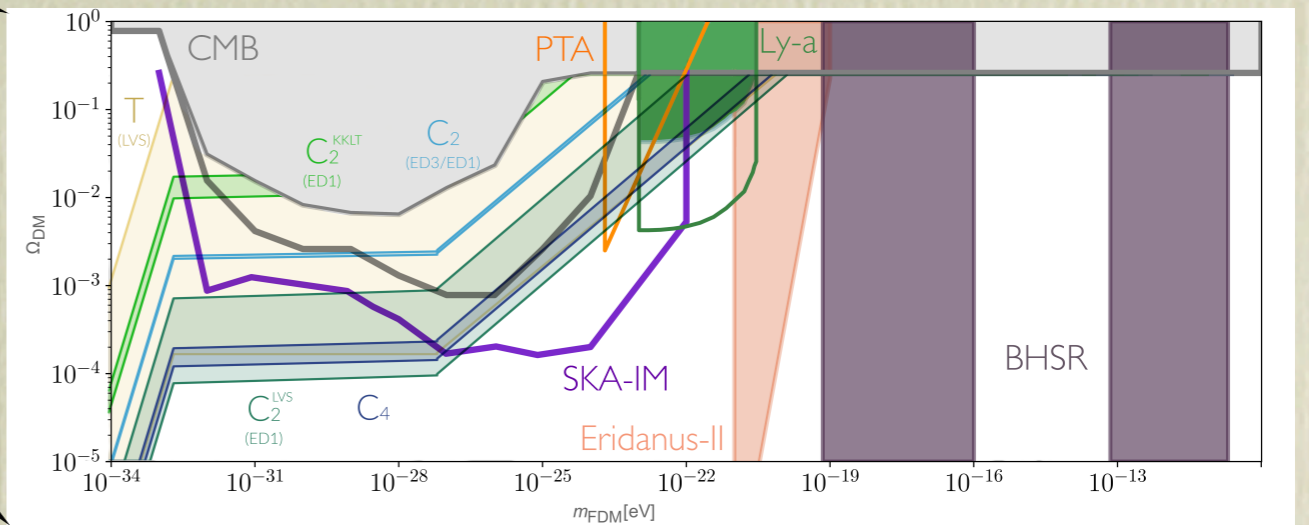
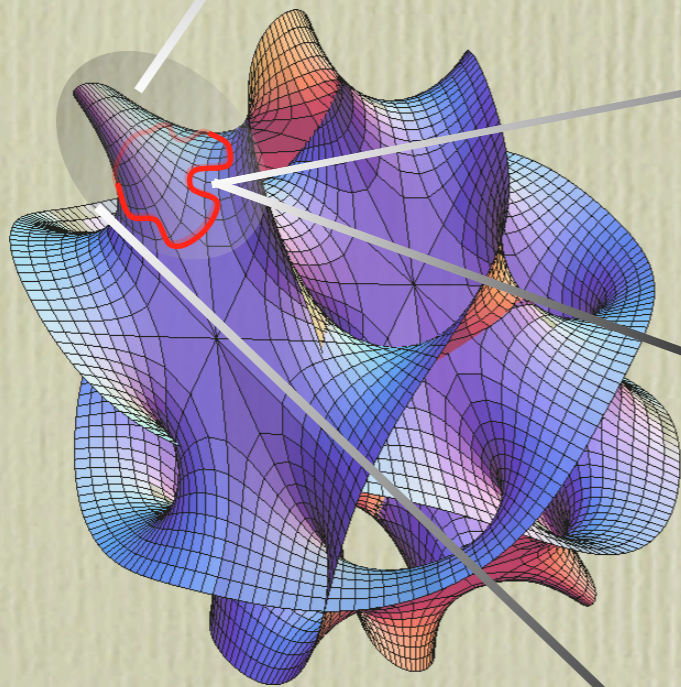
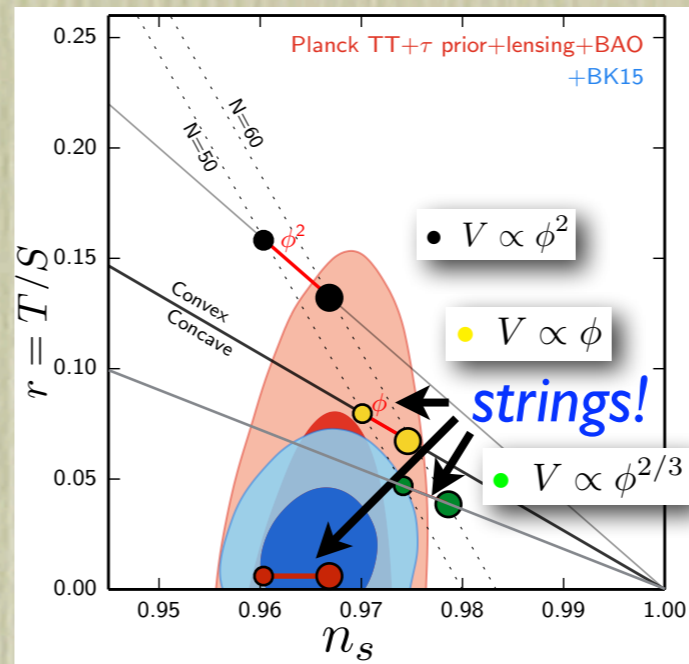


# The Top-Down Case for (String) Axions



Alexander Westphal  
(DESY)

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image credit: J. Leedom & Midjourney

Alexander Westphal  
(DESY)

# EFT of axion couplings

[Peccei & Quinn '77]  
[Weinberg '78]  
[Wilczek '78]

4d Lagrangian:  $\mathcal{L} = \frac{1}{2} f^2 \partial_\mu a \partial^\mu a + a \frac{g^2}{32\pi^2} \text{tr} G_{\mu\nu} \tilde{G}^{\mu\nu}$

non-perturbative effects:  
instantons of action S  
generate scalar potential

$$\mathcal{L} = \frac{1}{2} f^2 (\partial a)^2 - M_p^4 A e^{-S} \cos(a)$$

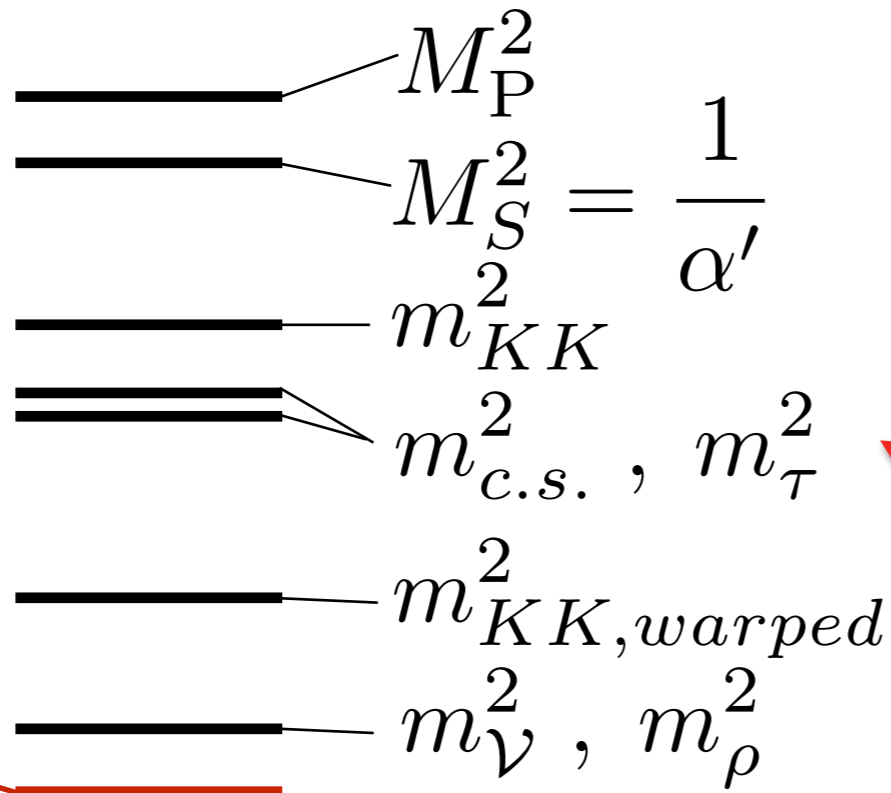
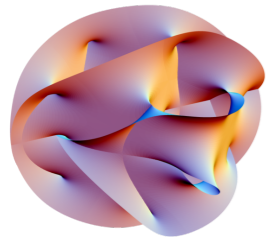
continuous shift symmetry  
broken to:

$$a \rightarrow a + 2\pi n, \quad n \in \mathbb{Z}$$

Axion mass:  $m_a^2 = M_p^4 \frac{A e^{-S}}{f^2}$

Axion-SM couplings:  $\mathcal{L} \supset -\frac{e^2}{32\pi^2} \frac{C_{i\gamma}}{f_{a_i}} a_i F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{C_{ie}}{2f_{a_i}} \bar{e} \gamma^\mu \gamma_5 e \partial_\mu a_i$

Chern-Simons (CS) term



$H^2$

$m_\phi^2$

**strings live in 10D**  
**-> extra dim.s:**  
**— many light scalars**  
**(moduli & axions)...**

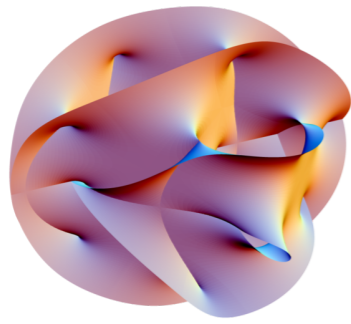
... unless **epicycles**:

- KK modes: ADD or RS
- clockwork ...
- little Higgs
- composite Higgs
- ...

desert ??

generically YES !!

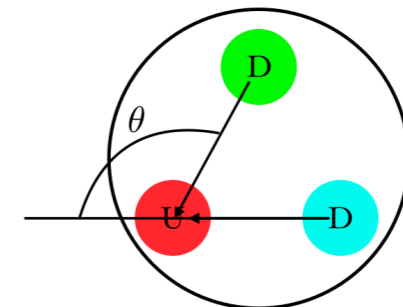




From Top-Down



And Bottom-Up



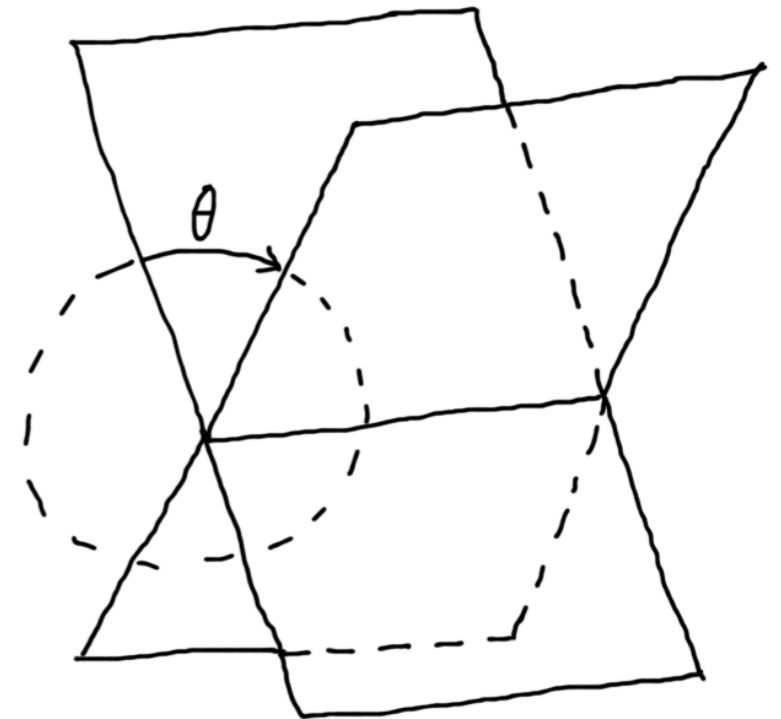
How can we find them?

# axions in string theory ...

[Svrcek & Witten '06]

- string theory:

- extra dimensions
- higher p-form gauge fields
- branes



- axions:

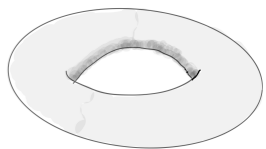
- Kaluza-Klein 0-modes of gauge fields
- angles  $\theta_a$  between branes
- ~~phases of open string matter fields~~

## type IIB closed string axions

$$\int_{\Sigma^{(p)}} C_p = A_0$$

$A_0$  = 0-form, i.e. an axion

$\Sigma^{(p)}$  = internal p-cycle of the Calabi-Yau



in particular, in type IIB we have :

$$\int_{\Sigma_i^{(4)}} C_4 = \theta_i$$

$$\int_{\Sigma_a^{(2)}} C_2 = c_a$$

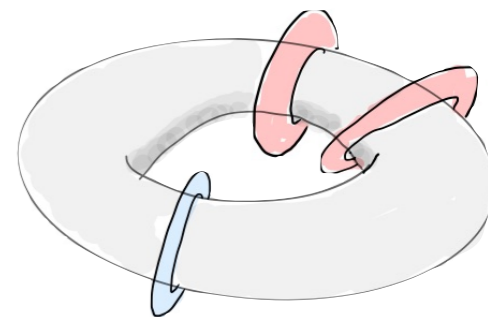
$$\int_{\Sigma_a^{(2)}} B_2 = b_a$$

we call them axions because:

- after compactification: continuous shift symmetry inherited from the 10d gauge invariance
- introduce branes: shift symmetry broken to a discrete one  
+ generate a potential (hence a mass) for the axions

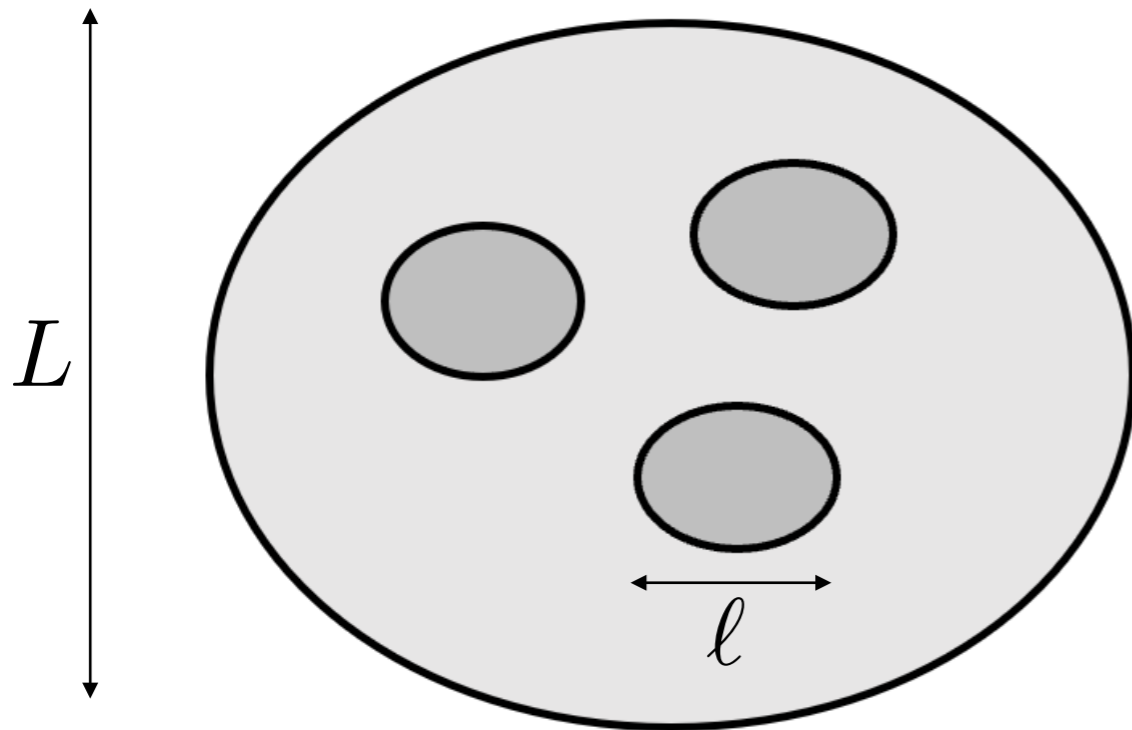


instantons of complex action  $S$



# string theory matching of axion EFT

[Banks, Dine, Fox & Gorbatov '03; Svrcek & Witten '06]



$$\mathcal{L}_{\text{kin.}} \sim M_{\text{P}}^2 \frac{\ell^{q'}}{L^p} \partial_\mu a \partial^\mu a, \quad q' < p, \quad \ell \lesssim L \quad \Rightarrow \quad f \lesssim M_{\text{P}}$$

$$\delta V \sim \text{Re}(e^{-S}), \quad \text{Re } S = \frac{\ell^q}{g_s^\#}$$

in most cases:

$$S f \lesssim M_{\text{P}}$$

(axionic WGC)

[Arkani-Hamed, Motl, Nicolis & Vafa '06]



- **consequence of string extra dimensions:**
  - many cycles —  $O(100)$
  - each cycle: a p-form 0-mode axion
- ★ string theory generically contains **many axions**
- ★ **decay constants** are **high**  
... **power-law** in extra-dim. size
- ★ **masses** distribute **exponentially wide**  
... **exponential** in extra-dim. size
- ★ couplings to SM: mostly no ...  
... exceptions highly model-dependent (e.g. kinetic mixing)

**a string theory axiverse !**

[Cicoli, Conlon & Quevedo '12]

[Higaki & Takahashi '12]

[Hebecker, Mangat,  
Rompineve & Witkowski '14]

⋮

[Cicoli, Sinha & Wiley Deal '22]

# • closed string axion pheno:

## - dark radiation

axion production from moduli decay in type IIB string models of moduli stabilization (LVS, KKLT ...)

$$\frac{K}{M_{\text{P}}^2} = -n_1 \ln(T_1 + \bar{T}_1) + \dots$$

type IIB on CY:

e.g. [Demirtas, Gendler, Long,  
McAllister & Moritz '21]

often has  $h^{1,1} > 1$

volume moduli &  $C_4$ -axions

$$\mathcal{L}_{kin.} = K_{i\bar{j}} \partial_\mu T^i \partial^\mu \bar{T}^{\bar{j}} \supset \frac{M_{\text{P}}^2}{4} \frac{n_1}{\tau_1^2} \partial_\mu \tau_1 \partial^\mu \tau_1 + \frac{M_{\text{P}}^2}{4} \frac{n_1}{\tau_1^2} \partial_\mu a_1 \partial^\mu a_1$$

mass & couplings from non-pert. effects:

### Cosmic Axion Background

[Conlon & Marsh '13]

[Dror, Murayama & Rodd '21]

$$V = \Lambda (\tau_i)^4 \cos(a)$$

heavy moduli decay

into relativistic axions “**dark radiation**”

non-linearity of couplings  
drives parametric resonance



- **closed string axion pheno:**

- **dark matter**

[Preskill, Wise & Wilczek '83]

[Abbott & Sikivie '83]

[Dine & Fischler '83]

high-scale decay constants:  $f > H$  (PQ broken) even during inflation

but exponentially light:  $m \ll H$  during inflation

population of non-relativistic axion matter density  $\rho$   
via misalignment:

random displacement of axion  $a$  during inflation from de Sitter vacuum fluctuations

every Hubble patch has different  $\rho$ , ours is selected anthropically

axion abundance:  
( $H < f$ )

$$\frac{\Omega_a h^2}{0.112} \simeq 2.2 \times \left( \frac{m_a}{10^{-22} \text{ eV}} \right)^{1/2} \left( \frac{f}{10^{17} \text{ GeV}} \right)^2 a_{\text{in}}^2$$

[Cicoli, Goodsell & Ringwald '12]

- closed string axion pheno:

- what dark matter?

- if  $m > 10^{-18}$  eV ... cold dark matter

- if  $10^{-25}$  eV  $< m < 10^{-19}$  eV ... fuzzy (or wave) dark matter

- [Hu, Barkana & Gruzinov '00]

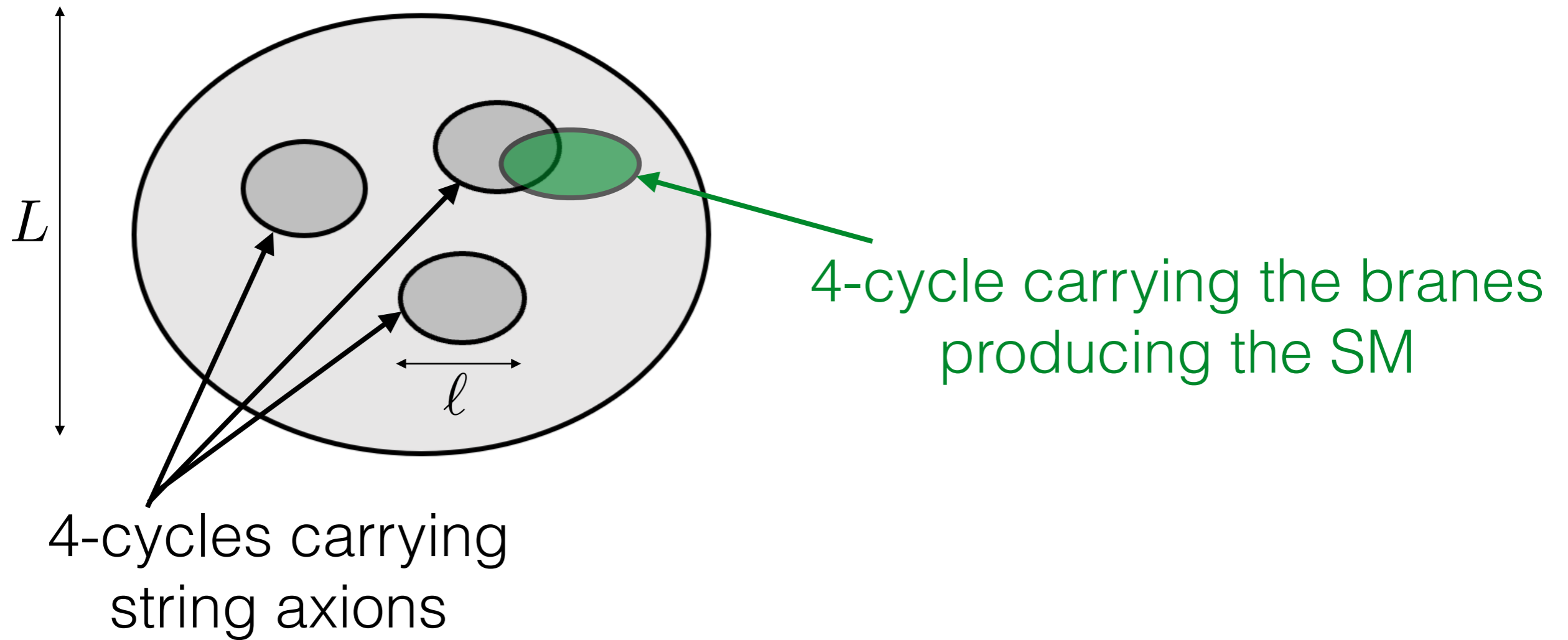
- [Hui, Ostriker, Tremaine & Witten '16]

- [Cicoli, Guidetti, Righi & AW '21]

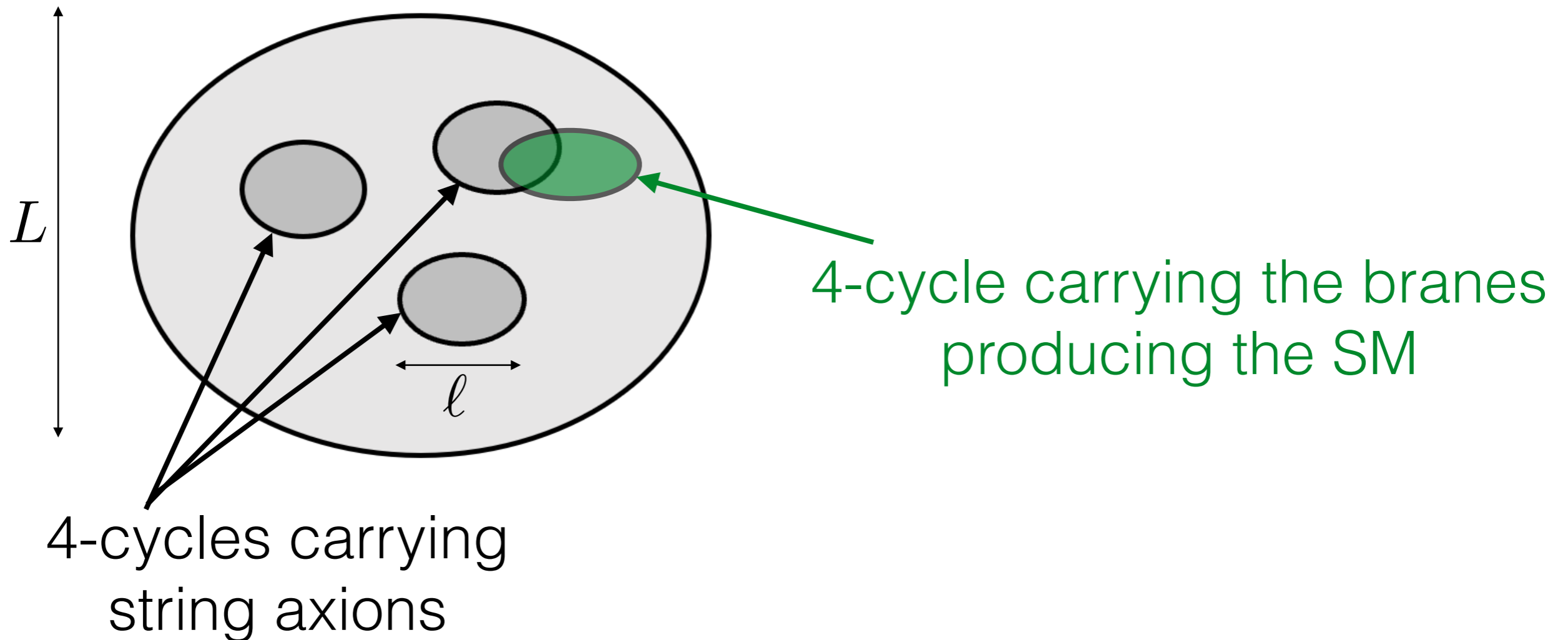
- other production mechanisms ...

- ... from topological defects, cosmic strings

# SM couplings of closed string axions ?



## SM couplings of closed string axions ?



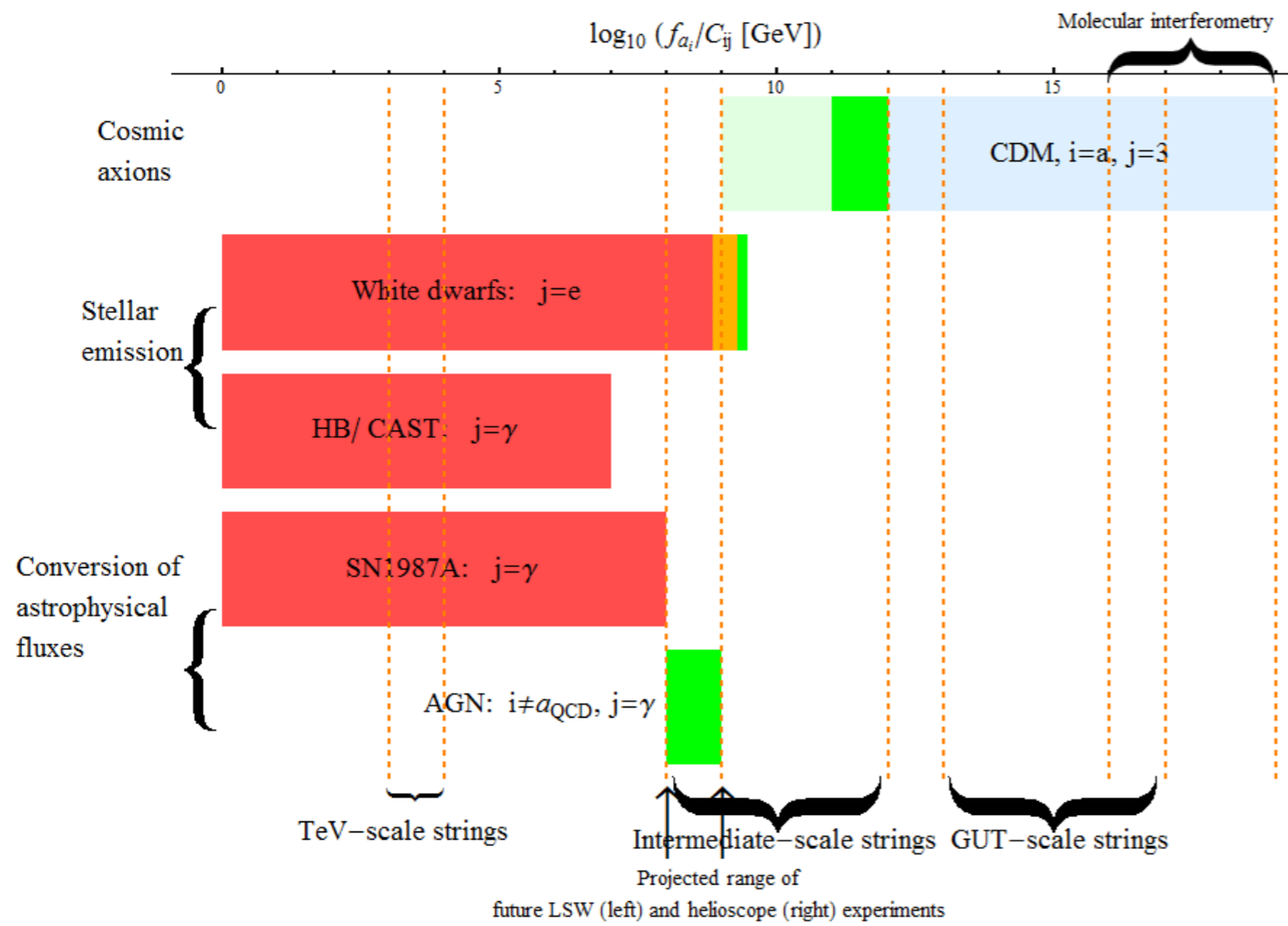
- **closed string axion pheno:**

[Gendler, Marsh, McAllister & Moritz '23]

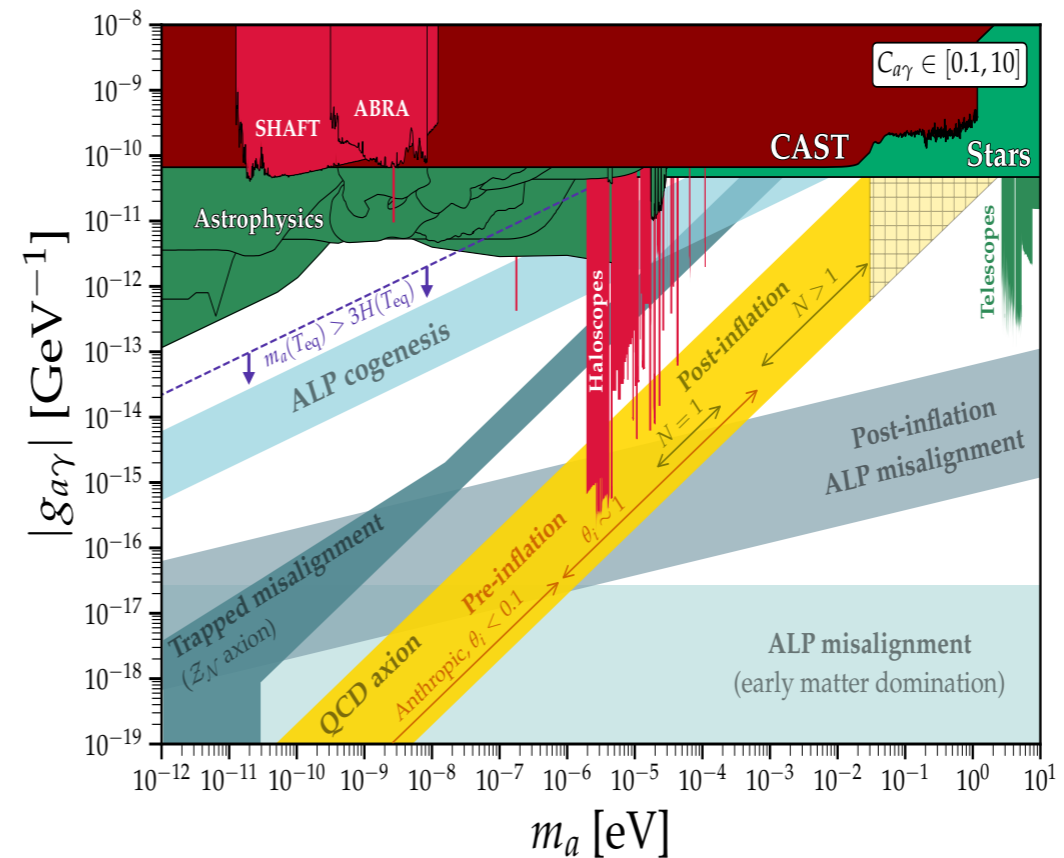
- handful of axions visible to SM, most are **DARK !!**
- SM-couplings depend on extra-dim. size & cycle intersection data
- kinetic mixing between U(1) gauge fields coupled to dark axions and SM U(1) gauge field often small

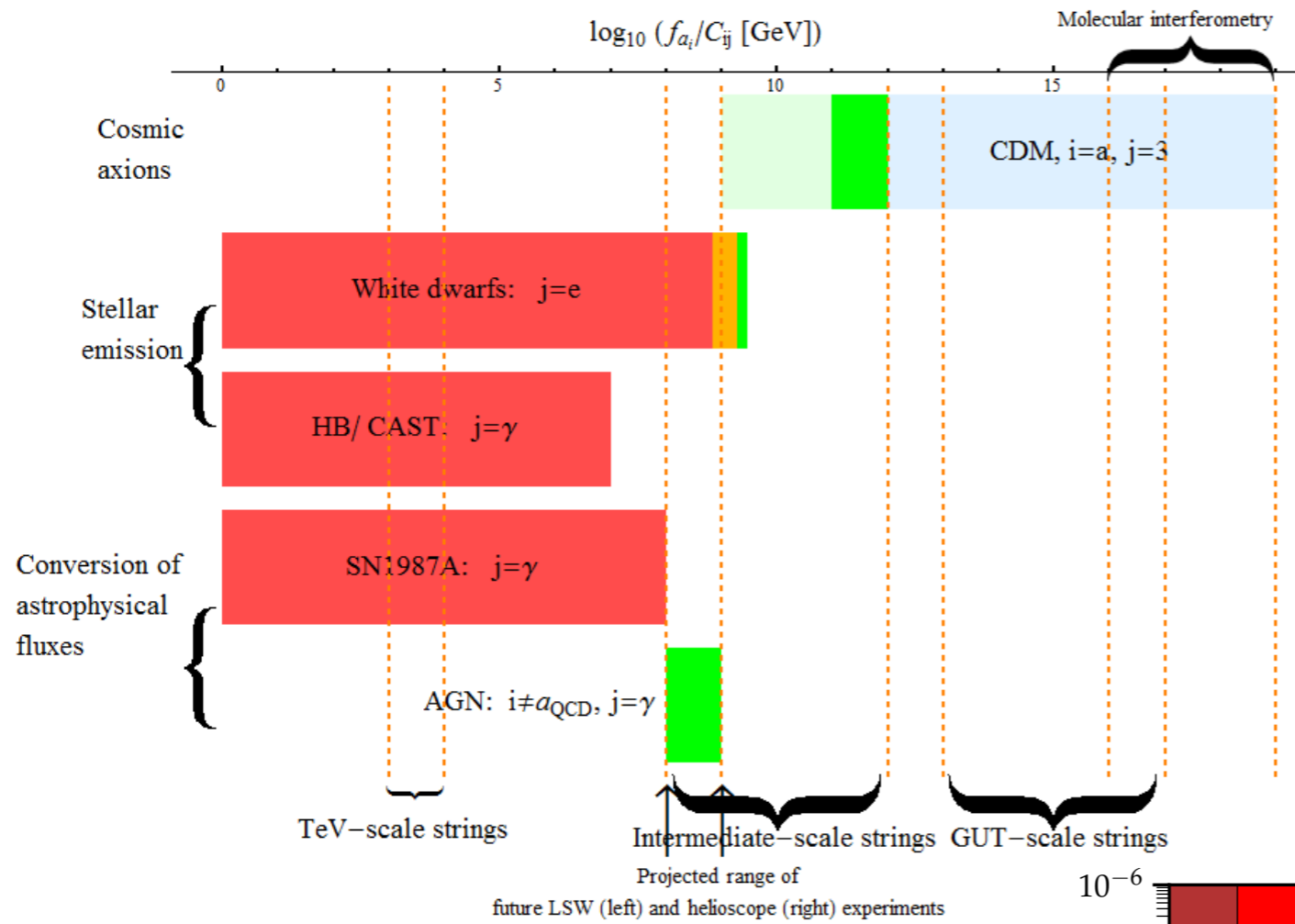
[Berg, Marsh, McAllister & Pajer '10]

[Hebecker, Jaeckel & Kuespert '23]

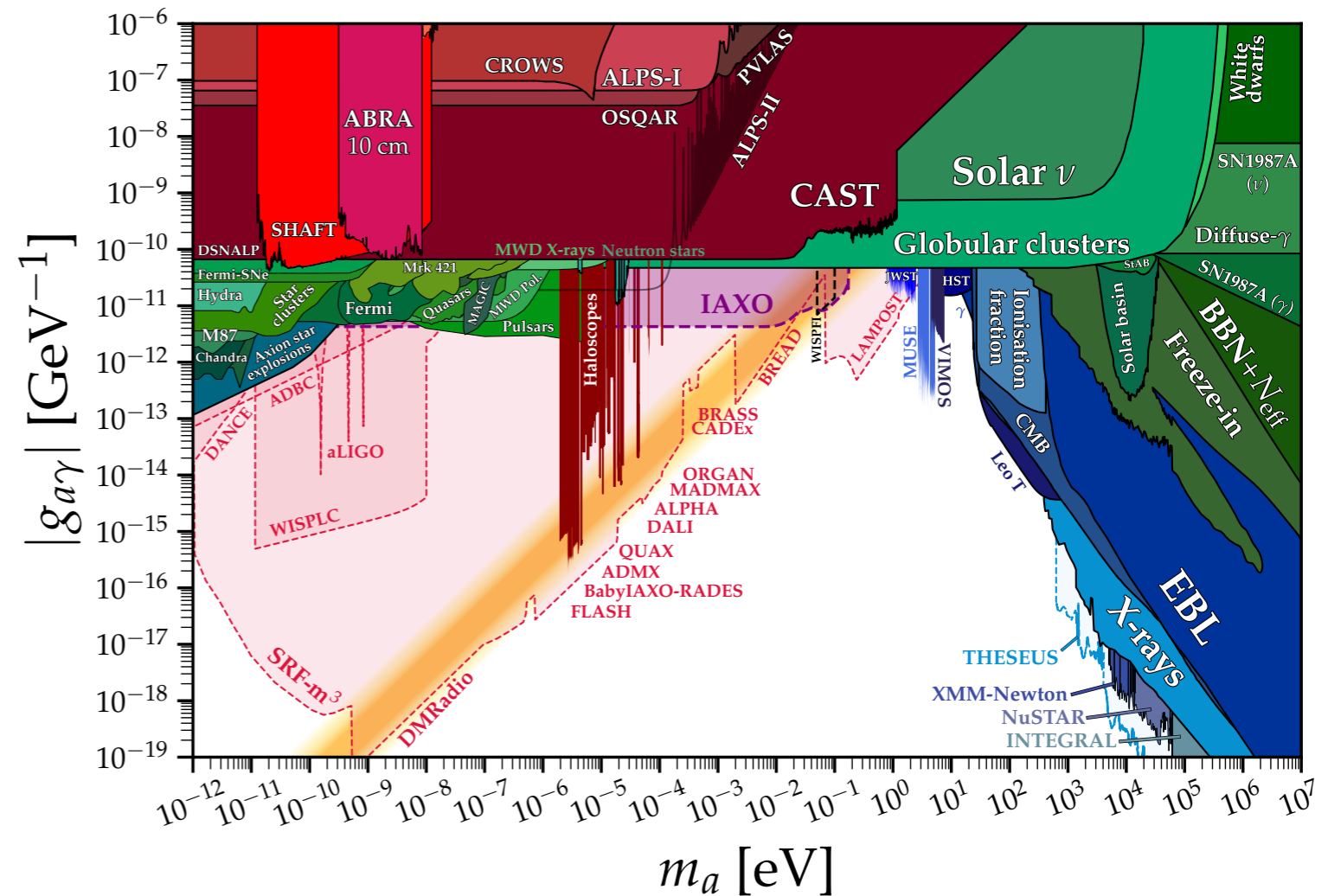


the 1...3 string axions with SM couplings





the 1...3 string axions with SM couplings





the many other string axions are dark ...

- couplings are either gravitational (Planck-suppressed):

- can be part of dark matter

- or dark radiation: visible e.g. in CMB as  $\Delta N_{\text{eff}}$  [Cicoli, Conlon & Quevedo '12]  
[Higaki & Takahashi '12]

- axion clouds around fast-spinning BHs deplete spin:  
detection via BH superradiance

[Arvanitaki, Dimopoulos, Dubovsky, Kaloper & March-Russell '09]

- and/or to dark sectors, e.g. dark U(1) gauge fields:

- spectator role during inflation [Biagetti, Dimastrogiovanni, Fasiello & Peloso '14]  
[Dimastrogiovanni, Fasiello & Fujita '16]  
[Obata & Soda '17]

- slow-rolling axion coupled to dark U(1) during inflation:  
production of dark gauge field & induced peaked GWs

[Anber & Sorbo '09, Sorbo '11] [Adshead & Wyman '12]

... [Namba, Peloso, Shiraishi, Sorbo & Unal '15] [Dimastrogiovanni, Fasiello & Fujita '16], [Domcke, Pieroni & Binetruy '16] ...

- for string axions quite generic & constrains axiverse

[Dimastrogiovanni, Fasiello, Leedom, Putti & AW '23]

# summary

- axion pheno to large part determined by couplings in kinetic term and NP scalar potential + matter & gauge field couplings
- these couplings are top-down determined by compactification data — e.g intersection #s, fluxes, or topological data (e.g. for thractions)
- axion-matter couplings depend on axion type and SM realization (7-branes on 4-cycle, 3-branes at CY singularity, thractions ?)

## need both:

- **explicit** string model **constructions** to study structure & parameter range of axion couplings
- **scans** over large sets of string vacua **to get** number frequency **distribution of axion EFT parameters**