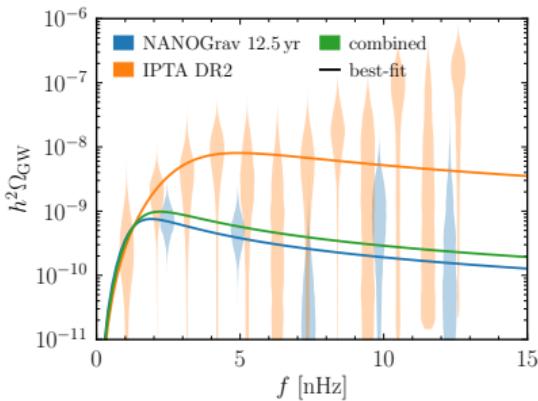
- wall tension: $\sigma = 8 \frac{m_a^2 f_a^2}{m_a}$ } size of potential barriers
} wall thickness $\sim \lambda_{\text{Coulomb}} \sim m_a^{-1}$

- additional breaking of $Z_N \implies V_b$

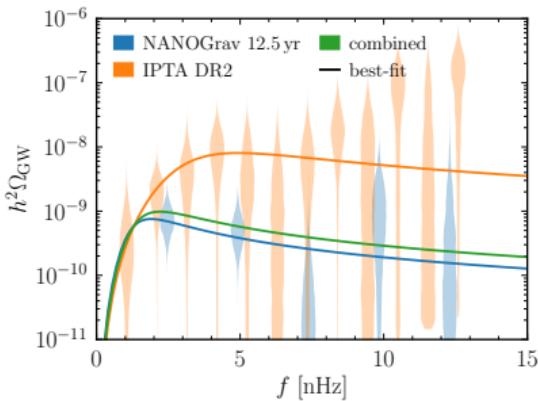
- DWs annihilate \longrightarrow ALPs \longrightarrow dark radiation

$$\Delta N_{\text{eff}} \sim \frac{\rho_{\text{DW}}}{T_{\text{ann}}^4} \sim \frac{m_a f_a^2}{M_{\text{Pl}} T_{\text{ann}}^2}$$

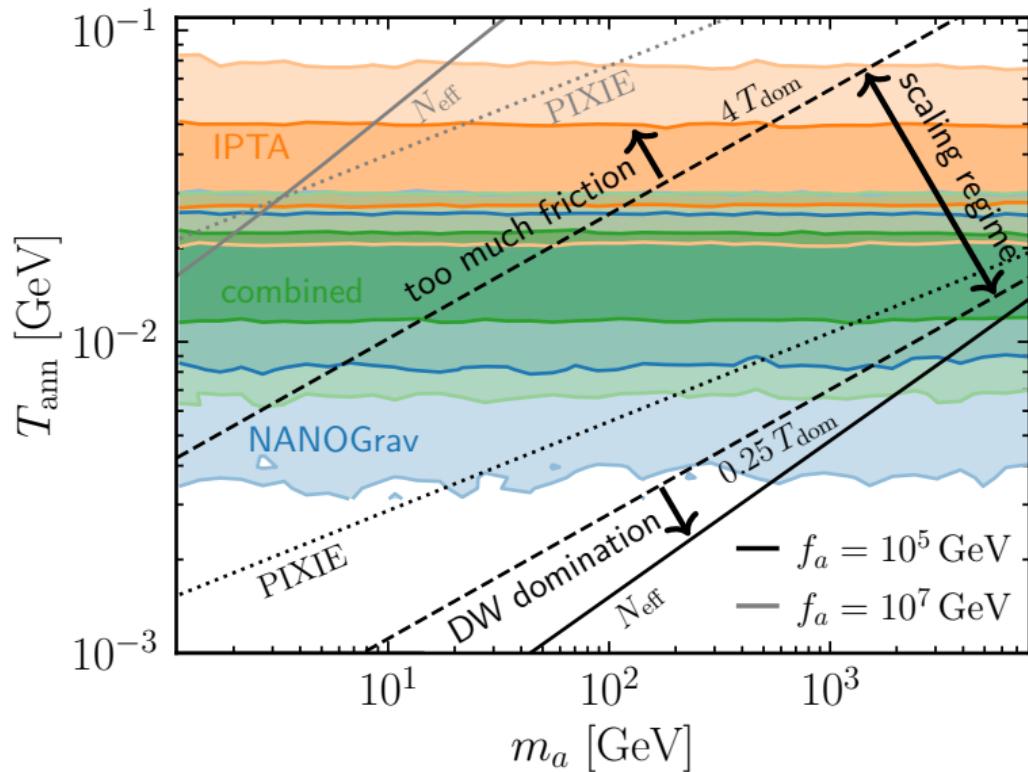
PTA Fit – ALP DWs



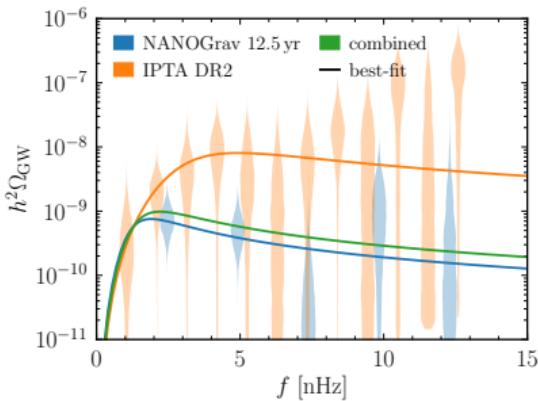
PTA Fit – ALP DWs



[see also Ferreira et al. (JCAP 2023)]



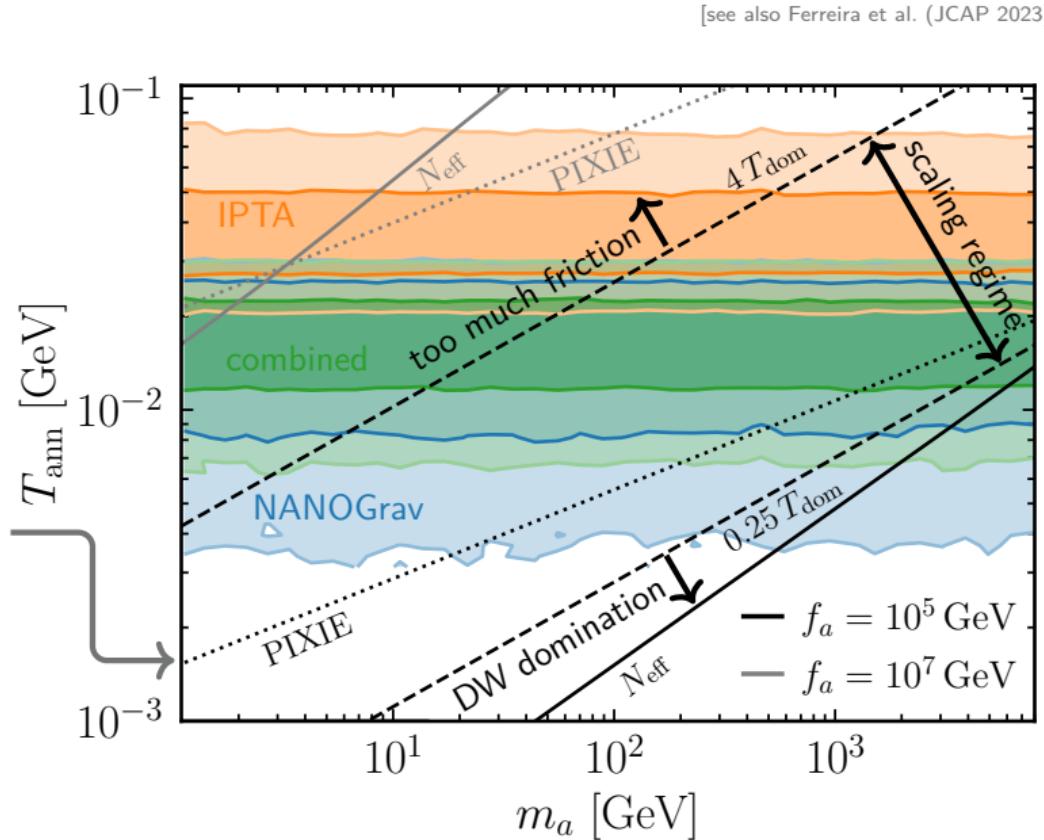
PTA Fit – ALP DWs



μ -distortions in CMB

[Ramberg, Ratzinger, Schwaller (JCAP 2023)]

from acoustic waves in baryon-photon fluid induced via gravity



Aligned Axion

[Higaki et al. (JHEP 2016)]

○ N axions:

$$V = \underbrace{\sum_{i=1}^N \lambda_i \left(\Phi_i^\dagger \Phi_i - \frac{f_i}{2} \right)^2}_{\text{breaks } U(1)^N \text{ spontaneously}} + \underbrace{\left[\sum_{i=1}^{N-1} \epsilon_i \Phi_i \Phi_{i+1}^3 + \text{h.c.} \right]}_{U(1)^{N-1} \text{ explicitly broken to discrete sym.}}$$

Aligned Axion

[Higaki et al. (JHEP 2016)]

- N axions:
$$V = \underbrace{\sum_{i=1}^N \lambda_i \left(\Phi_i^\dagger \Phi_i - \frac{f_i}{2} \right)^2}_{\text{breaks } U(1)^N \text{ spontaneously}} + \underbrace{\left[\sum_{i=1}^{N-1} \epsilon_i \Phi_i \Phi_{i+1}^3 + \text{h.c.} \right]}_{U(1)^{N-1} \text{ explicitly broken to discrete sym.}}$$

$\implies (N-1)$ heavy axions + massless QCD axion

\hookrightarrow DWs \hookrightarrow anomalous $\implies V_b \sim \Lambda_{\text{QCD}}^4$
- effective decay constant for QCD axion: $f_{\text{eff}} \sim \sqrt{N} 3^N f \gg f$

Aligned Axion

[Higaki et al. (JHEP 2016)]

- N axions:
$$V = \underbrace{\sum_{i=1}^N \lambda_i \left(\Phi_i^\dagger \Phi_i - \frac{f_i}{2} \right)^2}_{\text{breaks } U(1)^N \text{ spontaneously}} + \underbrace{\left[\sum_{i=1}^{N-1} \epsilon_i \Phi_i \Phi_{i+1}^3 + \text{h.c.} \right]}_{U(1)^{N-1} \text{ explicitly broken to discrete sym.}}$$

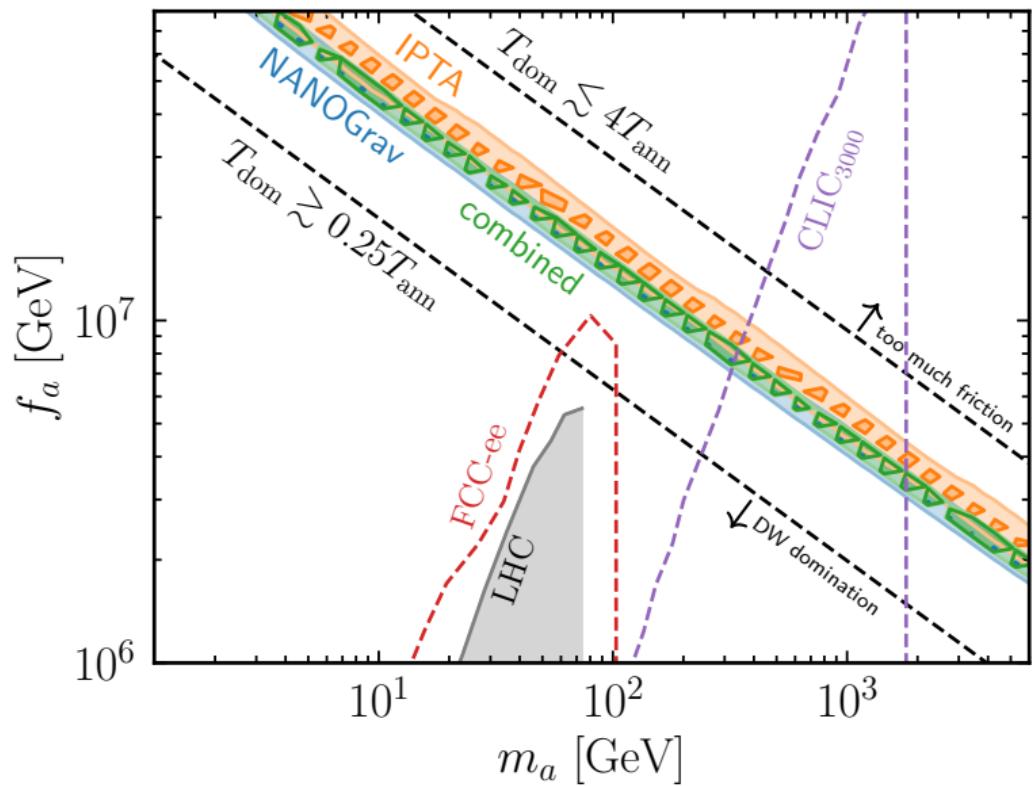
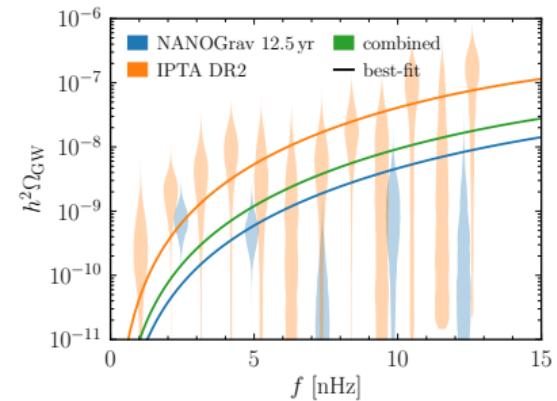
$\implies (N-1)$ heavy axions + massless QCD axion
 \downarrow DWs \downarrow anomalous $\implies V_b \sim \Lambda_{\text{QCD}}^4$
- effective decay constant for QCD axion: $f_{\text{eff}} \sim \sqrt{N} 3^N f \gg f$
- annihilation: $\sigma \sim m_a f_a^2, V_b \sim \Lambda_{\text{QCD}}^4$ $\implies T_{\text{ann}} \sim 1 \text{ GeV} \left(\frac{10^7 \text{ GeV}}{f_a} \right)^4 \sqrt{\frac{10 \text{ GeV}}{m_a}}$

Aligned Axion

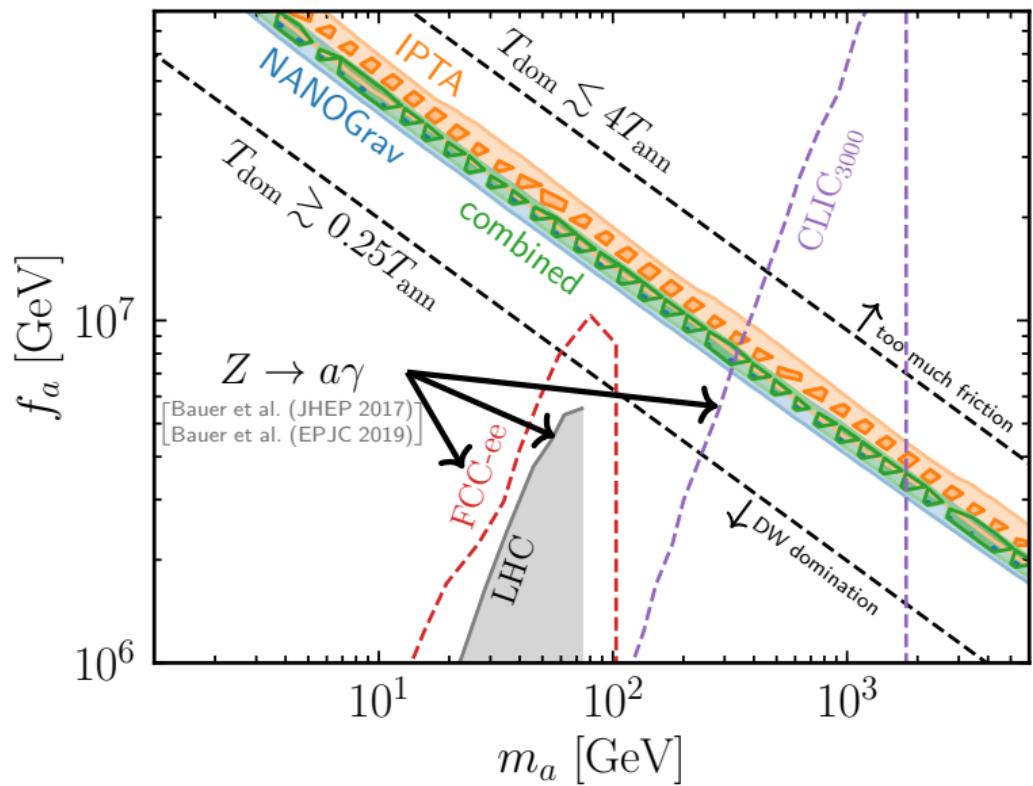
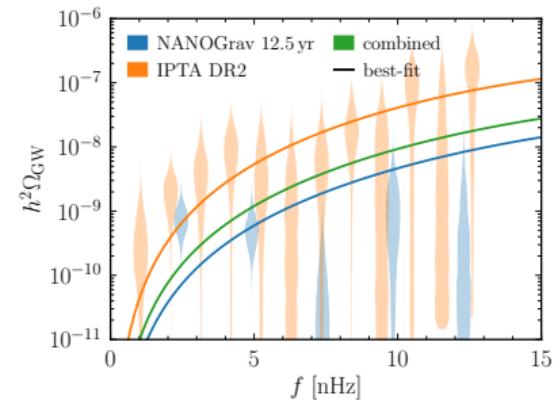
[Higaki et al. (JHEP 2016)]

- N axions:
$$V = \underbrace{\sum_{i=1}^N \lambda_i \left(\Phi_i^\dagger \Phi_i - \frac{f_i}{2} \right)^2}_{\text{breaks } U(1)^N \text{ spontaneously}} + \underbrace{\left[\sum_{i=1}^{N-1} \epsilon_i \Phi_i \Phi_{i+1}^3 + \text{h.c.} \right]}_{U(1)^{N-1} \text{ explicitly broken to discrete sym.}}$$
- $\Rightarrow (N-1)$ heavy axions + massless QCD axion
 $\quad \quad \quad \downarrow \text{DWs} \quad \quad \quad \downarrow \text{anomalous} \Rightarrow V_b \sim \Lambda_{\text{QCD}}^4$
- effective decay constant for QCD axion: $f_{\text{eff}} \sim \sqrt{N} 3^N f \gg f$
- annihilation: $\sigma \sim m_a f_a^2, V_b \sim \Lambda_{\text{QCD}}^4 \quad \Rightarrow \quad T_{\text{ann}} \sim 1 \text{ GeV} \left(\frac{10^7 \text{ GeV}}{f_a} \right)^4 \sqrt{\frac{10 \text{ GeV}}{m_a}}$
- heavy axion decay to SM: $\Gamma_{a \rightarrow gg} \sim 10^{11} \text{ s}^{-1} \left(\frac{m_a}{10 \text{ GeV}} \right)^3 \left(\frac{10^7 \text{ GeV}}{f_a} \right)^2$
 \Rightarrow decay before BBN, no N_{eff} constraints

PTA Fit – Aligned Axion DWs



PTA Fit – Aligned Axion DWs



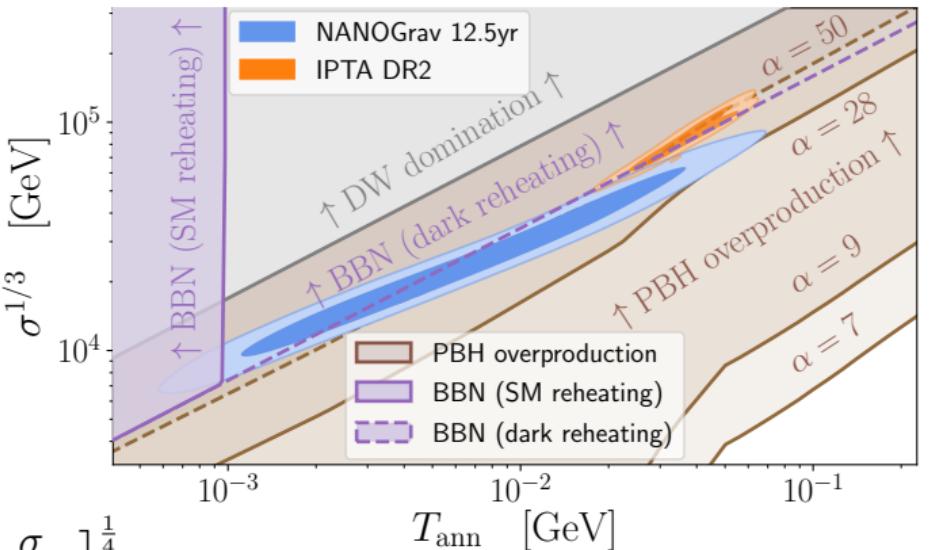
Primordial Black Hole Production

DWs can collapse to PBHs if

$$R < r_S = 2GM$$

$$\sim 2G \left(\frac{4\pi^2}{3} R^3 V_b + 4\pi R^2 \sigma \right)$$

$$\implies T_{\text{PBH}} \sim 100 \text{ MeV} \sqrt{\frac{T_{\text{ann}}}{10 \text{ MeV}}} \left[\frac{\sigma}{\text{TeV}^2} \right]^{\frac{1}{4}}$$



[Gouttenoire, Vitagliano, arXiv:2306.17841 [gr-qc]]

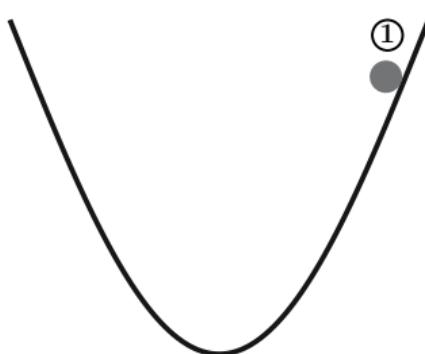
Bosonic Instabilities

1. Pulsar Timing Arrays
2. Cosmological Phase Transitions
3. Domain Walls
4. **Bosonic Instabilities**
5. Conclusions

Misalignment Mechanism

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi)$$

$$\implies \ddot{\phi} + 3H\dot{\phi} + m_a^2\phi = 0$$

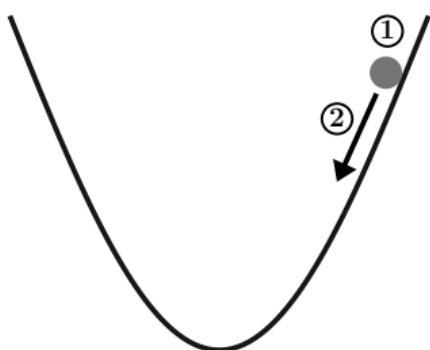


1. $H \gg m_a$: axion pinned by Hubble friction.

Misalignment Mechanism

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi)$$

$$\implies \ddot{\phi} + 3H\dot{\phi} + m_a^2\phi = 0$$

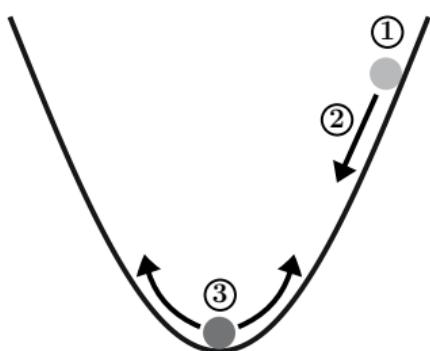


1. $H \gg m_a$: axion pinned by Hubble friction.
2. $H \sim m_a$: axion starts to roll

Misalignment Mechanism

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi)$$

$$\implies \ddot{\phi} + 3H\dot{\phi} + m_a^2\phi = 0$$



1. $H \gg m_a$: axion pinned by Hubble friction.
2. $H \sim m_a$: axion starts to roll
3. $H \ll m_a$: axion oscillates

Bosonic Instabilities: Audible Axions

Machado, Ratzinger, Schwaller, Stefanek (JHEP 2019, PRD 2020)
 Ratzinger, Schwaller, Stefanek (SciPost Phys. 2021)

Misalignment Mechanism

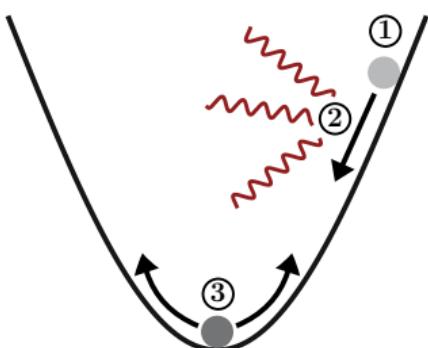
+

coupling to dark photon

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \frac{\alpha}{4 f_a} X_{\mu\nu} \tilde{X}^{\mu\nu}$$

$$\implies \ddot{\phi} + 3H\dot{\phi} + m_a^2 \phi = -\frac{\alpha}{f_a} \langle X_{\mu\nu} \tilde{X}^{\mu\nu} \rangle$$

e.g. to deplete axion abundance
 Agrawal, Marques-Tavares, Xue (JHEP, 2018)
 Kitajima, Sekiguchi, Takahashi (PLB 2018)
 or for dark-photon dark-matter
 Dror, Harigaya, Narayan (PRD 2019);
 Co, Pierce, Zhang, Zhao (PRD 2019);
 Bastero-Gil, Santiago, Ubaldi, Vega-Morales
 (JCAP 2019); Agrawal, Kitajima, Reece,
 Sekiguchi, Takahashi (PLB 2019)



1. $H \gg m_a$: axion pinned by Hubble friction.
 2. $H \sim m_a$: axion starts to roll
 3. $H \ll m_a$: axion oscillates
- \implies dark photon production during phase 2 (and 3)

Bosonic Instabilities: Audible Axions

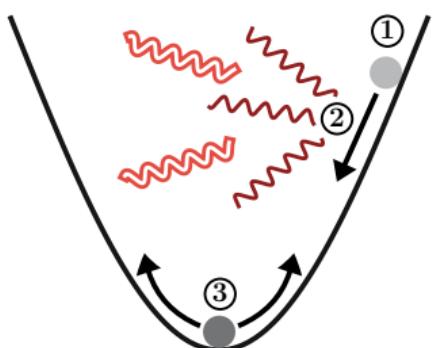
Machado, Ratzinger, Schwaller, Stefanek (JHEP 2019, PRD 2020)
 Ratzinger, Schwaller, Stefanek (SciPost Phys. 2021)

Misalignment Mechanism + coupling to dark photon

$$\mathcal{L} \supset \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) - \frac{1}{4} X_{\mu\nu} X^{\mu\nu} - \frac{\alpha}{4 f_a} X_{\mu\nu} \tilde{X}^{\mu\nu}$$

$$\implies \ddot{\phi} + 3H\dot{\phi} + m_a^2 \phi = -\frac{\alpha}{f_a} \langle X_{\mu\nu} \tilde{X}^{\mu\nu} \rangle$$

e.g. to deplete axion abundance
 Agrawal, Marques-Tavares, Xue (JHEP, 2018)
 Kitajima, Sekiguchi, Takahashi (PLB 2018)
 or for dark-photon dark-matter
 Dror, Harigaya, Narayan (PRD 2019);
 Co, Pierce, Zhang, Zhao (PRD 2019);
 Bastero-Gil, Santiago, Ubaldi, Vega-Morales
 (JCAP 2019); Agrawal, Kitajima, Reece,
 Sekiguchi, Takahashi (PLB 2019)



1. $H \gg m_a$: axion pinned by Hubble friction.
 2. $H \sim m_a$: axion starts to roll
 3. $H \ll m_a$: axion oscillates
- \implies dark photon production during phase 2 (and 3)
- \implies gravitational wave emission from dark photons

Dark photon production

dark photon EoM: $X''_{\pm}(\tau, k) + \omega_{\pm}^2(k)X_{\pm}(\tau, k) = 0$ $\omega_{\pm}^2(k) = \left(k^2 \mp k \frac{\alpha\phi'(\tau)}{f_a} \right)$

- modes with $k < \left| \frac{\alpha\phi'(\tau)}{f_a} \right|$ experience tachyonic instability in one helicity

Dark photon production

$$\text{dark photon EoM: } X''_{\pm}(\tau, k) + \omega_{\pm}^2(k)X_{\pm}(\tau, k) = 0 \quad \omega_{\pm}^2(k) = \left(k^2 \mp k \frac{\alpha\phi'(\tau)}{f_a} \right)$$

- modes with $k < \left| \frac{\alpha\phi'(\tau)}{f_a} \right|$ experience tachyonic instability in one helicity
- largest growth for $\tilde{k} = \left| \frac{\alpha\phi'(\tau)}{2f_a} \right|$ with $|\omega^2(\tilde{k})| = \tilde{k}^2$

Dark photon production

$$\text{dark photon EoM: } X''_{\pm}(\tau, k) + \omega_{\pm}^2(k)X_{\pm}(\tau, k) = 0 \quad \omega_{\pm}^2(k) = \left(k^2 \mp k \frac{\alpha\phi'(\tau)}{f_a} \right)$$

- modes with $k < \left| \frac{\alpha\phi'(\tau)}{f_a} \right|$ experience tachyonic instability in one helicity
- largest growth for $\tilde{k} = \left| \frac{\alpha\phi'(\tau)}{2f_a} \right|$ with $|\omega^2(\tilde{k})| = \tilde{k}^2$
- for oscillating axion: $\tilde{k} \sim a^{-1/2}$

Dark photon production

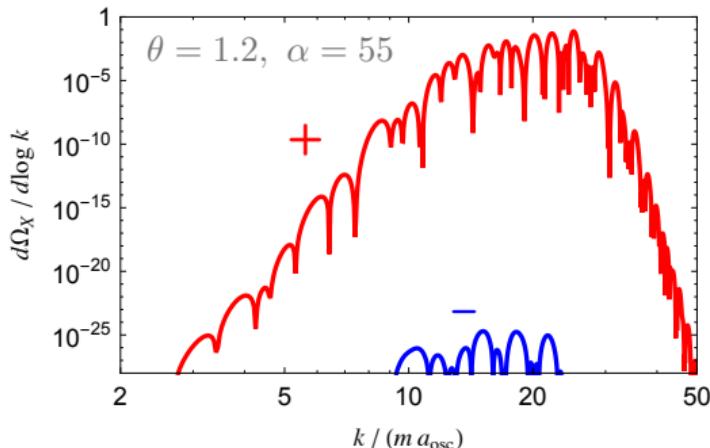
$$\text{dark photon EoM: } X''_{\pm}(\tau, k) + \omega_{\pm}^2(k)X_{\pm}(\tau, k) = 0 \quad \omega_{\pm}^2(k) = \left(k^2 \mp k \frac{\alpha\phi'(\tau)}{f_a} \right)$$

- modes with $k < \left| \frac{\alpha\phi'(\tau)}{f_a} \right|$ experience tachyonic instability in one helicity
- largest growth for $\tilde{k} = \left| \frac{\alpha\phi'(\tau)}{2f_a} \right|$ with $|\omega^2(\tilde{k})| = \tilde{k}^2$
- for oscillating axion: $\tilde{k} \sim a^{-1/2}$
- tachyonic band closes when $|\omega(\tilde{k})| < a m_a$
 $\implies \sim \text{energy transfer stops}$

Dark photon production

dark photon EoM: $X''_{\pm}(\tau, k) + \omega_{\pm}^2(k)X_{\pm}(\tau, k) = 0$ $\omega_{\pm}^2(k) = \left(k^2 \mp k \frac{\alpha\phi'(\tau)}{f_a} \right)$

- modes with $k < \left| \frac{\alpha\phi'(\tau)}{f_a} \right|$ experience tachyonic instability in one helicity
- largest growth for $\tilde{k} = \left| \frac{\alpha\phi'(\tau)}{2f_a} \right|$ with $|\omega^2(\tilde{k})| = \tilde{k}^2$
- for oscillating axion: $\tilde{k} \sim a^{-1/2}$
- tachyonic band closes when $|\omega(\tilde{k})| < a m_a$
 $\implies \sim$ energy transfer stops



[Machado, Ratzinger, Schwaller, Stefanek (JHEP 2019)]

dark photon spectrum:

- peaked around $\tilde{k} \approx a_{\text{osc}} m_a (\alpha\theta/2)^{2/3}$
- first tachyonic helicity dominates

Gravitational Wave Spectrum

GWs generated at t_* around the time when the tachyonic band closes:

- peak frequency:

$$f_{\text{peak}} \sim 2 \frac{\tilde{k}_*}{a_0} \sim 4 \text{ nHz} \left(\frac{\alpha \theta}{100} \right)^{\frac{2}{3}} \left(\frac{m_a}{10^{-15} \text{ eV}} \right)^{\frac{1}{2}}$$

Gravitational Wave Spectrum

GWs generated at t_* around the time when the tachyonic band closes:

- peak frequency:

$$f_{\text{peak}} \sim 2 \frac{\tilde{k}_*}{a_0} \sim 4 \text{ nHz} \left(\frac{\alpha \theta}{100} \right)^{\frac{2}{3}} \left(\frac{m_a}{10^{-15} \text{ eV}} \right)^{\frac{1}{2}}$$

- peak amplitude:

$$\Omega_{\text{GW}}^{\text{peak}} \sim \frac{\left(\rho_X^*/f_{\text{peak}}^* \right)^2}{\rho_c M_{\text{pl}}^2} \left(\frac{a_*}{a_0} \right)^4 \sim 10^{-7} \left(\frac{f_a}{M_{\text{pl}}} \right)^4 \left(100 \frac{\theta^2}{\alpha} \right)^{\frac{4}{3}}$$

Gravitational Wave Spectrum

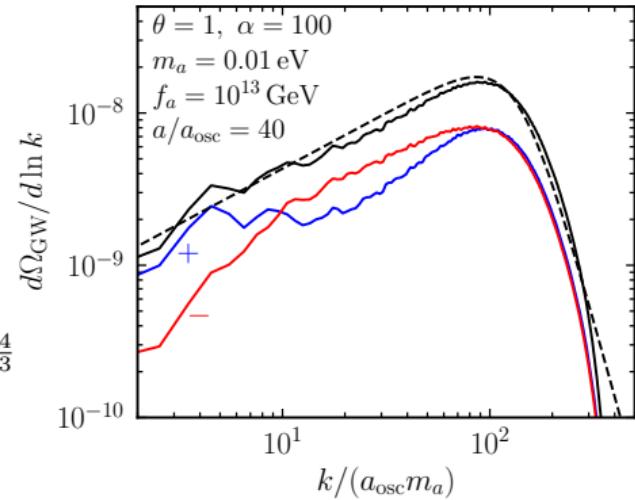
GWs generated at t_* around the time when the tachyonic band closes:

- peak frequency:

$$f_{\text{peak}} \sim 2 \frac{\tilde{k}_*}{a_0} \sim 4 \text{ nHz} \left(\frac{\alpha \theta}{100} \right)^{\frac{2}{3}} \left(\frac{m_a}{10^{-15} \text{ eV}} \right)^{\frac{1}{2}}$$

- peak amplitude:

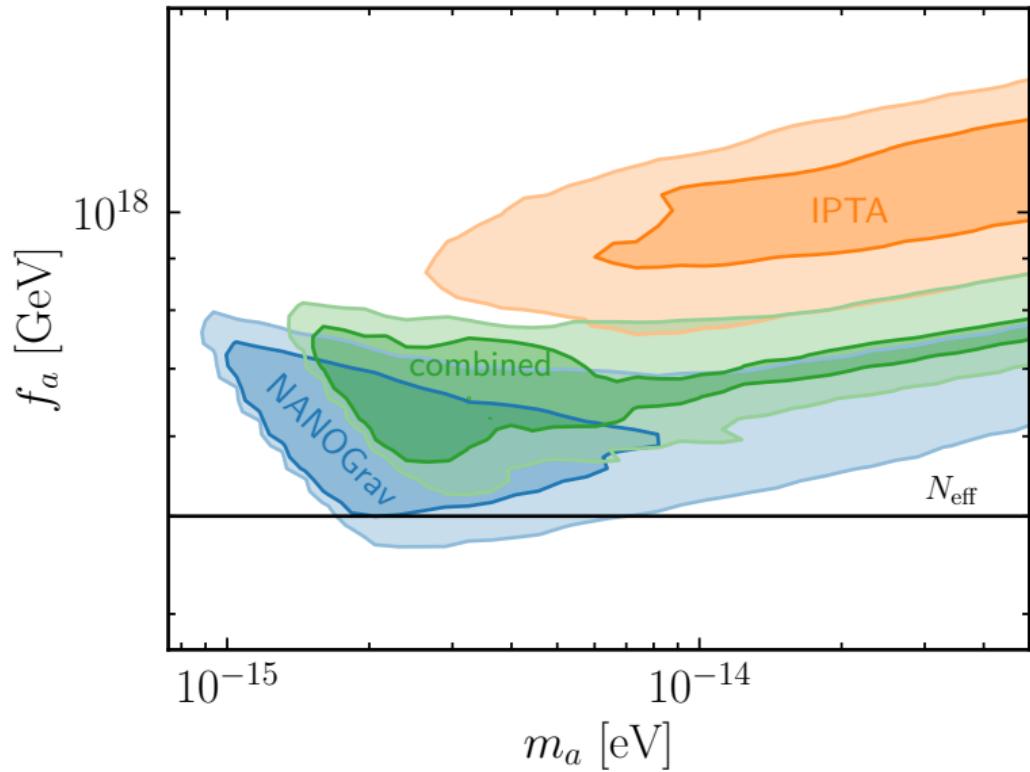
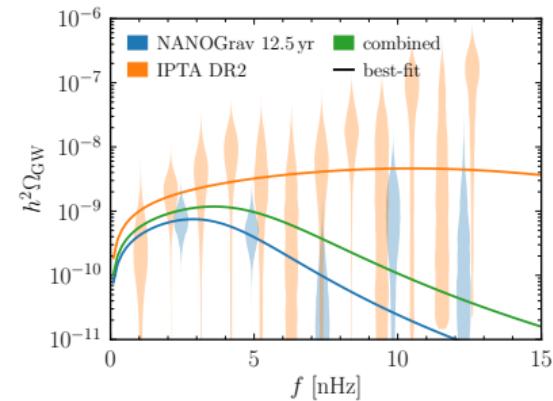
$$\Omega_{\text{GW}}^{\text{peak}} \sim \frac{\left(\rho_X^*/f_{\text{peak}}^* \right)^2}{\rho_c M_{\text{pl}}^2} \left(\frac{a_*}{a_0} \right)^4 \sim 10^{-7} \left(\frac{f_a}{M_{\text{pl}}} \right)^4 \left(100 \frac{\theta^2}{\alpha} \right)^{\frac{4}{3}}$$



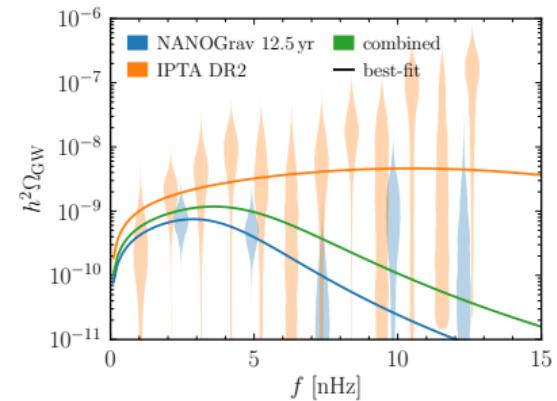
- spectral shape from lattice: [Ratzinger, Schwaller, Stefanek (SciPost Phys. 2022)]

$$\Omega_{\text{GW}} = \Omega_{\text{GW}}^{\text{peak}} \mathcal{S}(f/f_{\text{peak}}), \quad \mathcal{S}(x) = x^{0.73} \left[\frac{1}{2} (1 + x^{4.2}) \right]^{\frac{-4.96 - 0.73}{4.2}}$$

PTA Fit – Audible Axions

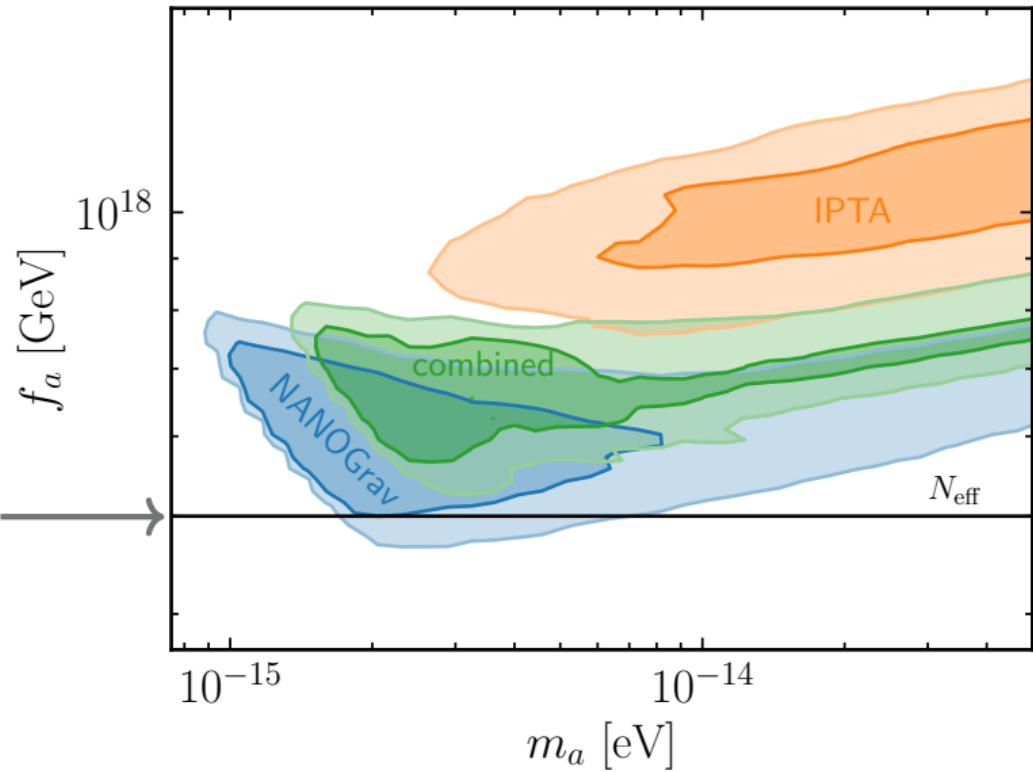


PTA Fit – Audible Axions



$$\Delta N_{\text{eff}} \sim 9.1 \left(\frac{\theta f_a}{M_{\text{Pl}}} \right)^2$$

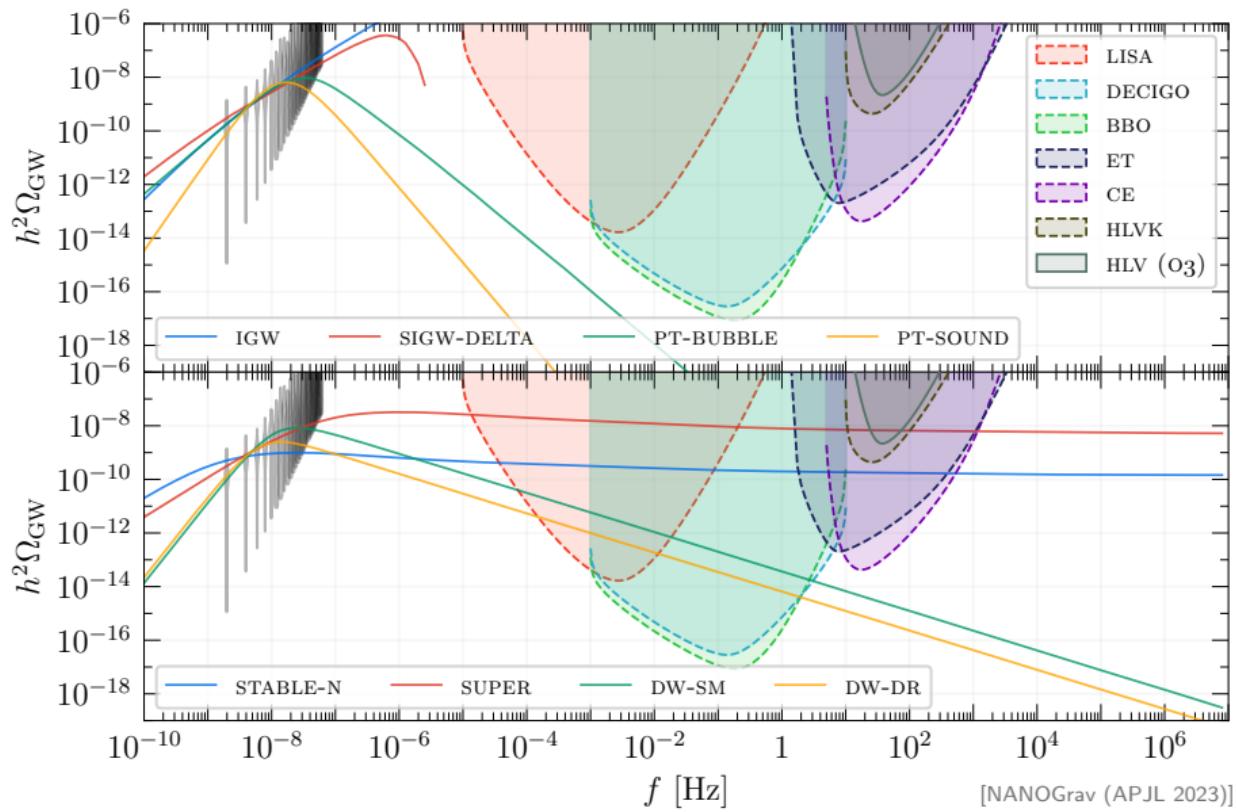
(also: $\Omega_a > \Omega_{\text{DM}}$)



Conclusions

1. Pulsar Timing Arrays
2. Cosmological Phase Transitions
3. Domain Walls
4. Bosonic Instabilities
5. Conclusions

Fits to NANOGrav 15-year dataset



Conclusion

Cosmological sources of a SGWB in PTA data:

- phase transition: Coleman-Weinberg model  (at least for 12.5-year dataset)

Conclusion

Cosmological sources of a SGWB in PTA data:

- phase transition: Coleman-Weinberg model ✓ (at least for 12.5-year dataset)
- domain walls: ALP or aligned axion ✗ (PBH overproduction)

Conclusion

Cosmological sources of a SGWB in PTA data:

- phase transition: Coleman-Weinberg model ✓ (at least for 12.5-year dataset)
- domain walls: ALP or aligned axion ✗ (PBH overproduction)
- bosonic inst.: audible axion ✗ (ΔN_{eff} , except for tiny region)

Conclusion

Cosmological sources of a SGWB in PTA data:

- **phase transition:** Coleman-Weinberg model ✓ (at least for 12.5-year dataset)
- **domain walls:** ALP or aligned axion ✗ (PBH overproduction)
- **bosonic inst.:** audible axion ✗ (ΔN_{eff} , except for tiny region)
- **cosmic strings:** global ALP strings ✗ (ΔN_{eff})
- **scalar-induced:** single-field w/ inflection point ✗ (CMB)

Conclusion

Cosmological sources of a SGWB in PTA data:

- phase transition: Coleman-Weinberg model ✓ (at least for 12.5-year dataset)
- domain walls: ALP or aligned axion ✗ (PBH overproduction)
- bosonic inst.: audible axion ✗ (ΔN_{eff} , except for tiny region)
- cosmic strings: global ALP strings ✗ (ΔN_{eff})
- scalar-induced: single-field w/ inflection point ✗ (CMB)

Thank you for your attention!



WEIZMANN
INSTITUTE
OF SCIENCE

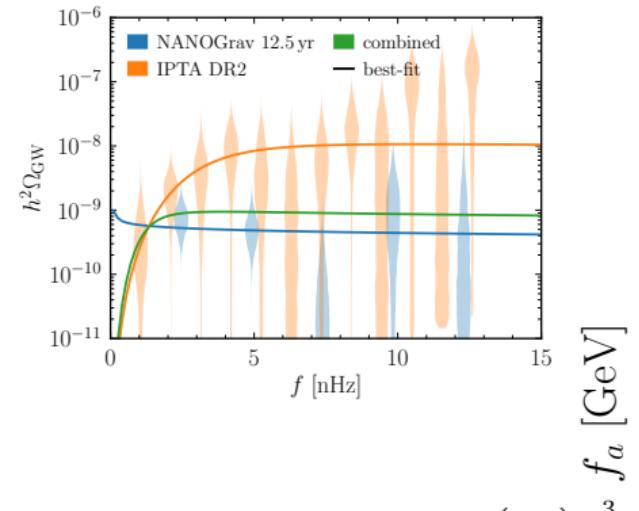
Primordial Gravitational Waves and Pulsar Timing Array Data

E.M., Morgante, Puchades Ibáñez, Ramberg, Ratzinger, Schenk, and Schwaller

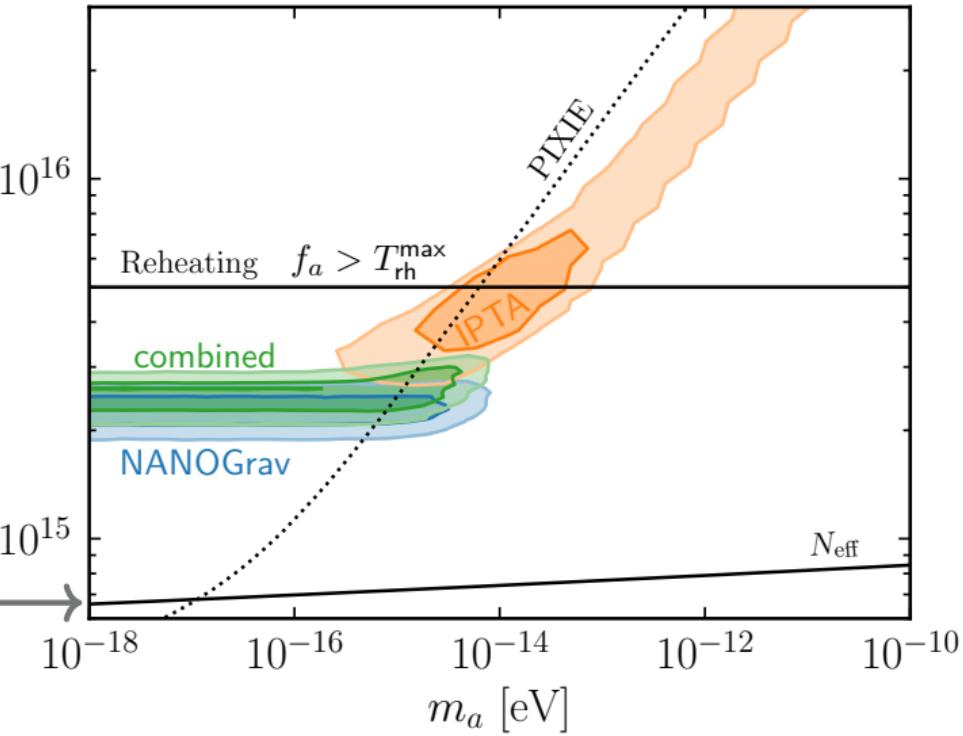
JHEP **10** (2023), 171 [arXiv:2306.14856 [hep-ph]]

backup slides

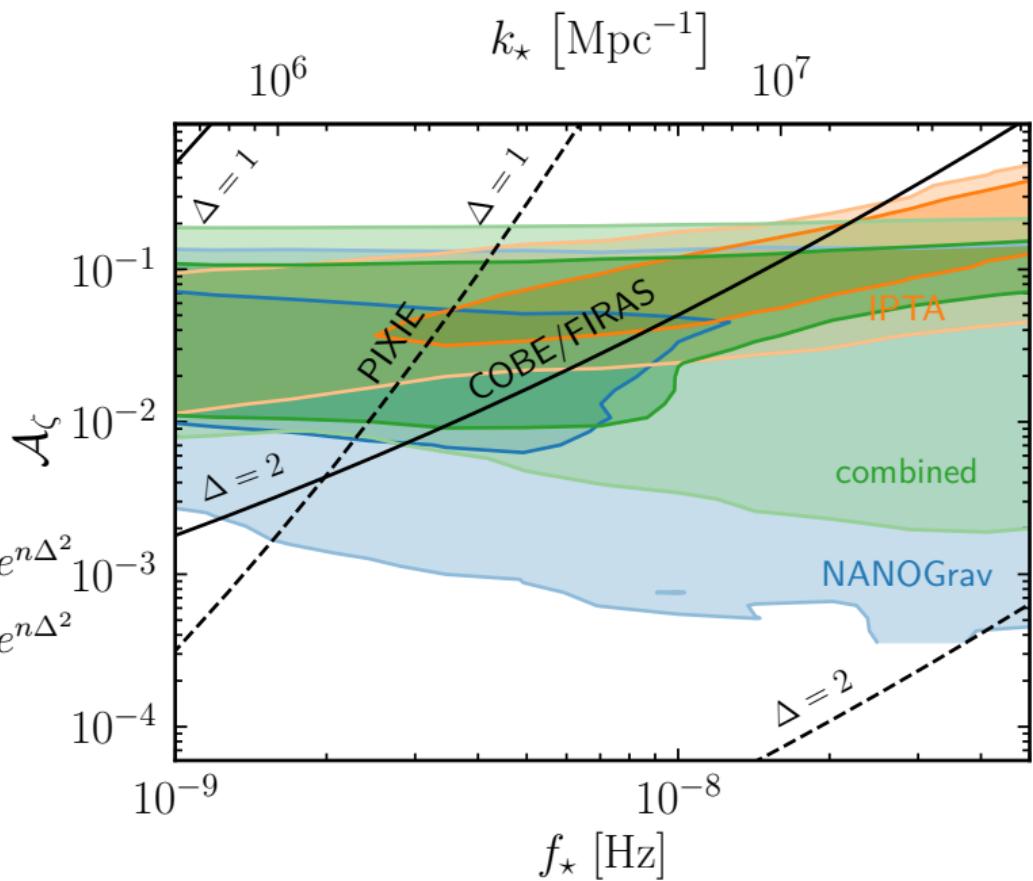
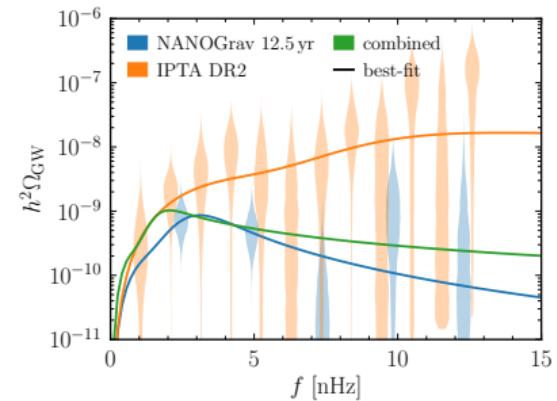
PTA Fit – Global ALP Strings



$$\Delta N_{\text{eff}} = 0.4 \left(\frac{f_a}{10^{15} \text{ GeV}} \right)^2 \left(\frac{\log \left(\frac{f_a}{m_a} \right)}{80} \right)^3$$

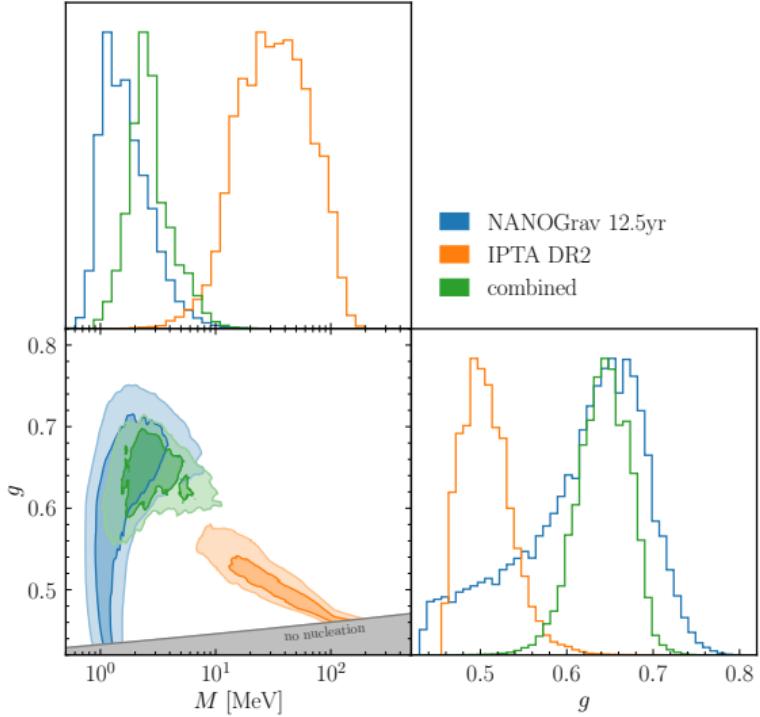


PTA Fit – Scalar-Induced GWs

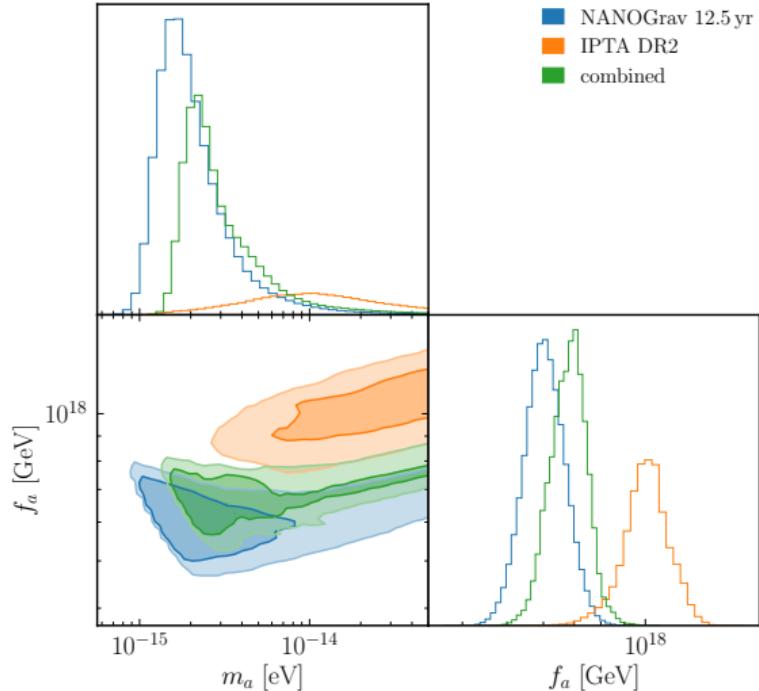


$$P_\zeta(k) = \begin{cases} \frac{A_\zeta e^{-\frac{\log^2(k/k_*)}{2\Delta^2}}}{\sqrt{2\pi}\Delta}, & \frac{k}{k_*} < e^{n\Delta^2} \\ \frac{A_\zeta e^{n\Delta^2}}{\sqrt{2\pi}\Delta} \left(\frac{k}{k_*}\right)^{-n}, & \frac{k}{k_*} > e^{n\Delta^2} \end{cases} 10^{-3}$$

Triangle Plots I

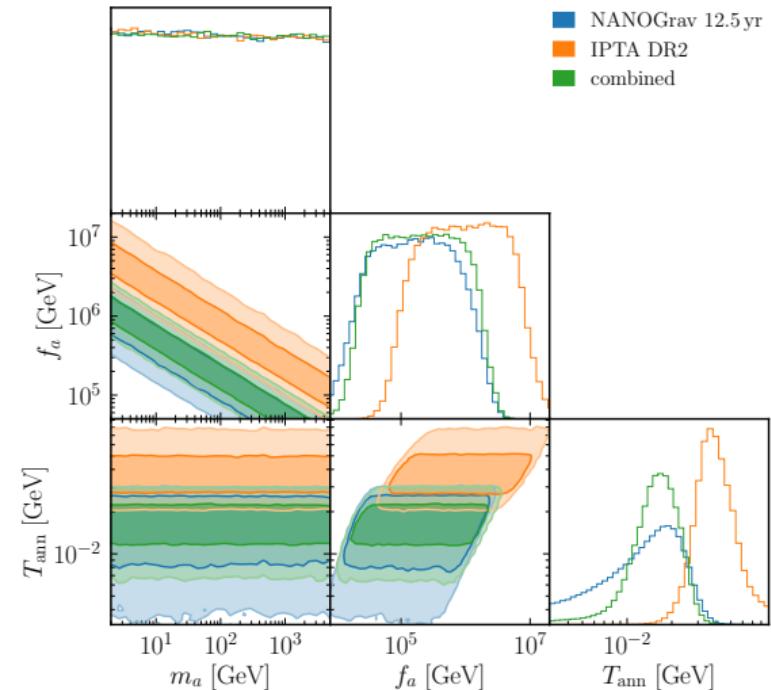


PhT in Coleman-Weinberg model

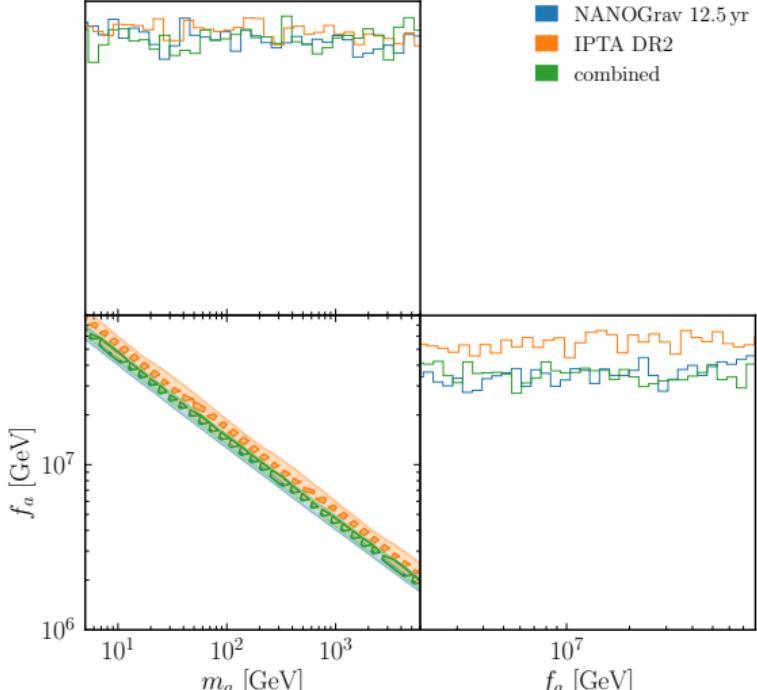


audible axions

Triangle Plots II

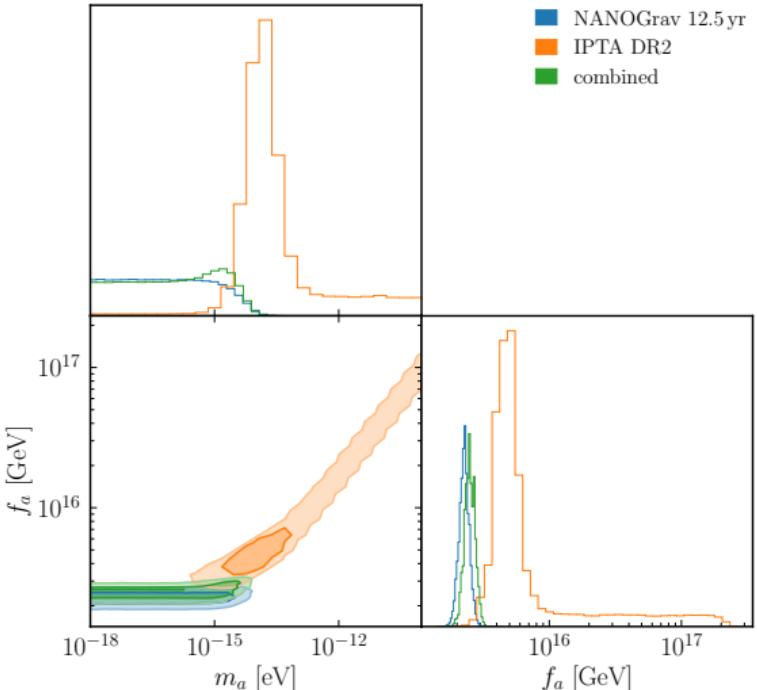


ALP domain walls

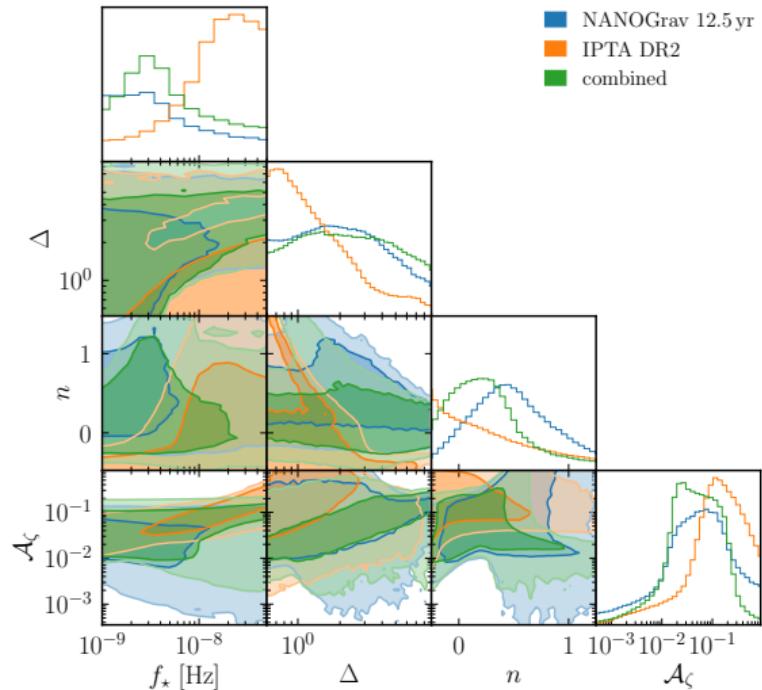


aligned axion domain walls

Triangle Plots III

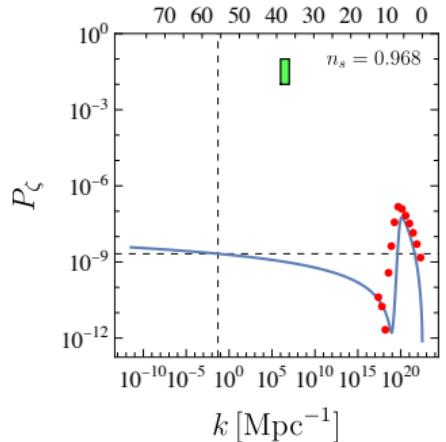
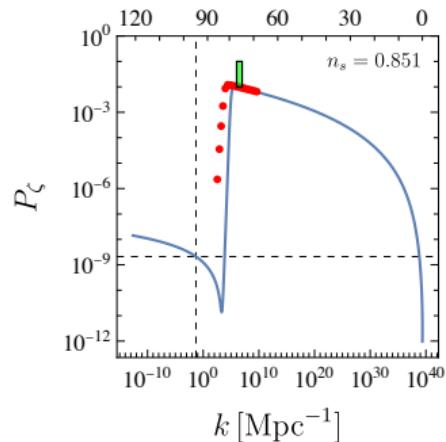
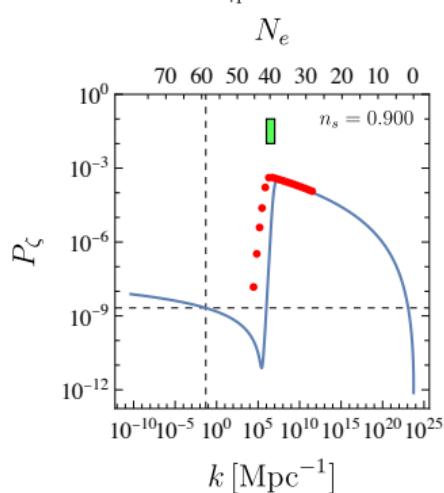
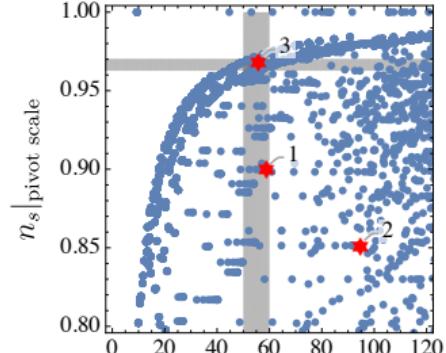
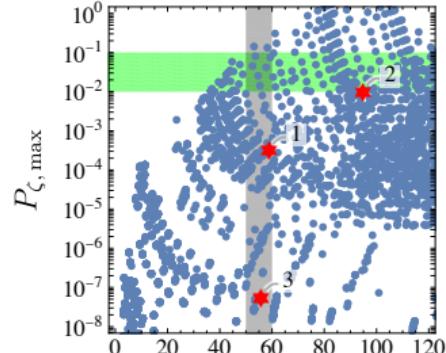
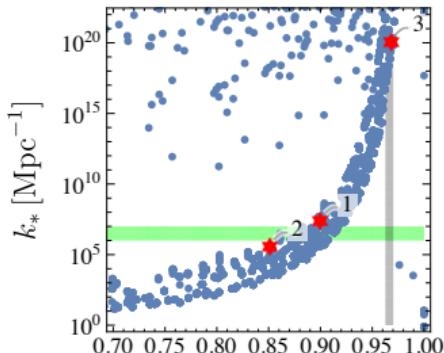


global ALP strings

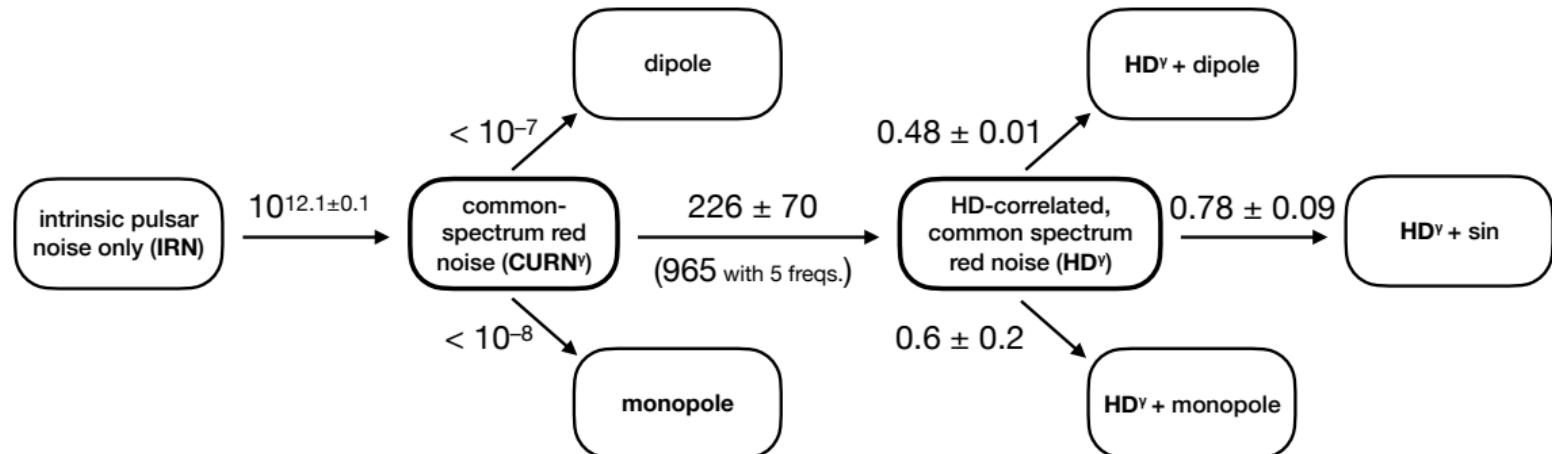


scalar-induced GWs

Single-field Inflation with Inflection Point

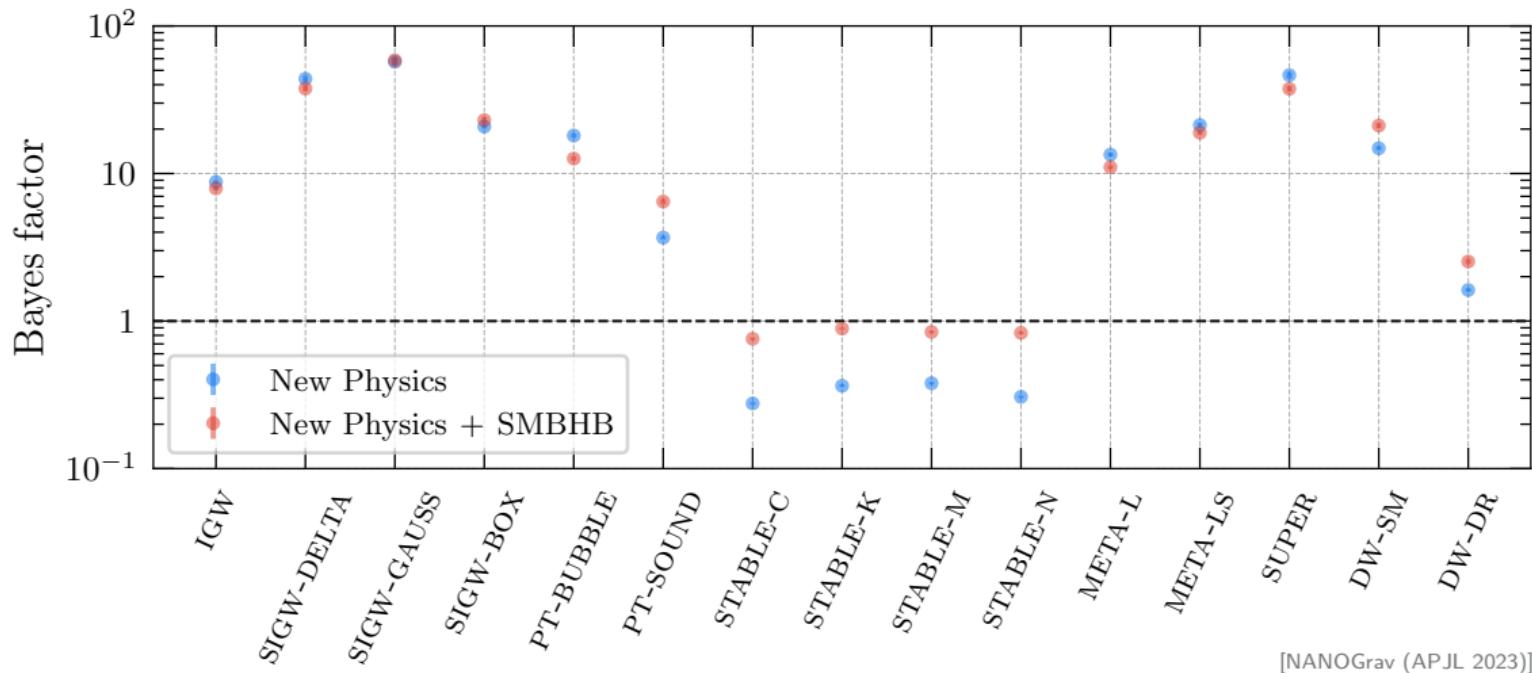


NANOGrav 15-year dataset: Bayesian Evidence



[NANOGrav (APJL 2023)]

NANOGrav 15-year dataset: Bayes Factors



[NANOGrav (APJL 2023)]