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# Asymptotic Superluminality

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(Ongoing work)

23 July 2025, Cargese 2025 International Summer School, IESC

# Motivation

## Causality Bounds on EFTs

- An EFT taken beyond its regime of validity displays time advance in a gravitational background.
- This time advance can be used to bound—or rule out—EFTs.

## Subtleties of Gravitational Scattering

- Notion of time delay/advance requires an asymptotic Minkowski background: well-defined in  $D > 4$ , ambiguous in  $D = 4$ .
- Gauge dependence, ambiguous asymptotic structure.



# Gao-Wald Theorem

## Criterion for Promptness (GR)

Let  $(M, g_{ab})$  be a **null geodesically complete spacetime** satisfying the **null energy condition**:

$$R_{ab}k^ak^b \geq 0,$$

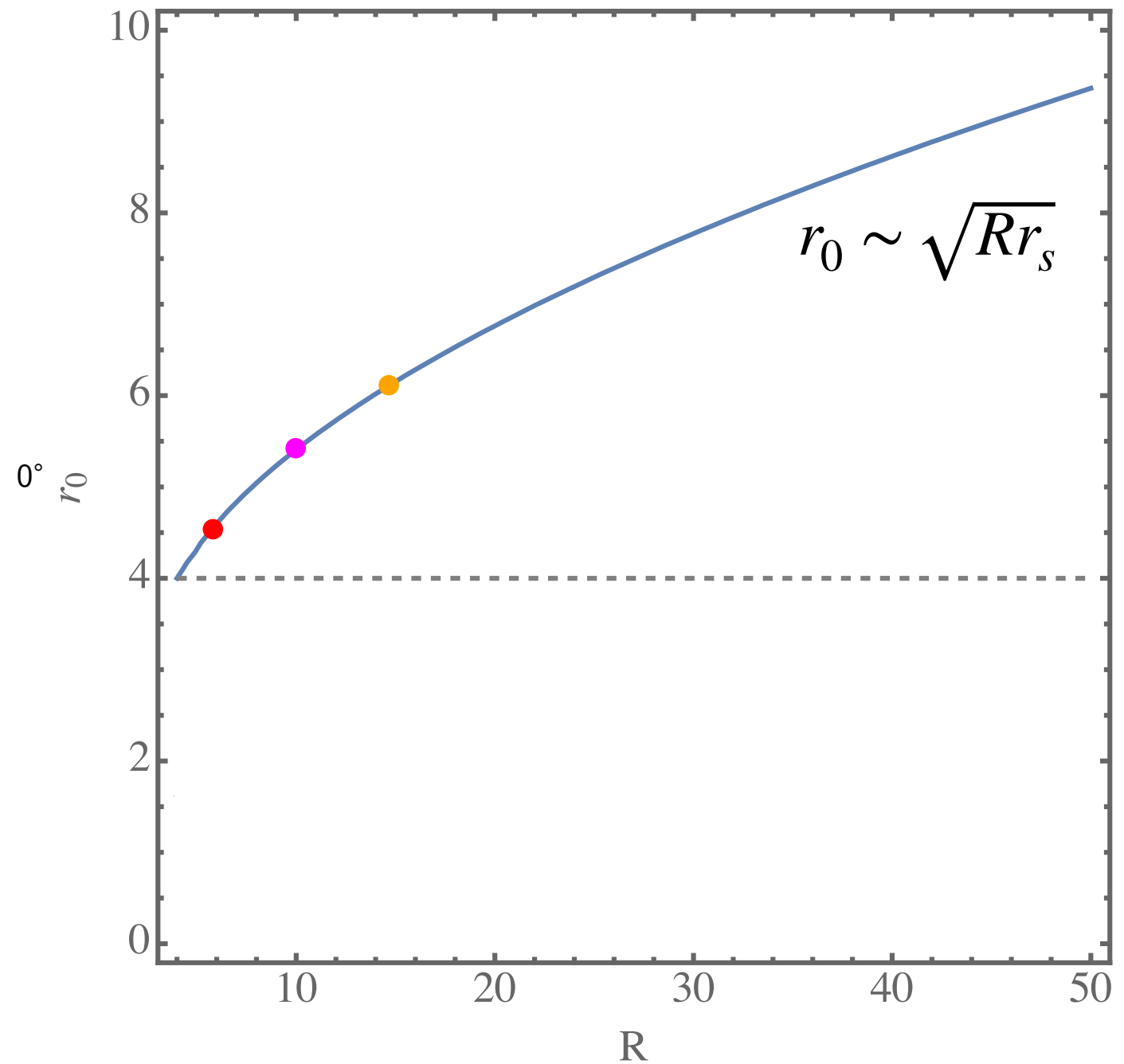
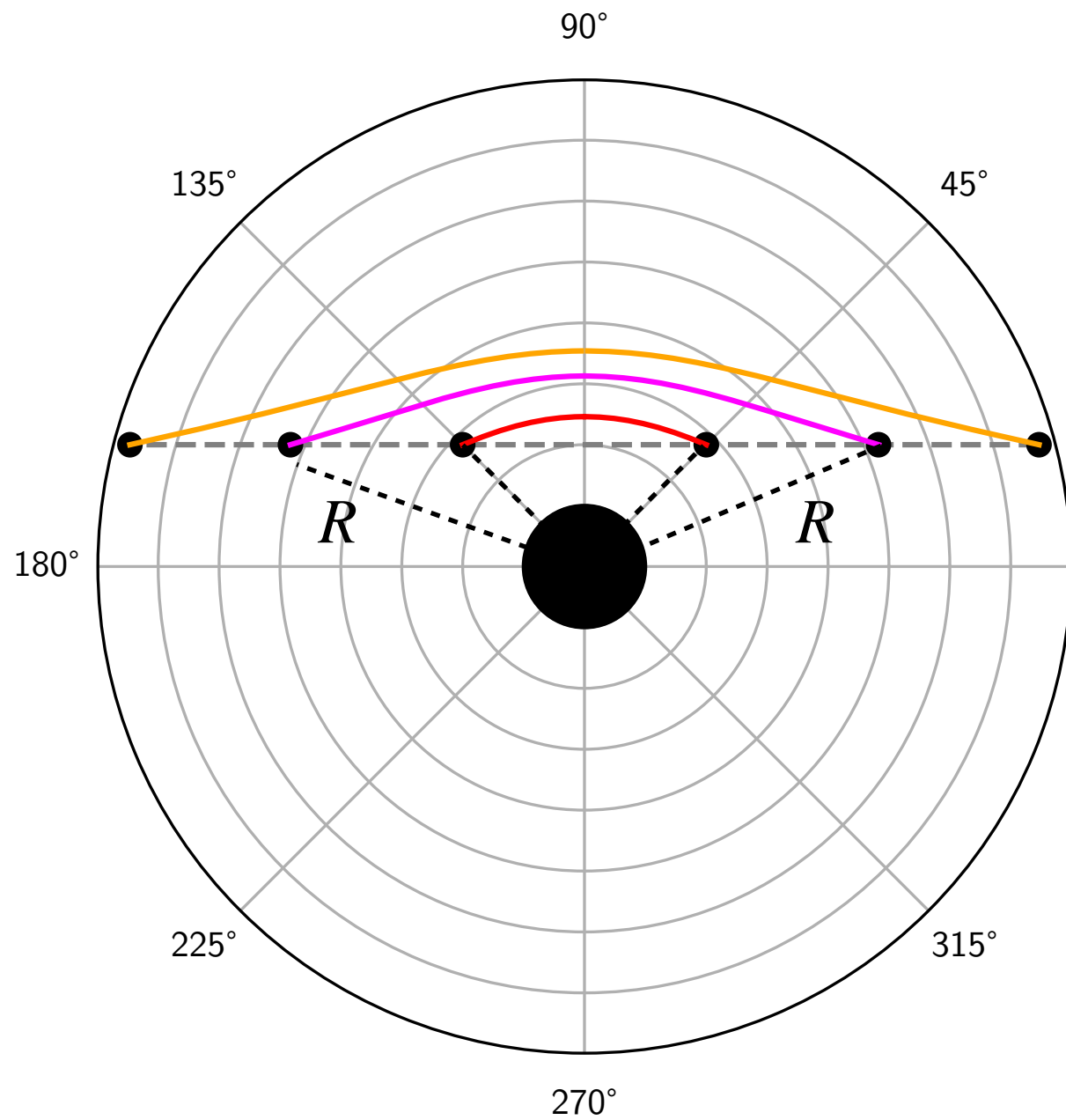
and the **null generic condition** i.e. each null geodesic contains a point in which

$$k_{[a}R_{b]cd[e}k_{f]}k^ck^d \neq 0,$$

then given any compact region  $K \subset M$ , there exists another compact region  $K'$  such that if  $q, p \notin K'$  and  $q \in J^+(p) - I^+(p)$ , then any causal curve  $\gamma$  connecting  $p$  to  $q$  cannot intersect the region  $K$ .

# Gao-Wald Theorem

## 4D Schwarzschild



# FFR Theory

## EFT of a Maxwell Field

- The first higher-order interaction relevant in Ricci-flat ( $R_{\mu\nu} = 0$ ) background is given by

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} R - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \alpha R_{\mu\nu\rho\sigma} F^{\mu\nu} F^{\rho\sigma} \right]$$

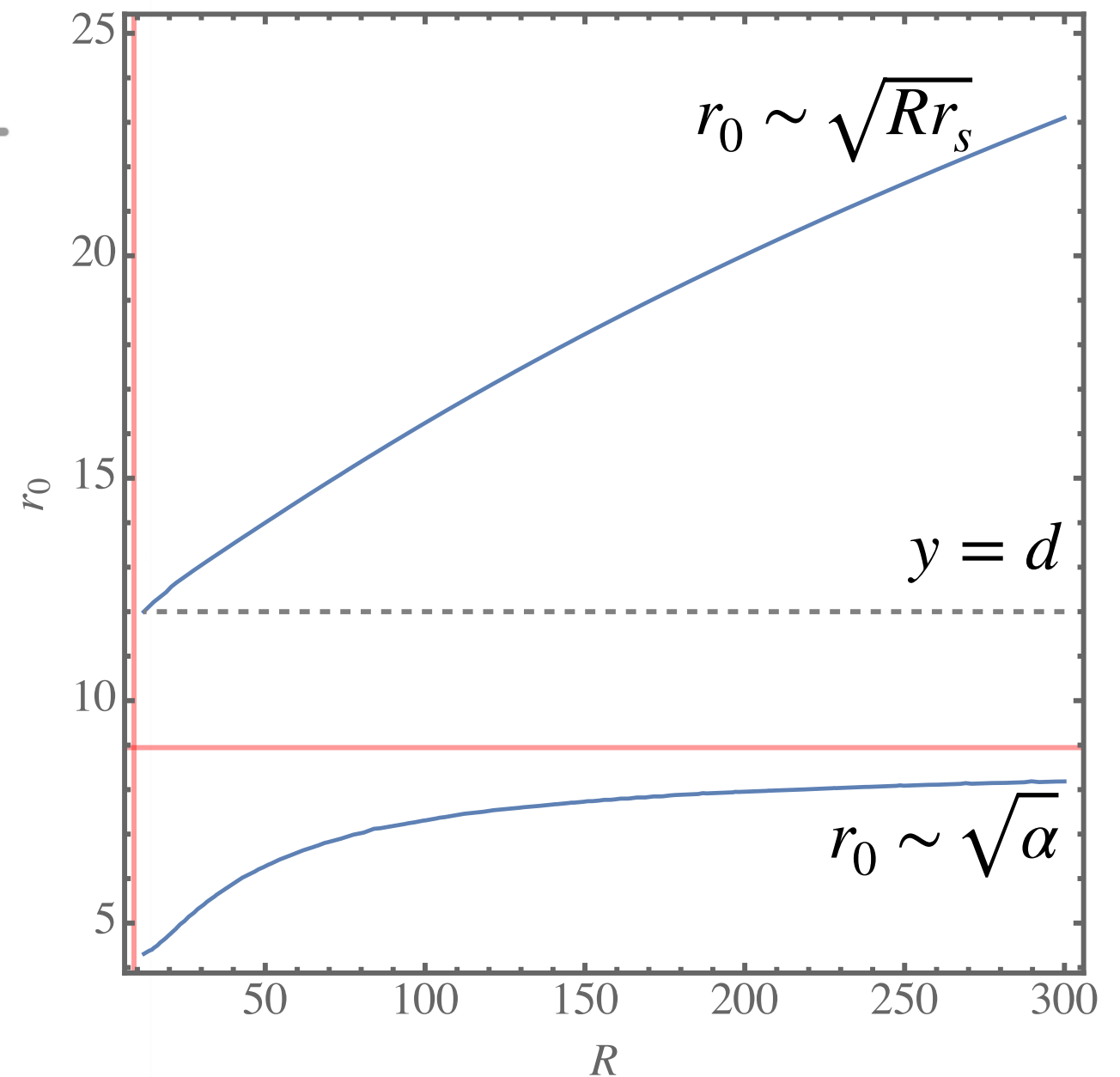
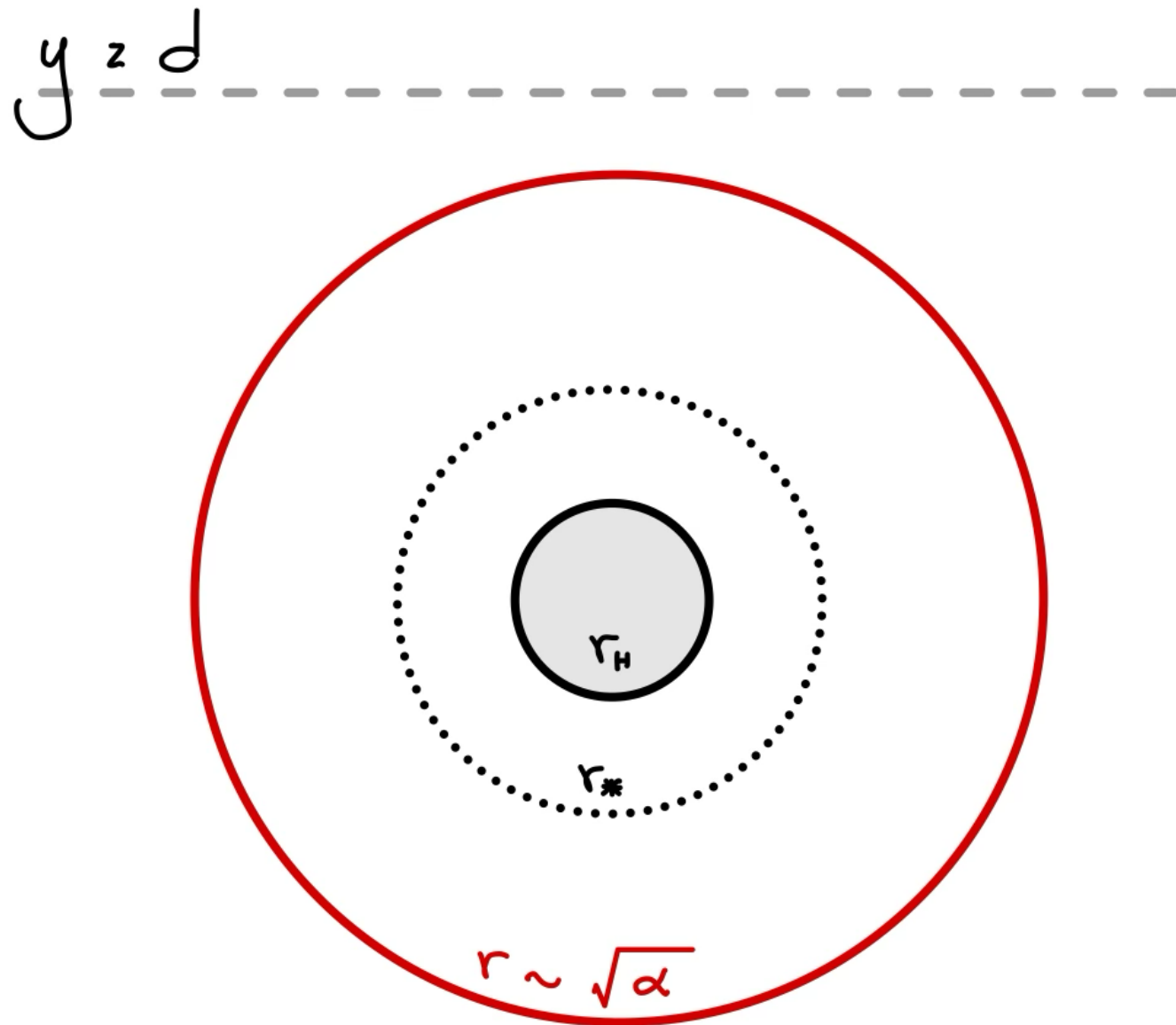
Photon Equation of Motion:

$$\nabla_\nu F^{\mu\nu} + \alpha R^{\mu\nu}{}_{\rho\sigma} \nabla_\nu F^{\rho\sigma} = 0$$

- Well known EFT that exhibits superluminality, cutoff of theory is  $\alpha \gtrsim r_s^2$ .

# FFR Theory

## Limit of Causal Curves



# FFR Theory

## Time Delay ( $D = 4$ )

- The gravitational time delay is log-divergent

$$\Delta t_{\text{GR}} \propto \ln\left(\frac{2R}{r_0}\right)$$

- Comparing the coordinate time of the two geodesics, we have

$$\Delta t_1 - \Delta t_2 \sim \ln\left(\frac{r_{0,2}}{r_{0,1}}\right) \sim \ln\left(\frac{\sqrt{\alpha}}{\sqrt{R}}\right) \xrightarrow{R \rightarrow \infty} \Delta t_1 < \Delta t_2$$

- The geodesic that goes to infinity is prompt i.e. the asymptotic structure of this (pathological) theory is still Schwarzschild!

# Towards a General Statement?

**Assertion 1:** This effect is general for any pathological EFTs on a 4D Schwarzschild background *provided that the time advance is bounded*.

(Can we formalise this into a ‘Theorem’?)

**Assertion 2:** Pathologies of the theory cannot be probed by scattering at infinity. They should be probed through other local means e.g. constructing CTCs.



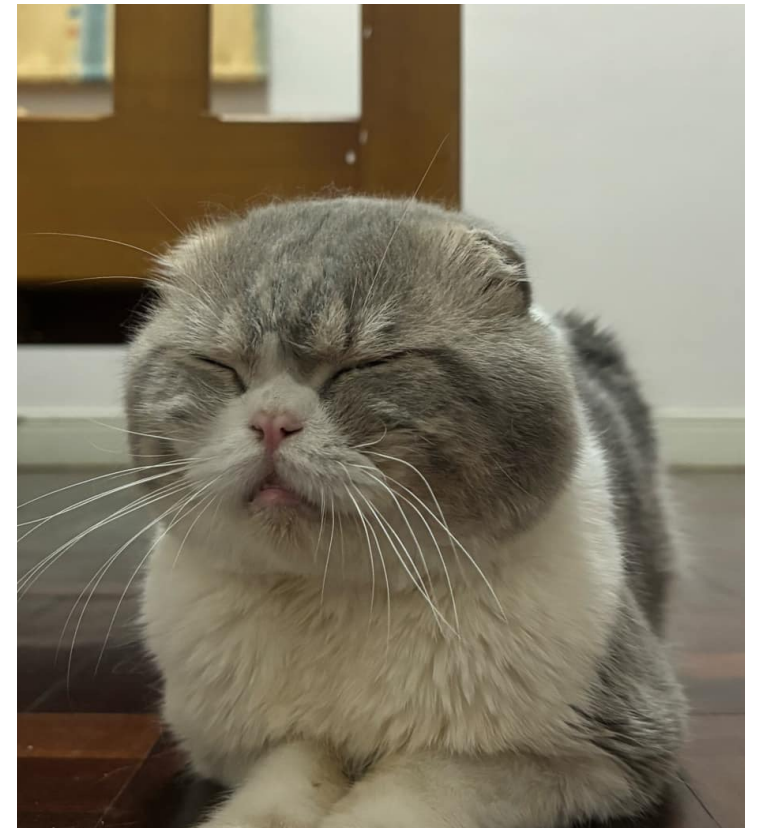
# Warin Patrick McBlain





Interests in Physics: Cosmology (in general)

Interests outside Physics: Eating, DRINKING & **C A T S**





# Q&A Session

# Motivation

## Causality Bounds on EFTs

- An EFT taken beyond its regime of validity displays time advance in a gravitational background.

Example: Modification of BHs—BH with scalar hair

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} R - \frac{1}{2} (\nabla_\mu \phi)^2 + M_{pl} \alpha \phi \mathcal{R}_{GB}^2 + M_{pl} \tilde{\alpha} \phi R_{\mu\nu\rho\sigma} \tilde{R}^{\mu\nu\rho\sigma} \right]$$

- The time delay of scalar-graviton scattering is given by

$$\Delta t_{\pm} = 2r_s \left( \ln \frac{b_0}{b} \pm \sqrt{2} \frac{|\hat{\alpha}|}{b^2} \right); \quad \hat{\alpha} = \alpha + i\tilde{\alpha}$$

- At small enough impact parameter  $b$ ,  $\Delta t_-$  can be negative i.e. a time advance—violation of causality.

# Motivation

## Causality Bounds on EFTs

- This time advance can be used to bound—or rule out—EFTs.

Example: Modification of BHs—BH with scalar hair

- The time delay of scattering implies

$$\Delta t_{\pm} = 2r_s \left( \ln \frac{b_0}{b} \pm \sqrt{2} \frac{|\hat{\alpha}|}{b^2} \right) \longrightarrow |\hat{\alpha}| \gtrsim b^2 \ln \frac{b_0}{b}$$

- To avoid time advances at low energies, the EFT must break down at distance scales such that this condition cannot be satisfied i.e.

$$\frac{1}{\Lambda} > b \longrightarrow |\hat{\alpha}| \gtrsim \frac{\ln(b_0 \Lambda)}{\Lambda^2}$$

# Subtlety I: Gauge Dependency

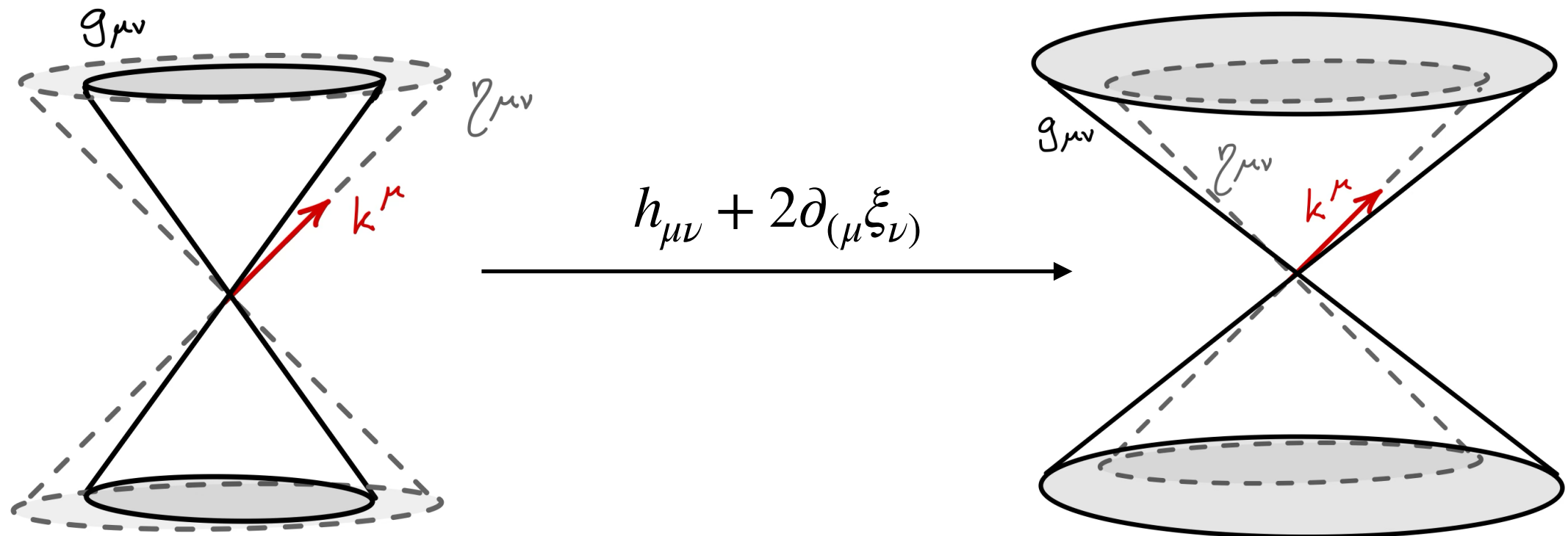
- In linearised gravity that satisfies the Null Energy Condition

$$R_{\mu\nu}k^\mu k^\nu \geq 0, \quad \eta_{\mu\nu}k^\mu k^\nu = 0,$$

the perturbation satisfies (assuming Lorenz gauge)

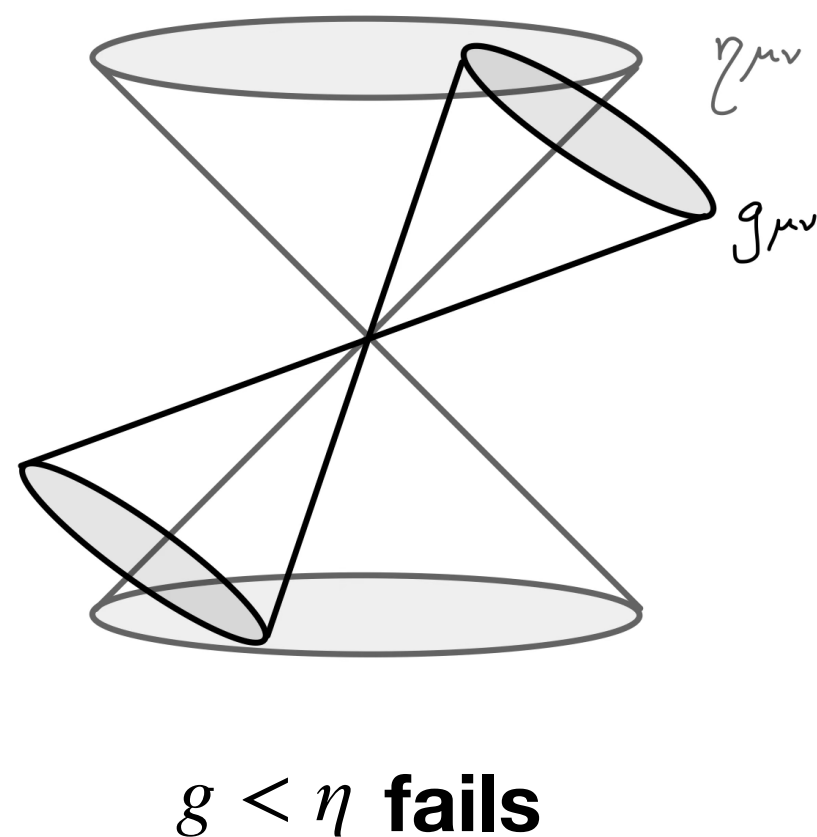
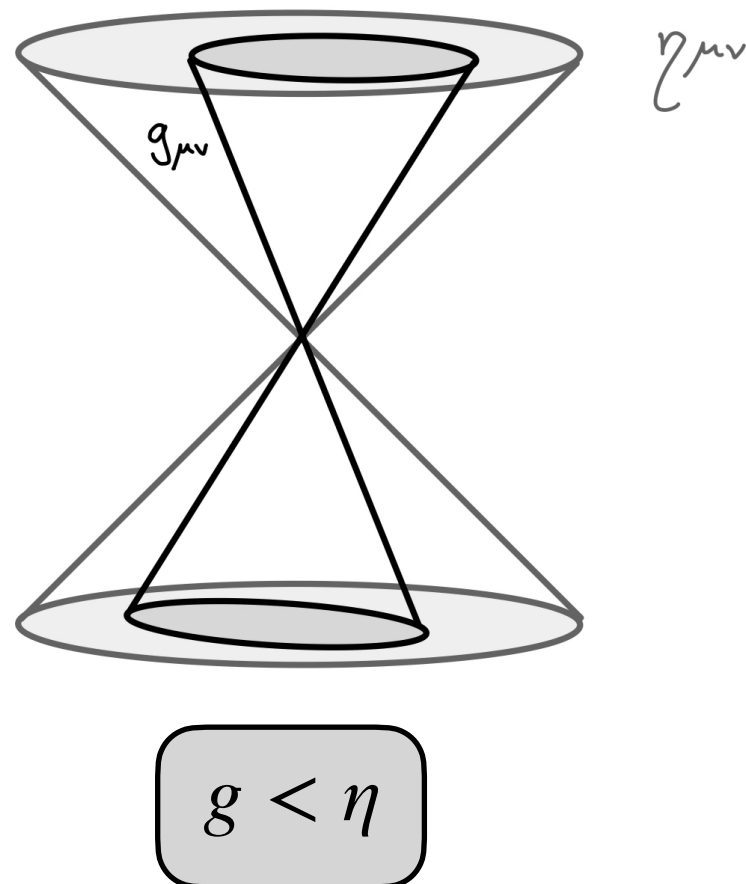
$$h_{\mu\nu}k^\mu k^\nu \geq 0$$

- A gauge transformation on  $h_{\mu\nu}$  can reverse the above conclusion!



# Subtlety II: Asymptotic Structure

- One could imagine that in a sensible theory, the fields in the theory should propagate according to its effective metric  $g_{\mu\nu}$  in a way that also obeys causality set by  $\eta_{\mu\nu}$ .



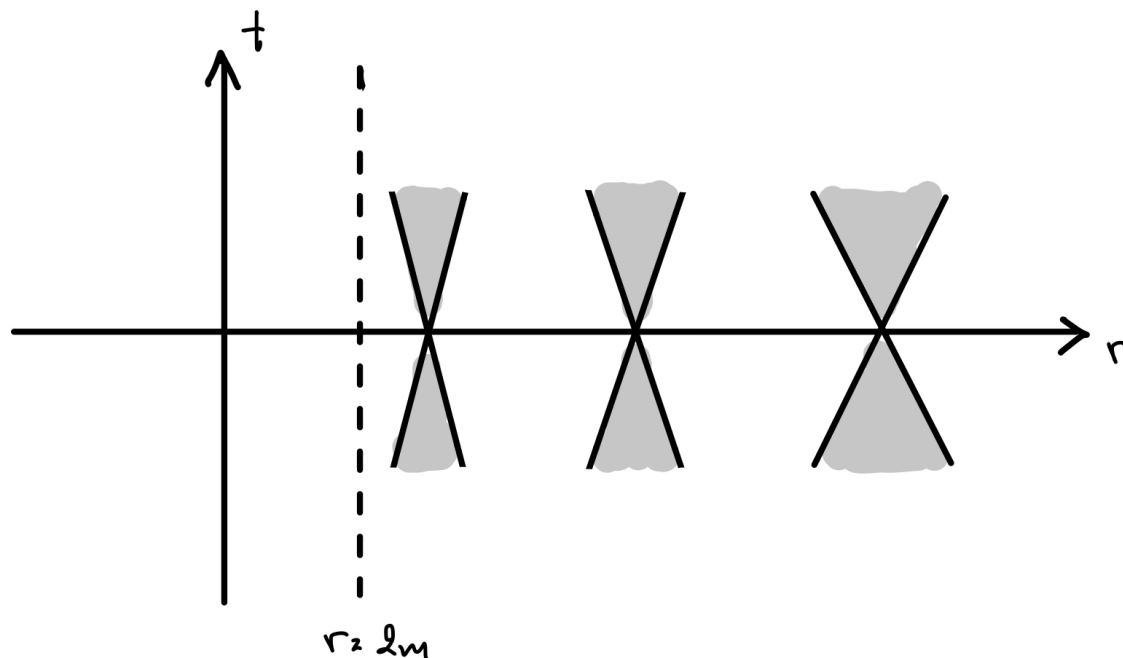
- However, this condition fails for many examples in GR.

# Subtlety II: Asymptotic Structure

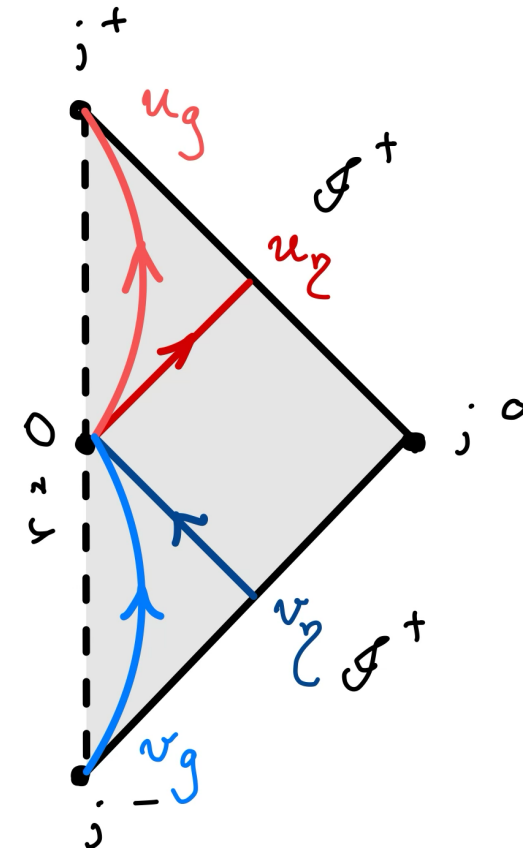
4D Schwarzschild, standard coordinates:

$$ds^2 = - \left( 1 - \frac{r_s}{r} \right) dt^2 + \left( 1 - \frac{r_s}{r} \right)^{-1} dr^2 + r^2 d\Omega^2$$

$$d\sigma^2 = - dt^2 + dr^2 + r^2 d\Omega^2$$



$$\frac{dt}{dr} = \pm \frac{1}{1 - r_s/r}, \quad g < \eta$$



$$u_g = t - r - r_s \ln(r - r_s)$$

$$v_g = t + r + r_s \ln(r - r_s)$$

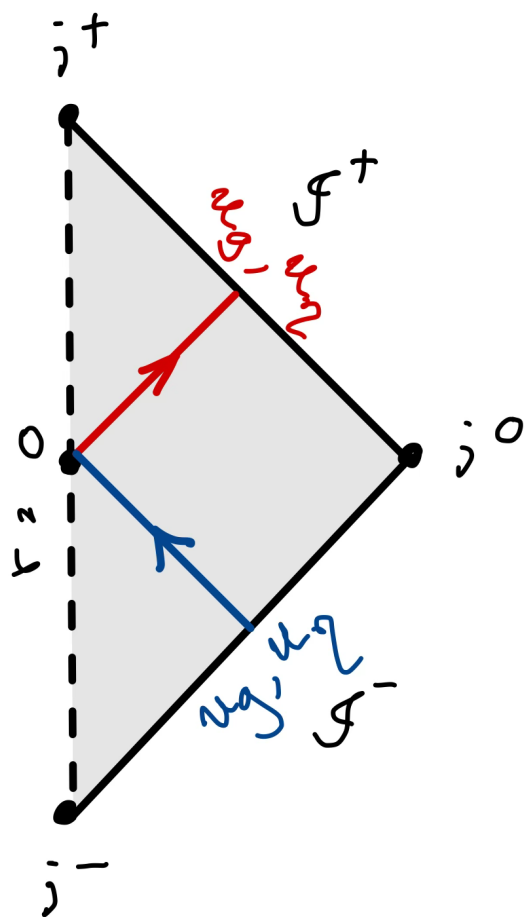
- Not optimal for scattering problems in Schwarzschild spacetimes.

# Subtlety II: Asymptotic Structure

4D Schwarzschild, tortoise coordinates:

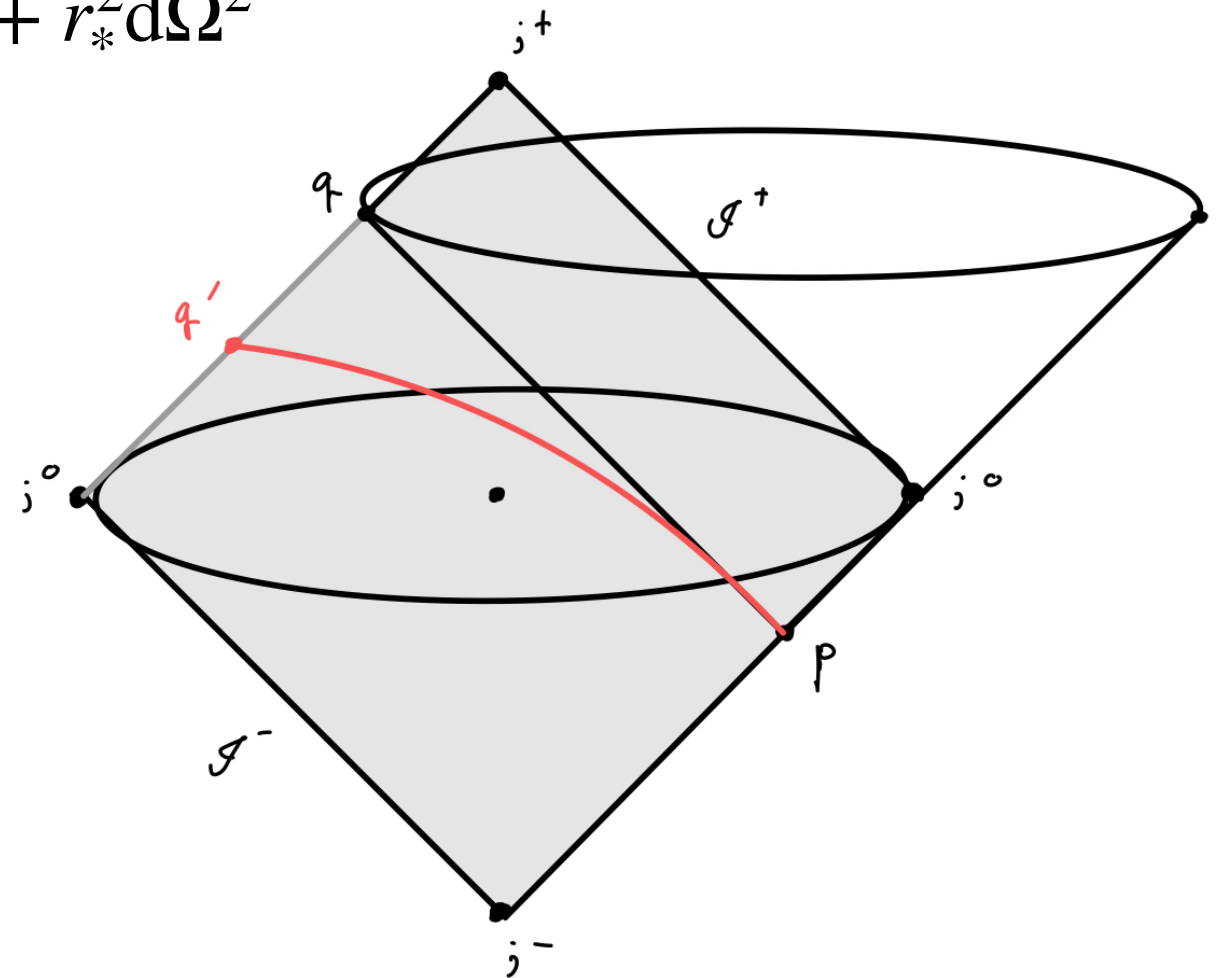
$$ds^2 = - \left( 1 - \frac{r_s}{r} \right) [dt^2 + dr_*^2] + r^2 d\Omega^2$$

$$d\sigma^2 = - dt^2 + dr_*^2 + r_*^2 d\Omega^2$$



$$u_g = u_\eta = t - r_*$$

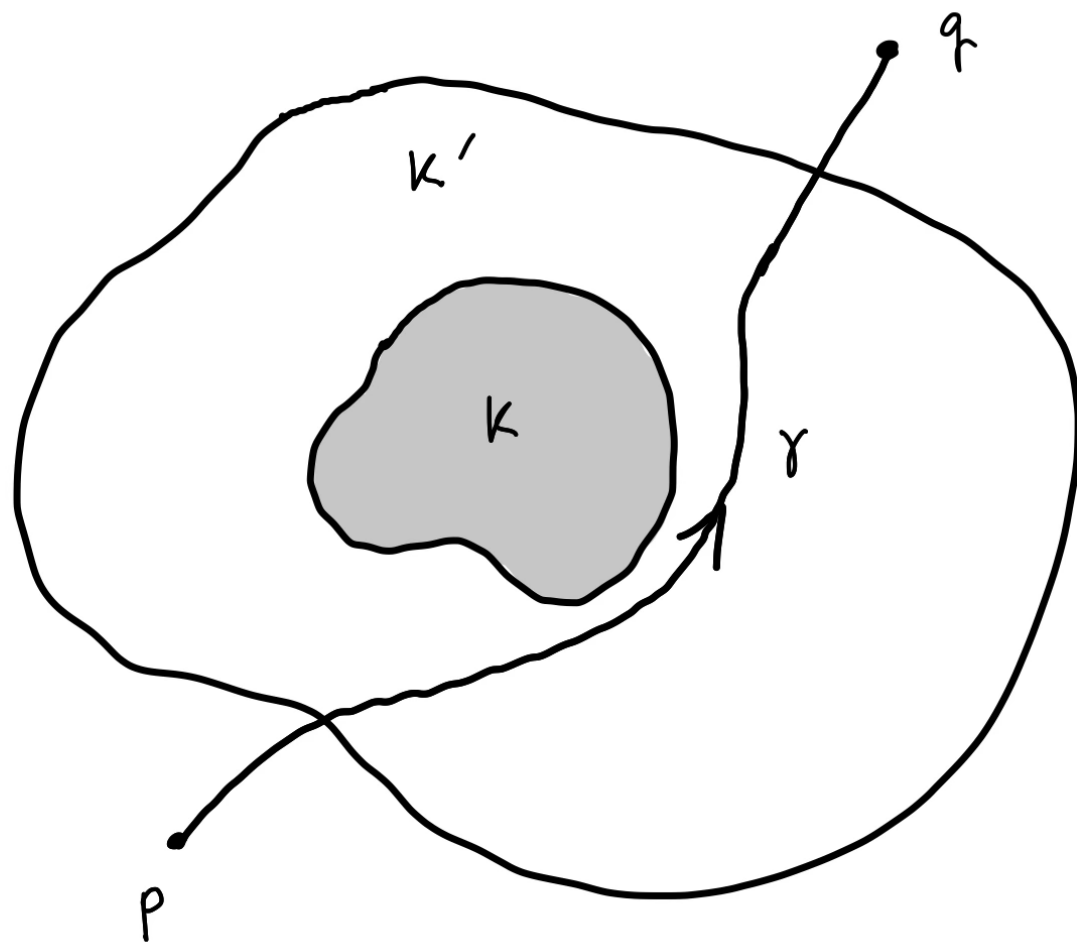
$$v_g = v_\eta = t + r_*$$



$g < \eta$  **fails**—a causal curve that reaches points outside Minkowski can be constructed.

# Gao-Wald Theorem

## Criterion for Promptness



- Gauge-invariant, dimension independent criterion for promptness in GR.
- For asymptotically flat spacetimes, this theory implies that if  $p, q \rightarrow \infty$  then  $\gamma \rightarrow \infty$  as well.
- If ‘time advance’ can be produced in  $K$  then one expects the prompt null geodesic to enter  $K$ .



# FFR Theory

## ‘Drummond&Hathrell’ Problem

- Start from QED in curved spacetimes.

$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} R - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \bar{\psi} (\not{D} + m) \psi \right]$$

Integrating out fermions to 1-loop order on Ricci-flat backgrounds ( $R_{\mu\nu} = R = 0$ ):

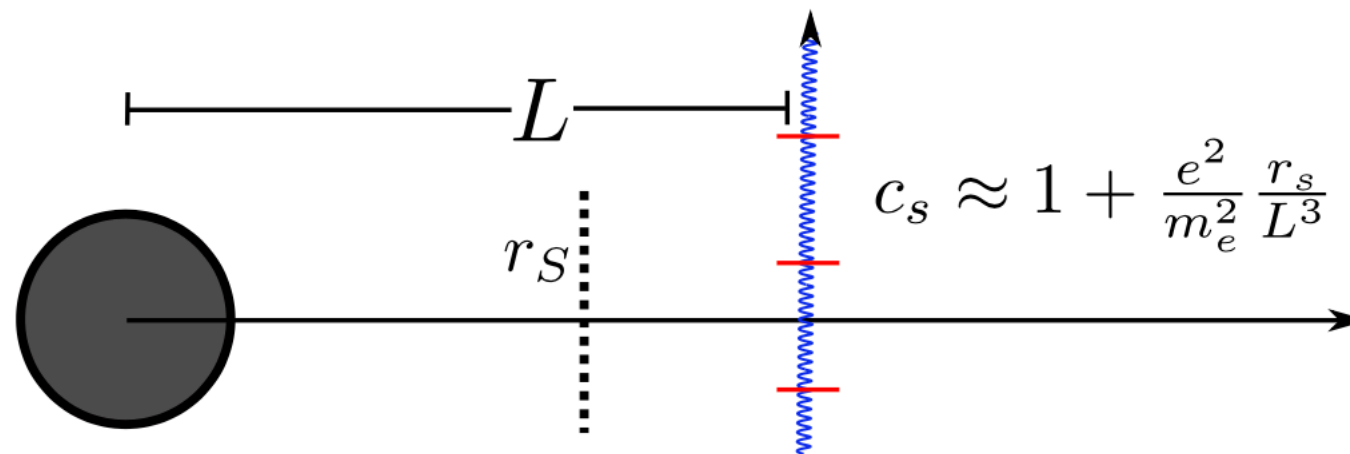
$$S = \int d^4x \sqrt{-g} \left[ \frac{M_{pl}^2}{2} R - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \alpha R_{\mu\nu\rho\sigma} F^{\mu\nu} F^{\rho\sigma} \right]; \quad \alpha \sim \frac{e^2}{m_e^2}$$

Photon Equation of Motion:

$$\nabla_\nu F^{\mu\nu} + \alpha R_{\mu\nu\rho\sigma} F^{\mu\nu} F^{\rho\sigma} = 0$$

# FFR Theory

## Resolution to ‘Drummond&Hathrell’ Problem



- One resolution of this problem is that the effect is small and parametrically smaller than the cutoff of the EFT.
- A superluminal photon will generate a cumulative effect of order

$$\Delta d \sim \delta c_s \times L \sim m_e^{-1} \left( \frac{e^2}{m_e} \frac{r_s^2}{L} \right) \ll m_e^{-1}$$

- For SM electron and solar mass BH we have

$$m_e^{-1} \sim 10^{-13} \text{ m}$$

$$\Delta d \lesssim 10^{-31} \text{ m}$$

# FFR Theory

## Time Delay ( $D > 4$ )

Working in the perturbative regime [ $r_0 \gg r_s$ ,  $r_0 \gg (\alpha r_s^{D-3})^{\frac{1}{D-1}}$ ]

- The usual gravitational time delay has no divergences

$$\Delta t_{\text{Sch}} \propto \frac{1}{r_0^{D-4}}$$

- With *FFR* included, calculation shows that there is a region where time advance is possible— $r_0 \sim \sqrt{\alpha}$ .
- Expectation: the asymptotic structure in this theory should be modified.