

Introduction to Quantum Computing

Lecture 2

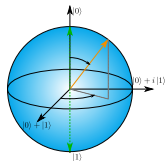
Stefan Kühn

DESY Summer Student Program, 01.08.2023

Recap of Lecture I

Qubits

- > Quantum mechanical two level systems $\mathcal{H} = \{|0\rangle, |1\rangle\}$
- > Can be in superposition
- > Qubits can be entangled



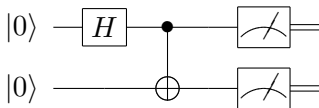
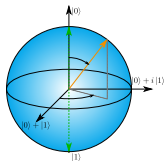
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Qubits

- > Quantum mechanical two level systems $\mathcal{H} = \{|0\rangle, |1\rangle\}$
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Quantum gates and quantum circuits

- > Quantum gates: unitary operations on a single/few qubits
- > Combining quantum gates we can express any unitary operation



Recap of Lecture I

Qubits

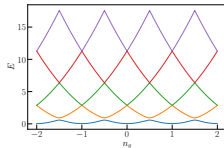
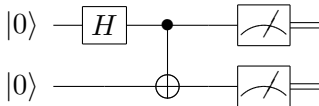
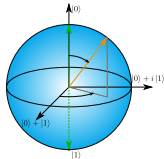
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Quantum gates and quantum circuits

- > Quantum gates: unitary operations on a single/few qubits
- > Combining quantum gates we can express any unitary operation

Realization of qubits in physical systems

- > Neutral atoms
- > Trapped ions
- > Superconducting qubits



Outline

Superdense coding

The Deutsch-Josza algorithm

Complexity theory

Hybrid quantum-classical algorithms

Superdense coding

Setting

- > Alice wants to communicate two classical bits to Bob



- > Classically we have to send (at least) two bits
- > Can we do better using quantum bits?

Superdense coding

Bell states

- > There are four Bell states

$$\begin{aligned} |\Phi^+\rangle &= \frac{1}{\sqrt{2}} (|0\rangle \otimes |0\rangle + |1\rangle \otimes |1\rangle), & |\Psi^+\rangle &= \frac{1}{\sqrt{2}} (|0\rangle \otimes |1\rangle + |1\rangle \otimes |0\rangle) \\ |\Phi^-\rangle &= \frac{1}{\sqrt{2}} (|0\rangle \otimes |0\rangle - |1\rangle \otimes |1\rangle), & |\Psi^-\rangle &= \frac{1}{\sqrt{2}} (|0\rangle \otimes |1\rangle - |1\rangle \otimes |0\rangle) \end{aligned}$$

- > All Bell states represent a maximally entangled pair of qubits
- > They form a basis for the Hilbert space of two qubits

Superdense coding

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- > All Bell states represent a maximally entangled pair of qubits
- > They form a basis for the Hilbert space of two qubits
- > Applying one of the Pauli gates X , Y or Z to one of the qubits transforms them into each other (up to global phases)

$$X \otimes \mathbb{1} |\Phi^+\rangle = |\Psi^+\rangle, \quad Y \otimes \mathbb{1} |\Phi^+\rangle = -i |\Psi^-\rangle, \quad Z \otimes \mathbb{1} |\Phi^+\rangle = |\Phi^-\rangle$$

Superdense coding

Superdense coding

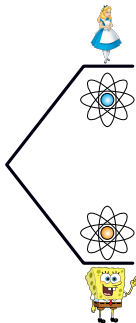
- > Prepare a Bell state, e.g. $|\Phi^+\rangle$



Superdense coding

Superdense coding

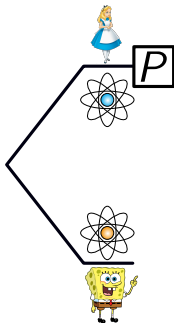
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Superdense coding

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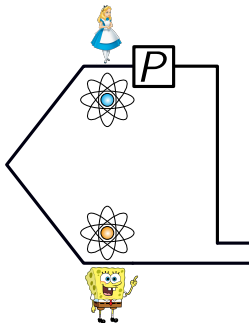
- > Prepare a Bell state, e.g. $|\Phi^+\rangle$
- > Alice and Bob each get one qubit of the Bell state
- > Alice applies a Pauli gate on her qubit depending on the classical bitstring she wants to send
 - $\mathbb{1}$ for 00
 - X for 01
 - Y for 10
 - Z for 11



Superdense coding

Superdense coding

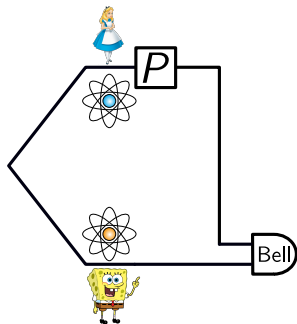
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- > Alice sends her qubit to Bob



Superdense coding

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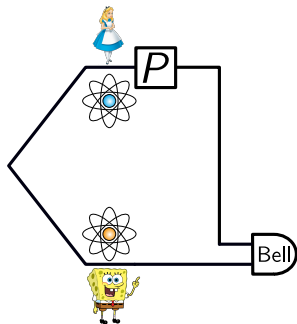
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 - $\mathbb{1}$ for 00
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- > Alice sends her qubit to Bob
- > Bob performs a Bell measurement which tells him Alice's bit string
- ⇒ We only need a single qubit!



Superdense coding

Superdense coding

- > The entangled Bell pair is a resource
- > It can be shared long before the communication should take place

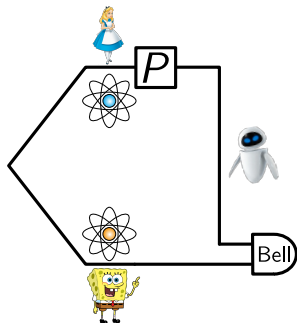
Superdense coding

Superdense coding

- > The entangled Bell pair is a resource
- > It can be shared long before the communication should take place
- > Communication is secure
 - If an eavesdropper gets access to Alice's qubit they only have one of the qubits
 - As we have seen the reduced density operator for a single qubit of Bell state is maximally mixed

$$\rho_1 = \frac{1}{2} (|0\rangle\langle 0| + |1\rangle\langle 1|)$$

- ⇒ The eavesdropper cannot infer information about the classical bit string



2.

Superdense coding

The Deutsch-Josza algorithm

Complexity theory

Hybrid quantum-classical algorithms

The Deutsch-Josza algorithm

Setting

- > Given: a function $f : \mathbb{Z}_2^n \rightarrow \mathbb{Z}_2$ that is promised to be constant or balanced
- > Task: find out if f is constant or balanced

The Deutsch-Josza algorithm

Setting

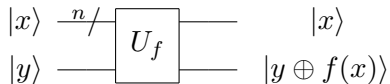
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- > Classical computer: try more than half of the possible inputs
 $\Rightarrow \frac{1}{2} \times 2^n + 1 = 2^{n-1} + 1$ function calls

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The Deutsch-Josza algorithm

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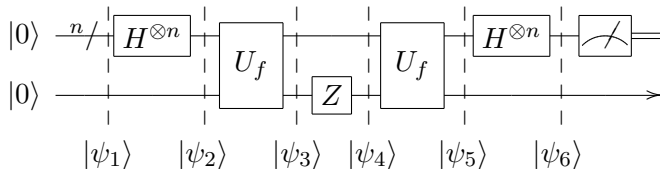
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 $\Rightarrow \frac{1}{2} \times 2^n + 1 = 2^{n-1} + 1$ function calls
- > Let us assume we have a unitary $U_f |x\rangle |y\rangle = |x\rangle |y \oplus f(x)\rangle$



- > U_f is called an **oracle**

The Deutsch-Josza algorithm

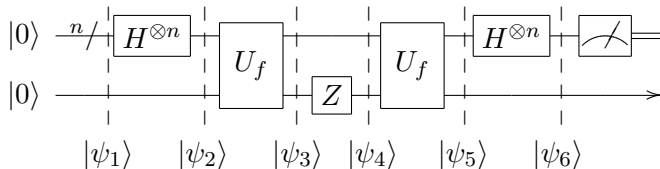
Deutsch-Josza algorithm



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The Deutsch-Josza algorithm

Deutsch-Josza algorithm

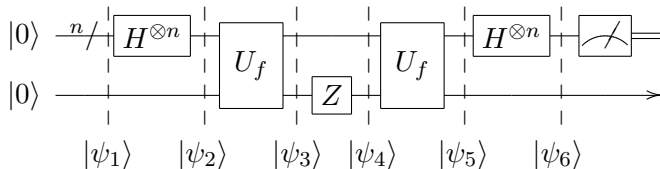


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Deutsch-Josza algorithm



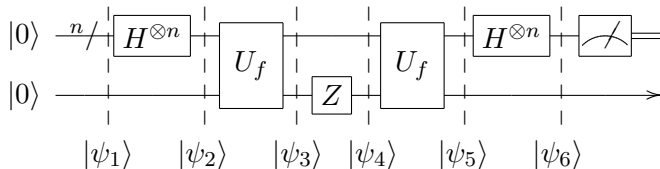
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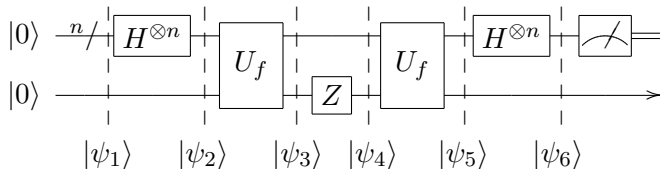
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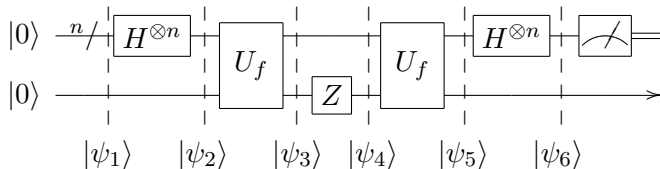
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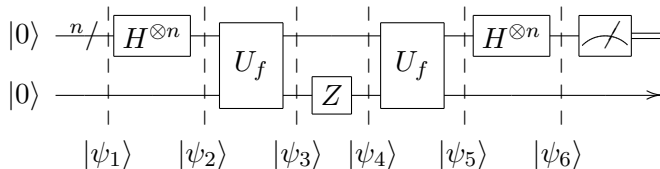
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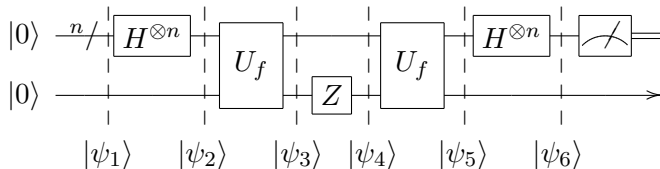
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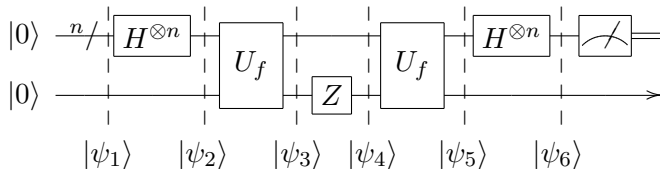
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The Deutsch-Josza algorithm

Deutsch-Josza algorithm

- > Quantum algorithm allows for deciding whether f is balanced or not with **two calls to the oracle** (independent of n)
- > Query the oracle in **superposition**
- > Constructive **interference** (destructive interference) yields an unity (zero) amplitude in the constant (balanced) case

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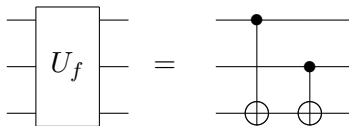
The Deutsch-Josza algorithm needs exponentially fewer calls to the oracle than the classical algorithm.

The Deutsch-Josza algorithm

Deutsch-Josza algorithm on quantum hardware

- > Example for $n = 2$ input bits and the following Boolean function

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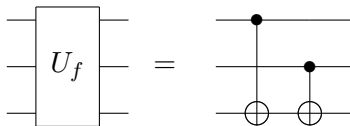


The Deutsch-Josza algorithm

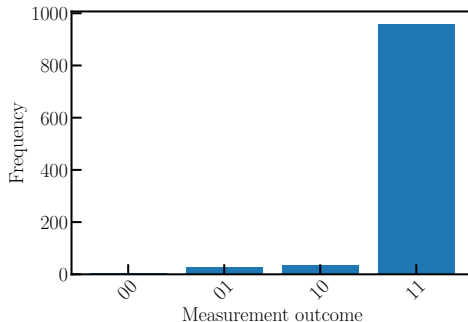
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> Results on actual quantum hardware (ibmq_lagos)



3.

Superdense coding

The Deutsch-Josza algorithm

Complexity theory

Hybrid quantum-classical algorithms

Complexity theory

Solving problems on a quantum computer

- > Many more known quantum algorithms that (might) perform better than the best known classical algorithms
 - ▶ Shor's factoring algorithm
 - ▶ Grover's search algorithm
 - ▶ HHL algorithm for linear equations
 - ▶ Quantum Simulation
 - ▶ Bernstein–Vazirani algorithm
 - ▶ ...
- > Exploiting quantum features such as superposition and entanglement these algorithms can outperform the best known classical algorithms

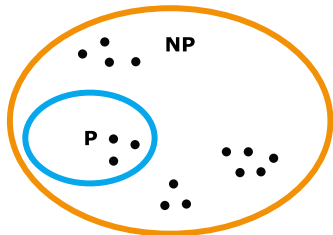
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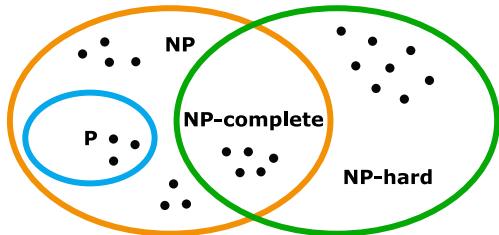
Which problems can be solved efficiently on quantum computers?

Complexity theory



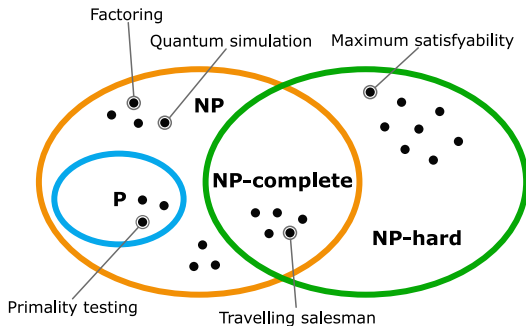
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- > NP: decision problems solvable by a non-deterministic Turing machine in polynomial time
 - “Hard problems”
 - **Solution can be checked** on a deterministic Turing machine in **polynomial time**

Complexity theory



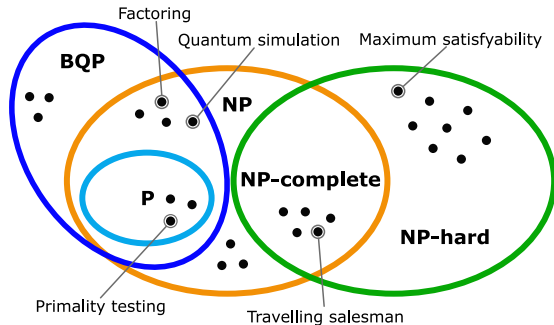
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Complexity theory



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 - “Hard problems”
 - **Solution can be checked** on a deterministic Turing machine in **polynomial time**

BQP (bounded-error quantum polynomial time):

- > Decision problems solvable by a quantum computer in **polynomial time**
- > Quantum equivalent to P, “easy problems”

Complexity theory

The Church-Turing thesis

All **physically reasonable models of computation** have the **same set of computable functions**.

⇒ Quantum computers cannot compute functions that are uncomputable on a classical computer

Complexity theory

The Church-Turing thesis

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The extended Church-Turing thesis

All **physically reasonable models of computation differ** in complexity by **at most polynomial factors**.

- ⇒ Extended Church-Turing thesis would no longer hold if **quantum supremacy** is demonstrated

Quantum supremacy using a programmable superconducting processor

<https://doi.org/10.1038/s41586-019-1666-5>

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RESEARCH

QUANTUM COMPUTING

Quantum computational advantage using photons

Han-Sen Zhong^{1,2*}, Hui Wang^{1,2*}, Yu-Hao Deng^{1,2*}, Ming-Cheng Chen^{1,2*}, Li-Chao Peng^{1,2}, Yi-Han Luo^{1,2}, Jian Qin^{1,2}, Dian Wu^{1,2}, Xing Ding^{1,2}, Yi Hu^{1,2}, Peng Hu³, Xiao-Yan Yang³, Wei-Jun Zhang³, Hao Li³, Yuxuan Li⁴, Xiao Jiang^{1,2}, Lin Gan⁴, Guangwen Yang⁴, Lixing You³, Zhen Wang³, Li Li^{1,2}, Nai-Le Liu^{1,2}, Chao-Yang Lu^{1,2}, Jian-Wei Pan^{1,2}

4.

Superdense coding

The Deutsch-Josza algorithm

Complexity theory

Hybrid quantum-classical algorithms

Hybrid quantum-classical algorithms

Current NISQ devices

- > Small or intermediate scale
- > Considerable amount of noise
- > Only shallow circuits can be executed faithfully/no error correction
- > Quantum advantage demonstrated

Hybrid quantum-classical algorithms

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Solving “useful” problems

- > Large number of qubits
- > Deep circuits
- > Quantum error correction necessary
- > So far only proof of principle demonstrations

Hybrid quantum-classical algorithms



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Solving “useful” problems

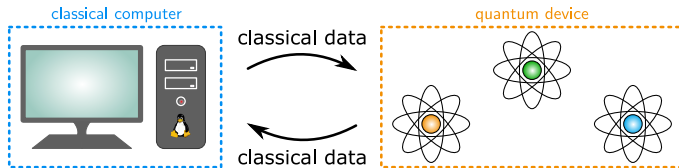
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How can we utilize existing quantum hardware in a beneficial way?

Hybrid quantum-classical algorithms

Hybrid quantum-classical algorithms

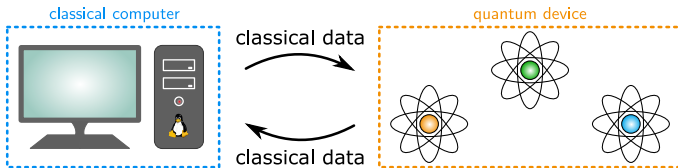
- > Combine classical and quantum devices
- > Rely on classical computing where possible
- > Use the quantum device as a coprocessor
 - Tackle the classically hard/intractable part of the problem
 - Feed the classical data obtained from a measurement back to the classical computer



Hybrid quantum-classical algorithms

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Even modest quantum hardware can yield advantages

Hybrid quantum-classical algorithms

Hybrid quantum-classical algorithms

- > Focus on optimization problems

$$\min_{\vec{\theta}} \mathcal{C}(\vec{\theta}), \quad \vec{\theta} = \mathbb{R}^n$$

- > Solve them iteratively using a parametric ansatz
 - Quantum coprocessor: efficiently evaluate the cost function $\mathcal{C}(\vec{\theta}_i)$ for given $\vec{\theta}_i$
 - Classical computer: given $\mathcal{C}(\vec{\theta}_i)$, find optimized $\vec{\theta}_{i+1}$

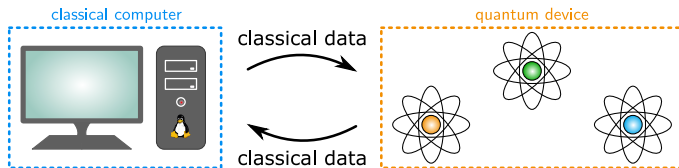
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⇒ **Feedback loop** between the classical computer and the quantum device

Quantum Approximate Optimization Algorithm (QAOA)

Hybrid quantum-classical algorithms

Quantum Approximate Optimization Algorithm (QAOA)

- > Algorithm for approximating (binary) combinatorial optimization problems

$$\begin{aligned} & \min_{x \in V} C(x) \\ & \text{subject to } x \in S \end{aligned}$$

- > x : binary string in $V = \{0, 1\}^n$ encoding a solution
- > $S \subseteq V$: feasible solutions
- > $C : V \rightarrow \mathbb{R}$ cost function
- > Objective is to find the optimal solution

Hybrid quantum-classical algorithms

The Max-Cut problem

Max-Cut

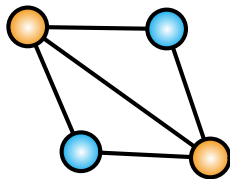
- > **Input:** undirected graph $G = (V, E)$
- > **Task:** find a bipartition of $V = A \cup B$ such that the number of edges between A and B is maximal

Hybrid quantum-classical algorithms

The Max-Cut problem

Max-Cut

- > **Input:** undirected graph $G = (V, E)$
- > **Task:** find a bipartition of $V = A \cup B$ such that the number of edges between A and B is maximal
- > Max-Cut is NP-hard
- ⇒ We cannot find a (quantum) algorithm which solves it polynomial time
- > We can however try to find a good approximation to the exact solution in polynomial time

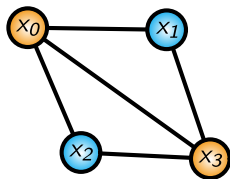


Hybrid quantum-classical algorithms

Max-Cut as combinatorial optimization problem

- > Max-Cut on a Graph $G = (V, E)$ can be expressed as combinatorial optimization problem
- > Label the vertices as x_i define a function w_{ij}

$$x_i = \begin{cases} 0 & \text{if } i \in A \\ 1 & \text{if } i \in B \end{cases} \quad w_{ij} = \begin{cases} 1 & \text{iff } (i, j) \in E \\ 0 & \text{otherwise} \end{cases}$$



Hybrid quantum-classical algorithms

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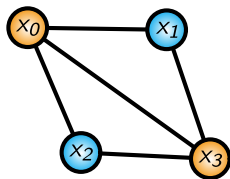
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- > Cost function

$$C(x) = \sum_{i,j=0}^{n-1} w_{ij} x_i (x_j - 1) = \sum_{(i,j) \in E} (x_i (x_j - 1) + x_j (x_i - 1))$$

- ⇒ Contribution of -1 iff endpoints of edge (i, j) belong to different subsets
- > Finding the Max-Cut for G is equivalent to minimizing $C(x)$



Hybrid quantum-classical algorithms

Max-Cut as Hamiltonian problem

- > Cost function can be turned into a Hamiltonian using the mapping $x_i \rightarrow \frac{1}{2}(1 - Z_i)$

$$H_c = \frac{1}{2} \sum_{(i,j) \in E} (Z_i Z_j - 1)$$



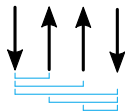
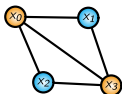
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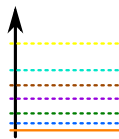
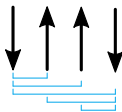
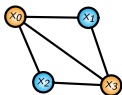
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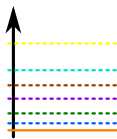
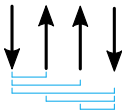
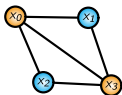
- > Diagonal Hamiltonian of Ising type, summands commute
- > The eigenstates of H are computational basis states encode graph cuts
- > The lower the energy, the larger the number of edges between the subsets
- > The ground state encodes the bit string of the optimal solution x^*

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Max-Cut as Hamiltonian problem

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How to obtain a low energy state of H ?

Hybrid quantum-classical algorithms

The Quantum Approximate Optimization Algorithm (QAOA)

- > We want to find a parametric quantum state $|\psi_p(\vec{\gamma}, \vec{\beta})\rangle$, $\vec{\gamma}, \vec{\beta} \in \mathbb{R}^p$ which minimizes

$$\mathcal{C}(\vec{\gamma}, \vec{\beta}) = \langle \psi_p(\vec{\gamma}, \vec{\beta}) | H_c | \psi_p(\vec{\gamma}, \vec{\beta}) \rangle$$

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- > Transverse field Hamiltonian

$$H_x = \sum_i X_i$$

- > Ansatz for $|\psi_p(\vec{\gamma}, \vec{\beta})\rangle$

$$|\psi_p(\vec{\gamma}, \vec{\beta})\rangle = e^{-i\beta_p H_x} e^{-i\gamma_p H_c} \dots e^{-i\beta_1 H_x} e^{-i\gamma_1 H_c} |+\rangle^{\otimes n}$$

- > $|+\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$ is an eigenstate of X , $X|+\rangle = +1|+\rangle$

E. Farhi, J. Goldstone, S. Gutmann, arXiv:1411.4028

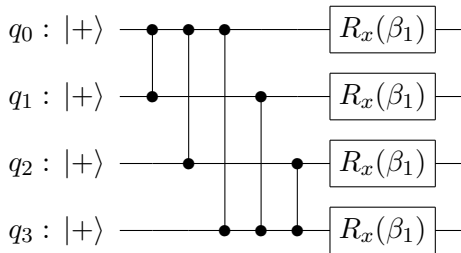
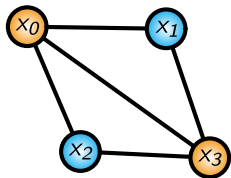
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> Circuit for $p = 1$

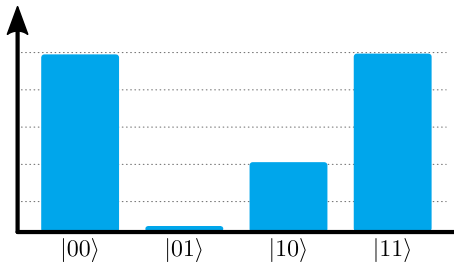


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Hybrid quantum-classical algorithms

The Quantum Approximate Optimization Algorithm (QAOA)

- > $|\psi_p(\vec{\gamma}, \vec{\beta})\rangle$ can be an (entangled) superposition of basis states
- > Measuring $|\psi_p(\vec{\gamma}, \vec{\beta})\rangle$ gives us distribution of bit strings x



Hybrid quantum-classical algorithms

The Quantum Approximate Optimization Algorithm (QAOA)

- > After minimizing $\mathcal{C}(\vec{\gamma}, \vec{\beta})$ the wave function $|\psi_p(\vec{\gamma}, \vec{\beta})\rangle$ has dominant component(s) of low energy states of H_c
- ⇒ Measuring $|\psi_p(\vec{\gamma}, \vec{\beta})\rangle$ gives us a bit string(s) x with high approximation ratio

$$\alpha = \frac{C(x)}{C(x^*)}$$

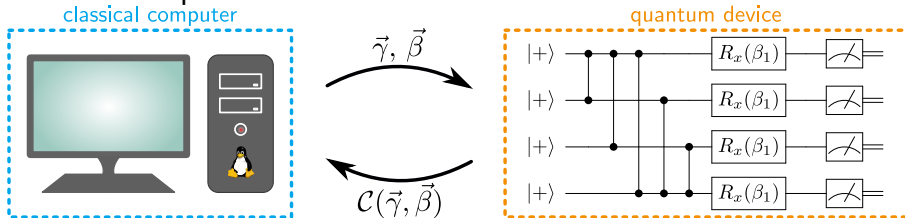
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- > Optimization of the parameters

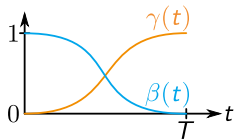


Hybrid quantum-classical algorithms

The Quantum Approximate Optimization Algorithm (QAOA)

- > Ansatz is inspired by trotterized adiabatic time evolution
- > Choose functions $\gamma(t)$, $\beta(t)$ such that

$$\gamma(t) \rightarrow \begin{cases} 0 & \text{for } t \rightarrow 0 \\ 1 & \text{for } t \rightarrow T \end{cases} \quad \beta(t) \rightarrow \begin{cases} 1 & \text{for } t \rightarrow 0 \\ 0 & \text{for } t \rightarrow T \end{cases}$$



and set

$$\gamma_p = \gamma(p\Delta t)\Delta t, \quad \beta_p = \beta(p\Delta t)\Delta t$$

- > With these conventions each time steps corresponds to

$$e^{-i\beta_p H_x} e^{-i\gamma_p H_c} \approx e^{-i(\beta(p\Delta t)H_x + \gamma(p\Delta t)H_c)\Delta t}$$

- > In the limit $p \rightarrow \infty$ the QAOA ansatz can be seen as adiabatic transformation of the initial eigenstate of H_x to the ground state of H_c

Hybrid quantum-classical algorithms

The Quantum Approximate Optimization Algorithm (QAOA)

- > Even $p = 1$ can in general not be simulated on a classical computer efficiently
- > For some problems classical algorithms got a better approximation ratio α
- > Performance depends on the ratio of variables to clauses
- ⇒ Theoretically it is not entirely clear how QAOA performs

B. Barak et al., arXiv:2106.05900
E. Farhi, A. Harrow, arXiv:1602.07674

V. Akshay, H. Philathong, M. E. S. Morales, J. D. Biamonte, Phys. Rev. Lett 124, 090504 (2020)

Variational Quantum Eigensolver (VQE)

Hybrid quantum-classical algorithms

Variational Quantum Algorithms

- > Similar principle can be used to find ground states of quantum Hamiltonians
- > Define a cost function

$$\mathcal{C}(\vec{\theta}) = \langle \psi(\vec{\theta}) | H | \psi(\vec{\theta}) \rangle$$

- > $|\psi(\vec{\theta})\rangle$ ansatz realized by a parametric quantum circuit
- > Provided $|\psi(\vec{\theta})\rangle$ is expressive enough the minimum of $\mathcal{C}(\vec{\theta})$ is obtained for the ground state of H

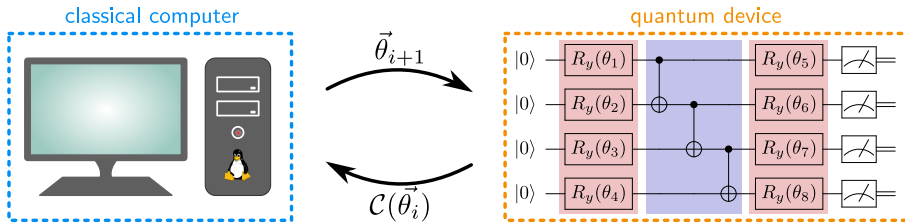
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Peruzzo et al., Nat. Commun. 5, 1 (2014)
J. R. McClean et al., New J. Phys. 18, 023023 (2016)

Variational Quantum Algorithms

Variational Quantum Eigensolver (VQE)

- > For Hamiltonians H that are the sum of a polynomial number of terms the cost function can be measured efficiently on a quantum device, for example

$$H = \sum_{i=1}^N h_{i,i+1} \quad \Leftrightarrow \quad \mathcal{C}(\theta) = \sum_i \langle \psi(\vec{\theta}) | h_{i,i+1} | \psi(\vec{\theta}) \rangle$$

- > In general $\langle \psi(\vec{\theta}) | h_{i,i+1} | \psi(\vec{\theta}) \rangle$ cannot be efficiently evaluated on a classical computer

A. Kandala et al., Nature 549, 242 (2017)
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- > In general $\langle \psi(\vec{\theta}) | h_{i,i+1} | \psi(\vec{\theta}) \rangle$ cannot be efficiently evaluated on a classical computer
- > Example: the Ising model

$$H_{\text{Ising}} = \sum_{i=1}^{N-1} Z_i Z_{i+1} + h \sum_{i=1}^N X_i$$

A. Kandala et al., Nature 549, 242 (2017)
J. R. McClean et al., New J. Phys. 18, 023023 (2016)
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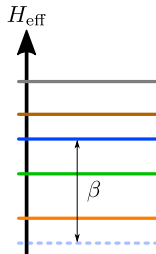
Variational Quantum Algorithms

Variational Quantum Deflation (VQD)

- > VQE can also be extended to higher excited states

$$H_{\text{eff}} = H + \beta \sum_{k=0}^n |\psi_k\rangle\langle\psi_k|$$

- > β has to be chosen large enough



O. Higgot, D. Wang, S. Brierley, Quantum 3, 156 (2019)

Variational Quantum Algorithms

Variational Quantum Deflation (VQD)

- > In practice we never explicitly construct the projector
- > Rather we have to compute overlaps

$$\mathcal{C}(\vec{\theta}) = \langle \psi(\vec{\theta}) | H_{\text{eff}} | \psi(\vec{\theta}) \rangle = \langle \psi(\vec{\theta}) | H | \psi(\vec{\theta}) \rangle + \beta \sum_k |\langle \psi_k | \psi(\vec{\theta}) \rangle|^2$$

- > There are several ways to compute the overlap
 - “Reversing the circuit”: double the depth but keep same number of qubits
 - SWAP test: same circuit depth but double the number of qubits

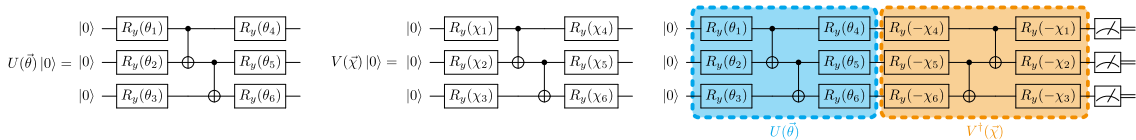
Variational Quantum Algorithms

Reversing the circuit for computing absolute values of overlaps

- > Since both $|\psi_k\rangle = V(\vec{\chi})|0\rangle$ and $|\psi\rangle = U(\vec{\theta})|0\rangle$ are variational circuits the absolute value of the overlap is given by

$$|\langle\psi_k|\psi\rangle| = |\langle 0|V^\dagger(\vec{\chi})U(\vec{\theta})|0\rangle|$$

- > The absolute value can thus be estimated by preparing the state $V^\dagger(\vec{\chi})U(\vec{\theta})|0\rangle$ and recording the number of times one measures $|0\rangle$ at the end
- > For commonly used quantum gates the inverse is often trivial

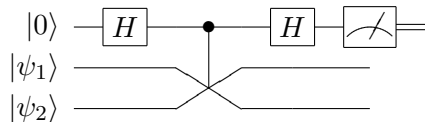


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Variational Quantum Algorithms

SWAP test for computing absolute values of overlaps

- > The SWAP test allows for computing the overlap of arbitrary states in registers

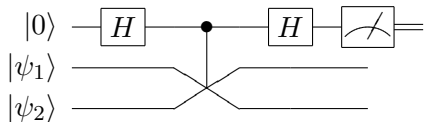


A. Barenco *et al.*, SIAM Journal on Computing 26, 1541 (1997)

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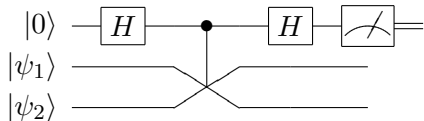


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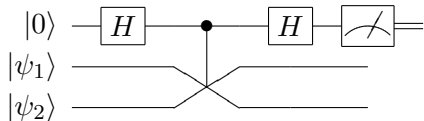


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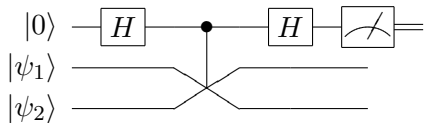


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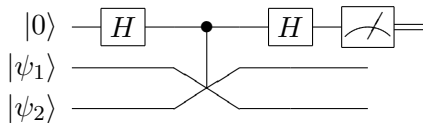
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- 4 Hadamard creates $\frac{1}{2}|0\rangle(|\psi, \phi\rangle + |\psi, \phi\rangle) + \frac{1}{2}|1\rangle(|\psi, \phi\rangle - |\phi, \psi\rangle)$



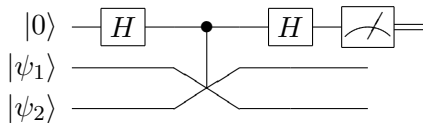
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- 4 Hadamard creates $\frac{1}{2}|0\rangle(|\psi, \phi\rangle + |\psi, \phi\rangle) + \frac{1}{2}|1\rangle(|\psi, \phi\rangle - |\phi, \psi\rangle)$
- 5 Probability of measuring the ancilla in state $|0\rangle$:
$$p_0 = \frac{1}{2} + \frac{1}{2}|\langle\psi, \phi\rangle|^2$$

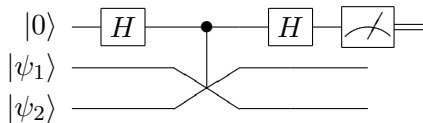


A. Barenco *et al.*, SIAM Journal on Computing 26, 1541 (1997)

Variational Quantum Algorithms

SWAP test for computing absolute values of overlaps

- > The SWAP test allows for computing the overlap of arbitrary states in registers
- > Based on the controlled SWAP operation/Fredkin gate
 - 1 Start from $|0, \psi, \phi\rangle$
 - 2 Hadamard creates $\frac{1}{2}(|0, \psi, \phi\rangle + |1, \psi, \phi\rangle)$
 - 3 Fredkin gate yields $\frac{1}{2}(|0, \psi, \phi\rangle + |1, \phi, \psi\rangle)$
 - 4 Hadamard creates $\frac{1}{2}|0\rangle(|\psi, \phi\rangle + |\psi, \phi\rangle) + \frac{1}{2}|1\rangle(|\psi, \phi\rangle - |\phi, \psi\rangle)$
 - 5 Probability of measuring the ancilla in state $|0\rangle$:
$$p_0 = \frac{1}{2} + \frac{1}{2}|\langle\psi, \phi\rangle|^2$$
- > More efficient ways for VQD: destructive SWAP test



A. Barenco *et al.*, SIAM Journal on Computing 26, 1541 (1997)

L. Cincio *et al.*, New J. Phys. 20, 113022 (2018)

J. C. Garcia-Escartin, P. Chamorro-Posada Phys. Rev. A 87, 052330 (2013)

Variational Quantum Algorithms

QAOA

- > Combinatorial optimization problems
- > Problem Hamiltonian is diagonal in the computational basis
- > Circuit structure is fixed
- > In the limit of infinite layers provably converges to the exact solution

Variational Quantum Algorithms

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VQE

- > Ground states/low-lying excitations
- > Efficient as long as H has only a polynomial number of terms
- > Hamiltonian exists only as a measurement
- > Great freedom choosing the circuit
 - Problem requirements
 - Available hardware
 - Expressiveness

Variational Quantum Algorithms

QAOA

- > Combinatorial optimization problems
- > Problem Hamiltonian is diagonal in the computational basis
- > Circuit structure is fixed
- > In the limit of infinite layers provably converges to the exact solution

- > Best answer for the given set of resources
- > Largely resilient to systematic errors of the device

VQE

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Example: VQE for the Schwinger model

Variational Quantum Algorithms

VQE for the Schwinger model

- > Hamiltonian in spin language after integrating out the gauge fields

$$H = w \sum_{j=1}^{N-1} \left(\sigma_j^+ \sigma_{j+1}^- + \sigma_j^- \sigma_{j+1}^+ \right) + \frac{m}{2} \sum_{j=1}^N (-1)^j Z_j + g \sum_{j=1}^N \left(\varepsilon_0 - \frac{1}{2} \sum_{l=1}^j (Z_l + (-1)^l) \right)^2$$

- > Symmetries of the problem
 - Conservation of total charge $Q_{\text{tot}} = \sum_n Q_n = \sum_n Z_n$
 - CP symmetry within a charge sector: reflecting around the center and flipping the spins
- > Symmetries can be incorporated in the VQE ansatz

Variational Quantum Algorithms

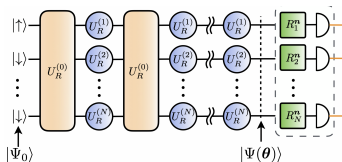
VQE for the Schwinger model

- > Resource Hamiltonians used for gate operations $U = \exp(-i\theta H_R^k)$

$$H_R^0 = \sum_{i=1}^{N-1} \sum_{j=i+1}^N J_{ij} (\sigma_i^+ \sigma_j^- + \sigma_i^- \sigma_j^+) + B \sum_{i=1}^N Z_i,$$

$$H_R^j = \frac{\Delta_0}{2} Z_j$$

- > Resource Hamiltonians respect charge conservation
 - > Restrict the parameters of the single-qubit gates as $\theta^j = -\theta^{N+1-j}$ to enforce CP symmetry
 - > For large systems the bulk should approximately be translation invariant
- ⇒ Restrict parameters in bulk to the same value
- > Initial wave function compatible with the symmetries: Neel state $|\psi_0\rangle = |\uparrow\downarrow \cdots \uparrow\downarrow\rangle$

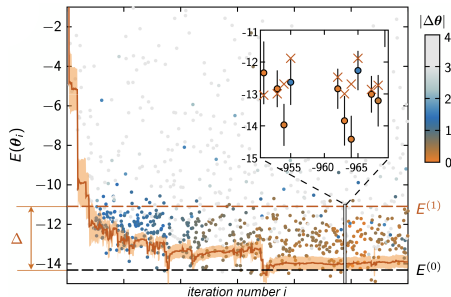


C. Kokail et al., Nature 569, 355 (2019)

Variational Quantum Algorithms

VQE for the Schwinger model

- > Results on a trapped ion quantum computer with 20 qubits formed by $^{40}\text{Ca}^+$ ions
- > Simulation involves 6 layers and 15 parameters



⇒ Ground state can be determined with good accuracy

Variational Quantum Algorithms

VQE for the Schwinger model

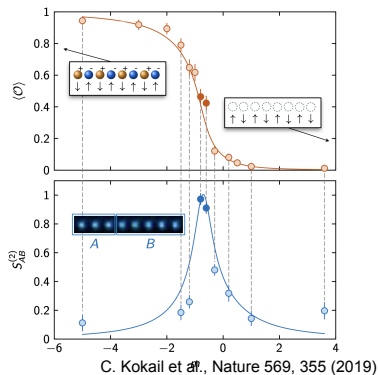
- > As we have seen in the previous lecture, a negative mass in the Schwinger model corresponds to the case of $\theta = \pi$
- > For this value of θ , the model undergoes a second order phase transition for
- > The phase structure with negative masses can also be explored with VQE

Variational Quantum Algorithms

VQE for the Schwinger model

- > As we have seen in the previous lecture, a negative mass in the Schwinger model corresponds to the case of $\theta = \pi$
- > For this value of θ , the model undergoes a second order phase transition for
- > The phase structure with negative masses can also be explored with VQE
- > The detect the phase transition monitor

$$O = \frac{1}{2N(N-1)} \sum_{i,j>i} \left((1 + (-1)^i Z_i)(1 + (-1)^j Z_j) \right)$$



Thank you for your attention!

Questions?