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# Flavour Deconstructing the Composite Higgs

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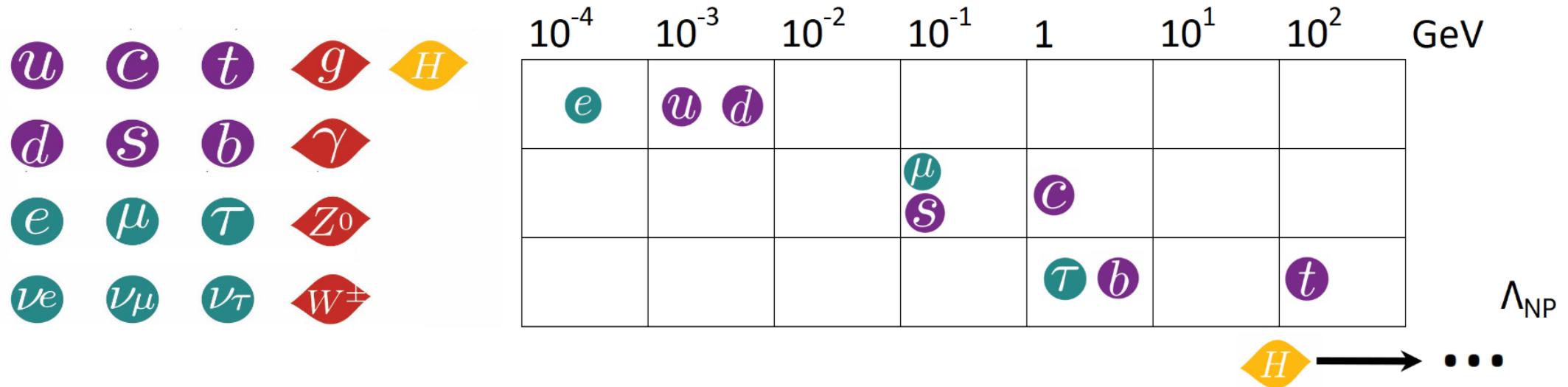
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# Higgs Hierarchy Problem

- Heavy NP (scale  $\Lambda$ ) : corrections to Higgs mass very far from measured value;
- Precise cancellation between bare mass and corrections (in principle unrelated!)
- Fine tuning ( $\Delta$ ) = ratio between corrections and observed mass;
- Acceptable tuning ( $O(1\%)$ ) if  $\Lambda \sim \text{TeV}$ .

$$H \quad -y_t \quad \text{---} \quad \text{---} \quad H$$
$$\text{---} \quad \text{---} \quad \text{---} \quad E$$
$$\text{---} \quad \text{---} \quad \text{---} \quad y_t$$
$$\delta_{\text{SM}} m_H^2 = \frac{3 y_t^2}{8 \pi^2} \Lambda_{\text{SM}}^2$$
$$\Delta \geq \frac{\delta_{\text{SM}} m_H^2}{m_H^2} = \frac{3 y_t^2}{8 \pi^2} \left( \frac{\Lambda_{\text{SM}}}{m_H} \right)^2 \simeq \left( \frac{\Lambda_{\text{SM}}}{450 \text{ GeV}} \right)^2$$

# Flavour puzzle



- Strongly hierarchical Yukawa sector (fermion masses and mixings) : Flavour Puzzle.
- Possible solution: introducing New Physics (NP) at high energy scales.
- Consequences on the Higgs hierarchy problem.

# Hints from the Standard Model

$$\mathcal{L}_{\text{SM}} = \mathcal{L}_{\text{gauge}} + \mathcal{L}_{\text{Higgs}}$$

- Gauge sector does not distinguish fermion flavour;
- The Higgs sector breaks this symmetry, introducing a hierarchy among families;
- Third family looks special!

$$\frac{m_t}{m_u} \approx \frac{173 \text{ GeV}}{2 \text{ MeV}} \approx 10^5$$

$$V_{\text{CKM}} \approx \begin{pmatrix} 1 & 0.2 & 0.004 \\ 0.2 & 1 & 0.04 \\ 0.009 & 0.04 & 1 \end{pmatrix}$$

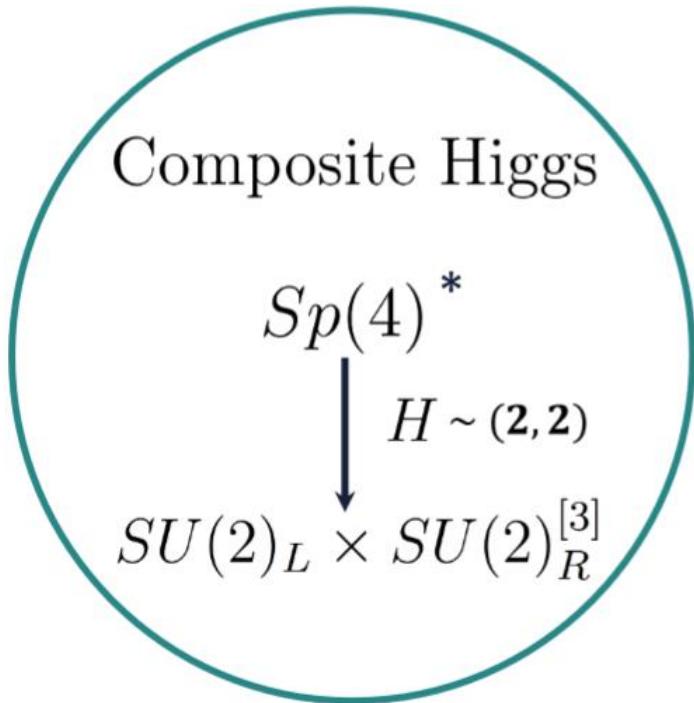


Maybe the SM is only flavour universal at low energies!

At high E, G = G<sub>12</sub> × G<sub>3</sub>?

Heavy NP coupled to 3rd gen?

# Non-universal dynamics and composite Higgs



Non-universal dynamics + composite Higgs

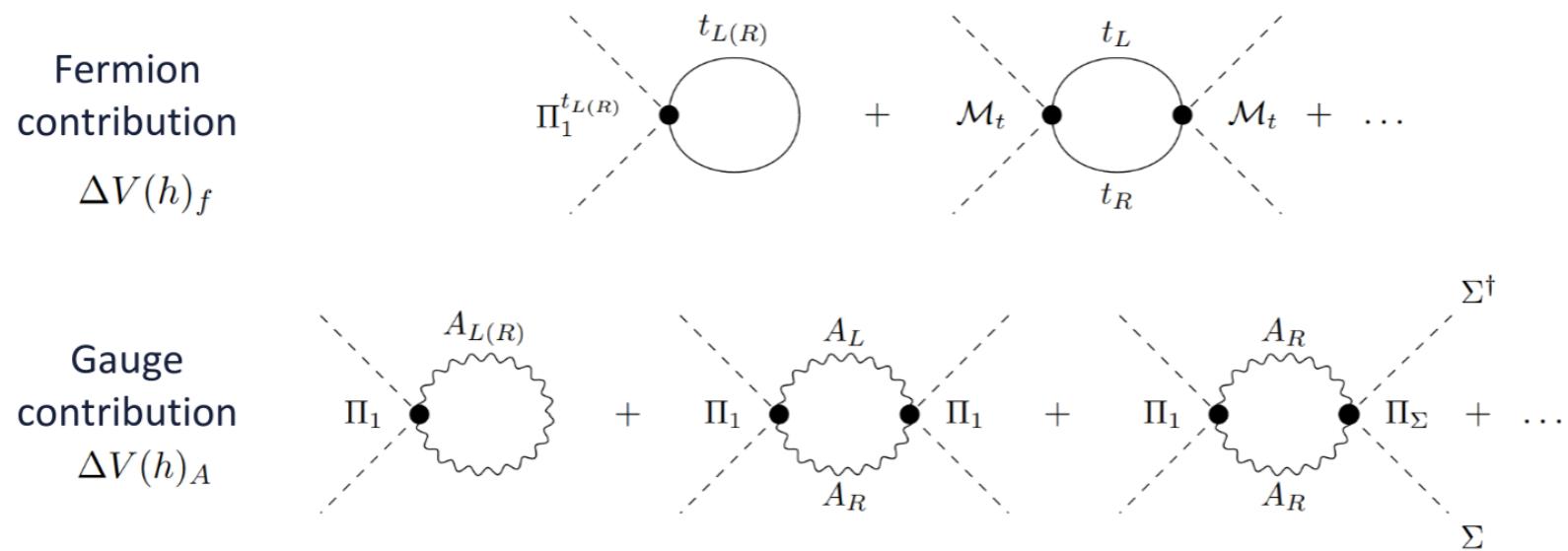


Yukawa hierarchies + stable m<sub>H</sub>

- Higgs : pNGB from global Sp(4) breaking .
- Higgs only couples directly to 3rd gen.
- Shift symmetry: no Higgs mass at LO.
- Potential generated at one-loop. Reduced tuning

# Higgs potential

- Sp(4) explicitly broken by Gauge interactions and mixing w/ composite sector;
- Radiatively generated potential



# Conclusions and outlook

- TeV scale NP coupled to 3rd generation is compatible with current exp. bounds.
- Higgs compositeness + non-universality → Predictive BSM model
- Flavour non-universality lowers the scale of NP

Thank you for your attention!

# Backup slides

Elementary fields		$U(1)_{B-L}^{[3]}$	$U(1)_Y^{[12]}$	$SU(2)_L$	$SU(2)_R^{[3]}$
chiral light quarks	$q_L^{[12]}$	0	1/6	<b>2</b>	<b>1</b>
	$u_R^{[12]}$	0	2/3	<b>1</b>	<b>1</b>
	$d_R^{[12]}$	0	-1/3	<b>1</b>	<b>1</b>
chiral 3 <sup>rd</sup> gen. quarks	$q_L^{[3]}$	1/6	0	<b>2</b>	<b>1</b>
	$q_R^{[3]}$	1/6	0	<b>1</b>	<b>2</b>
vector-like quarks	$F_L^q$	1/6	0	<b>2</b>	<b>1</b>
	$F_R^q$	0	1/6	<b>1</b>	<b>2</b>
scalar link fields	$\Sigma_R$	0	1/2	<b>1</b>	<b>2</b>
	$\Omega_q$	-1/6	1/6	<b>1</b>	<b>1</b>
	$\Omega_\ell$	1/2	-1/2	<b>1</b>	<b>1</b>

# Higgs potential

- Potential is periodic in  $h$ , comprising 2 independent trig. functions.

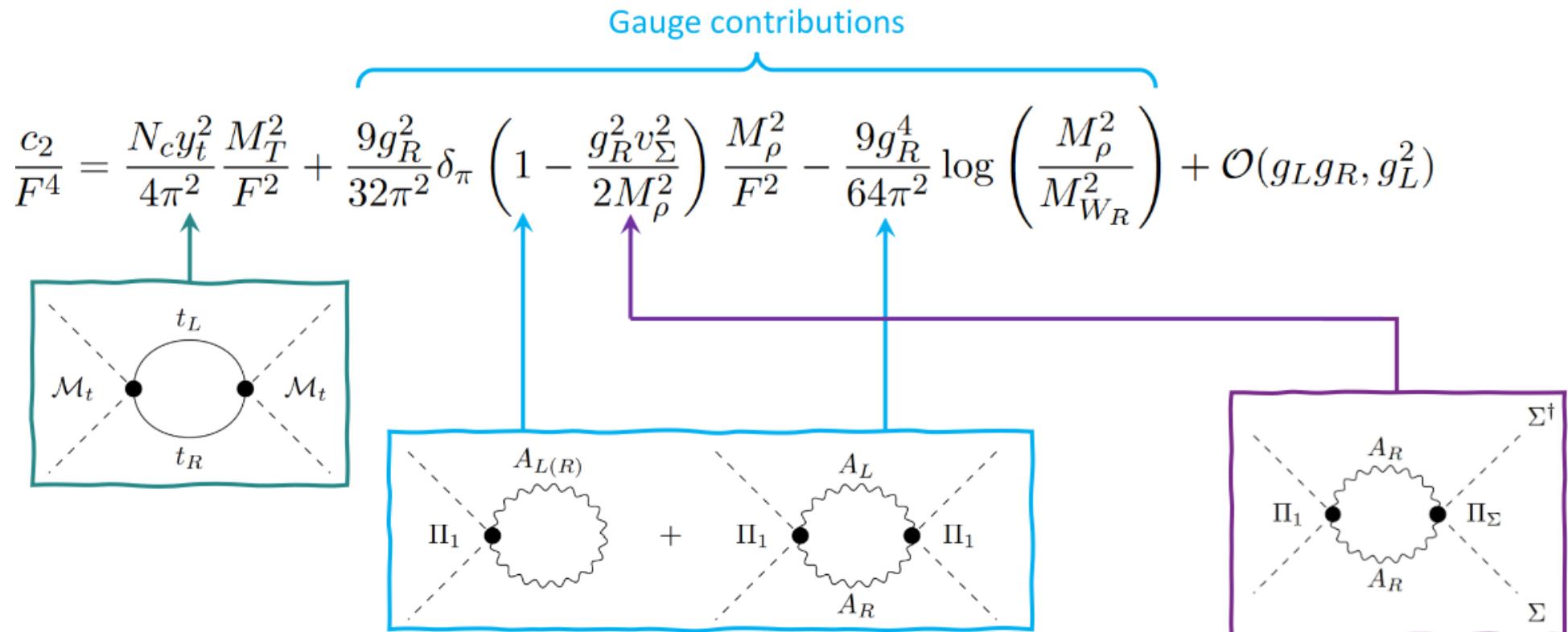
$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$

- Tuning requirements from matching to SM.

$$\frac{c_1}{F^4} \Big|_{\text{phys.}} = \frac{m_h^2}{F^2} \longrightarrow \text{Tuning} \left( m_h^2/F^2 \lesssim 0.03 \right) \quad \quad \quad \frac{c_2}{F^4} \Big|_{\text{phys.}} = \frac{2m_h^2}{v^2} \approx \frac{1}{2}$$

- Flavour non-universality: freedom to choose suitable SU(2) $R$  coupling.

$$V(h) = \Delta V_f(h) + \Delta V_A(h) \approx c_0 - c_1 \sin^2\left(\frac{h}{2F}\right) + c_2 \sin^4\left(\frac{h}{2F}\right)$$



# Higgs potential: quadratic term

- Tuning only stems from  $c_1$
- To stabilize Higgs potential, we need:
  - Introduction of L-R symmetry.
  - Suitable choice of  $g_R$  to suppress top-partner contribution

$$\left. \frac{c_1}{F^4} \right|_{\text{phys.}} = \frac{m_h^2}{F^2} \longrightarrow \text{Tuning } (m_h^2/F^2 \lesssim 0.03)$$

$$\frac{c_1}{F^4} = \frac{N_c}{8\pi^2} \underbrace{\left[ (\lambda_R^t)^2 \kappa_R^t - (\lambda_L^t)^2 \kappa_L^t \right]}_{\text{L-R symmetry}} \frac{M_f^2}{F^2} + \boxed{\frac{N_c y_t^2}{4\pi^2} \frac{M_T^2}{F^2}} \underbrace{- \frac{9g_R^2}{32\pi^2} \left( 1 - \frac{g_R^2 v_\Sigma^2}{2M_\rho^2} \right) \frac{M_\rho^2}{F^2}}_{\text{Gauge contribution}} + \mathcal{O}(g_L g_R, g_L^2)$$

# Potenziale di Higgs: termine quartico

$c_2$  naturale (0(1)) dal matching con il MS

$$\frac{c_2}{F^4} \Big|_{\text{phys.}} = \frac{2m_h^2}{v^2} \approx \frac{1}{2}$$

Constraint sul rapporto  $M_T/F$

$$\frac{c_2}{F^4} = \frac{N_c y_t^2}{4\pi^2} \frac{M_T^2}{F^2} + \text{Gauge contributions (suppressed)}$$

↑  
Top partner →  $M_T \approx 2.5F$

# Experimental constraints: super-strong sector

- Effects on VVh and VVhh couplings
- Top partners, TeV scale resonances
- EWPO, S parameter



$$\Lambda_{\text{HC}} \gtrsim 12 \text{ TeV (95\% CL)}$$



$$M_T \gtrsim 1.5 \text{ TeV} \longrightarrow \Lambda_{\text{HC}} \gtrsim 15 \text{ TeV}$$
$$M_\rho \gtrsim 5 \text{ TeV}$$



$$g_{L,R}^2 \frac{v^2}{M_\rho^2} \lesssim 10^{-3}$$

# Experimental constraints: Gauge bosons

- Flavour (e.g.  $B \rightarrow X_s + \gamma$ ) and Z-pole

$$v_\Sigma \gtrsim 3 \text{ TeV}$$

- Drell-Yan from LHC

$$v_\Sigma \gtrsim 2.7 \text{ TeV}$$

- $B_s$  mixing

$$v_\Sigma \gtrsim 2 \text{ TeV}$$

# Breaking pattern

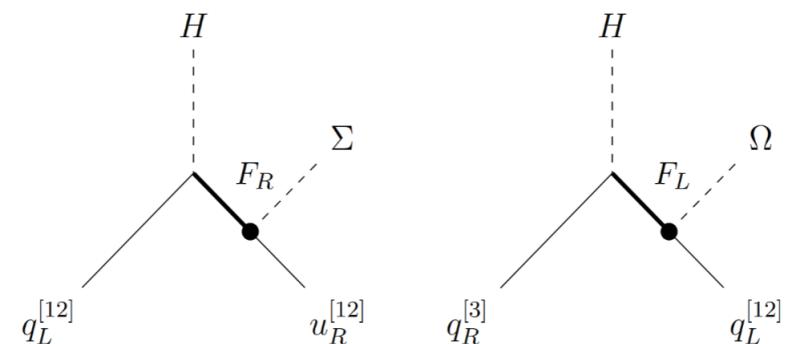
$$SU(3)_c \times SU(2)_L \times SU(2)_R^{[3]} \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$



$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

$$\begin{array}{c} SU(2)_R^{[3]} \times U(1)_{T_R^3}^{[12]} \xrightarrow{\langle \Sigma \rangle} U(1)_{T_R^{[3]}} \\ U(1)_{B-L}^{[3]} \times U(1)_{B-L}^{[12]} \xrightarrow{\langle \Omega \rangle} U(1)_{B-L} \end{array} \longrightarrow \begin{array}{l} \epsilon_R = \frac{\langle \Sigma \rangle}{M_F} \\ \epsilon_\Omega = \frac{\langle \Omega \rangle}{M_F} \end{array} \longrightarrow Y_{u,d,e} \sim \begin{pmatrix} \textcolor{purple}{U(2)^5} & \\ \begin{pmatrix} \epsilon_R & \epsilon_\Omega \\ \epsilon_R \epsilon_\Omega & 1 \end{pmatrix} & \end{pmatrix}$$

- VEV of elementary scalars generate breaking to  $G_{\text{SM}}$  at low energy ;
- Yukawa coupling effectively generated through insertion of heavy vector-like fermions (100 TeV).



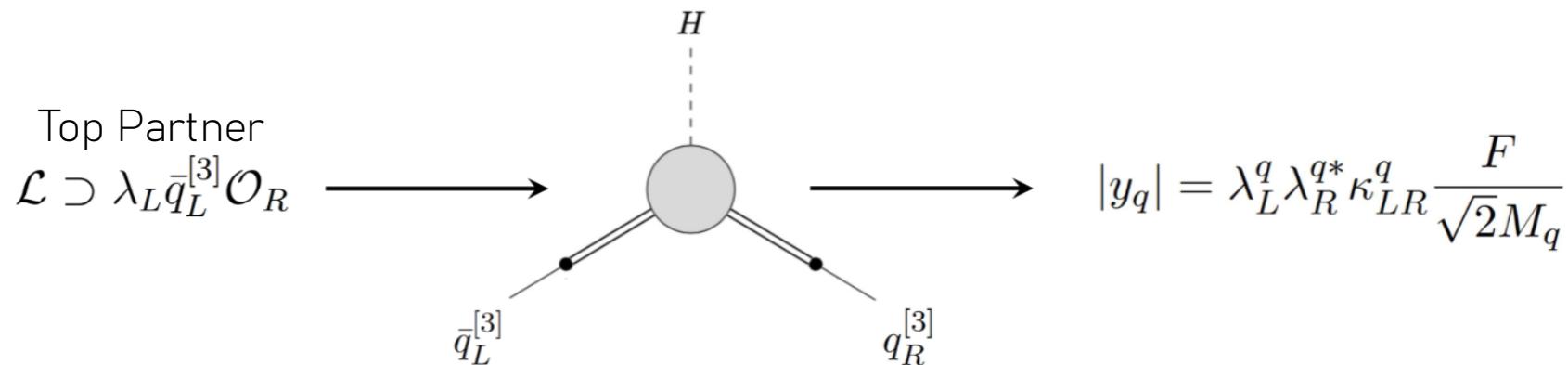
# Breaking pattern

$$SU(3)_c \times Sp(4) \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$

Non-perturbative dynamics

$$\Lambda_{\text{HC}} \downarrow \quad H \sim (\mathbf{2}, \mathbf{2})$$
$$SU(3)_c \times SU(2)_L \times SU(2)_R^{[3]} \times U(1)_{B-L}^{[3]} \times U(1)_Y^{[12]}$$

Partial compositeness: 3rd family fermions mix w/ composite sector.



# Yukawa sector

- Observed Yukawa matrices are obtained by tuning the ratio of the scalar VEV and the VLF mass

$$\epsilon_R = \frac{\langle \Sigma \rangle}{M_F} \quad \epsilon_\Omega = \frac{\langle \Omega \rangle}{M_F}$$

- Tuning requirements on the Higgs potential do not uniquely fix the scale of flavour non universality !

$$Y_{u,d,e} \sim \begin{pmatrix} \textcolor{purple}{\epsilon_R} & \epsilon_\Omega \\ \epsilon_R \epsilon_\Omega & 1 \end{pmatrix}^{\textcolor{purple}{U(2)^5}}$$

$$\epsilon_\Omega = O(|V_{cb}|) = O(10^{-1})$$

$$\epsilon_R = O(m_c/m_t) = O(10^{-2})$$

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# Higgs potential: quadratic term

$$\frac{c_1}{F^4} \supset \frac{N_c y_t^2}{4\pi^2} \frac{M_T^2}{F^2} - \frac{9g_R^2}{32\pi^2} \left(1 - \frac{g_R^2 v_\Sigma^2}{2M_\rho^2}\right) \frac{M_\rho^2}{F^2}$$

- Potential stabilized by appropriate  $g_R$  choice  
(allowed by non-universality)
- Large enough to be of the same order of  $y_t$  contribution
- Small enough not to generate a sign inversion in the Gauge contribution

$$g_{R,3} = O(1) \gg g_{R,12} \approx g_Y^{\text{SM}}$$

$$M_{W_R}^2 = \frac{1}{4} g_R^2 v_\Sigma^2 < \frac{1}{2} M_\rho^2$$

# Summary of model

- $O(1)$   $SU(2)_R$  gauge coupling from third family  • 3% fine tuning
- Light top-partner (2 TeV)  •  $O(1\%)$  modifications to Higgs couplings
- Strong SSB scale  $\Lambda \sim 20$  TeV and  $v_\Sigma \sim 3$  TeV  •  $O(10^{-3})$  modifications to EW sector