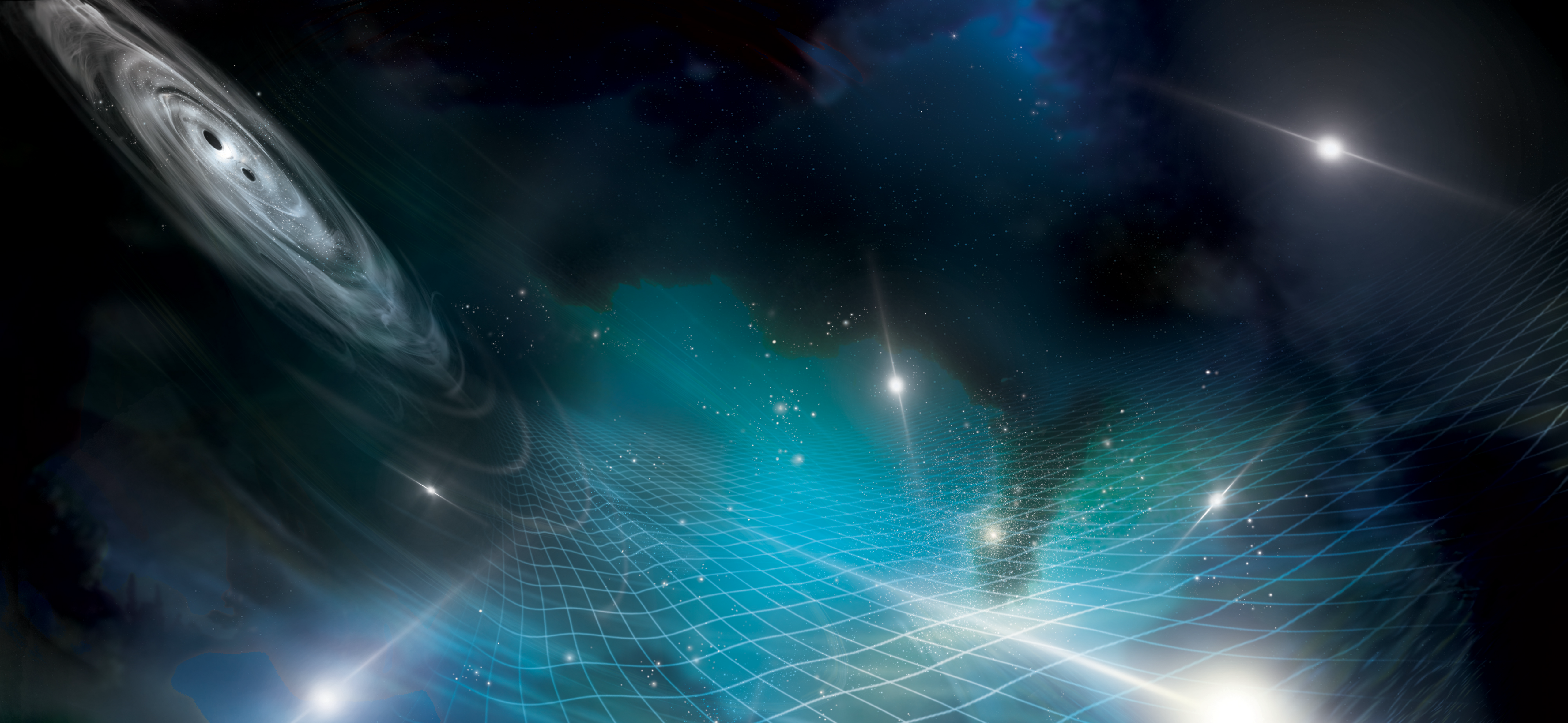




Opening a new window into the gravitational-wave Universe with PTAs

Andrea Mitridate



SOME ASTRONOMY MILESTONES

first documented observations
1800 - 1600 BC



telescope is invented
(1608)

confirmation of heliocentric theory

Kepler's laws (1609 - 1619)

Newton's laws (1687)

...

Hubble discovery of Universe expansion (1929)

first radio telescope
(1937)

CMB discovery (1964)

first binary pulsar is discovered (1974)

first millisecond pulsar is discovered (1983)



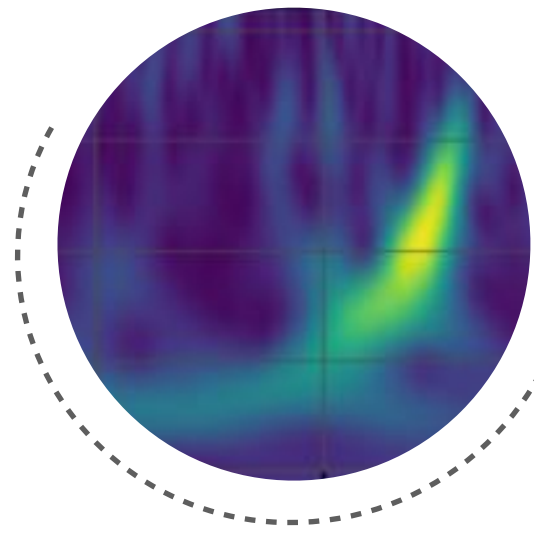
first space-based telescope
(1968)

full access to the electromagnetic spectrum

test of DM self-interactions (bullet cluster)

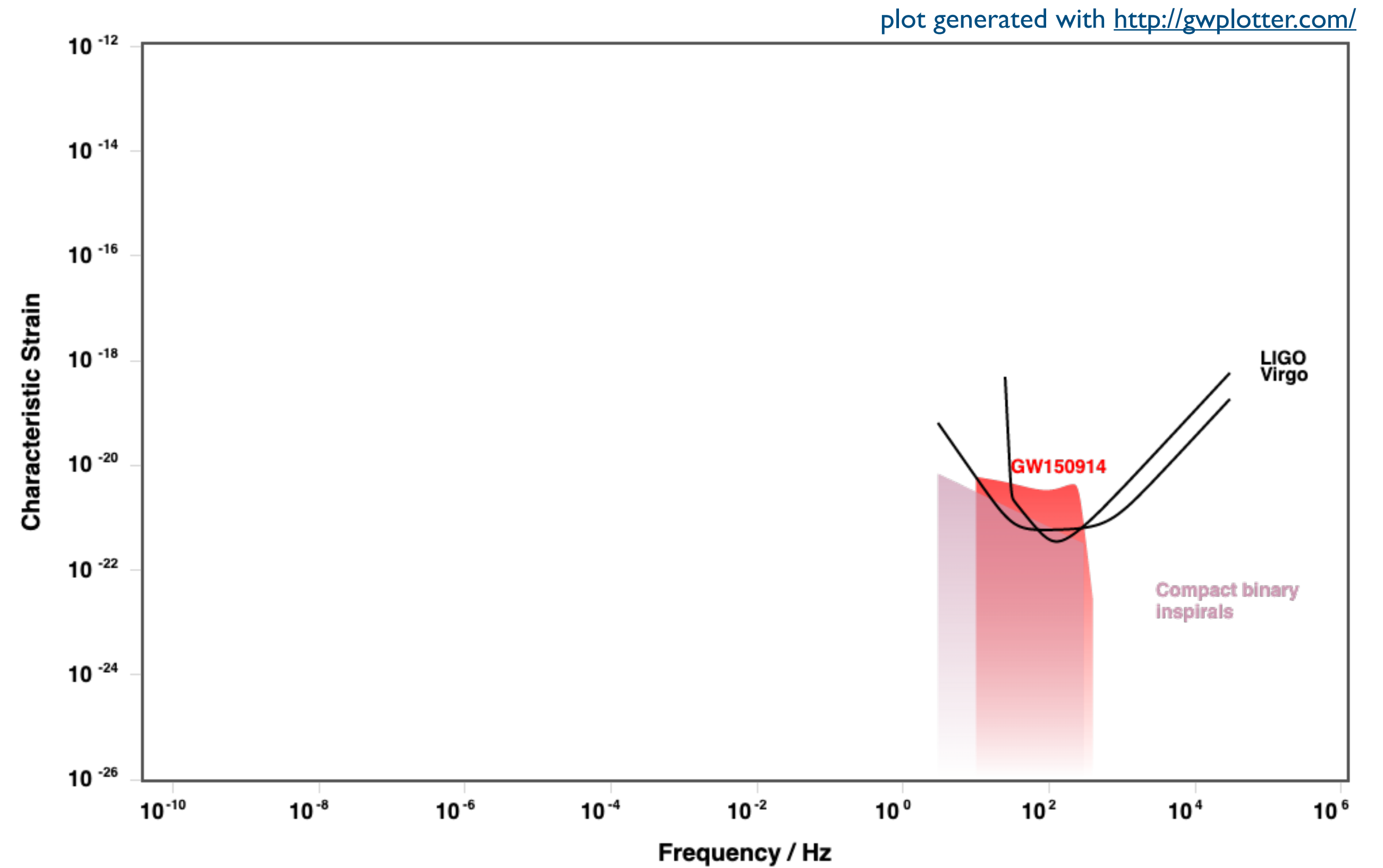
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THE BEGINNING OF A NEW JOURNEY

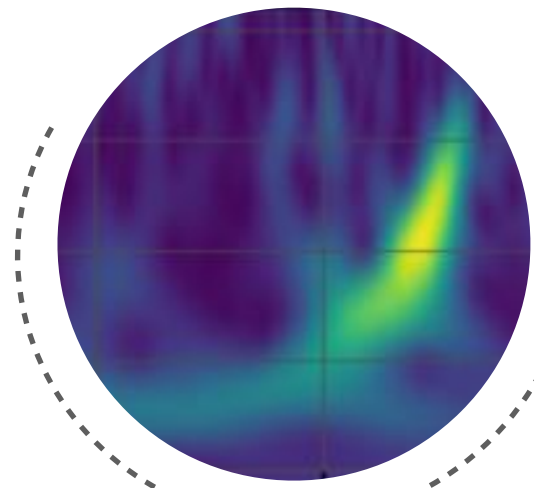


first GW observation (2015)

?

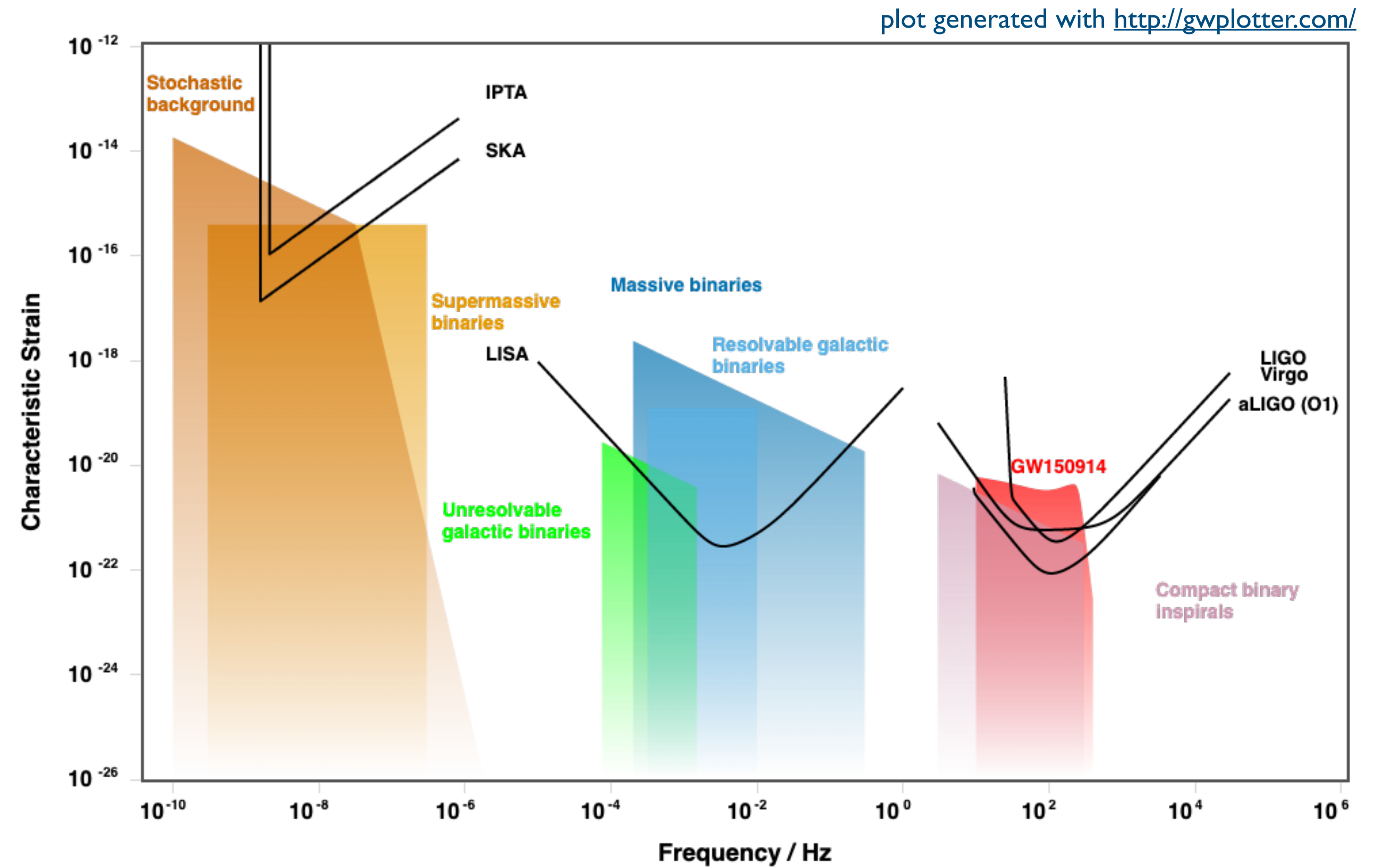


THE BEGINNING OF A NEW JOURNEY



first GW observation (2015)

?



PULSARS

**Rotation
Axis**



**Magnetic
Field Axis**

Radiation Beams

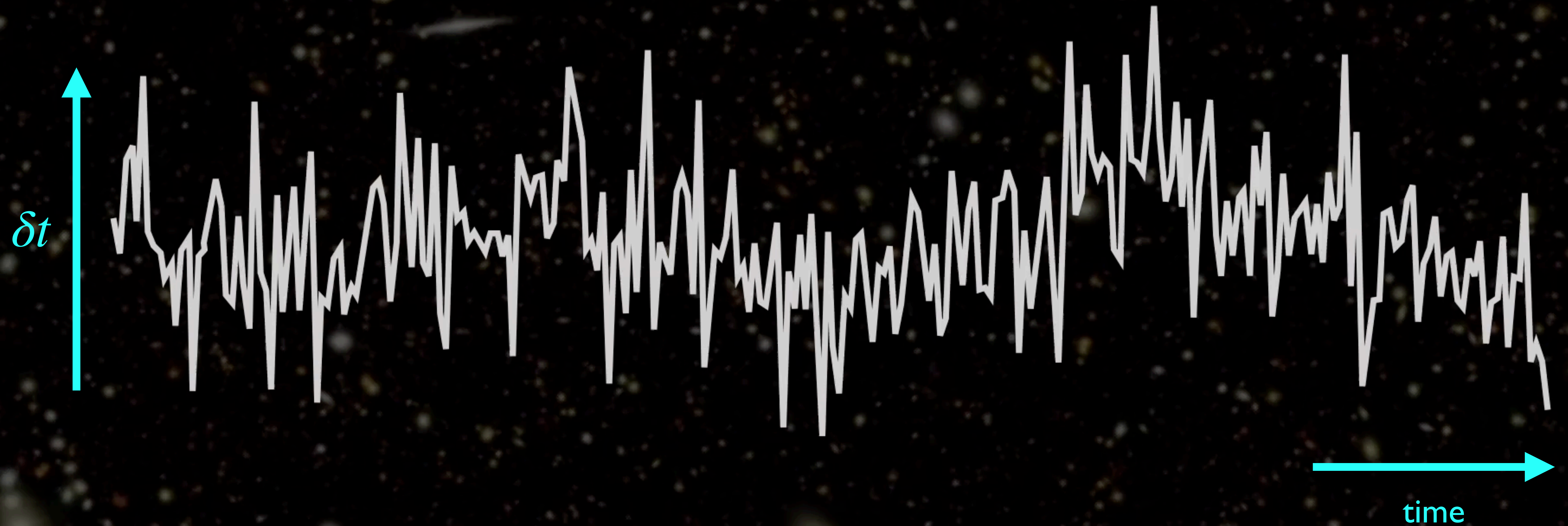
PULSARS AS CLOCKS



A GALAXY-SIZED DETECTOR FOR GW



NOISE vs SIGNAL



NOISE vs SIGNAL

NOISE

pulsar dependent
spatially uncorrelated

SIGNAL

common across pulsars
spatially **correlated**

The NANOGrav 12.5-year Data Set: Search For An Isotropic Stochastic Gravitational-Wave Background

ZAVEN ARZOUMANIAN,¹ PAUL T. BAKER,² HARSHA BLUMER,^{3,4} BENECÉ BÉCSY,⁵ ADAM BRAZIER,^{6,7} PAUL R. BROOK,^{3,4}
 SARAH BURKE-SPOLAOR,^{3,4,8} SHAMI CHATTERJEE,⁶ SIYUAN CHEN,^{9,10,11} JAMES M. CORDES,⁶ NEIL J. CORNISH,⁵
 FRONEFIELD CRAWFORD,¹² H. THANKFUL CROMARTIE,⁶ MEGAN E. DECESAR,^{13,14,*} PAUL B. DEMOREST,¹⁵
 TIMOTHY DOLCH,¹⁶ JUSTIN A. ELLIS,¹⁷ ELIZABETH C. FERRARA,¹⁸ WILLIAM FIORE,^{3,4} EMMANUEL FONSECA,¹⁹
 NATHAN GARVER-DANIELS,^{3,4} PETER A. GENTILE,^{3,4} DEBORAH C. GOOD,²⁰ JEFFREY S. HAZBOUN,^{21,*}
 A. MIGUEL HOLGADO,^{22,23} KRISTINA ISLO,²⁴ ROSS J. JENNINGS,⁶ MEGAN L. JONES,²⁴ ANDREW R. KAISER,^{3,4}
 DAVID L. KAPLAN,²⁴ LUKE ZOLTAN KELLEY,²⁵ JOEY SHAPIRO KEY,²¹ NIMA LAAL,²⁶ MICHAEL T. LAM,^{27,28}
 T. JOSEPH W. LAZIO,²⁹ DUNCAN R. LORIMER,^{3,4} JING LUO,³⁰ RYAN S. LYNCH,³¹ DUSTIN R. MADISON,^{3,4,*}
 MAURA A. MCLAUGHLIN,^{3,4} CHIARA M. F. MINGARELLI,^{32,33} CHERRY NG,³⁴ DAVID J. NICE,¹³
 TIMOTHY T. PENNUCCI,^{35,36,*} NIHAN S. POL,^{3,4,37} SCOTT M. RANSOM,³⁵ PAUL S. RAY,³⁸ BRENT J. SHAPIRO-ALBERT,^{3,4}
 XAVIER SIEMENS,^{26,24} JOSEPH SIMON,^{29,39} RENÉE SPIEWAK,⁴⁰ INGRID H. STAIRS,²⁰ DANIEL R. STINEBRING,⁴¹
 KEVIN STOVALL,¹⁵ JERRY P. SUN,²⁶ JOSEPH K. SWIGGUM,^{13,*} STEPHEN R. TAYLOR,³⁷ JACOB E. TURNER,^{3,4}
 MICHELE VALLISNERI,²⁹ SARAH J. VIGELAND,²⁴ CAITLIN A. WITT,^{3,4}

THE NANOGrAV COLLABORATION

ABSTRACT

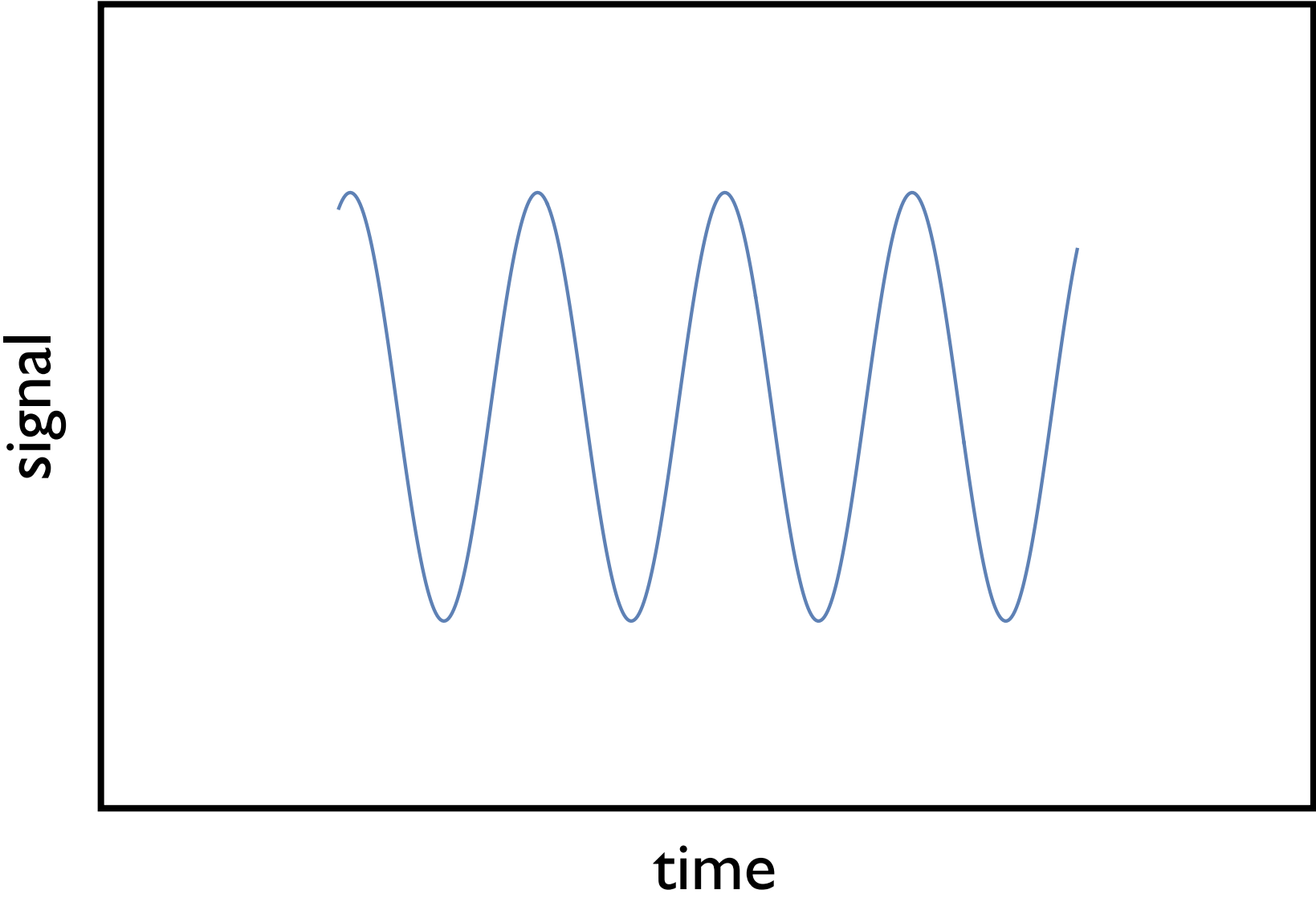
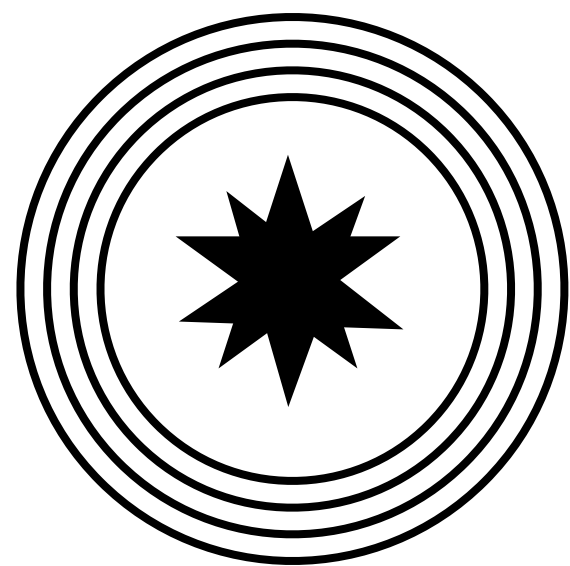
We search for an isotropic stochastic gravitational-wave background (GWB) in the 12.5-year pulsar-timing data set collected by the North American Nanohertz Observatory for Gravitational Waves. Our analysis finds strong evidence of a stochastic process, modeled as a power-law, with common amplitude and spectral slope across pulsars. Under our fiducial model, the Bayesian posterior of the amplitude for an $f^{-2/3}$ power-law spectrum, expressed as the characteristic GW strain, has median 1.92×10^{-15} and 5%–95% quantiles of $1.37\text{--}2.67 \times 10^{-15}$ at a reference frequency of $f_{\text{yr}} = 1 \text{ yr}^{-1}$; the Bayes factor in favor of the common-spectrum process versus independent red-noise processes in each pulsar exceeds 10,000. However, we find no statistically significant evidence that this process has quadrupolar spatial correlations, which we would consider necessary to claim a GWB detection consistent with general relativity. We find that the process has neither monopolar nor dipolar correlations, which may arise from, for example, reference clock or solar system ephemeris systematics, respectively. The amplitude posterior has significant support above previously reported upper limits; we explain this in terms of the Bayesian priors assumed for intrinsic pulsar red noise. We examine potential implications for the supermassive black hole binary population under the hypothesis that the signal is indeed astrophysical in nature.

EVIDENCE FOR A COMMON SIGNAL

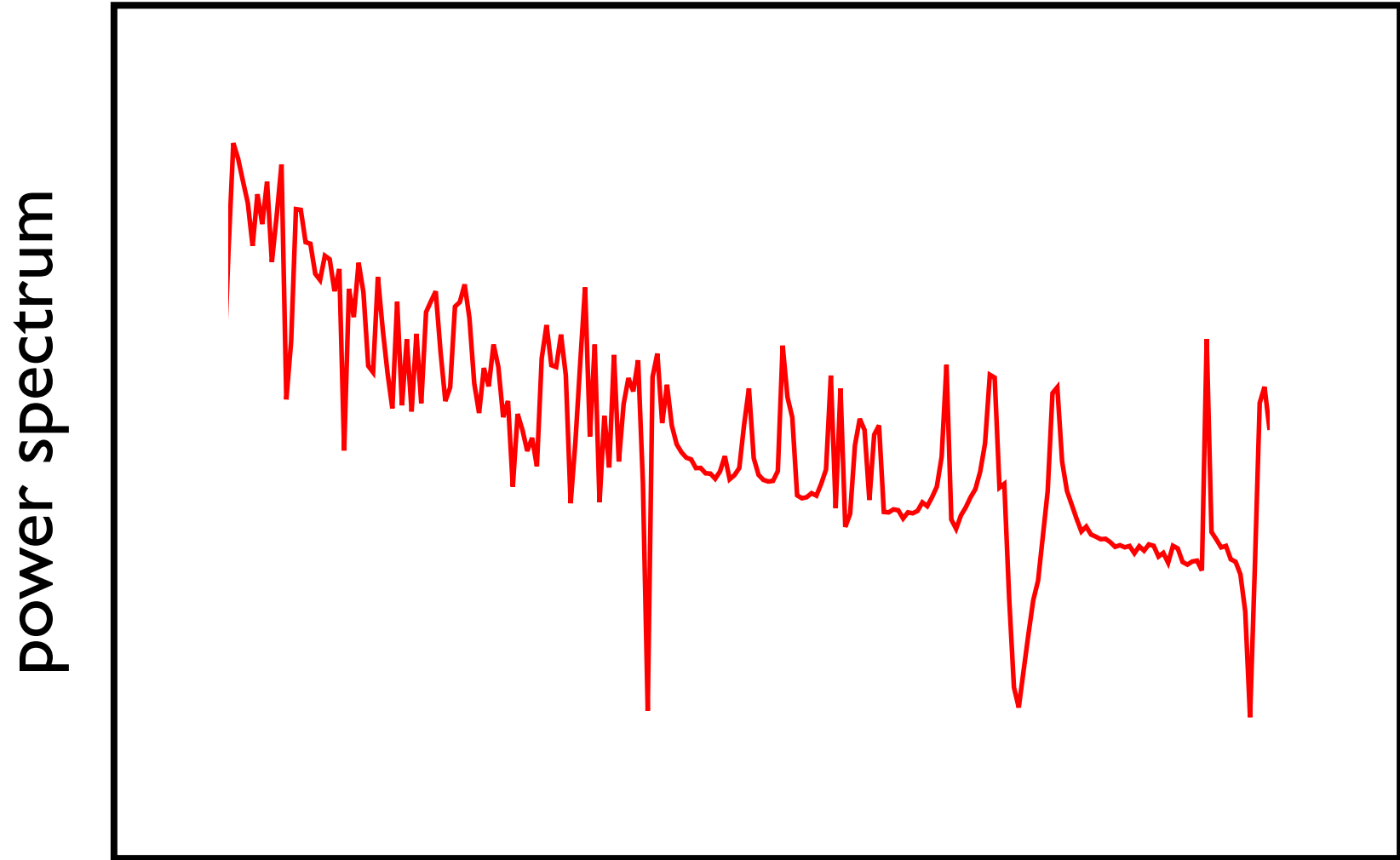
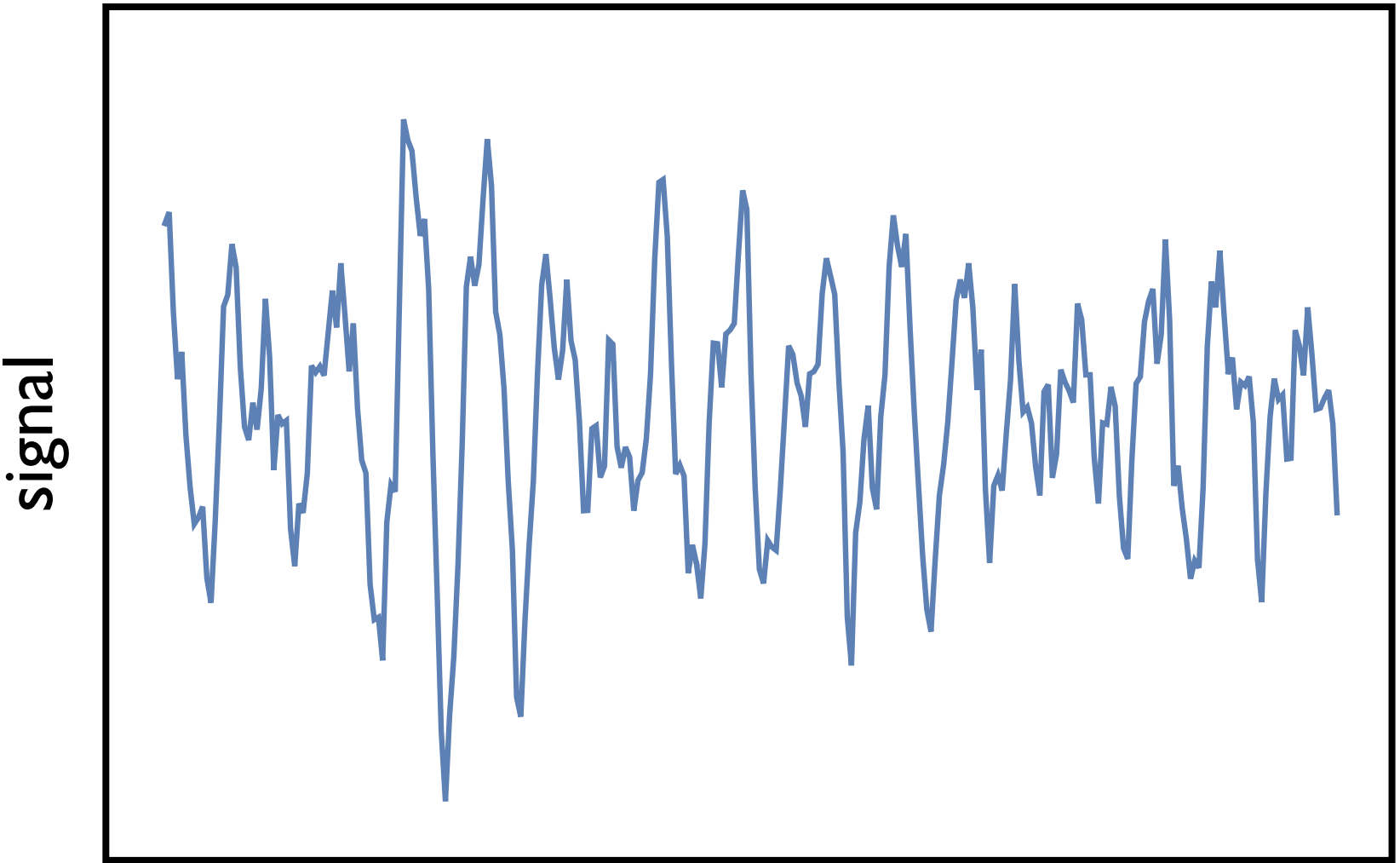
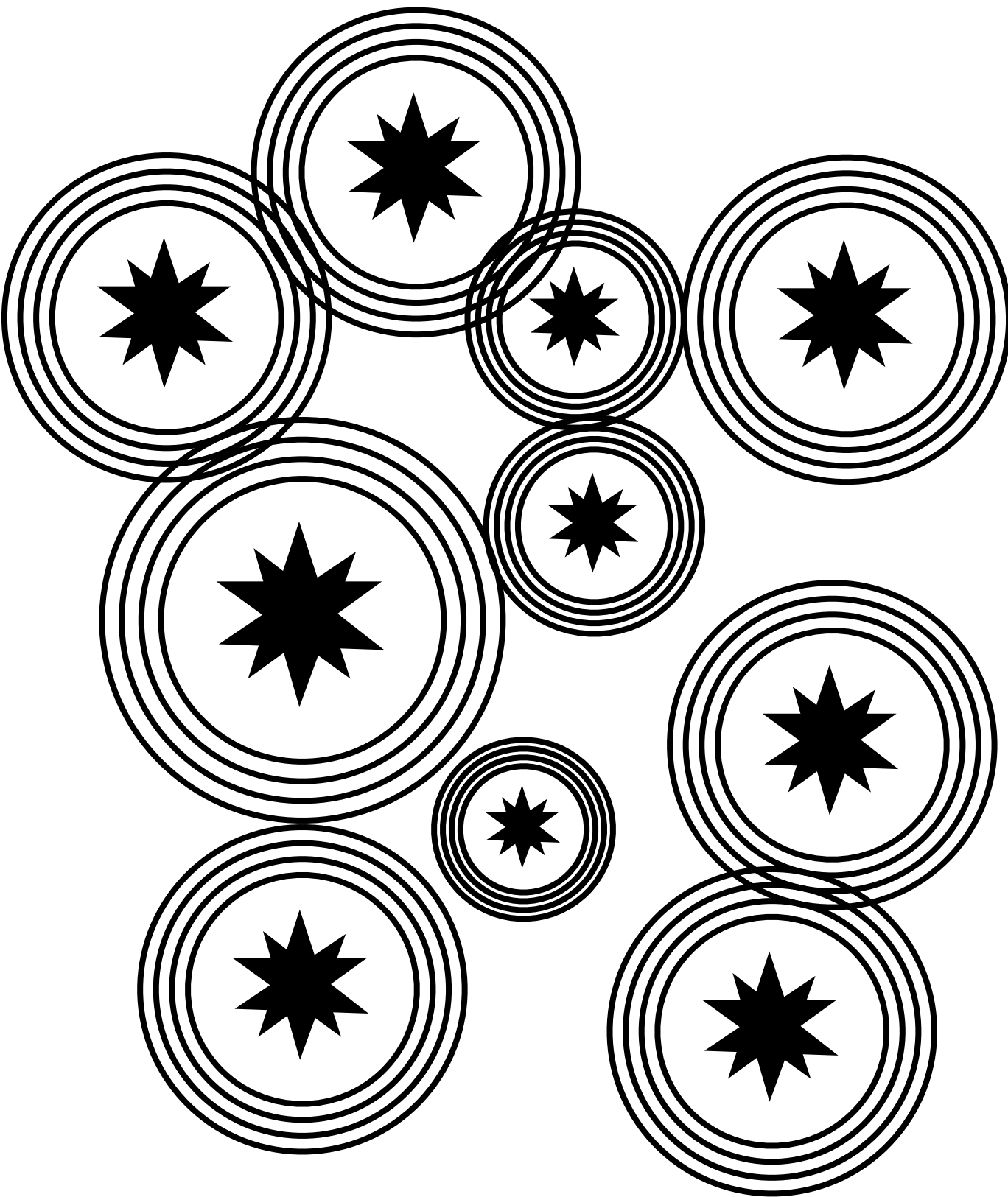
ABSTRACT

We search for an isotropic stochastic gravitational-wave background (GWB) in the 12.5-year pulsar-timing data set collected by the North American Nanohertz Observatory for Gravitational Waves. Our analysis finds strong evidence of a stochastic process, modeled as a power-law, with common amplitude and spectral slope across pulsars. Under our fiducial model, the Bayesian posterior of the amplitude for an $f^{-2/3}$ power-law spectrum, expressed as the characteristic GW strain, has median 1.92×10^{-15} and 5%–95% quantiles of $1.37\text{--}2.67 \times 10^{-15}$ at a reference frequency of $f_{\text{yr}} = 1 \text{ yr}^{-1}$; the Bayes factor in favor of the common-spectrum process versus independent red-noise processes in each pulsar exceeds 10,000. However, we find no statistically significant evidence that this process has quadrupolar spatial correlations, which we would consider necessary to claim a GWB detection consistent with general relativity. We find that the process has neither monopolar nor dipolar correlations, which may arise from, for example, reference clock or solar system ephemeris systematics, respectively. The amplitude posterior has significant support above previously reported upper limits; we explain this in terms of the Bayesian priors assumed for intrinsic pulsar red noise. We examine potential implications for the supermassive black hole binary population under the hypothesis that the signal is indeed astrophysical in nature.

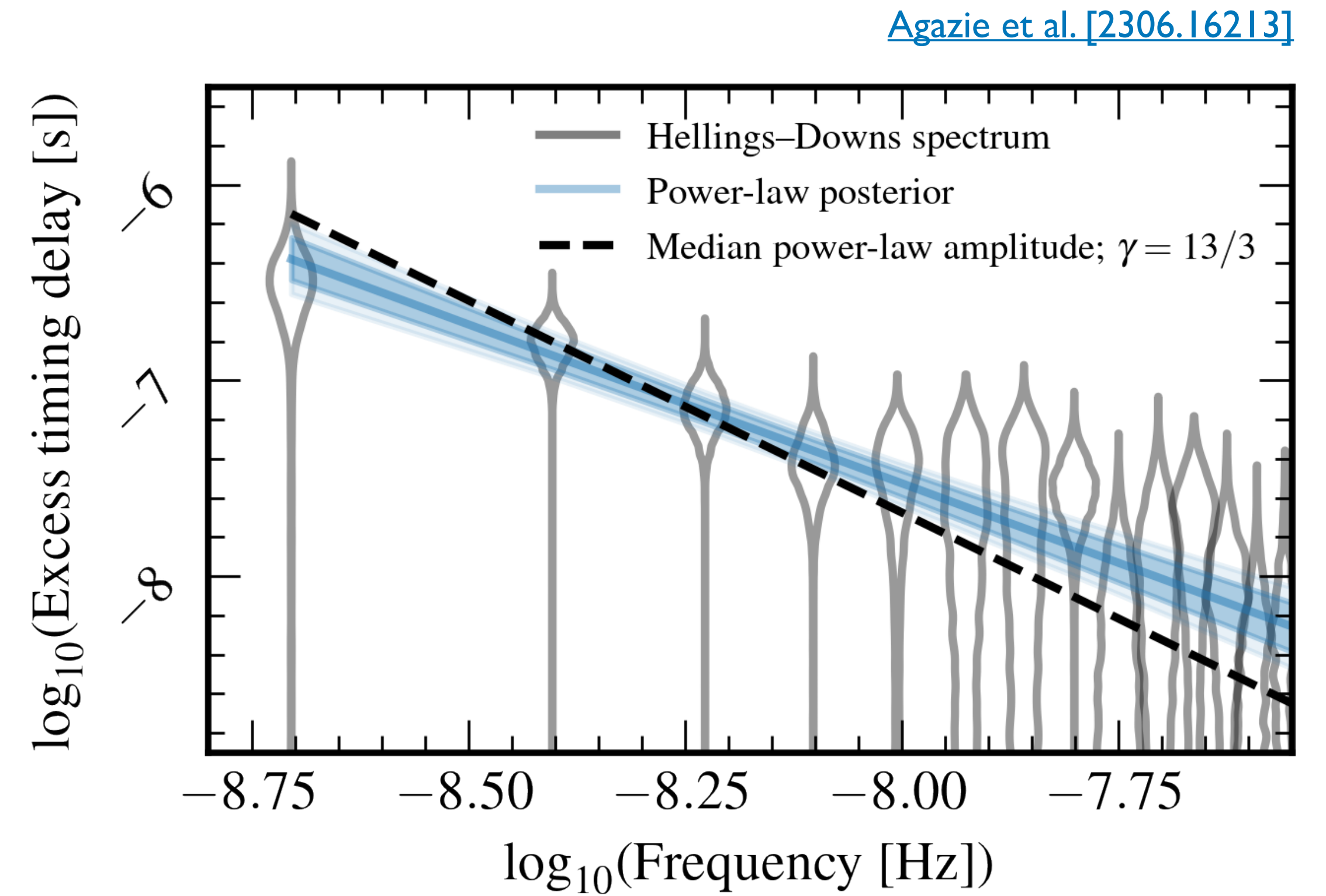
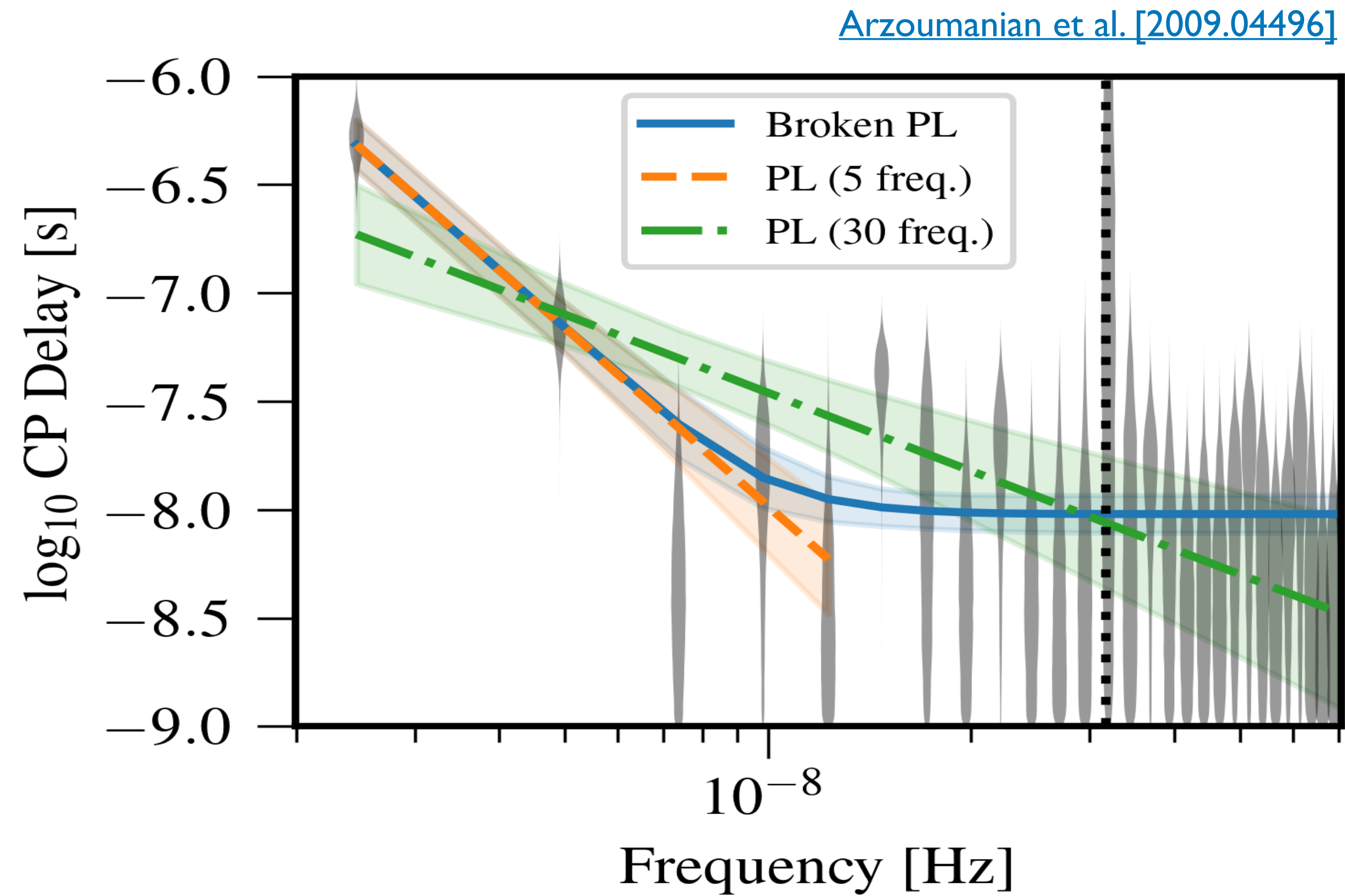
DETERMINISTIC vs STOCHASTIC



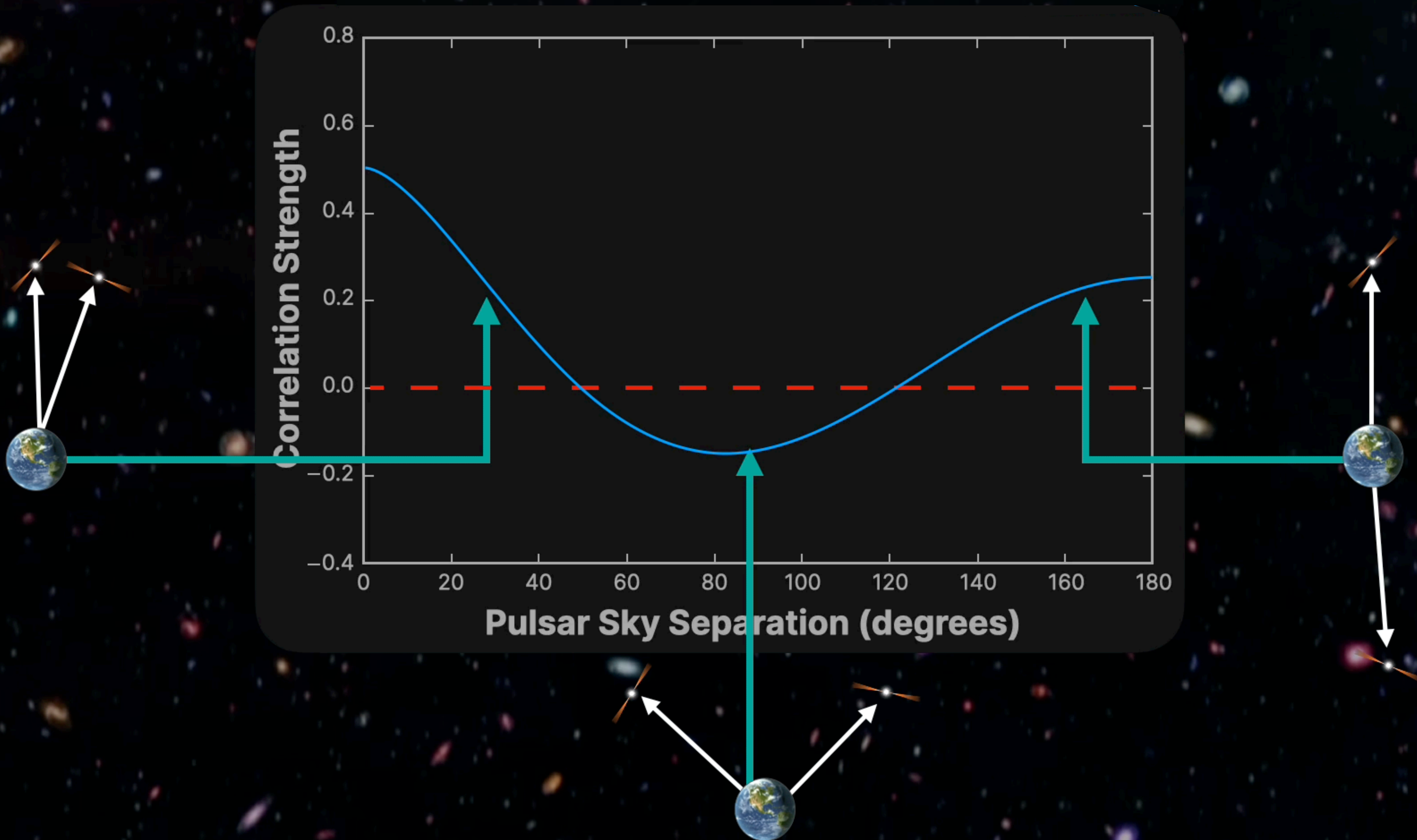
DETERMINISTIC vs STOCHASTIC



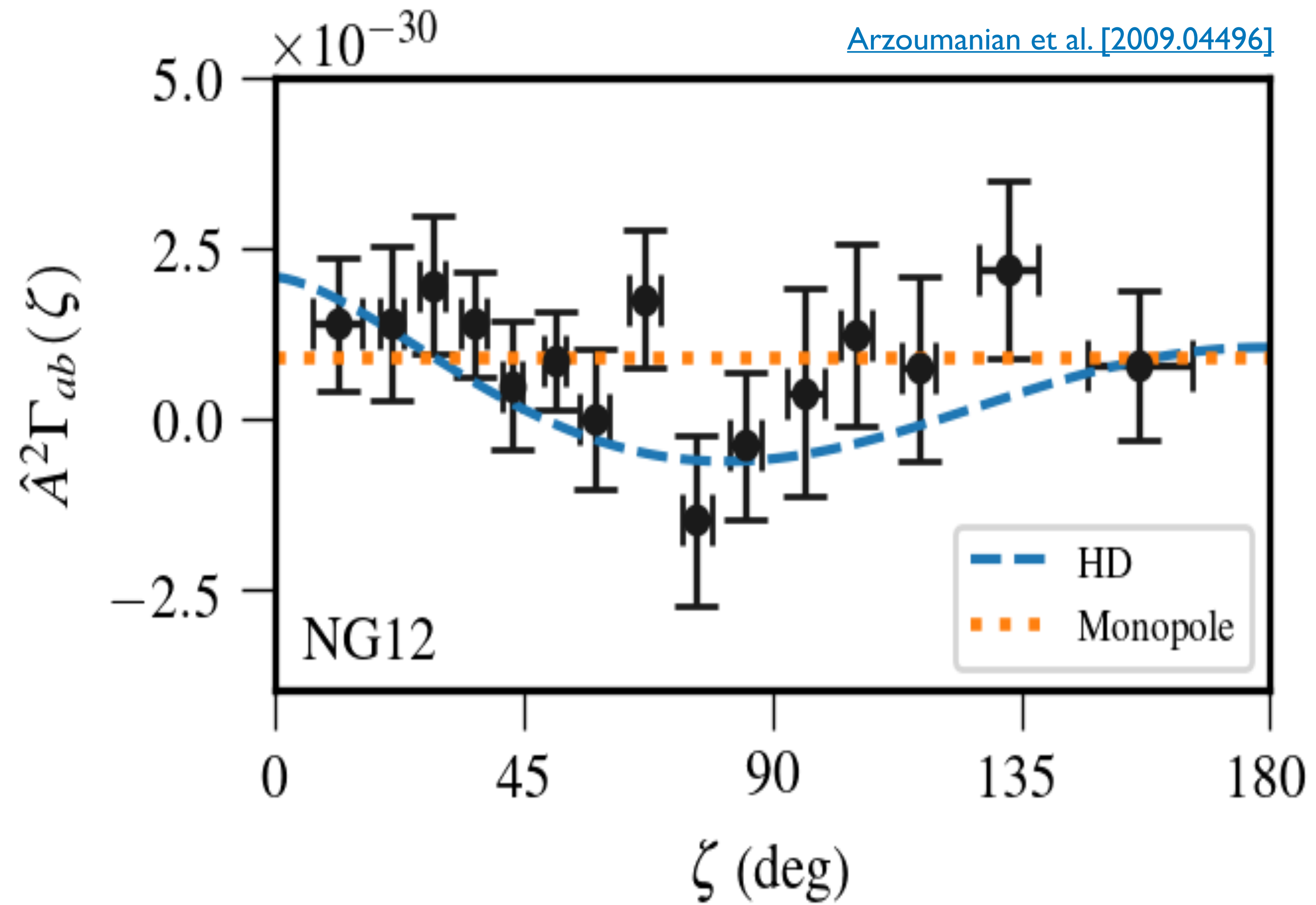
EVIDENCE FOR A COMMON SIGNAL



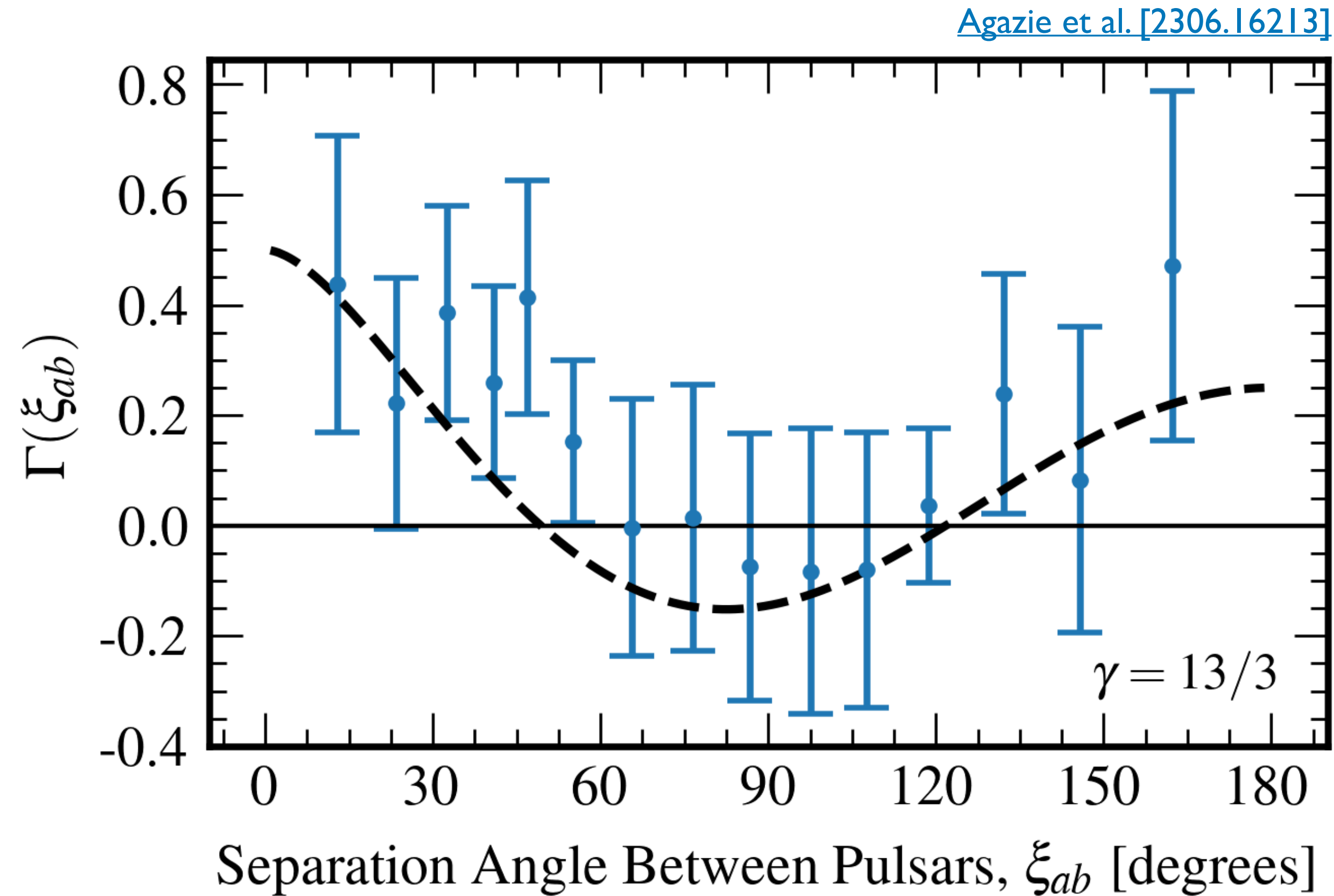
WHAT ABOUT CORRELATIONS?



WHAT ABOUT CORRELATIONS?



WHAT ABOUT CORRELATIONS?



we find evidence for Hellings & Downs correlation with a p -value of $5 \times 10^{-5} - 1.9 \times 10^{-4}$ (approx. $3.5 - 4\sigma$)

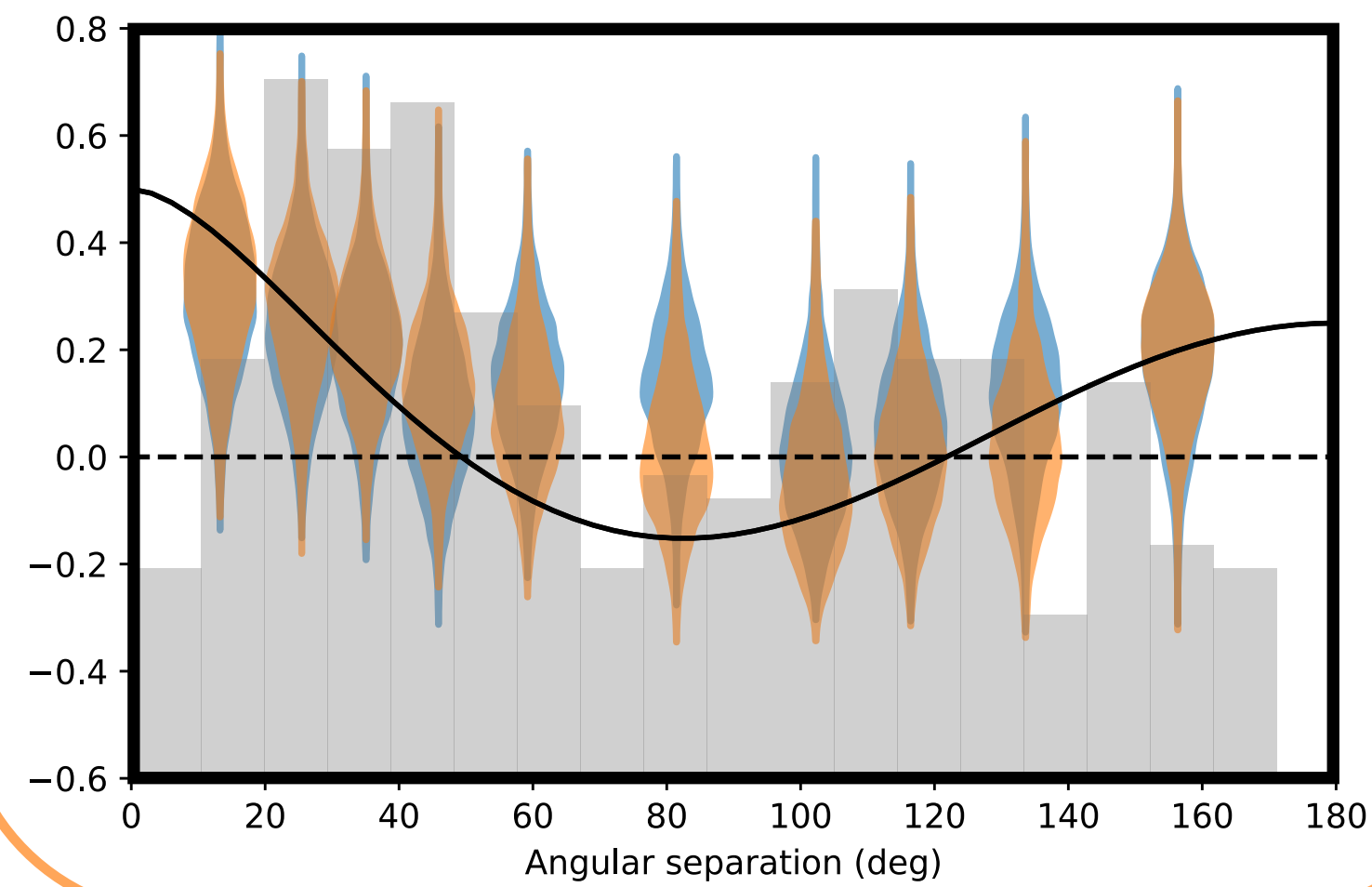
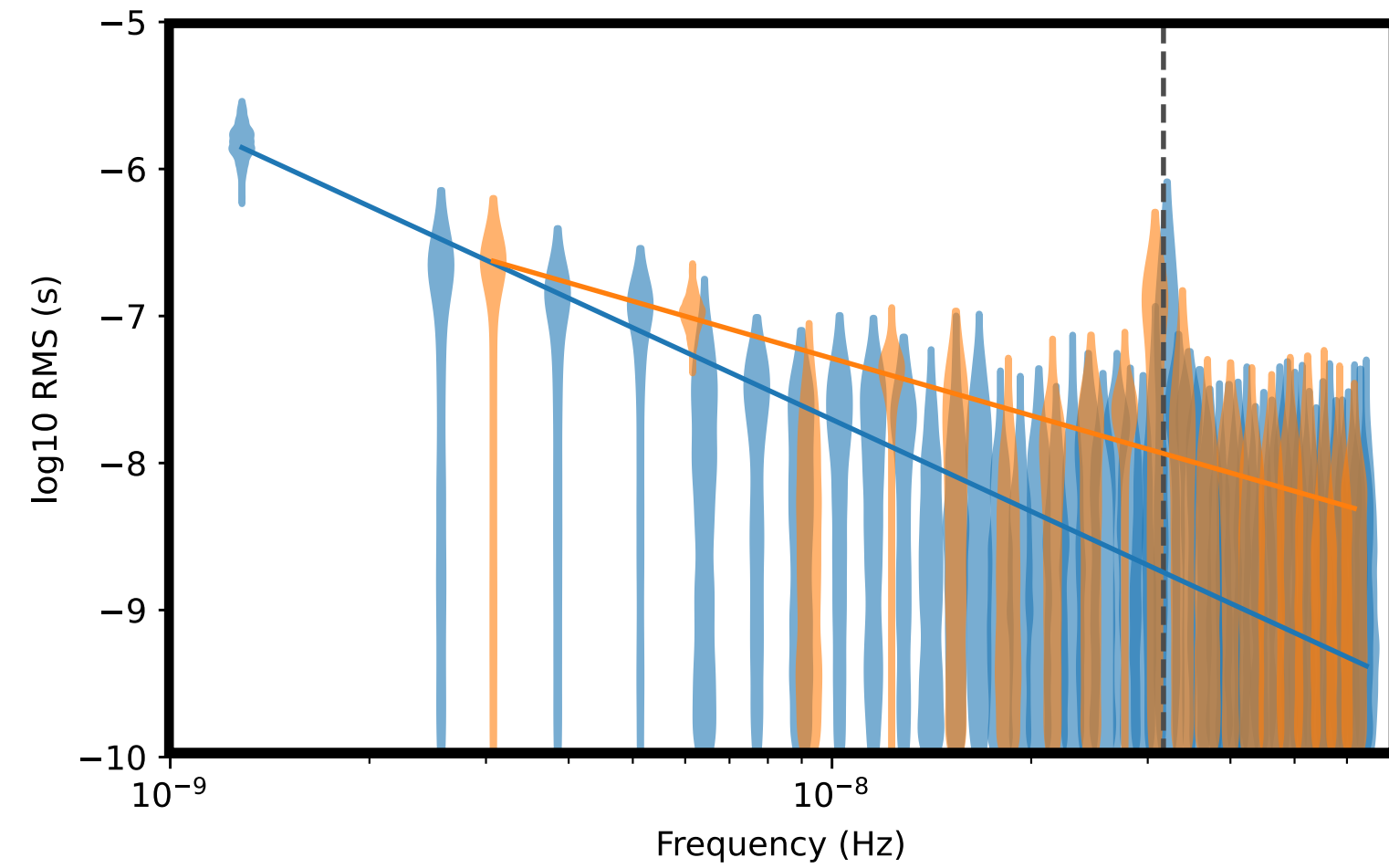
OTHER PTAs



OTHER PTAs

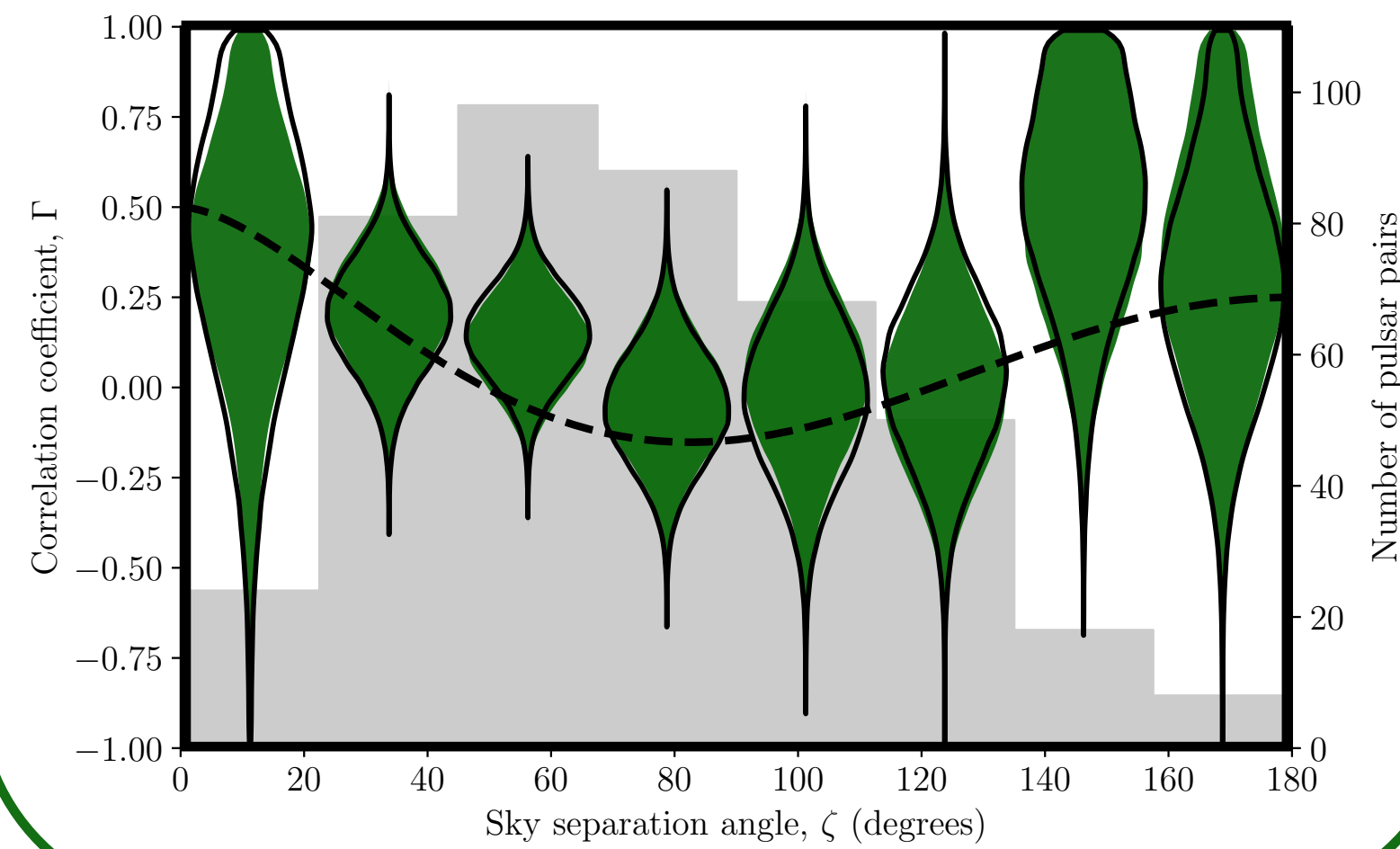
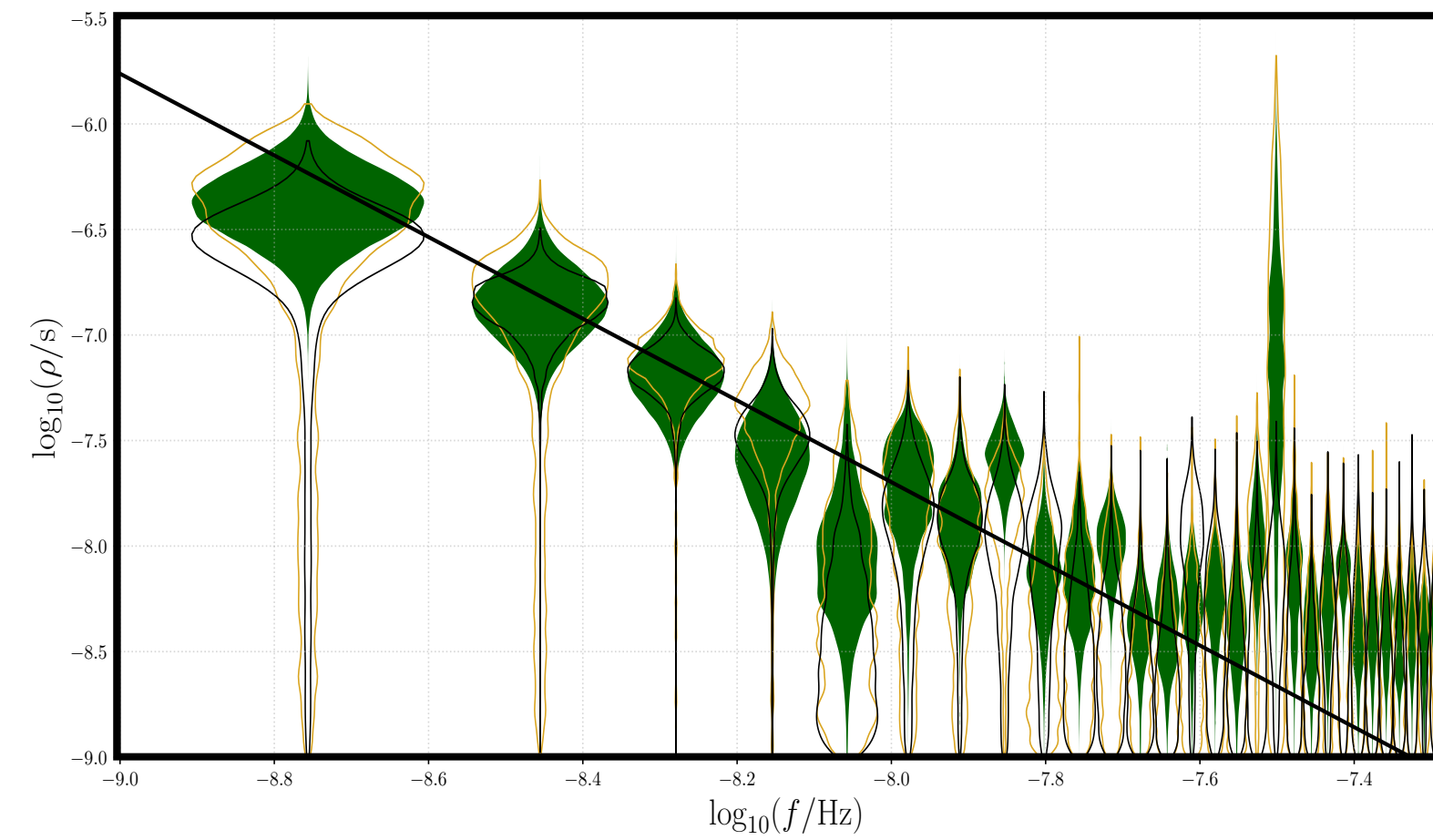
EPTA: 3σ

[Antoniadis et al. \[2306.16214\]](#)



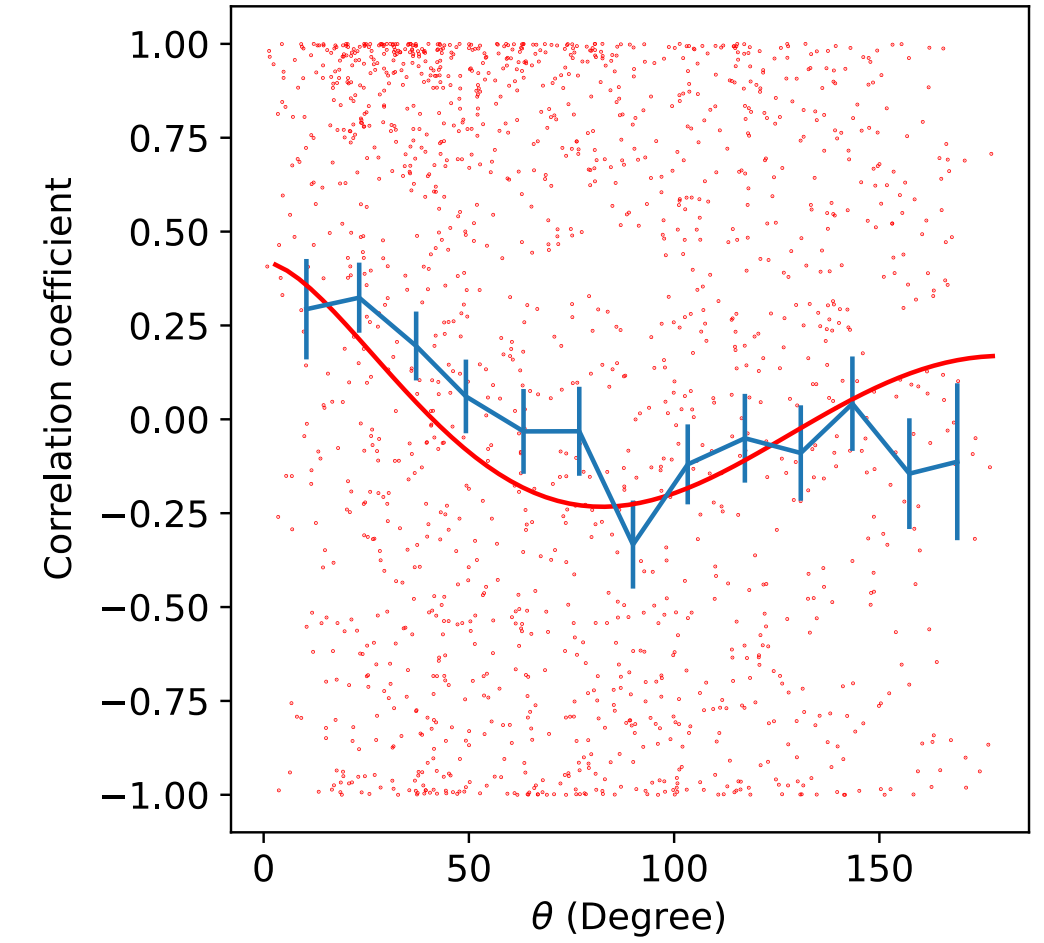
PPTA

[Reardon et al. \[2306.16215\]](#)

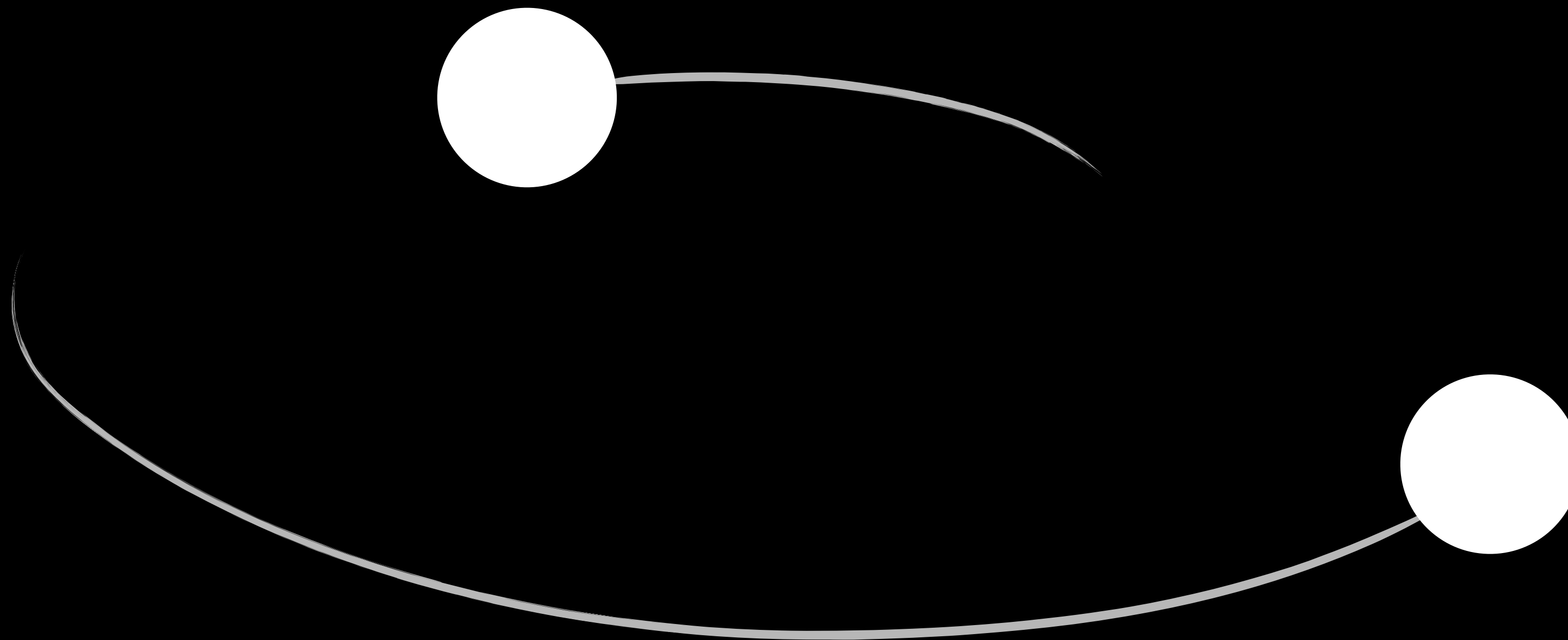


CPTA: 4.6σ

[Xu et al. \[2306.16216\]](#)



what is the source?



Supermassive Black Holes Binaries







GWB FROM SMBHB

$$h_c^2(f) = \int dM dq dz \frac{\partial^4 N}{\partial M \partial q \partial z \partial \ln f_p} h_s^2(f_p)$$

Phinney 2001, Wyithe & Loeb 2003

GWB FROM SMBHB

GW signal from individual SMBHB

$$h_c^2(f) = \int dM dq dz \frac{\partial^4 N}{\partial M \partial q \partial z \partial \ln f_p} h_s^2(f_p)$$

Phinney 2001, Wyithe & Loeb 2003

averaged strain for a circular
SMBHB

$$h_s^2(f) = \frac{32}{5} \frac{(G\mathcal{M})^{10/3}}{d_c^2} (2\pi f_p)^{4/3}$$

Finn & Thorne 2000

GWB FROM SMBHB

GW signal from individual SMBHB

$$h_c^2(f) = \int dM dq dz \frac{\partial^4 N}{\partial M \partial q \partial z \partial \ln f_p} h_s^2(f_p)$$

Phinney 2001, Wyithe & Loeb 2003

number density of SMBHB binaries

averaged strain for a circular
SMBHB

$$h_s^2(f) = \frac{32}{5} \frac{(GM)^{10/3}}{d_c^2} (2\pi f_p)^{4/3}$$

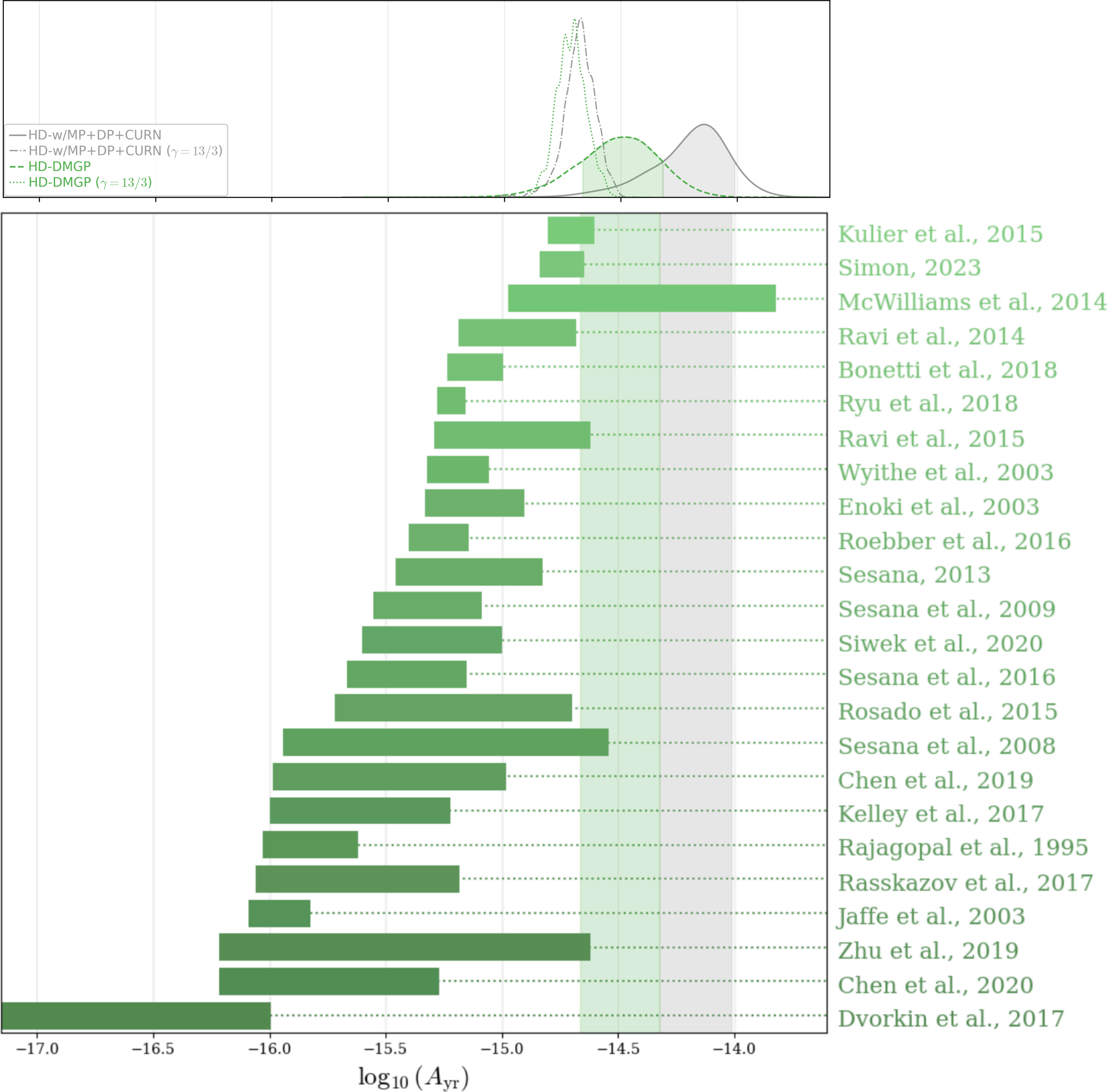
Finn & Thorne 2000

the SMBHB density depends on

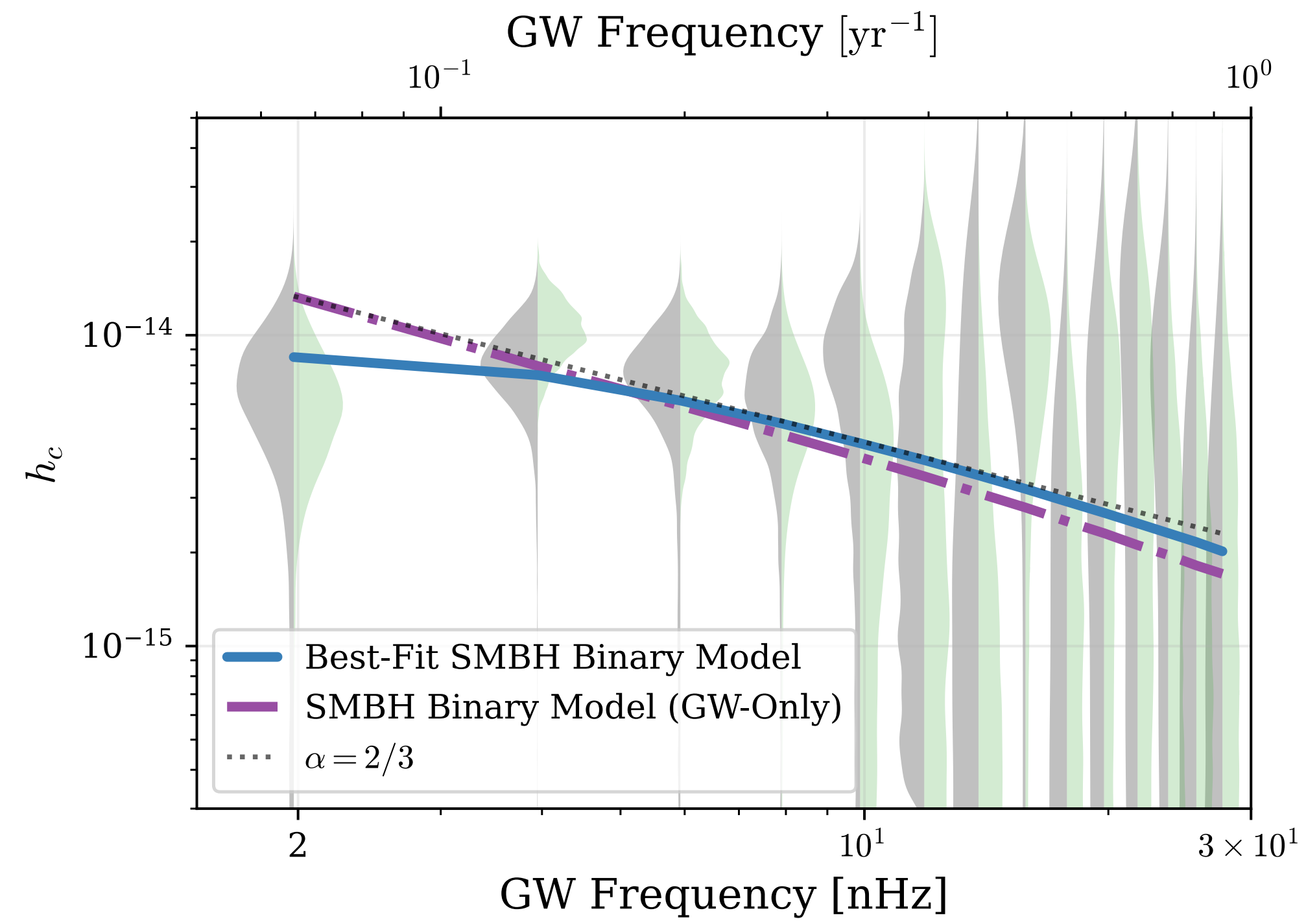
1. galaxies merger rate
2. SMBHB - galaxy mass relation
3. SMBHB binary evolution

EXPECTATIONS

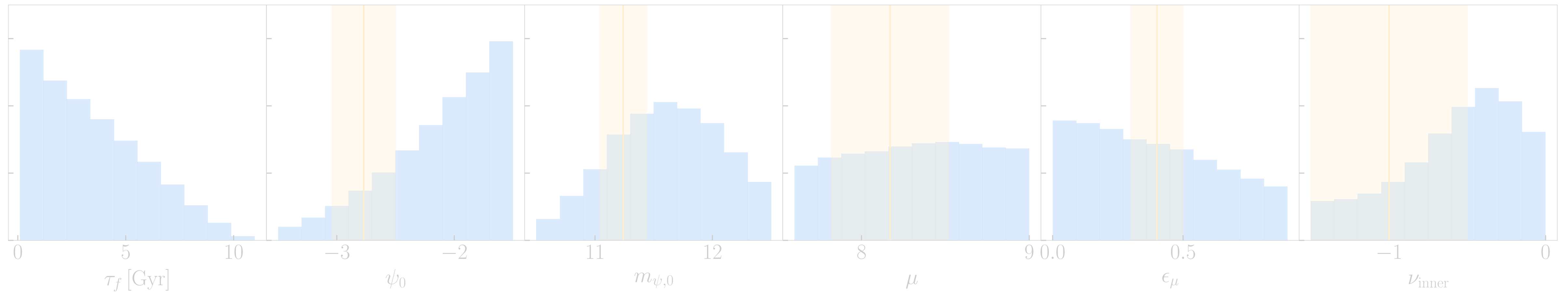
[Agazie et al. \[2306.16220\]](#)



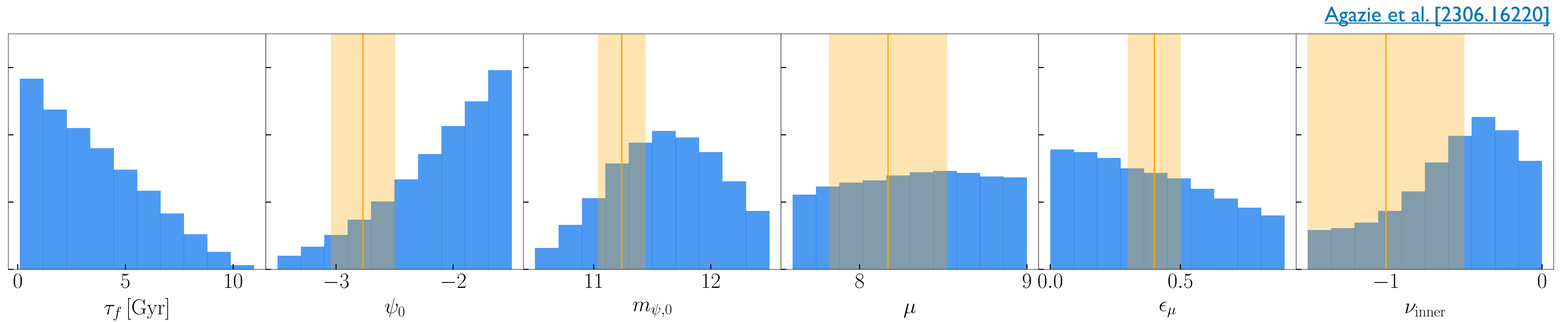
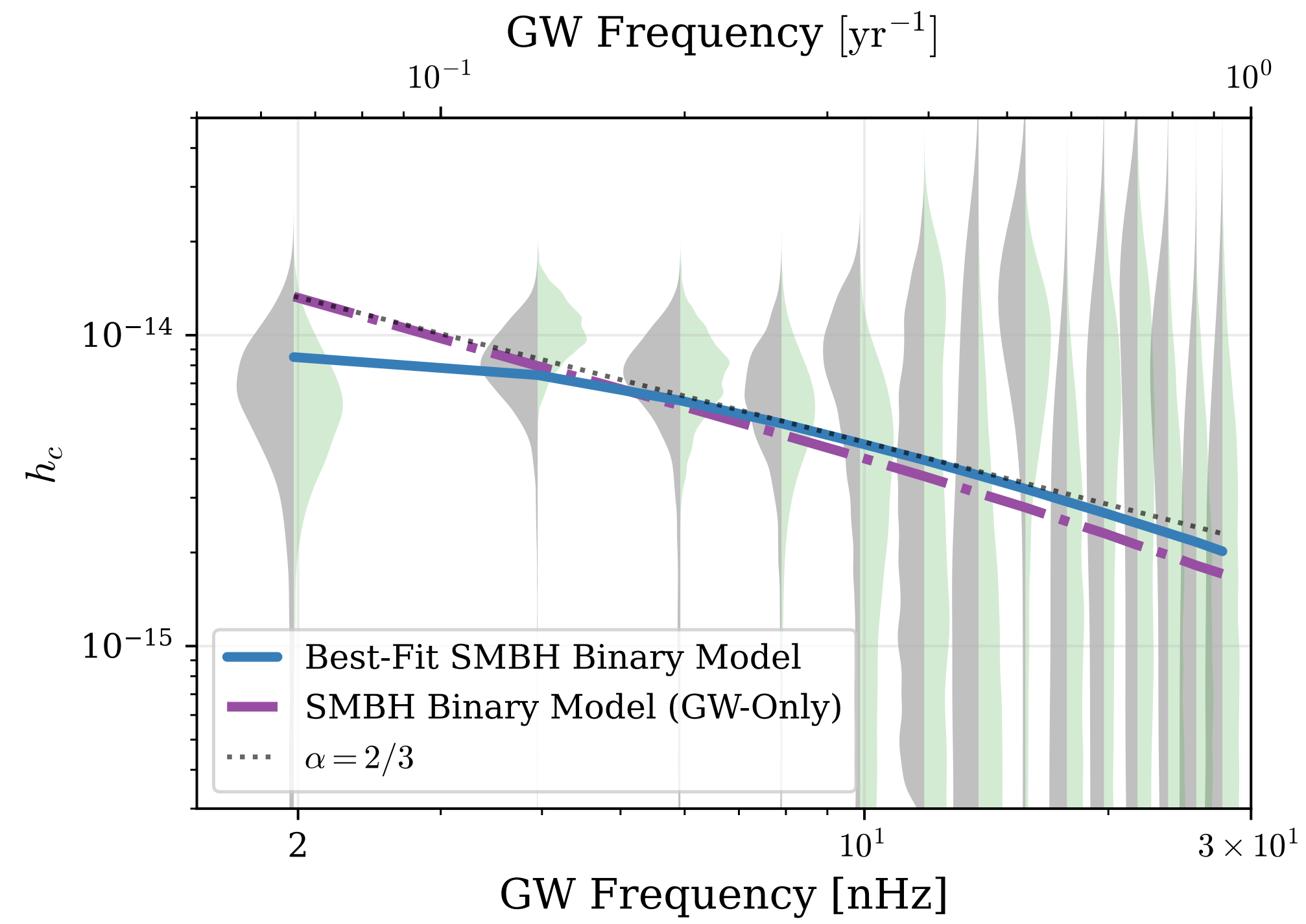
ADJUSTING EXPECTATIONS



[Agazie et al. \[2306.16220\]](#)

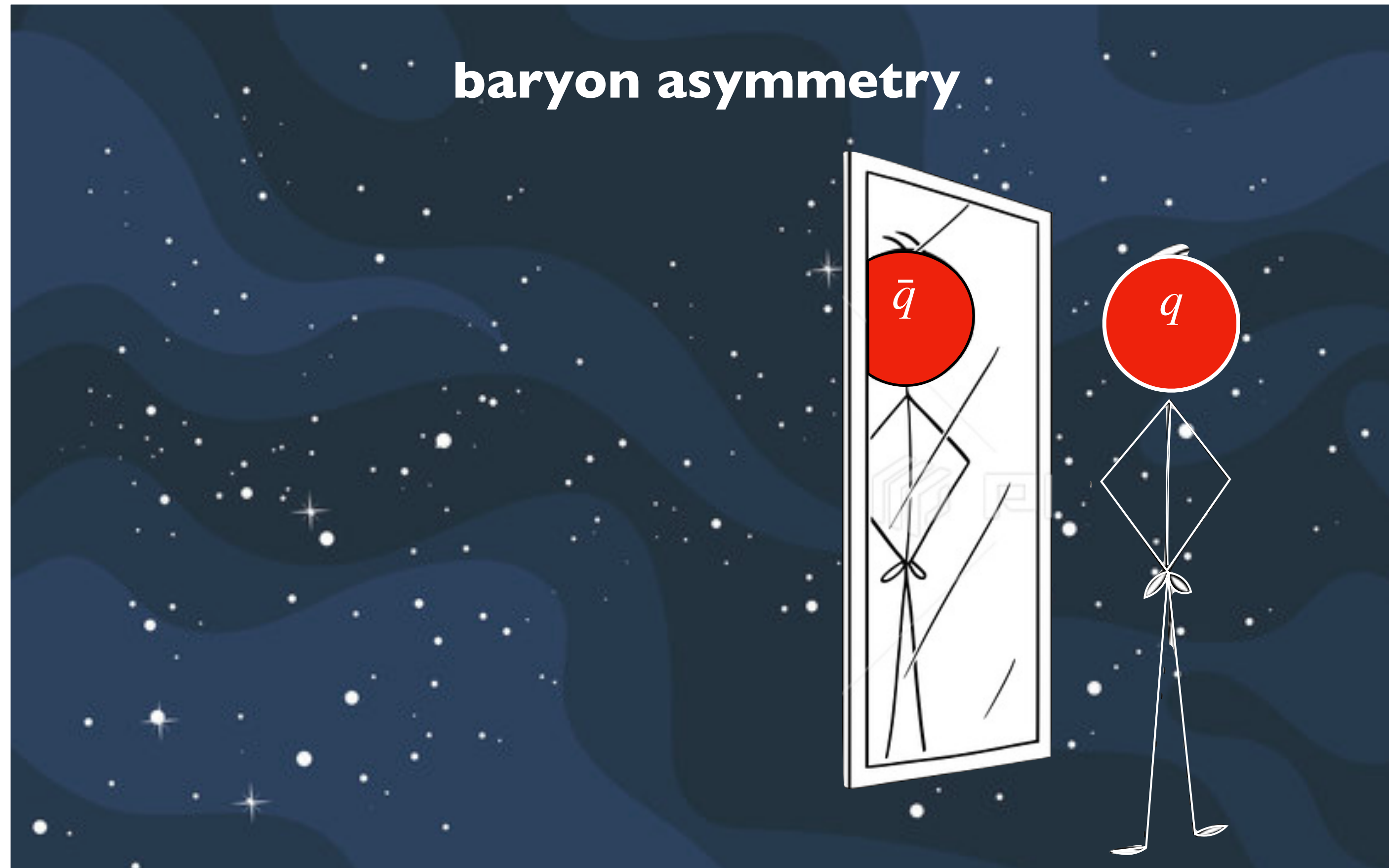


ADJUSTING EXPECTATIONS

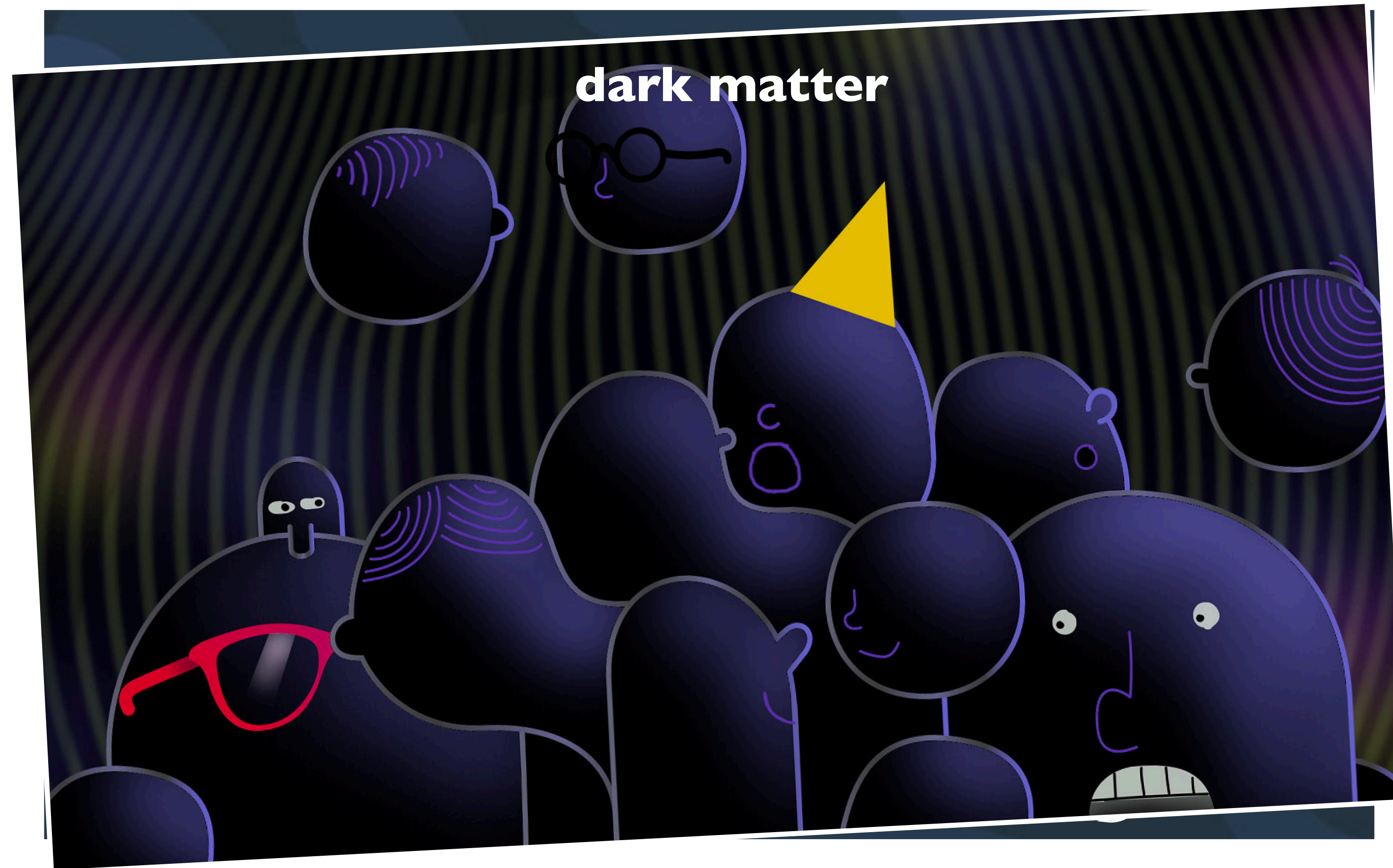


can it be something else?

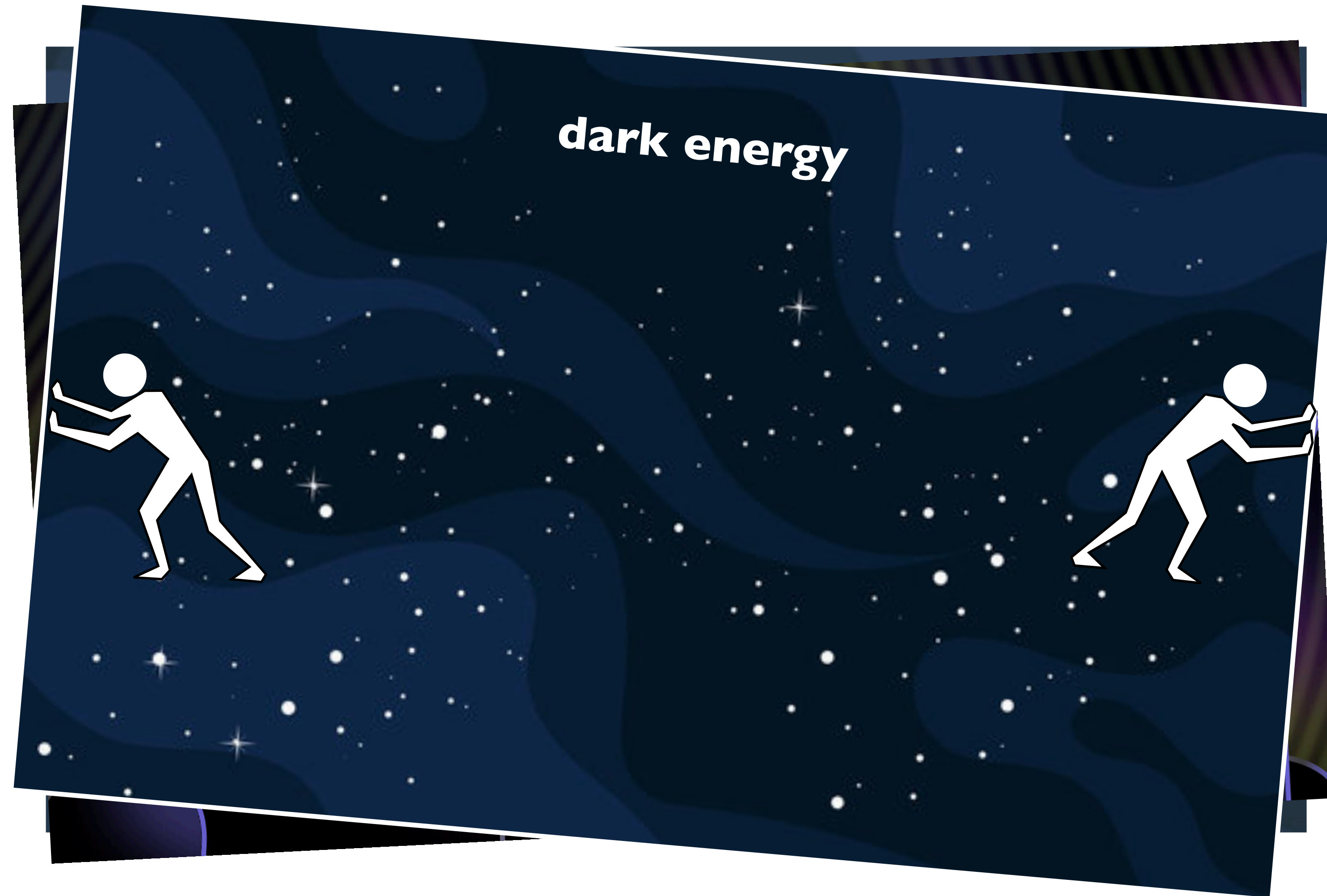
THE BIG PUZZLES



THE BIG PUZZLES



THE BIG PUZZLES



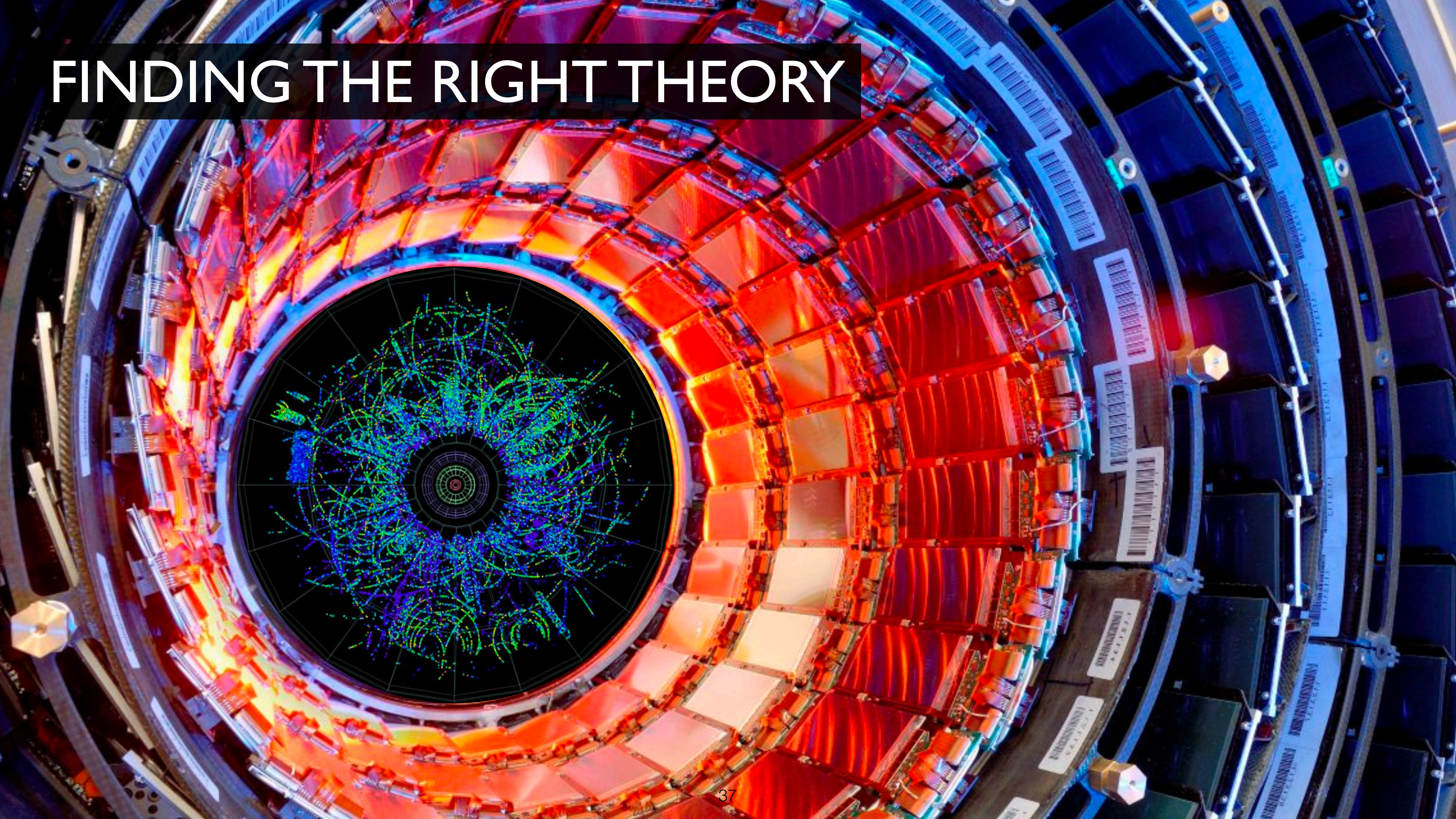
THE BIG PUZZLES

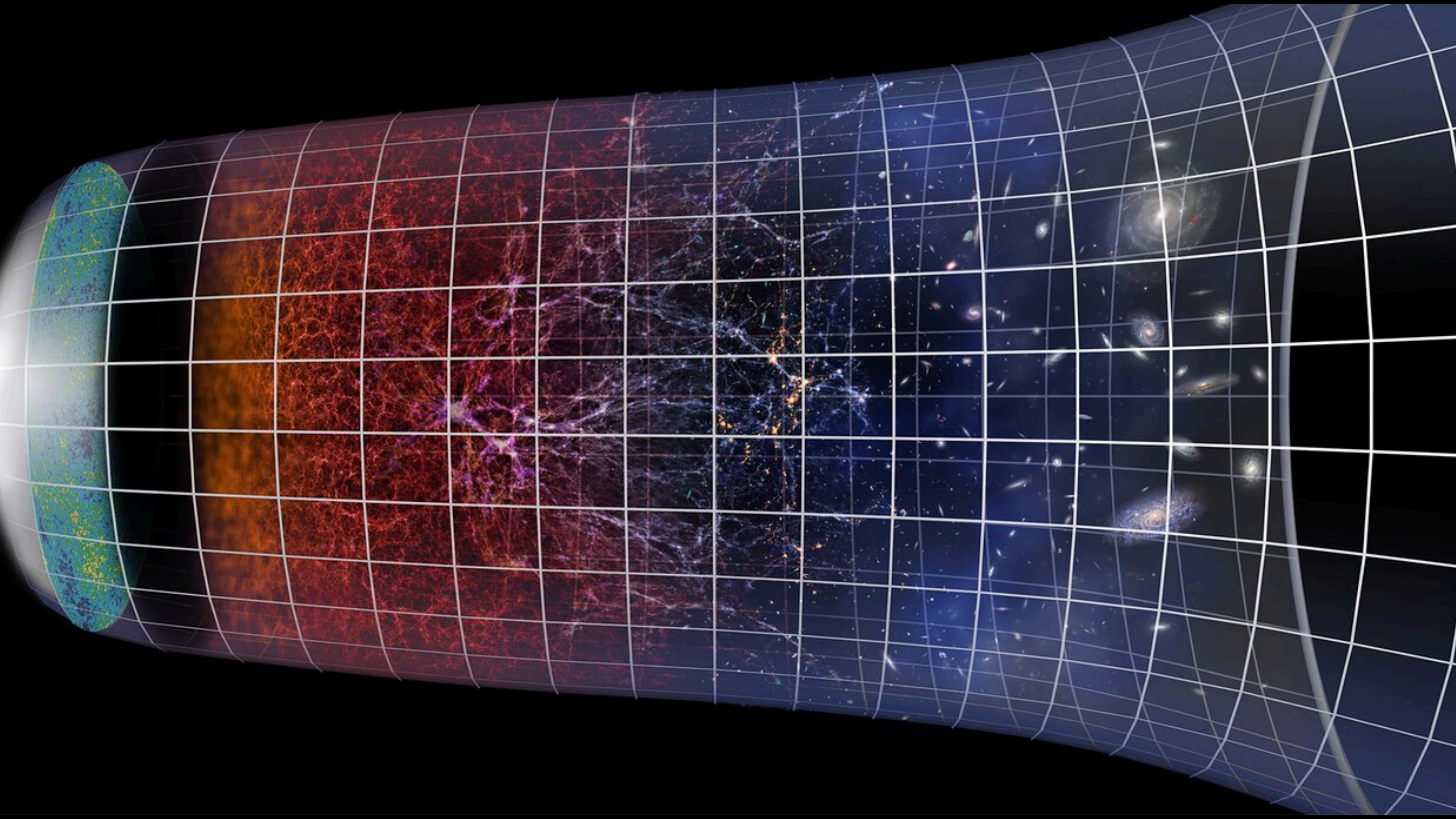


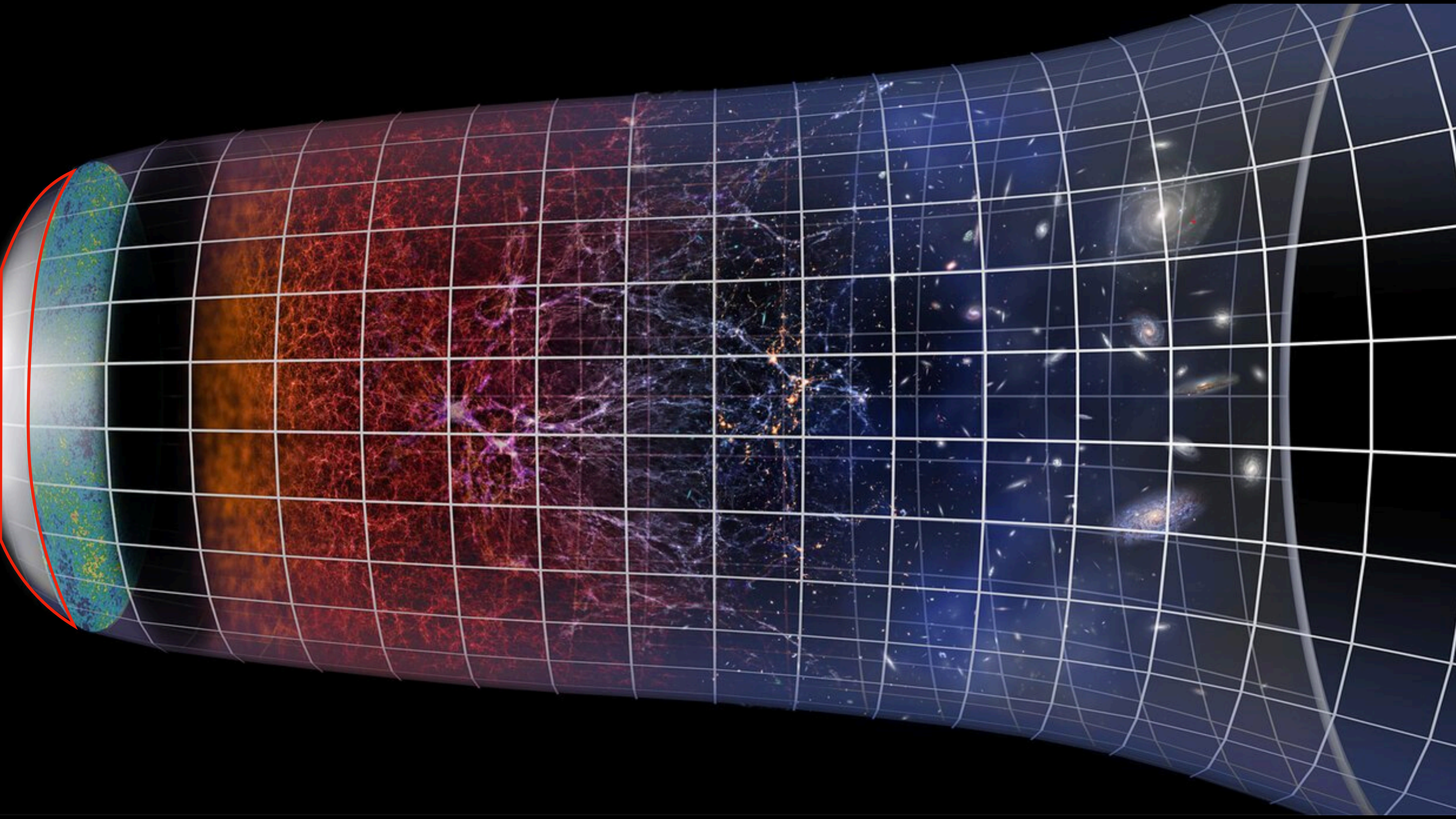
THE BIG PUZZLES

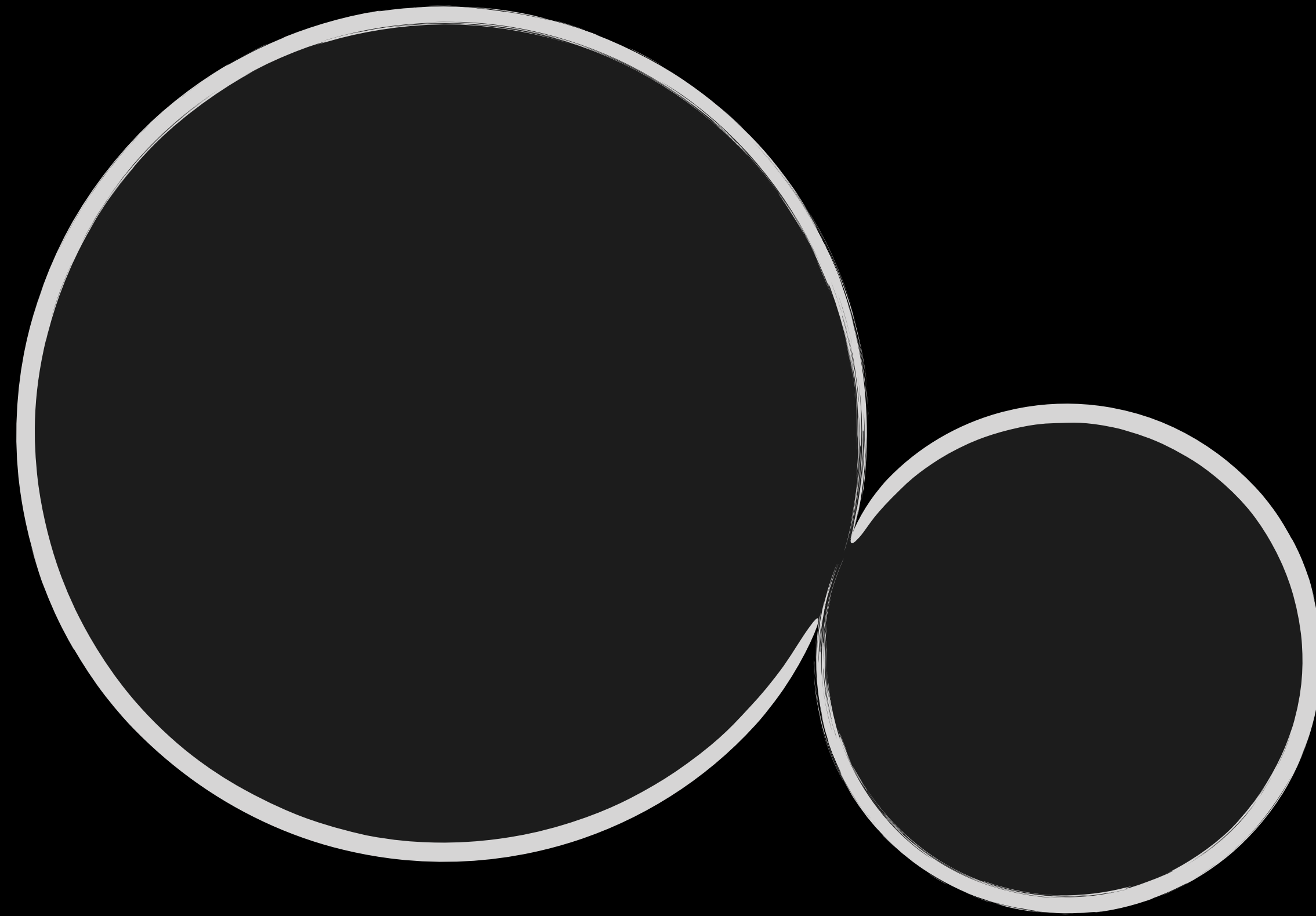


FINDING THE RIGHT THEORY







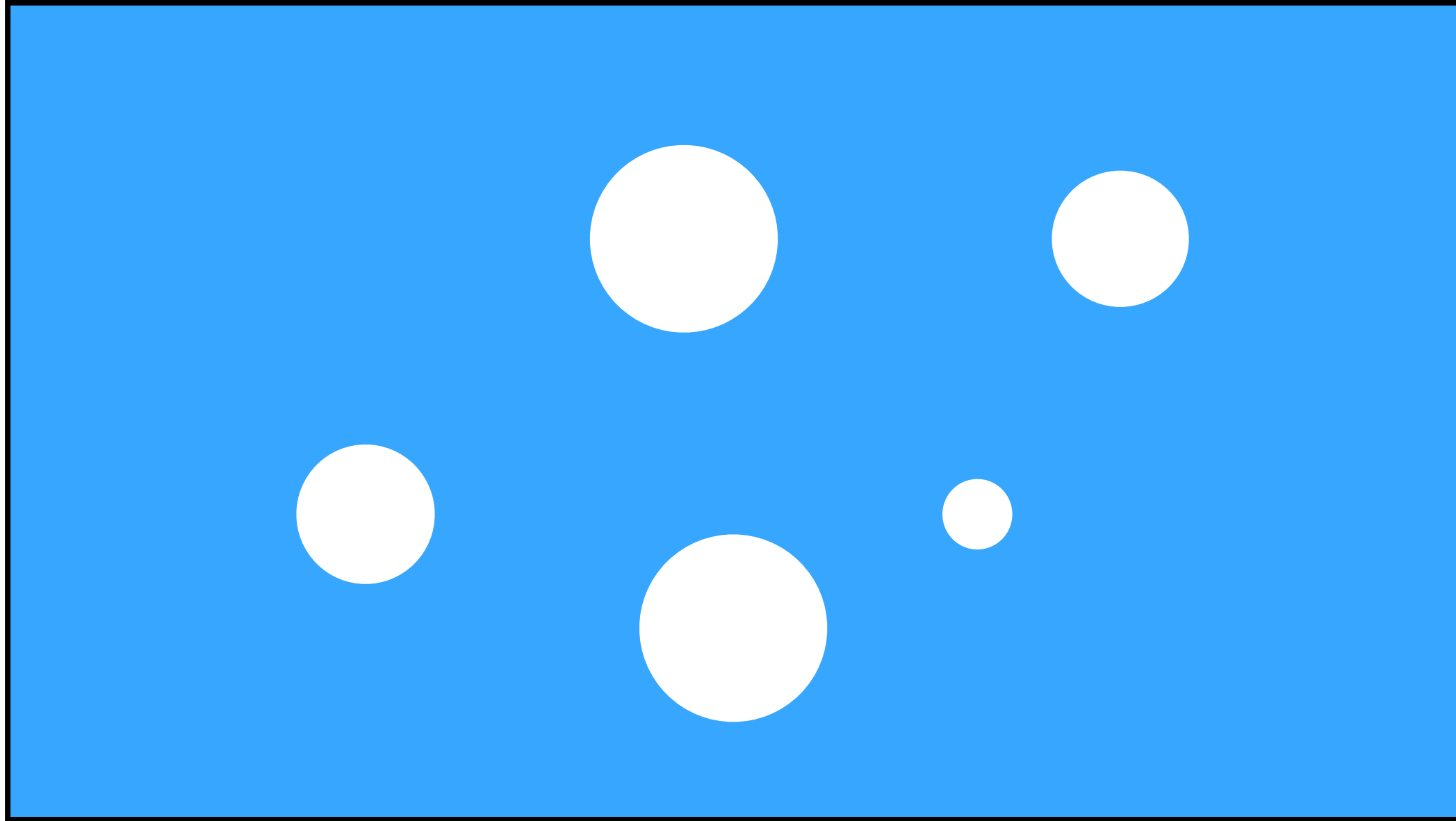


GWB from cosmological phase transitions

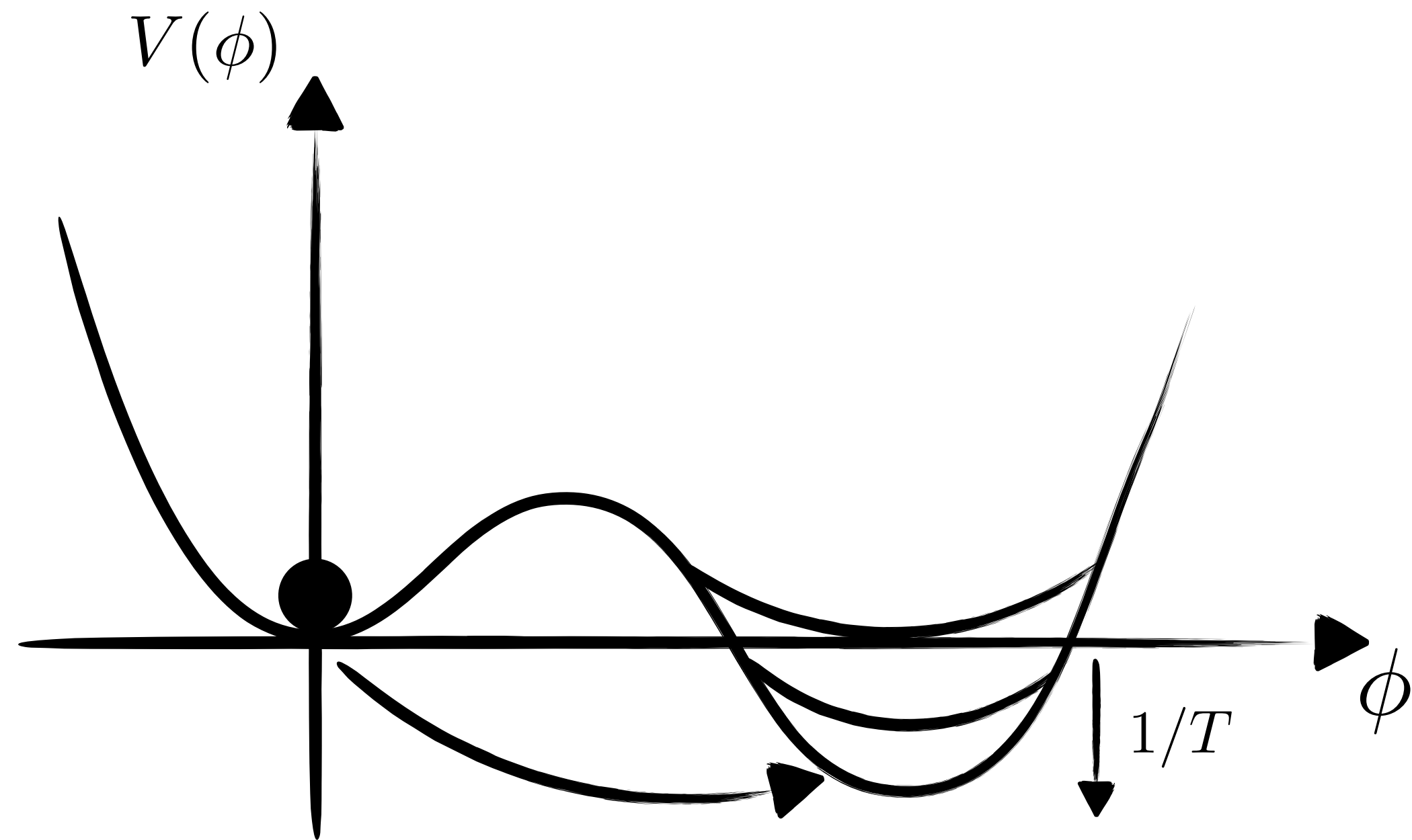
PHASE TRANSITIONS



PHASE TRANSITIONS

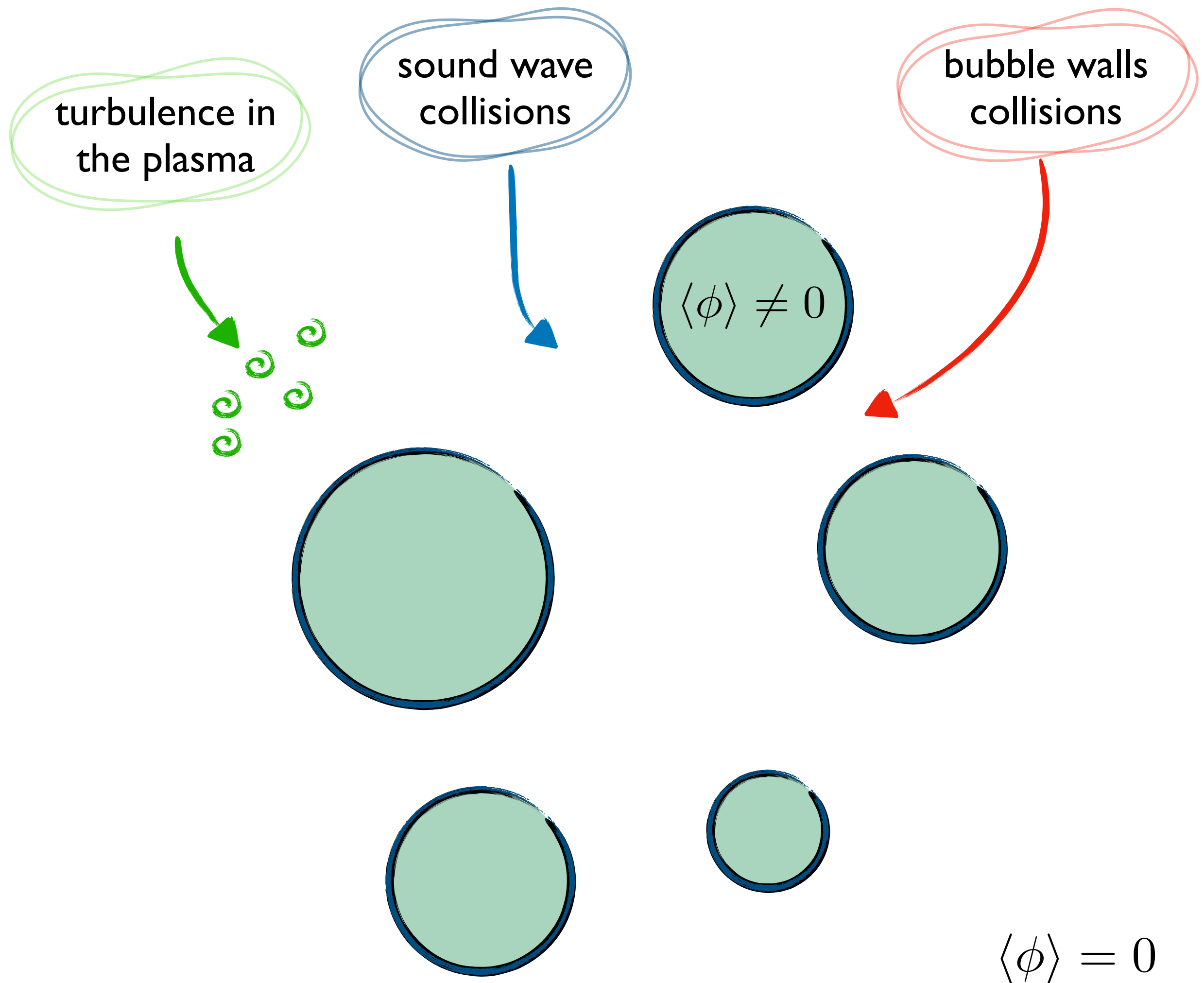
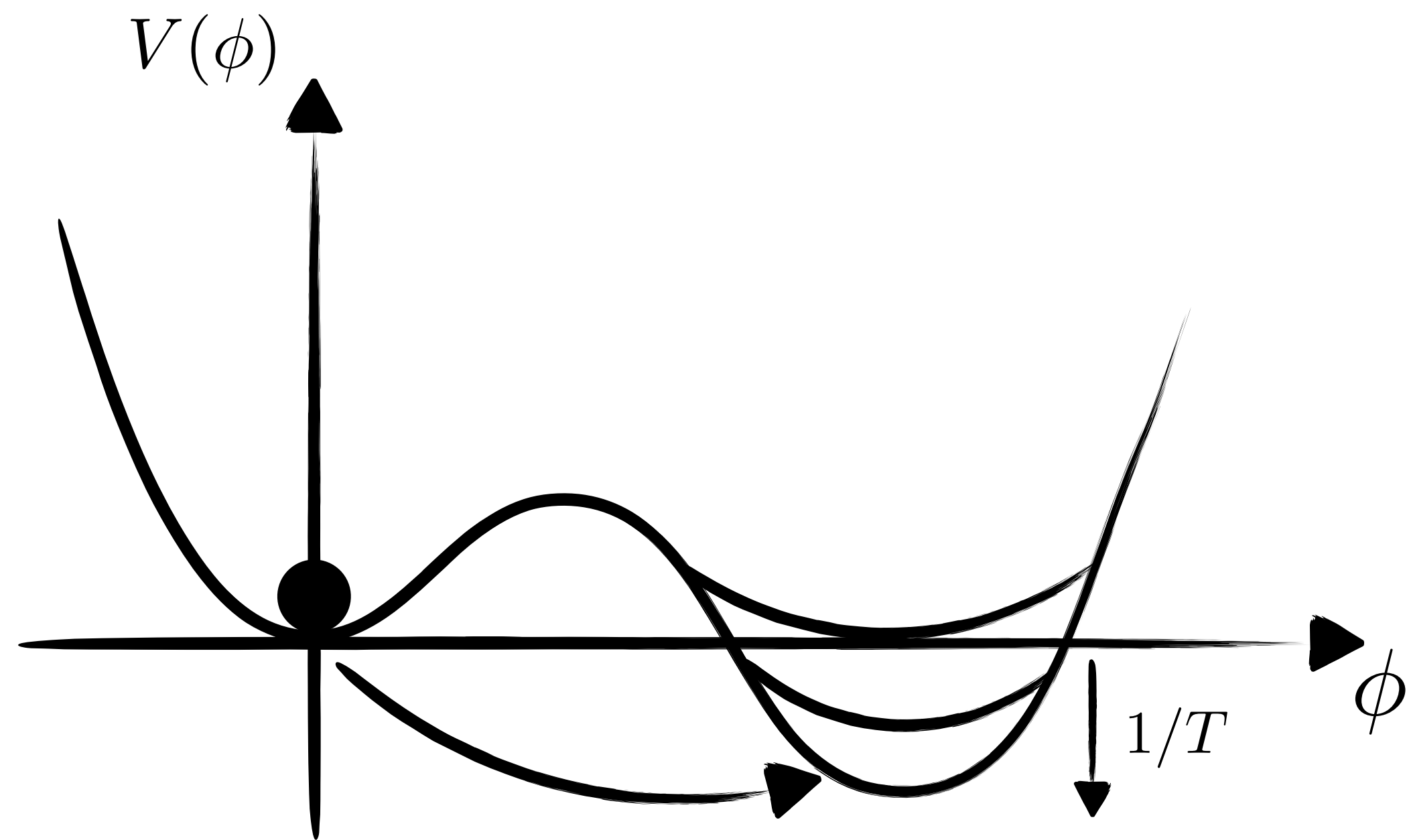


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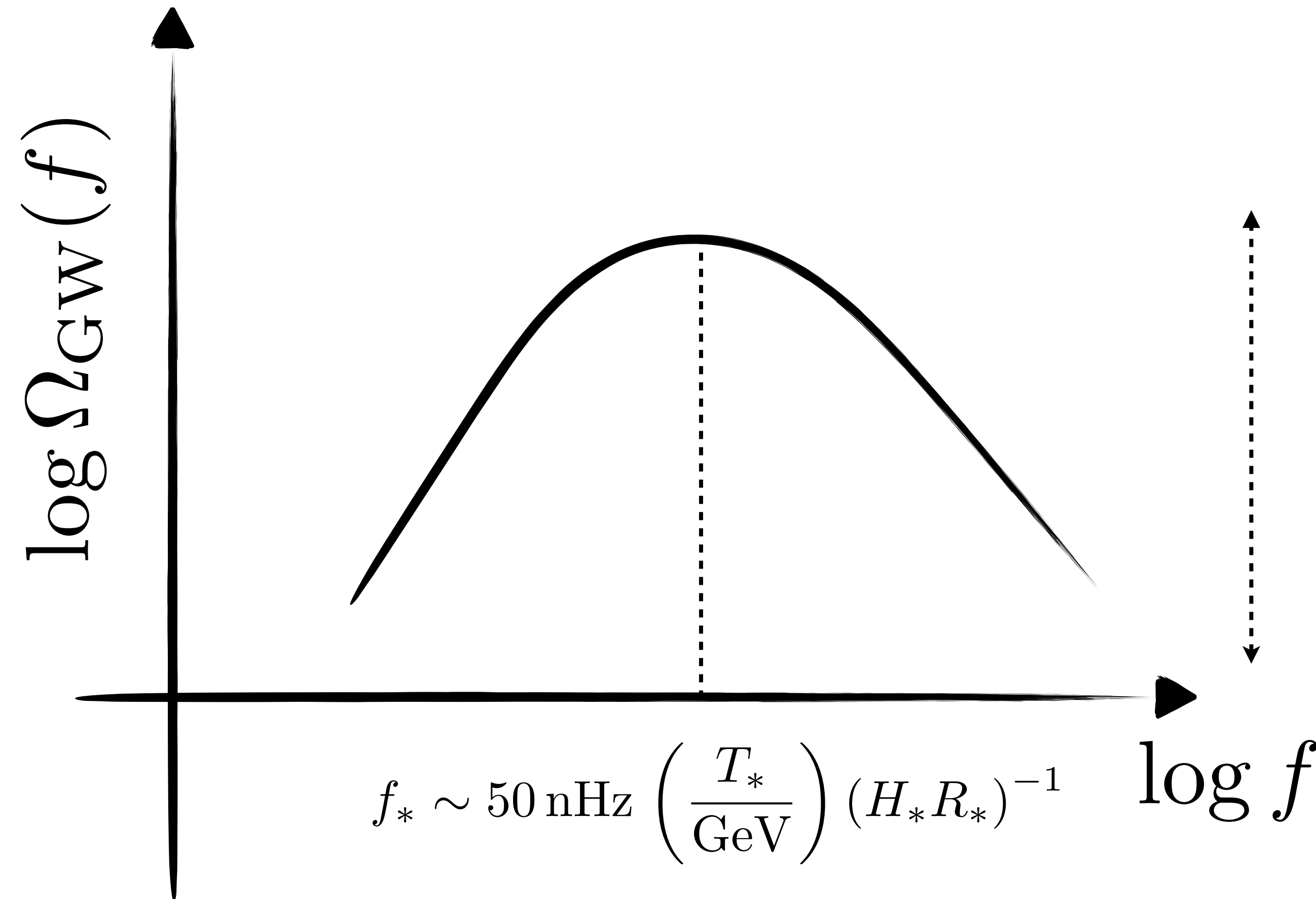


$$\langle \phi \rangle = 0$$

PHASE TRANSITIONS

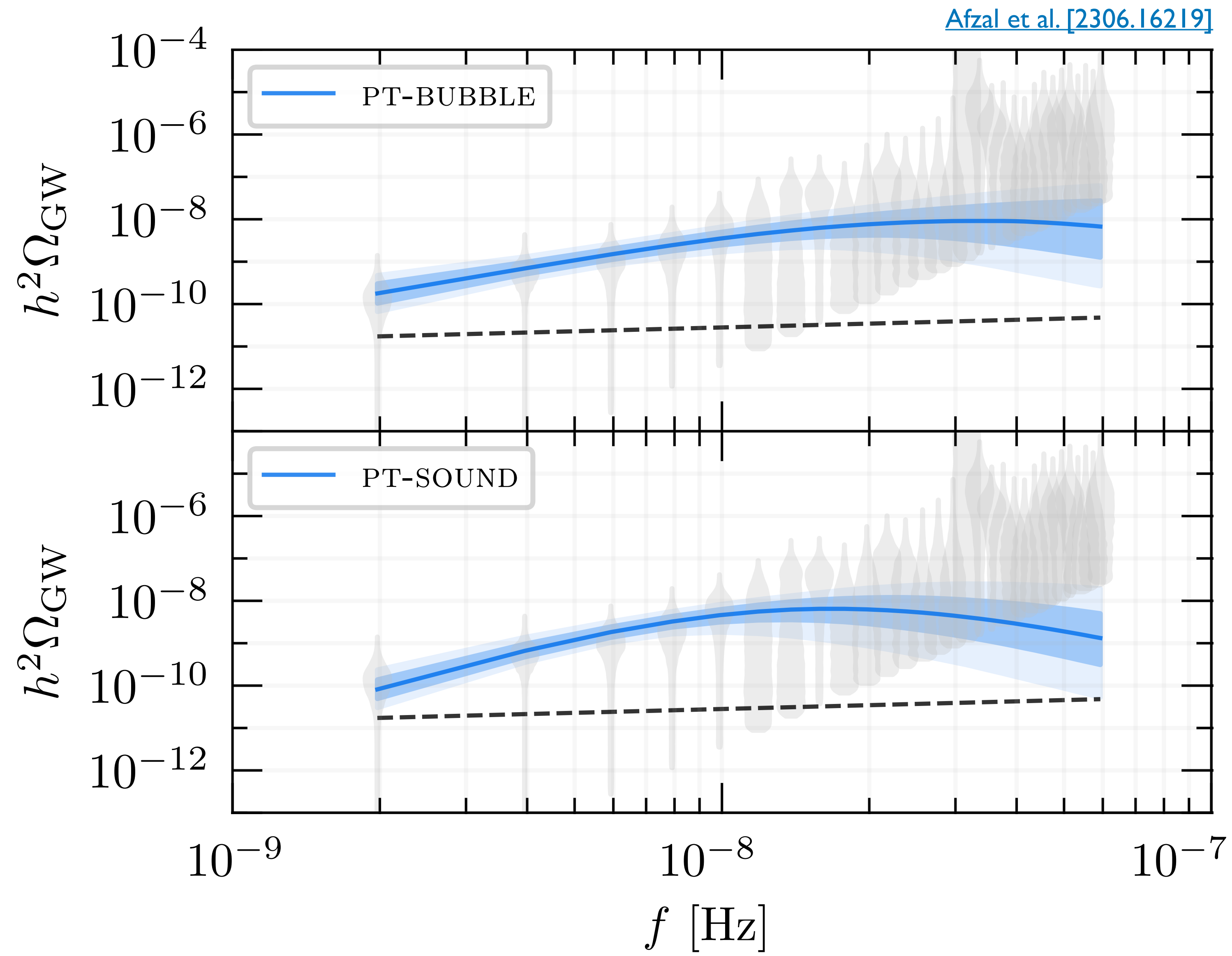


PHASE TRANSITIONS

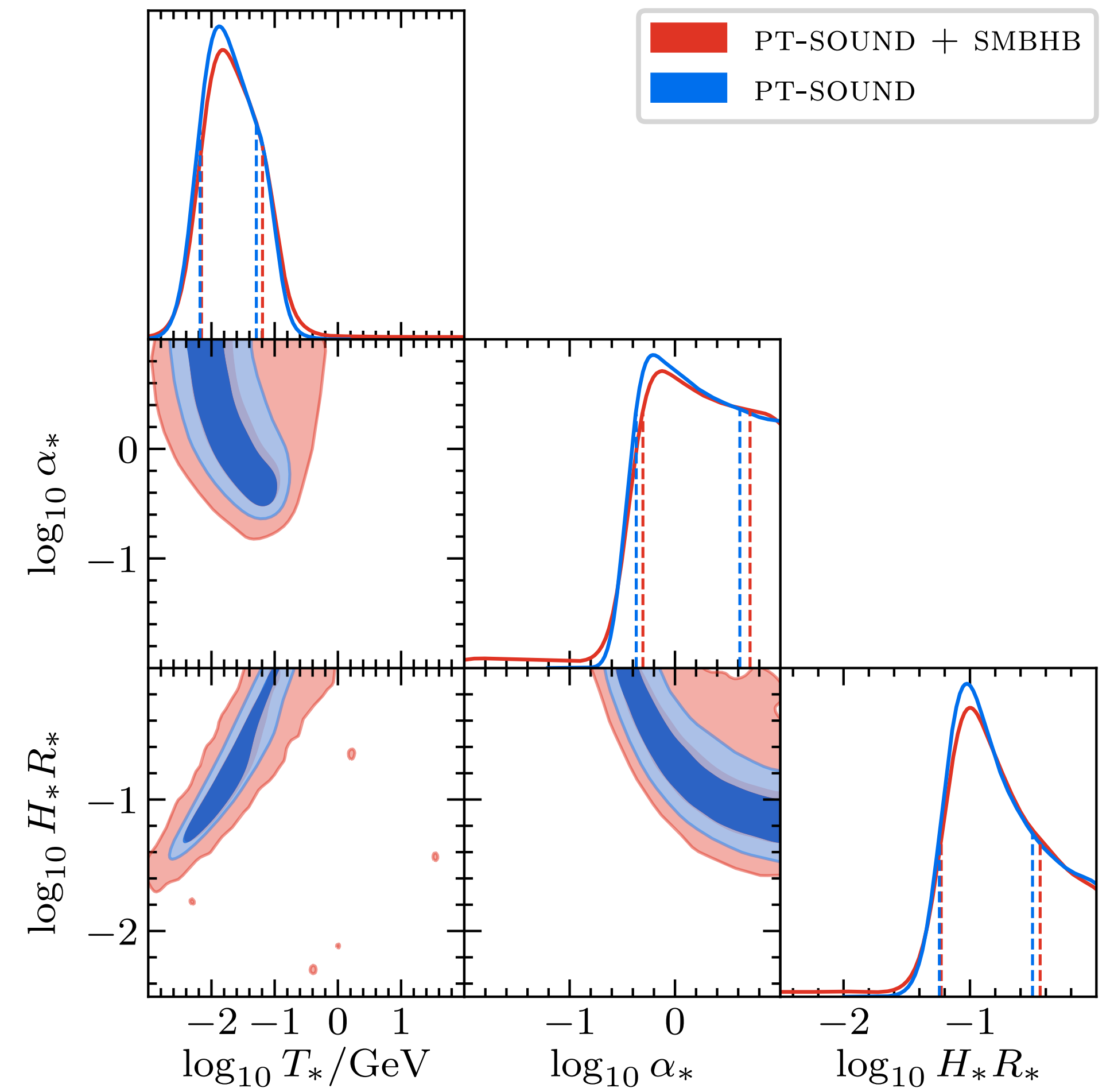
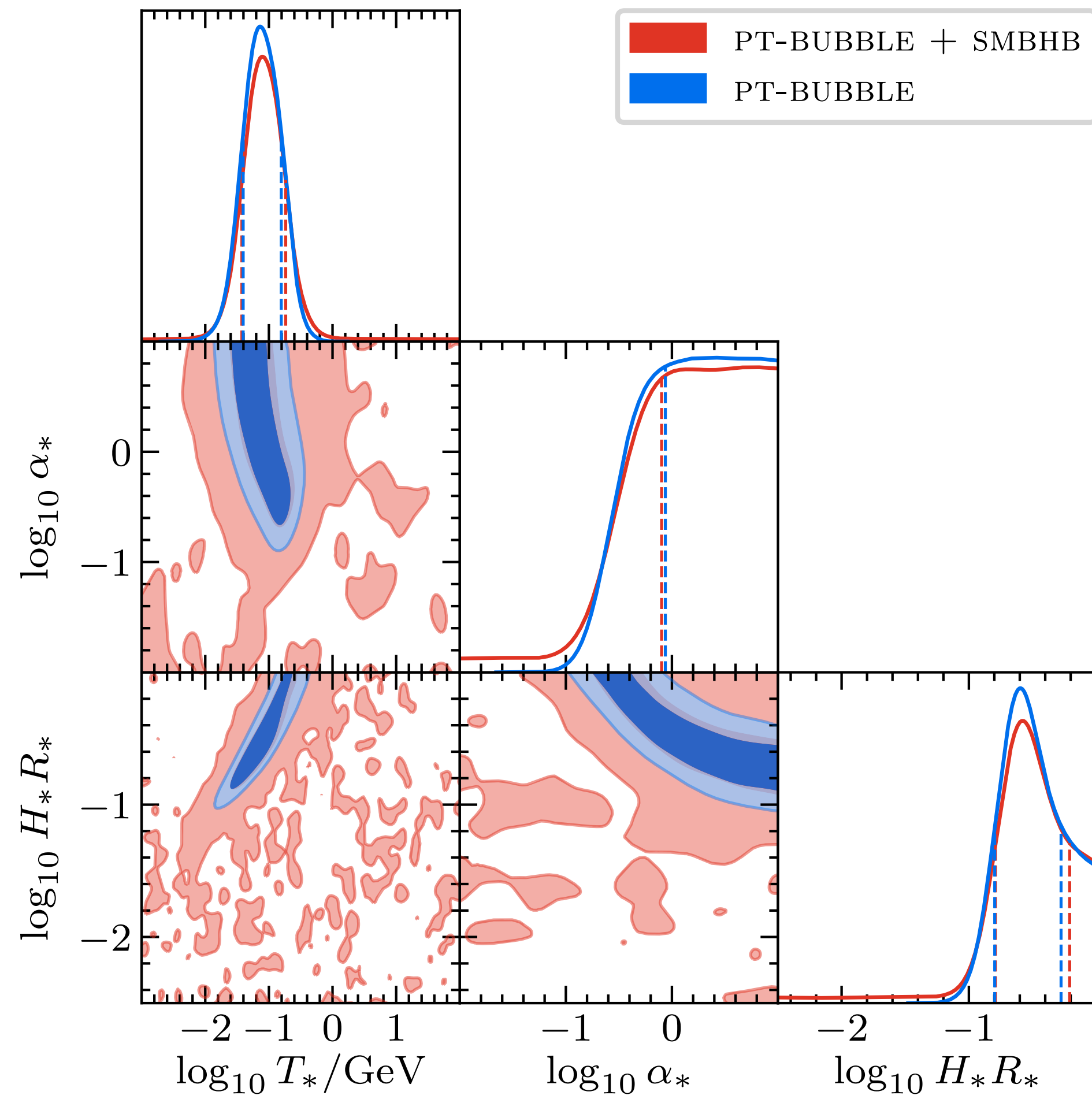


$$\sim 10^{-5} \left(\frac{\alpha_*}{1 + \alpha_*} \right)^2 (H_* R_*)^p$$

PHASE TRANSITIONS



PHASE TRANSITIONS



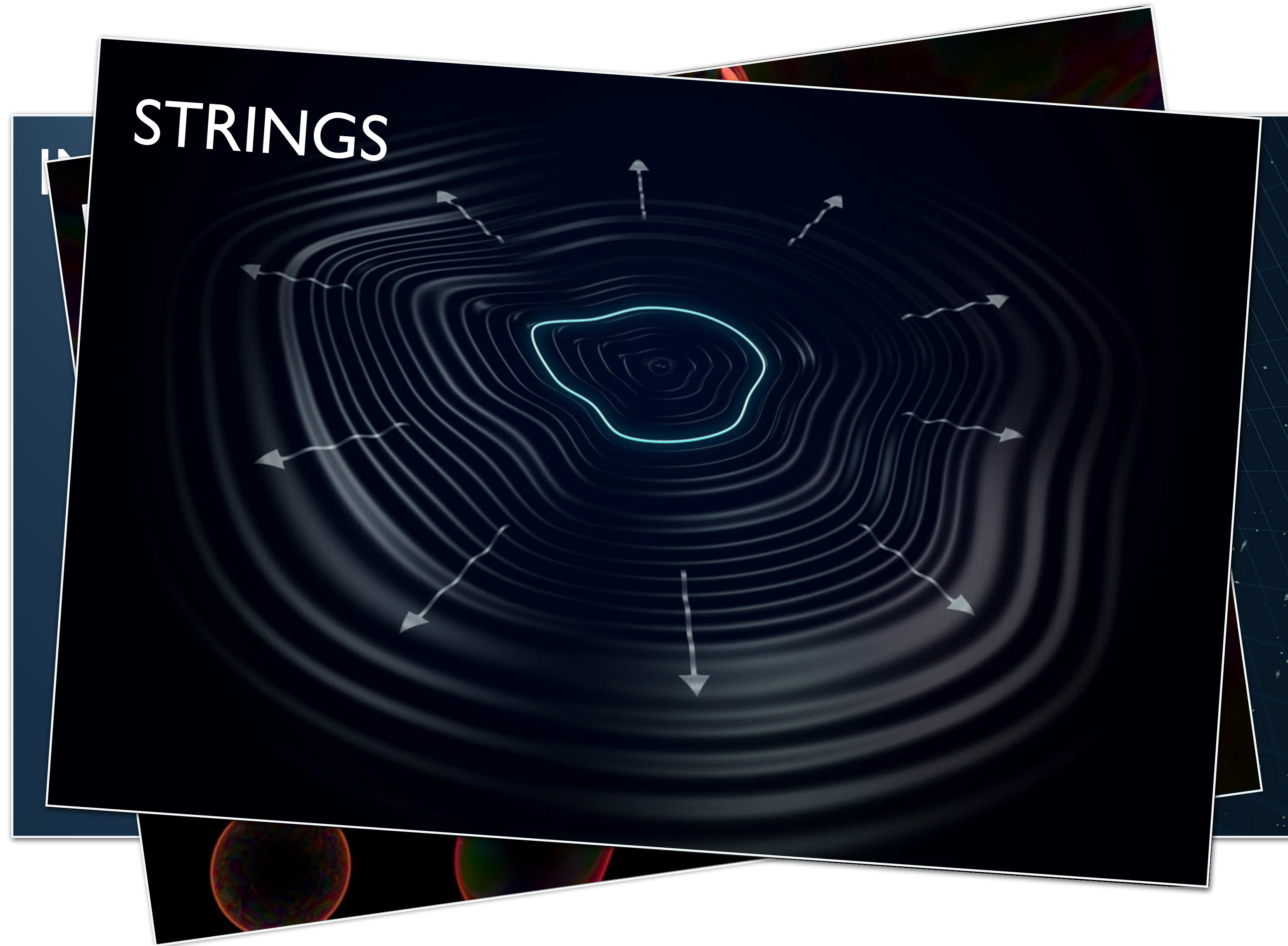
OTHER SUSPECTS



OTHER SUSPECTS



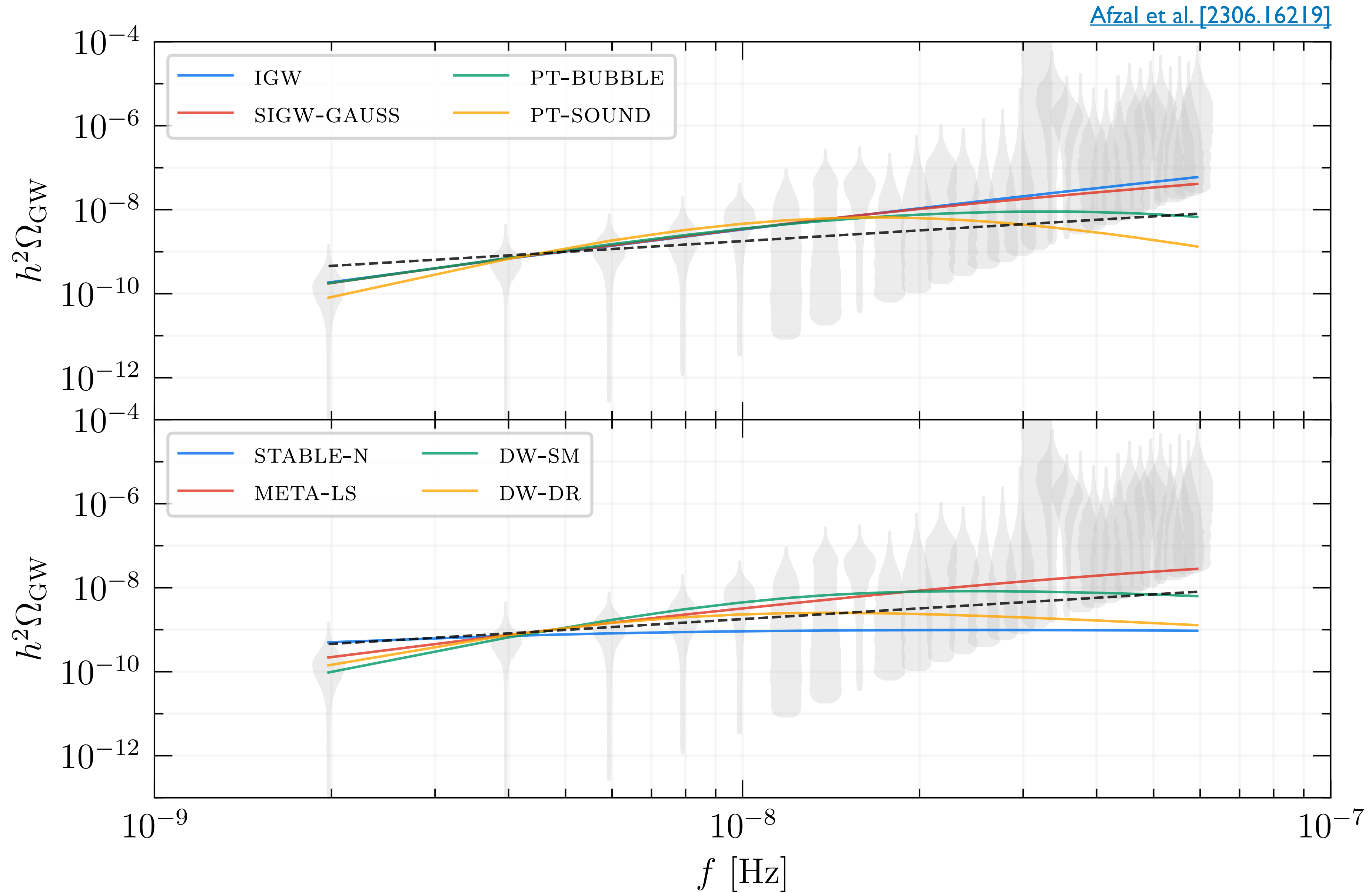
OTHER SUSPECTS



OTHER SUSPECTS



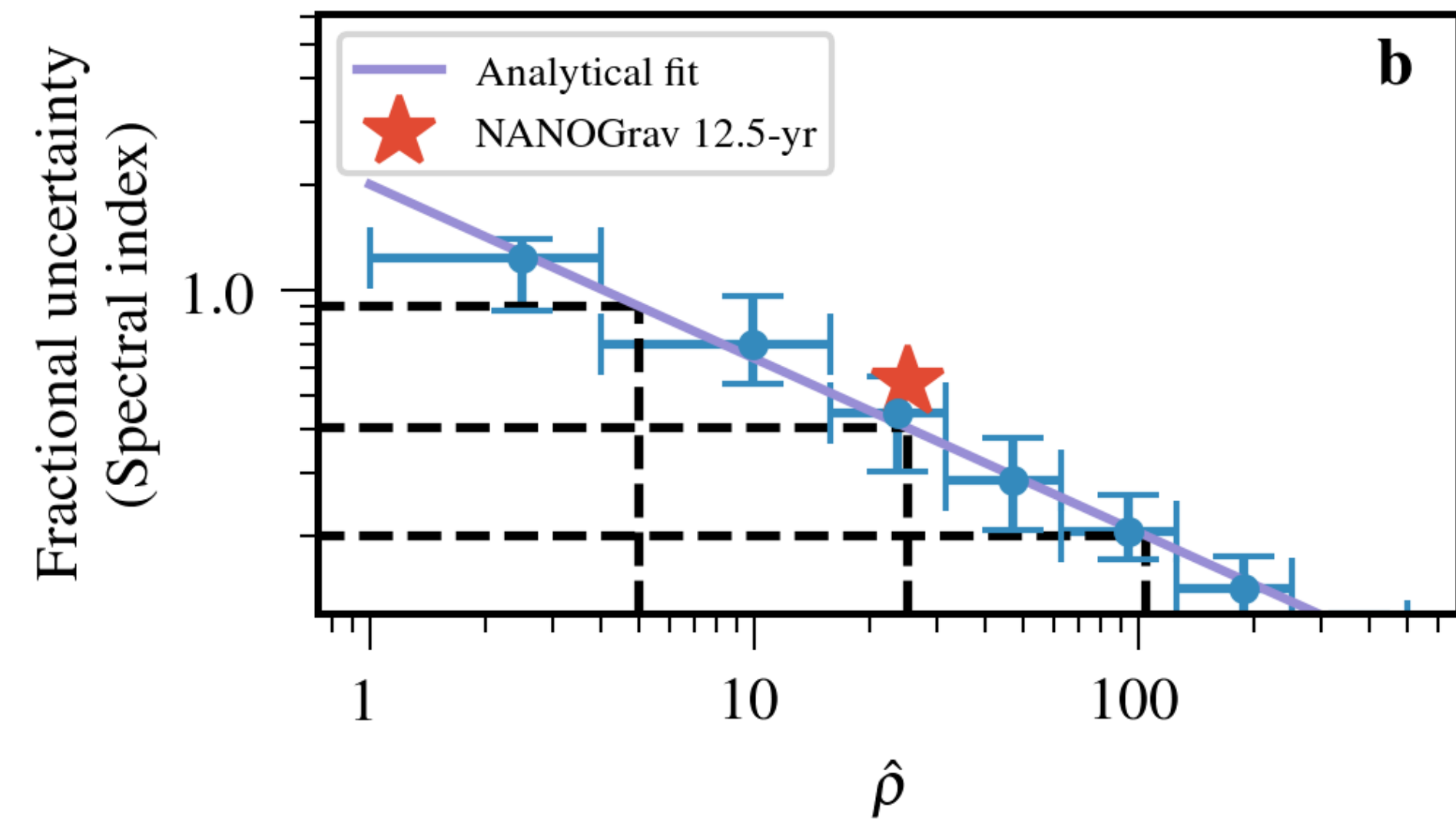
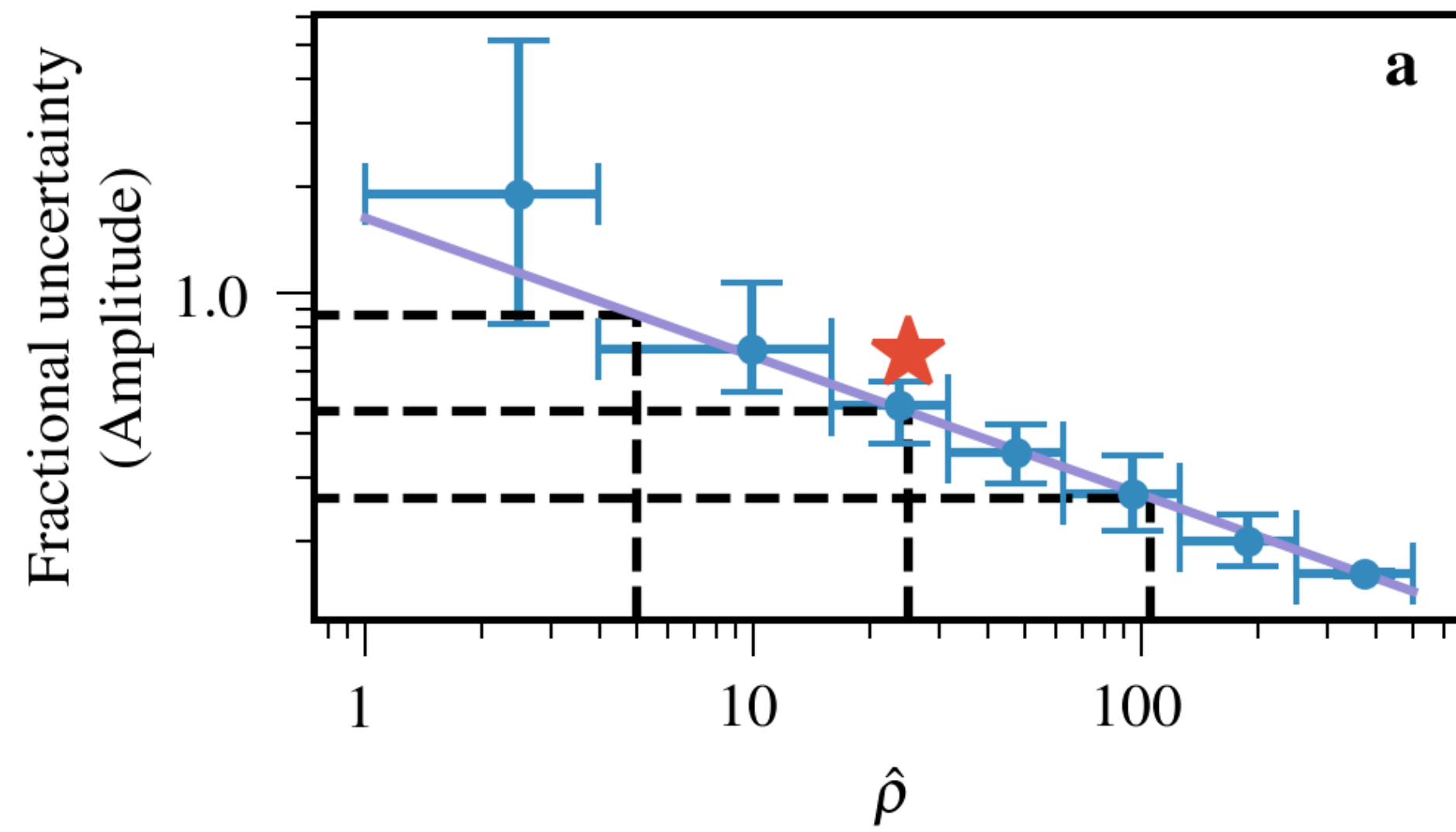
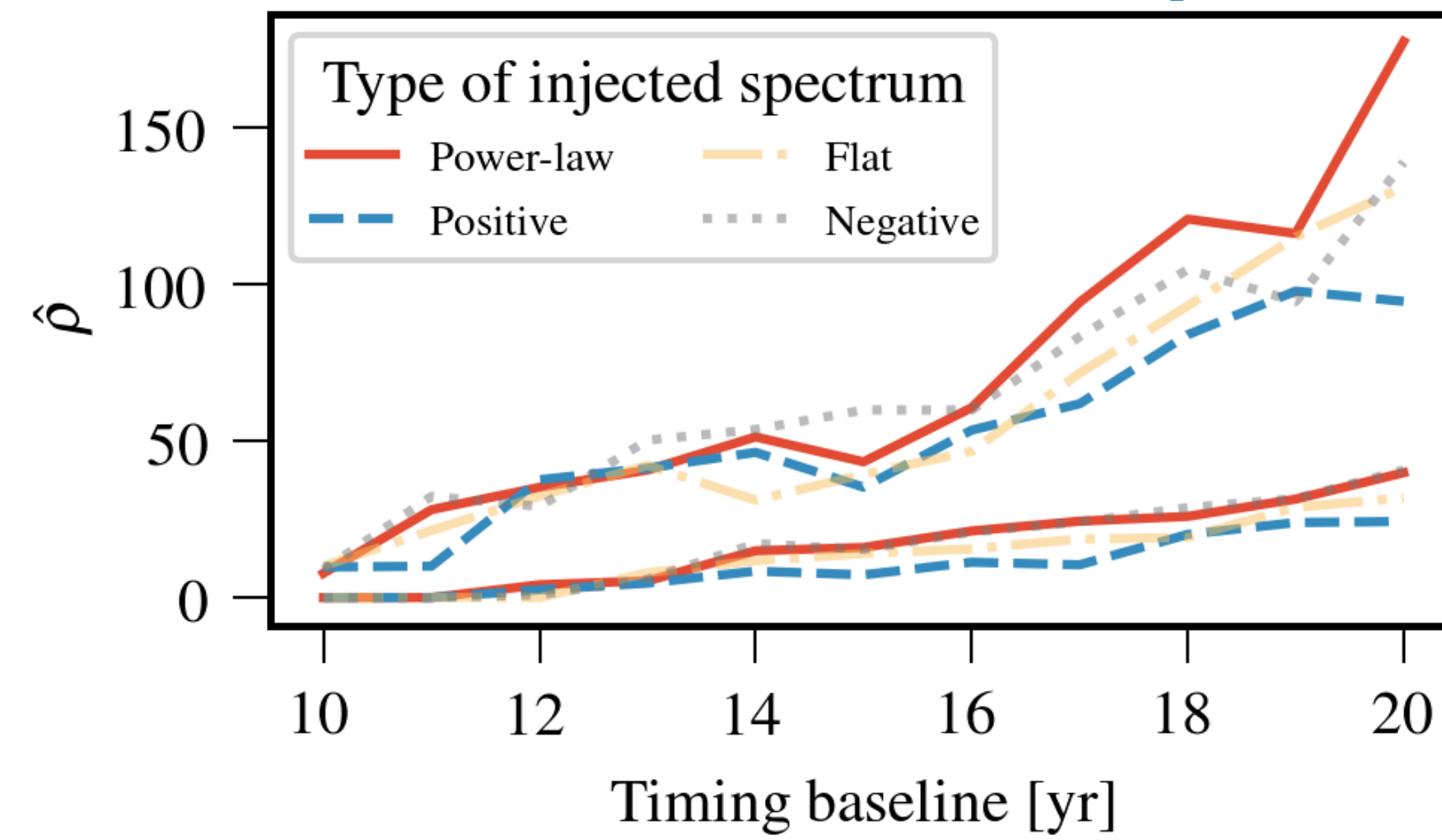
OTHER SUSPECTS



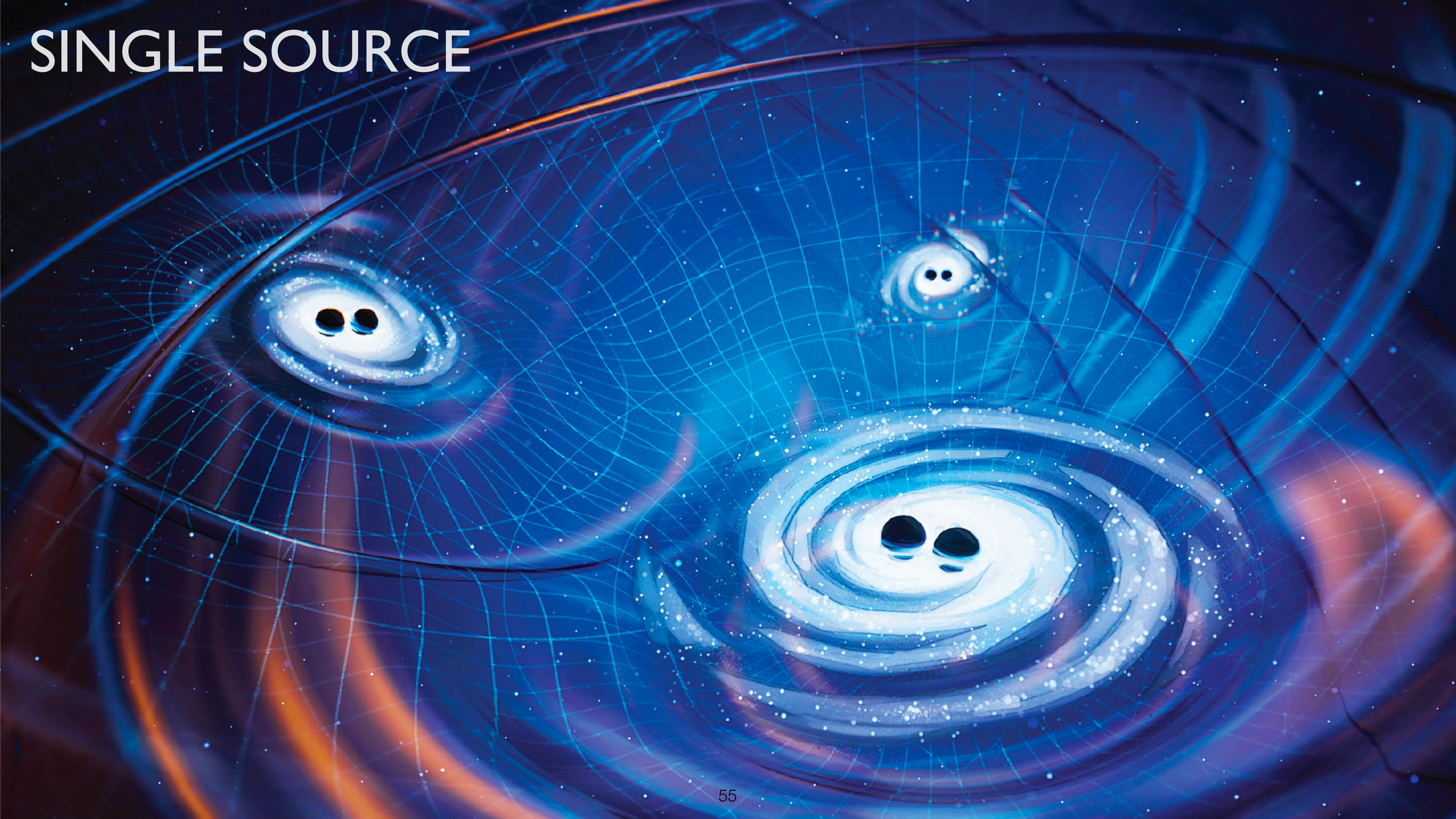
astrophysics or new physics?

SPECTRAL SHAPE

[Pol et al. \[2010.11950\]](#)

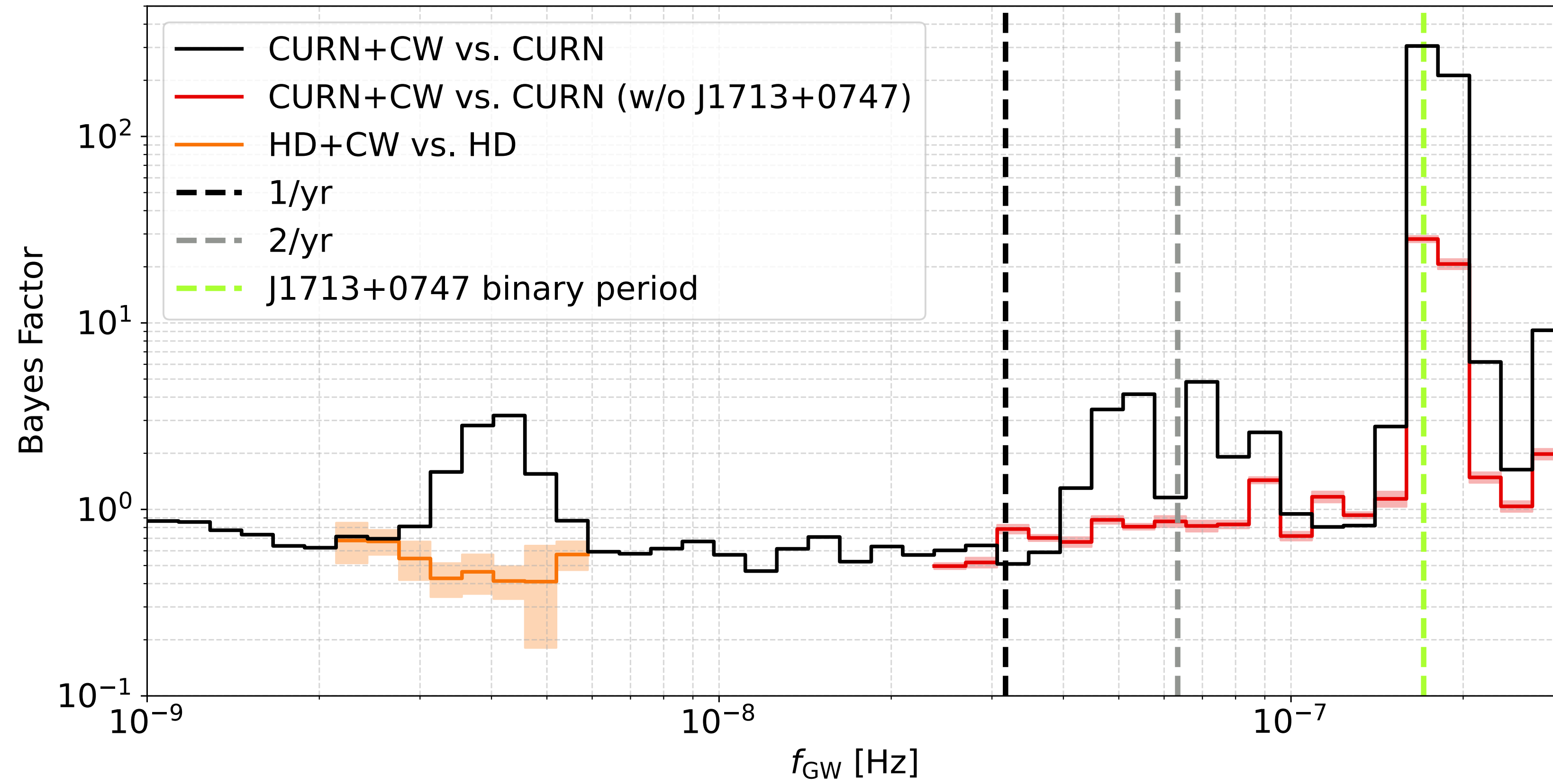


SINGLE SOURCE



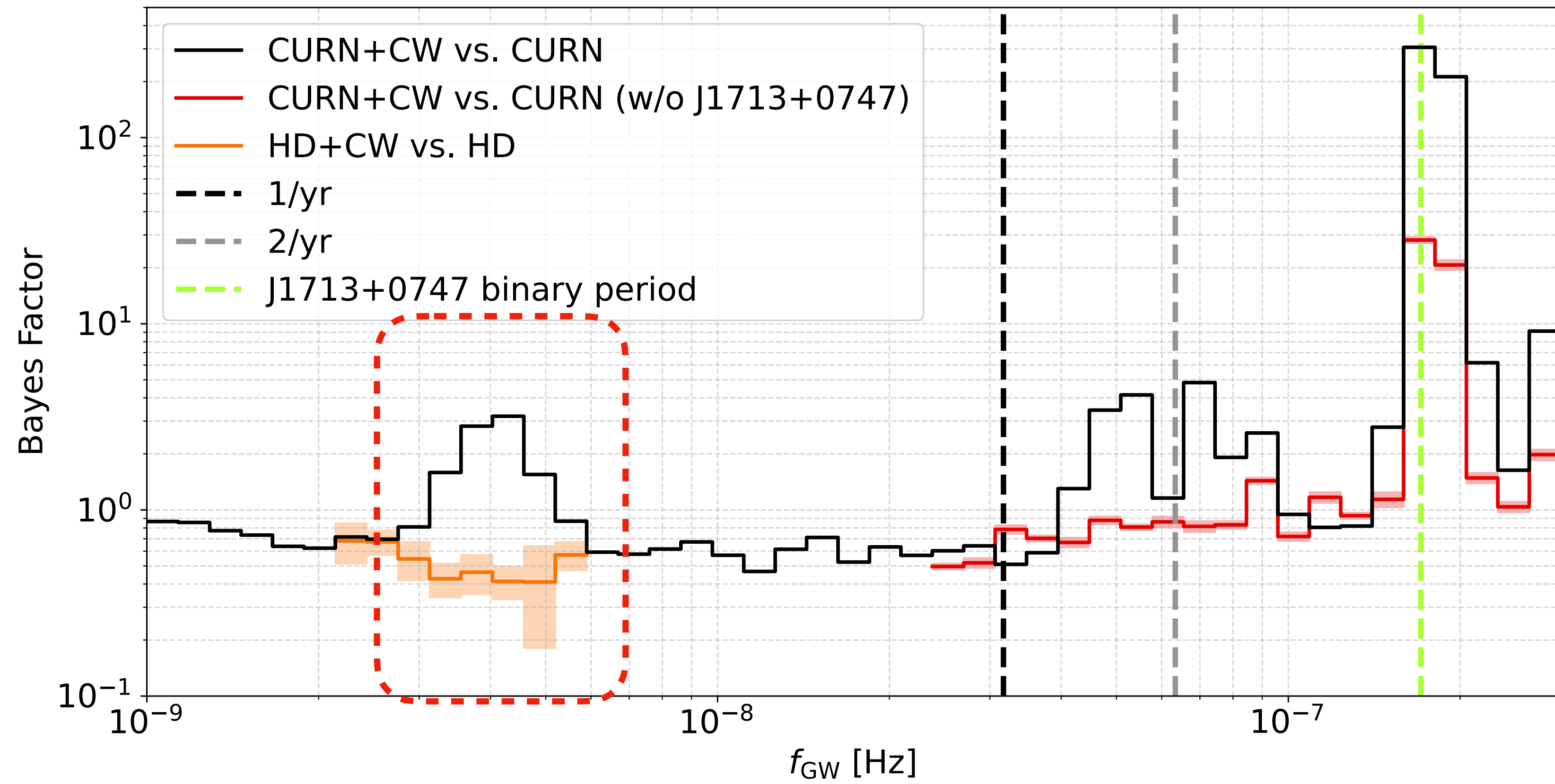
SINGLE SOURCE

[Agazie et al. \[2306.16222\]](#)

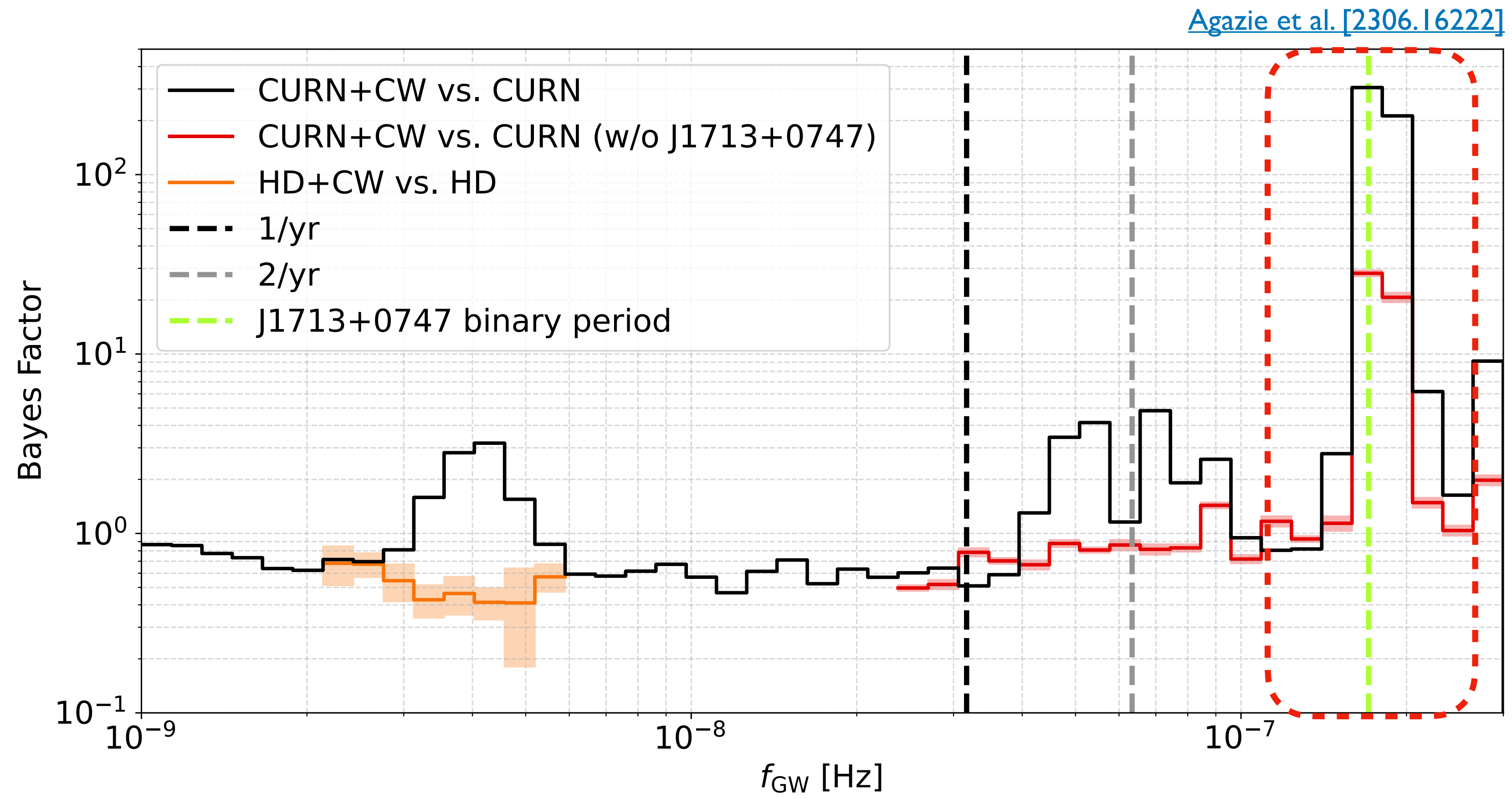


SINGLE SOURCE

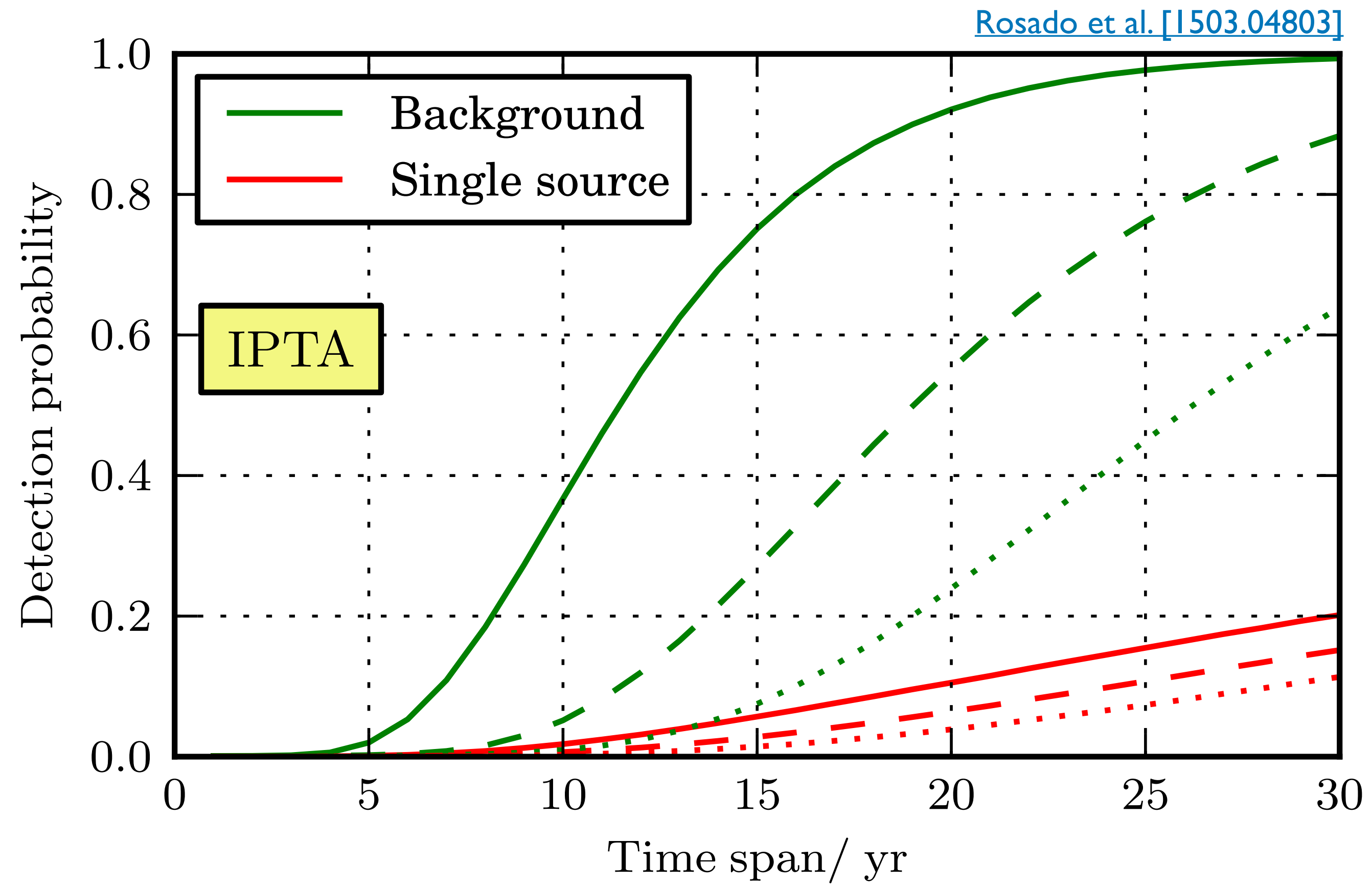
[Agazie et al. \[2306.16222\]](#)



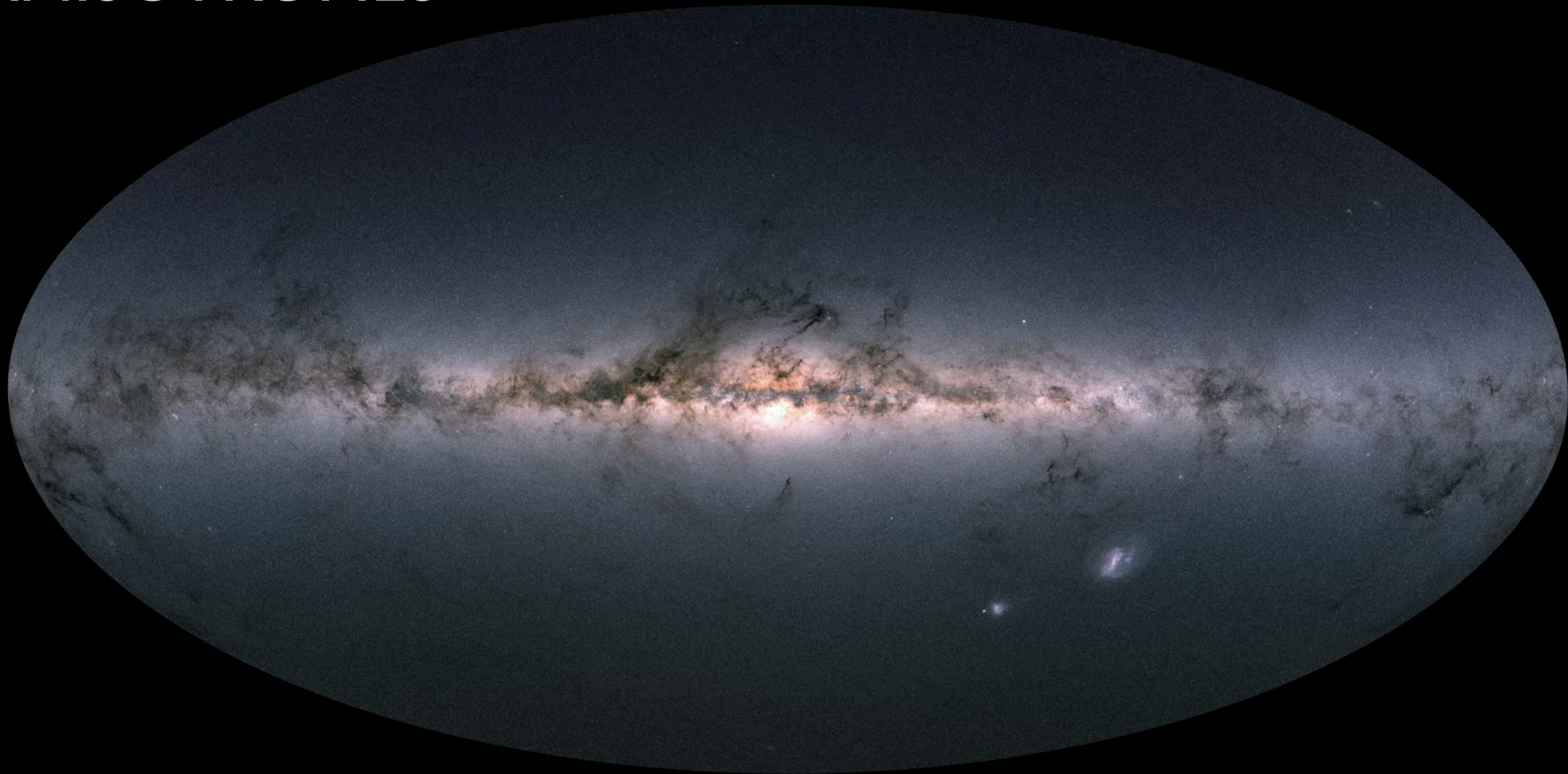
SINGLE SOURCE



SINGLE SOURCE

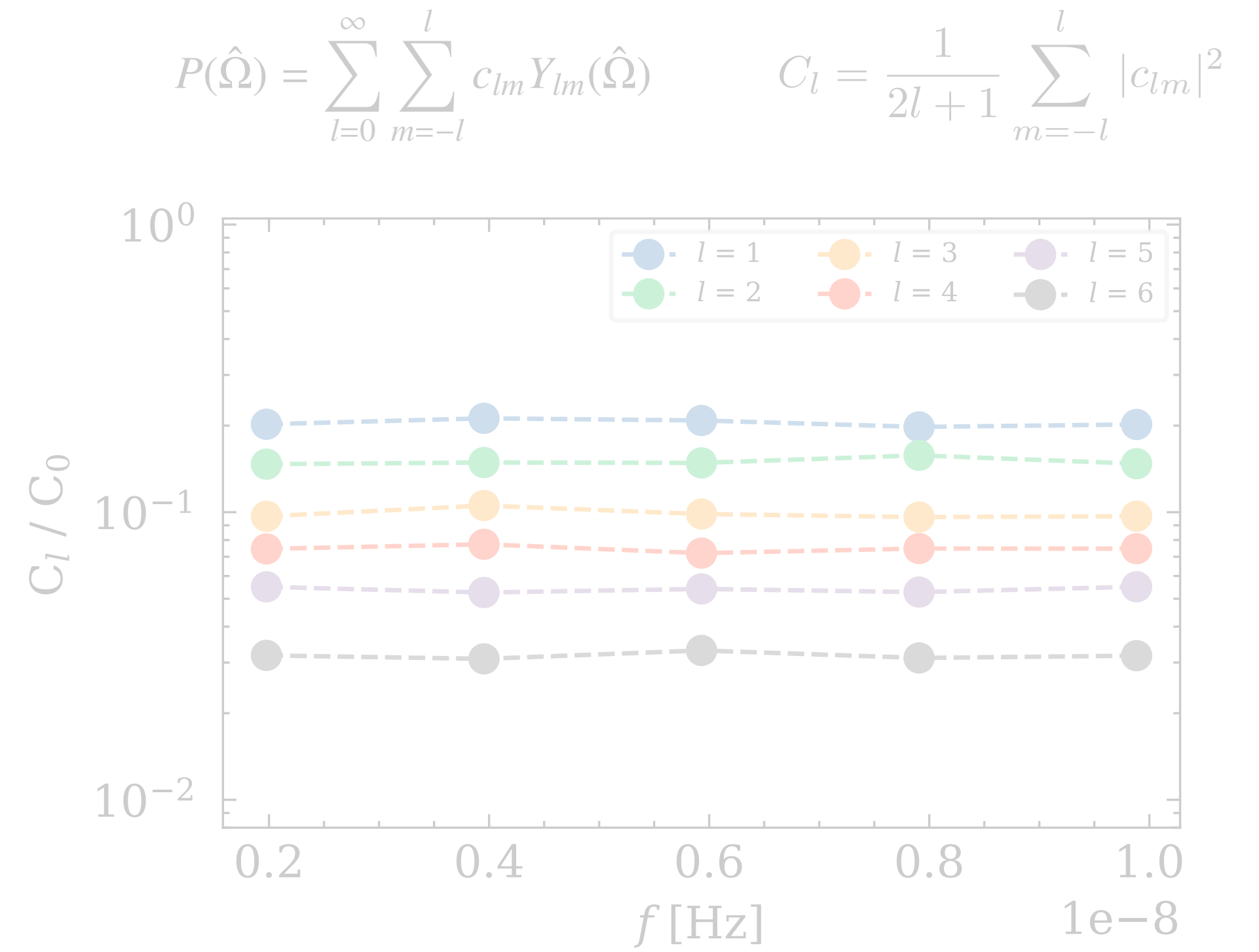
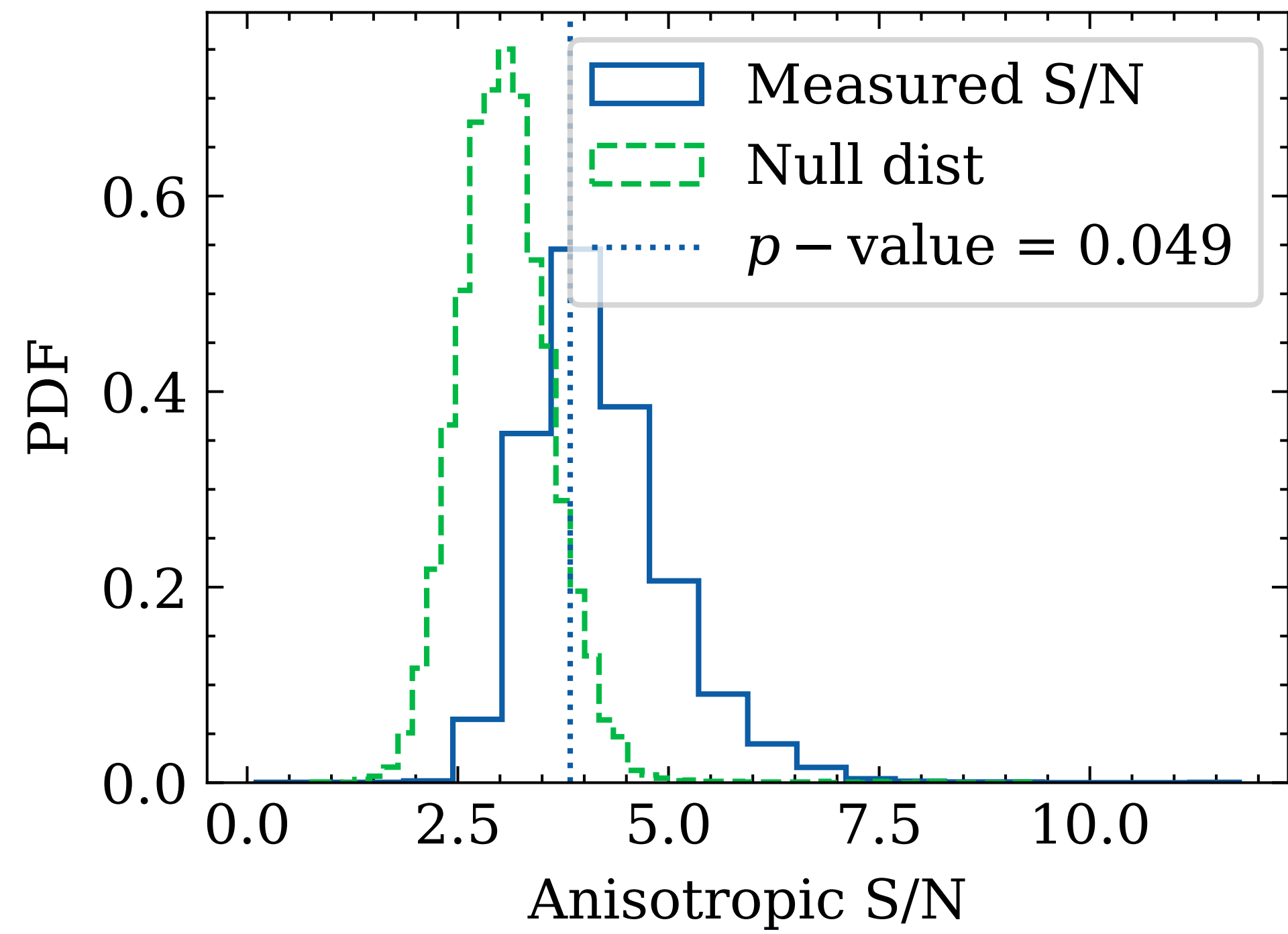


ANISOTROPIES

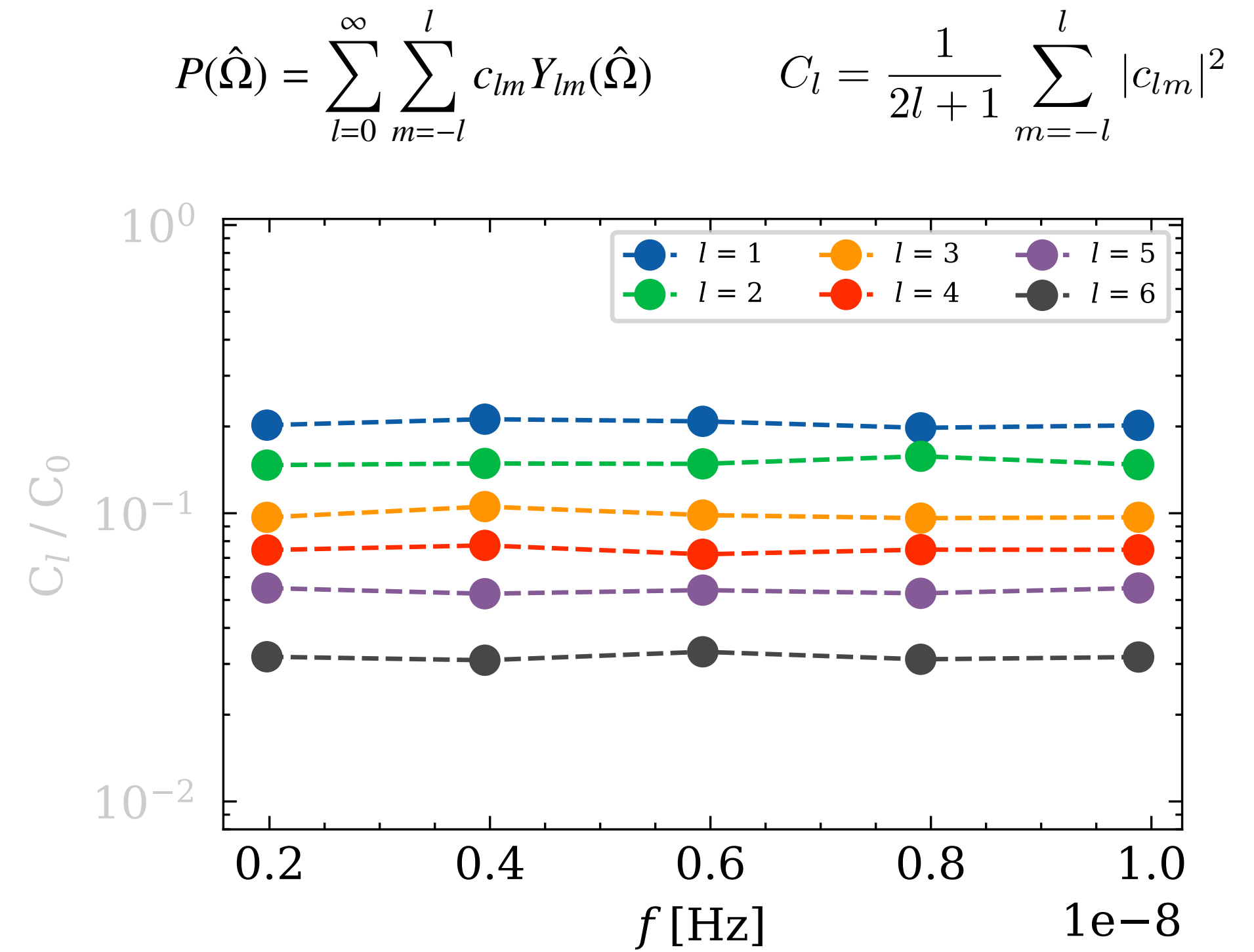
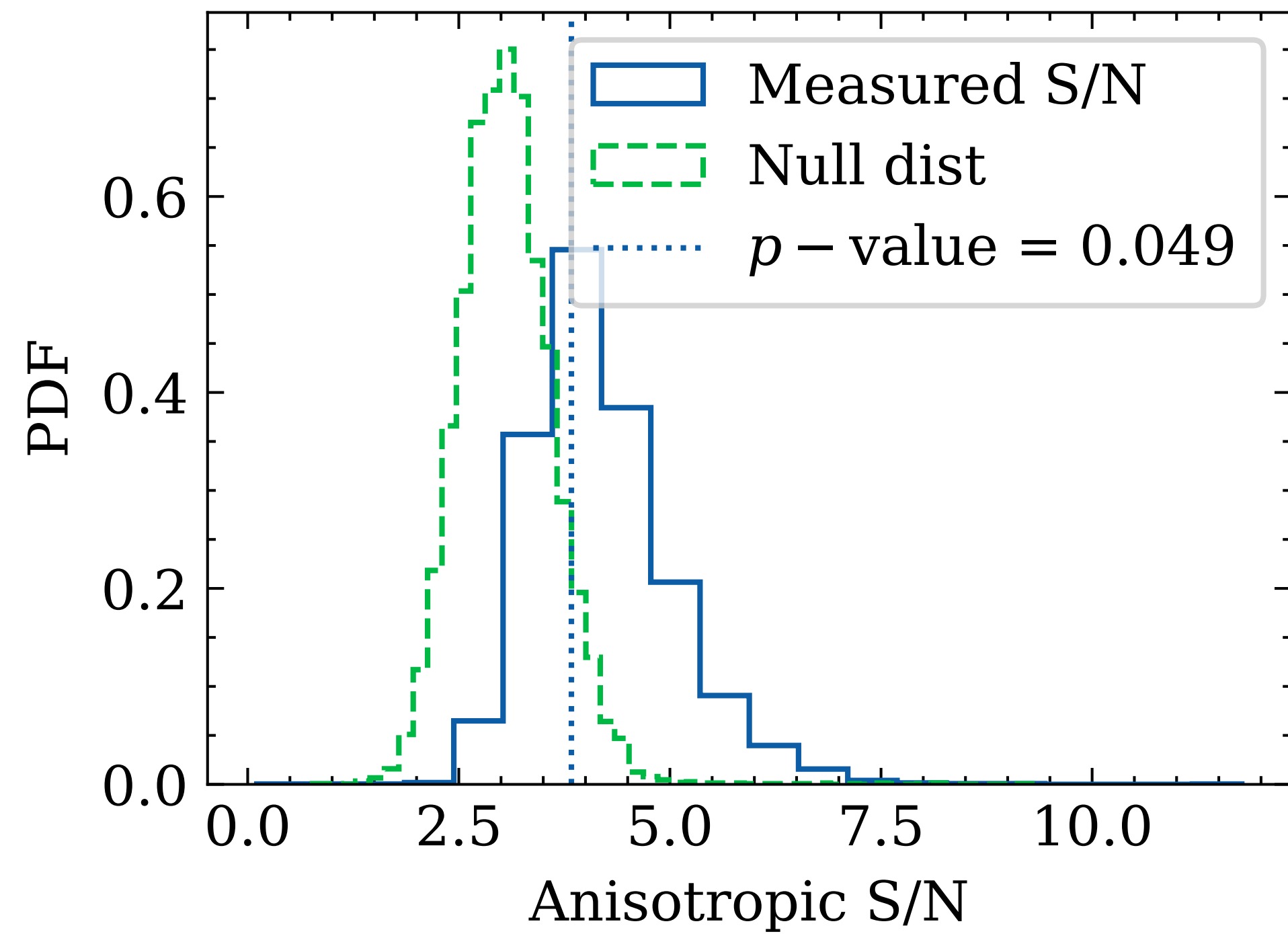


Credit: ESA/Gaia/DPAC

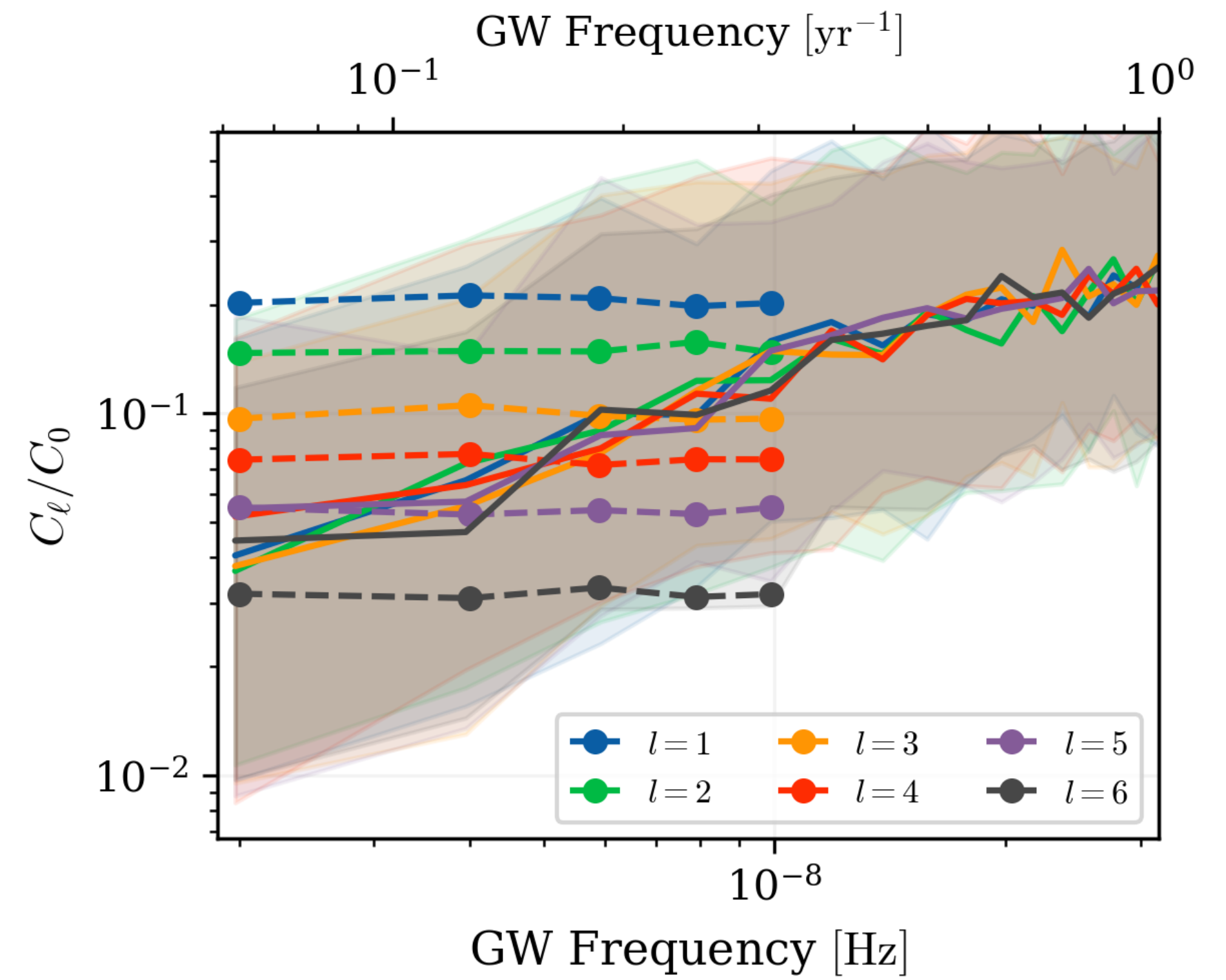
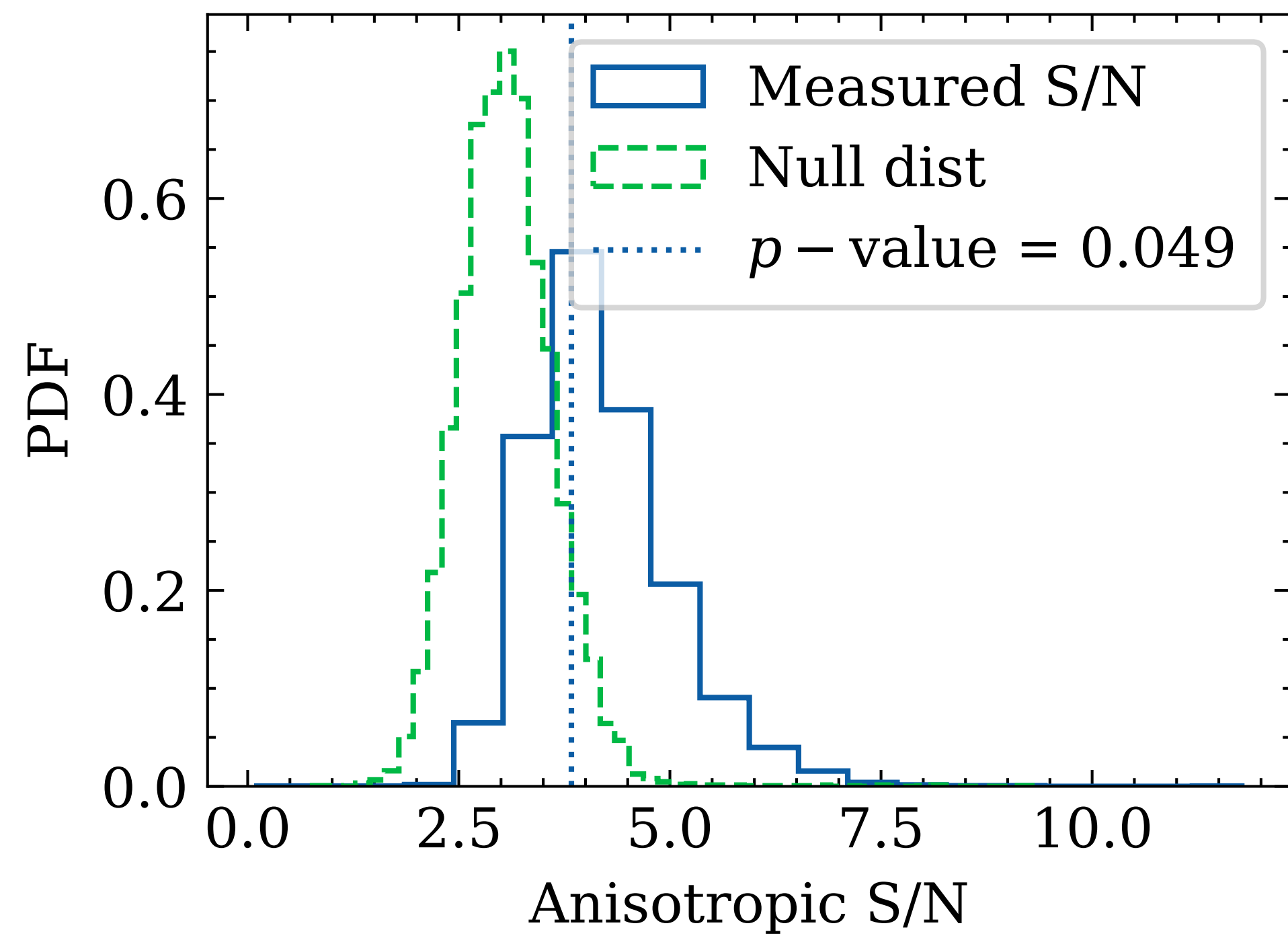
ANISOTROPIES



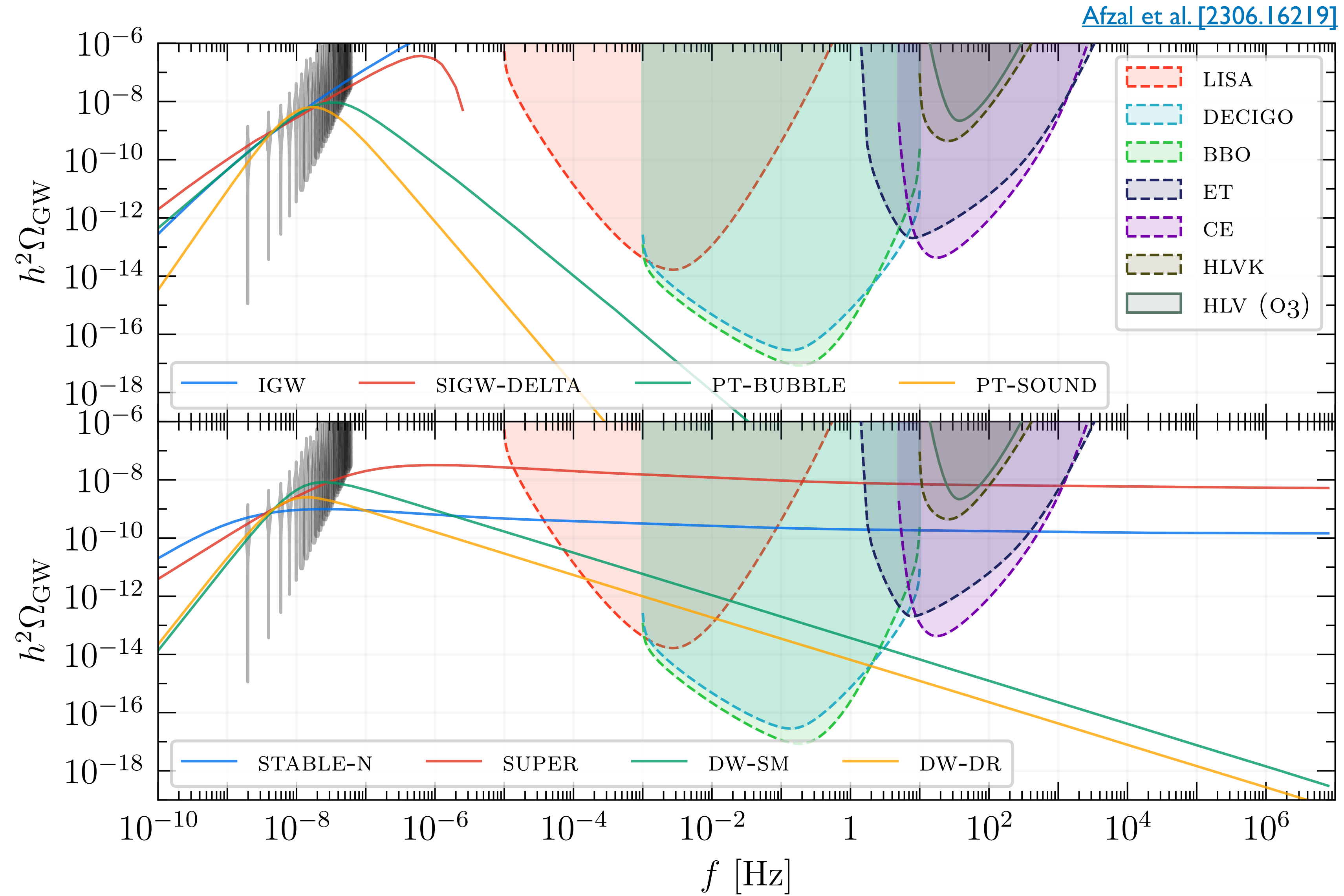
ANISOTROPIES



ANISOTROPIES

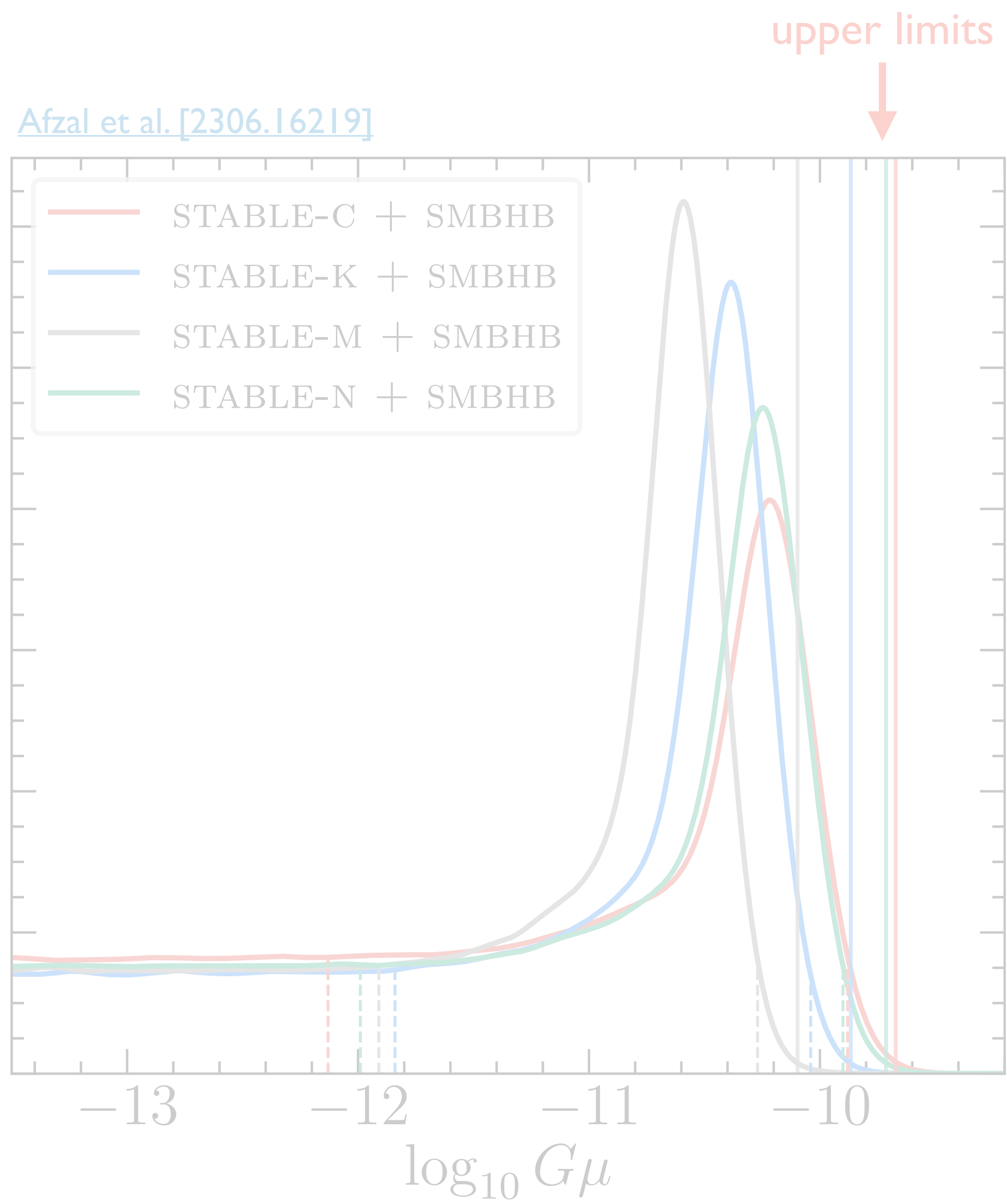
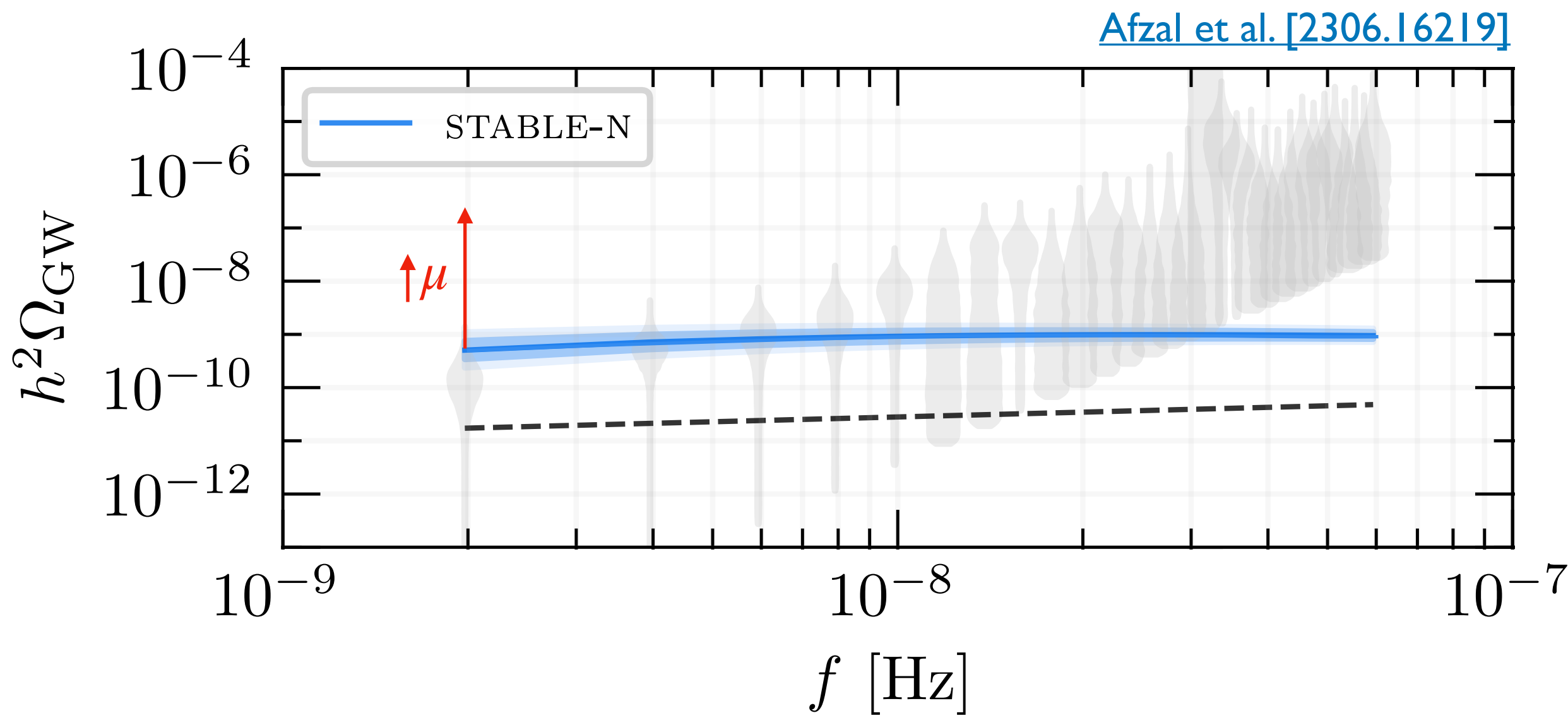
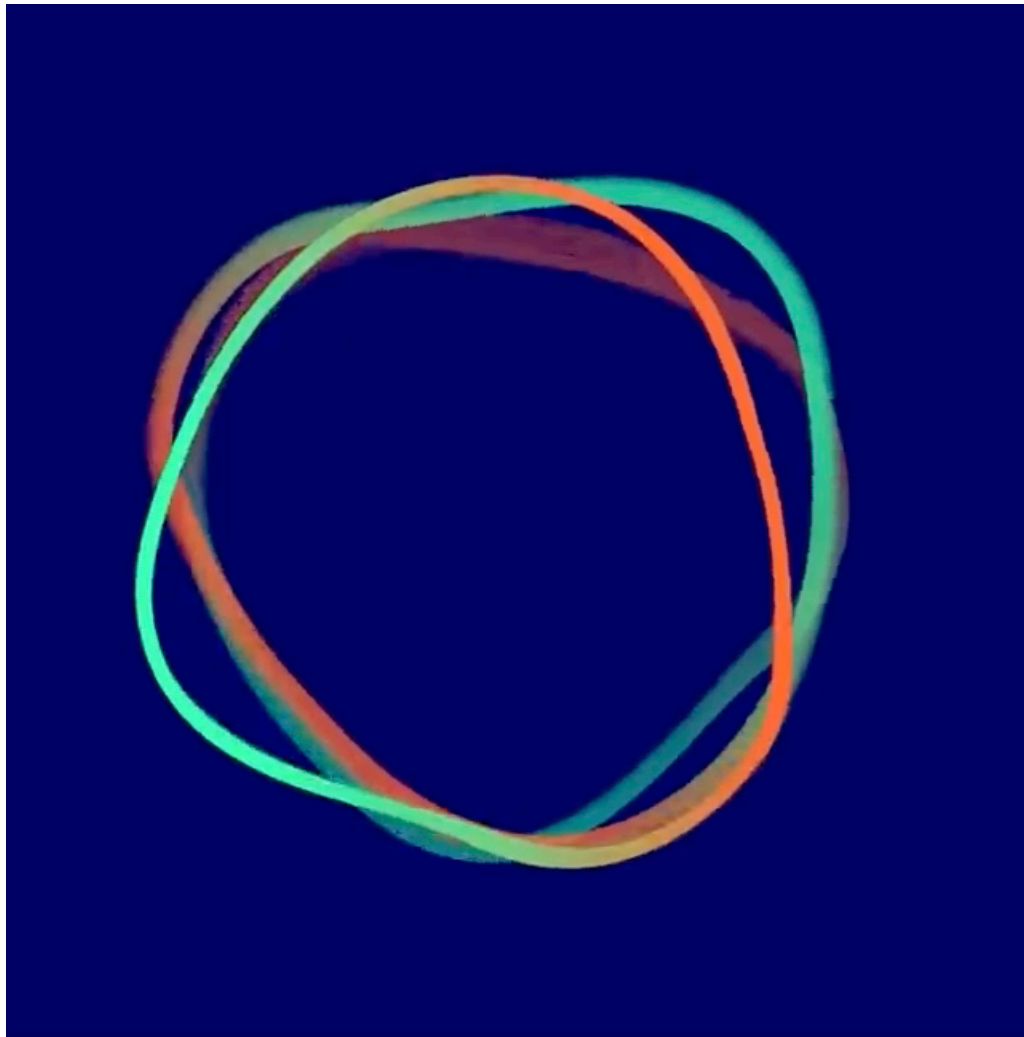


OTHER GWs EXPERIMENTS

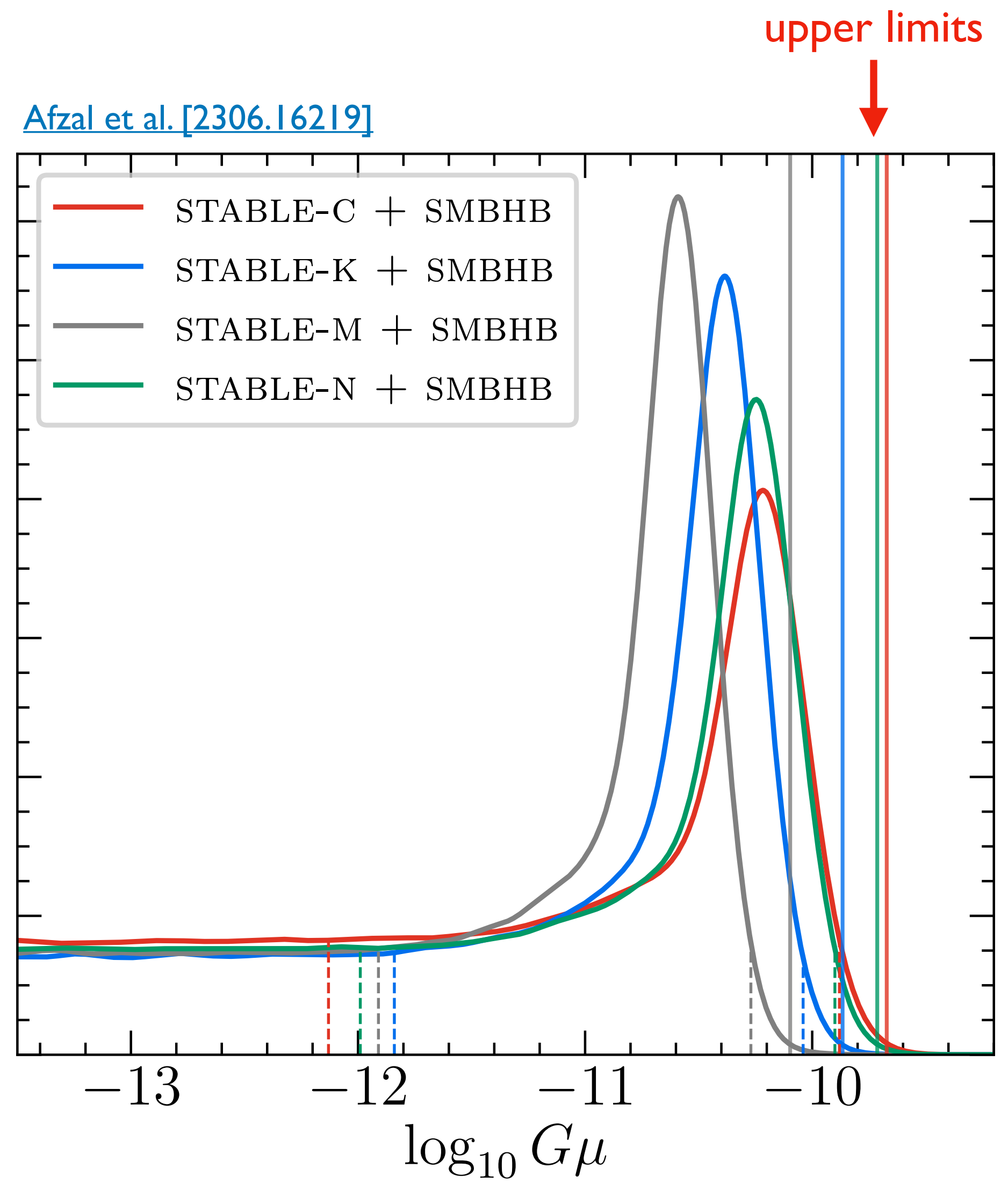
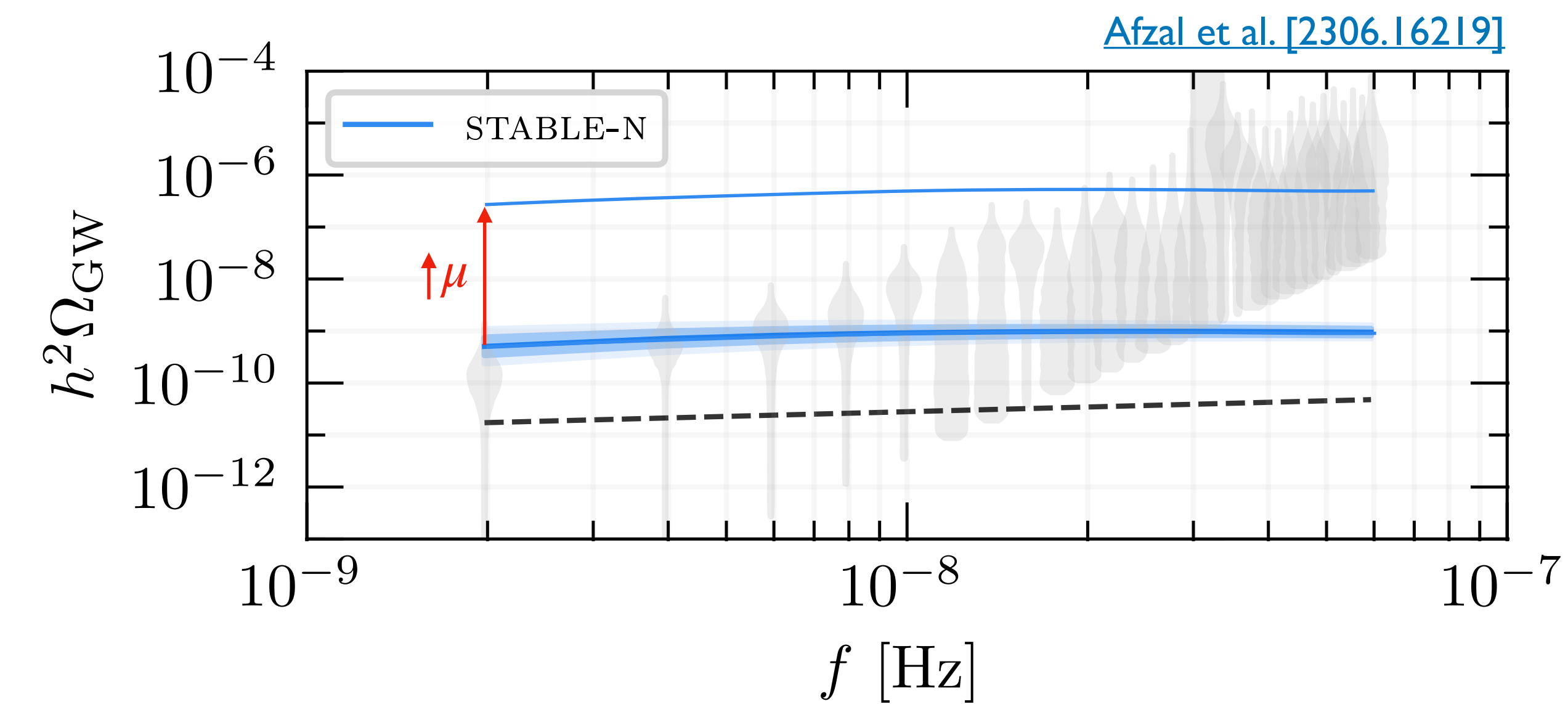
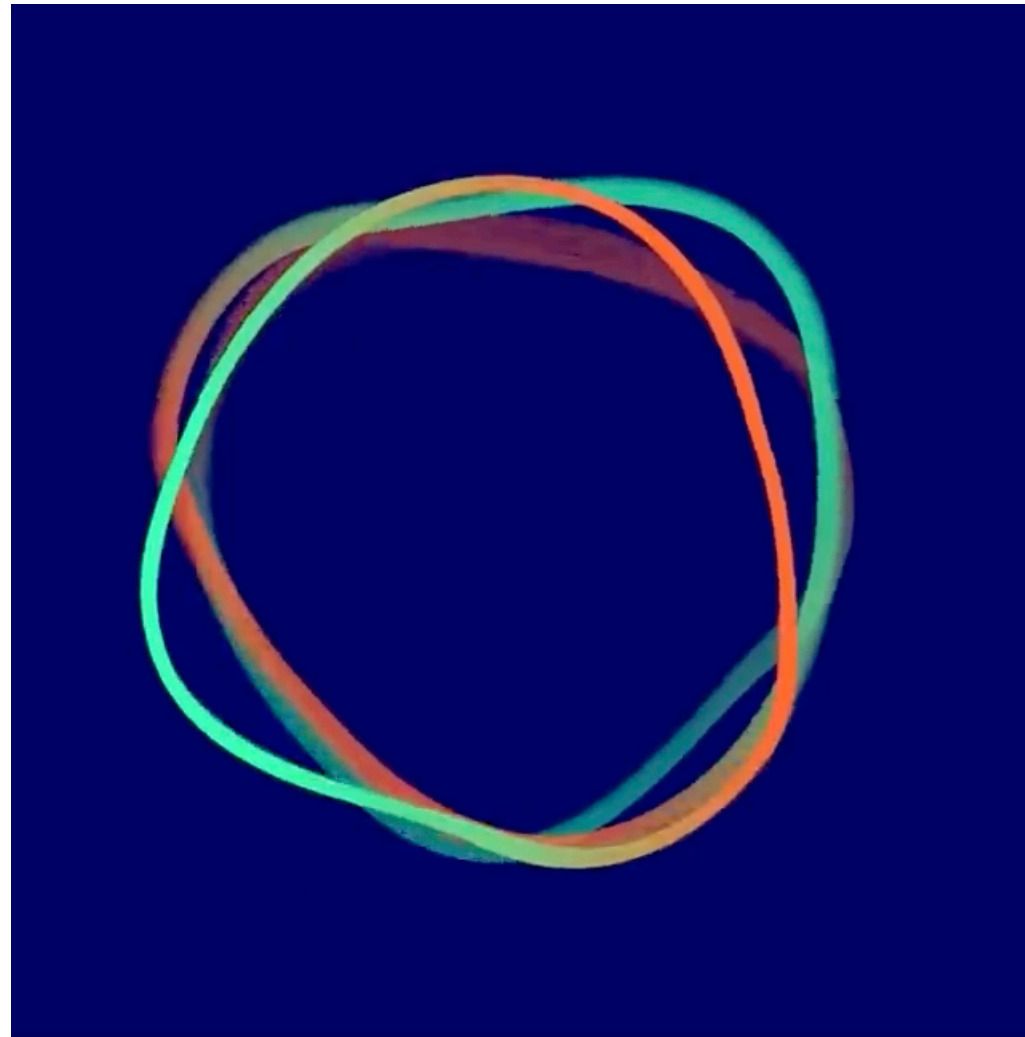


PTAs are not only a discovery tool!

COSMIC STRINGS



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$$\phi(\vec{x}, t) = \frac{\sqrt{2\rho_\phi}}{m_\phi} \hat{\phi}(\vec{x}) \cos(m_\phi t + \gamma(\vec{x}))$$

DM density

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DM mass

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gravitational signal

$$s(t) \sim \frac{G\rho_\phi}{m_\phi^3} \sin(2m_\phi t)$$

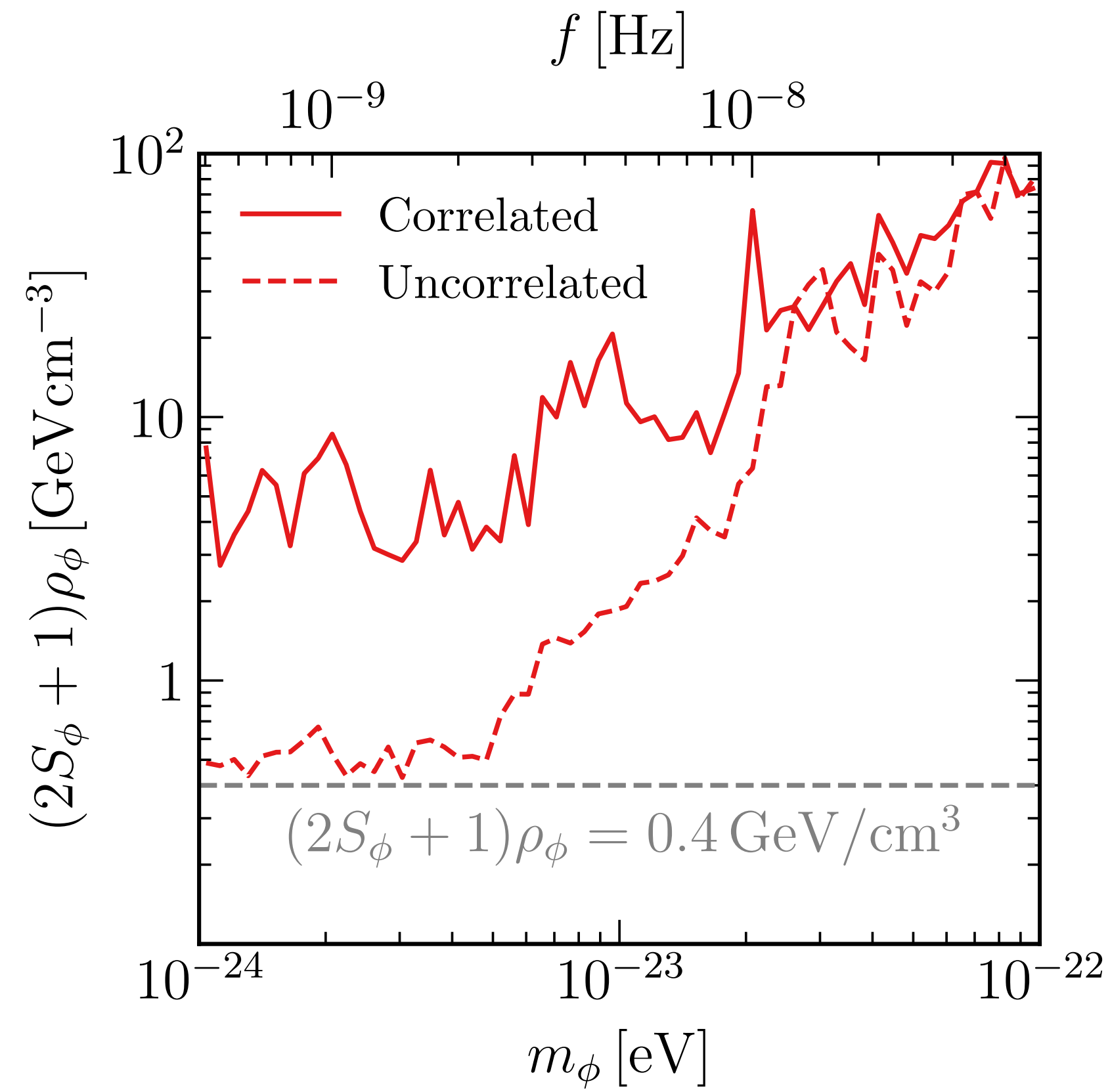
[Khmelnitsky, Rubakov \[1309.5888\]](#)

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[Afzal et al. \[2306.16219\]](#)

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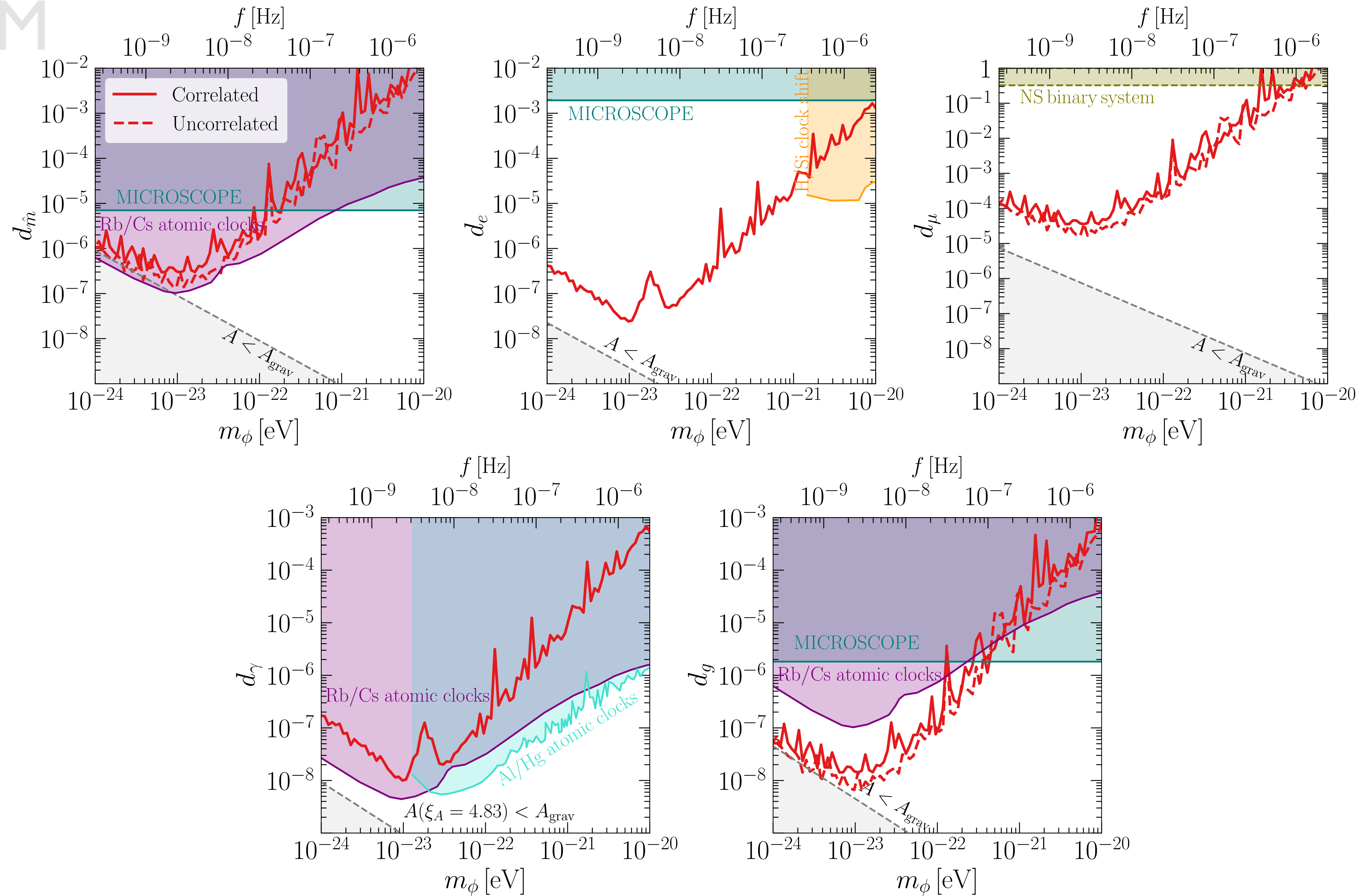
$$s(t) \sim \frac{G\rho_\phi}{m_\phi^3} \sin(2m_\phi t)$$

[Khmelnitsky, Rubakov \[1309.5888\]](#)

direct coupling signals

$$s(t) \sim d \frac{\sqrt{\rho_\phi}}{m_\phi^2 \Lambda} \sin(m_\phi t)$$

[Kaplan, AM, Trickle \[2205.06817\]](#)



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