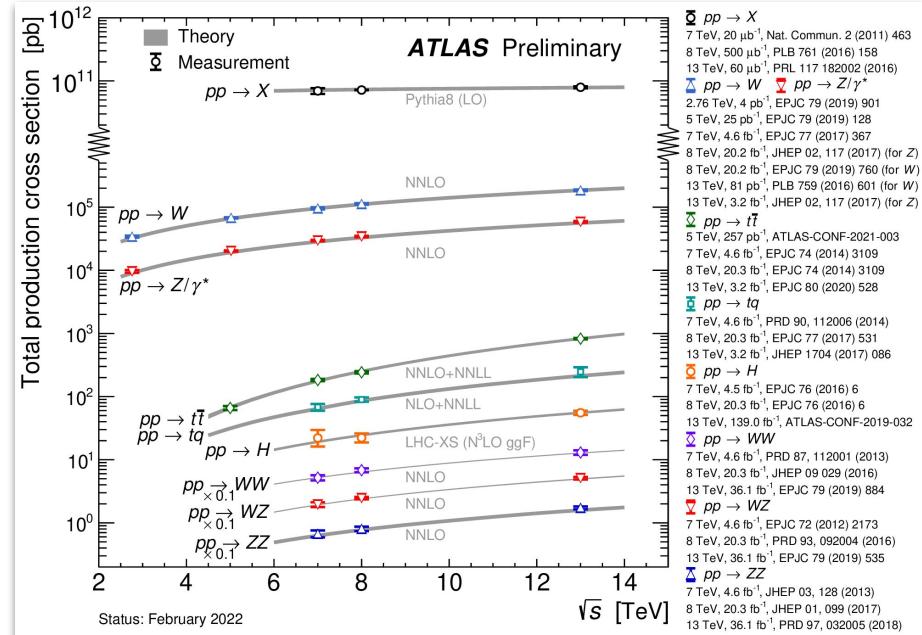
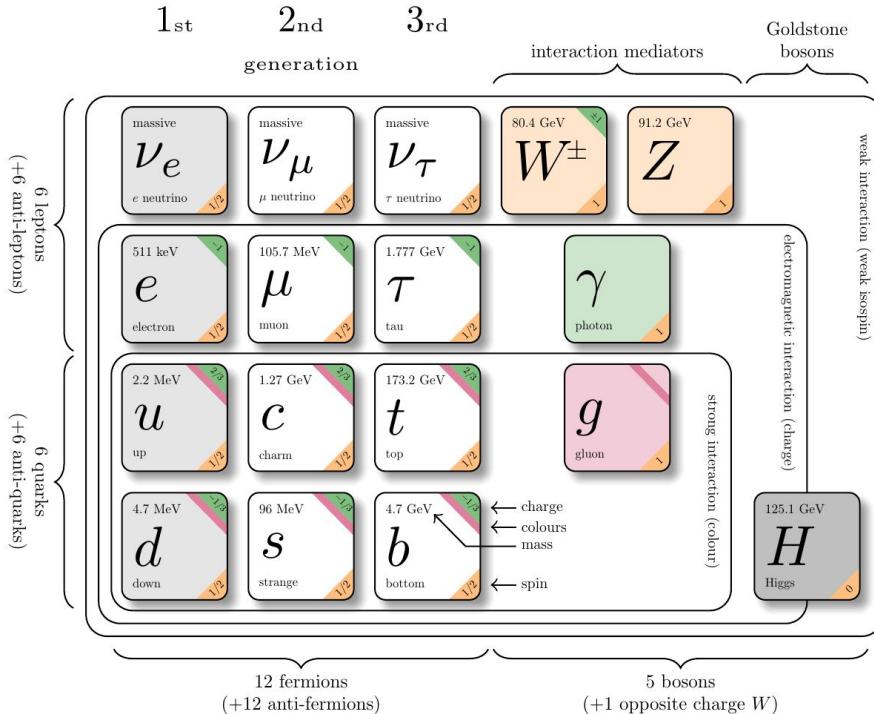


# Physics beyond the standard model II

Ben Brüers  
DESY Zeuthen, 30.08.2023



# Last time...



<http://cds.cern.ch/record/2804061>

**The Standard model is very precise, but...**

# Last time (2)...

- ...there are a bunch of things it cannot explain:

9 fermion masses ( $m_u, m_d, m_c, m_s, m_b, m_t; m_e, m_\mu, m_\tau$ )

- + 2 Higgs boson parameters: the mass & VEV ( $m_H, v$ )
- + 3 coupling parameters ( $g_W, g', g_s$ )
- + 4 CKM parameters (3 mixing angles + 1 CP violating phase)
- + 1 CP violating phase in QCD (see later)

19 free parameters

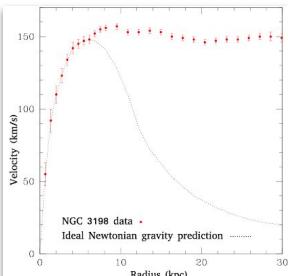
$$\delta M_H^2 = \frac{G_F \Lambda^2}{4\pi^2 \sqrt{2}} (6M_W^2 + 3M_Z^2 + M_H^2 - 12m_t^2)$$

$$\delta M_H^2|_{t\text{-loop}} \approx -\frac{3G_F}{\pi^2 \sqrt{2}} m_t^2 \Lambda^2 \approx -0.075 \Lambda^2$$

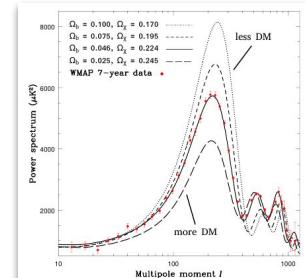
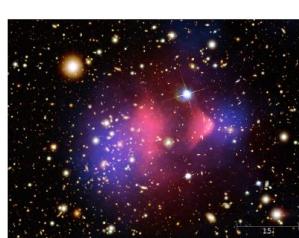
“bare mass” tuned very finely,  $O(10^{-4})$ - $O(10^{-34})$  GeV!

$$\mathcal{L}_{\text{QCD}} = \bar{\psi}_i (\underbrace{i\gamma^\mu (D_\mu)_{ij}}_{\text{quark dynamics}} - \underbrace{m \delta_{ij}}_{\text{quark mass}}) \psi_j - \underbrace{\frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}}_{\text{gluon dynamics}} - \underbrace{\Theta \frac{\alpha_s}{8\pi} G^{\mu\nu a} \tilde{G}_{\mu\nu}^a}_{\text{CP-violating term}}$$

? ? ?



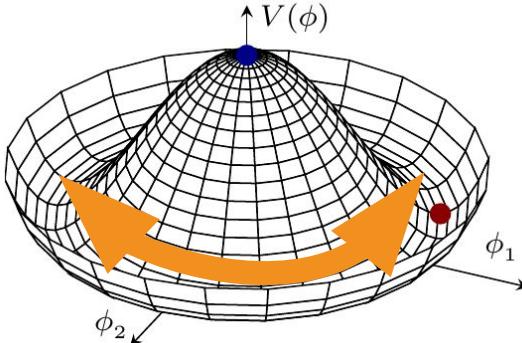
K. Begeman, Astron. Astrophys. 223 (1989), pp. 47–60  
[doi:10.1086/381970](https://doi.org/10.1086/381970)



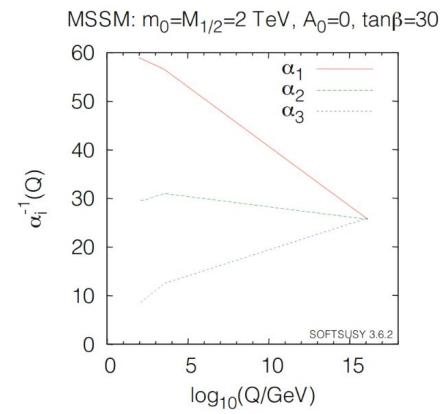
[arxiv:1001.4635](https://arxiv.org/abs/1001.4635) [arxiv:1006.2483](https://arxiv.org/abs/1006.2483)

# Last time (3)...

- Discussed multiple SM extension:
  - Axions + ALPs
  - 2HDM(+a)
  - SUSY
  - GUTs
  - Extra dimensions



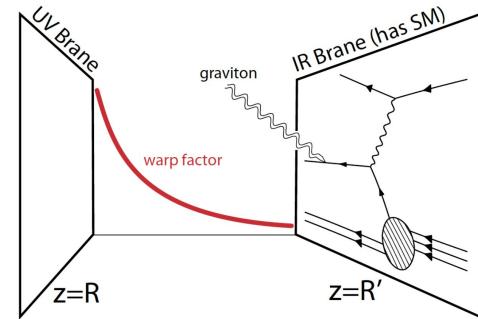
Drawing inspired by  
[https://indico.scc.kit.edu/event/477/contributions/4854/attachments/2575/3683/KSEITA\\_Durbach\\_2019\\_pargner.pdf](https://indico.scc.kit.edu/event/477/contributions/4854/attachments/2575/3683/KSEITA_Durbach_2019_pargner.pdf)



[https://www.zauthen.desy.de/students/2019/lectures/Pueschel\\_Lecture\\_BeyondSM\\_2019.pdf](https://www.zauthen.desy.de/students/2019/lectures/Pueschel_Lecture_BeyondSM_2019.pdf)

$$Q |\text{fermion}\rangle = |\text{boson}\rangle$$

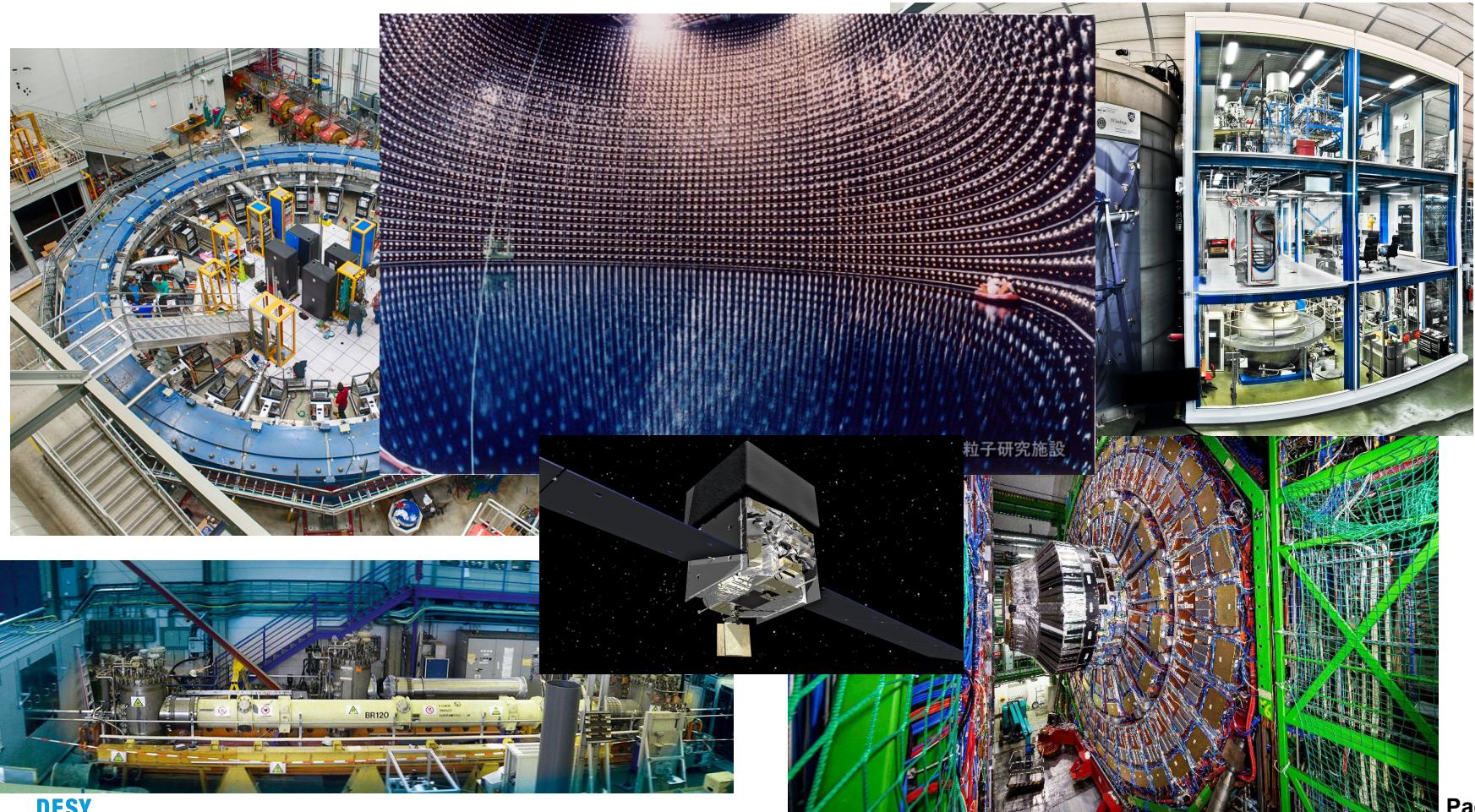
$$Q |\text{boson}\rangle = |\text{fermion}\rangle$$



Model	Dark Matter	Hierarchy problem	Strong CP problem	Unification	Gravity
Axions	✓	-	✓	-	-
2HDM	✓	-	-	-	-
SUSY	✓	✓	possible	✓	e.g. mSUGRA
GUTs	-	-	-	✓	-
Extra dims.	✓	✓	-	possible	✓

# Today:

- Will discuss experiments searching for physics beyond the Standard Model
- Will cover collider-based and other experiments



# The LHC

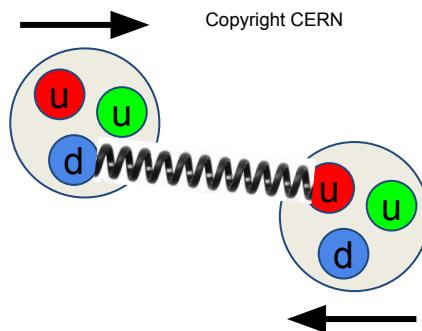
## LHC

→ circumference  
27 km

→ proton-proton  
collisions with a  
CME of  
 $\sqrt{s} = 13 \text{ TeV}$

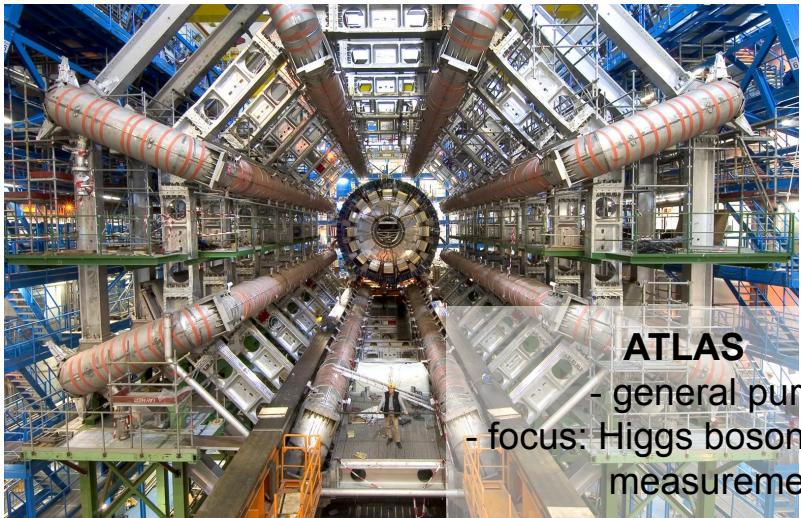
→ interacting  
particles are  
quarks, so their  
CME often  
smaller

→ detectors  
surround points  
of collision

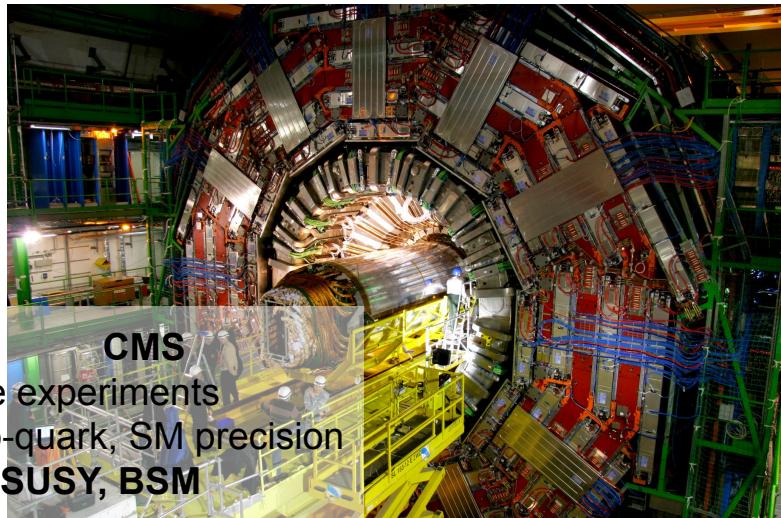


<https://previews.123rf.com/images/denisnata/denisnata1003/denisnata100300048/6555910-black-spiral-telephone-cable-isolated-on-white-background.jpg>

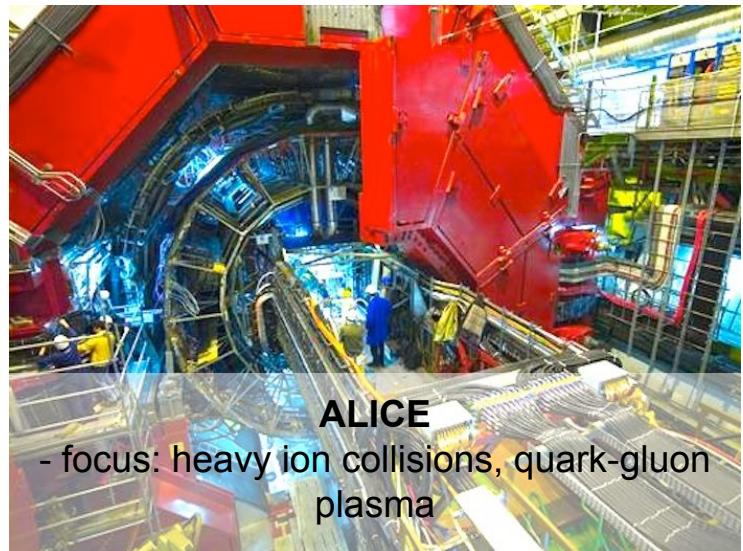
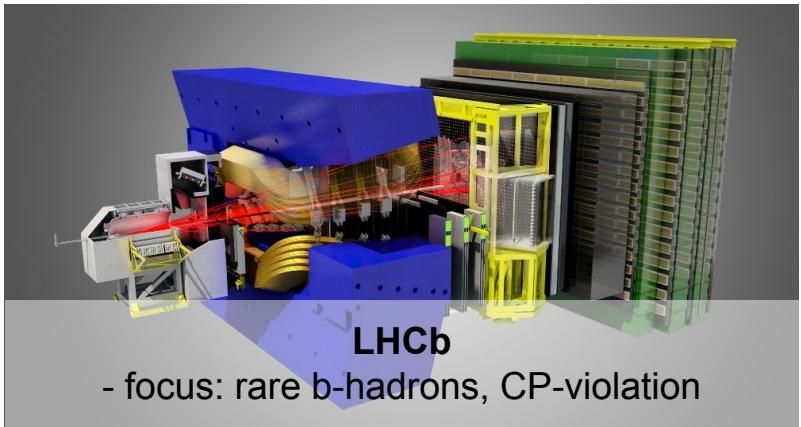
# Experiments at the LHC



[https://www.weltmaschine.de/sites/sites\\_custom/site\\_weltmaschine/content/e28861/e36564/e36588/e36608/0511013\\_01-A4-at-144-dpi.jpg](https://www.weltmaschine.de/sites/sites_custom/site_weltmaschine/content/e28861/e36564/e36588/e36608/0511013_01-A4-at-144-dpi.jpg)

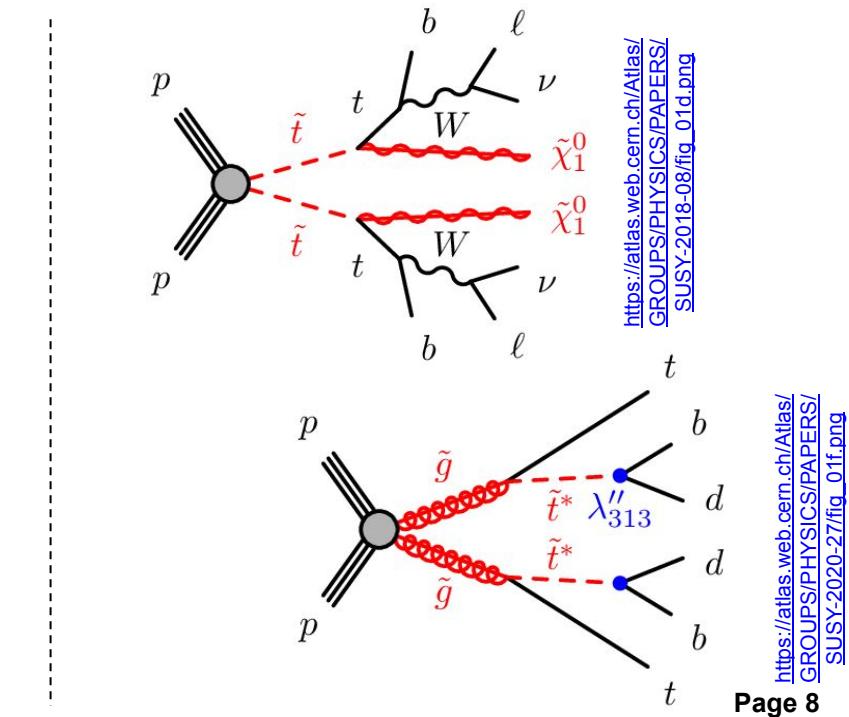
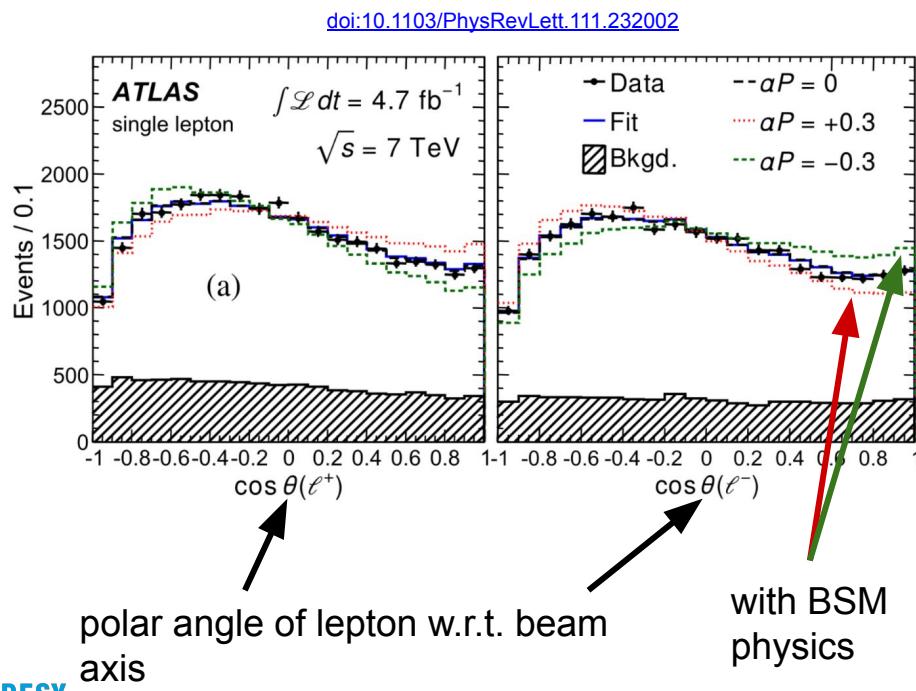


[https://cms.cern/sites/default/files/field/image/cds-record-1275108-hoch-20071215\\_721-nice.jpg](https://cms.cern/sites/default/files/field/image/cds-record-1275108-hoch-20071215_721-nice.jpg)



# How to search for BSM physics at colliders

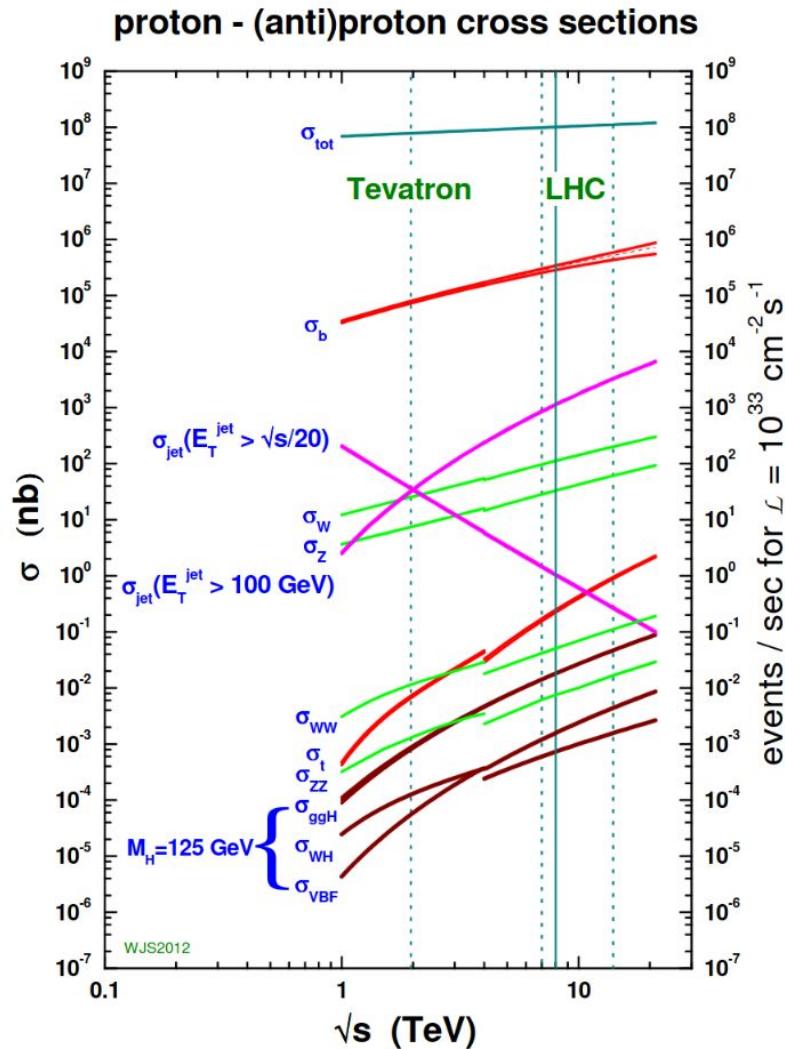
- BSM physics must be rare and/or involve heavy particles (else discovered)
- Two search principles:
  - BSM physics slightly modifies masses, couplings, etc.  
→ investigate these effects in precision measurements
  - **BSM physics particles produced in proton collisions**  
→ **search for the particles / reconstruct their decays, etc.**  
→ **ideally investigate signatures that are rare in the SM**



# Metrics to keep in mind

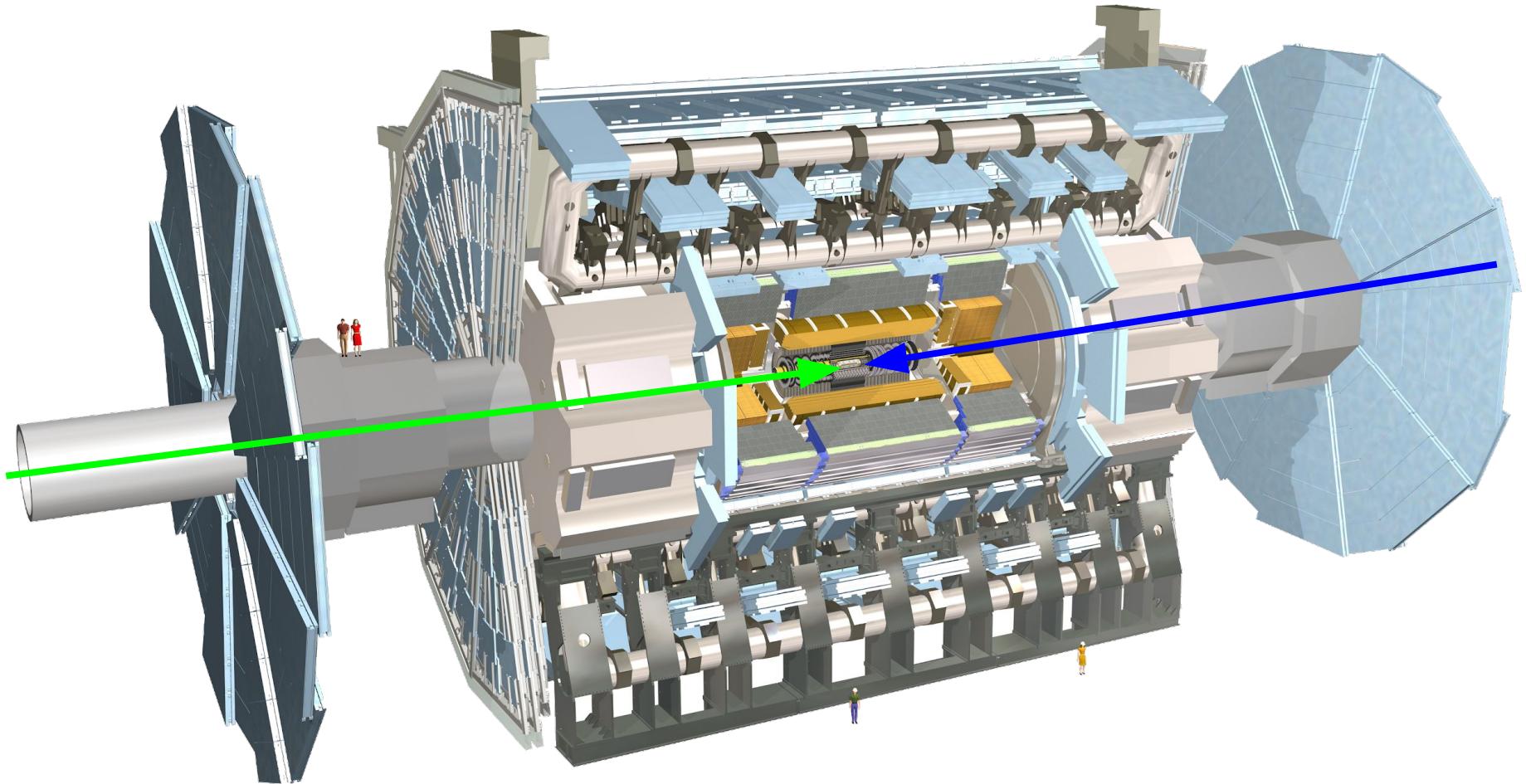
- BSM physics must be rare and/or involve heavy particles (else discovered)
- Cross-section must not be too low, else will not be able to discover anything  
→ **important model metric**
- Too investigate very rare signatures need many collision events
- The mass of the particles must not be heavier than CME, else not produced
- Can only detect particles produced in area of detector (detector not infinite) and if detector sufficiently efficient  
→ **important collider metrics: CME, luminosity, detector efficiency**

$$N_{\text{obs}} = \int \mathcal{L}(t) \sigma \varepsilon A dt$$



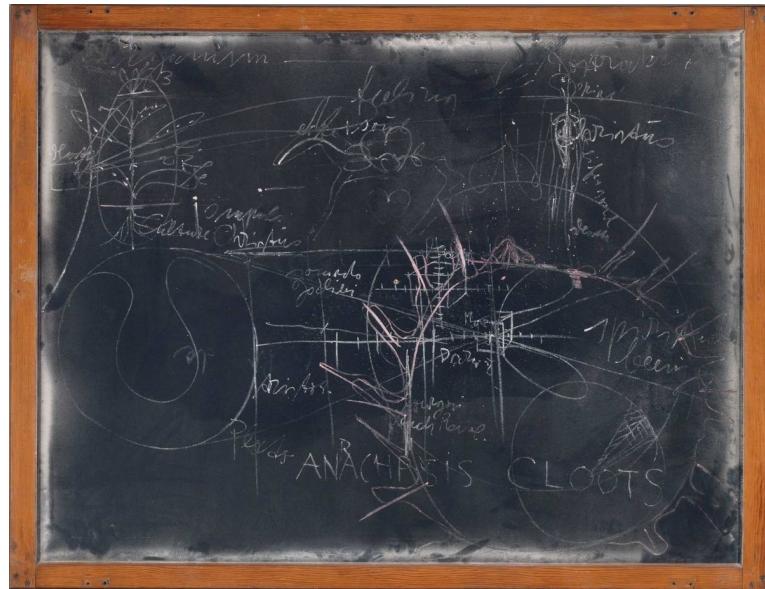
[http://www.hep.ph.ic.ac.uk/~wstirlin/plots/crosssections2012\\_v5.pdf](http://www.hep.ph.ic.ac.uk/~wstirlin/plots/crosssections2012_v5.pdf)

# The ATLAS detector



# How do we measure particles with the ATLAS detector?

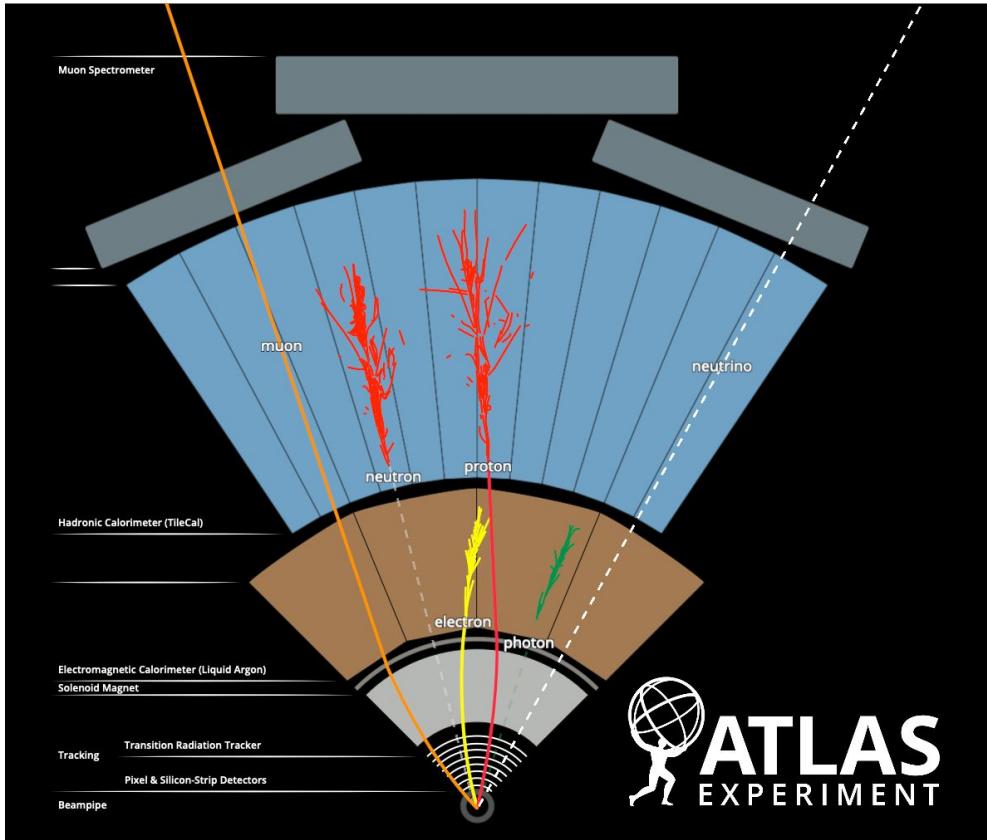
It's blackboard time!



<https://cameo.mfa.org/images/b/ba/2000.979-CR9834-d1.jpg>

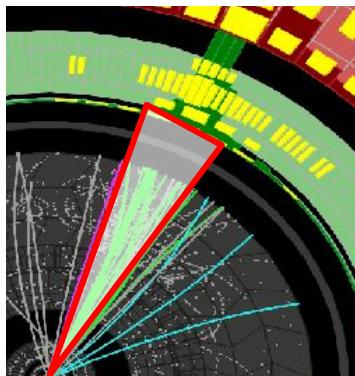
(summary on the next slide)

# Particle reconstruction: ATLAS detector



$q$  →  
confinement

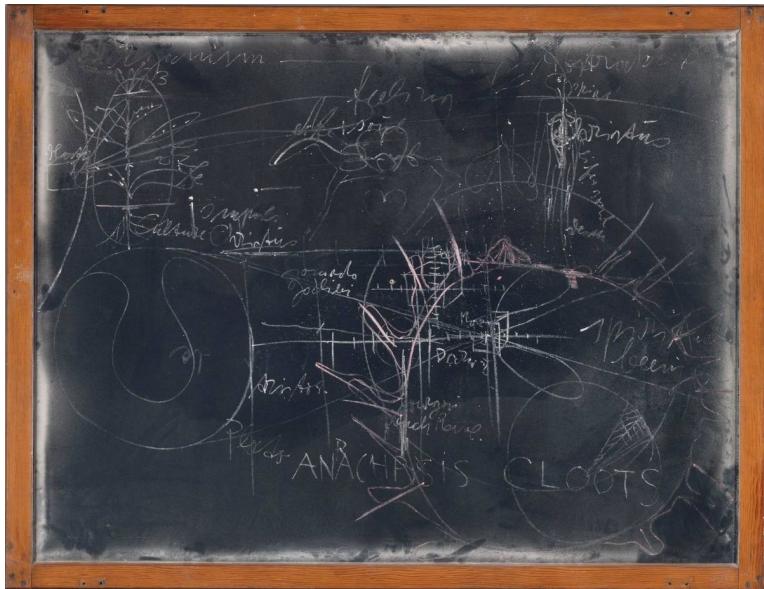
had. 1  
had. 2  
had. 3  
...  
had. n



- Multi-layer detector, measure particles via their interaction with the detector
- Inner-most detector: measure tracks of charged particles (electron, muon, charged hadron)
- ECAL / HCAL: measure energy of EM-interacting / hadronically interacting particle  
→ concept: make particles lose all their energy and measure the loss
- Muon chambers for high precision muon momentum measurements
- NB: due to quark confinement, quarks cannot exist alone  
→ form collimated hadron sprays, we call these “jets”

# What about invisible particles?

It's blackboard time!



<https://cameo.mfa.org/images/b/ba/2000.979-CR9834-d1.jpg>

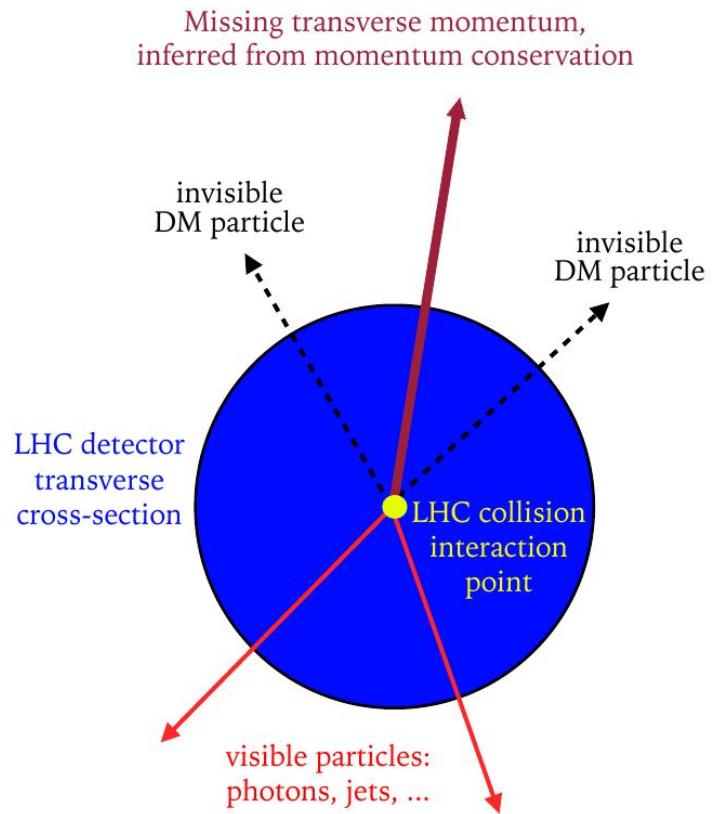
(summary on the next slide)

# How do we measure invisible particles?

- Neutrinos and potential BSM particles interact weakly with detector → no signal
- But: if produced with particles that produce a signal: use **momentum conservation to infer** on them
- Protons collide heads on: momentum of interacting particles in direction orthogonal to the beam axis is ~0
- Momentum is conserved → momenta of all particles in the plane transverse to the beam axis must sum to zero  
→ can infer on the total momentum of the particles escaping detection and the direction in the transverse plane  
→ **this is called Missing Transverse Momentum and the magnitude Missing Transverse Energy (MET)**

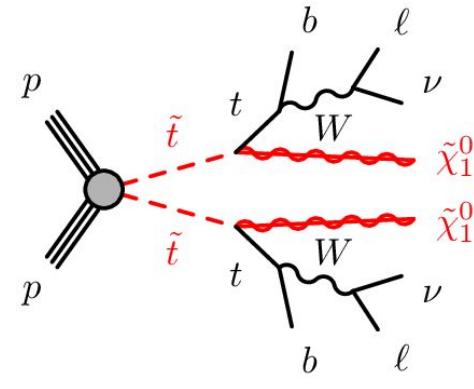
$$0 = \sum_{\text{all}} \vec{p}_{T,i} = \sum_{\text{visible}} \vec{p}_{T,i} + \sum_{\text{in-visible}} \vec{p}_{T,j}$$

**Missing transverse momentum**



# How do we search for the BSM particles?

- In the data, select those collision events with:
  - 2 jets compatible with a bottom-quark
  - 2 leptons
  - MET compatible with  $2x \chi_1^0$  and  $2x \nu$
  - Maybe an invariant mass requirement on the system of  $b-W-\chi_1^0$
- **Caveat: we might not have saved these events in ATLAS!!**
  - Protons collide every 25 ns in ATLAS
    - 1 Mb of data / event or 40 Tb of data per second!!
    - cannot store this!! & most not of interest (low energy)
  - Solution: coarsely analyse all collision event as they happen
  - Only store events fulfilling certain criteria, e.g. 6 jets, 2 leps → “triggering”
  - **So: the BSM physics you can find depends on your trigger criteria!**



[https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/SUSY-2018-08/fig\\_01d.png](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/SUSY-2018-08/fig_01d.png)



# Triggering criteria in ATLAS (in 2018)

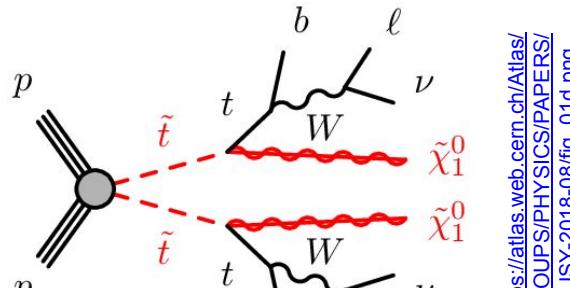
- Lots of triggers there! But sometimes need to add one...

[ATL-DAQ-PUB-2019-001/](#)

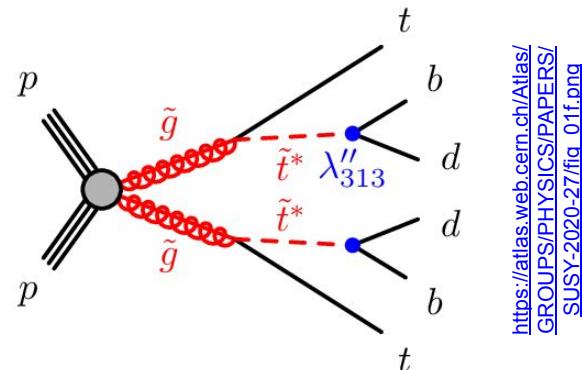
Trigger	Typical offline selection	Trigger Selection		L1 Peak Rate [kHz] $L=2.0 \times 10^{34} \text{ cm}^{-2} \text{s}^{-1}$	HLT Peak Rate [Hz]
		L1 [GeV]	HLT [GeV]		
Single leptons	Single isolated $\mu$ , $p_T > 27 \text{ GeV}$	20	26 (i)	16	218
	Single isolated tight $e$ , $p_T > 27 \text{ GeV}$	22 (i)	26 (i)	31	195
	Single $\mu$ , $p_T > 52 \text{ GeV}$	20	50	16	70
	Single $e$ , $p_T > 61 \text{ GeV}$	22 (i)	60	28	20
	Single $\tau$ , $p_T > 170 \text{ GeV}$	100	160	1.4	42
Two leptons	Two $\mu$ , each $p_T > 15 \text{ GeV}$	$2 \times 10$	$2 \times 14$	2.2	30
	Two $\mu$ , $p_T > 23, 9 \text{ GeV}$	20	22, 8	16	47
	Two very loose $e$ , each $p_T > 18 \text{ GeV}$	$2 \times 15$ (i)	$2 \times 17$	2.0	13
	One $e$ & one $\mu$ , $p_T > 8, 25 \text{ GeV}$	20 ( $\mu$ )	7, 24	16	6
	One loose $e$ & one $\mu$ , $p_T > 18, 15 \text{ GeV}$	15, 10	17, 14	2.6	5
	One $e$ & one $\mu$ , $p_T > 27, 9 \text{ GeV}$	22 (e, i)	26, 8	21	4
	Two $\tau$ , $p_T > 40, 30 \text{ GeV}$	20 (i), 12 (i) (+jets, topo)	35, 25	5.7	93
	One $\tau$ & one isolated $\mu$ , $p_T > 30, 15 \text{ GeV}$	12 (i), 10 (+jets)	25, 14 (i)	2.4	17
	One $\tau$ & one isolated $e$ , $p_T > 30, 18 \text{ GeV}$	12 (i), 15 (i) (+jets)	25, 17 (i)	4.6	19
Three leptons	Three very loose $e$ , $p_T > 25, 13, 13 \text{ GeV}$	20, $2 \times 10$	$24, 2 \times 12$	1.6	0.1
	Three $\mu$ , each $p_T > 7 \text{ GeV}$	$3 \times 6$	$3 \times 6$	0.2	7
	Three $\mu$ , $p_T > 21, 2 \times 5 \text{ GeV}$	20	$20, 2 \times 4$	16	9
	Two $\mu$ & one loose $e$ , $p_T > 2 \times 11, 13 \text{ GeV}$	$2 \times 10$ ( $\mu$ )	$2 \times 10, 12$	2.2	0.5
	Two loose $e$ & one $\mu$ , $p_T > 2 \times 13, 11 \text{ GeV}$	$2 \times 8, 10$	$2 \times 12, 10$	2.3	0.1
Single photon	One loose $\gamma$ , $p_T > 145 \text{ GeV}$	24 (i)	140	24	47
Two photons	Two loose $\gamma$ , each $p_T > 55 \text{ GeV}$	$2 \times 20$	$2 \times 50$	3.0	7
	Two $\gamma$ , $p_T > 40, 30 \text{ GeV}$	$2 \times 20$	35, 25	3.0	21
	Two isolated tight $\gamma$ , each $p_T > 25 \text{ GeV}$	$2 \times 15$ (i)	$2 \times 20$ (i)	2.0	15
Single jet	Jet ( $R = 0.4$ ), $p_T > 435 \text{ GeV}$	100	420	3.7	35
	Jet ( $R = 1.0$ ), $p_T > 480 \text{ GeV}$	111 (topo: $R = 1.0$ )	460	2.6	42
	Jet ( $R = 1.0$ ), $p_T > 450 \text{ GeV}$ , $m_{\text{jet}} > 45 \text{ GeV}$	111 (topo: $R = 1.0$ )	420, $m_{\text{jet}} > 35$	2.6	36
$b$ -jets	One $b$ ( $\epsilon = 60\%$ ), $p_T > 285 \text{ GeV}$	100	275	3.6	15
	Two $b$ ( $\epsilon = 60\%$ ), $p_T > 185, 70 \text{ GeV}$	100	175, 60	3.6	11
	One $b$ ( $\epsilon = 40\%$ ) & three jets, each $p_T > 85 \text{ GeV}$	$4 \times 15$	$4 \times 75$	1.5	14
	Two $b$ ( $\epsilon = 70\%$ ) & one jet, $p_T > 65, 65, 160 \text{ GeV}$	$2 \times 30, 85$	$2 \times 55, 150$	1.3	17
	Two $b$ ( $\epsilon = 60\%$ ) & two jets, each $p_T > 65 \text{ GeV}$	$4 \times 15,  \eta  < 2.5$	$4 \times 55$	3.2	15
Multijets	Four jets, each $p_T > 125 \text{ GeV}$	$3 \times 50$	$4 \times 115$	0.5	16
	Five jets, each $p_T > 95 \text{ GeV}$	$4 \times 15$	$5 \times 85$	4.8	10
	Six jets, each $p_T > 80 \text{ GeV}$	$4 \times 15$	$6 \times 70$	4.8	4

# Model dependence of collider searches

- Disadvantage of triggering: need to know what you are looking for...
- Need to select a model to design your search!  
→ BSM searches at collider depend on the model you are investigating!!!!



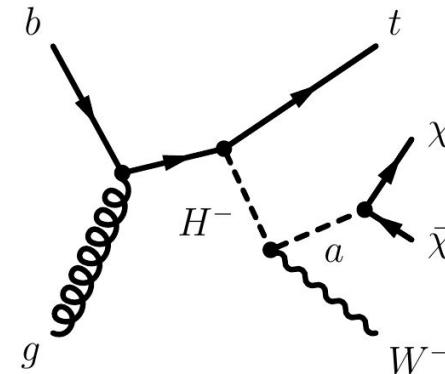
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[https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/SUSY-2020-27/fig\\_01f.png](https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/SUSY-2020-27/fig_01f.png)

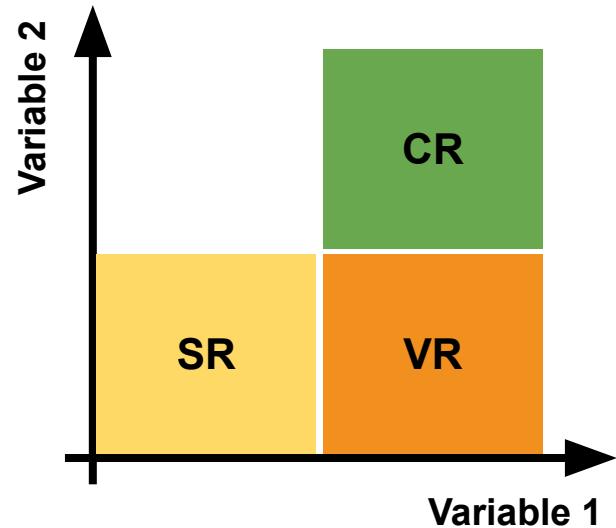
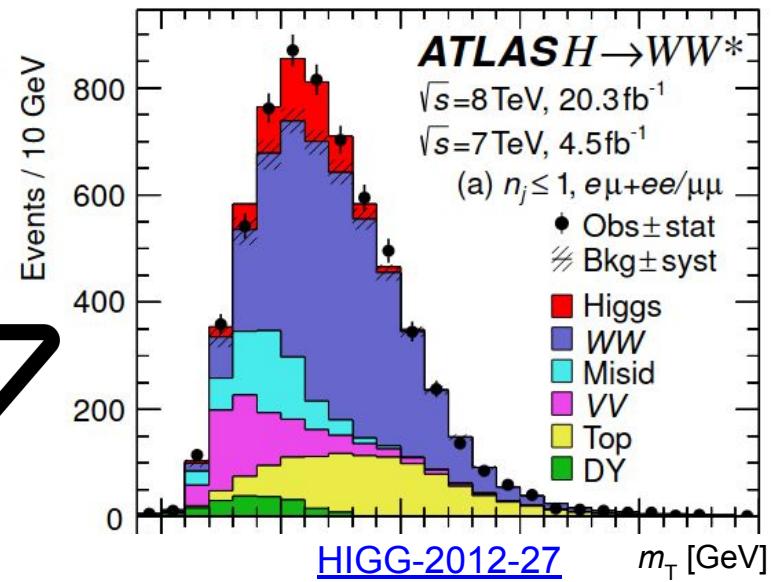
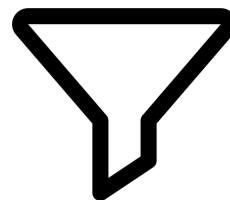


[https://www.ataxia.org/wp-content/uploads/2022/07/needle-in-a-haystack-1752846\\_960\\_720.jpg](https://www.ataxia.org/wp-content/uploads/2022/07/needle-in-a-haystack-1752846_960_720.jpg)



# Ok fine, we triggered, we reconstructed our particles, what's next?

- The entire dataset still contains very many collision events
- BSM physics is rare → need to filter data to “see” our BSM physics
- Use e.g. number of reconstructed electrons, jet momenta, etc.
- Filtered data comprised of our signal & Standard Model background
- Define multiple “filters” = “regions”
  - **Signal regions:** much signal, little bkg.
  - **Control regions:** estimate bkg.
  - **Validation regions:** verify bkg. estimate

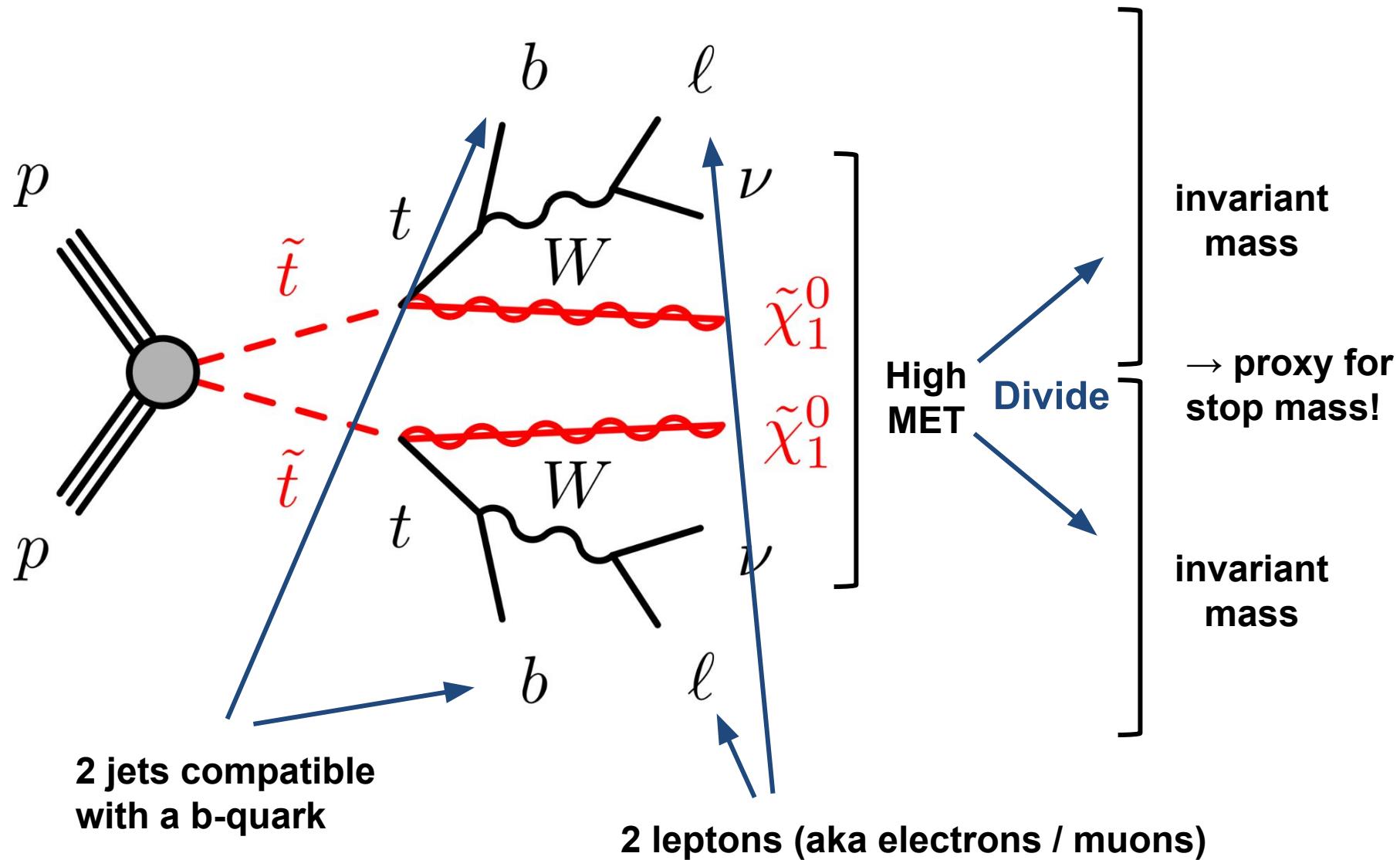


# Searching for SUSY

- Recall: there are many SUSY models, e.g. MSSM, pMSSM, CMSSM, etc.
- In practise: study simplified versions of these models to reduce number of parameters to study
- Often also reduce number of production channels / decay channels
- Result give an **indication** for what would happen in the full model

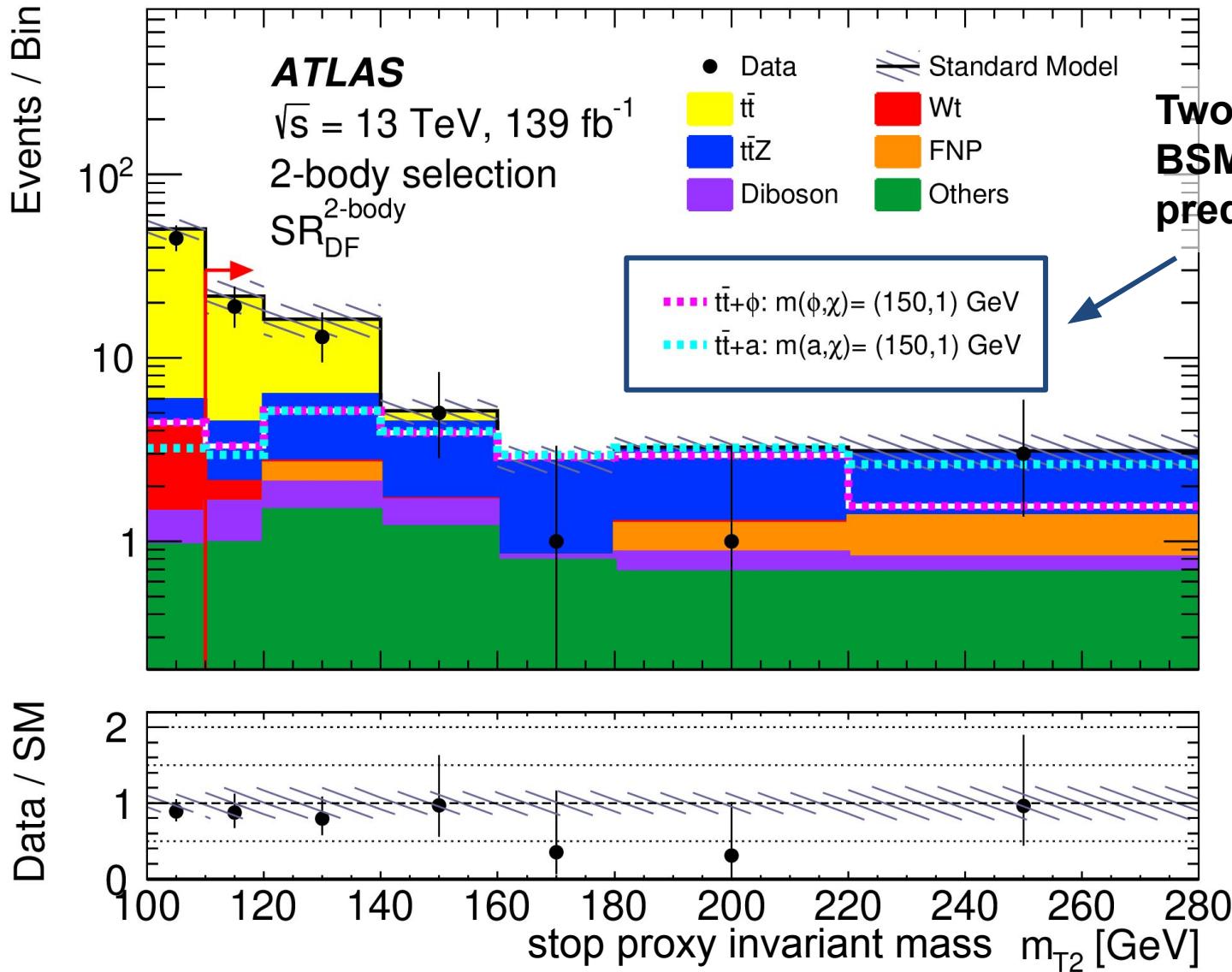
# SUSY example: stop-2L

[SUSY-2018-08](#)



# Stop-2L: final distribution

[SUSY-2018-08](#)

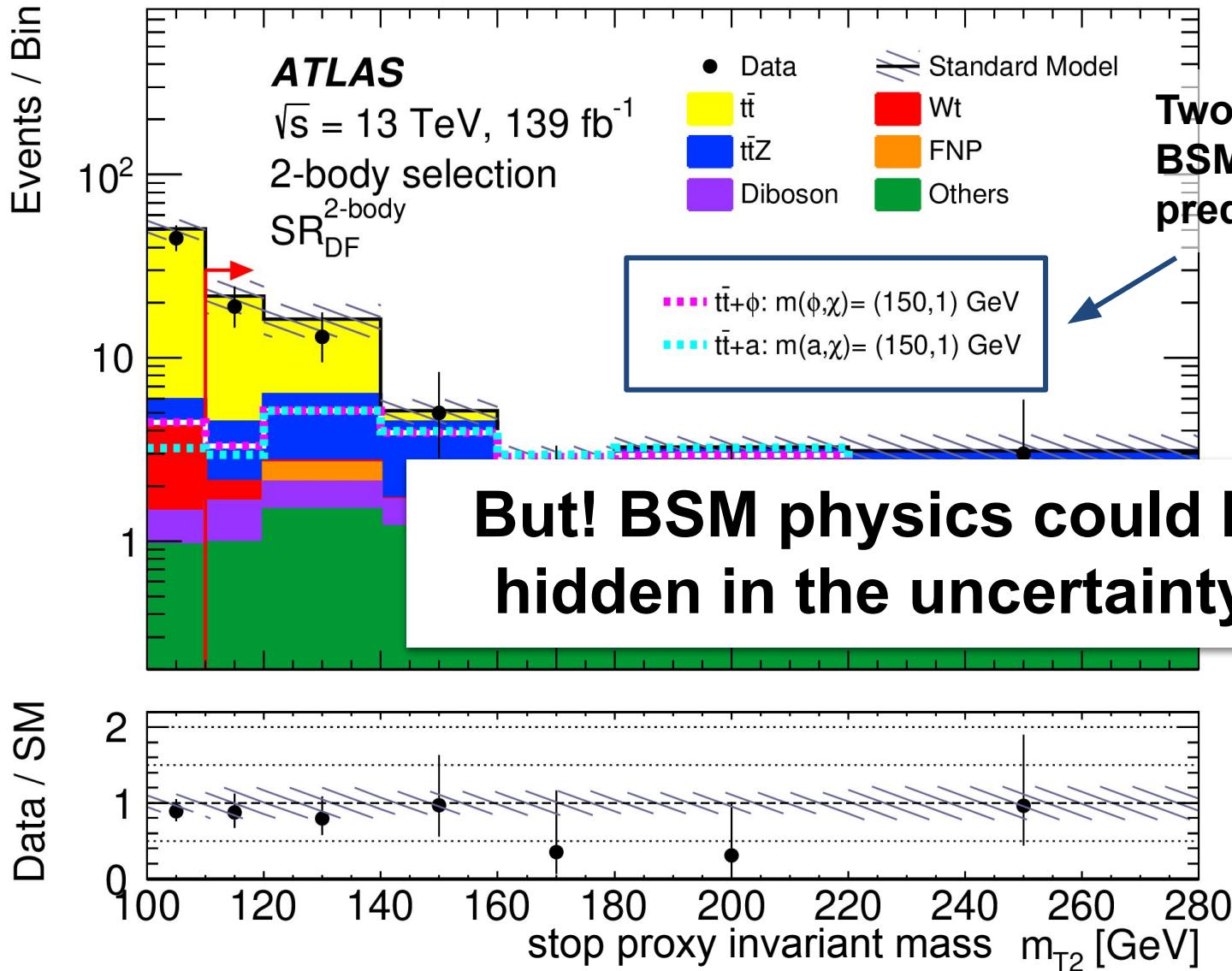


**Two example  
BSM  
predictions!**

Good agreement  
of data & SM  
prediction  
→ didn't find BSM  
physics **within**  
**precision** (i.e.  
cyan / move  
models excluded)

# Stop-2L: final distribution

[SUSY-2018-08](#)

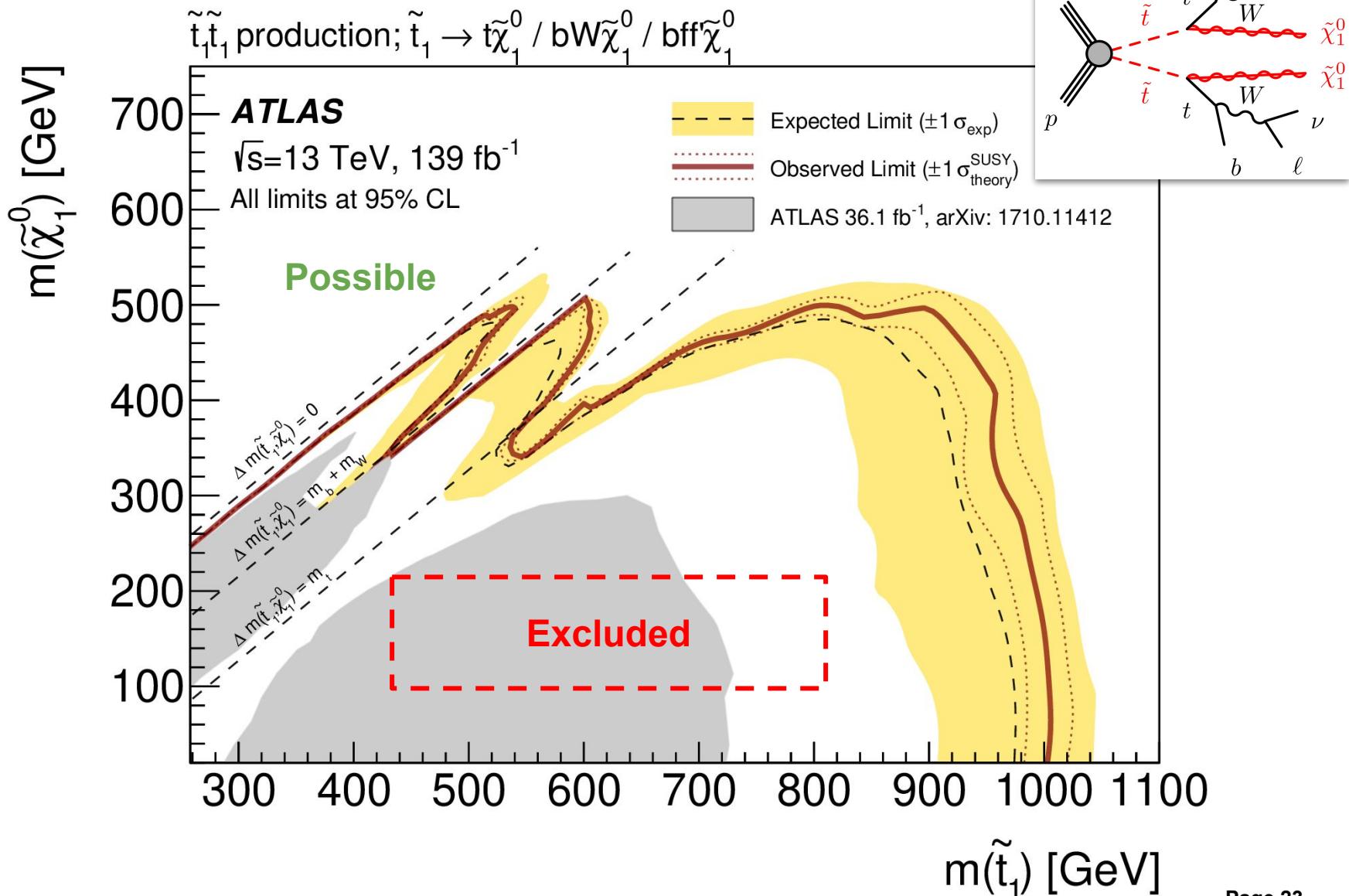


Two example  
BSM  
predictions!

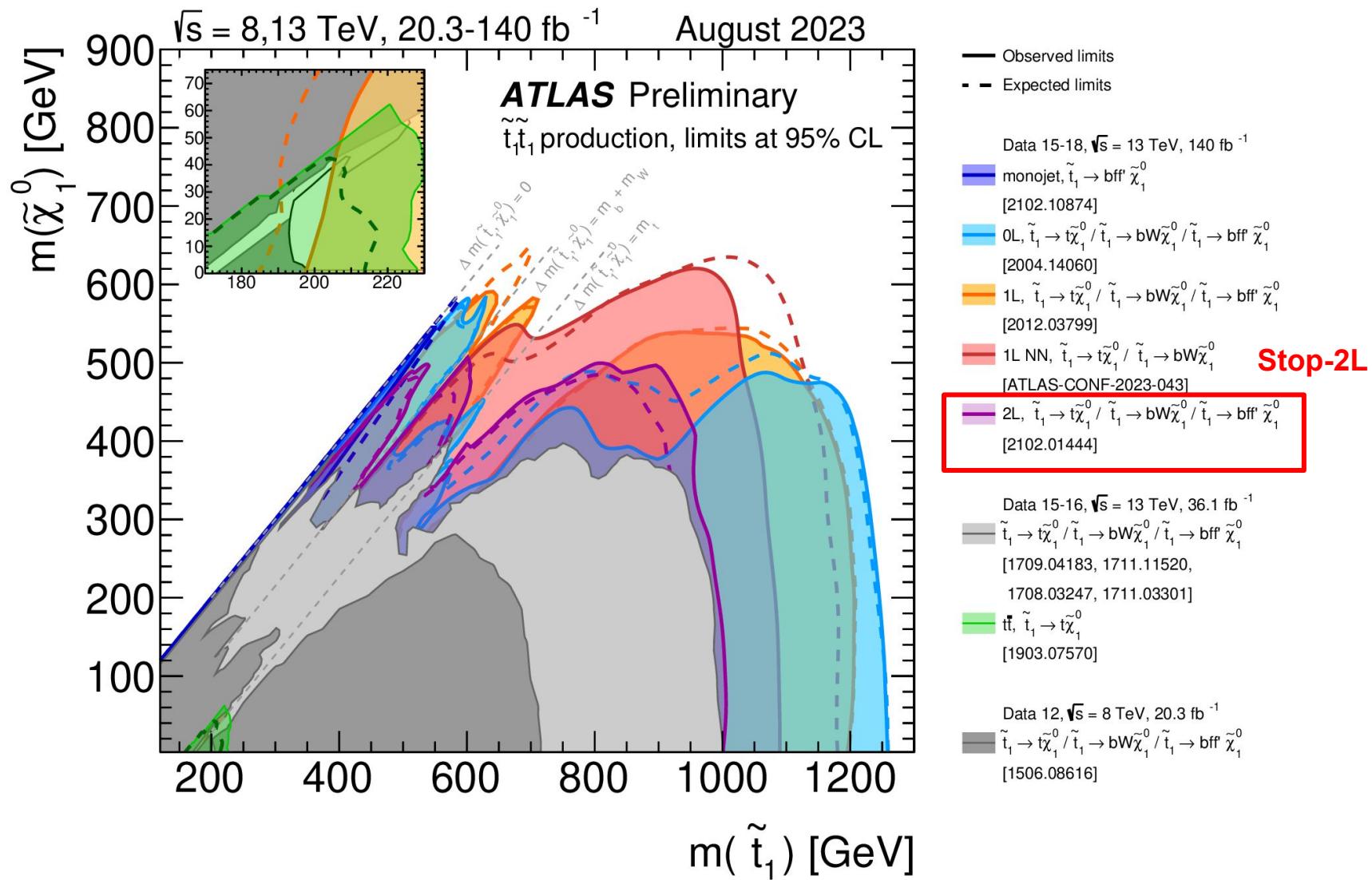
agreement  
of data & SM  
prediction  
→ didn't find BSM  
physics **within**  
**precision** (i.e.  
cyan / move  
models excluded)

# Stop-2L: results

Cannot exclude a model if it is hidden in the uncertainty\* → exclusion limits



# There are many SUSY searches out there



# There are many SUSY searches out there

- SUSY particle masses between 100 GeV and 1 TeV largely excluded
  - Some models under pressure; hierarchy problem only solved if masses “light”

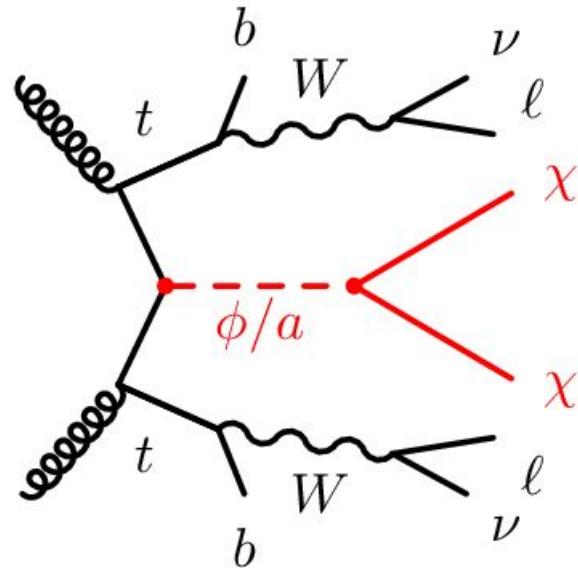
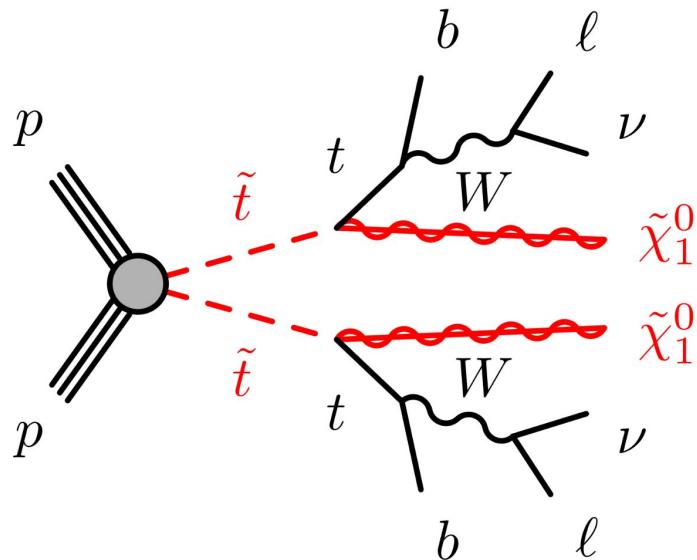
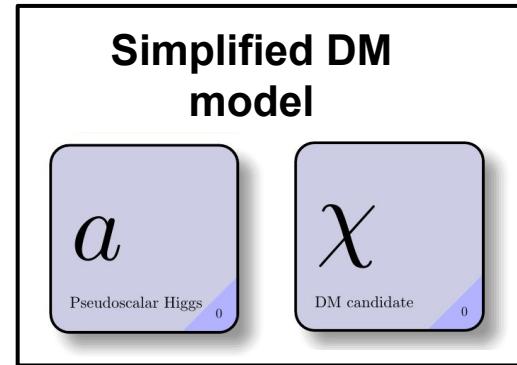
Model	Signature	$\int \mathcal{L} dt [fb^{-1}]$	Mass limit
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{g}\rightarrow q\tilde{\chi}_1^0$	0 e, $\mu$ mono-jet 2-6 jets 1-3 jets	$E_T^{\text{miss}}$ $E_T^{\text{miss}}$ 140 140
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, $\mu$	2-6 jets
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}W\tilde{\chi}_1^0$	1 e, $\mu$	2-6 jets
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	ee, $\mu\mu$	2 jets
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow q\tilde{q}WZ\tilde{\chi}_1^0$	0 e, $\mu$ SS e, $\mu$	7-11 jets 6 jets
	$\tilde{g}\tilde{g}, \tilde{g}\rightarrow t\bar{t}\tilde{\chi}_1^0$	0-1 e, $\mu$ SS e, $\mu$	3 b 6 jets
		$E_T^{\text{miss}}$ 140 140	
3 <sup>rd</sup> gen. squarks direct production	$\tilde{b}_1\tilde{b}_1$	0 e, $\mu$	2 b
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{\chi}_2^0 \rightarrow b h\tilde{\chi}_1^0$	0 e, $\mu$	6 b 2 $\tau$ 2 b
	$\tilde{l}_1\tilde{l}_1, \tilde{l}_1\rightarrow\tilde{\chi}_1^0$	0-1 e, $\mu$	$\geq 1$ jet
	$\tilde{l}_1\tilde{l}_1, \tilde{l}_1\rightarrow W\tilde{\chi}_1^0$	1 e, $\mu$	3 jets/1 b
	$\tilde{l}_1\tilde{l}_1, \tilde{l}_1\rightarrow\tilde{\tau}_1 bv, \tilde{\tau}_1\rightarrow\tau\tilde{G}$	1-2 $\tau$	2 jets/1 b
	$\tilde{l}_1\tilde{l}_1, \tilde{l}_1\rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c}\rightarrow c\tilde{\chi}_1^0$	0 e, $\mu$ 0 e, $\mu$	2 c mono-jet
	$\tilde{l}_1\tilde{l}_1, \tilde{l}_1\rightarrow\tilde{\chi}_2^0, \tilde{\chi}_2^0\rightarrow Z/h\tilde{\chi}_1^0$	1-2 e, $\mu$	1-4 b
	$\tilde{l}_2\tilde{l}_2, \tilde{l}_2\rightarrow\tilde{l}_1 + Z$	3 e, $\mu$	1 b
		$E_T^{\text{miss}}$ 140 140	
EW direct	$\tilde{\chi}_1^+\tilde{\chi}_2^0$ via WZ	Multiple $\ell/jets$ ee, $\mu\mu$	$E_T^{\text{miss}}$ $E_T^{\text{miss}}$ 140 140
	$\tilde{\chi}_1^+\tilde{\chi}_2^\mp$ via WW	2 e, $\mu$	$E_T^{\text{miss}}$ 140
	$\tilde{\chi}_1^+\tilde{\chi}_2^0$ via Wh	Multiple $\ell/jets$	$E_T^{\text{miss}}$ $E_T^{\text{miss}}$ 140 140
	$\tilde{\chi}_1^+\tilde{\chi}_1^-$ via $\tilde{\ell}_L/\tilde{\nu}$	2 e, $\mu$	$E_T^{\text{miss}}$ 140
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau}\rightarrow\tau\tilde{\chi}_1^0$	2 $\tau$	$E_T^{\text{miss}}$ 140
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell}\rightarrow\ell\tilde{\chi}_1^0$	2 e, $\mu$ ee, $\mu\mu$	0 jets $\geq 1$ jet
	$H\tilde{H}, \tilde{H}\rightarrow h\tilde{G}/Z\tilde{G}$	0 e, $\mu$ 4 e, $\mu$ 0 e, $\mu$ 2 e, $\mu$	$\geq 3$ b 0 jets $\geq 2$ large jets $\geq 2$ jets
		$E_T^{\text{miss}}$ 140 140 140 140	

100 GeV

1 TeV

# Searching for Dark Matter with Stop-2L

- In a simplified Dark Matter model, can produce Dark Matter from a mediator radiating off an inner top-quark line
- This mediator decays into 2x DM particle
- Final state the same as in Stop-2L!!!
- Simultaneously search for SUSY and DM!!!

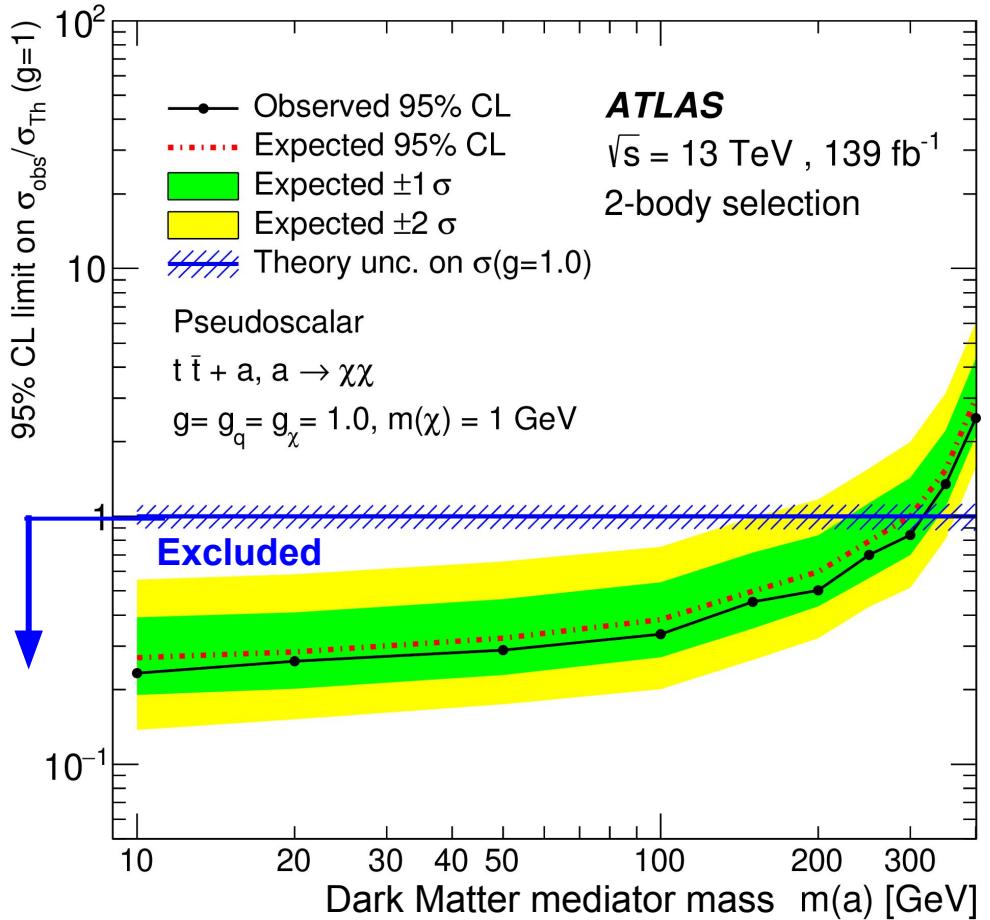


# Stop-2L Dark Matter results

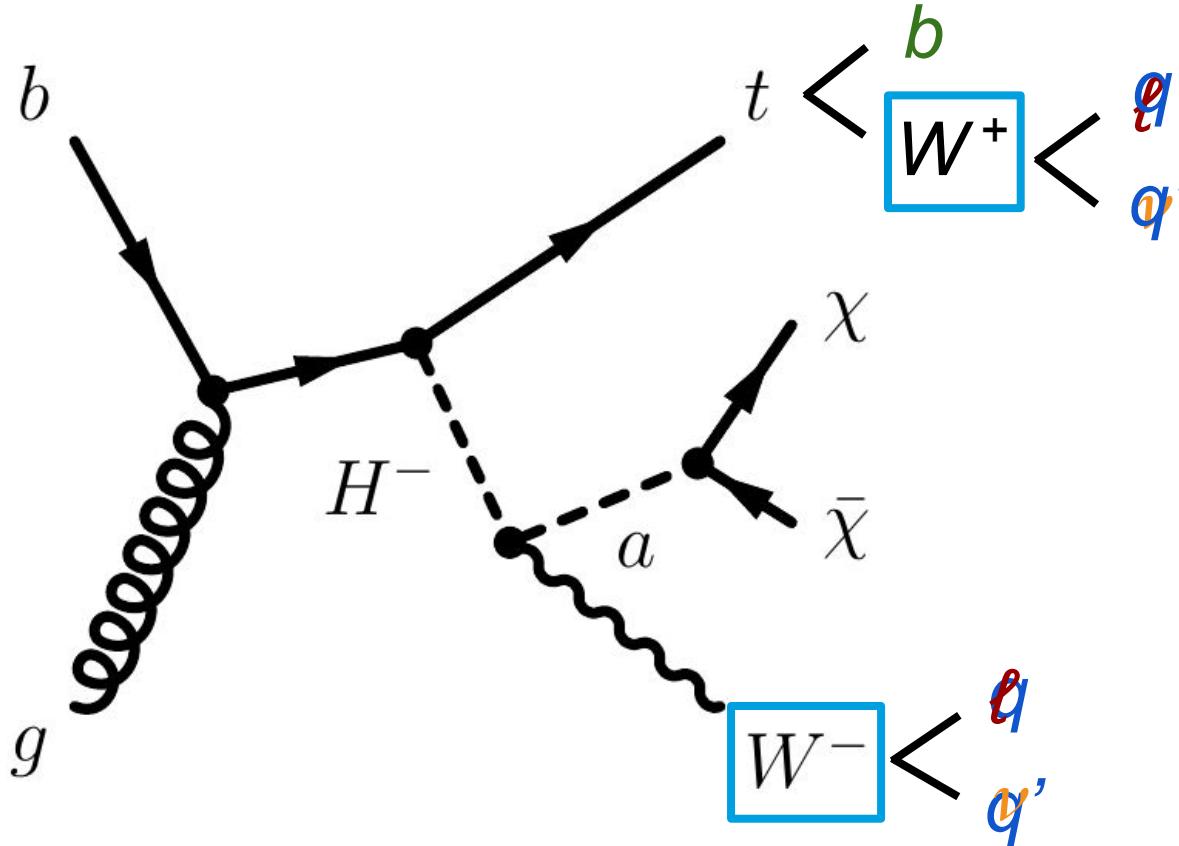
- Exclude DM mediator masses up to 300 GeV

= “maximum factor with which I can multiply the BSM prediction so it is hidden in the uncertainty”\*

If below 1: not hidden in uncertainty!  
So excluded!



# Searching for Dark Matter in tW+DM

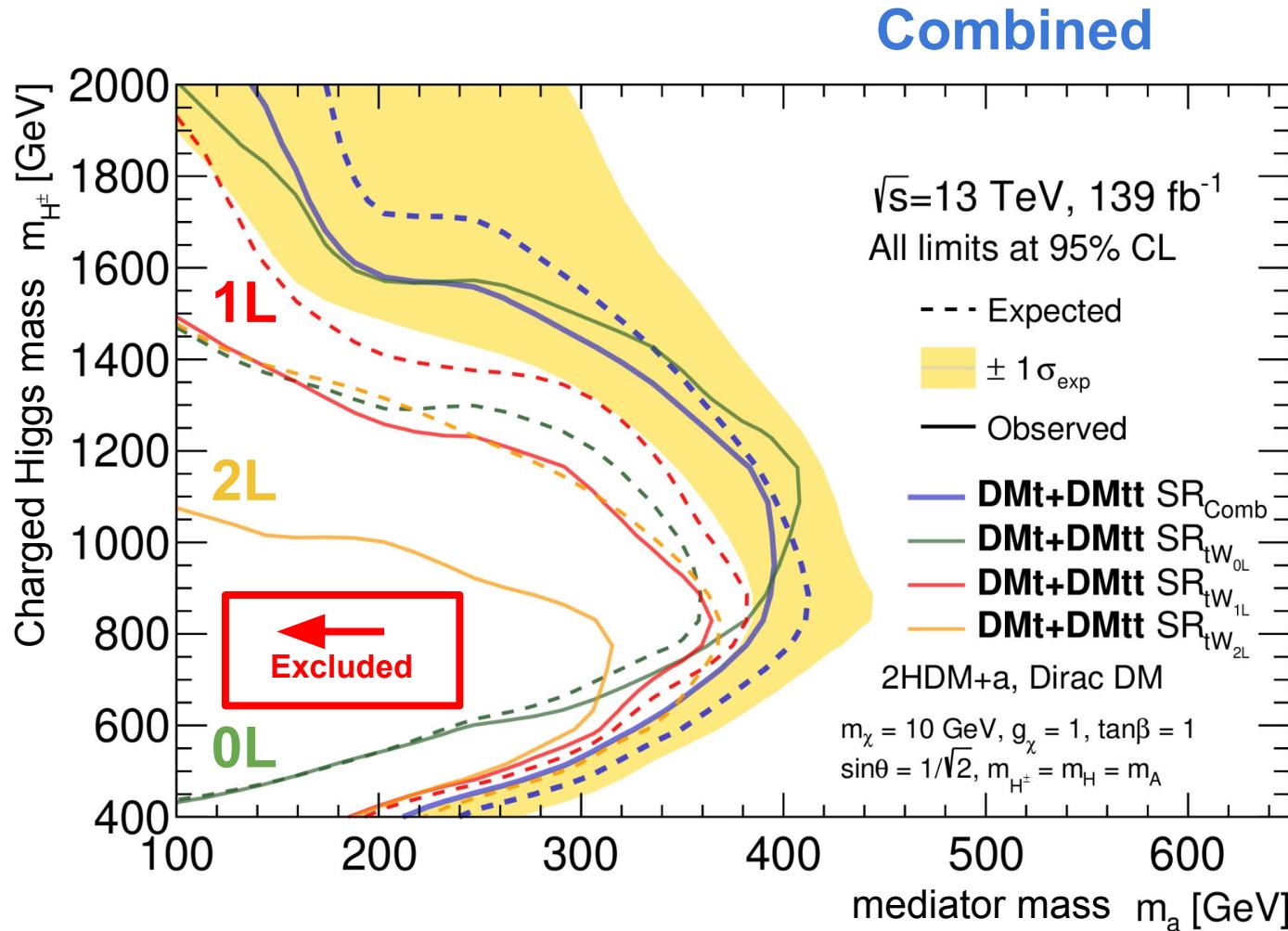


2HDM+a		
$h^0$ Scalar Higgs 125.1 GeV	$H^\pm$ Charged Higgs unknown $\pm 1$	$\chi$ DM candidate
$H^0$ Scalar Higgs	$A$ Pseudoscalar Higgs	$a$ Pseudoscalar Higgs

- tW+DM high x-sec
- Two **W-bosons**: zero, one or two visible leptons ( $e, \mu$ )
- **0L, 1L, 2L channel**  
→ basically 3 separate searches
- To get maximum sensitivity  
→ combine the three searches

# tW+DM: results

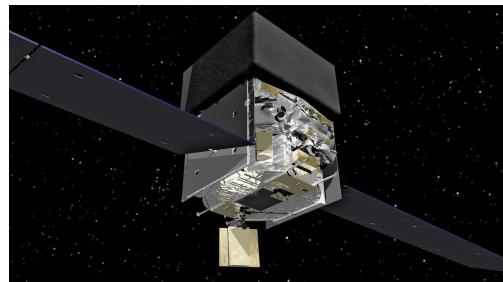
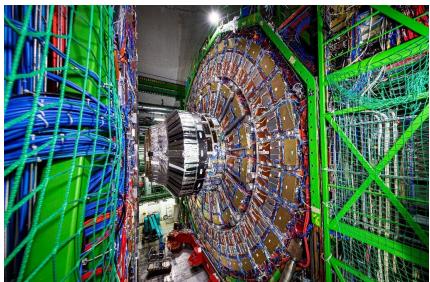
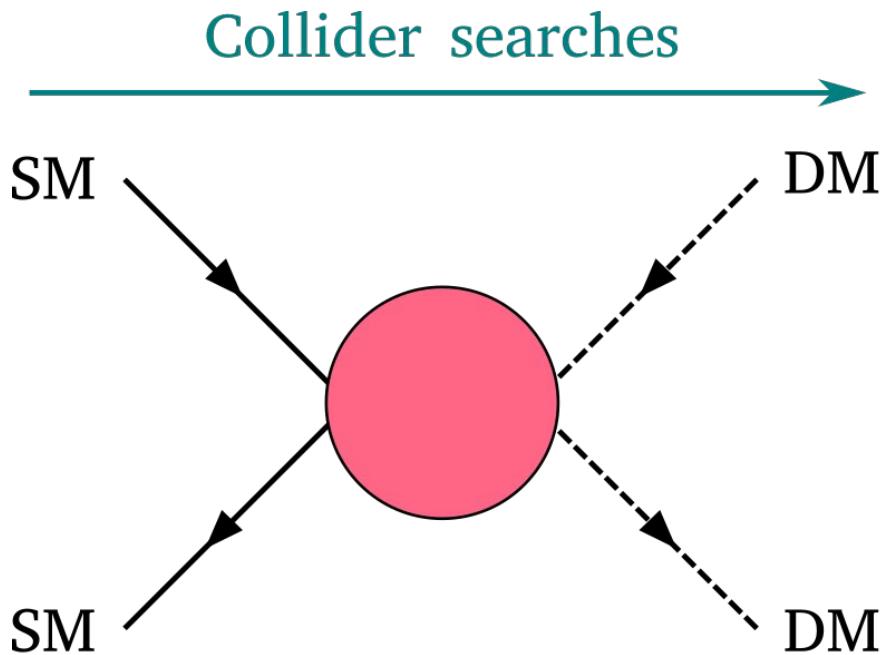
- Exclude DM mediator masses up to 300 GeV



# BREAK

(5-10 mins)

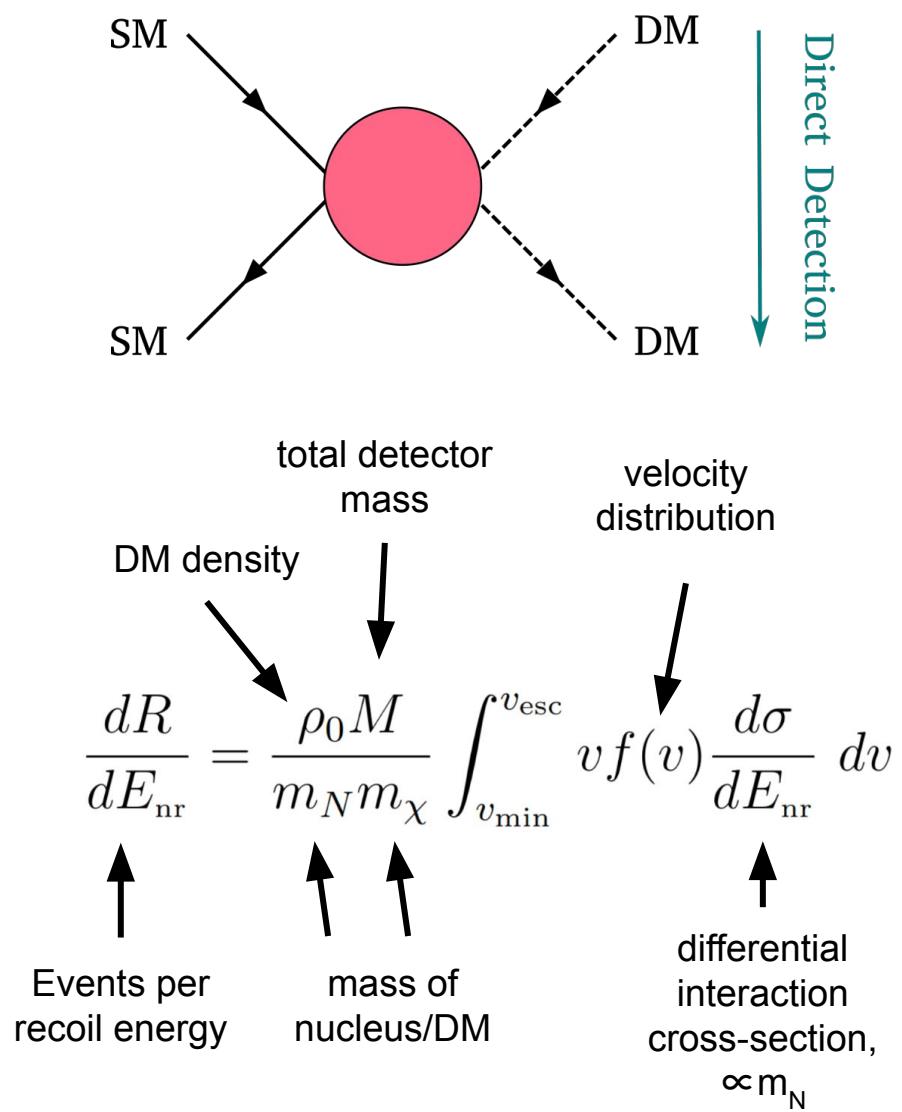
# Other ways to search for Dark Matter



# Direct detection Dark Matter searches

[arxiv:1903.03026](https://arxiv.org/abs/1903.03026)

- Assume: permanent flow of DM particles through planet earth
- **DM scatters with nuclei** (as no EM interaction) → **measurable**
- Number of scattered DM particles depends on:
  - Density of DM,  $\rho_0 = 0.3 \text{ GeV}/\text{c}^2/\text{cm}^3$
  - DM velocity w.r.t. earth  $\sim 220 \text{ km/s}$
  - DM-nuclei cross-section
  - Mass of DM & interacting nucleus
  - **Mass of detector → higher better**
  - **Spin of DM particle and spin sensitivity of target material**
  - **Minimum recoil energy sensitivity of detector → lower better**

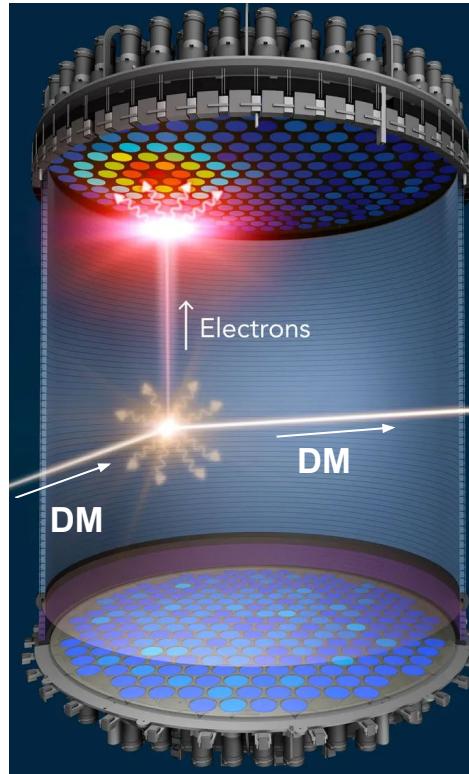
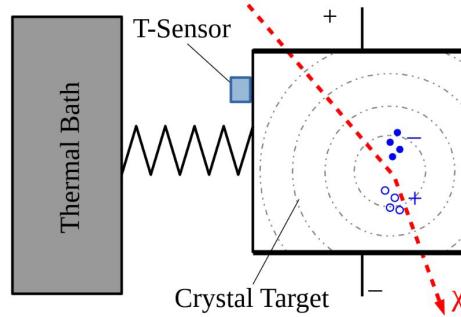


# Experimental setups of DD experiments

- Measure recoil of nuclei by DM → phonons/heat
- Backgrounds:
  - nuclear recoils by neutrons, cosmic muons
  - particles ( $\gamma$ ,  $\beta^\pm$ ) from radioactive decays → often interact via EM
- Use very pure, non-activated, well-selected materials & shielding
- Measure two signals: one from EM recoil (e.g. charge; higher for bkg.); one from nuclear recoil (e.g. heat)

Image credit bottom:

[https://th-thumbnailer.cdn-si-edu.com/3Qxw0k7nJ0cC7BUwbFSP5zHO6w8=/fit-in/1600x0/https%3A%2F%2Ftf-cmsv2-smithsonianmag-media.s3.amazonaws.com%2Ffile%2F42%2Fd5%2F42d5a303-b006-4972-9efa-92797a25ba9c%2F31667821088\\_f762d2c200\\_o.jpg](https://th-thumbnailer.cdn-si-edu.com/3Qxw0k7nJ0cC7BUwbFSP5zHO6w8=/fit-in/1600x0/https%3A%2F%2Ftf-cmsv2-smithsonianmag-media.s3.amazonaws.com%2Ffile%2F42%2Fd5%2F42d5a303-b006-4972-9efa-92797a25ba9c%2F31667821088_f762d2c200_o.jpg)



## Cryogenic detectors:

- At very low T (mK)
- Scattering DM increases T
- Bkg. ionises material → separate signal
- Materials: Ge, Si

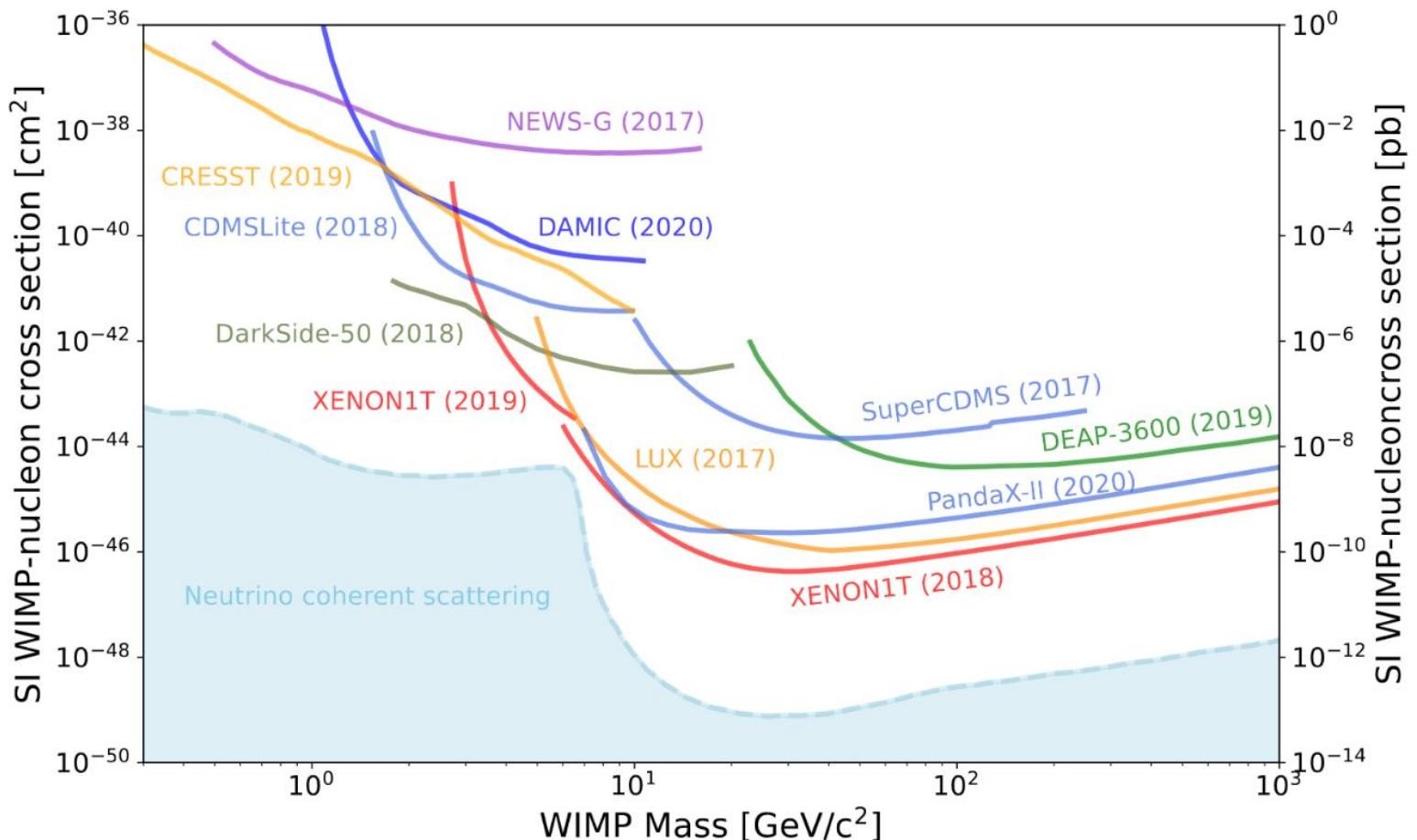
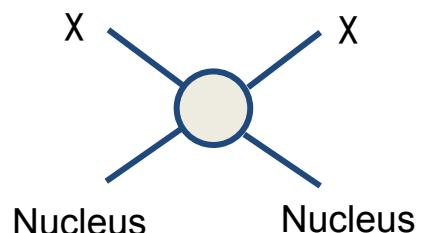
## Noble liquid detectors:

- E.g. liquid Ar, Xe
- Phonons from DM → photons + ionisation
- Example XENON experiment
- Detect photons with PMTs at the side
- Detect ionisation by E-field + scintillation signal at the top
- Ratio of photons/ionisation different for DM / bkg.

# DD experiments – results

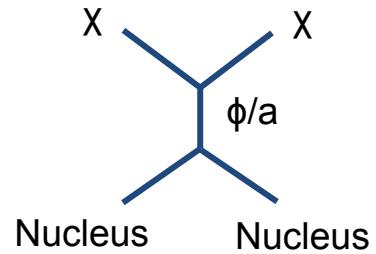
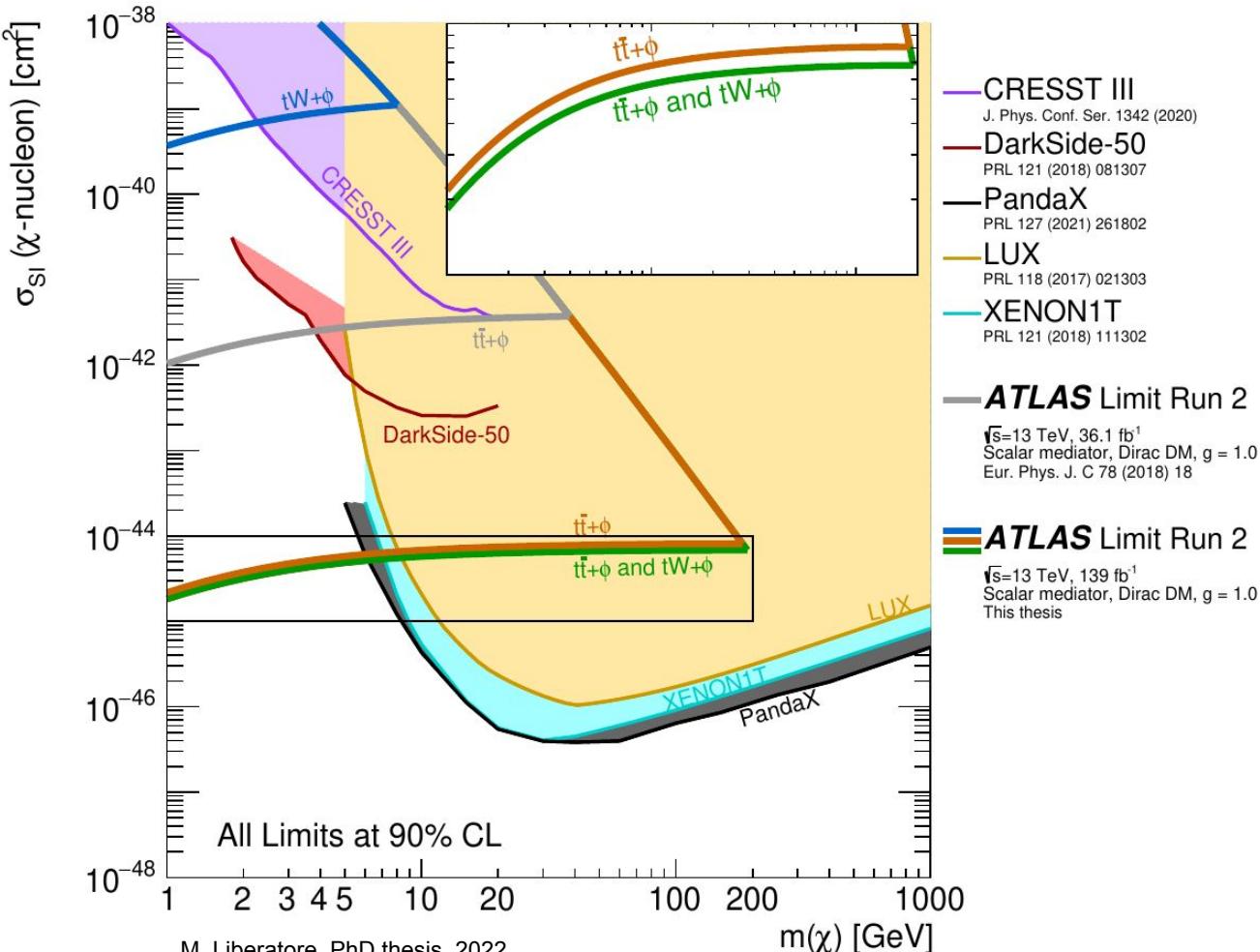
[rpp2022-rev-dark-matter](#)

- Direct detection experiment results interpreted in terms of Effective Field Theories, no assumption on interaction mechanism of DM and nucleus
- The neutrino floor is not that far!



# DD experiments – comparison to collider

- Comparing to collider: need to use a “simplified DM model” → additional mediator  $\phi/a$
- Collider searches dominate at low masses!

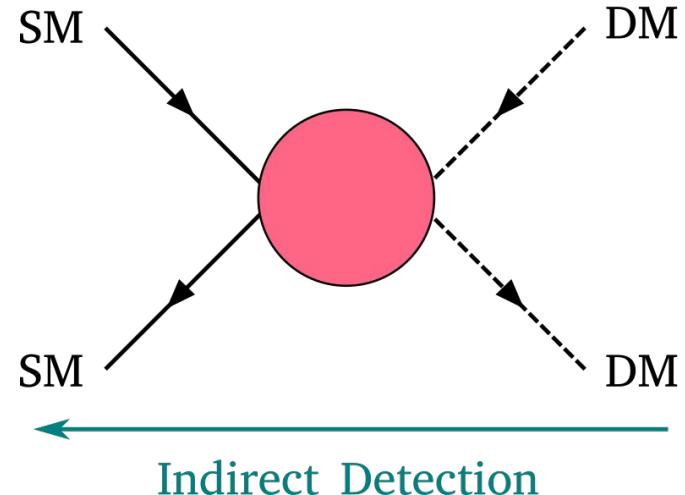


# Indirect detection Dark Matter searches

[rpp2022-rev-dark-matter](#)

- In areas of high mass density: Dark Matter will annihilate → indirect detection
- Annihilation yields signals such as photon pairs, neutrinos, baryons
- High mass density areas: sun, galaxy centers, ...
- Production rate of IDD events depends on:
  - Annihilation rate / cross-section
  - Dark Matter density
  - The number of final state particles

$$\Gamma_f^A = c \frac{\rho_{\text{DM}}^2}{m_{\text{DM}}^2} \langle \sigma v \rangle N_f^A$$

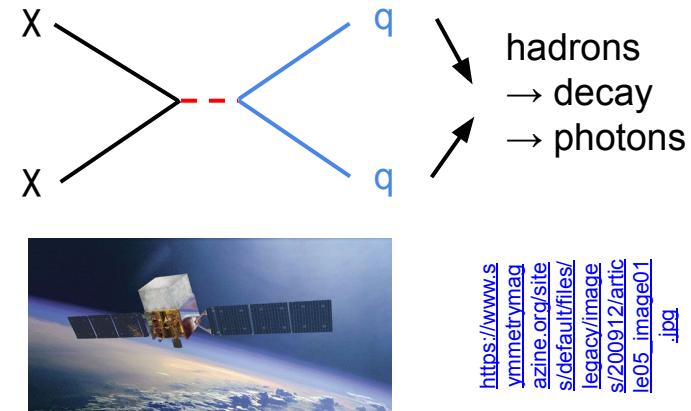


[https://www.youcanseethemilkyway.com/wp-content/uploads/2023/01/Massive\\_Black\\_Hole\\_at\\_the\\_Center\\_of\\_the\\_Milky\\_Way.jpg](https://www.youcanseethemilkyway.com/wp-content/uploads/2023/01/Massive_Black_Hole_at_the_Center_of_the_Milky_Way.jpg)

# Experimental techniques of indirect detection

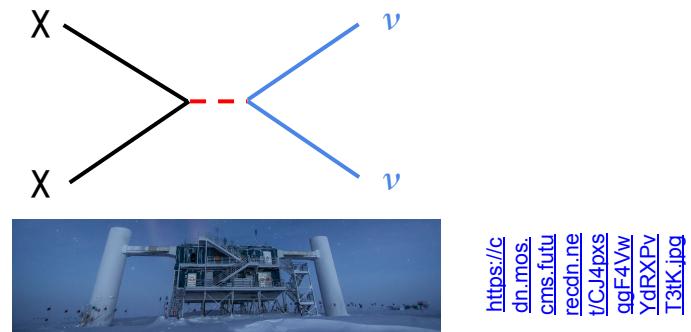
- **Photons**

- produced e.g. when  $\text{DM} + \text{DM} \rightarrow \text{quarks}$
- search for high energy photons (“gamma rays”) e.g. from the galactic center
- experiments: e.g. FERMI-LAT



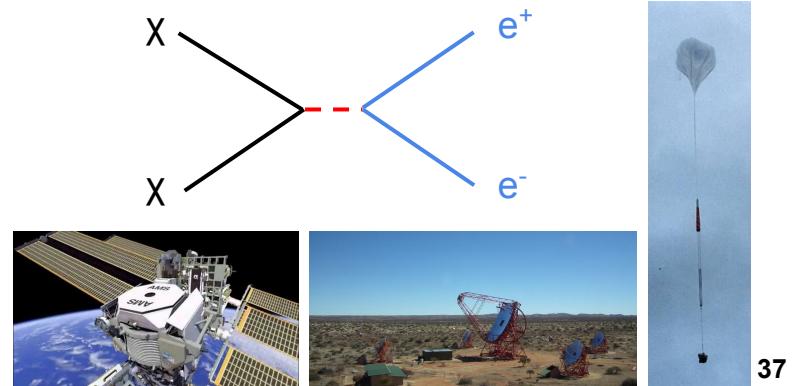
- **Neutrinos**

- produced e.g. in the sun
- search for high energy neutrinos from sun
- experiments: e.g. IceCube, Kamiokande



- **Antiparticles**

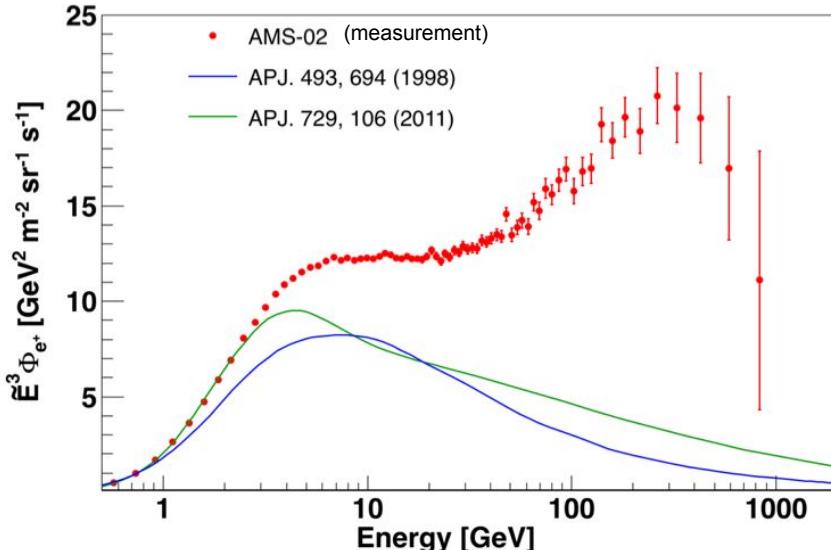
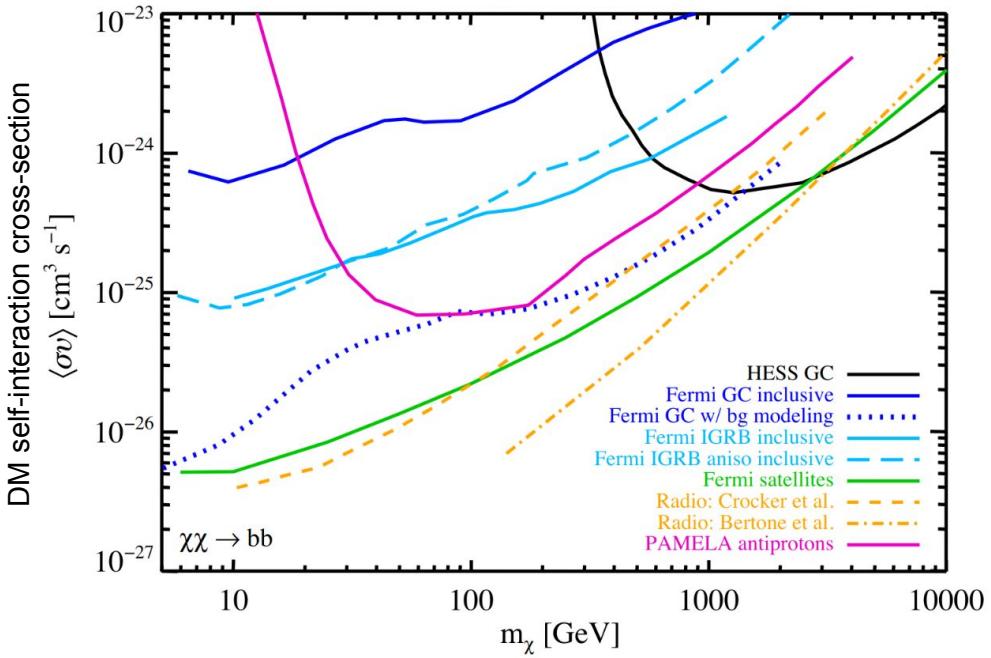
- produced e.g. if  $\text{DM} + \text{DM} \rightarrow e^+e^-$
- charged particles deflected in universe  
→ generally search for antiparticles
- experiments: e.g. HEAT, AMS, HESS



# Indirect detection: results

[arxiv:1604.00014](https://arxiv.org/abs/1604.00014)

- Covered mass range comparable to DD
- Exciting result: excess of positrons observed by multiple experiments (AMS, FERMI, PAMELA)
- Yet unclear if it is due to Dark Matter, astrophysical origin investigated

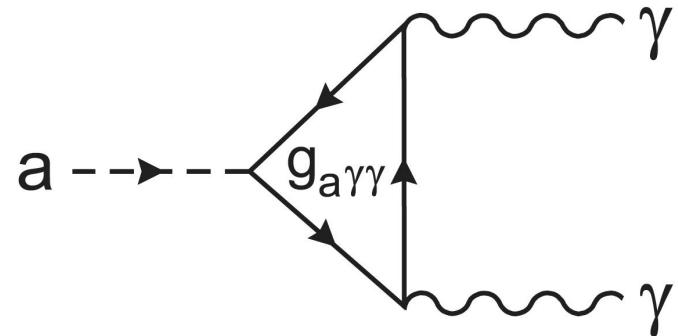
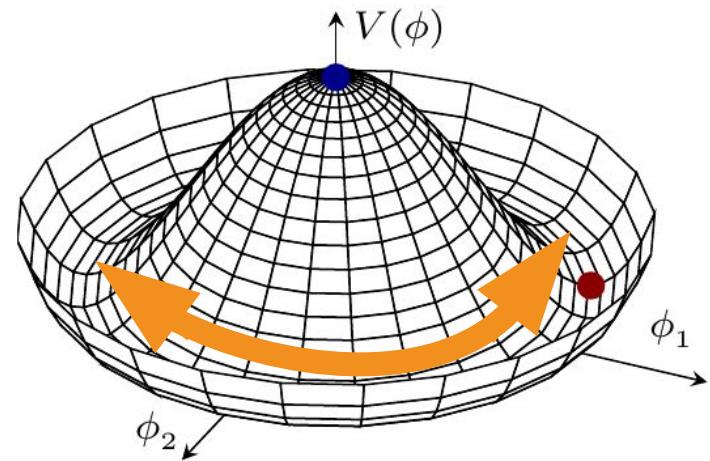
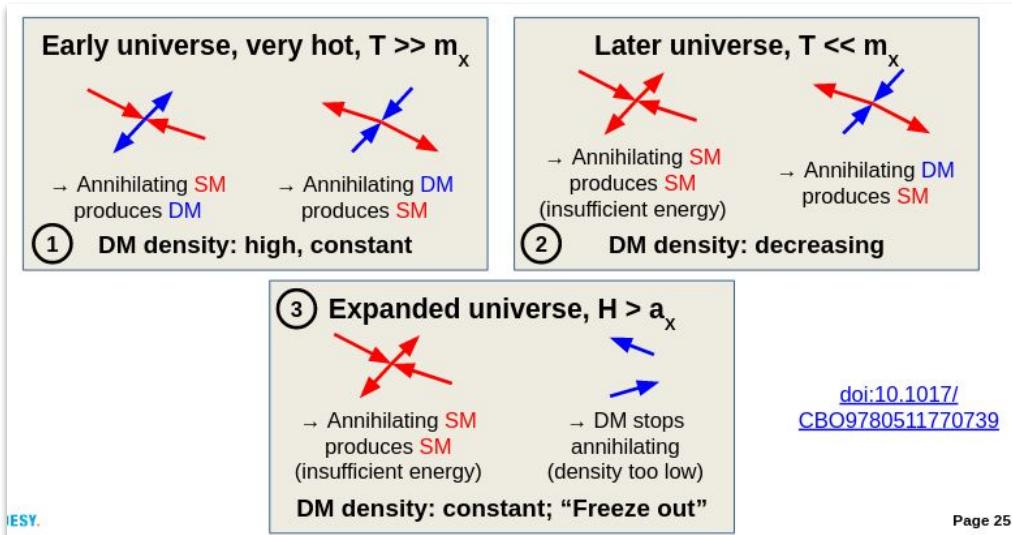


<https://ams02.space/sites/default/files/inline-images/20PhysRevLett.122.041102.Fig5.png>

# Axions

(Lecture 1)

- Searches so far considered WIMP DM  
→ mass & x-sec. similar to weak bosons
- Also covered Axion Dark Matter last time  
→ much lighter DM candidates  
→ origin from solution of strong CP probl.
- Interact with SM matter by 2-photon vertex

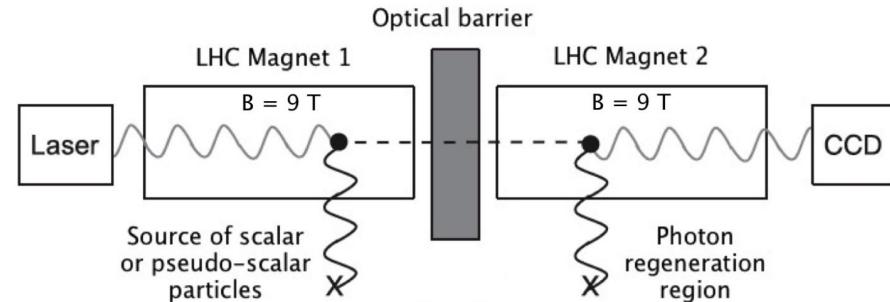


# Axions search experiments

- Idea of Axion experiments:
  - can interact with Axions by exposing them to strong magnetic fields (=high photon flux)
  - can produce Axions by using strong lasers and strong magnetic fields
- ALPs experiment: shine laser into a strong magnetic field
  - photons converted to Axions
  - photons stopped by wall, Axions pass through
  - on the other side: another strong magnet → Axions converted to photons



[https://particle-physics.desy.de/sites/site\\_particle-physics/content/e221990/e222445/e228223/e228225/e228229/ALPS\\_qer.jpg](https://particle-physics.desy.de/sites/site_particle-physics/content/e221990/e222445/e228223/e228225/e228229/ALPS_qer.jpg)



[arxiv:1410.2566](https://arxiv.org/abs/1410.2566)

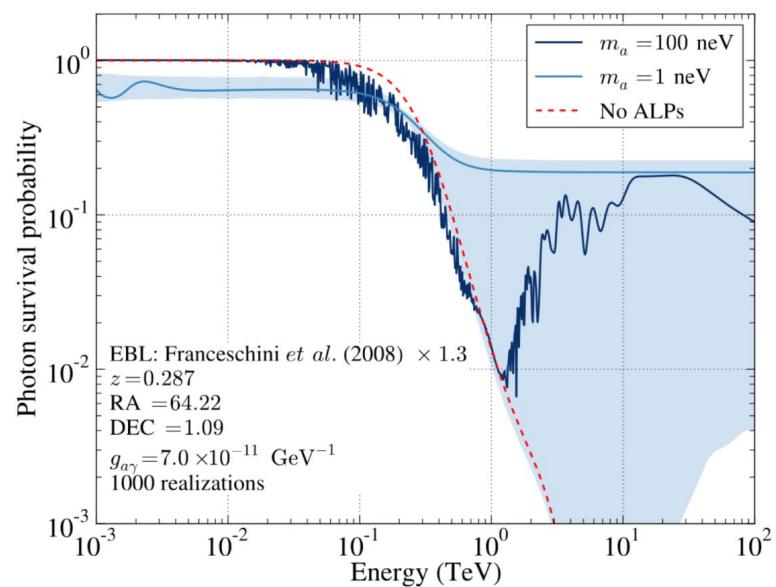
# Searching for Axions from & in the universe

- CAST / (Baby-)IAXO experiments: search for Axions produced in sun
  - use a large magnet to point at the sun: Axions converted to photons
  - measure these photons
- The magnetic fields converting (photons → Axions) & (Axions → photons) can also be of astrophysical origin
  - reduced attenuation of high energy photons



[https://mediastream.cern.ch/MediaArchive/Photo/Public/2002\\_0209017/0209017\\_01/0209017\\_01-A4-at-144-dpi.jpg](https://mediastream.cern.ch/MediaArchive/Photo/Public/2002_0209017/0209017_01/0209017_01-A4-at-144-dpi.jpg)

DESY.

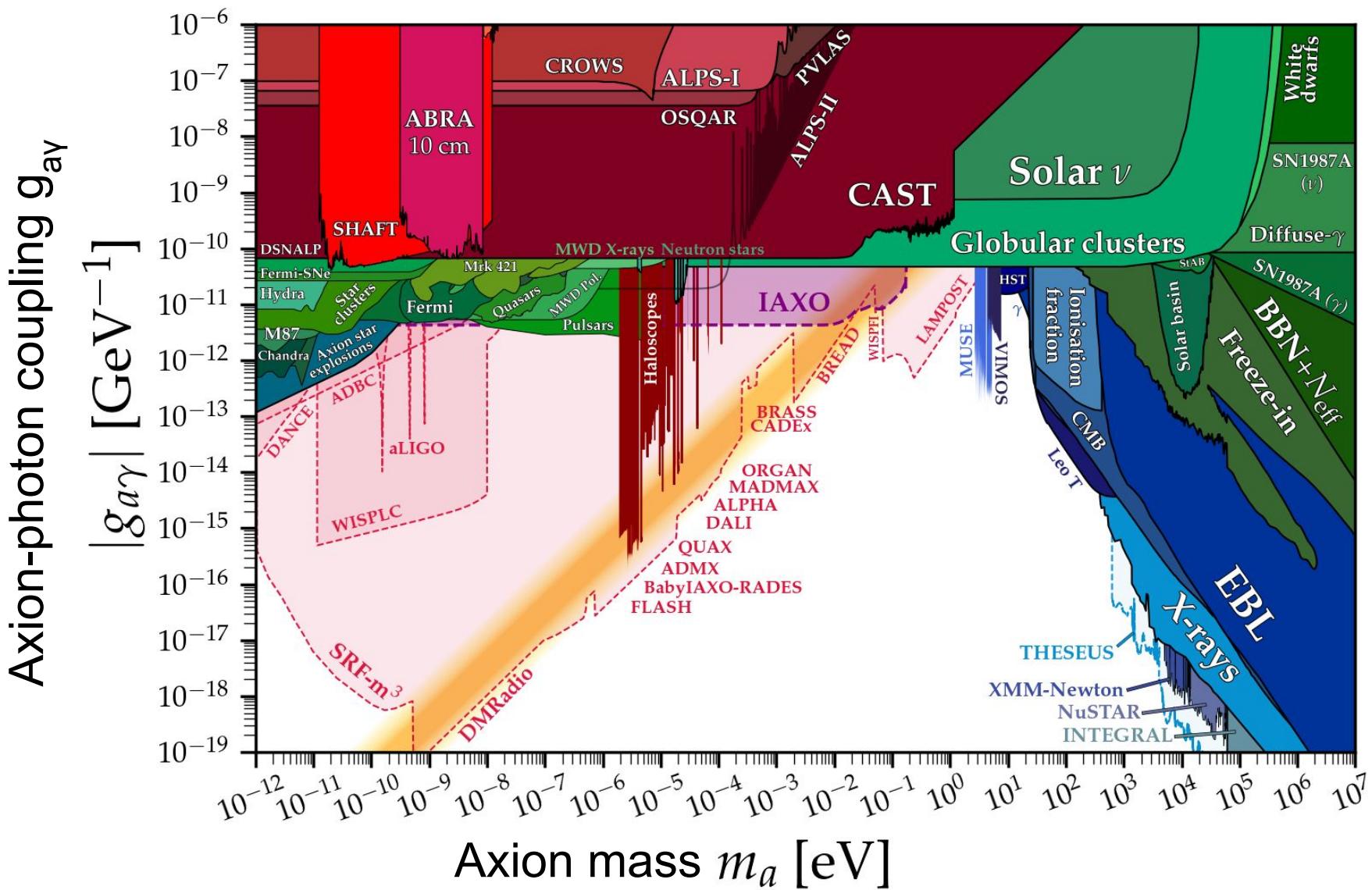


M. Meyer, The Opacity of the Universe for High and Very High Energy  $\gamma$ -Rays,  
<https://inspirehep.net/literature/1254304>

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# Axions: exclusion limits

Ciaran O'Hare, AxionLimits,  
doi:10.5281/zenodo.3932430, [Link](#)



# The magnetic moment of the muon

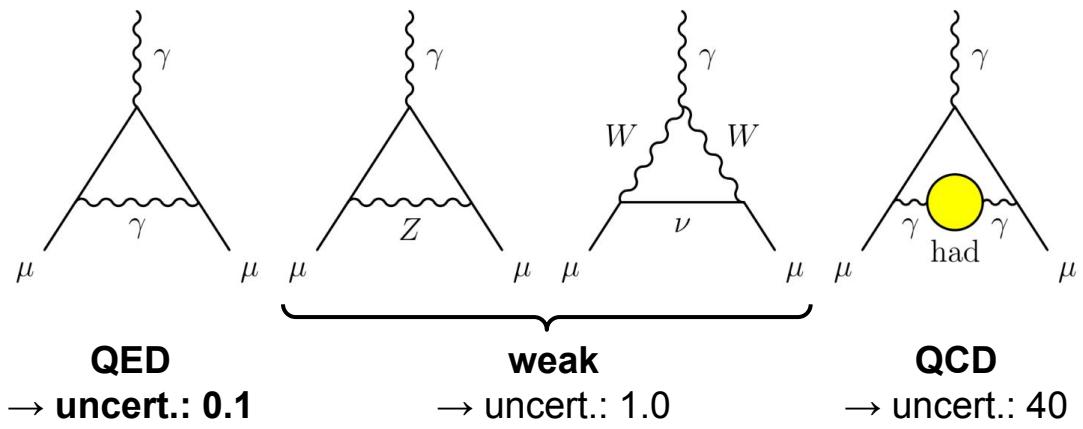
[rpp2022-rev-q-2-muon-anom-mag-moment](#)

- Recall: circular current = magnetic field → can be attributed to a magnetic moment
- Similarly: charged particles on circular orbits & spinning around themselves exhibit magnetic moment
- Magnetic moment basically indicates how strongly a particle is affected by a magnetic field

$$\mathbf{m}_S = -\frac{g_S \mu_B \mathbf{S}}{\hbar}, \quad \mu_B = \frac{e\hbar}{2m_e}$$

spin-vector  
constant

- For the muon:  $g_\mu \approx 2$ , loop effects →  $g > 2$
  - Measure the difference:  
 $a_\mu = (g_\mu - 2) / 2$   
→ **stress-test**
- Standard Model**  
(the loops it predicts)

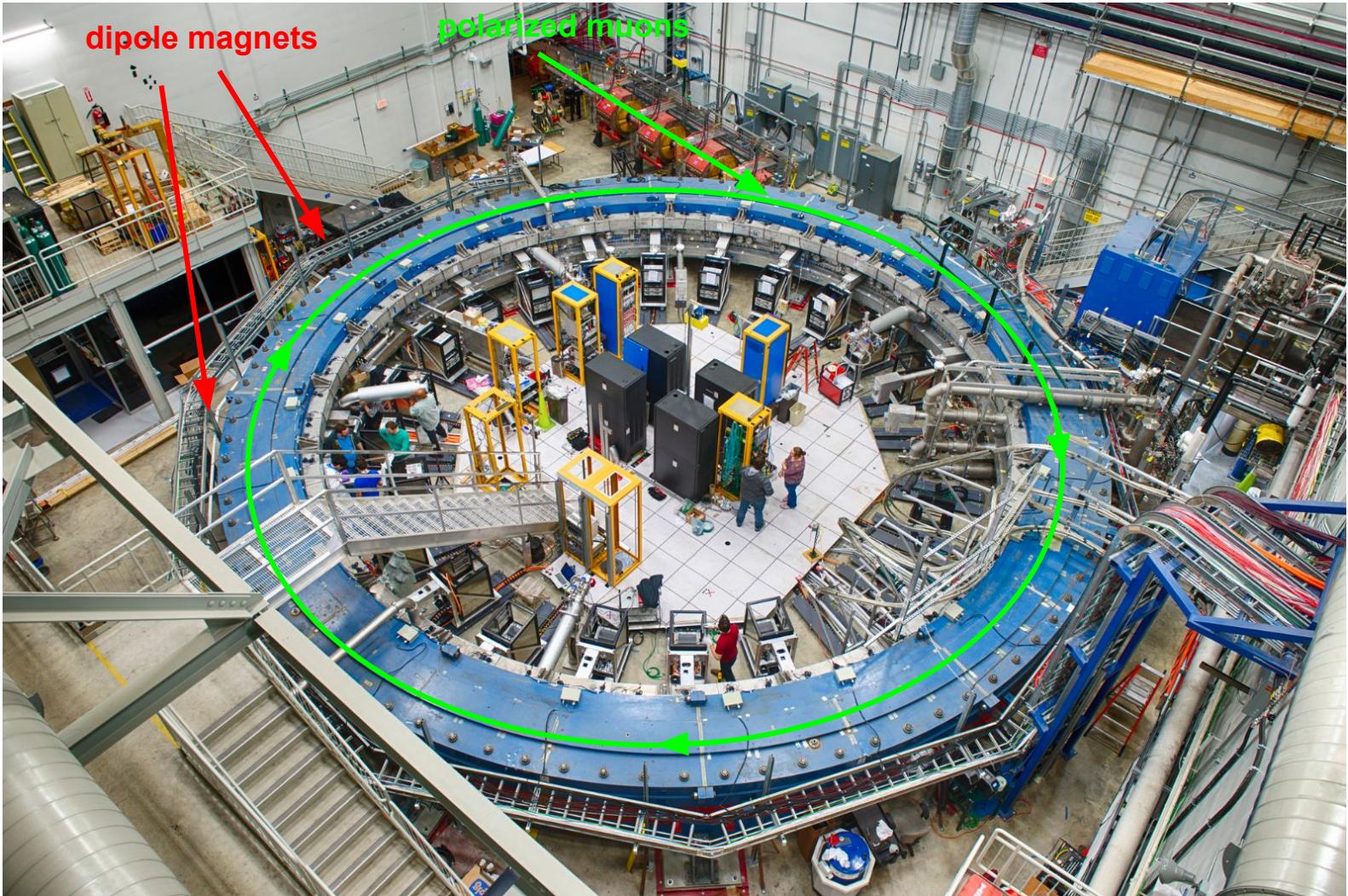


<https://www.particlebites.com/?p=8972>

# How to measure g-2

$$\mathbf{F} = q \mathbf{v} \times \mathbf{B}$$

- Produce muons, store ring via dipole magnets



[https://upload.wikimedia.org/wikipedia/commons/7/7c/Fermilab\\_g-2\\_%28E989%29\\_ring.jpg](https://upload.wikimedia.org/wikipedia/commons/7/7c/Fermilab_g-2_%28E989%29_ring.jpg)

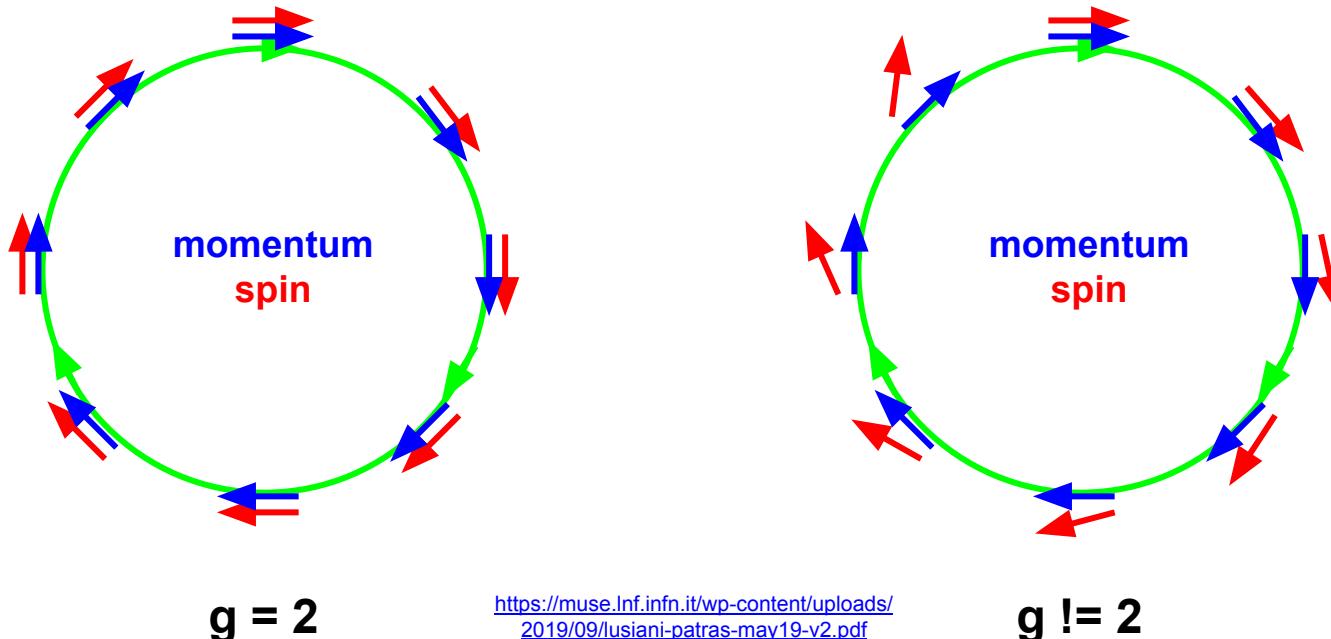
# How to measure g-2 (2)

[Sweigart, PhD thesis](#)

- Two frequencies:
  - rotation frequency of muons (cyclotron frequency)  $\omega_c$
  - precision frequency due to spin  $\omega_s$
- Their difference depends on g-2!!!  
→ measure  $\omega_a$  &  $\mathbf{B}$  → g-2 (\*)!!

$$\omega_a \equiv \omega_s - \omega_c = -a_\mu \frac{eB}{m_\mu}$$

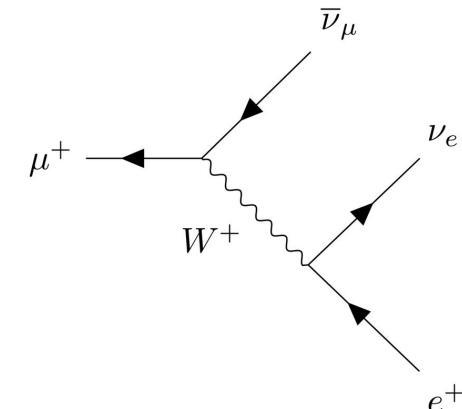
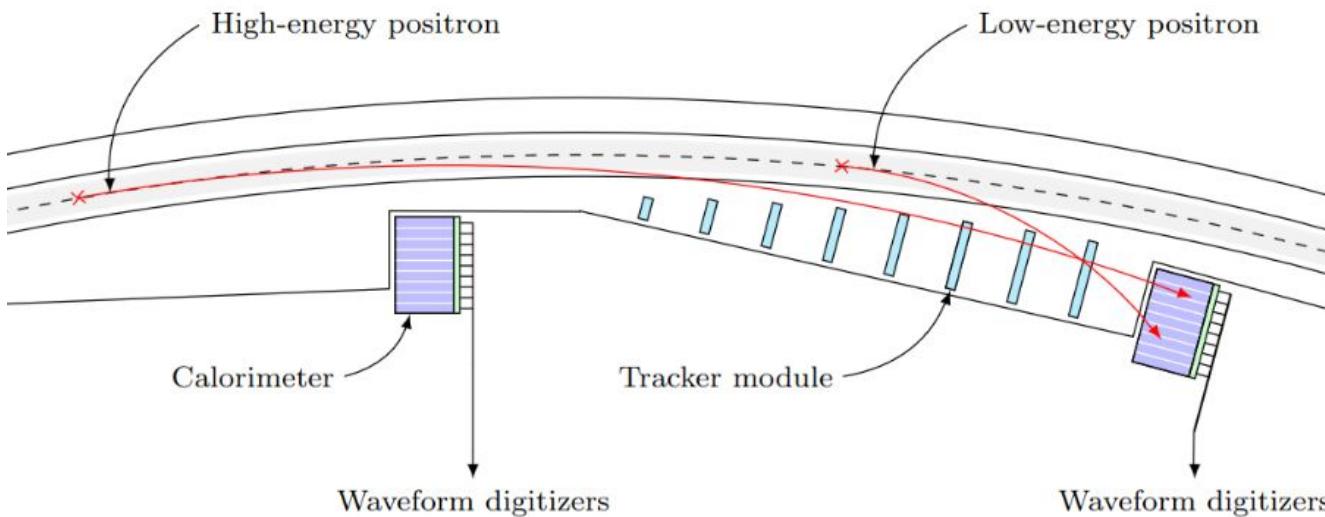
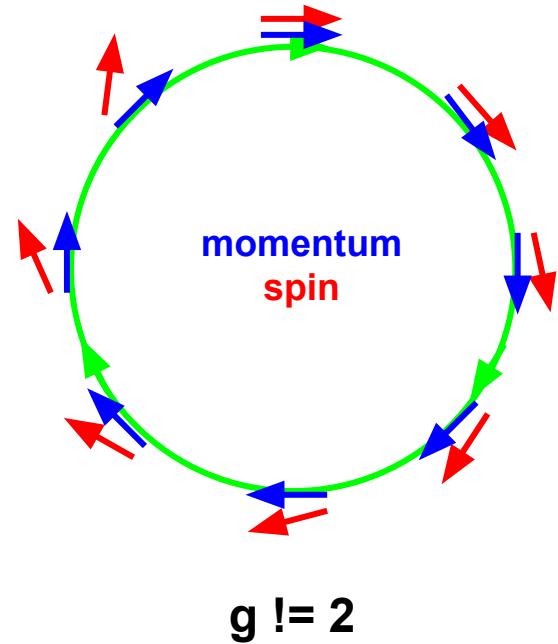
\*in reality you measure  $\omega_P$  - the Lamor frequency of the free proton - instead of  $\mathbf{B}$



<https://muse.lnf.infn.it/wp-content/uploads/2019/09/lusiani-patras-may19-v2.pdf>

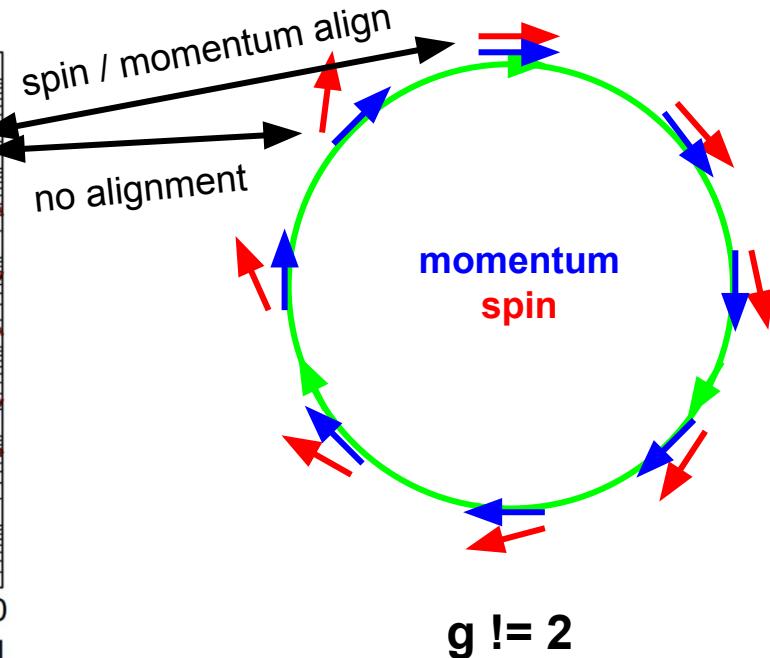
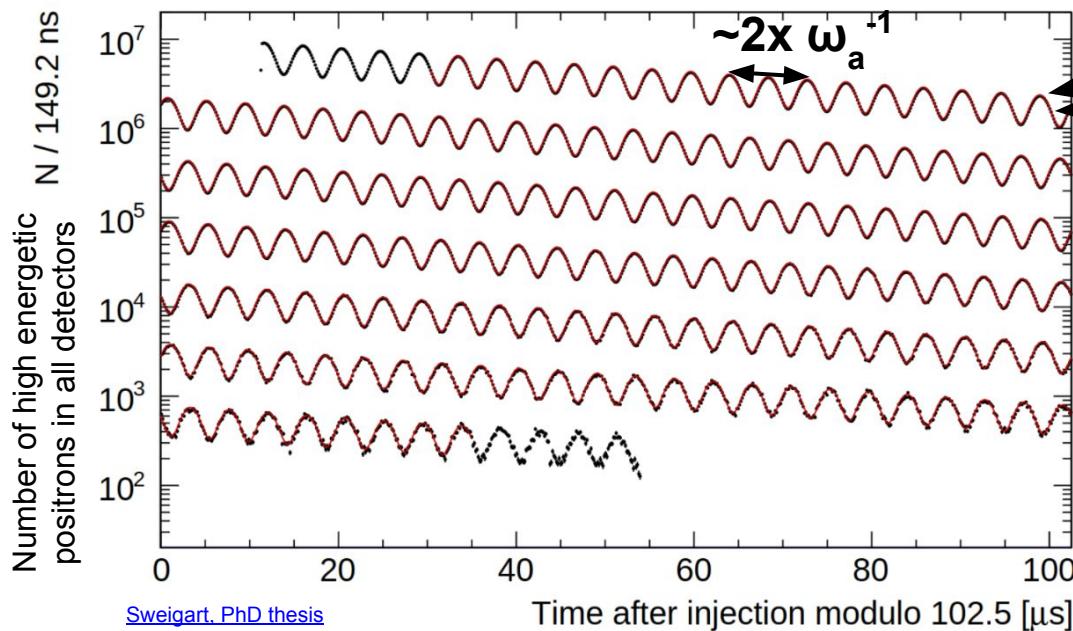
# How to measure g-2 (3)

- Muon is not stable and decays
- Measure energy of decay positrons  
→ Most energetic if muon momentum & spin align!
- If filter out most energetic positrons, see modulation with frequency  $\omega_a$ !!!

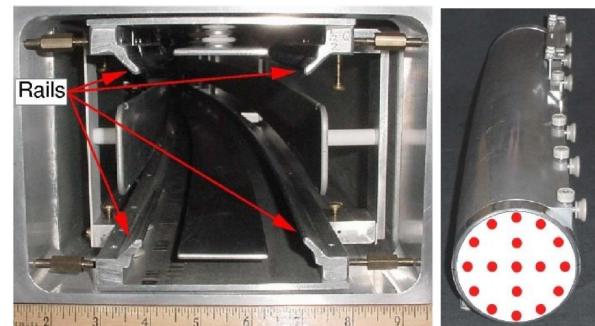


# How to measure g-2 (4)

<https://www.particlebites.com/?p=8999>

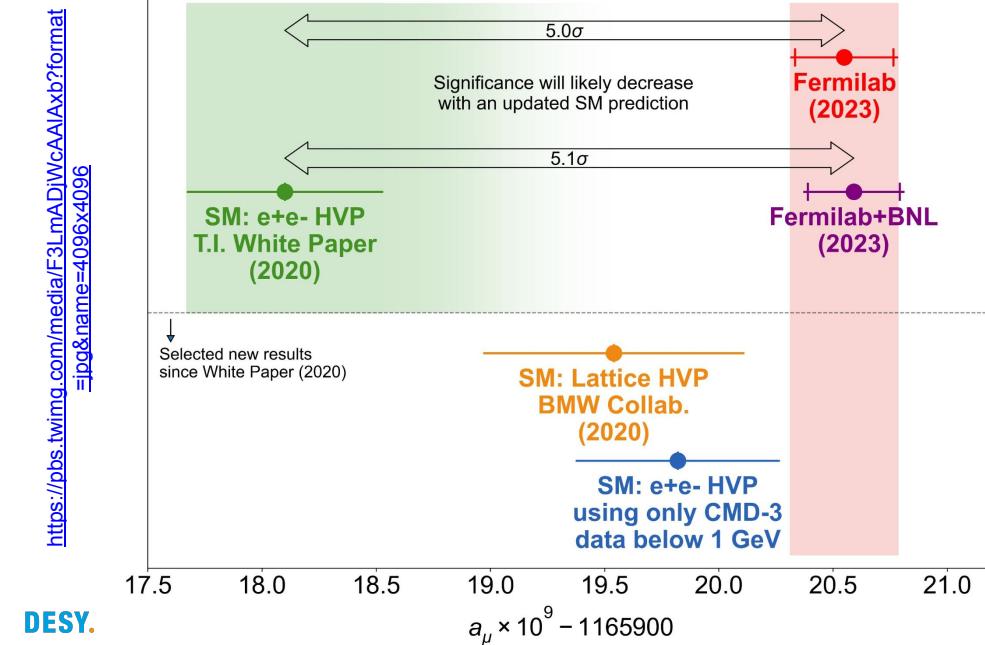
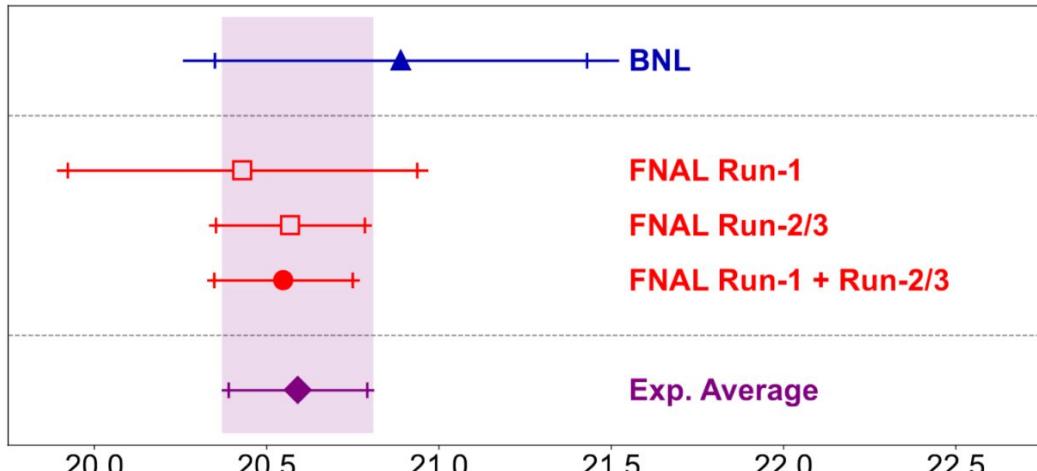


- If filter out most energetic positrons, see modulation with frequency  $\omega_a !!!$
- Besides  $\omega_a$ , must measure **B**!
- Measure magnetic field with special “trolley” measuring the field very precisely
- DESY every three days & many other tools



[https://conference.ippp.dur.ac.uk/event/999/contributions/5220/attachments/4218/4977/Schreckenberger\\_Planck2021.pdf](https://conference.ippp.dur.ac.uk/event/999/contributions/5220/attachments/4218/4977/Schreckenberger_Planck2021.pdf)

# Results of g-2



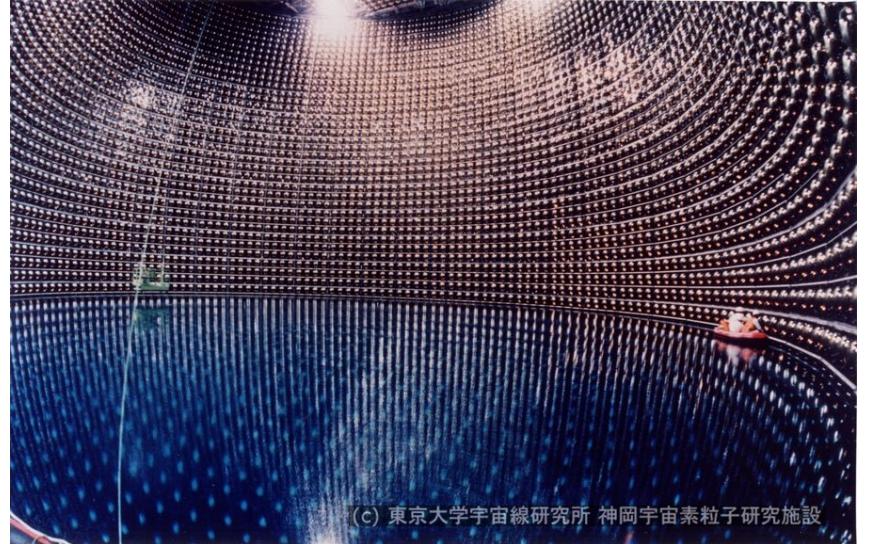
- g-2 measured extremely precisely to ~10 digits!
- 2 SM predictions: different QCD loop calculations
- Data-driven approach: an input is data from  $e^+e^-$ -colliders →  $5.1\sigma$  discrepancy to meas.
- The other is from lattice QCD → agrees with measurement! → but not fully reproduced
- Data-driven approach: new result which is compatible with measurement!?

**Exciting times for g-2!!**

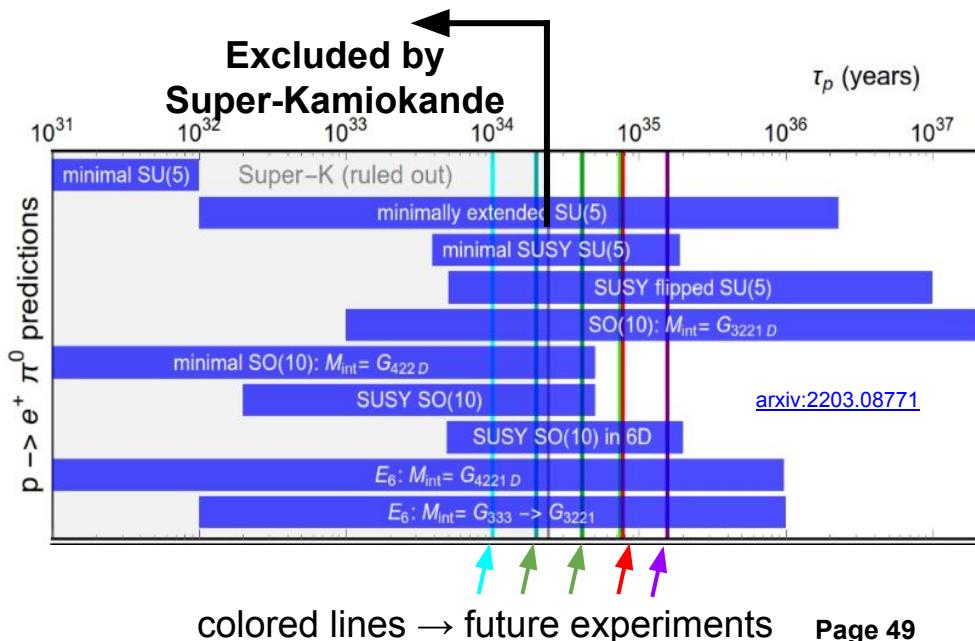
# Proton stability

[arxiv:2010.16098](https://arxiv.org/abs/2010.16098)

- Proton is stable in SM (Baryon number conservation)
- Grand Unified Theories can predict that the proton is not stable
- “Not stable” → lifetimes  $> 10^{30}$  y
- Search for proton decay e.g. with Super-Kamiokande
  - 27.5 T of Water → many protons
  - Search for  $p \rightarrow \ell^+ \pi^0$
  - Proton lifetime  $> 10^{34}$  y
- Multiple models excluded, but still room for SUSY ;)



<https://wechseltelekt.de/uploads/Super-K2-7d5e2bcc48802c113419e3c24de77cd03662a919.jpg>



# Summary

- Many, many searches / experiments targeting BSM physics
- Constraints on many model parameters, yet not clear what BSM physics is
- Many exciting experiments/results coming up investigating BSM physics further, e.g.
  - HL-LHC
  - XENON n-ton
  - ALPS II
  - BabyLAXO
  - g-2 full data set
  - Hyper Kamiokande
  - ...and many more!
- **What will we find?**



<https://thumbs.dreamstime.com/b/llama-closeup-funny-expression-face-mouth-open-appears-to-be-talking-humorous-96704071.jpg>

# Thank you

# Backup slides

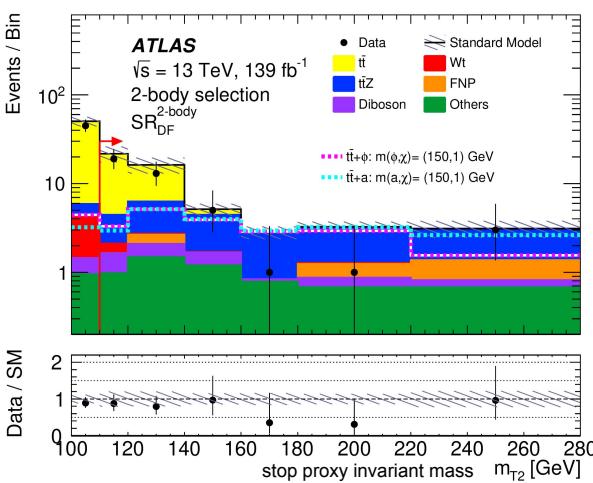
# Stop-2L: statistical analysis

SUSY-2018-08

## 1. Triggering



## 2. Event selection & background estimation

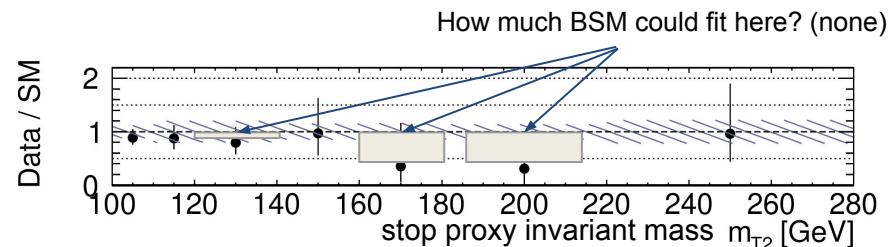


## 3. Statistical analysis

How compatible is the data with the Standard Model prediction? ([Data Analysis lecture](#))

$$\ell = -2 \ln \left( \frac{\mathcal{L}(\text{data}, H_1)}{\mathcal{L}(\text{data}, H_0)} \right)$$

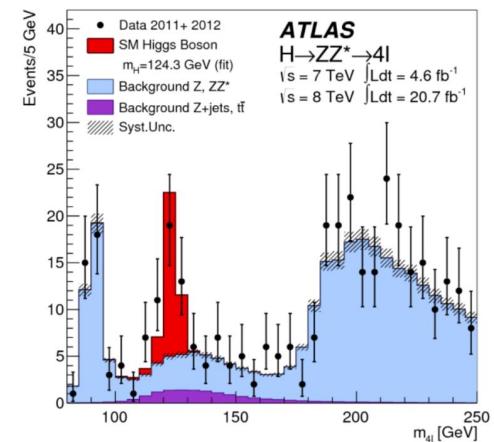
**3a) If very compatible** → up to which x-sec. can I exclude BSM physics? → *exclusion limits*



**3b) If not compatible**  
→ check  $10^{35}\times$  for mistakes  
→ **CELEBRATE!!!**



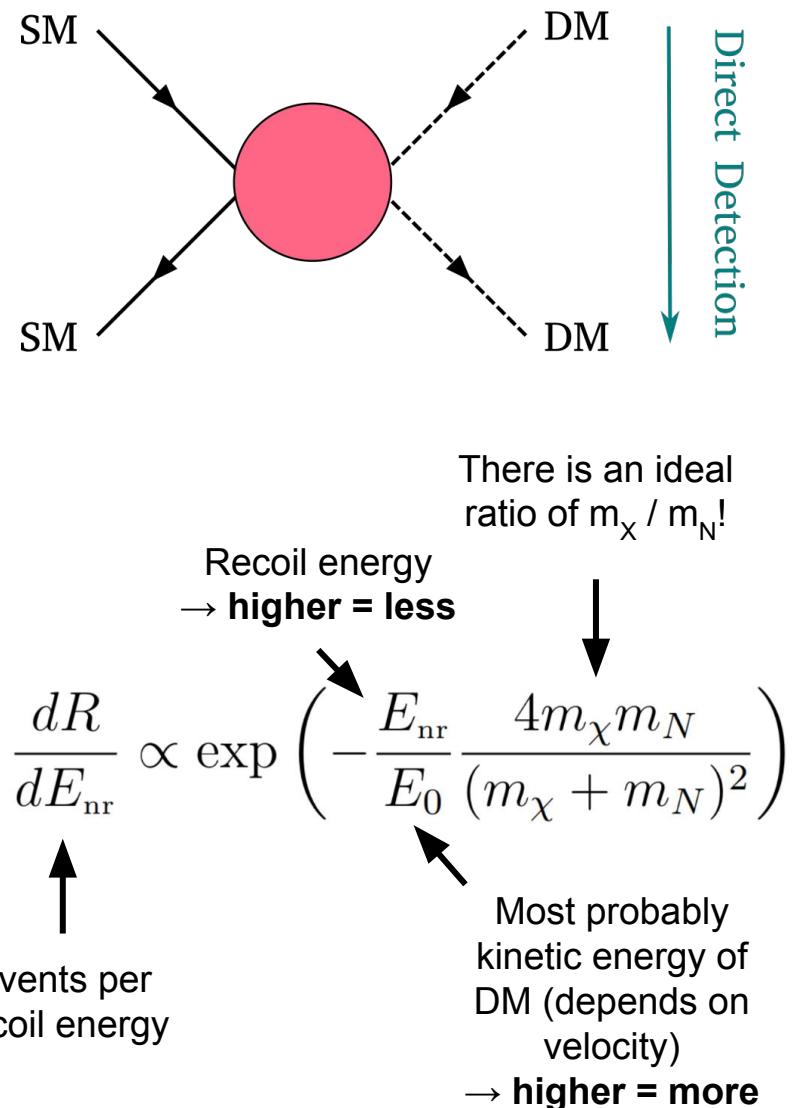
[https://media.istockphoto.com/id/13338482/d1efoto/dlampainen.jpg2?w=60&k=2&ke=70667&as=1&mt=151514GV701012\\_0&QICGIVMyPBeZ-8e](https://media.istockphoto.com/id/13338482/d1efoto/dlampainen.jpg2?w=60&k=2&ke=70667&as=1&mt=151514GV701012_0&QICGIVMyPBeZ-8e)



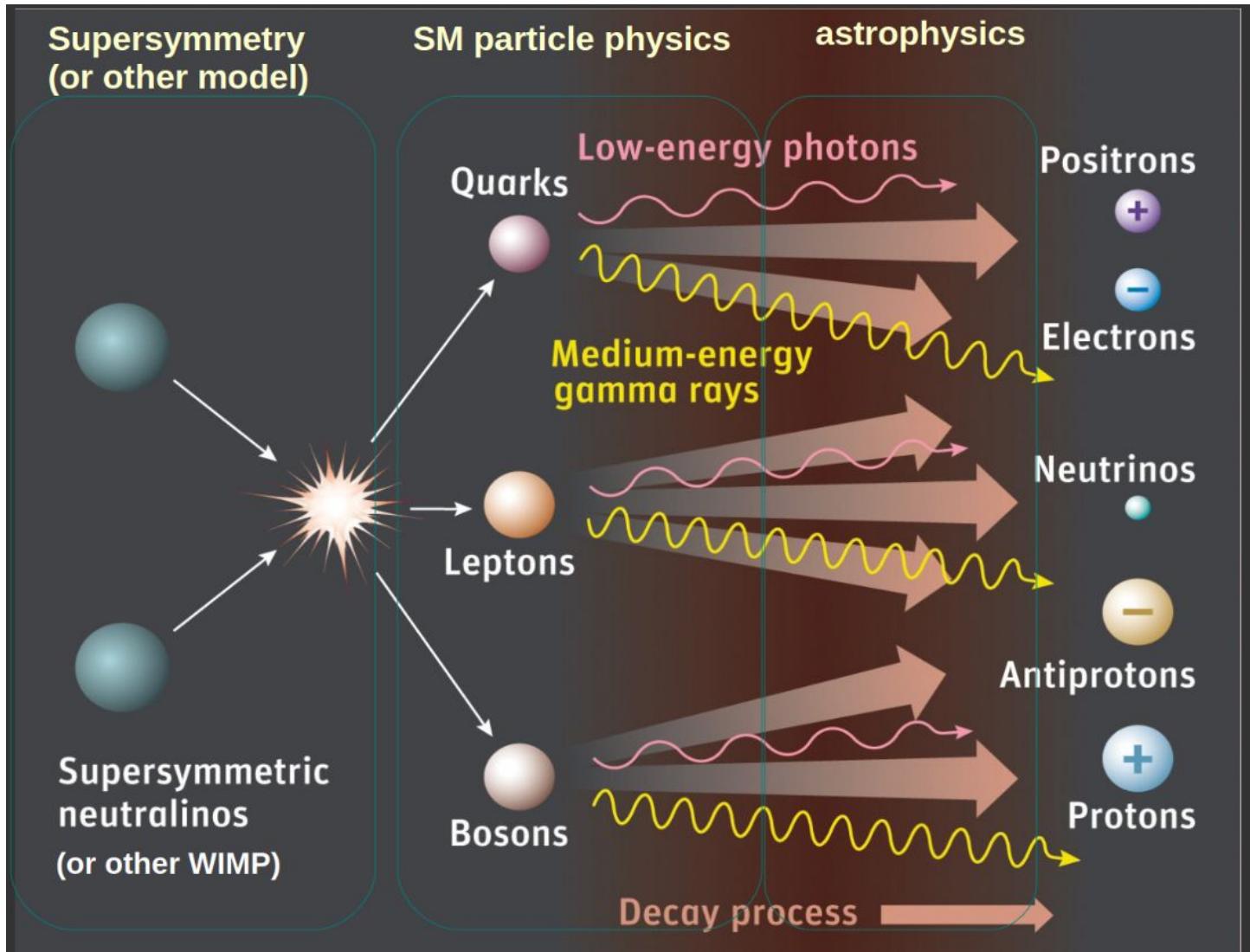
# Direct detection Dark Matter searches

[1903.03026](#)

- Assume: permanent flow of DM particles through planet earth
- DM scatters with SM particles → **measurable!**
- DM scatters with nuclei as no EM int.
- No. scattered DM particles depends:
  - Density of DM → from rotation curve:  $\rho_0 = 0.3 \text{ GeV}/c^2/\text{cm}^3$  ( $\pm 50\%$ )
  - DM velocity w.r.t. planet earth,  $\sim 220 \text{ km/s}$  with some variation
  - DM-nuclei cross-section
  - Mass of DM & interacting nuclei
  - Spin of DM particle and sensitivity of target material



# Indirect detection all in one slide



[https://www.mpi-hd.mpg.de/lin/events/isapp2011/pages/lectures/de\\_los\\_Heros.pdf](https://www.mpi-hd.mpg.de/lin/events/isapp2011/pages/lectures/de_los_Heros.pdf)

## Contact

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