ALPHA PLASMA EXPERIMENT

See 2210.00017

Pierluca Carenza, Stockholm University - OKC

AXION-PHOTON CONVERSION

$$\mathcal{L} = -\frac{g_{a\gamma}}{4} F_{\mu\nu} \tilde{F}^{\mu\nu} a = g_{a\gamma} \vec{E} \cdot \vec{B} a,$$

- Important for detection
- Generic feature of ALPs
- $P_{a\gamma}$ drops with the ALP energy

THE CONVERSION PROBABILITY

$$\mathcal{A}_{\gamma_i \to a} = -i \frac{g_{a\gamma}}{2} \int_{-\infty}^{\infty} d\rho \, \mathbf{u}_{T,i} \cdot \mathbf{B}(\rho \mathbf{u}) e^{-i\Phi(\rho)}$$
$$\Phi(\rho) = \int_{0}^{\rho} d\rho' \, \frac{m_{a}^{2} - \omega_{\mathrm{pl}}^{2}}{2\omega}$$

- This is the conversion probability at the lowest order
- Magnetic field and plasma frequency impact the conversion probabilty

ENHANCING THE CONVERSION PROBABILITY

- Oscillating B can resonate with the ALP mass and increase the conversion probability
- Breaking translation invariance makes the conversion possible

$$\mathcal{A}_{\gamma_i \to a} = -i \frac{g_{a\gamma}}{2} \int_{-\infty}^{\infty} d\rho \, B_T \cos(\Delta_0 \rho) e^{-i \frac{m_a^2}{2\omega} \rho}$$

$$\Phi(\rho) = \frac{m_a^2}{2\omega}\rho$$

ENHANCING THE CONVERSION PROBABILITY

- Match plasma frequency and ALP mass
- We need a tunable plasma frequency

$$\mathcal{A}_{\gamma_i \to a} = -i \frac{g_{a\gamma}}{2} \int_{-\infty}^{\infty} d\rho \, B_T e^{-i \frac{m_a^2 - \omega_{\rm pl}^2}{2\omega} \rho}$$

WIRE ARRAY METAMATERIAL

See 2203.13945



Plasma frequency given in terms of

$$\omega_{\mathrm{pl}}^2 \sim rac{1}{a^2 \ln\left(rac{a}{r}
ight)}$$

Where a is the spacing and r the tickness of the wires

THE PROTOTYPE

- Anisotropic material, vacuum for modes propagating along x and y
- Boundary conditions: closed by a conductor

THE PROTOTYPE



$$\begin{aligned} \hat{\mathbf{z}} \cdot \nabla_t \times \mathbf{B}_t &= -i\omega\epsilon_z E_z - i\omega g_{a\gamma} aB_e \\ \hat{\mathbf{z}} \cdot \nabla_t \times \mathbf{E}_t &= i\omega B_z , \end{aligned}$$
$$\hat{\mathbf{z}} \times \frac{\partial \mathbf{B}_t}{\partial z} + \nabla_t B_z \times \hat{\mathbf{z}} &= -i\omega \mathbf{E}_t , \\ \hat{\mathbf{z}} \times \frac{\partial \mathbf{E}_t}{\partial z} + \nabla_t E_z \times \hat{\mathbf{z}} &= i\omega \mathbf{B}_t . \end{aligned}$$



CAVITY TUNING

- Move two parallel sets of wires to change the spacing
- Simple and efficient to span a factor of a few GHz in frequency

Small E proportional to g_ag

$$E_z = -\frac{ag_{a\gamma}B_{\rm e}}{\epsilon_z} + \frac{ag_{a\gamma}B_{\rm e}}{\epsilon_z} \frac{J_0(\sqrt{\epsilon_z}r\omega)}{J_0(\sqrt{\epsilon_z}R\omega)},$$

$$\epsilon = 1 - \frac{\omega_p^2}{\omega^2 + i\omega\Gamma_p} \simeq 1 - \frac{\omega_p^2}{\omega^2} + i\frac{\Gamma_p\omega_p^2}{\omega^3},$$

Energy
deposited
by axions

$$P \sim \frac{\Gamma_a}{\Gamma_t} \epsilon_z'' \int ||\mathbf{E}|^2 dV$$
Energy in the antenna

THE READOUT

• How to detect the induced electric field?



HOW IS THE INDUCED E-FIELD?

Standing wave in a cavity

Modes in a lossy medium





RESONANT MODES

Geometry factor for a cylindrical plasma



$$P_s = \kappa \mathcal{G} V \frac{Q}{m_a} \rho_a g_{a\gamma}^2 B_{\rm e}^2 \,,$$

- Resonance moved away from the plasma frequency
- Some modes overlap, careful with axion coupling to all the modes

STRUCTURE OF THE MODES

$$\begin{split} & B_T \quad E_Z \\ & \nabla \times \mathbf{B}_a - \dot{\mathbf{E}}_a = -g_{a\gamma\gamma} \left(\mathbf{E}_0 \times \nabla a - \dot{a} \mathbf{B}_0 \right) \,, \\ & \nabla \times \mathbf{E}_a + \dot{\mathbf{B}}_a = 0 \,, \\ & \nabla \cdot \mathbf{B}_a = 0 \,, \end{split}$$

$$\boldsymbol{\nabla} \cdot \mathbf{E}_a = -g_{a\gamma\gamma} \, \mathbf{B}_0 \cdot \boldsymbol{\nabla} a \,,$$

- B_T = TM modes, excited by axions
- *E_z*=0= TE modes, not excited by axions

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SIMULATION VS REALITY



• The agreement can be improved by better simulations

See 2203.10083

THE REACH

 Reaching the QCD axion band in 2yr of data taking

 Stage II upgraded detector design reaching the SQL



THANKS FOR YOUR ATTENTION!



Department of Physics

Postdoctoral Fellow in Experimental Astroparticle Physics

https://www.su.se/department-of-physics/about-the-department/work-with-us/postdoctoral-fellowin-experimental-astroparticle-physics-1.698316