PROBING SCALAR PARTICLES AND FORCES WITH COMPACT OBJECTS

Mateja Bošković

GWAT, DESY Theory Fellows meeting 2023 06.12.2023.



(Averaged) trajectory



- BSc, MSc in Theoretical Physics @ Uni Belgrade; MSc thesis w Vitor Cardoso at Tecnico Lisboa
- PhD in Astroparticle Physics w Enrico Barausse at SISSA

Overview

- ► Understanding the nature of DM and DE → search for new particles; we focus on massive and massless, force-mediating, scalars
- GW astronomy allows one to probe some of these models by studying the coalescence of compact object binaries
- Alternative motivation: Scalars are accessible models of matter under strong gravity conditions

Overview

- Gravitational probes of ultra-light bosons
- Solitons in strong gravity
- Gravitational probes of scalar fifth forces

(i) Gravitational probes of ultra-light bosons (1/3)

- Compton/De Broglie wavelengths on astrophysical scales from stellar BHs to galaxies
- Large occupancies in the case of DM halos, superradiant clouds etc.
- e.g. Motion in time-periodic background generated by axion DM MB+ [1806.07331]



Reviews: Hui [2101.11735], Brito, Cardoso, Pani [1501.06570]; Fig: Baumann+ [1912.04932]

(i) Gravitational probes of ultra-light bosons (2/3)

 Axion-photon coupling and BH superradiance (parametric instability) MB+ [1811.04945]

•
$$\mathscr{L} \supset \frac{\beta}{f_a} a^* FF$$
 , $R > \lambda_*^{-1}$, $\lambda_* \sim \frac{\beta}{f_a} \mu \langle a \rangle_c$

• Typically plasma-blocked for $\omega_{
m pl} \stackrel{\scriptstyle >}{\scriptstyle \sim} \mu$



Refs: Rosa, Kephart [1709.06581], Sen [1805.06471], MB+ [1811.04945]; Fig/further development: Spieksma+ [2306.16447]

(i) Gravitational probes of ultra-light bosons (3/3)

- Resonantly enhanced level transitions; ionization
- Significantly enriches phenomenology; cloud survival entangled with the orbit formation
- Ongoing work [w. Koschnitzke, Porto]



Refs: Baumann+ [1804.03208, 1912.04932, 2112.14777], Tomaselli, Spieksma, Bertone [2305.15460], Brito, Shah [2307.16093]; Fig: Baumann+ [1912.04932]

(ii) (Pseudo-)soliton stars: what and why

- (Pseudo-)soliton stars: localized, finite-energy and stable (long living) solutions of the EoM of a field theory incld. gravity
- Motivation 1: connection with dark matter and EU models Bertone+ [1907.10610]
- Motivation 2: ECO paradigm ("no stone unturned")
 Giudice, McCullough, Urbano [1605.01209], Cardoso, Pani [1904.05363]
 - * Consistent with known & tested physics? Formation mechanism? Stable (on astro/cosmo scales)?
- Motivation 3: toy model of matter in strong gravity
 - $\star\,$ Everything is in the action

(ii) Buchdahl bound and beyond

- WEC* + micro stability** \implies Buchdahl bound $C_B \le 0.44$ (constant density star); $C = GM/(Rc^2)$
- Subluminal condition $c_{\rm s} = \sqrt{\partial P / \partial \rho} \le 1$ lowers the Buchdahl bound:
 - * saturated by LinEoS $\rho \propto P$: $C_{B+C} = 0.354$ Urbano, Veermäe [1810.07137]
 - radially stable elastic objects must satisfy C_{EOmax} < 0.376 Alho+ [2107.12272, 2202.00043]
- What microphysics can saturate these conditions?



Fig: Alho+ [2202.00043] * $\rho \ge 0 \land \rho + P \ge 0$, ** $P \ge 0 \land dP/d\rho \ge 0$

(ii) (Soliton) boson stars

- Complex scalars w. U(1): $\mathscr{L}_{\Phi} = -\partial_{\mu} \Phi^{\dagger} \partial^{\mu} \Phi V(|\Phi|)$
- Parametrized deviation from the degenerate vacuum (gravitating Q-balls)

$$V_6 = \phi_0^2 \Big[(\mu^2 - \omega_0^2) \varphi^2 (1 - \varphi^2)^2 + \omega_0^2 \varphi^2 \Big], \varphi = \phi / \phi_0$$

Thin wall regime: bulk of the star is in the degenerate vacuum



Review: Liebling, Palenzuela [1202.5809]

(ii) SBSs are maximally stiff and compact

- Degenerate vacuum $\omega_0 = 0$
 - * Effective LinEoS in the bulk $\varphi \approx 1 \rightarrow \varphi' \approx V \approx 0$ $\rightarrow P \approx \rho$

$$\star (c_s)_a \approx 1 - 4(\varphi_c - 1)$$

- * Parameter space scanned w. $\Lambda = \sigma_0/M_{\rm Pl}$; thin wall realizable in the ultra-compact subspace: $\Lambda \lesssim 0.25$
- ► False vacuum $\omega_0 \neq 0$ * $(c_s^2)_a \approx \frac{2 - (\omega_0/\mu)^2}{2 + (\omega_0/\mu)^2}$ * $C_{\text{max}} \lesssim C_{\text{B+C}} - 0.06(\omega_0/\mu)^2$





Ref: MB, Barausse [2111.03870]; Consistent with the subsequent results of Cardoso+ [2112.05750], Collodel, Doneva [2203.08203]

(ii) (S)BSs abhor angular momentum [1/2]

- ▶ BS have quantized angular momentum J = kQ, $k \in \mathbb{N}$
- Rotating BSs generically suffer from non-axisymmetric instability Sanchis-Gual+ [1907.12565] ...
- ... which can be quenched w. sufficiently strong self-interactions, incl. SBS Siemonsen, East [2011.08247], Dmitriev+ [2104.00962]
- Can rotating SBS form from the binary inspiral of the non-rotating ones?



Fig: Siemonsen, East [2011.08247]

(ii) (S)BSs abhor angular momentum [2/2]

- If $M < M_{max}$ BS will form; else BH
- Parameterized condition for the rotating remnant J_{c,K} ≥ N(M₁) + N(M₂) & C > C_{NAI}
- Instead of rotating remnants, in two cases excess angular momentum is emitted in the form of blobs
- For q > 1, blobs can induce superkicks v ~ 0.05c
- Subsequent work: formation of rotating BS possible but not generic Siemonsen, East [2302.06627]

 Binary SBS simulations from Palenzuela+ [1710.09432], Bezares, MB+ [2201.06113]





(iii) Fifth force screening

- ▶ Fifth force screening necessary to evade Solar System bounds
- Is screening effective beyond staticity and spherical symmetry?
- General scalar theory in the decoupling limit $\varphi = \bar{\varphi} + \bar{\pi}$ $\mathscr{L} = -\frac{1}{2} \bar{Z}^{\mu\nu} \partial_{\mu} \pi \partial_{\nu} \pi - \frac{1}{2} \bar{m}^2 \pi^2 + \bar{g} T + \dots$, $\bar{Z} = \bar{Z} (\partial \bar{\phi}, \partial^2 \bar{\phi}, \dots)$
- Fifth-force potential of a point source (all functions of $\bar{\varphi}$): $\hat{\pi} \sim \frac{\bar{g}}{\sqrt{\bar{z}}} \frac{\exp\left(-\frac{\bar{m}}{\sqrt{\bar{z}}}r\right)}{r}$
- Varieties of screening:
 - $\star \ \hat{\pi}$ as a trigger (via potential): weak coupling \bar{g} (symmetron), large mass \bar{m} (chameleon)
 - * $\partial \hat{\pi}$ as a trigger (via acceleration) \bar{Z} : kinetic screening
 - $\star \ \partial^2 \hat{\pi}$ as a trigger (via curvature) \bar{Z} : Vainshtein mechanism

Review: Joyce+ [1407.0059]

(iii) k-essence

• k-essence action $\mathscr{L} = K(X)$

$$\mathcal{K} = -\frac{1}{2}X + \Lambda^4 \sum_{n=2}^{N} \frac{c_n}{2n} \left(\frac{X}{\Lambda^4}\right)^n, X = g^{\mu\nu} \partial_{\mu} \varphi \partial_{\nu} \varphi$$

- Cosmological context $\Lambda \sim \sqrt{H_0 M_{\text{Pl}}} \sim \text{meV}$
 - ★ Only unconstrained sector of Horndeski after GW170817 and requiring GW → DE Creminelli+ ['17, '18, '19]
- k-mouflage K(X): turns off the fifth force when X ≥ Λ⁴
- Radiative stability for large X de Rham, Ribeiro [1405.5213]

Review: Joyce+ [1407.0059]



(iii) Screening in a binary

- At the Newtonian order $\partial_i(K_X \partial^i \varphi) = \frac{\alpha}{2M_{\text{Pl}}} T$
- Helmholtz decomposition:
 - $egin{aligned} &\chi\equiv \mathcal{K}_Xoldsymbol{
 aligned} arphi\ ,\ &\chi=-rac{1}{2}oldsymbol{
 aligned} \psi+oldsymbol{B} \end{aligned}$
- ▶ Identified the regime where **B** can be ignored: $q \gg 1$, $\left(\frac{r_{sc}}{D}\right)^2 \ll 1$
- Checked numerically; qualitatively fine when q ≈ 1
- Saddle point: attractive forces cancel, breakdown of screening $\frac{\delta}{D} \simeq \frac{1}{\kappa} \frac{q^{3/2}}{(1+\sqrt{q})^4}$
- e.g. Earth-Moon $(\delta \approx 0.2 \text{km})$, Sun-Earth $(\delta \approx 1 \text{km})$ [$\Lambda \sim \text{meV}$]



Ref: MB, Barausse [2305.07725]

On-going interests

Phenomenology with gravitational two-body problem

- $\star\,$ Are scalar clouds robust? New signatures?
- $\star\,$ Beyond boson clouds and ECOs
- Classical non-linearities and EFTs
 - $\star\,$ Screening vs. positivity bounds
 - $\star\,$ Breakdown of the well-posedness of the Cauchy problem