Flavors from time-dependent D7 embeddings

J. Große

Institute of Physics Jagiellonian University



DESY Theory Workshop - QCD: string theory meets collider physics

Hamburg, September 25-28, 2007

arXiv:0709.3910







3



2 D7 embedding



Image: A matrix

э

• Choose proper time/rapidity coordinates (τ, y, x^2, x^3)

$$x^0 = \tau \cosh y$$
 $x^1 = \tau \sinh y$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance (y)
 - infinitely large nuclei $(\not\!\!\!\!/^2,\not\!\!\!/^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance (y)
 - infinitely large nuclei $(\not z^2, \not z^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance (y)
 - infinitely large nuclei $(\not x^2, \not x^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance ($\not\!\!\!\!/$
 - infinitely large nuclei $(\not\!\!\!\! x^2,\not\!\!\!\! x^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$

$$T_{\tau\tau} = \varepsilon(\tau) \quad T_{yy} = -\tau^2 \left(\varepsilon(\tau) + \tau \varepsilon'(\tau)\right) \quad T_{xx} = \varepsilon(\tau) + \frac{1}{2}\tau \varepsilon'(\tau)$$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance (y)
 - infinitely large nuclei $(\not\!\!\!\!/^2,\not\!\!\!/^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$
- Viscous hydrodynamics

$$T_{\mu\nu} = (\varepsilon + p)u_{\mu}u_{\nu} + p\eta_{\mu\nu} + \eta(\nabla_{\mu}u_{\nu} + \nabla_{\nu}u_{\mu}) \qquad u^2 = -1$$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance (y)
 - infinitely large nuclei $(\not\!\!\!\!/^2,\not\!\!\!/^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$
- Viscous hydrodynamics

$$T_{\mu\nu} = (\varepsilon + p)u_{\mu}u_{\nu} + p\eta_{\mu\nu} + \eta(\nabla_{\mu}u_{\nu} + \nabla_{\nu}u_{\mu}) \qquad u^2 = -1$$

$$\implies \varepsilon(\tau) \sim \frac{\varepsilon_0}{\tau^{4/3}} (1 - \frac{2\eta_0}{\varepsilon_0^{1/4} \tau^{2/3}} + \dots) \qquad \eta = \eta_0 T^3$$

- Choose proper time/rapidity coordinates (τ, y, x^2, x^3)
- Assume
 - boost invariance (y)
 - infinitely large nuclei $(\not\!\!\!\!/^2,\not\!\!\!/^3)$
 - energy momentum conservation $\partial_{\mu} T^{\mu\nu} = 0$
 - tracelessness $T_{\mu}{}^{\mu} = 0$
- Viscous hydrodynamics

$$T_{\mu\nu} = (\varepsilon + p)u_{\mu}u_{\nu} + p\eta_{\mu\nu} + \eta(\nabla_{\mu}u_{\nu} + \nabla_{\nu}u_{\mu}) \qquad u^2 = -1$$

$$\implies \varepsilon(\tau) \sim \frac{\varepsilon_0}{\tau^{4/3}} (1 - \frac{2\eta_0}{\varepsilon_0^{1/4} \tau^{2/3}} + \dots) \qquad \eta = \eta_0 T^3$$
$$\varepsilon \sim T^4$$

Dual geometry of $\mathcal{N} = 4$ SYM viscous fluid

 ${\it Janik/Peschanski, Nakamura/Sin, Janik/Heller}$

$$\begin{split} ds^2 &= \frac{r^2}{L^2} (-e^{\mathcal{A}(v,\tau)} d\tau^2 + e^{\mathcal{B}(v,\tau)} \tau^2 dy^2 + e^{\mathcal{C}(v,\tau)} dx_{\perp}^2) + \frac{L^2 dr^2}{r^2} \\ v &= \frac{\varepsilon_0^{1/4}}{r \, \tau^{1/3}} \end{split}$$

• *A*, *B*, *C* ...

• How to obtain them

• Expand the coefficients

$$\mathcal{A} = a_0(v) + a_1(v)\tau^{-2/3} + \dots$$

- Solve Einstein's equation order by order
- From horizon: $T^4 \sim \varepsilon$
- Regularity \implies Bjorken dynamics and $\eta/s = 1/4\pi$

Dual geometry of $\mathcal{N} = 4$ SYM viscous fluid

Janik/Peschanski, Nakamura/Sin, Janik/Heller

$$\begin{split} ds^2 &= \frac{r^2}{L^2} (-e^{\mathcal{A}(v,\tau)} d\tau^2 + e^{\mathcal{B}(v,\tau)} \tau^2 dy^2 + e^{\mathcal{C}(v,\tau)} dx_{\perp}^2) + \frac{L^2 dr^2}{r^2} \\ v &= \frac{\varepsilon_0^{1/4}}{r \, \tau^{1/3}} \end{split}$$

• *A*, *B*, *C* ...

• How to obtain them

• Expand the coefficients

$$\mathcal{A} = a_0(v) + a_1(v)\tau^{-2/3} + \dots$$

• Solve Einstein's equation order by order

- From horizon: $T^4 \sim \varepsilon$
- Regularity \implies Bjorken dynamics and $\eta/s = 1/4\pi$

Dual geometry of $\mathcal{N} = 4$ SYM viscous fluid

Janik/Peschanski, Nakamura/Sin, Janik/Heller

$$\begin{split} ds^2 &= \frac{r^2}{L^2} (-e^{\mathcal{A}(v,\tau)} d\tau^2 + e^{\mathcal{B}(v,\tau)} \tau^2 dy^2 + e^{\mathcal{C}(v,\tau)} dx_{\perp}^2) + \frac{L^2 dr^2}{r^2} \\ v &= \frac{\varepsilon_0^{1/4}}{r \tau^{1/3}} \end{split}$$

- *A*, *B*, *C* ...
- How to obtain them
 - Expand the coefficients

$$\mathcal{A} = a_0(v) + a_1(v)\tau^{-2/3} + \dots$$

- Solve Einstein's equation order by order
- From horizon: $T^4 \sim \varepsilon$
- Regularity \implies Bjorken dynamics and $\eta/s = 1/4\pi$

Dual geometry of $\mathcal{N} = 4$ SYM viscous fluid

 ${\it Janik/Peschanski, Nakamura/Sin, Janik/Heller}$

$$\begin{split} ds^2 &= \frac{r^2}{L^2} (-e^{\mathcal{A}(v,\tau)} d\tau^2 + e^{\mathcal{B}(v,\tau)} \tau^2 dy^2 + e^{\mathcal{C}(v,\tau)} dx_{\perp}^2) + \frac{L^2 dr^2}{r^2} \\ v &= \frac{\varepsilon_0^{1/4}}{r \, \tau^{1/3}} \end{split}$$

• *A*, *B*, *C* ...

• How to obtain them

• Expand the coefficients

 $\mathcal{A} = a_0(v) + a_1(v)\tau^{-2/3} + \dots$

• Solve Einstein's equation order by order

• From horizon: $T^4 \sim \varepsilon$

• Regularity \implies Bjorken dynamics and $\eta/s = 1/4\pi$







A = A = A = A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A
 A

э

• Full metric

$$ds^{2} = ds_{4}^{2} + \frac{1}{r^{2}} \left(d\rho^{2} + \rho^{2} d\Omega_{3}^{2} + (dX^{8})^{2} + (dX^{9})^{2} \right)$$
$$r^{2} = \rho^{2} + (X^{8})^{2} + (X^{9})^{2}$$



• Full metric

$$ds^{2} = ds_{4}^{2} + \frac{1}{r^{2}} \left(d\rho^{2} + \rho^{2} d\Omega_{3}^{2} + (dX^{8})^{2} + (dX^{9})^{2} \right)$$
$$r^{2} = \rho^{2} + (X^{8})^{2} + (X^{9})^{2}$$

• DBI action
$$(X^9 = \Phi(\rho, \tau), X^8 = 0)$$

$$S \sim \int d\tau d\rho \, \tau \rho^3 \, \# \sqrt{1 + \# (\partial_\rho \Phi)^2 - \# \frac{(\partial_\tau \Phi)^2}{r^4}}$$

J. Große (Institute of Physics, JU) Flavors from time-dependent D7 embeddings DESY Theory Workshop 7 / 18

• Full metric

$$ds^{2} = ds_{4}^{2} + \frac{1}{r^{2}} \left(d\rho^{2} + \rho^{2} d\Omega_{3}^{2} + (dX^{8})^{2} + (dX^{9})^{2} \right)$$
$$r^{2} = \rho^{2} + (X^{8})^{2} + (X^{9})^{2}$$

• DBI action
$$(X^9 = \Phi(\rho, \tau), X^8 = 0)$$

$$S \sim \int d\tau d\rho \, \tau \rho^3 \, \# \sqrt{1 + \# (\partial_\rho \Phi)^2 - \# \frac{(\partial_\tau \Phi)^2}{r^4}}$$

• Asymptotic behaviour as in AdS_5

$$\Phi \xrightarrow[\rho \to \infty]{} m + \frac{c}{\rho^2} + \dots$$

• Full metric

$$ds^{2} = ds_{4}^{2} + \frac{1}{r^{2}} \left(d\rho^{2} + \rho^{2} d\Omega_{3}^{2} + (dX^{8})^{2} + (dX^{9})^{2} \right)$$
$$r^{2} = \rho^{2} + (X^{8})^{2} + (X^{9})^{2}$$

• Asymptotic behaviour as in AdS_5

$$\Phi \xrightarrow[\rho \to \infty]{} m + \frac{c}{\rho^2} + \dots$$

• Ansatz

$$\Phi = m + \sum_{k} \phi_k(\rho) \tau^{-\frac{k}{3}}$$

• Asymptotic behaviour as in AdS_5

$$\Phi \xrightarrow[\rho \to \infty]{} m + \frac{c}{\rho^2} + \dots$$

$$\Phi = m + \sum_{k} \phi_k(\rho) \tau^{-\frac{k}{3}}$$

$$\Phi \xrightarrow[\rho \to \infty]{} (m + \sum m_k \tau^{-\frac{k}{3}}) + \frac{1}{\rho^2} (\sum c_k \tau^{-\frac{k}{3}}) + \dots$$

 \implies boundary condition $m_k = 0$

• Asymptotic behaviour as in AdS_5

$$\Phi \xrightarrow[\rho \to \infty]{} m + \frac{c}{\rho^2} + \dots$$

$$\Phi = m + \sum_{k} \phi_k(\rho) \tau^{-\frac{k}{3}}$$

$$\Phi \xrightarrow[\rho \to \infty]{} (m + \sum m_k \tau^{-\frac{k}{3}}) + \frac{1}{\rho^2} (\sum c_k \tau^{-\frac{k}{3}}) + \dots$$

 \implies boundary condition $m_k = 0$

• Moreover: *regularity* in the interior

• Equation of motion

$$\frac{1}{\rho^3}\partial_{\rho}(\rho^3\partial_{\rho}\phi_k(\rho)) = \#(\rho) \cdot \begin{cases} 1 & \text{if } k = 8\\ -4\eta_0\varepsilon_0^{-\frac{1}{4}} & \text{if } k = 11\\ 0 & \text{else; provided } k < 14 \end{cases}$$

• Homogeneous Solution

$$\frac{1}{\rho^3} \partial_\rho (\rho^3 \partial_\rho \phi_k(\rho)) = 0$$
$$\implies \phi_k = m_k + \frac{c_k}{\rho^2}$$

• Equation of motion

$$\frac{1}{\rho^3}\partial_{\rho}(\rho^3\partial_{\rho}\phi_k(\rho)) = \#(\rho) \cdot \begin{cases} 1 & \text{if } k = 8\\ -4\eta_0\varepsilon_0^{-\frac{1}{4}} & \text{if } k = 11\\ 0 & \text{else; provided } k < 14 \end{cases}$$

• Homogeneous Solution

$$\frac{1}{\rho^3} \partial_\rho (\rho^3 \partial_\rho \phi_k(\rho)) = 0$$
$$\implies \phi_k \equiv 0$$

• Equation of motion

$$\frac{1}{\rho^3}\partial_{\rho}(\rho^3\partial_{\rho}\phi_k(\rho)) = \#(\rho) \cdot \begin{cases} 1 & \text{if } k = 8\\ -4\eta_0\varepsilon_0^{-\frac{1}{4}} & \text{if } k = 11\\ 0 & \text{else; provided } k < 14 \end{cases}$$

Final Solution

$$\Phi(\rho,\tau) = m + c \, \frac{\rho^4 + 3\rho^2 m^2 + 3m^4}{(m^2 + \rho^2)^3},$$
$$c = -\frac{\varepsilon_0^2 L^{16}}{54m^5} \tau^{-\frac{8}{3}} \left(1 - 4\eta_0 \varepsilon_0^{-\frac{1}{4}} \tau^{-\frac{2}{3}} + \dots\right)$$

• Equation of motion

$$\frac{1}{\rho^3}\partial_{\rho}(\rho^3\partial_{\rho}\phi_k(\rho)) = \#(\rho) \cdot \begin{cases} 1 & \text{if } k = 8\\ -4\eta_0\varepsilon_0^{-\frac{1}{4}} & \text{if } k = 11\\ 0 & \text{else; provided } k < 14 \end{cases}$$

• Final Solution

$$\Phi(\rho,\tau) = m + c \, \frac{\rho^4 + 3\rho^2 m^2 + 3m^4}{(m^2 + \rho^2)^3},$$
$$c = -\frac{\varepsilon_0^2 L^{16}}{54m^5} \tau^{-\frac{8}{3}} \left(1 - 4\eta_0 \varepsilon_0^{-\frac{1}{4}} \tau^{-\frac{2}{3}} + \dots\right)$$

• Equation of motion

$$\frac{1}{\rho^3}\partial_{\rho}(\rho^3\partial_{\rho}\phi_k(\rho)) = \#(\rho) \cdot \begin{cases} 1 & \text{if } k = 8\\ -4\eta_0\varepsilon_0^{-\frac{1}{4}} & \text{if } k = 11\\ 0 & \text{else; provided } k < 14 \end{cases}$$

• Final Solution

$$\Phi(\rho,\tau) = m + c \, \frac{\rho^4 + 3\rho^2 m^2 + 3m^4}{(m^2 + \rho^2)^3},$$
$$\langle \mathcal{O} \rangle = \frac{1}{216\pi^4} \frac{N_f \lambda^3}{N_c^3} \frac{\varepsilon^2}{m_q^5} \sim N_c N_f \, T^8.$$

Time-dependent embedding



Time-dependent embedding



Time-dependent embedding



Break-down of perturbative solution

• Time to approach the horizon or bending backwards





2 D7 embedding



æ

Boost invariant mesons

• usual 5d ansatz: mesons from $\Phi = \Phi_{embedding} + \delta \Phi$

 $\delta \Phi(\rho) = \delta \phi(\rho) \exp(ikx) \mathcal{Y}^{\ell}(S^3)$

3

Boost invariant mesons

• usual 5d ansatz: mesons from $\Phi = \Phi_{embedding} + \delta \Phi$

$$\delta \Phi(\rho) = \delta \phi(\rho) \exp(ikx) \mathcal{Y}^{\ell}(S^3)$$

• 4d boost invariant mesons

$$\Box \phi = \left[-\frac{1}{\tau} \partial_{\tau} \tau \partial_{\tau} + \tau^{-2} \partial_{y}^{2} + \partial_{x}^{2} \right] \phi = M^{2} \phi$$
$$\implies \phi = J_{0}(\omega \tau) e^{\pm i k_{\perp} x_{\perp}}$$
$$M^{2} = \omega^{2} + k_{\perp}^{2} \qquad k_{\perp}^{2} = k_{3}^{2} + k_{4}^{2}$$

Boost invariant mesons

• usual 5d ansatz: mesons from $\Phi = \Phi_{embedding} + \delta \Phi$

$$\delta \Phi(\rho) = \delta \phi(\rho) \exp(ikx) \mathcal{Y}^{\ell}(S^3)$$

• 4d boost invariant mesons

$$\Box \phi = \left[-\frac{1}{\tau} \partial_{\tau} \tau \partial_{\tau} + \tau^{-2} \partial_{y}^{2} + \partial_{x}^{2} \right] \phi = M^{2} \phi$$
$$\implies \phi = J_{0}(\omega \tau) e^{\pm i k_{\perp} x_{\perp}}$$
$$M^{2} = \omega^{2} + k_{\perp}^{2} \qquad k_{\perp}^{2} = k_{3}^{2} + k_{4}^{2}$$

• 5d ansatz

$$\delta \Phi = \delta \phi(\rho, \tau) J_0(\omega(\tau) \tau) \mathcal{Y}^{\ell}(S^3)$$
$$\delta \phi(\rho, \tau) = \sum \delta \phi_k(\rho) \tau^{-\frac{k}{3}}$$
$$\omega(\tau) = \sum \omega_k \tau^{-\frac{k}{3}}$$

æ

-

• 5d ansatz

$$\delta \Phi = \delta \phi(\rho, \tau) J_0(\omega(\tau) \tau) \mathcal{Y}^{\ell}(S^3)$$
$$\delta \phi(\rho, \tau) = \sum \delta \phi_k(\rho) \tau^{-\frac{k}{3}}$$
$$\omega(\tau) = \sum \omega_k \tau^{-\frac{k}{3}}$$

• equation of motion $\mathbf{y} := -\rho^2/m^2$

$$\begin{bmatrix} \mathbf{y}(1-\mathbf{y})\partial_{\mathbf{y}}^{2} + (\mathbf{c} - (\mathbf{a} + \mathbf{b} - 1)\mathbf{y})\partial_{\mathbf{y}} - \mathbf{a}\mathbf{b} \end{bmatrix} \phi_{k}(\mathbf{y}) = \mathcal{I}_{k}$$
$$\mathbf{a} = -n - 1, \qquad \mathbf{b} = -n, \qquad \mathbf{c} = 2.$$

Kruczenski/Mateos/Myers/Winters

∃ ► < ∃ ►</p>

3

Image: A matrix



J. Große (Institute of Physics, JU) Flavors from time-dependent D7 embeddings DESY Theory Workshop 14 / 18

Mass spectrum



- scalars and pseudoscalars are degenerate
- vectors deviate about 3%

But there is something wrong...

J. Große (Institute of Physics, JU) Flavors from time-dependent D7 embeddings DESY Theory Workshop 14 / 18

Mass spectrum



-

$$M_{\phi}(n=0) = \frac{4\pi}{\sqrt{\lambda}} \cdot \left[m_q + \frac{9\lambda^2 \varepsilon_0}{80\pi^4 \tau^{4/3} m_q^3} \cdot \left(1 - \frac{2\eta_0}{3\tau^{\frac{2}{3}} \varepsilon_0^{\frac{1}{4}}} \right) \right]$$

Image: A matrix

æ

∃ ► < ∃</p>

Does not agree with AdS/Schwarzschild?



- horizontal lines: SUSY spectrum (LO)
- dots: numerical solution
- solid curve: NLO low temperature expansion (static)

• static meson mass

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q - \frac{9\lambda^2}{320} \frac{T^4}{m_q^3} \right]$$

• metric

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q + \frac{9\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• static meson mass

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q - \frac{3\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• metric

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q + \frac{9\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• static meson mass

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q - \frac{3\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• metric

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q + \frac{9\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• static meson mass

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q - \frac{3\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• metric

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q + \frac{9\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

agrees with perfect fluid

• static meson mass

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q - \frac{3\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$

• metric

R

$$M = \frac{4\pi}{\sqrt{\lambda}} \left[m_q + \frac{9\lambda^2}{80\pi^4} \frac{\varepsilon_0}{m_q^3} t^{-\frac{4}{3}} \right]$$
 Why?

• static meson mass

∂_t is $\mathcal{O}(1)$ in T expansion

• metric

 ∂_t is $\mathcal{O}(1/t)$ in 1/t expansion

Done

- time-dependent D7 embedding
- chiral condensate
- meson spectra

Surprise

• meson spectra not adiabatic

To do

- "black hole" solutions
- thermodynamics

Thank you

э

- (日)