

Axion Beyond Discoveries:

Measuring Axion Couplings

A very brief overview

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AxionFest
DESY, 28–31 January 2024

Wish List

→ *Getting the couplings*



→ *Getting info on the Universe*



→ *Getting info on the model*



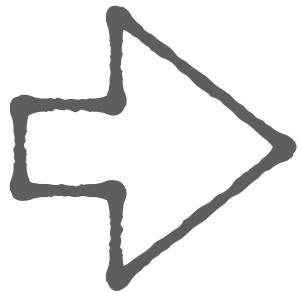
Measuring the Couplings

Best option: **Axion-Photon Coupling**



ALPS II

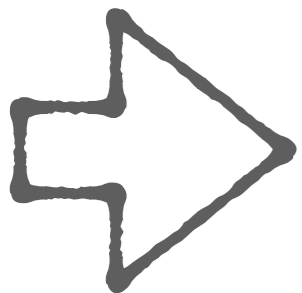
$$\text{Observed Flux} \sim |g_{a\gamma\gamma}|^4$$



Heliscopes

$$\text{Observed Flux} \sim \left(1 + C \frac{|g_{ae}|^2}{|g_{a\gamma}|^2} \right) |g_{a\gamma}|^4$$

Calculable



Haloscopes

Unknown

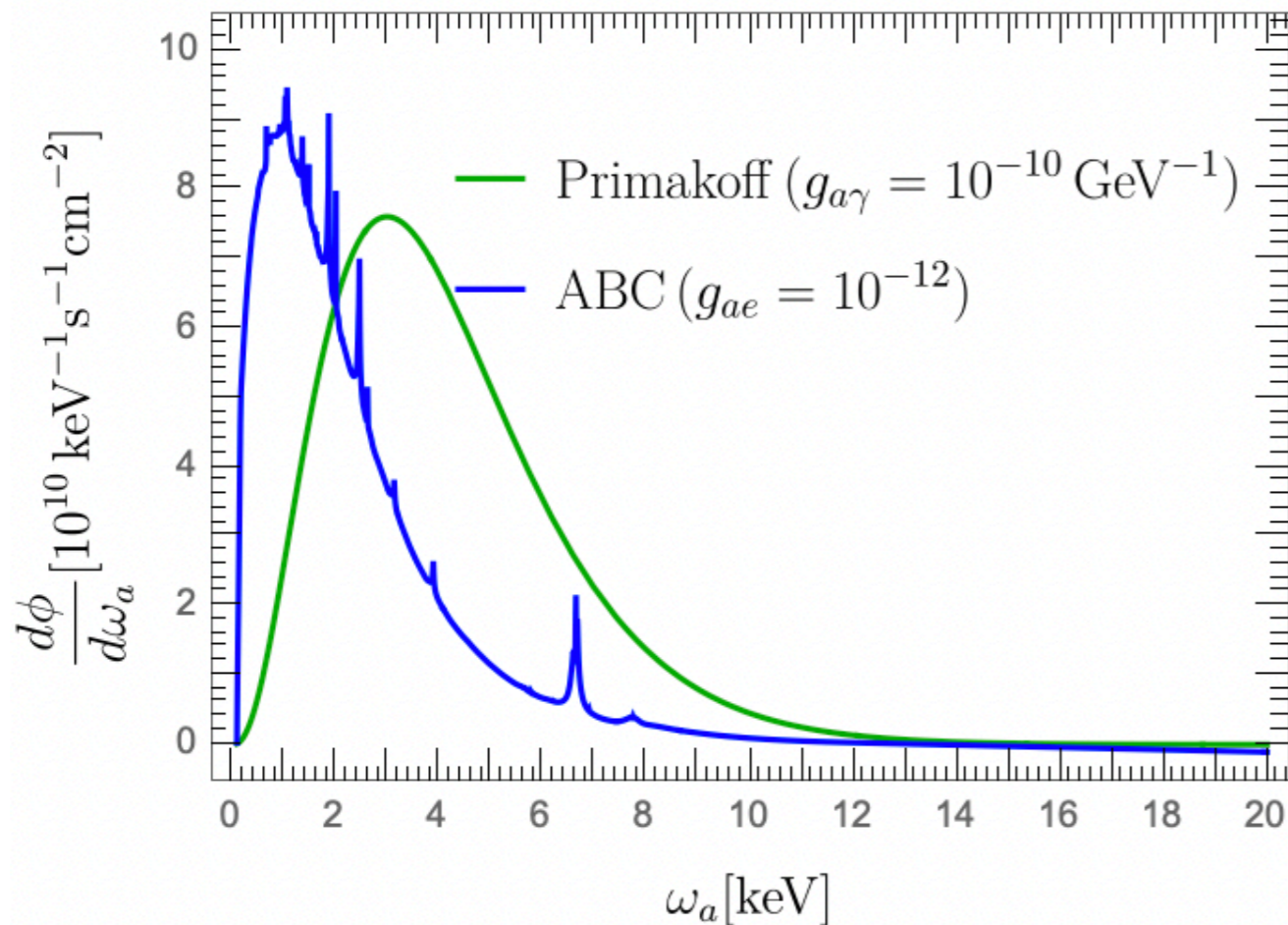
$$\text{Observed Flux} \sim \eta_{\text{DM}} |g_{a\gamma\gamma}|^2$$

Model dependent

+ Infer the mass in case of resonant detection

Solar Axions

Best option (for now) → **BabylAXO**



→ [J. Redondo, JCAP 1312 \(2013\)](#)

→ [S. Hoof, J. Jaeckel, L. J. Thormaehlen JCAP 09 \(2021\) 006](#)

→ find $g_{ae}/g_{a\gamma}$ from spectra?

Example: Threshold at 0.3 keV

$$\frac{\Phi_{(0.3-1)\text{keV}}}{\Phi_{\text{tot}}} \simeq 2.4 \% \quad (\text{KSVZ})$$

where $\Phi_{\text{tot}} = \Phi_{(0.3-10)\text{keV}}$.

DFSZ with typical couplings

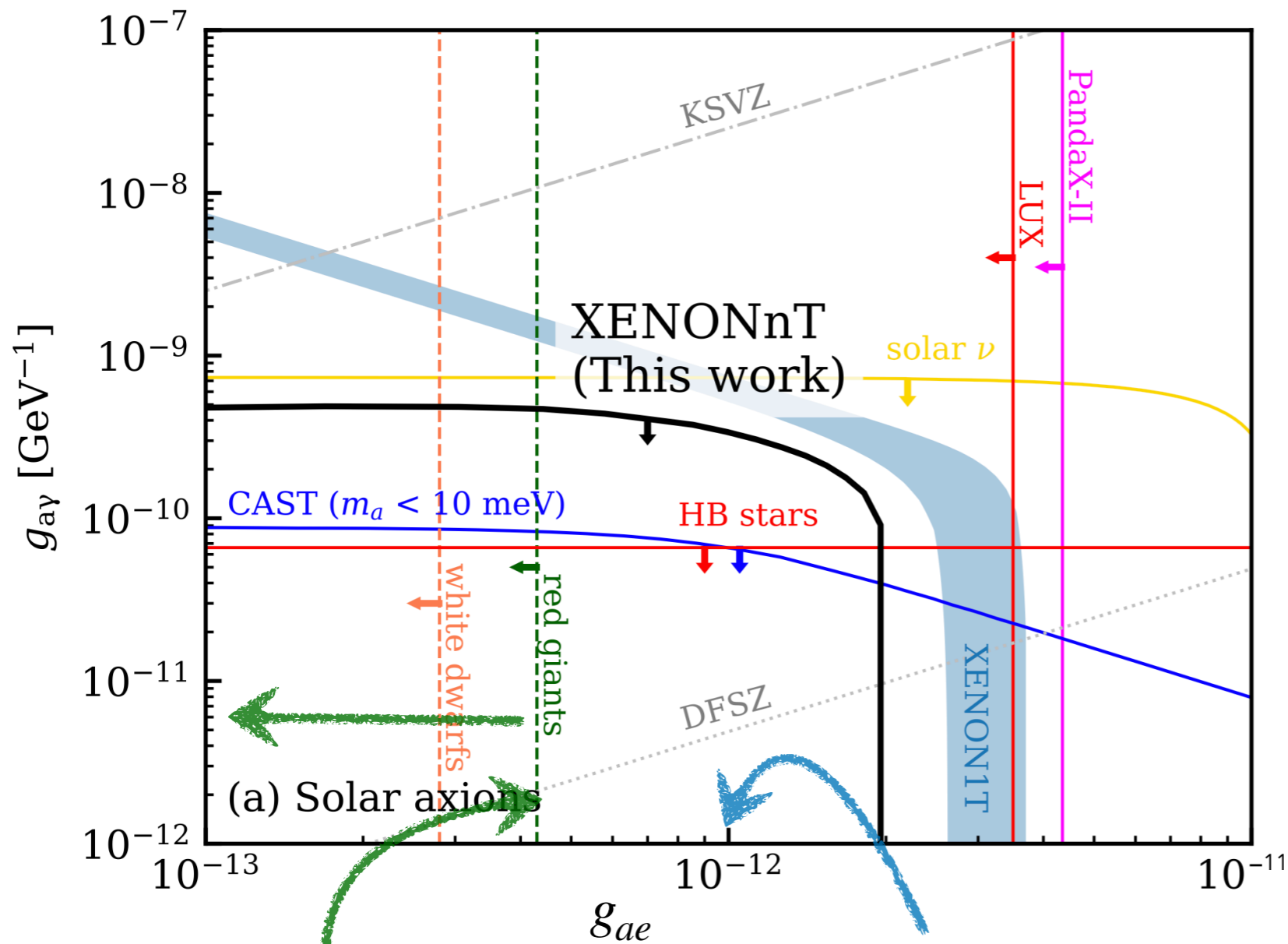
$$g_{ae}/g_{a\gamma} = 5 \times 10^{-2} \text{ GeV},$$

$$\frac{\Phi_{(0.3-1)\text{keV}}}{\Phi_{\text{tot}}} \simeq 22 \% \quad (\text{DFSZ})$$

Using other bins (e.g., 1-2 keV) is much less efficient.

No significant improvement in going down to 0.1 keV.

Direct detection of g_{ae}



Inferring g_{ae} directly not likely in near future because of strong astro bounds

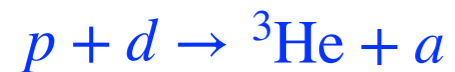
*E. Aprile et al.,
[Phys.Rev.Lett. 129 \(2022\)](#)*

RGB bound has been updated

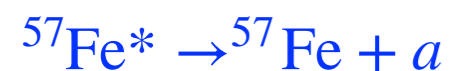
DARWIN
(Expected potential)

Axion Nucleon Couplings

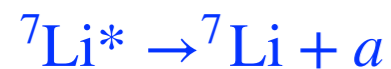
The Sun ...



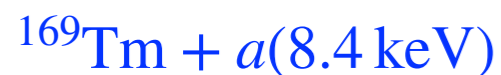
- Searched by CAST *JCAP* 03 (2010)
- Borexino *Phys.Rev.D* 85 (2012)
- and using previous SNO data *Phys.Rev.Lett.* 126 (2021)
- Recent analysis of the JUNO sensitivity shows potential to search in unexplored regions G. Lucente, N. Nath, F. Capozzi, MG, A. Mirizzi, *Phys.Rev.D* 106 (2022) 12
- Maybe accessible to IAXO (work in progress)



- Searched by CAST *JCAP* 12 (2009) + BabyIAXO *Eur.Phys.J.C* 82 (2022)
- New dedicated project under commissioning → [ISAI \(Investigating Solar Axion by Iron-57\)](#),



- Searched by Borexino *Eur.Phys.J.C* 54 (2008)
- CAST *JCAP* 03 (2010)



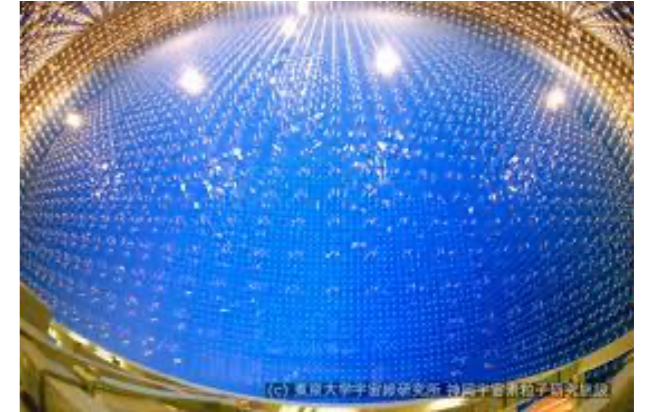
- Thulium garnet crystal as a bolometric detector, Derbin et al., (2023) [JETP Letters, Volume 118, Issue 3, p.160-164](#)

... + Supernovae

Direct Detection

→ Cherenkov

- A. Lella et al., [arXiv:2306.01048](#);
- Vonk, Guo, Meißner, [Phys.Rev.D 105 \(2022\)](#)
- Li, Hu, Guo, Meißner, [2312.02564](#)
- P. Carena et al., [arXiv:2306.17055](#)

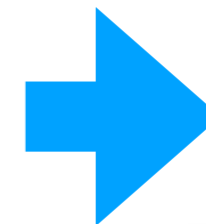


→ Colliders

- S. Asai, Y. Kanazawa, T. Moroi, T. Sichanugrist [Phys.Lett.B 829 \(2022\)](#)

→ Heliscopes

- Ge, Hamaguchi, Ichimura, Ishidoshiro, Kanazawa, [JCAP 11 \(2020\)](#);



new proposal by [Juan Anton Garcia Pascual](#)
(UNIZAR/CAPA)

Indirect detection

Through photon oscillations in B_{ext}

- F. Calore et al. e-Print: [2306.03925](#)
- A. Lella et al. In preparation
- Meyer et al. [Phys.Rev.Lett. 118 \(2017\)](#)



Post-Discovery: Axion Telescopes

Detecting stellar axions would allow to understand a lot about stars.

- Solar magnetic field

C. A. J. O'Hare, A. Caputo, A. J. Millar, E. Vitagliano [Phys.Rev.D 102 \(2020\) 4](#)

- Solar temperature profile

S. Hoof, J. Jaeckel, L. J. Thormaehlen, [arXiv:2306.00077](#)

- Solar chemical composition

J. Jaeckel, L. J. Thormaehlen, [Phys.Rev.D 100 \(2019\) 12](#)

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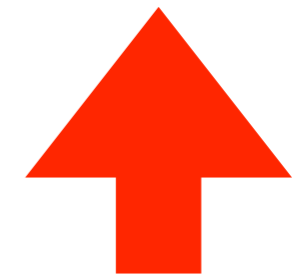


See → talk by
Sebastian Hoof

Post-Discovery: Axion Telescopes

Model	Phase	t_{cc} [yr]	$\log_{10} \frac{L_{eff}}{L_{\odot}}$	$\log_{10} \frac{T_{eff}}{K}$	Primakoff			Bremsstrahlung			Compton		
					C^P	E_0^P [keV]	β^P	C^B	E_0^B [keV]	β^B	C^C	E_0^C [keV]	β^C
0	He burning	155000	4.90	3.572	1.36	50	1.95	1.3E-3	35.26	1.16	1.39	77.86	3.15
1	before C burning	23000	5.06	3.552	4.0	80	2.0	2.3E-2	56.57	1.16	8.55	125.8	3.12
2	before C burning	13000	5.06	3.552	5.2	99	2.0	6.4E-2	70.77	1.09	17.39	156.9	3.09
3	before C burning	10000	5.09	3.549	5.7	110	2.0	8.9E-2	76.65	1.08	22.49	169.2	3.09
4	before C burning	6900	5.12	3.546	6.5	120	2.0	0.136	85.15	1.06	31.81	186.4	3.09
5	in C burning	3700	5.14	3.544	7.9	130	2.0	0.249	97.44	1.04	50.62	210.4	3.11
6	in C burning	730	5.16	3.542	12	170	2.0	0.827	129.17	1.02	138.6	269.1	3.17
7	in C burning	480	5.16	3.542	13	180	2.0	0.789	134.54	1.02	153.2	279.9	3.15
8	in C burning	110	5.16	3.542	16	210	2.0	1.79	151.46	1.02	252.7	316.8	3.17
9	in C burning	34	5.16	3.542	21	240	2.0	2.82	181.74	1.00	447.5	363.3	3.22
10	between C/Ne burning	7.2	5.16	3.542	28	280	2.0	3.77	207.84	0.99	729.2	415.7	3.23
11	in Ne burning	3.6	5.16	3.542	26	320	1.8	3.86	224.45	0.98	856.4	481.2	3.11

$$\frac{d\dot{N}_a}{dE} = \frac{10^{42}}{\text{keVs}} \left[C^P g_{11}^2 \left(\frac{E}{E_0^P} \right)^{\beta^P} e^{-(\beta^P + 1)E/E_0^P} + (P \rightarrow B, C; g_{11} \rightarrow g_{13}) \right]$$



Axion Flux very sensitive to evolutionary stage

Post-Discovery: Axion Telescopes

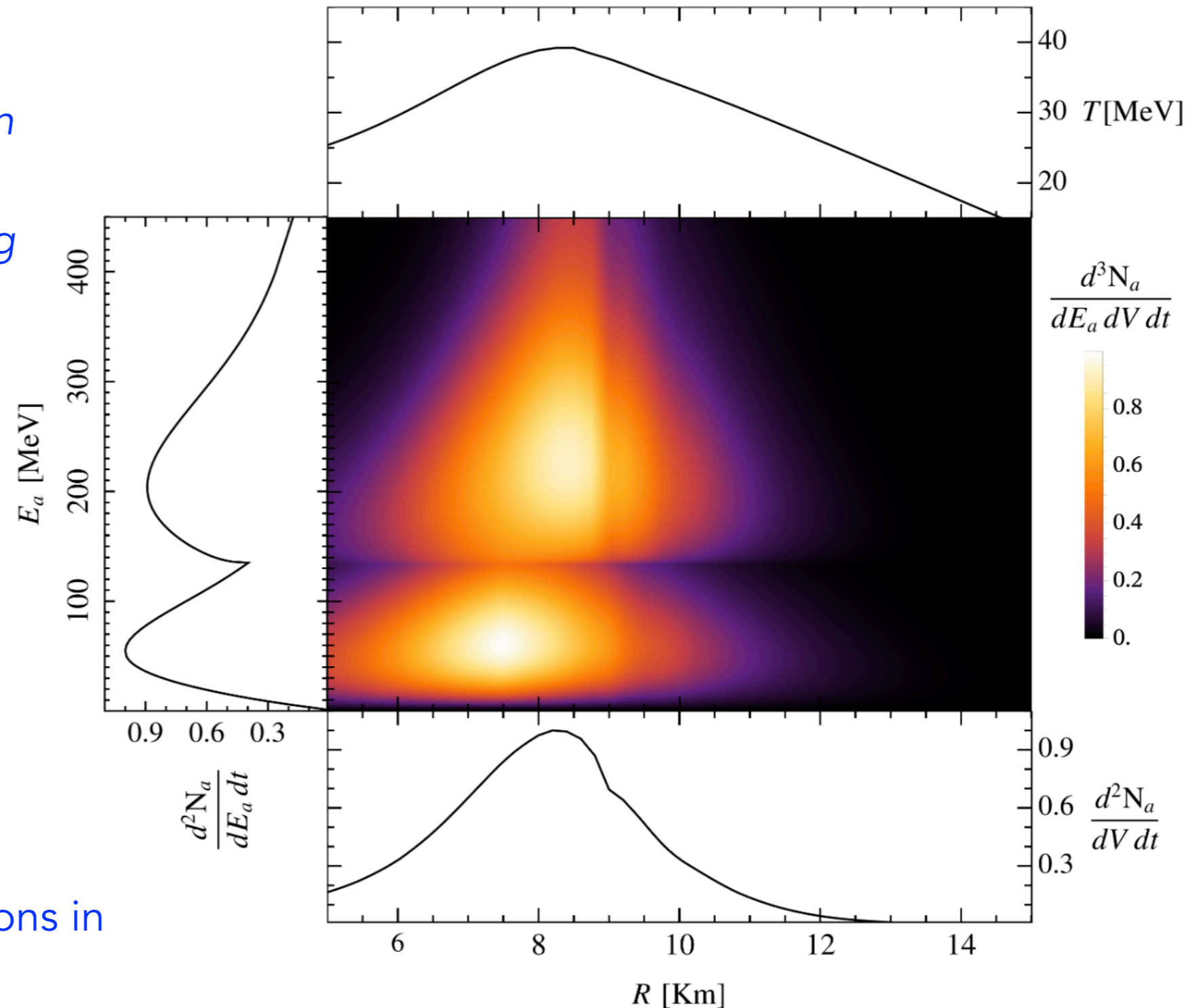
Looking into the SN Core

Contributions from *pion* processes and from nucleon bremsstrahlung

- Maybe accessible to IAXO (work in progress)

Can we distinguish the fluxes?

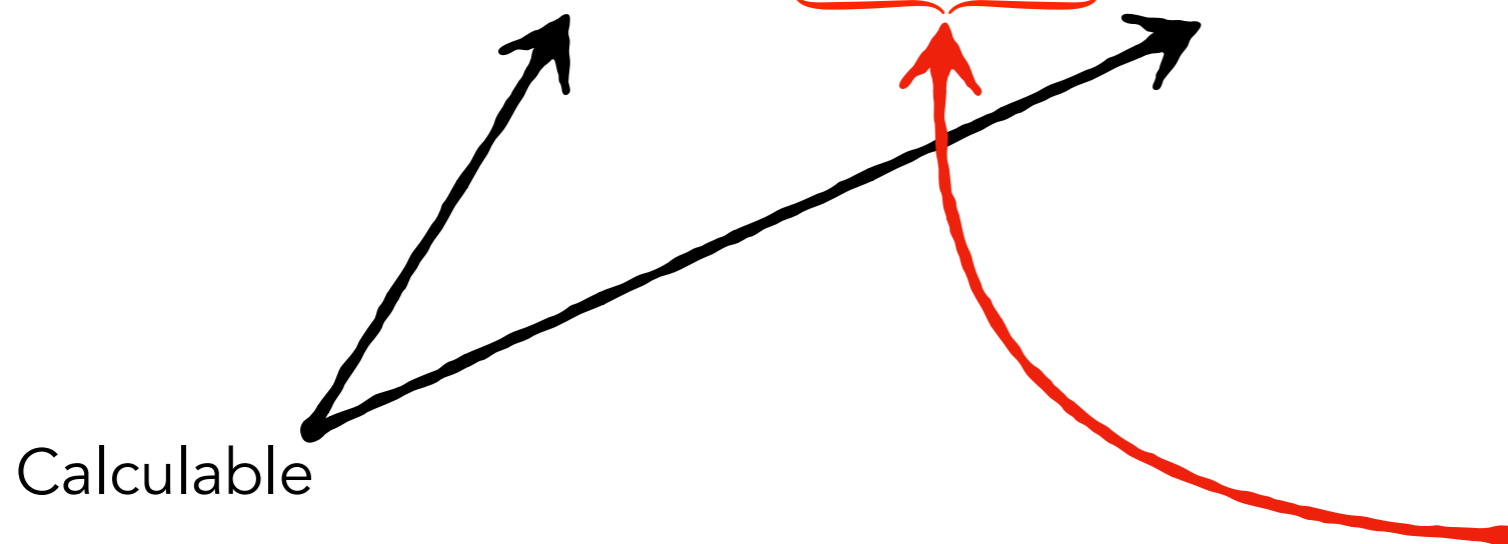
→ reveal the role of pions in SN



Inferring the UV completion?

The leading contribution to the running axion couplings arises from top loop diagrams induced by the axion-top coupling C_t

$$C_\Psi(2\text{GeV}) \simeq C_\Psi(f_a) + \underbrace{r_\Psi^t(m_{\text{BSM}})}_{\text{Depends on unknown high energy physics}} C_t(f_a) \quad (\Psi = u, d, e)$$



Depends on unknown high energy physics

\Rightarrow *unavoidable additional uncertainty*

\rightarrow Choi et al. [JHEP 08 \(2021\) 058](#)

\rightarrow Di Luzio et al. [Phys.Rev.D 108 \(2023\)](#)

UV Corrections to Couplings

$$C_{\Psi}(2\text{GeV}) \simeq C_{\Psi}(f_a) + r_{\Psi}^t(m_{\text{BSM}}) C_t(f_a)$$

Analytical Approximations

$$r_3^t(m_{\text{BSM}}) = r_u^t - r_d^t \simeq -0.54 \ln(\sqrt{x} - 0.52)$$

$$r_0^t(m_{\text{BSM}}) = r_u^t + r_d^t \simeq 3.8 \times 10^{-4} \ln^2(x - 1.25) \approx 0$$

$$r_e^t(m_{\text{BSM}}) \simeq -\frac{1}{2} r_3^t$$

$$\left. \begin{array}{l} r_3^t(m_{\text{BSM}}) = r_u^t - r_d^t \simeq -0.54 \ln(\sqrt{x} - 0.52) \\ r_0^t(m_{\text{BSM}}) = r_u^t + r_d^t \simeq 3.8 \times 10^{-4} \ln^2(x - 1.25) \approx 0 \\ r_e^t(m_{\text{BSM}}) \simeq -\frac{1}{2} r_3^t \end{array} \right\} \text{with } x = \log_{10} \left(\frac{m_{\text{BSM}}}{\text{GeV}} \right)$$

→ Di Luzio et al. [Phys.Rev.D 108 \(2023\)](#)

Less Uncertainty in Nuclear Couplings?

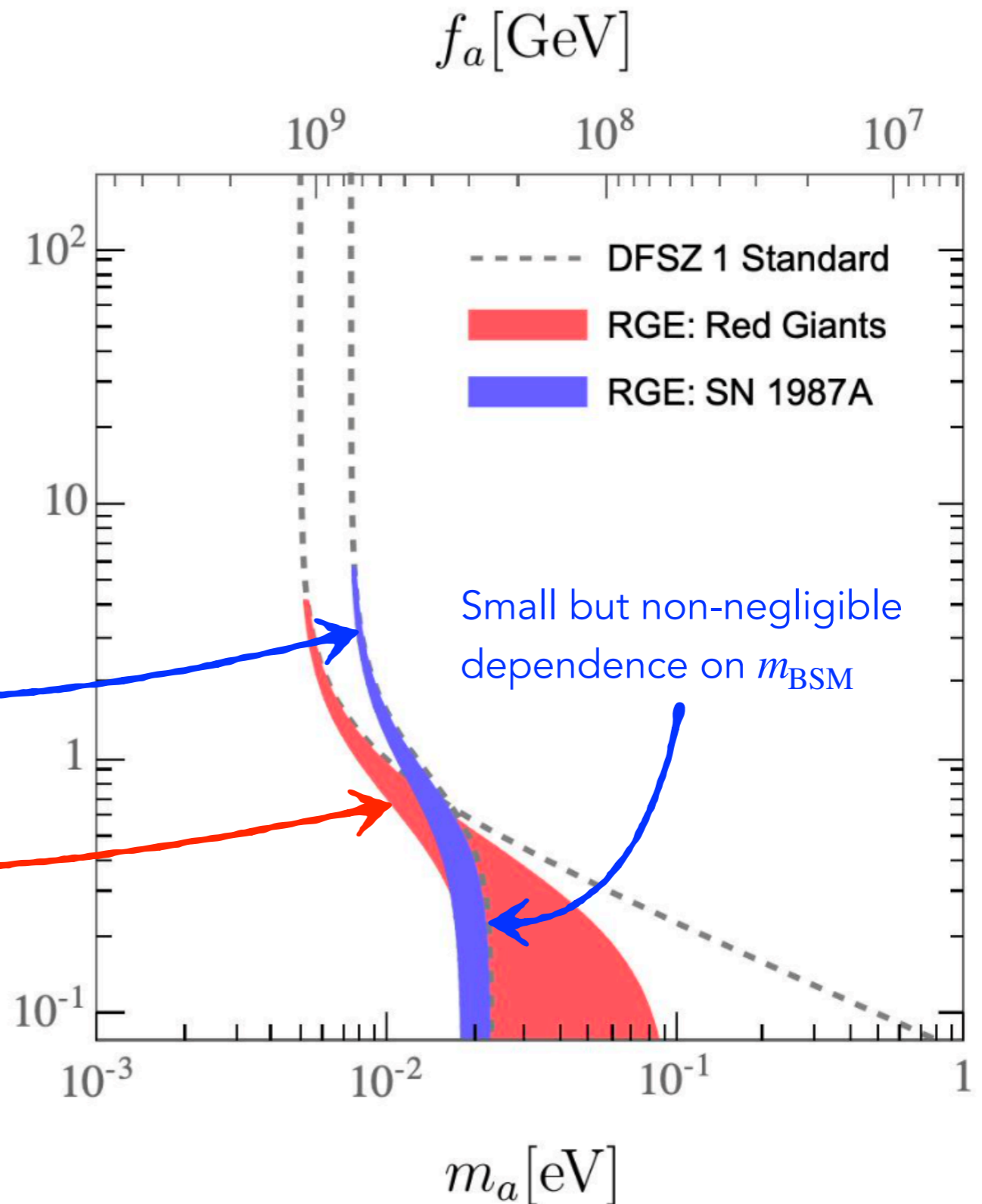
$$C_{\text{SN}} \simeq 1.4 (C_0^2 + 0.11 C_0 C_3 + 1.3 C_3^2)^{1/2}$$

This translates in the bound

$$m_a \leq \frac{\bar{m}}{C_{\text{SN}}}, \quad \text{with} \quad \bar{m} \sim 5 \text{ meV}$$

$$|C_e| \leq 1.65 \times 10^{-3} \left(\frac{m_a}{\text{eV}} \right)^{-1}$$

The C_3 dependence comes mostly (not exclusively) from pion interactions.

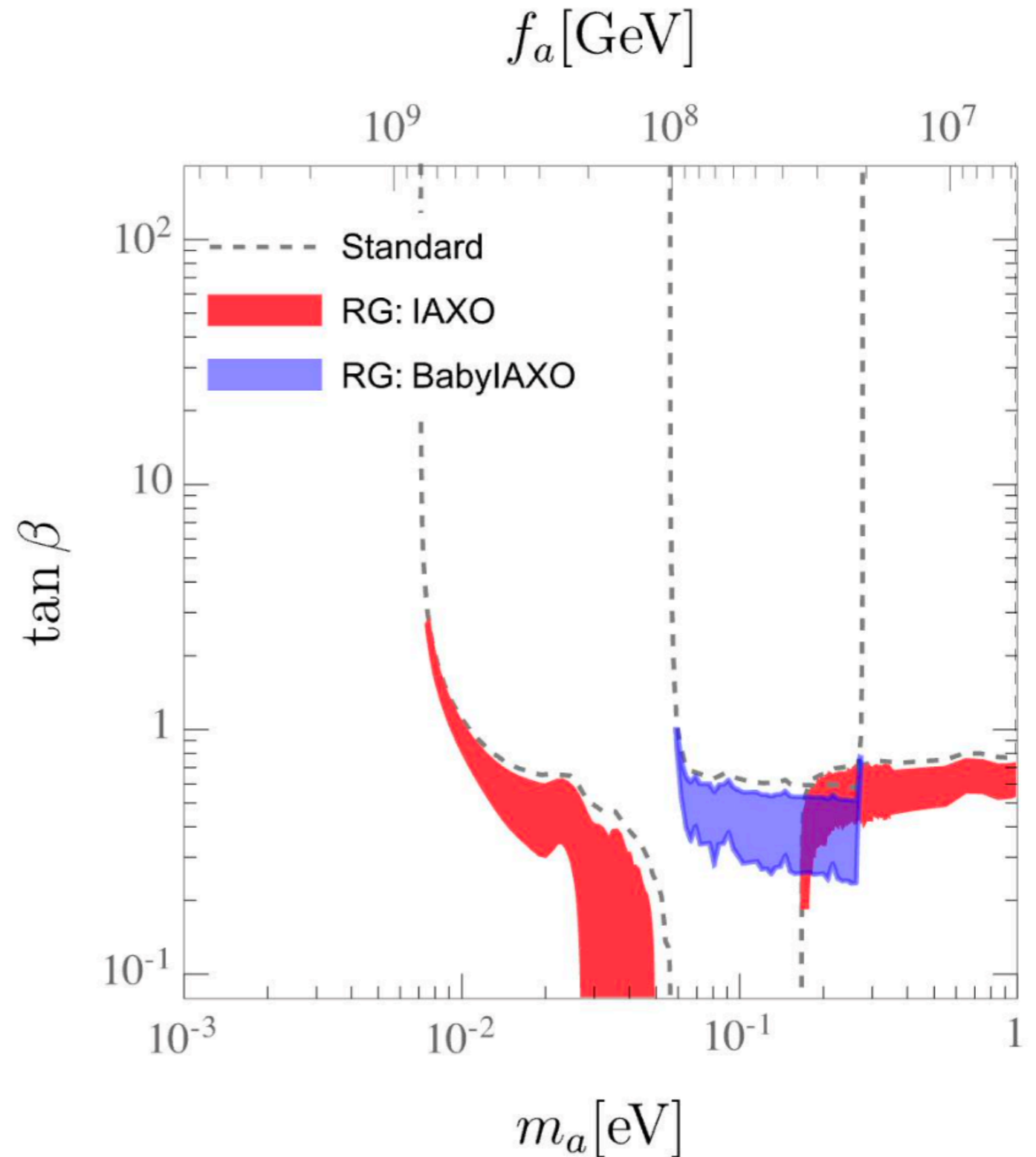


Detection Perspectives for DFSZ axions

The solar flux of DFSZ axions has always a g_{ae} component.

The **IAXO potential** for DFSZ parameter space is **higher** than naively expected.

However, **post-discovery uncertainties**



Conclusions

- Realistic options to find $g_{a\gamma}$ and perhaps $g_{a\gamma}/g_{ae}$.
- Several options also for some effective nuclear coupling g_{aN}^{eff}
- The door to the UV may be the **axion-photon coupling or isoscalar nuclear couplings**.
- **After we find the axions**, we can use them to study the sun and other stars. → That will be truly fun!! (See Sebastian talk and ask us questions)