



A colorful mirror solution to the strong CP problem

Quentin Bonnefoy
UC Berkeley & LBNL

DESY
27/06/2023

with L. Hall, C.A. Manzari & C. Scherb
2303.06156 [hep-ph]
+ w.i.p. with the same people & A. McCune

The strong CP problem

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Q_L	3	1	1	1	1
Y_u	3	$\bar{\mathbf{3}}$	1	1	1
Y_d	3	1	$\bar{\mathbf{3}}$	1	1
Y_e	1	1	1	3	$\bar{\mathbf{3}}$

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[Jarlskog '85]

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In particular, **neutron electric dipole moment** (EDM)

[Baluni '79, Crewther/Di Vecchia/Veneziano/Witten '79]

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[Pendlebury et al '15]

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Strong CP problem :

$$\bar{\theta} \lesssim 10^{-10}$$

moment (EDM)

[Bardeen '79, Crewther/Di Vecchia/Veneziano/Witten '79]

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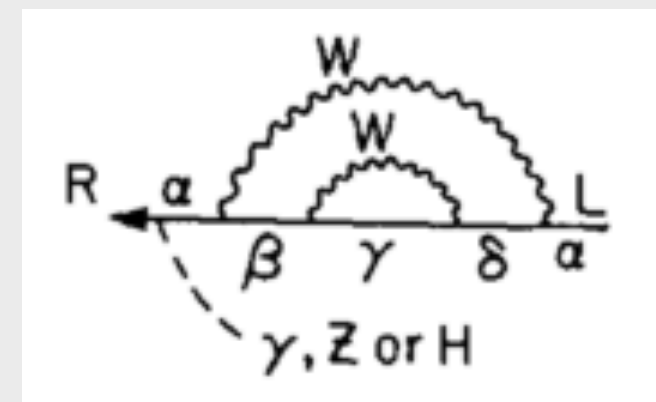
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[Ellis/Gaillard '79]



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By symmetry: **(C)P !**

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(C)P is not a symmetry of the SM ! Spontaneous breaking

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$$\text{CP} : \langle \phi \rangle \in \mathbb{C}$$

[Nelson '84, Barr '84]

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$SU(3) \times SU(2)_L \times SU(2)_R \times U(1)$

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↓

same + $\begin{array}{l} Q_R(\mathbf{3}, \mathbf{2}, 1/6) \\ d_L(\mathbf{3}, \mathbf{1}, -1/3) \\ u_L(\mathbf{3}, \mathbf{1}, 2/3) \end{array}$

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too light !

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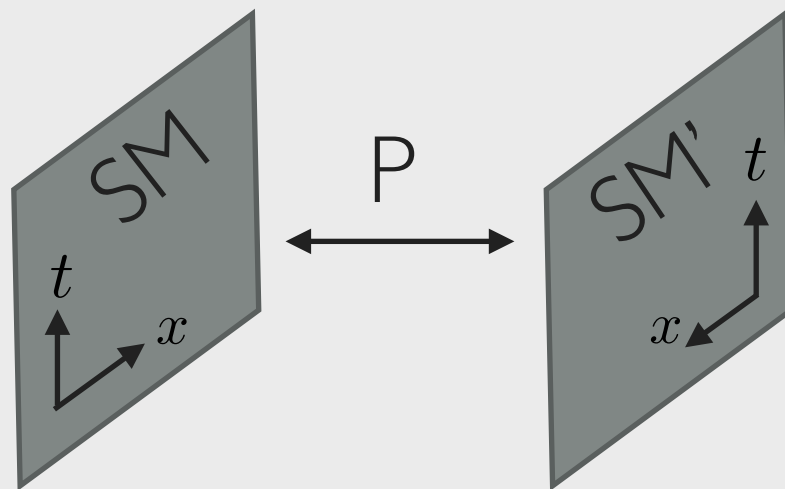
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In practice, both !

[Babu/Mohapatra '89, '90,
Barr/Chang/Senjanovic '91,
Hall/Harigaya '18, +Dunsky '18,
Craig/Garcia Garcia/Koszegi/
McCune '20, ...]

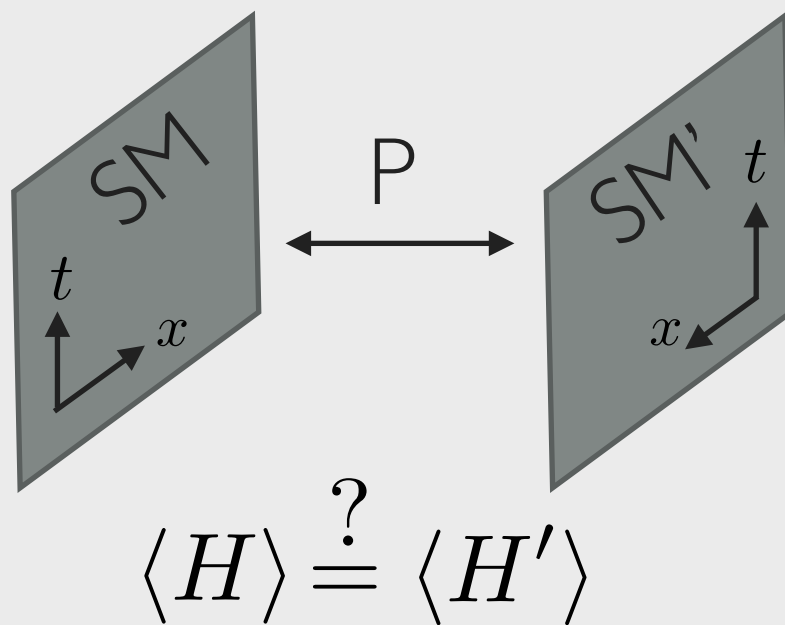
Parity solutions to the strong CP problem

Mirror world



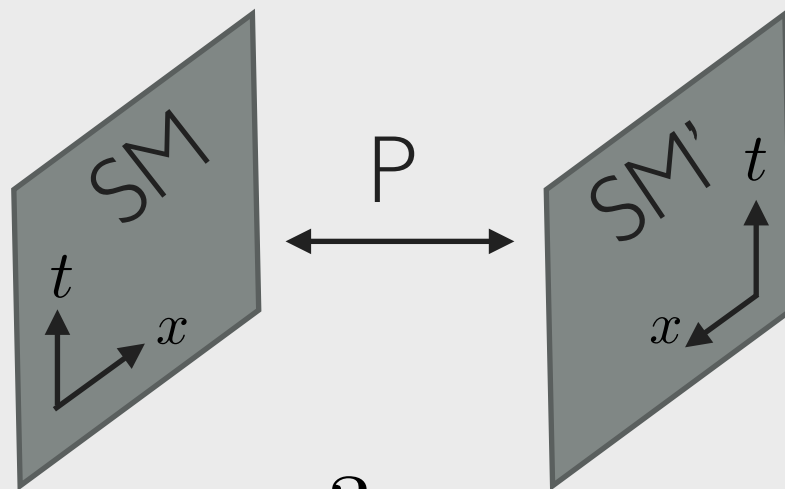
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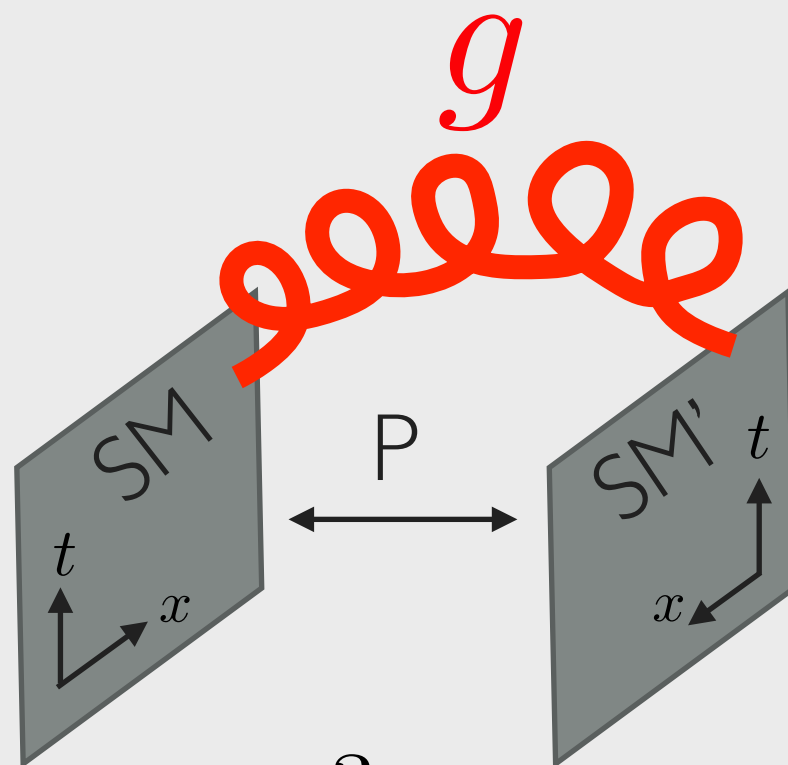
Mirror world and strong CP



$$\langle H \rangle \stackrel{?}{=} \langle H' \rangle$$
$$\bar{\theta}' = -\bar{\theta}$$

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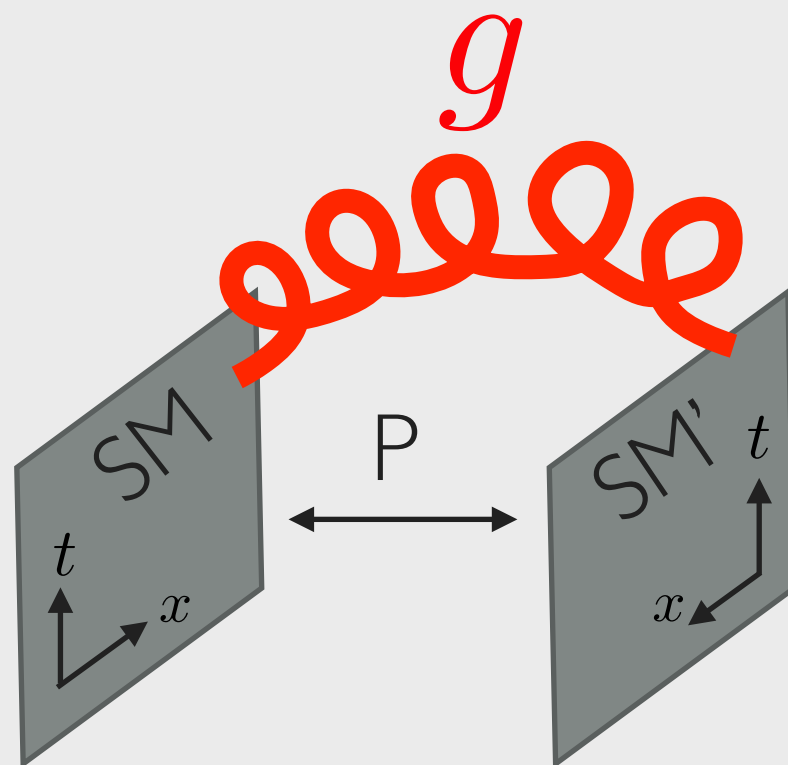
Mirror world and strong CP. Need shared color (P-invariant on its own) !



$$\langle H \rangle \stackrel{?}{=} \langle H' \rangle$$
$$\bar{\theta}_{\text{QCD}} = 0$$

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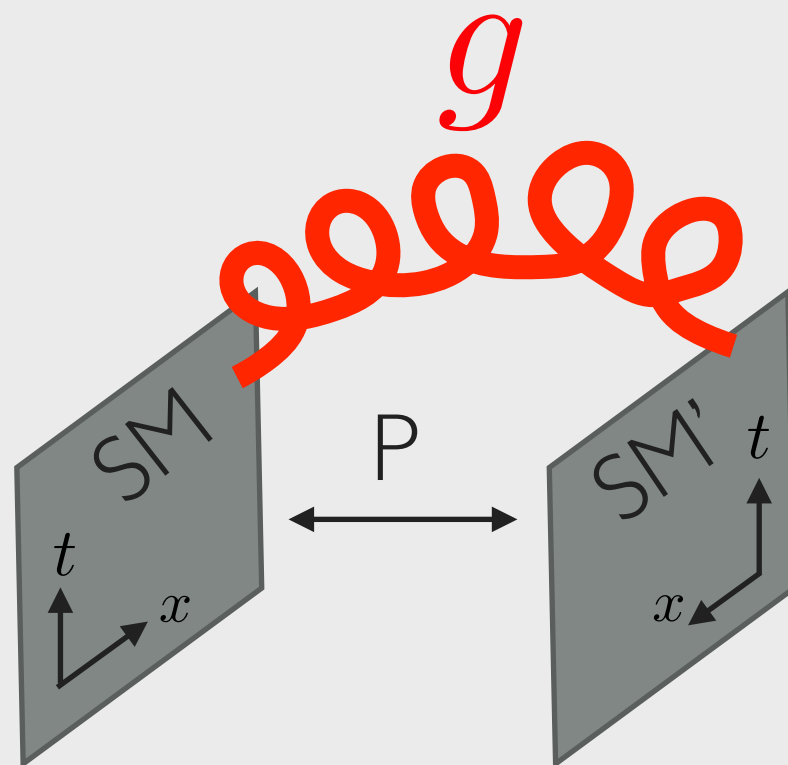


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Few parts of that landscape are explored !

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? Need $\bar{\theta} \approx 0$ even **below the scale of parity breaking**

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Similar to the SM case

[Ellis/Gaillard '79]

Parity solutions to the strong CP problem

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Not in all extensions

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Parity solutions to the strong CP problem

? Need $\bar{\theta} \approx 0$ even **below the scale of parity breaking**

Similar to the SM case

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Not in all extensions

[de Vries/Draper/Patel '21]

Easier in mirror models

[Barr/Chang/Senjanovic '91]

Our parity solution to the strong CP problem

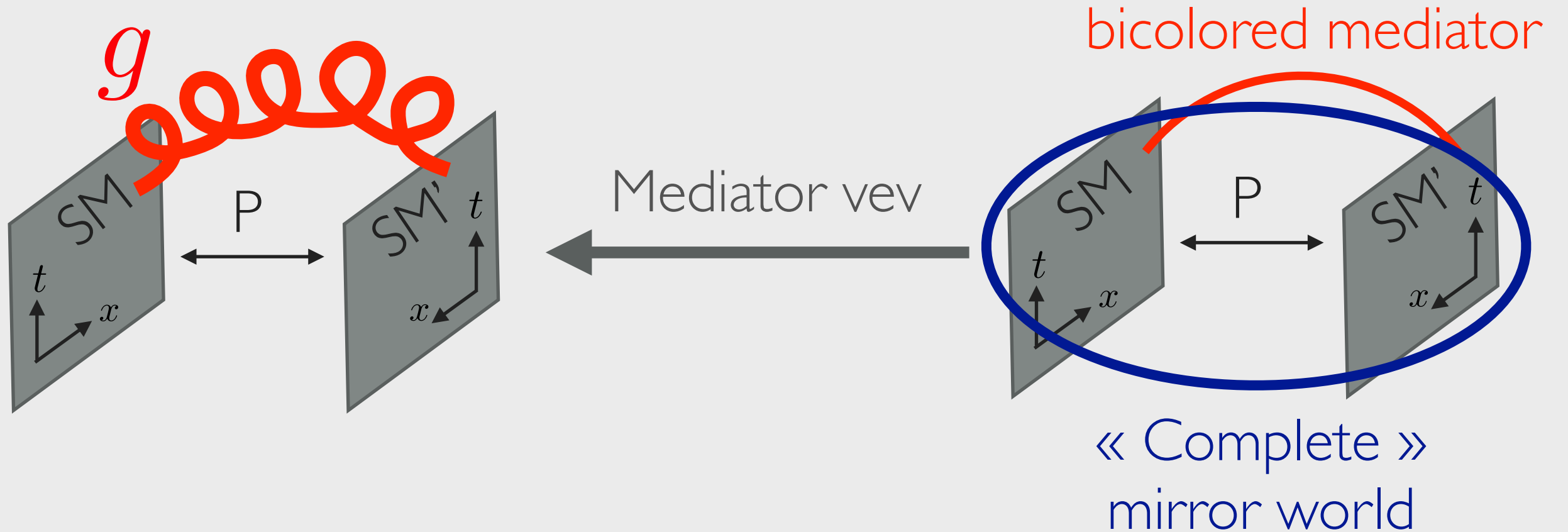
Our parity solution to the strong CP problem

We notice that



Our parity solution to the strong CP problem

We notice that

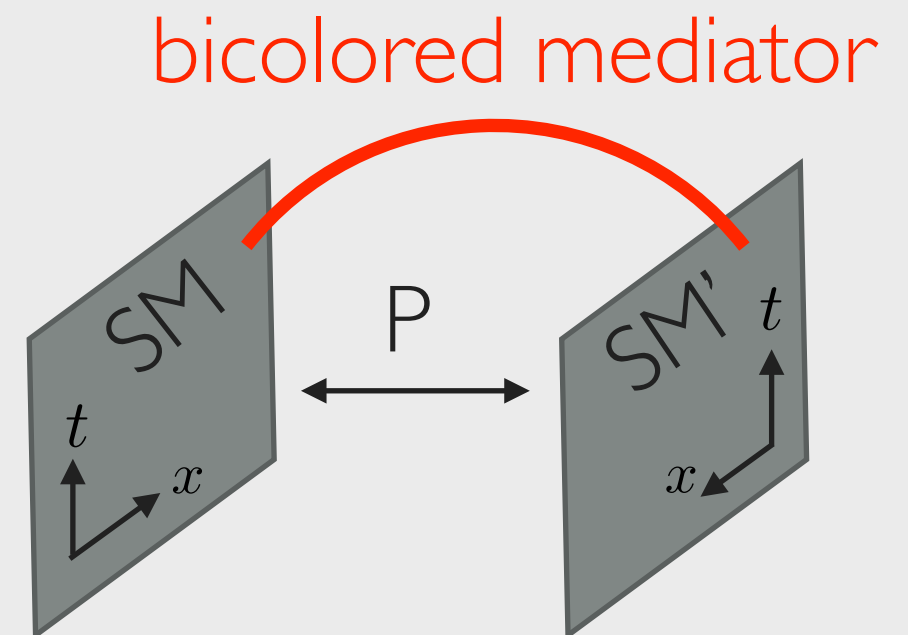


Our parity solution to the strong CP problem

We notice that



With the same starting point :

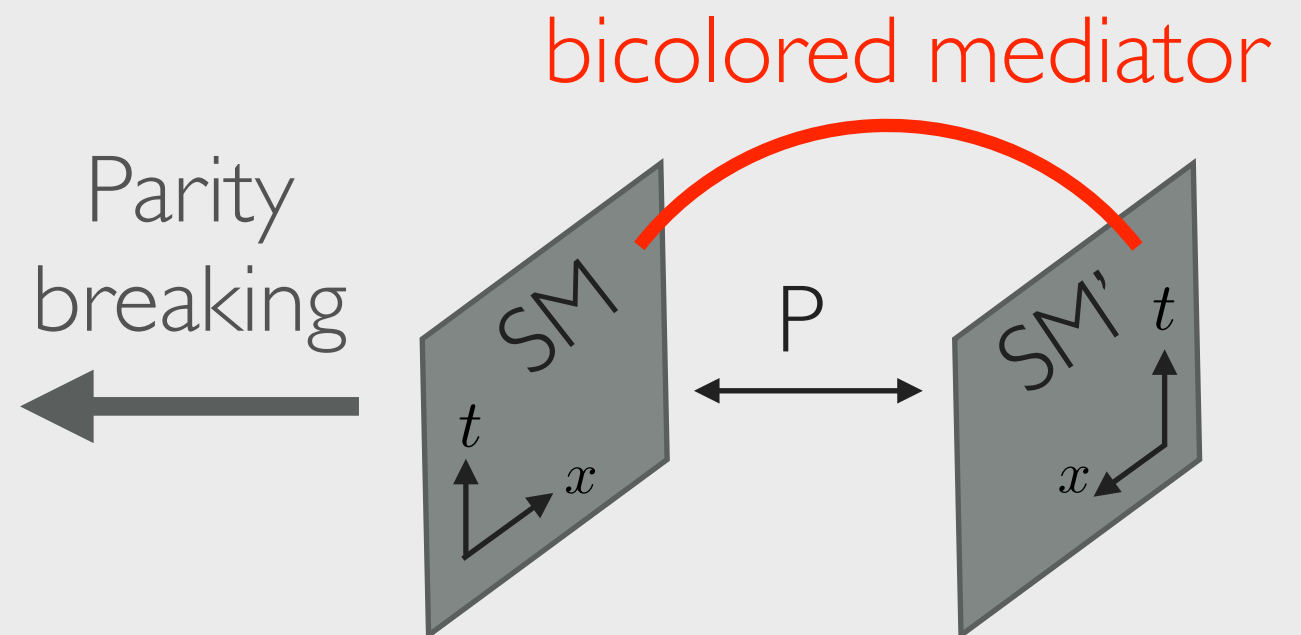


Our parity solution to the strong CP problem

We notice that



With the same starting point :

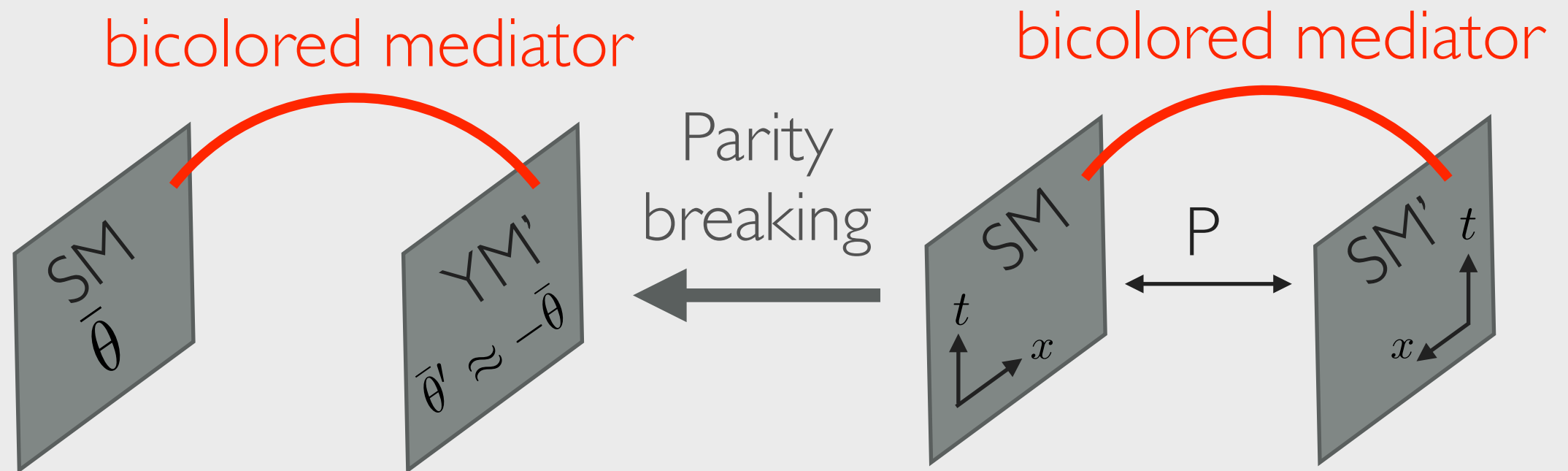


Our parity solution to the strong CP problem

We notice that



With the same starting point :

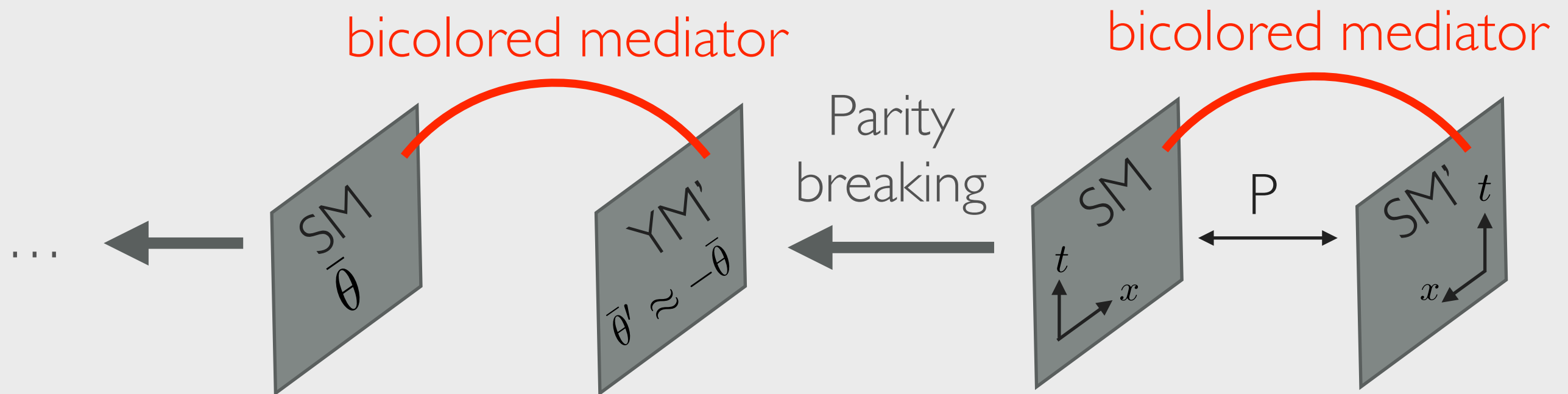


Our parity solution to the strong CP problem

We notice that

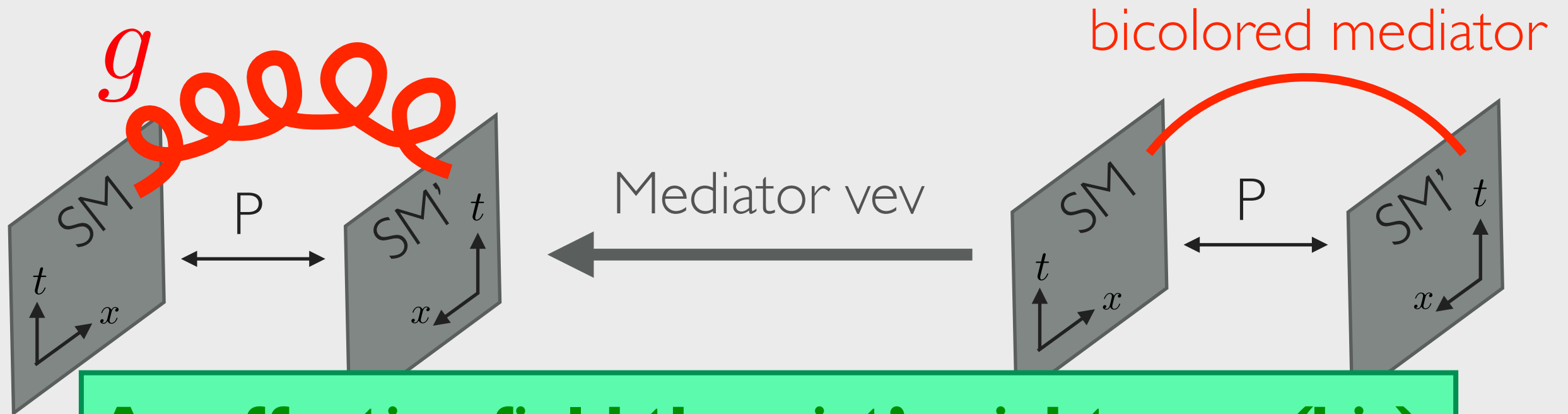


With the same starting point :



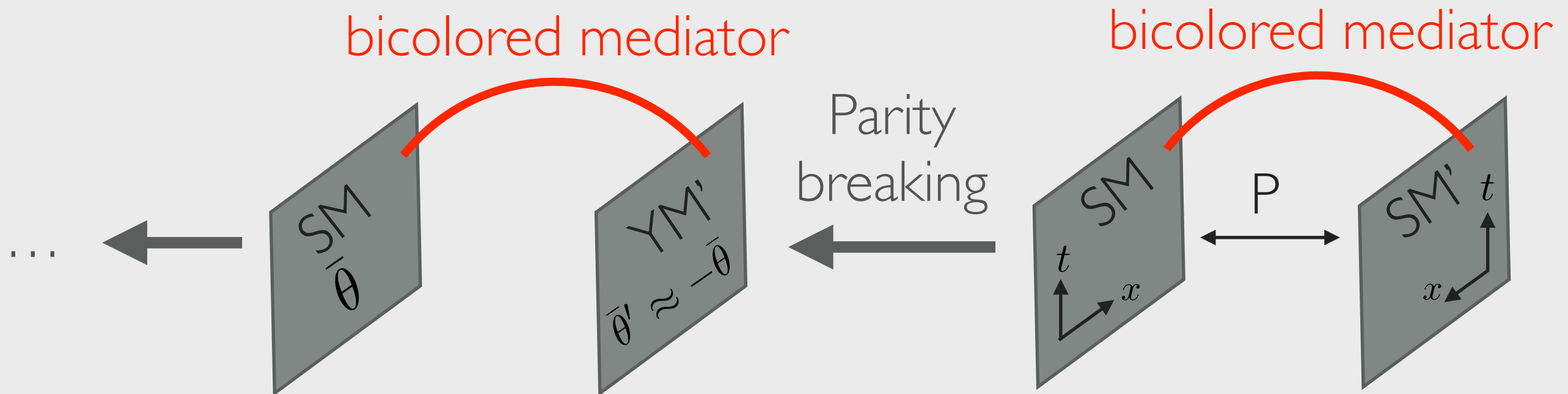
Our parity solution to the strong CP problem

We notice that

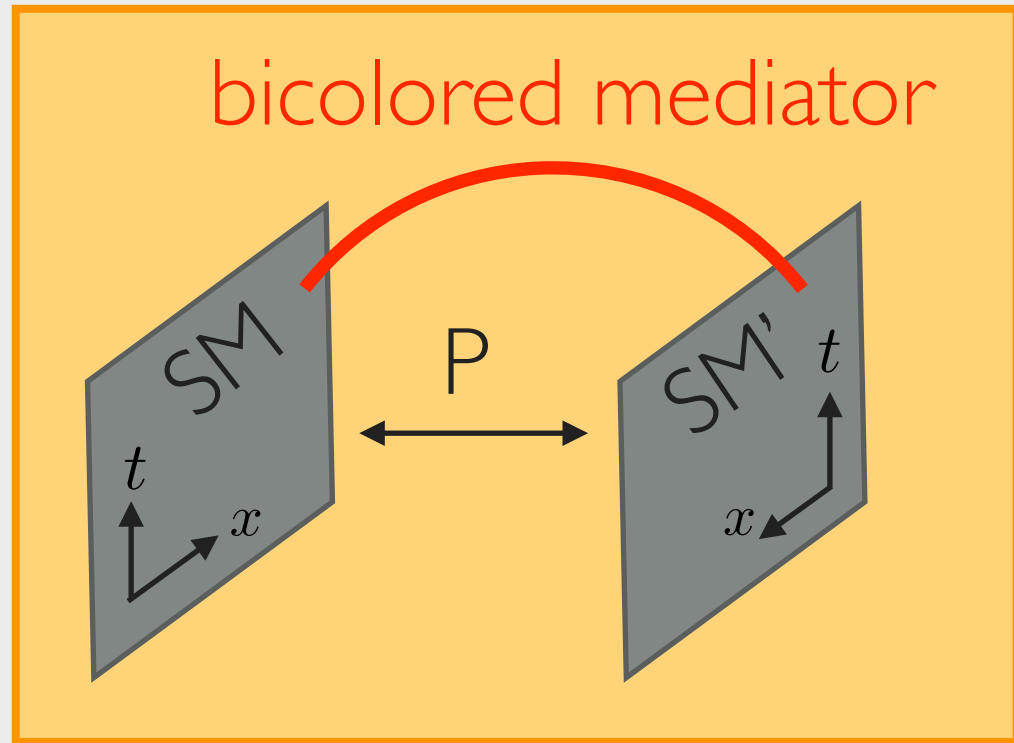


An effective field theorist's nightmare (bis)

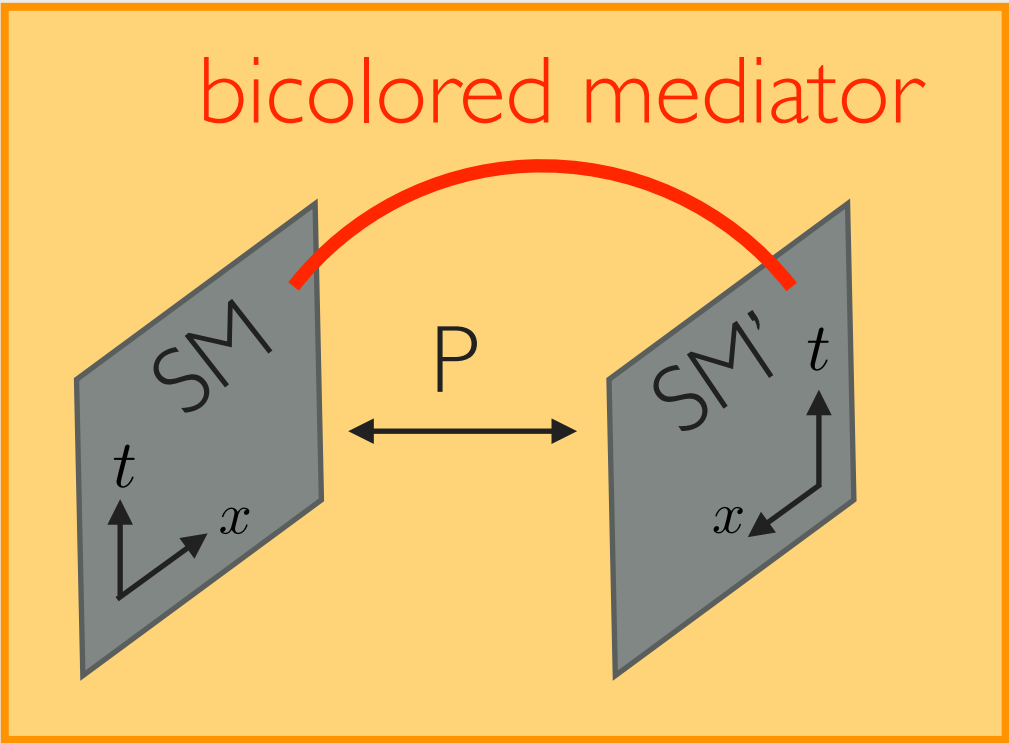
With the same starting point.



Our parity solution to the strong CP problem



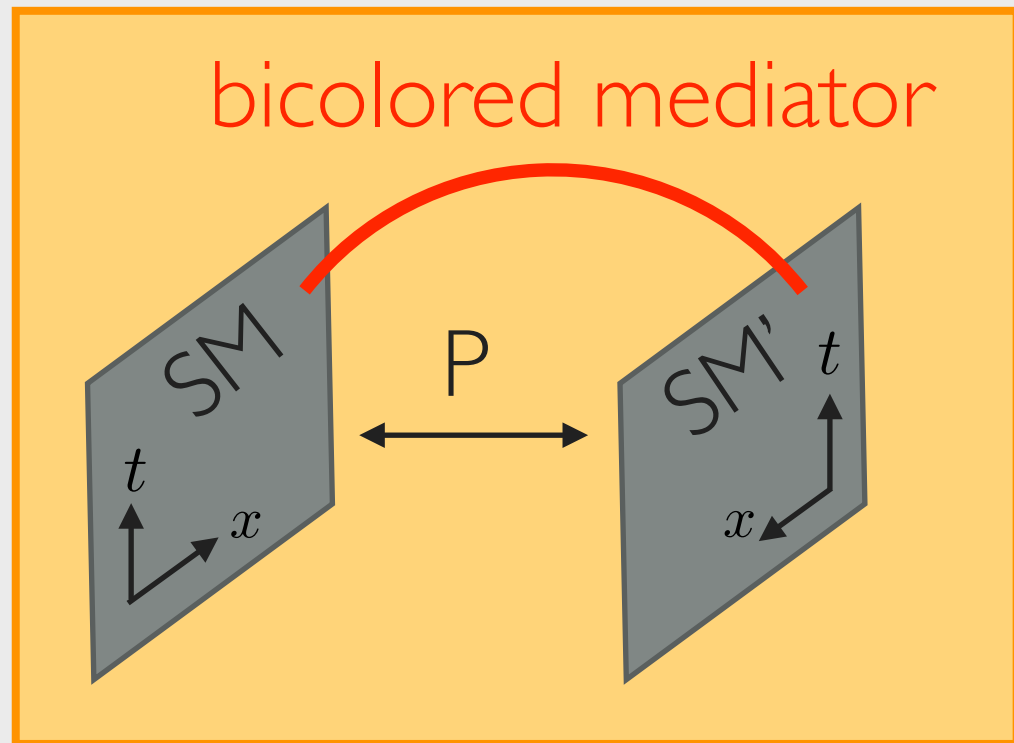
Our parity solution to the strong CP problem



	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
u^c	$\overline{\mathbf{3}}$	1	-2/3	1	1	0
d^c	$\overline{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\overline{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2



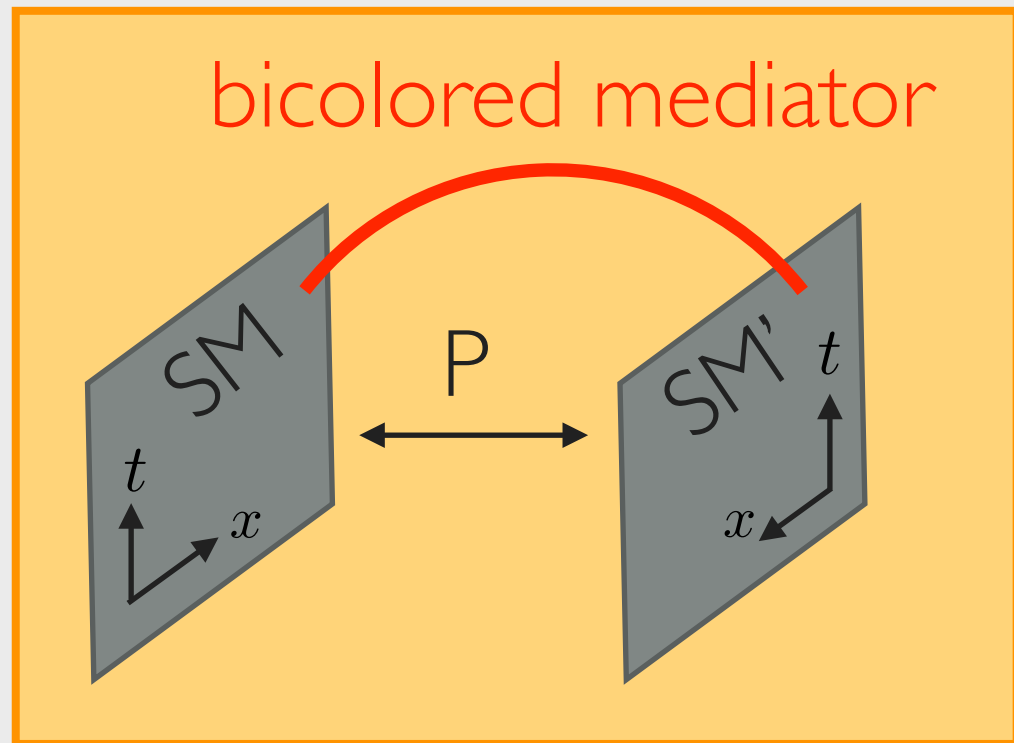
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	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
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L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

$$\left. \begin{array}{l} \theta' = -\theta \\ Y_q = Y_{q'}^\dagger \end{array} \right\} \implies \bar{\theta}' = -\bar{\theta}$$

Our parity solution to the strong CP problem

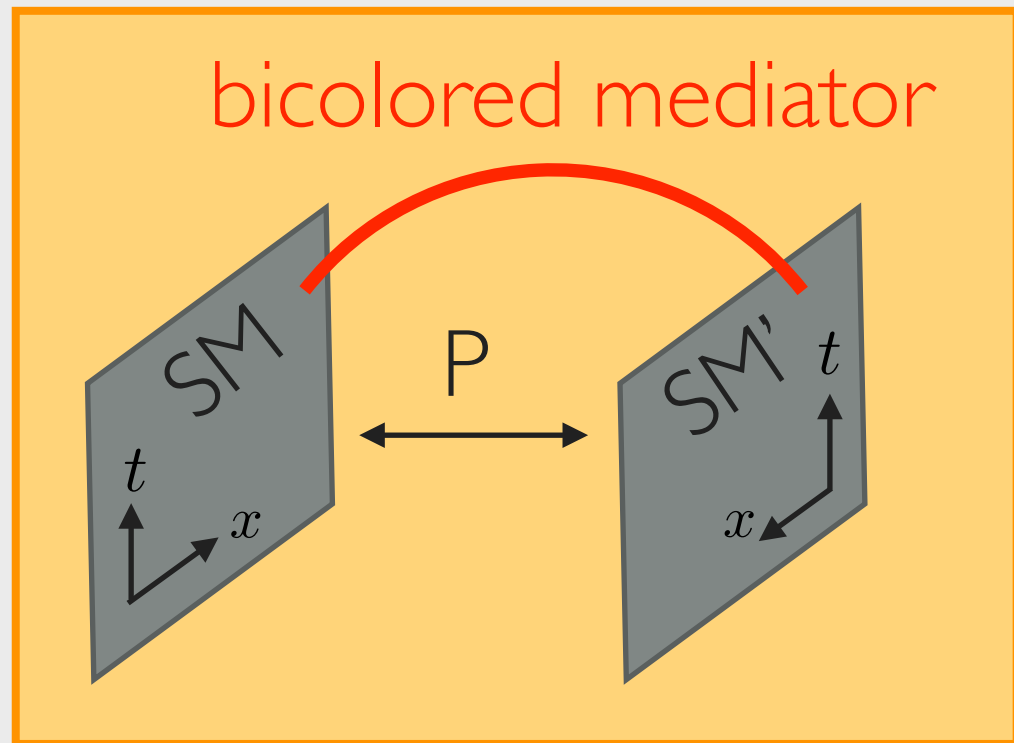


Bicolored mediator here:
bifundamental order
parameter $\langle \Sigma \rangle$

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
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e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

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Our parity solution to the strong CP problem



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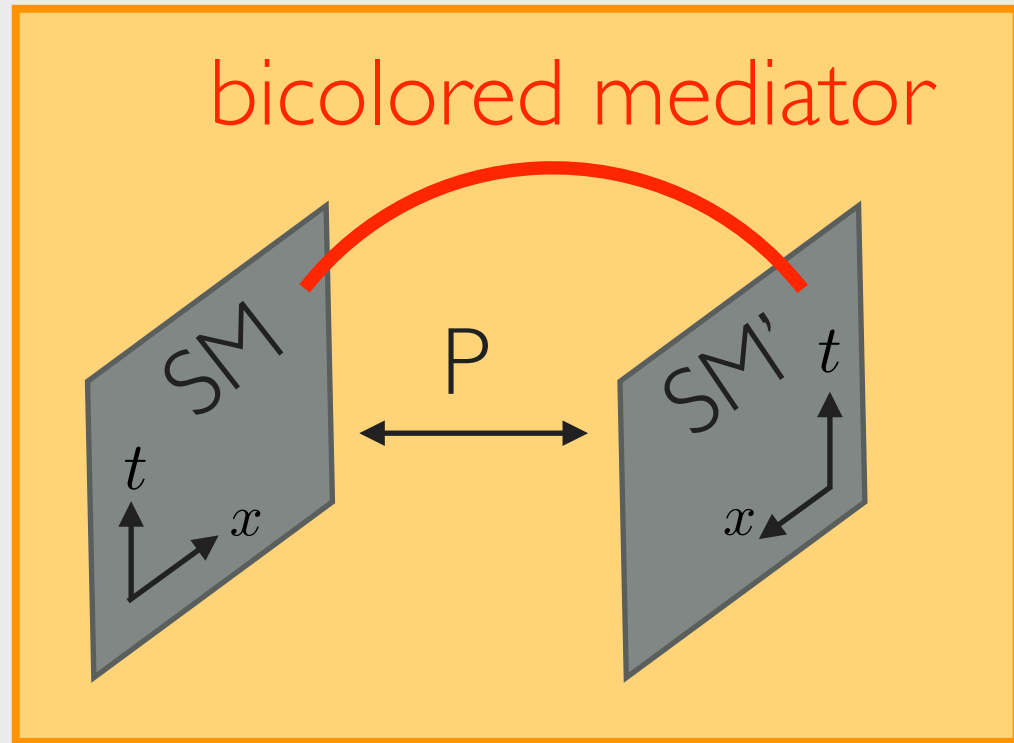
$$\langle \Sigma \rangle \propto v_3 \mathbf{1}$$

$$\Rightarrow g_3^2 G \tilde{G} = \big|_{\text{along QCD}} g_3'^2 G' \tilde{G}'$$

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
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d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

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Our parity solution to the strong CP problem



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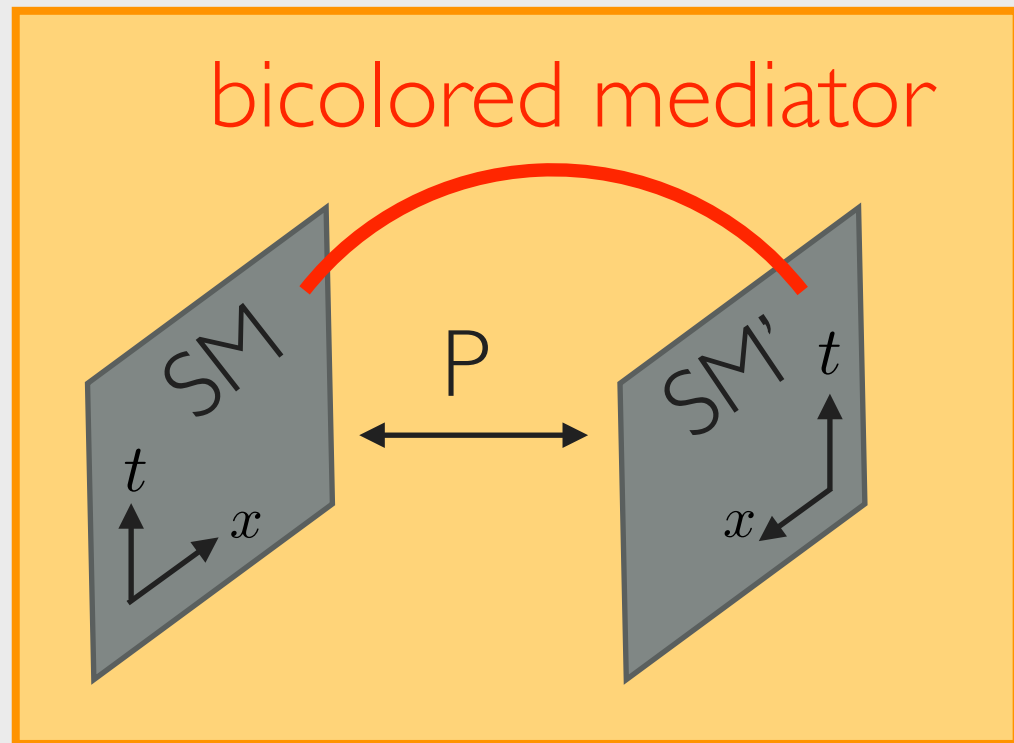
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d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2



$$\left. \begin{array}{l} \theta' = -\theta \\ Y_q = Y_{q'}^\dagger \end{array} \right\} \Rightarrow \bar{\theta}' = -\bar{\theta}$$

$$\Rightarrow \bar{\theta}_{\text{QCD}} = 0$$

Our parity solution to the strong CP problem



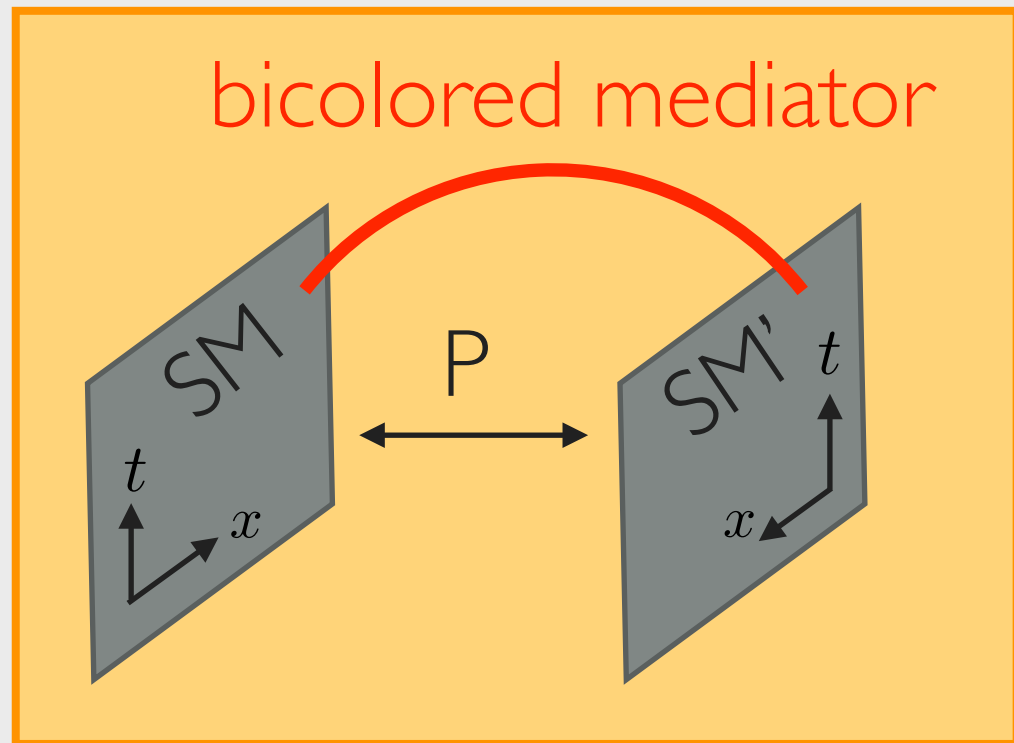
Below v_3 , colored mirror quarks

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
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d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

P

$$Y_q = Y_{q'}^\dagger$$

Our parity solution to the strong CP problem



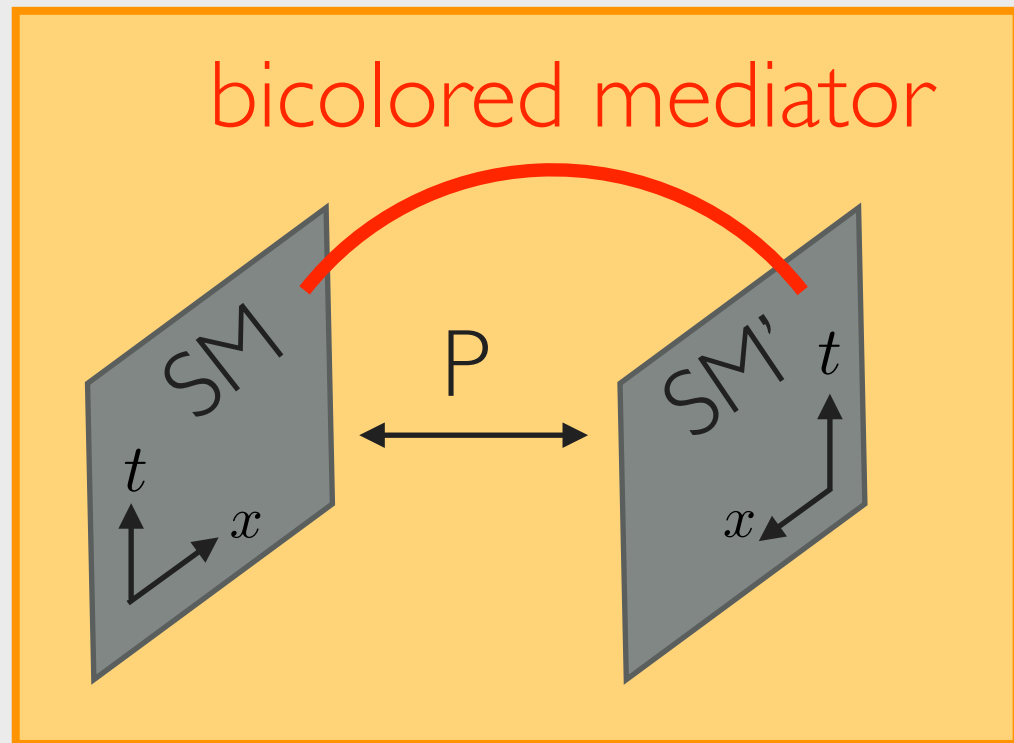
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e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

P

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem

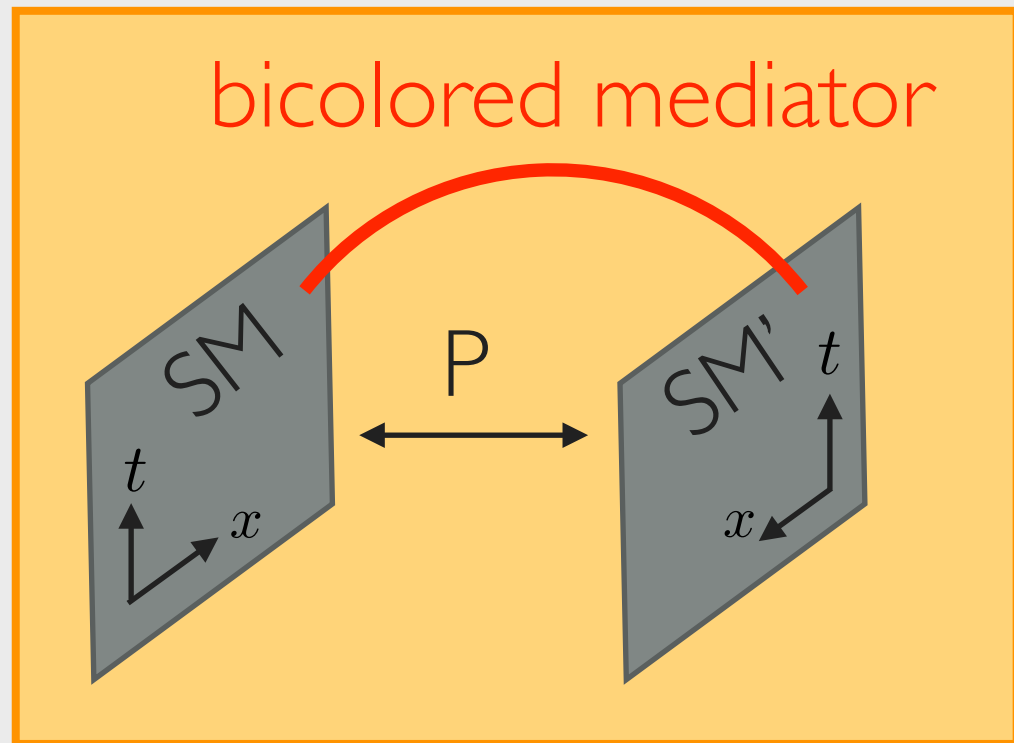


Below v_3 , colored mirror quarks : need $\langle H' \rangle \gg \langle H \rangle$ (hence P-breaking)

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
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d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem

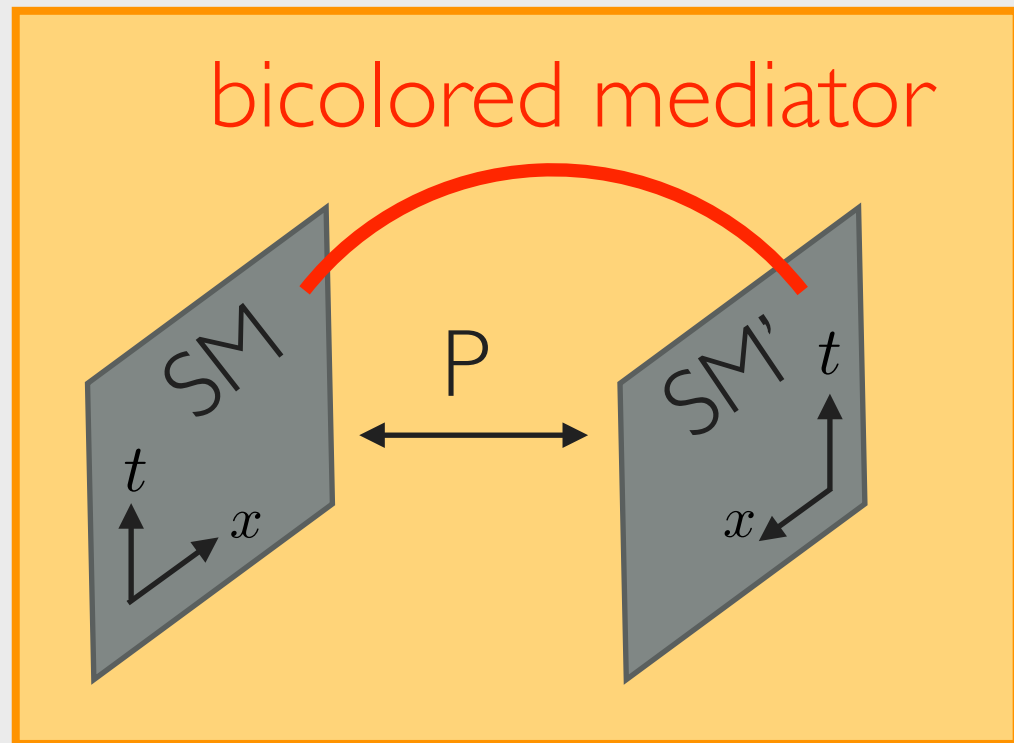


Below v_3 , colored mirror quarks : need $\langle H' \rangle \gg \langle H \rangle$ (hence P-breaking) $\equiv v_1$

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
u^c	$\bar{\mathbf{3}}$	1	-2/3	1	1	0
d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem



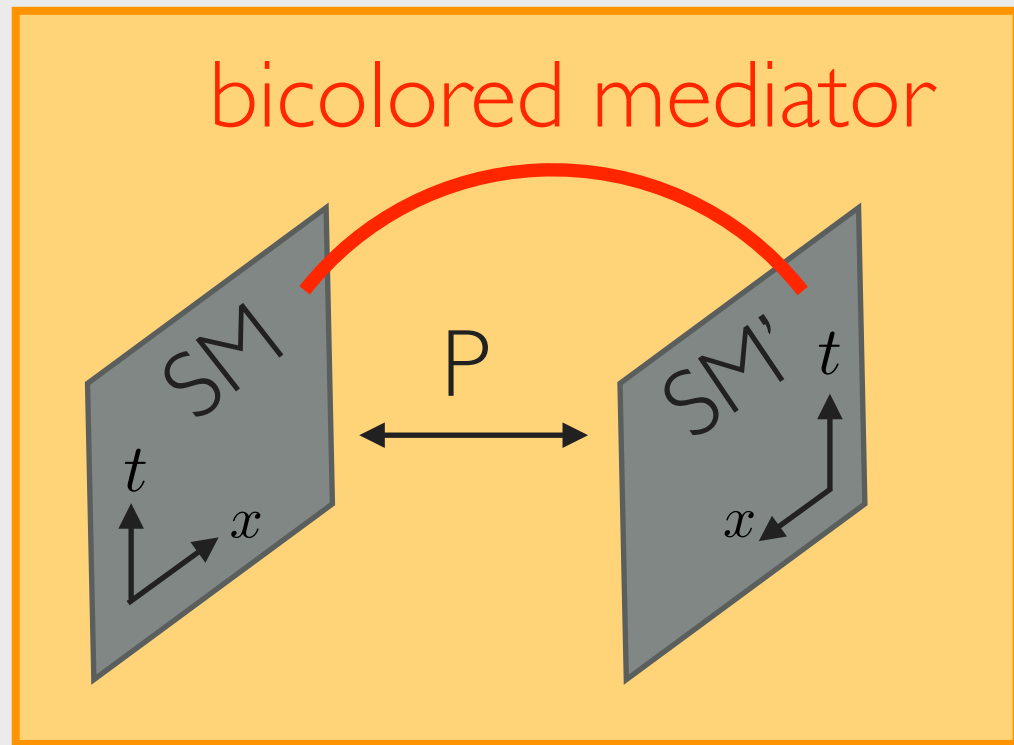
Below v_3 , colored mirror quarks : need $\langle H' \rangle \gg \langle H \rangle$ (hence P-breaking) $\equiv v'$

i.e. $v' \gtrsim 10^9 \text{ GeV}$

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
u^c	$\bar{\mathbf{3}}$	1	-2/3	1	1	0
d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem



Below v_3 , colored mirror quarks : need $\langle H' \rangle \gg \langle H \rangle$ (hence P-breaking) $\equiv v'$

Achieved through soft breaking or radiative corrections

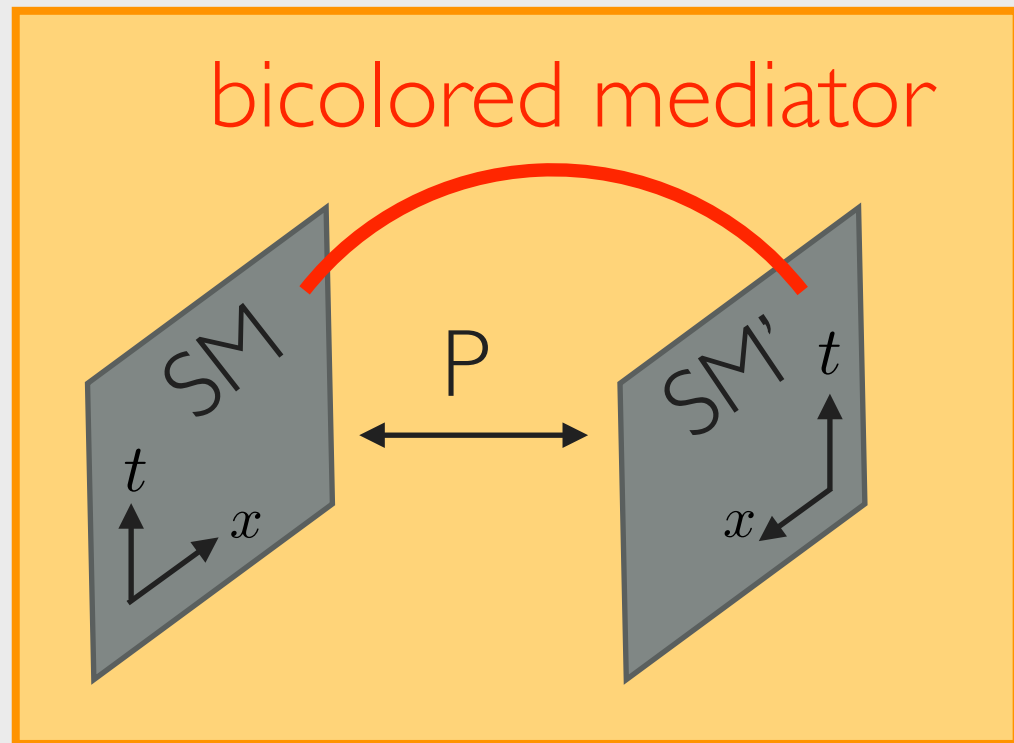
**[Babu/Mohapatra '89,
Hall/Harigaya '18]**

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	$\mathbf{3}$	$\mathbf{2}$	$1/6$	$\mathbf{1}$	$\mathbf{1}$	0
u^c	$\bar{\mathbf{3}}$	$\mathbf{1}$	$-2/3$	$\mathbf{1}$	$\mathbf{1}$	0
d^c	$\bar{\mathbf{3}}$	$\mathbf{1}$	$1/3$	$\mathbf{1}$	$\mathbf{1}$	0
L	$\mathbf{1}$	$\mathbf{2}$	$-1/2$	$\mathbf{1}$	$\mathbf{1}$	0
e^c	$\mathbf{1}$	$\mathbf{1}$	-1	$\mathbf{1}$	$\mathbf{1}$	0
H	$\mathbf{1}$	$\mathbf{2}$	$1/2$	$\mathbf{1}$	$\mathbf{1}$	0
Q'	$\mathbf{1}$	$\mathbf{1}$	0	$\bar{\mathbf{3}}$	$\mathbf{2}$	$-1/6$
u'^c	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{3}$	$\mathbf{1}$	$2/3$
d'^c	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{3}$	$\mathbf{1}$	$-1/3$
L'	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{1}$	$\mathbf{2}$	$1/2$
e'^c	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{1}$	$\mathbf{1}$	1
H'	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{1}$	$\mathbf{2}$	$-1/2$

P

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem



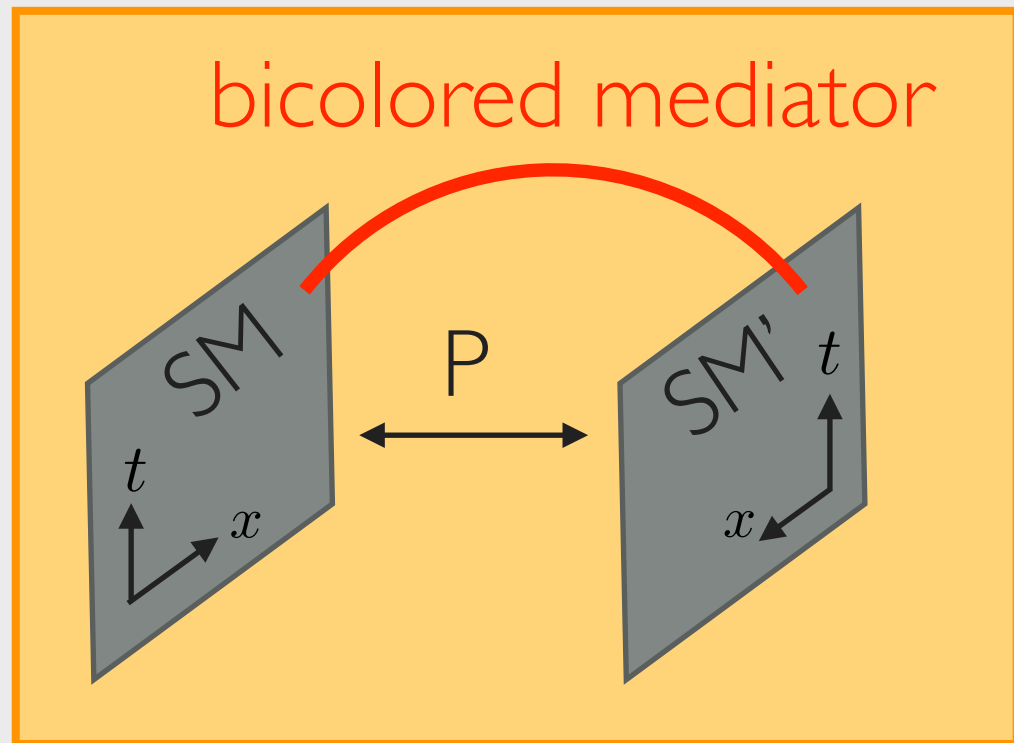
Very predictive model, two BSM scales : v_3 and v'

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	$\mathbf{3}$	$\mathbf{2}$	$1/6$	$\mathbf{1}$	$\mathbf{1}$	0
u^c	$\bar{\mathbf{3}}$	$\mathbf{1}$	$-2/3$	$\mathbf{1}$	$\mathbf{1}$	0
d^c	$\bar{\mathbf{3}}$	$\mathbf{1}$	$1/3$	$\mathbf{1}$	$\mathbf{1}$	0
L	$\mathbf{1}$	$\mathbf{2}$	$-1/2$	$\mathbf{1}$	$\mathbf{1}$	0
e^c	$\mathbf{1}$	$\mathbf{1}$	-1	$\mathbf{1}$	$\mathbf{1}$	0
H	$\mathbf{1}$	$\mathbf{2}$	$1/2$	$\mathbf{1}$	$\mathbf{1}$	0
Q'	$\mathbf{1}$	$\mathbf{1}$	0	$\bar{\mathbf{3}}$	$\mathbf{2}$	$-1/6$
u'^c	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{3}$	$\mathbf{1}$	$2/3$
d'^c	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{3}$	$\mathbf{1}$	$-1/3$
L'	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{1}$	$\mathbf{2}$	$1/2$
e'^c	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{1}$	$\mathbf{1}$	1
H'	$\mathbf{1}$	$\mathbf{1}$	0	$\mathbf{1}$	$\mathbf{2}$	$-1/2$

P

$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem



Very predictive model, two BSM scales : v_3 and v'

Different pheno on the parameter space. For $v_3 \ll v'$, **colored bosons** as lightest BSM states !

	$SU(3)$	$SU(2)_L$	$U(1)_Y$	$SU(3)'$	$SU(2)'$	$U(1)'$
Q	3	2	1/6	1	1	0
u^c	$\bar{\mathbf{3}}$	1	-2/3	1	1	0
d^c	$\bar{\mathbf{3}}$	1	1/3	1	1	0
L	1	2	-1/2	1	1	0
e^c	1	1	-1	1	1	0
H	1	2	1/2	1	1	0
Q'	1	1	0	$\bar{\mathbf{3}}$	2	-1/6
u'^c	1	1	0	3	1	2/3
d'^c	1	1	0	3	1	-1/3
L'	1	1	0	1	2	1/2
e'^c	1	1	0	1	1	1
H'	1	1	0	1	2	-1/2

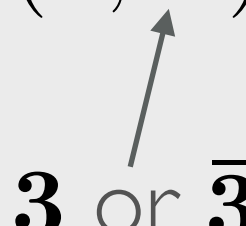
$$Y_q = Y_{q'}^\dagger \implies \frac{m_q}{m_{q'}} = \frac{\langle H \rangle}{\langle H' \rangle}$$

Our parity solution to the strong CP problem

Bicolored mediator : a scalar or strongly interacting fermions.

Our parity solution to the strong CP problem

Bicolored mediator : a scalar or strongly interacting fermions.

- Σ in $(\mathbf{3}, \mathbf{3}')$ of $SU(3) \times SU(3)'$

 $\mathbf{3}$ or $\bar{\mathbf{3}}$

Our parity solution to the strong CP problem

Bicolored mediator : a **scalar** or **strongly interacting fermions**.

- Σ in $(\mathbf{3}, \mathbf{3}')$ of $SU(3) \times SU(3)'$ with potential

$\mathbf{3}$ or $\bar{\mathbf{3}}$



$$V(\Sigma) = -m^2 \text{Tr}(\Sigma \Sigma^\dagger) + c \text{Tr}^2(\Sigma \Sigma^\dagger) \\ + \tilde{c} \text{Tr}(\Sigma \Sigma^\dagger)^2 + (\tilde{m} \det(\Sigma) + h.c.)$$

Our parity solution to the strong CP problem

Bicolored mediator : a **scalar** or **strongly interacting fermions**.

- Σ in $(\mathbf{3}, \mathbf{3}')$ of $SU(3) \times SU(3)'$ with potential
$$V(\Sigma) = -m^2 \text{Tr}(\Sigma \Sigma^\dagger) + c \text{Tr}^2(\Sigma \Sigma^\dagger) + \tilde{c} \text{Tr}(\Sigma \Sigma^\dagger)^2 + (\tilde{m} \det(\Sigma) + h.c.)$$

$\mathbf{3}$ or $\bar{\mathbf{3}}$

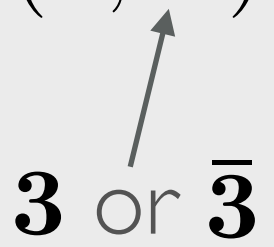


Breaking to the diagonal $SU(3)$ in a large fraction of parameter space

[Bai/Dobrescu '17]

Our parity solution to the strong CP problem

Bicolored mediator : a **scalar** or **strongly interacting fermions**.

- Σ in $(\mathbf{3}, \mathbf{3}')$ of $SU(3) \times SU(3)'$ with potential

$$V(\Sigma) = -m^2 \text{Tr}(\Sigma \Sigma^\dagger) + c \text{Tr}^2(\Sigma \Sigma^\dagger) + \tilde{c} \text{Tr}(\Sigma \Sigma^\dagger)^2 + (\tilde{m} \det(\Sigma) + h.c.)$$

Breaking to the diagonal $SU(3)$ in a large fraction of parameter space (but no (C)P breaking)

[Bai/Dobrescu '17]

Our parity solution to the strong CP problem

Bicolored mediator : a **scalar** or **strongly interacting fermions**.

- Σ in $(\mathbf{3}, \mathbf{3}')$ of $SU(3) \times SU(3)'$ with potential
$$V(\Sigma) = -m^2 \text{Tr}(\Sigma \Sigma^\dagger) + c \text{Tr}^2(\Sigma \Sigma^\dagger) + \tilde{c} \text{Tr}(\Sigma \Sigma^\dagger)^2 + (\tilde{m} \det(\Sigma) + h.c.)$$

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Breaking to the diagonal $SU(3)$ in a large fraction of parameter space (but no (C)P breaking)

[Bai/Dobrescu '17]



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[Weinberg '76, Susskind '78]

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|-----------|--------------|--------------|--------------------|---------------------|
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Our parity solution to the strong CP problem

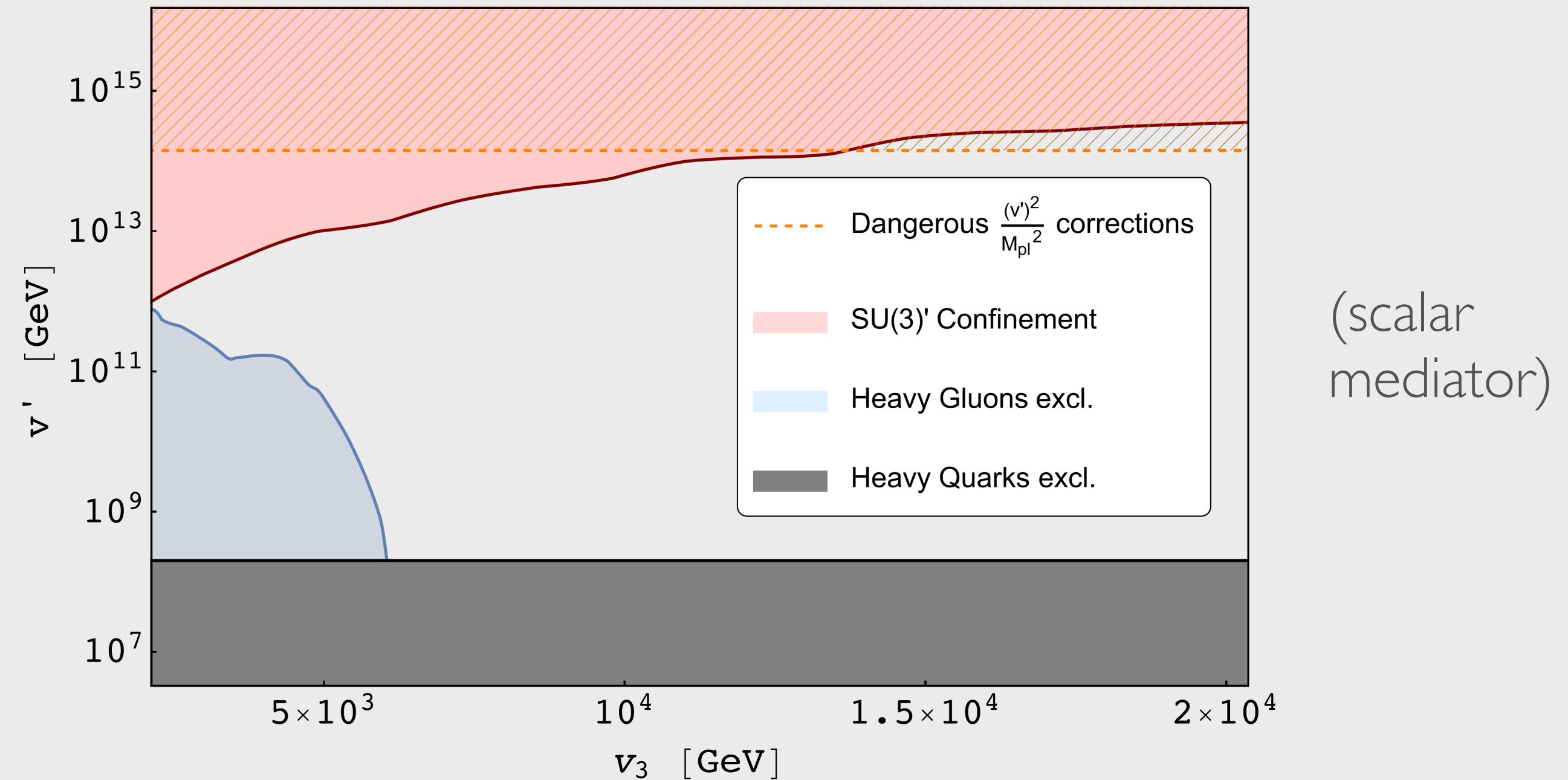
? Need $\bar{\theta} \approx 0$ even **below the scale of parity breaking**

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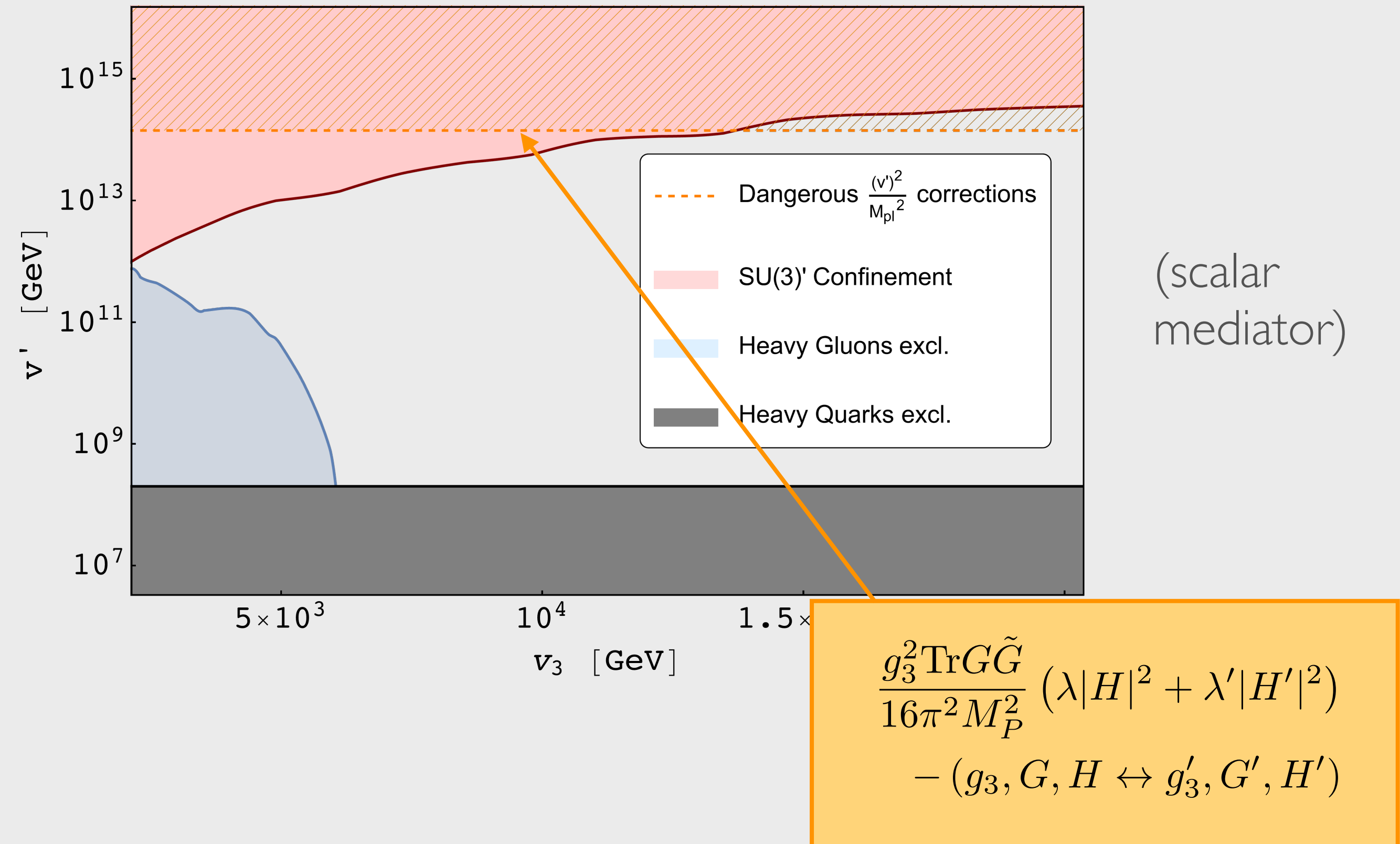
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Only mediators: gluons, bicolored mediator or heavy Higgs. Only CP phase: CKM. **Very small contributions** (at least 3-loops) **to** $\bar{\theta}_{\text{QCD}}$. Effect of small instantons also suppressed

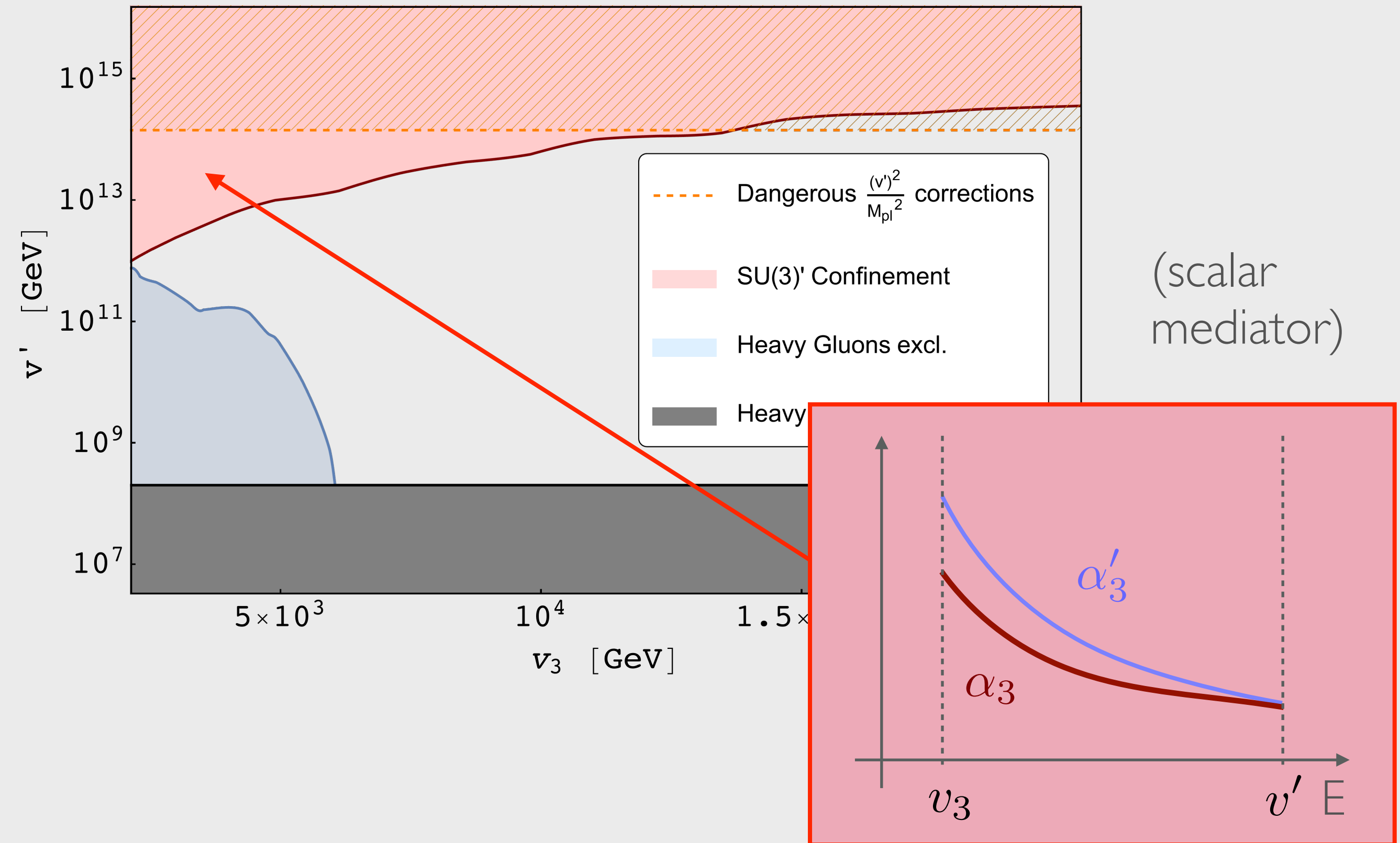
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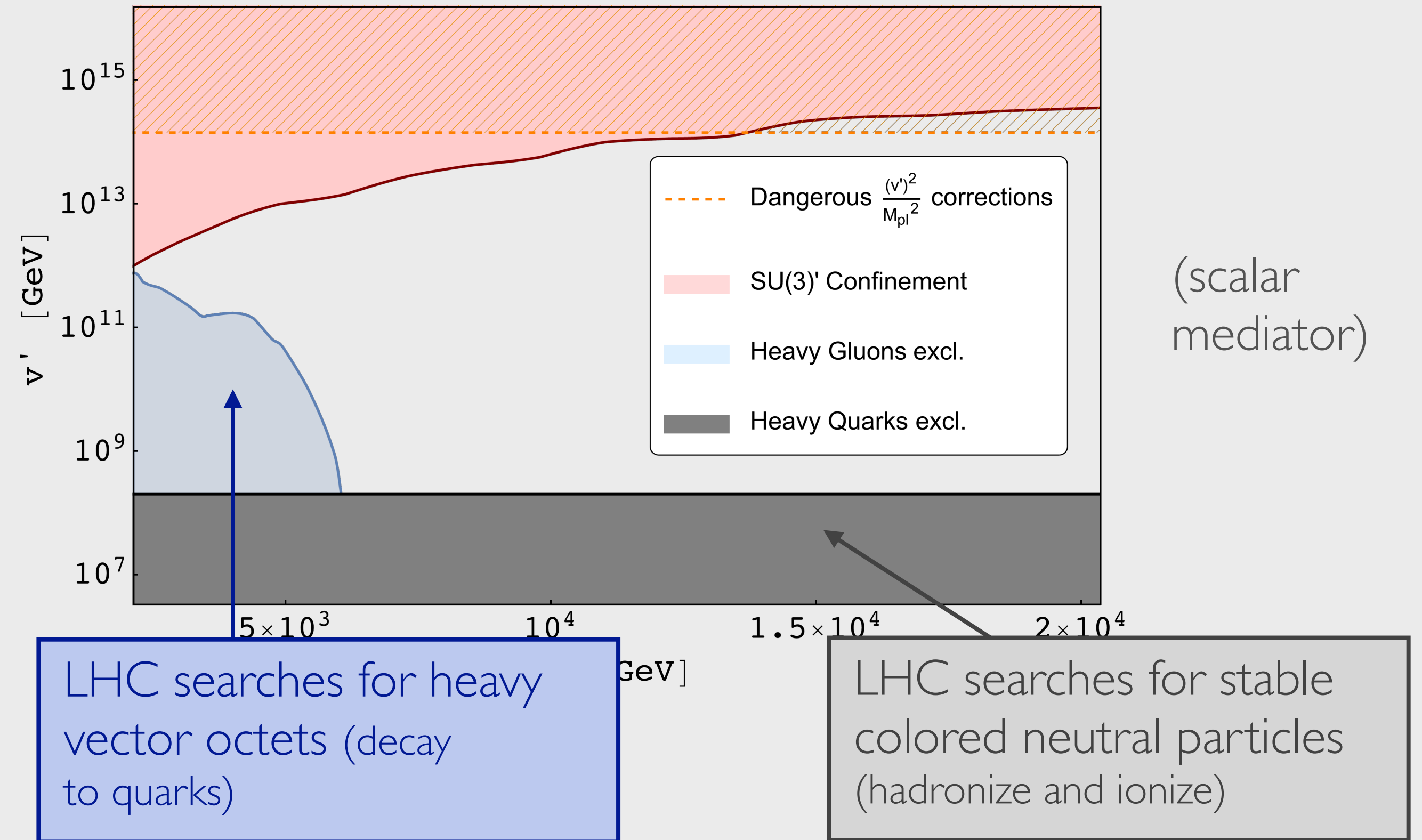
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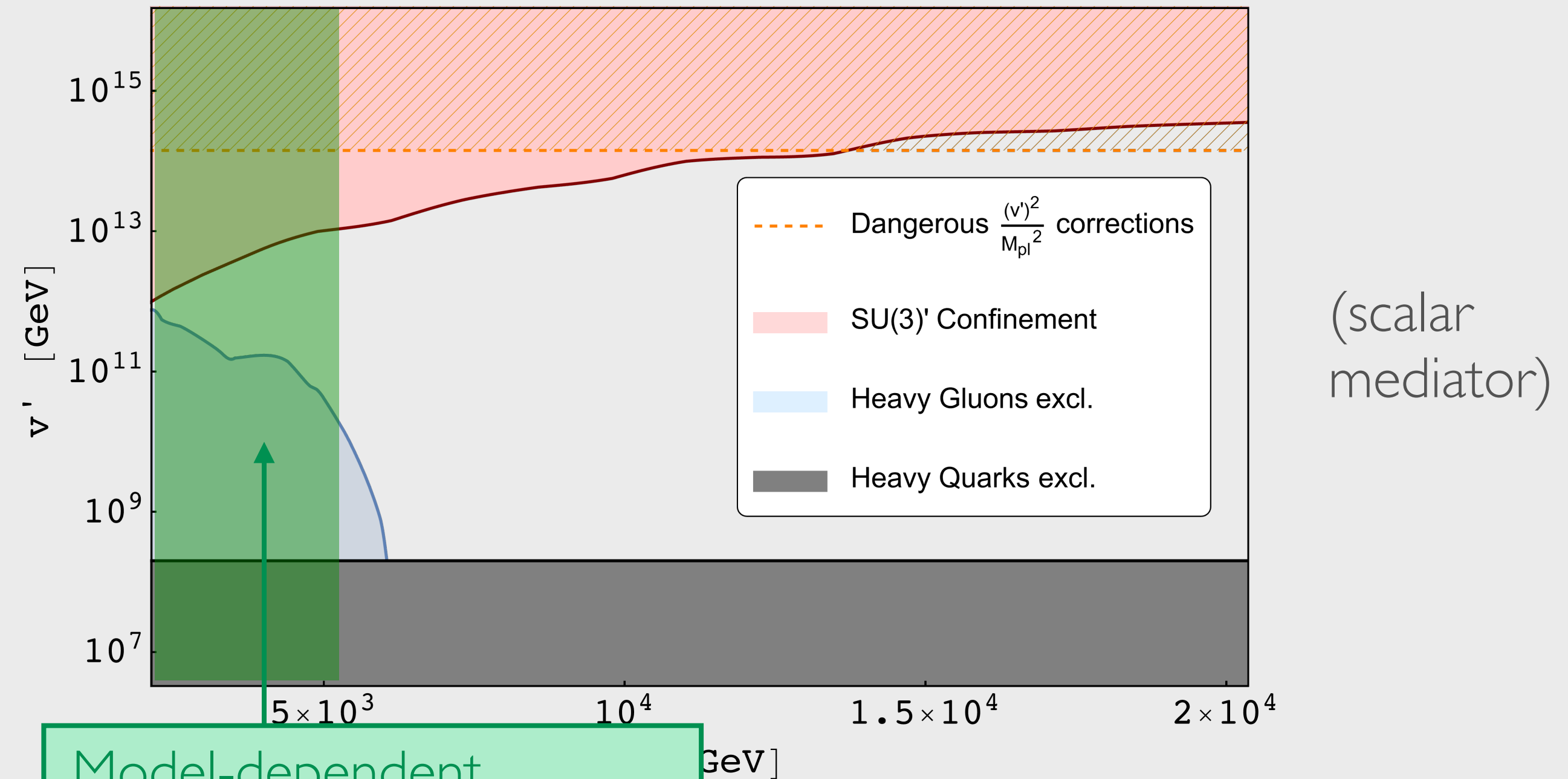
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Model-dependent
bounds on colored
scalars

Dark matter from the mirror world

Dark matter from the mirror world

w.i.p.

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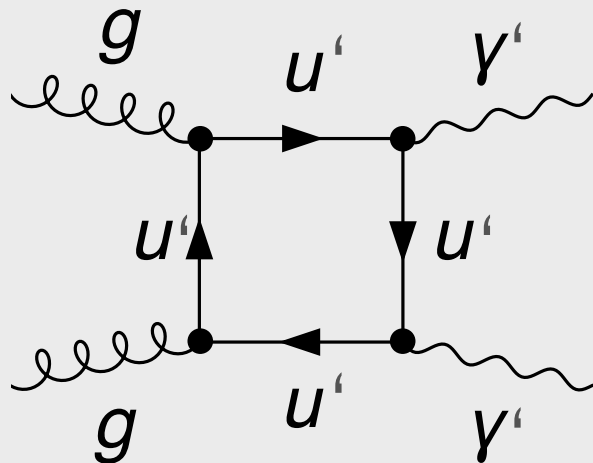
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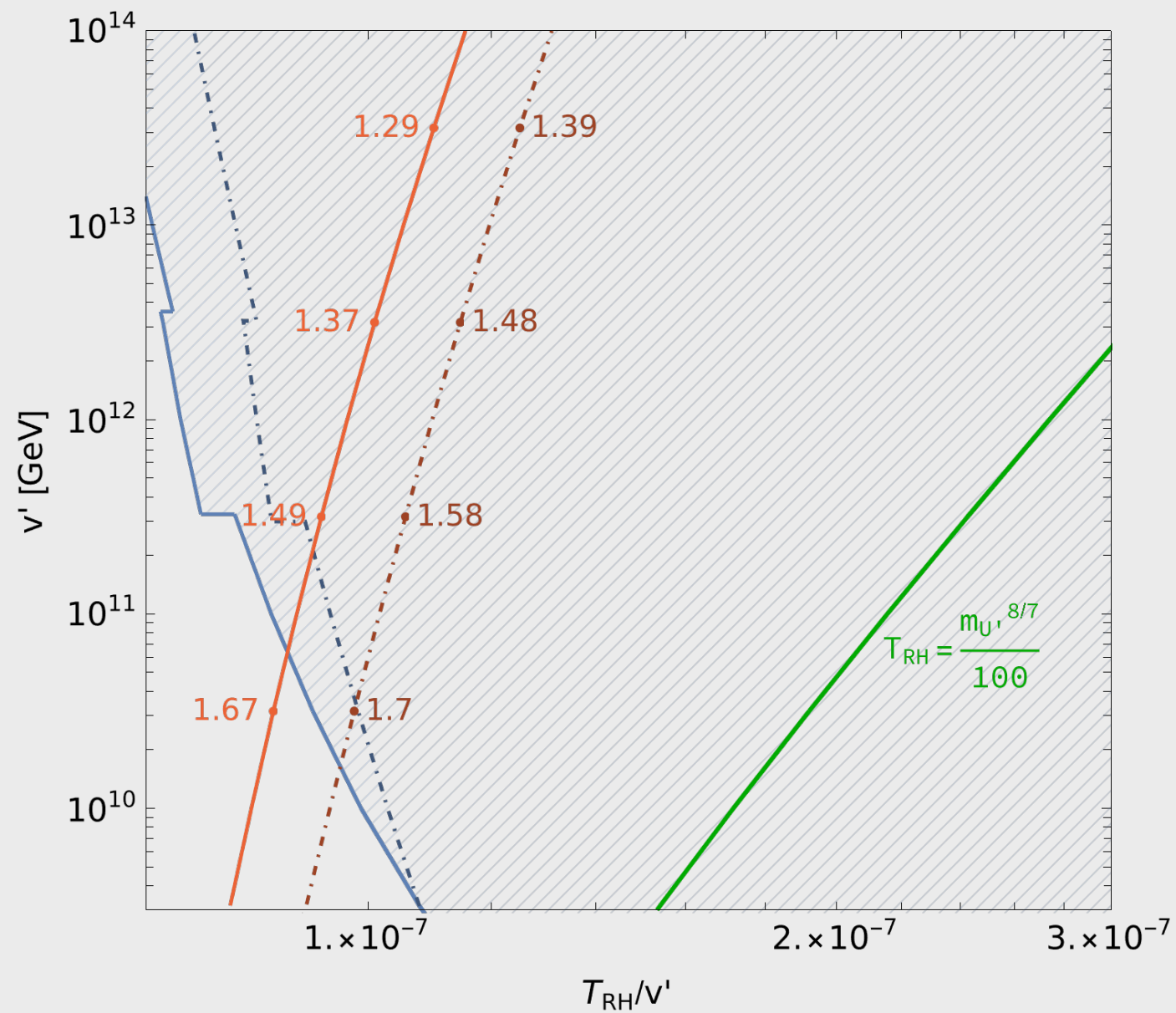
Sequential freeze-in from the mirror photon

[Hambye/Tytgat/Vandecasteele/Vanderheyden '18,
Bélanger/Delaunay/Pukhov/Zaldivar '19]

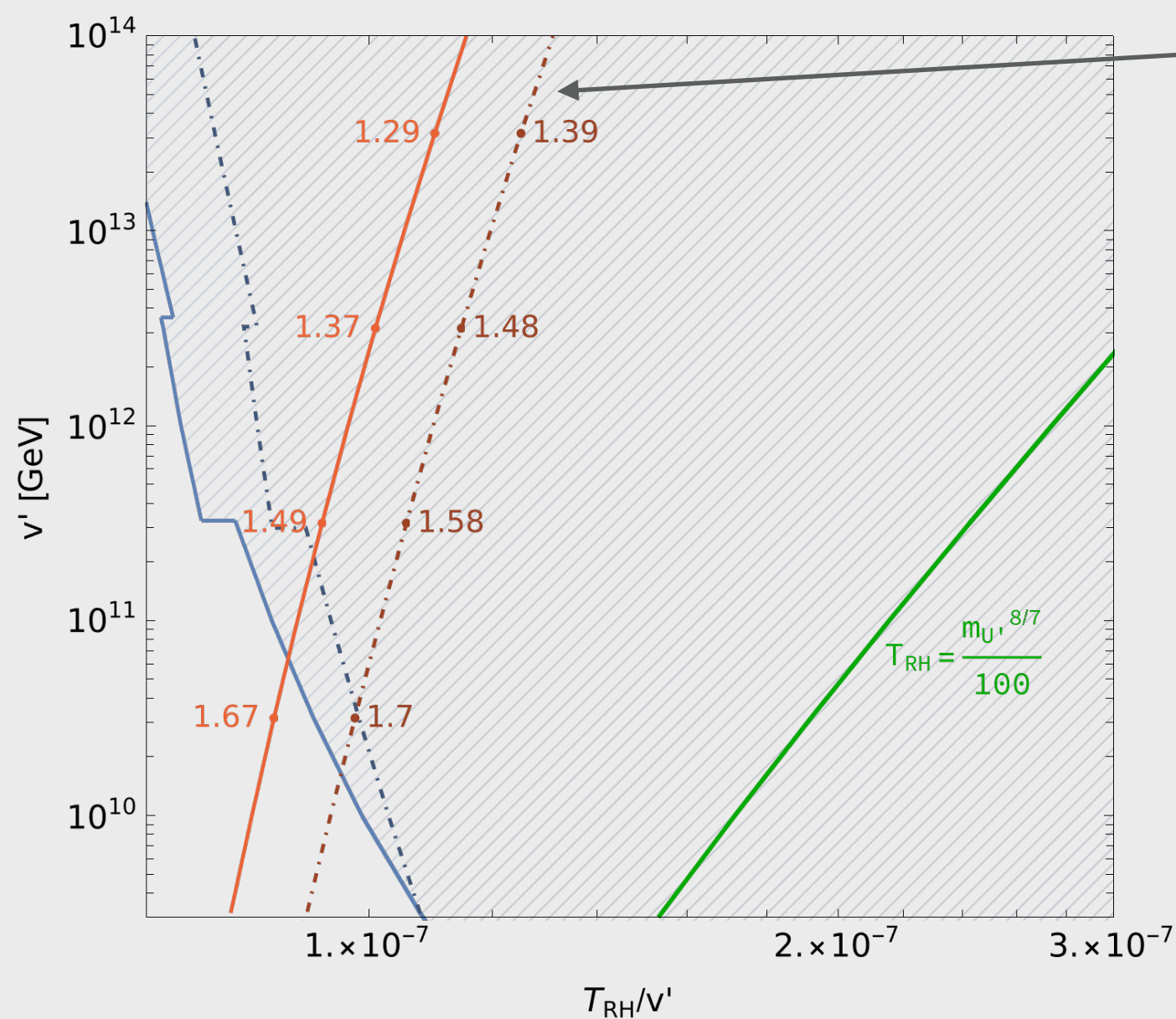


then $\gamma'\gamma' \rightarrow e'\bar{e}'$

Dark matter from the mirror world

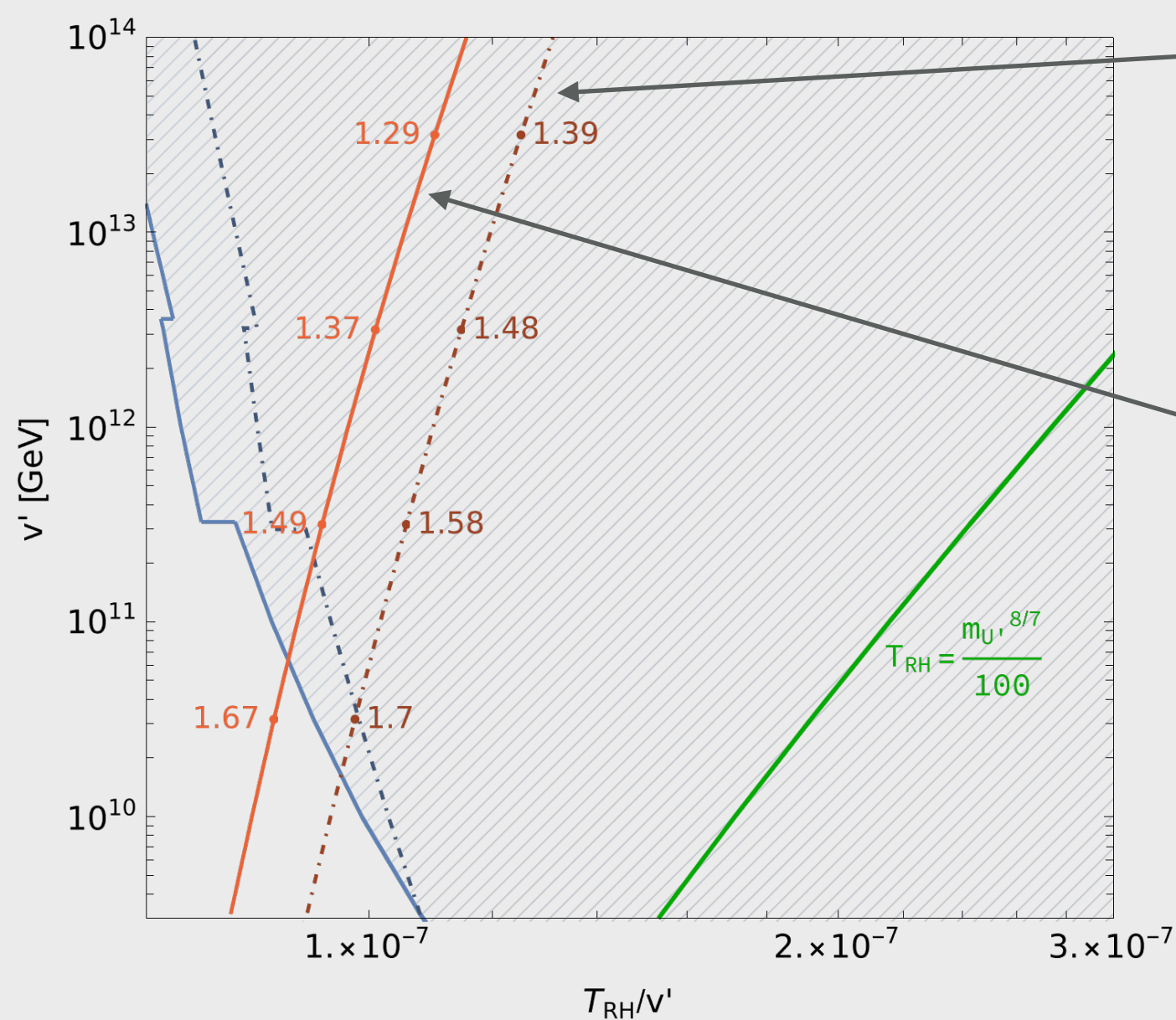


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Successful mirror electron freeze-in with color breaking before parity breaking

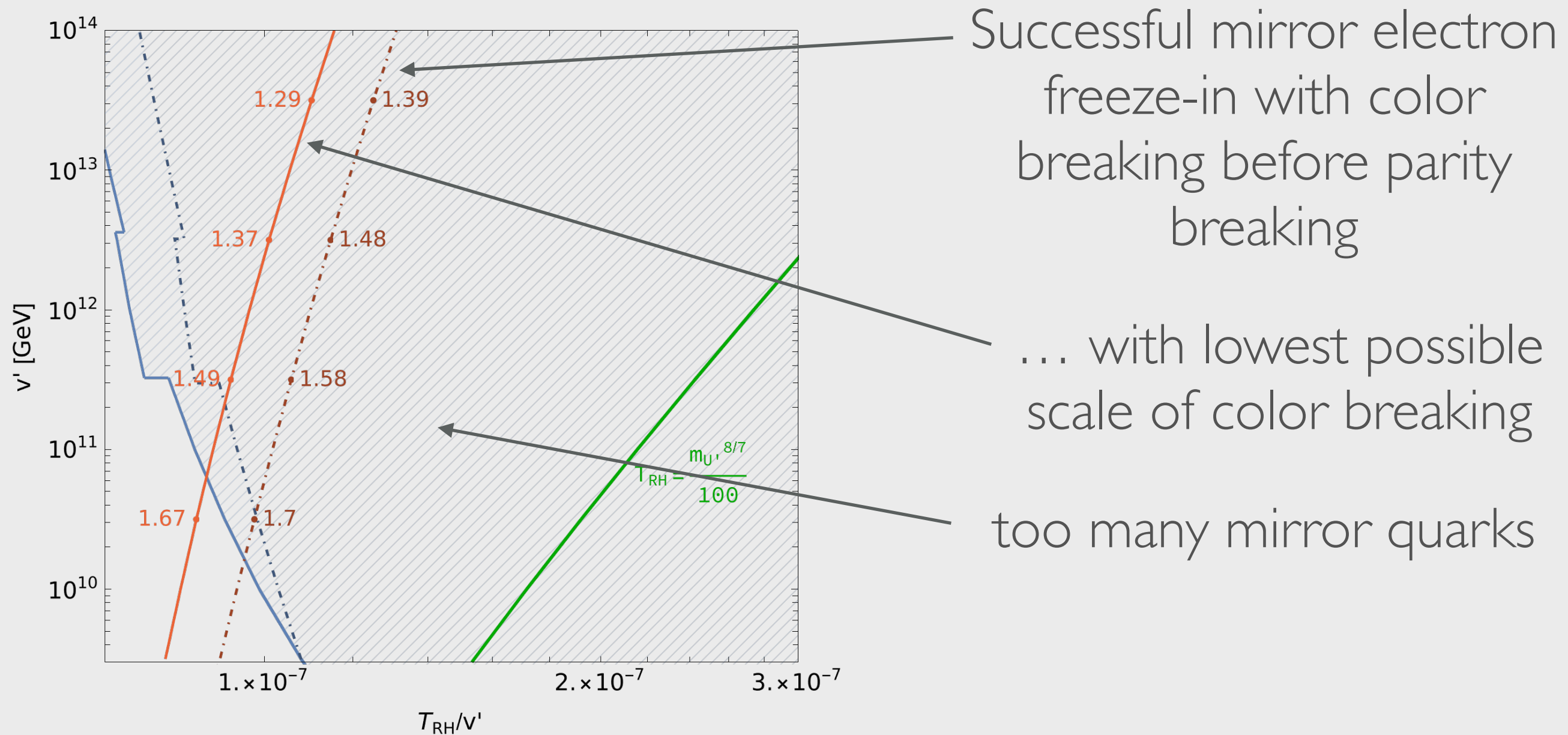
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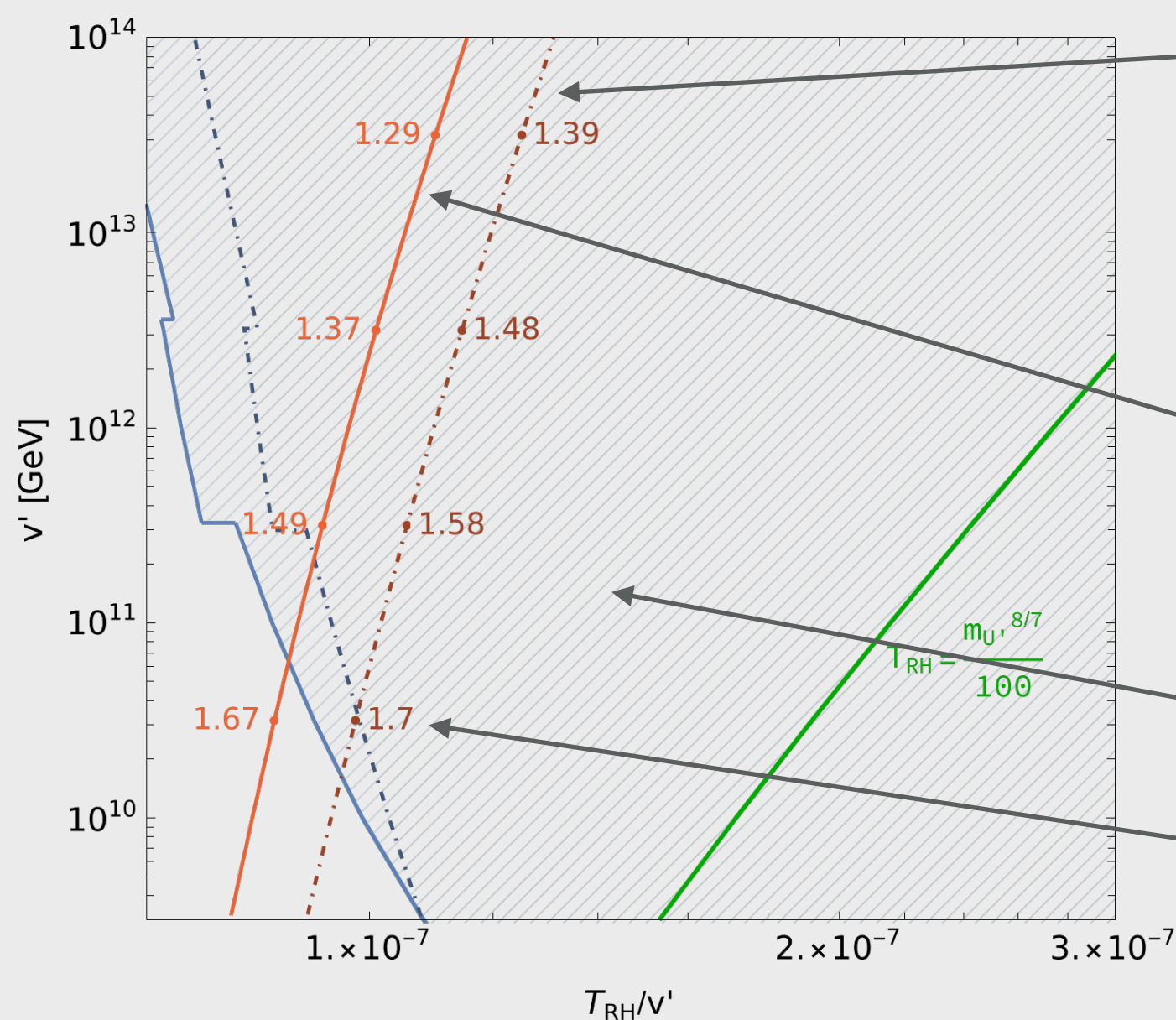
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... with lowest possible scale of color breaking

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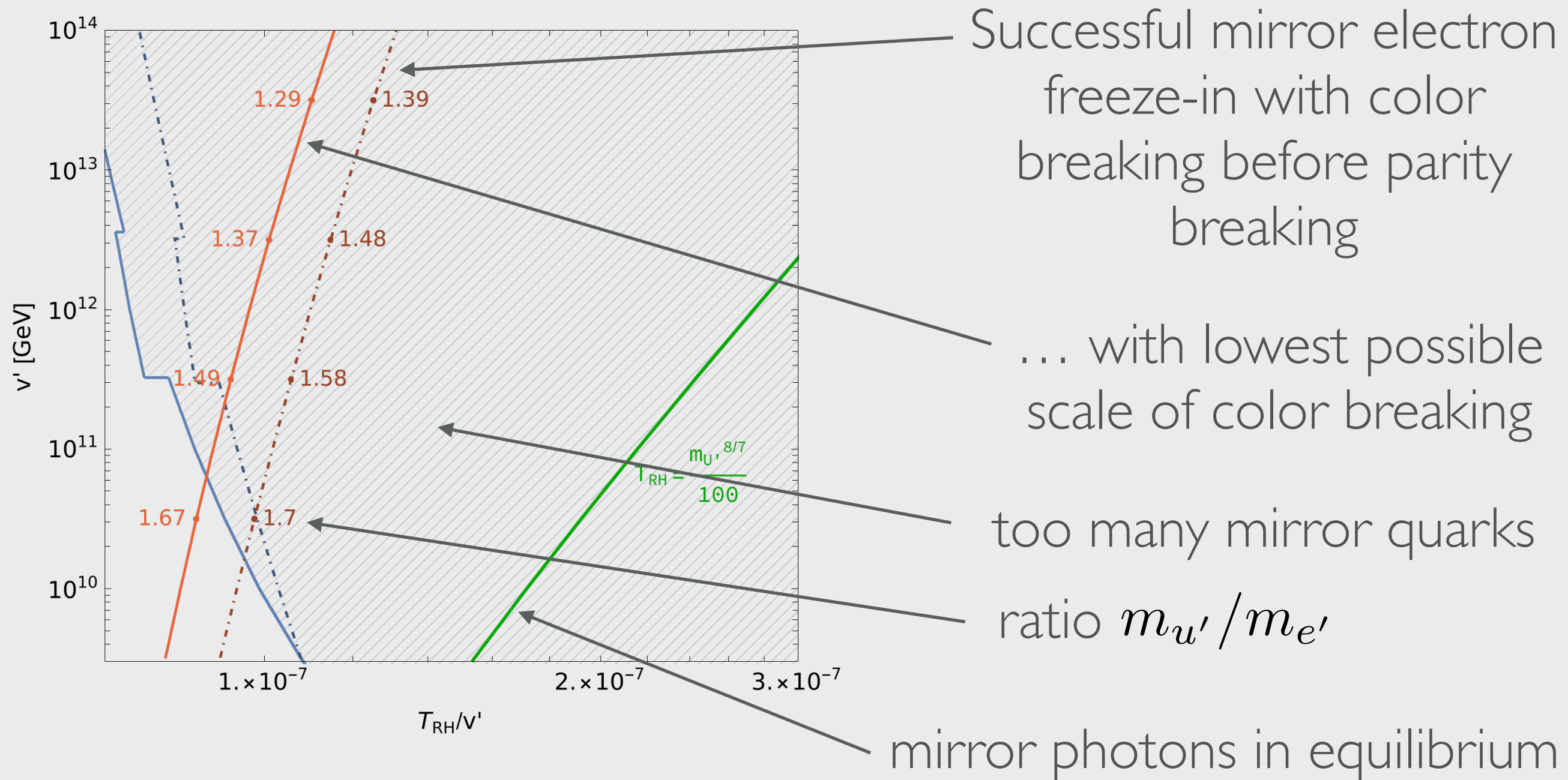
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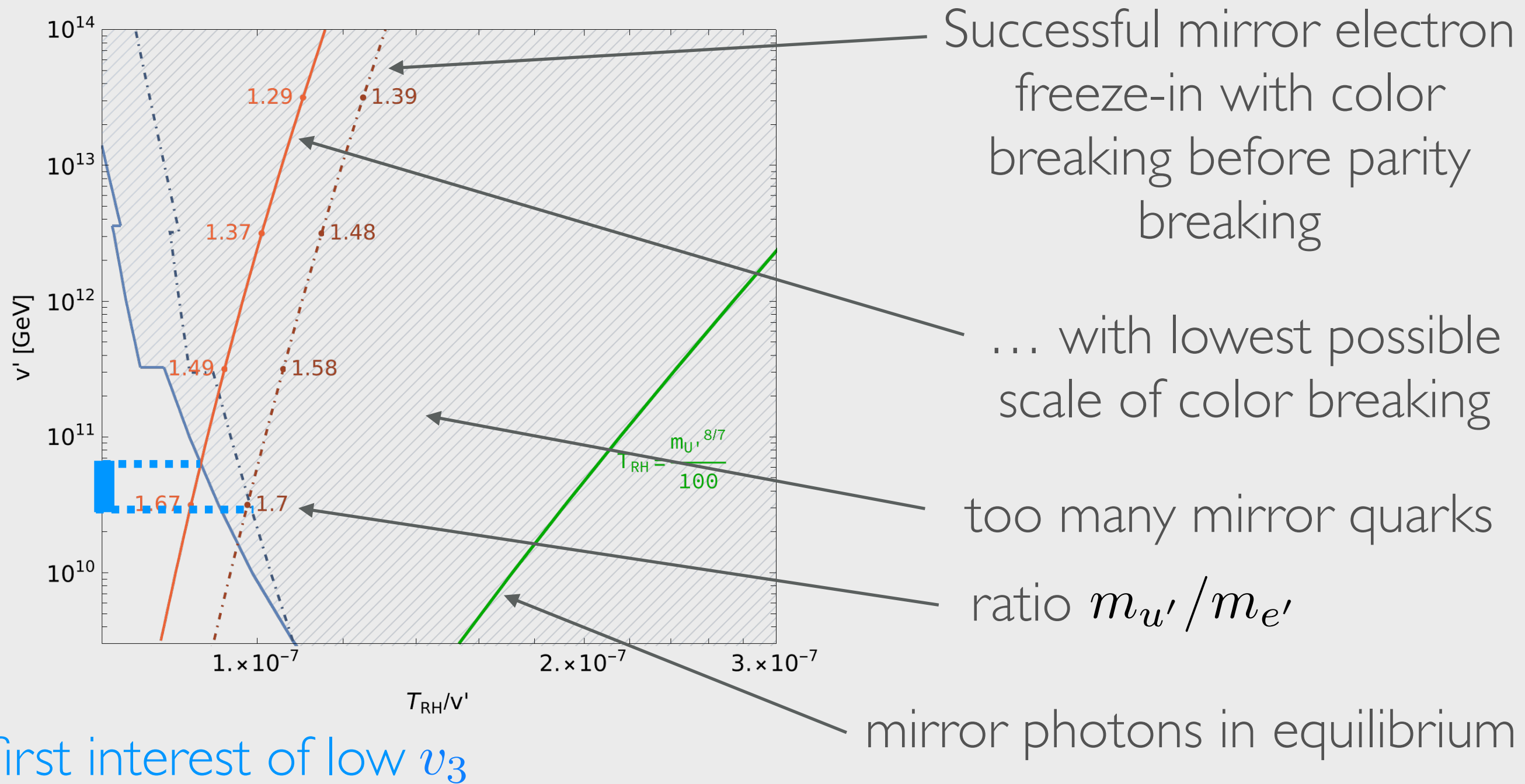
too many mirror quarks

ratio $m_{u'}/m_{e'}$

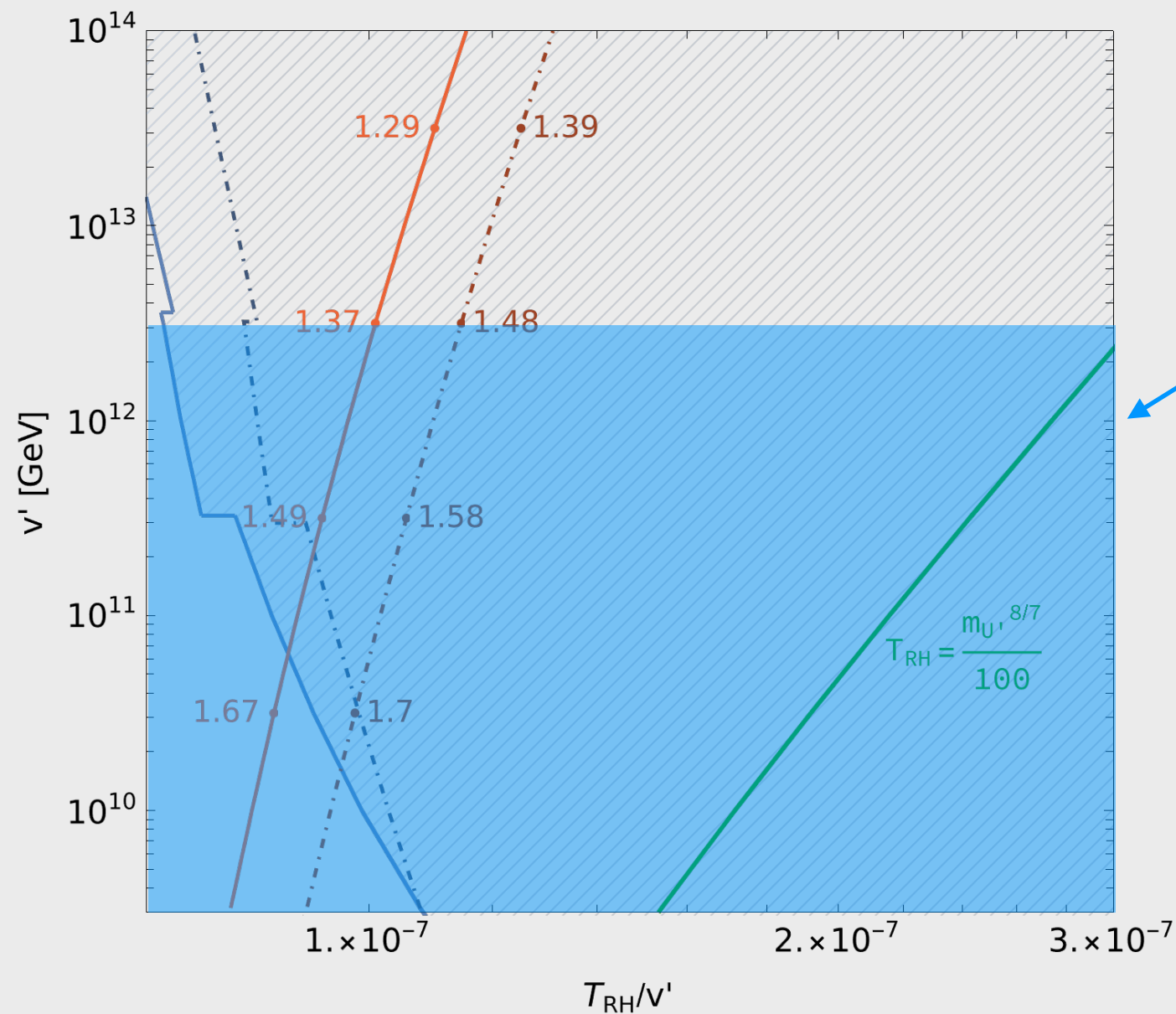
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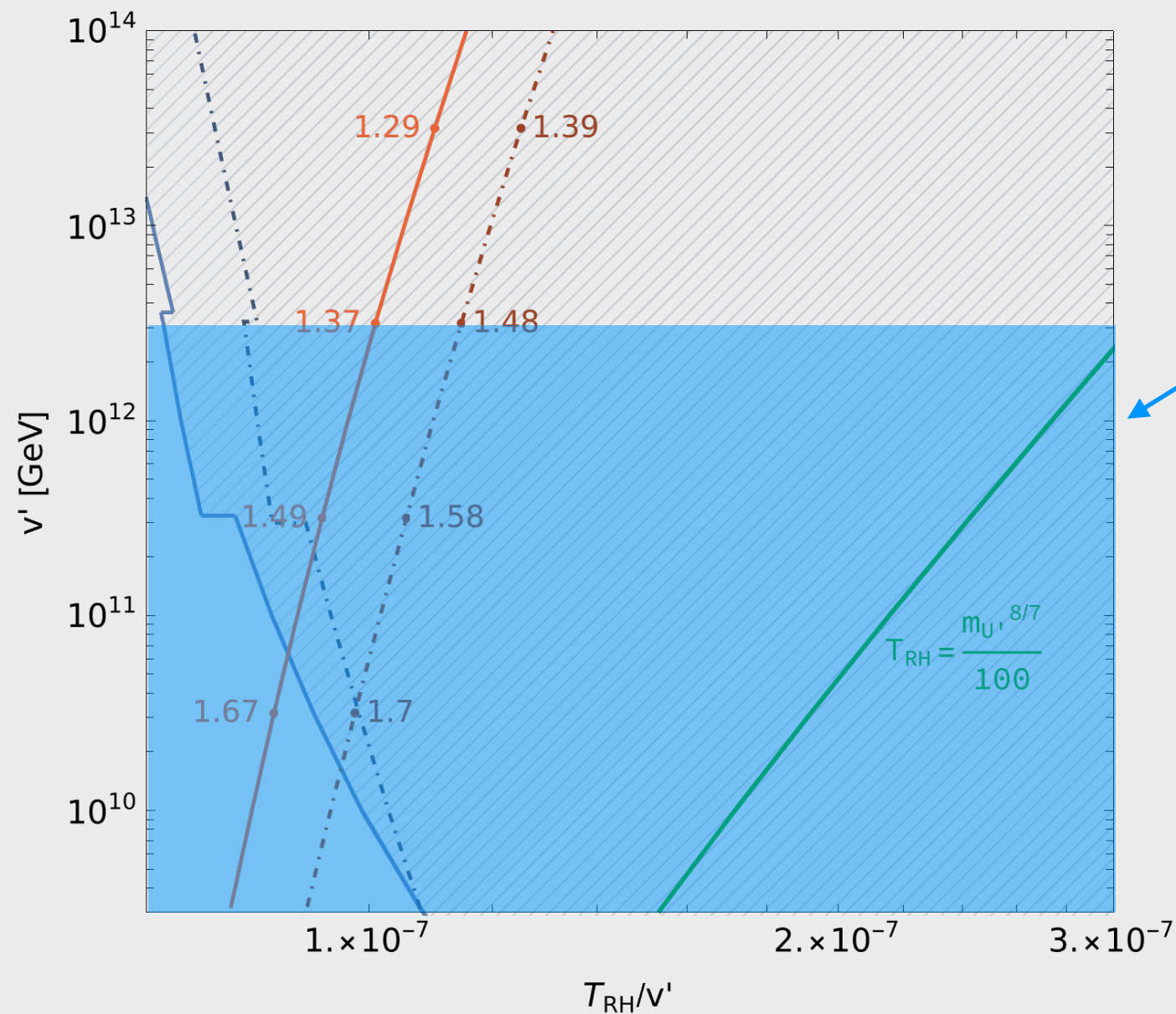
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excluded by irreducible
kinetic mixing and direct
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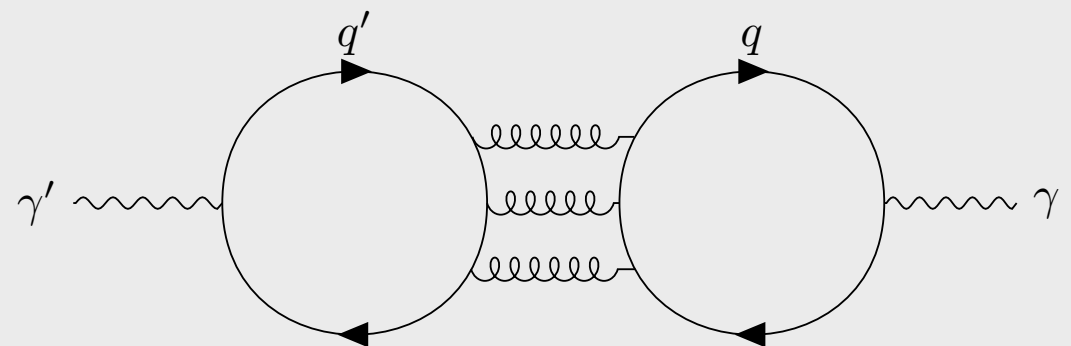
main interest of low v_3

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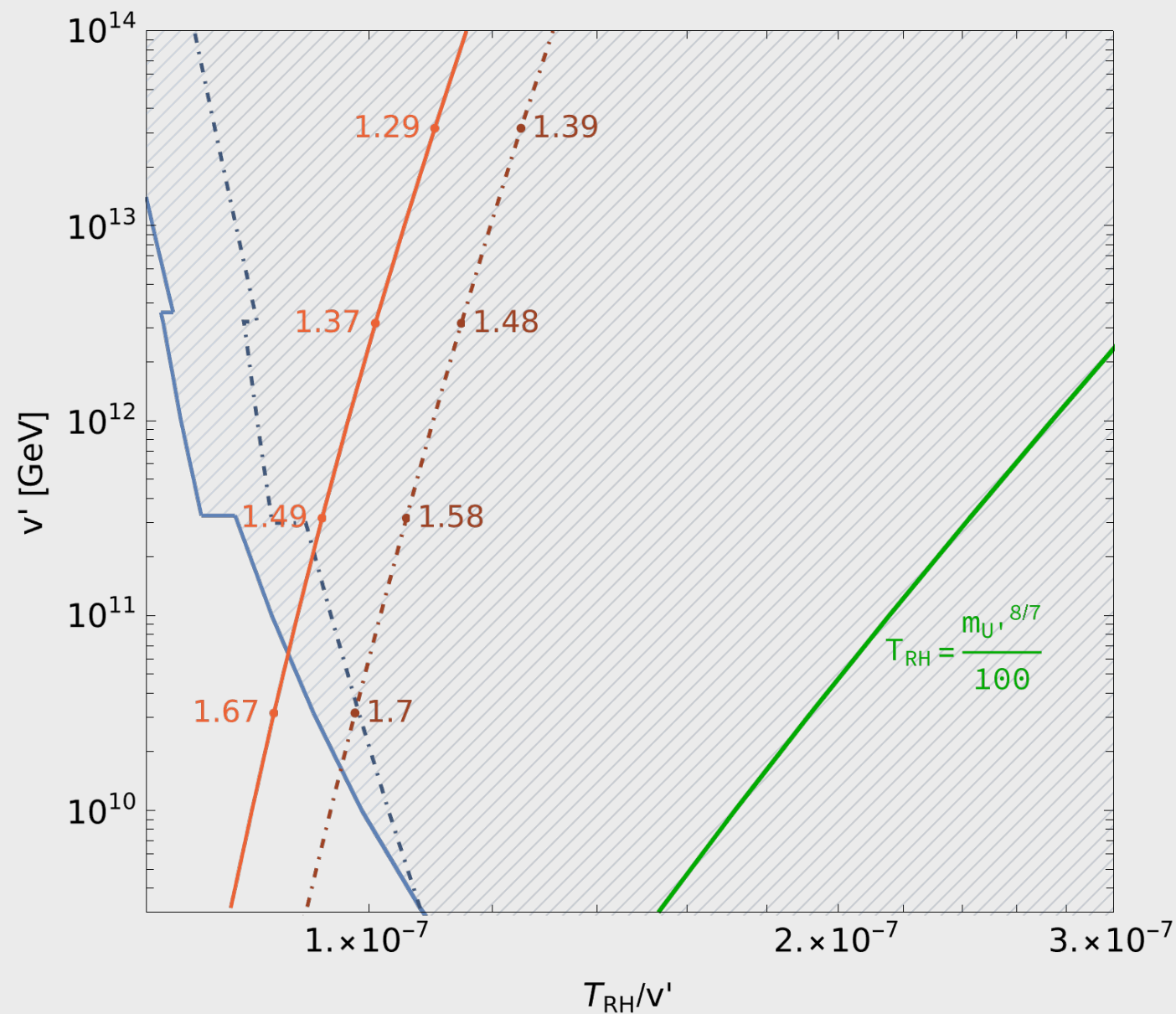
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$$m_{u'} < v_3$$

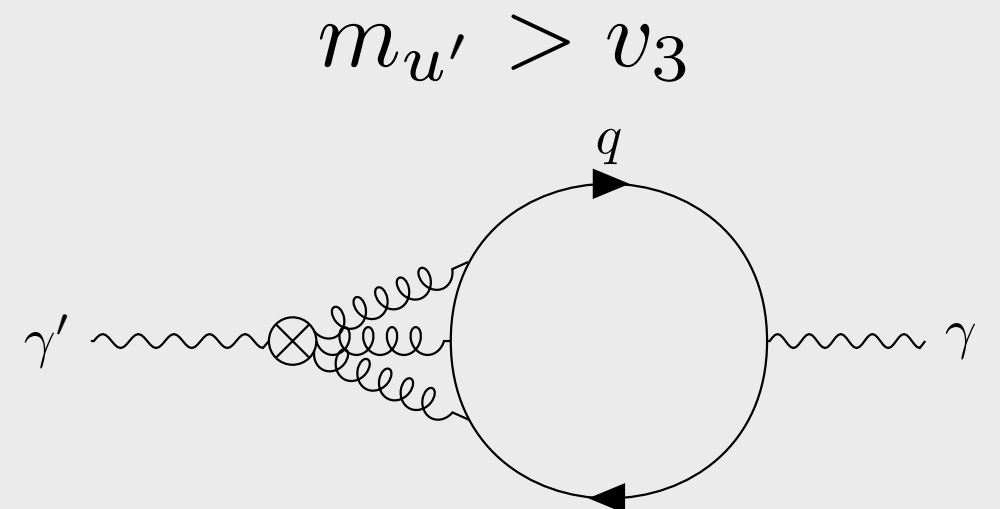
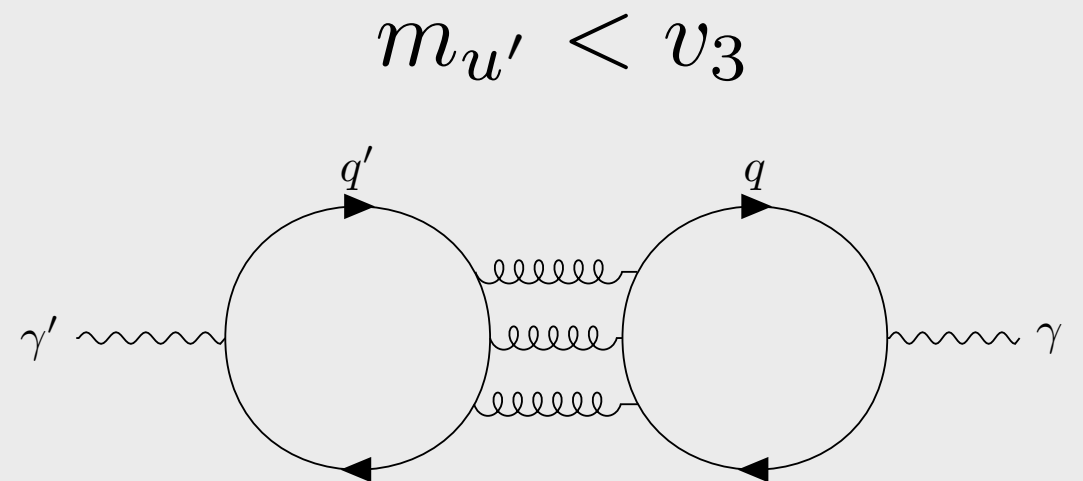


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Outlook

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Solutions to the strong CP problem

Need for UV input for the axion solution :

axion quality problem

[Georgi/Hall/Wise '81, Barr/Seckel '92,
Holman/Hsu/Kolb/Watkins/Widrow '92,
Kamionkowski/March-Russell '92]

No global symmetries :

$$\mathcal{L}_{PQ} = \frac{\delta m_a^2}{2} a^2 \quad \text{with} \quad \delta m_a^2 \xrightarrow{M_P \rightarrow \infty} 0$$

Need

$$\frac{\delta m_a^2}{m_a^2} \lesssim 10^{-10}$$

$$\mathcal{L}_{PQ} = \frac{c_{PQ}}{M_P^{n-4}} \phi^n + h.c.$$

$\phi \sim f_a e^{ia/f_a}$

could be
 $M_P^m \Lambda_h^{n-m-4}$

[QB '22]

Need to understand physics all the way to M_P
... without IR impact

An effective field theorist's nightmare (bis)