



PHASING OUT OF DARKNESS

**From Sterile Neutrino Dark Matter to Neutrino Masses
via Time-Dependent Mixing**

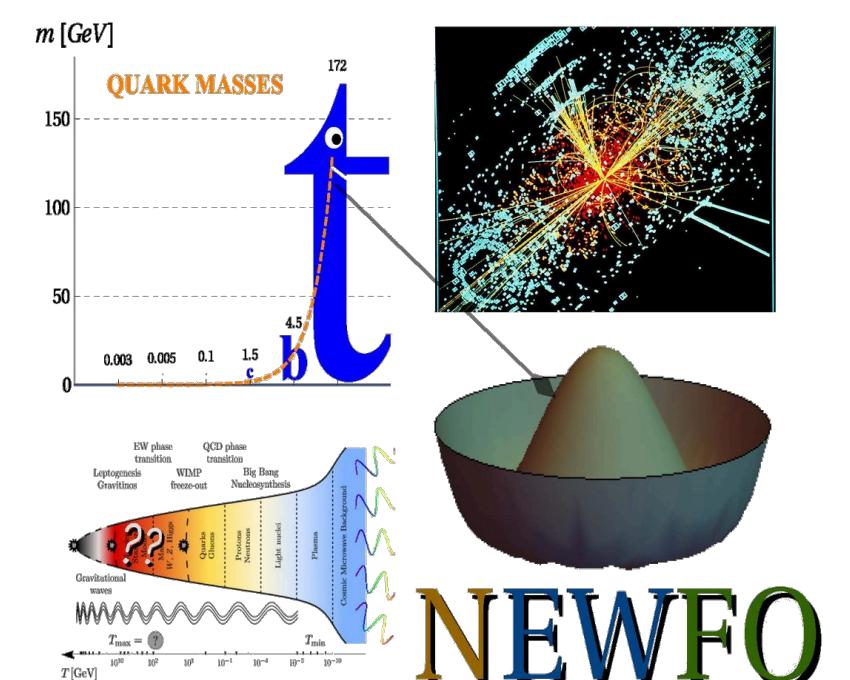
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July 29, 2025

Cargese Summer School

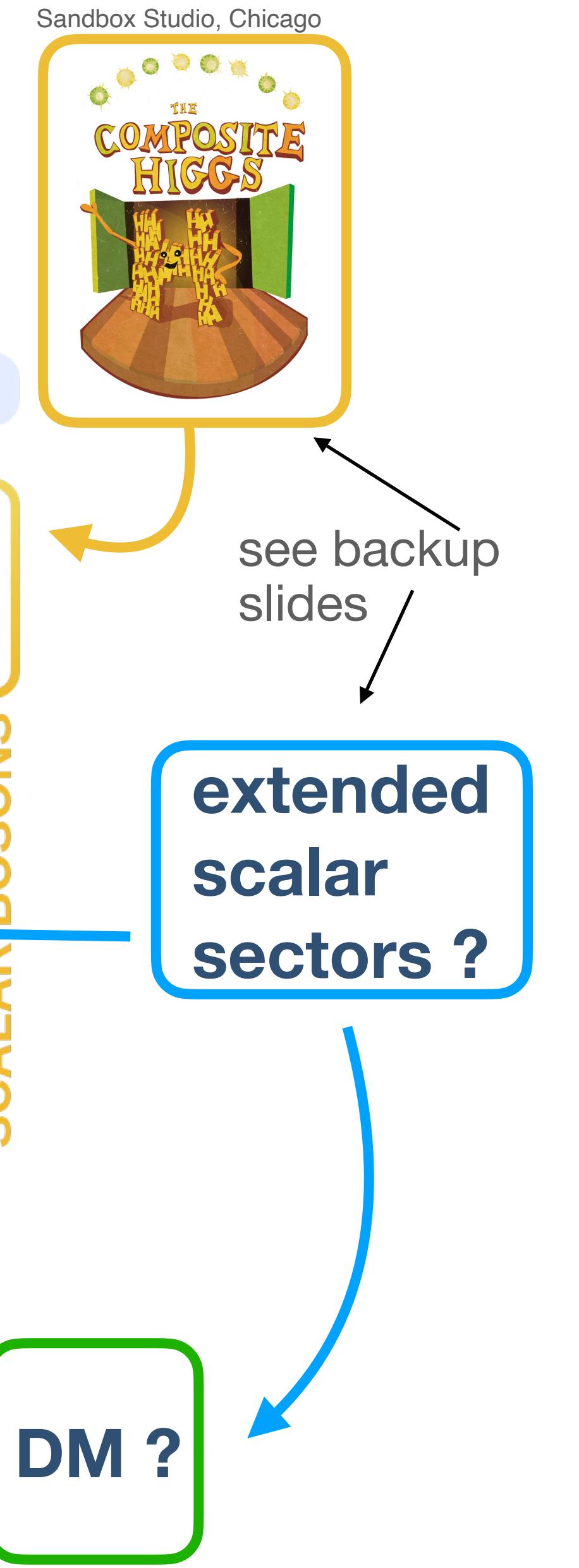
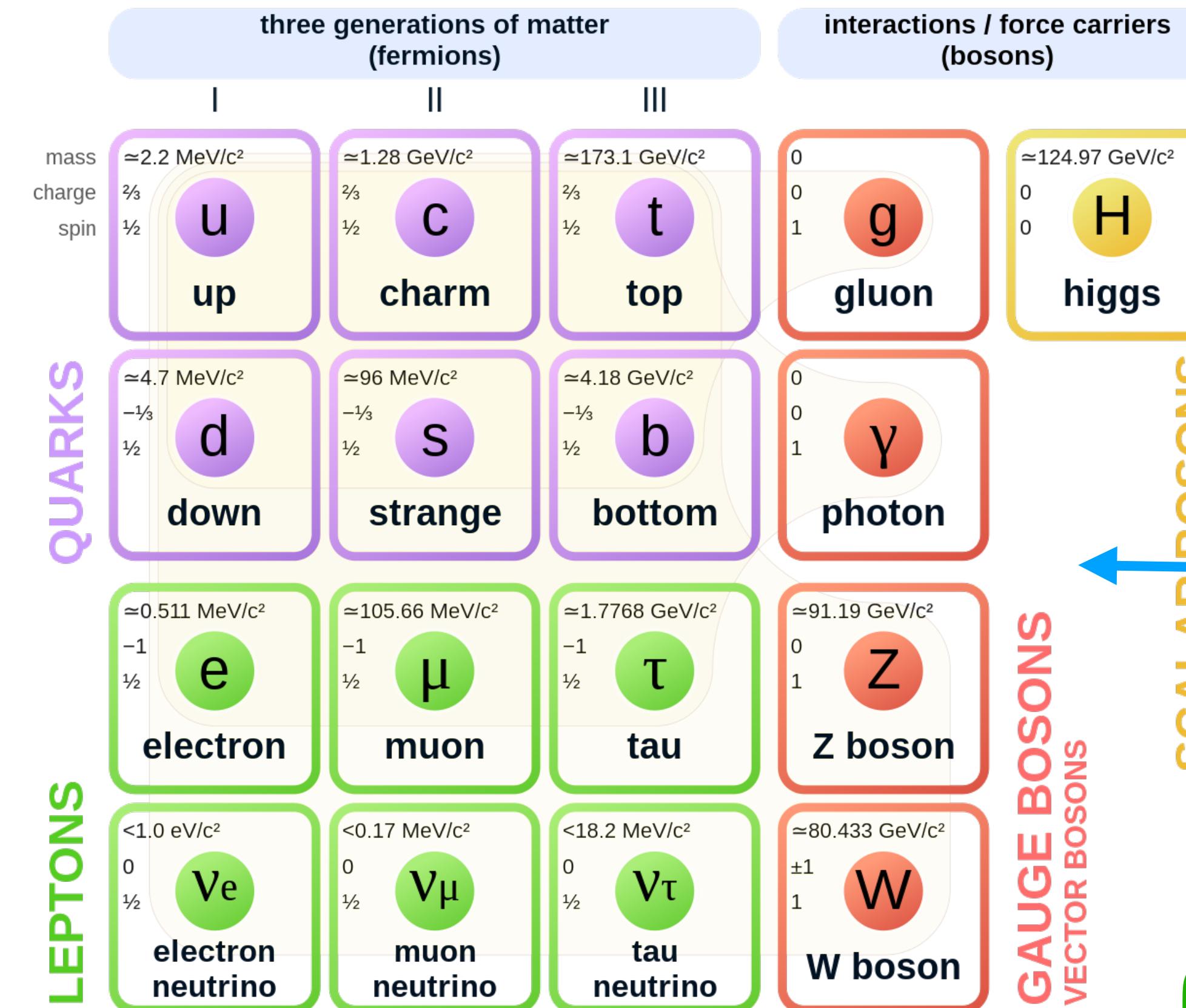


Content...

...of my research:

- Late Phase Transition to Reconcile Sterile Neutrino DM with X-Ray Bounds
- Majoron DM production via non-minimal coupling to gravity (ongoing)
- Characterising New Resonances at LHC in EFT (ongoing)
- Decreasing Tuning in CH Models with unexplored LHC signature

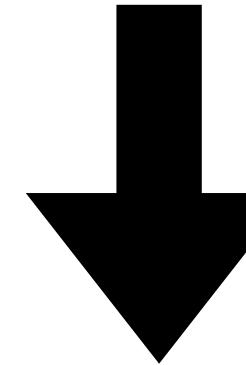
...of the Standard Model:



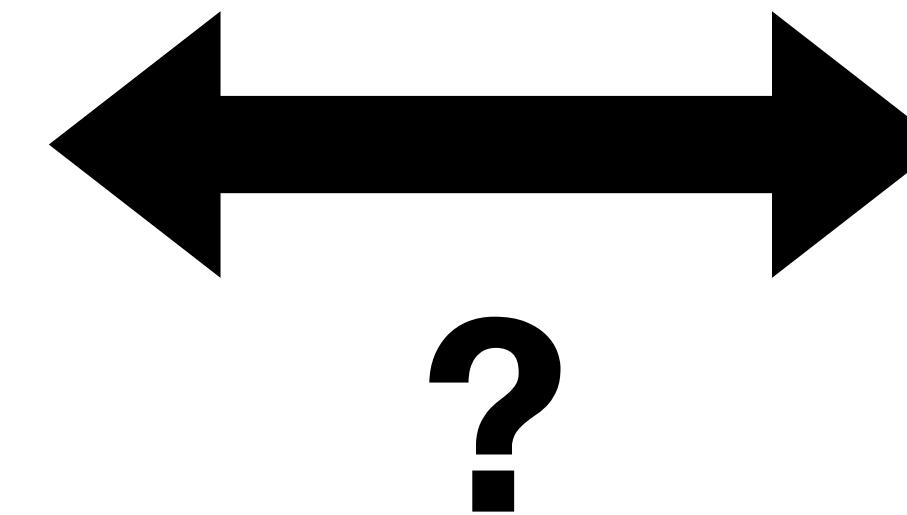


Two open questions of the SM

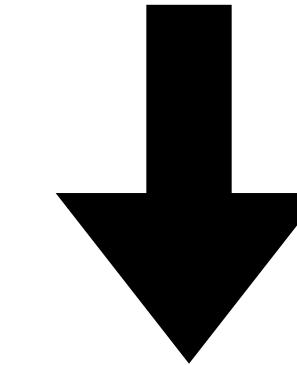
Why are neutrinos massive?



Seesaw Mechanism
(via heavy ν_R)



What is the nature of dark matter?



Sterile Neutrinos
(keV masses)

+ many more ideas...



can they be one and the same?

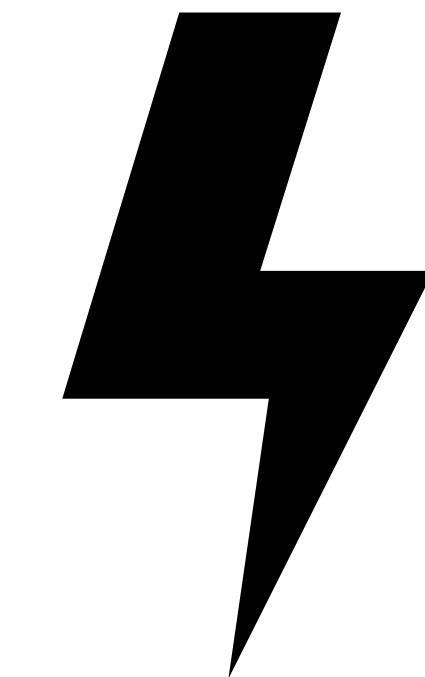


Conflicting bounds

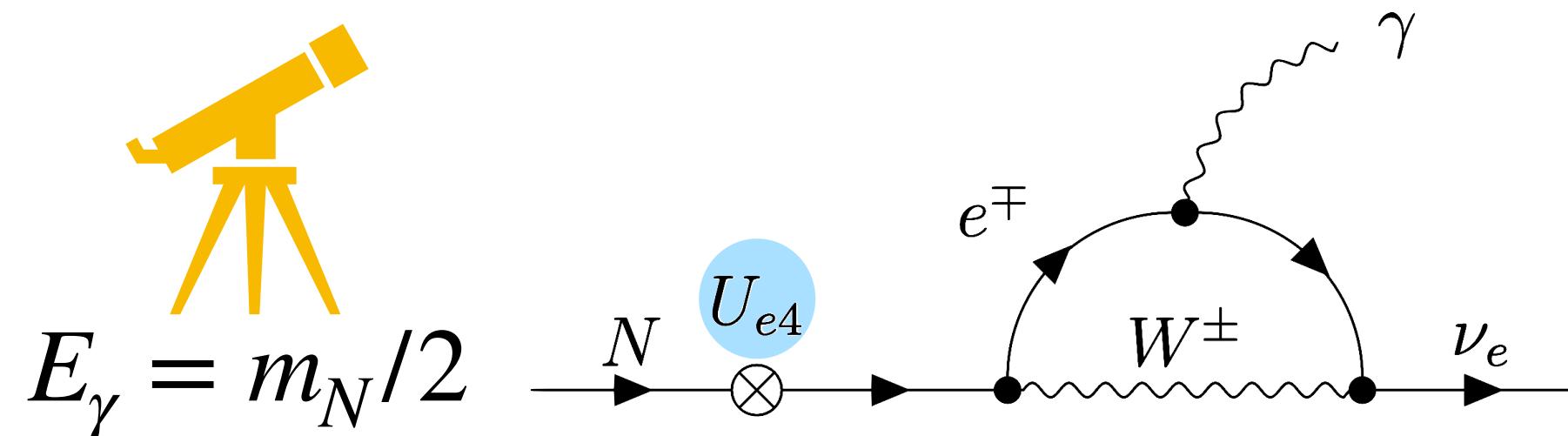
- neutrino masses
constrained by oscillation measurements
→ lower bound on mixing (in type-I seesaw)



$$m_\nu \sim \frac{M_D^2}{M_R} \sim U_{\alpha I}^2 M_R$$



- radiative DM decay $N \rightarrow \nu \gamma$
constrained by X-ray observations
→ upper bound on mixing (if 100% DM)



How to reconcile the bounds (here)

- neutrino masses constrained by oscillation measurements

→ **lower bound** on mixing



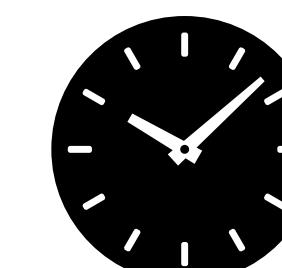
- radiative DM decay $N \rightarrow \nu \gamma$ constrained by X-ray observations

→ **upper bound** on mixing

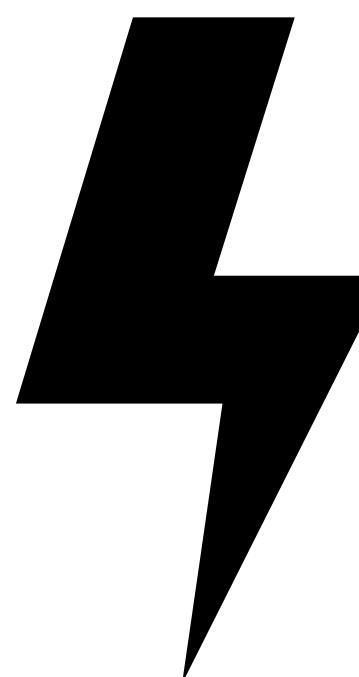


from DM dense objects
(10^4 to 10^{10} ly away)

add time dependence
to mixing!



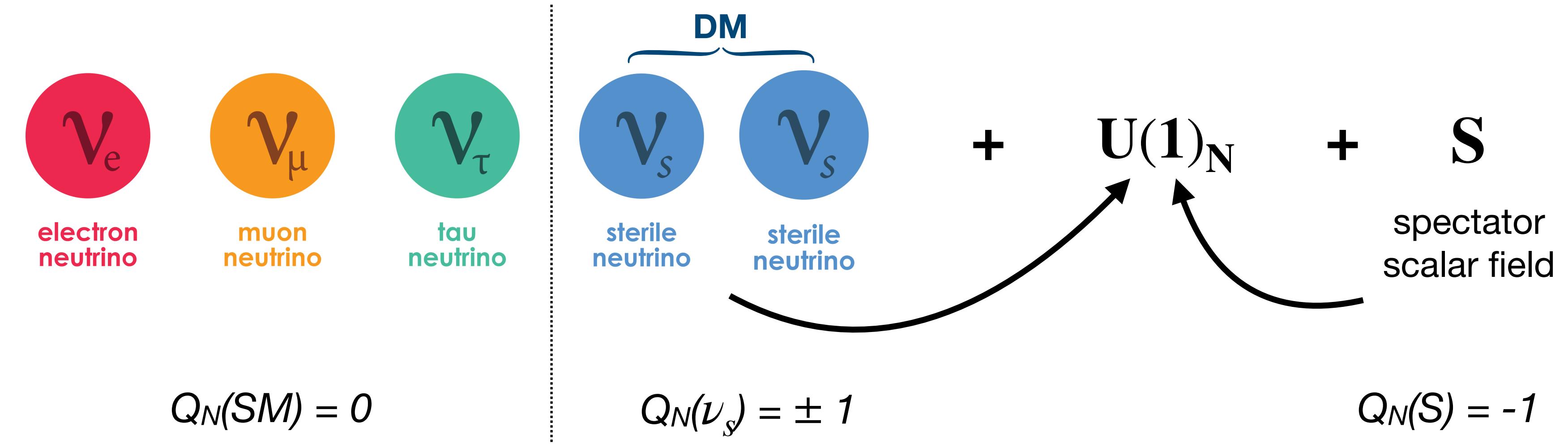
$\epsilon(t)$



→ **reconciles
keV sterile neutrinos
as DM candidate
& as type-I seesaw neutrinos**



Model



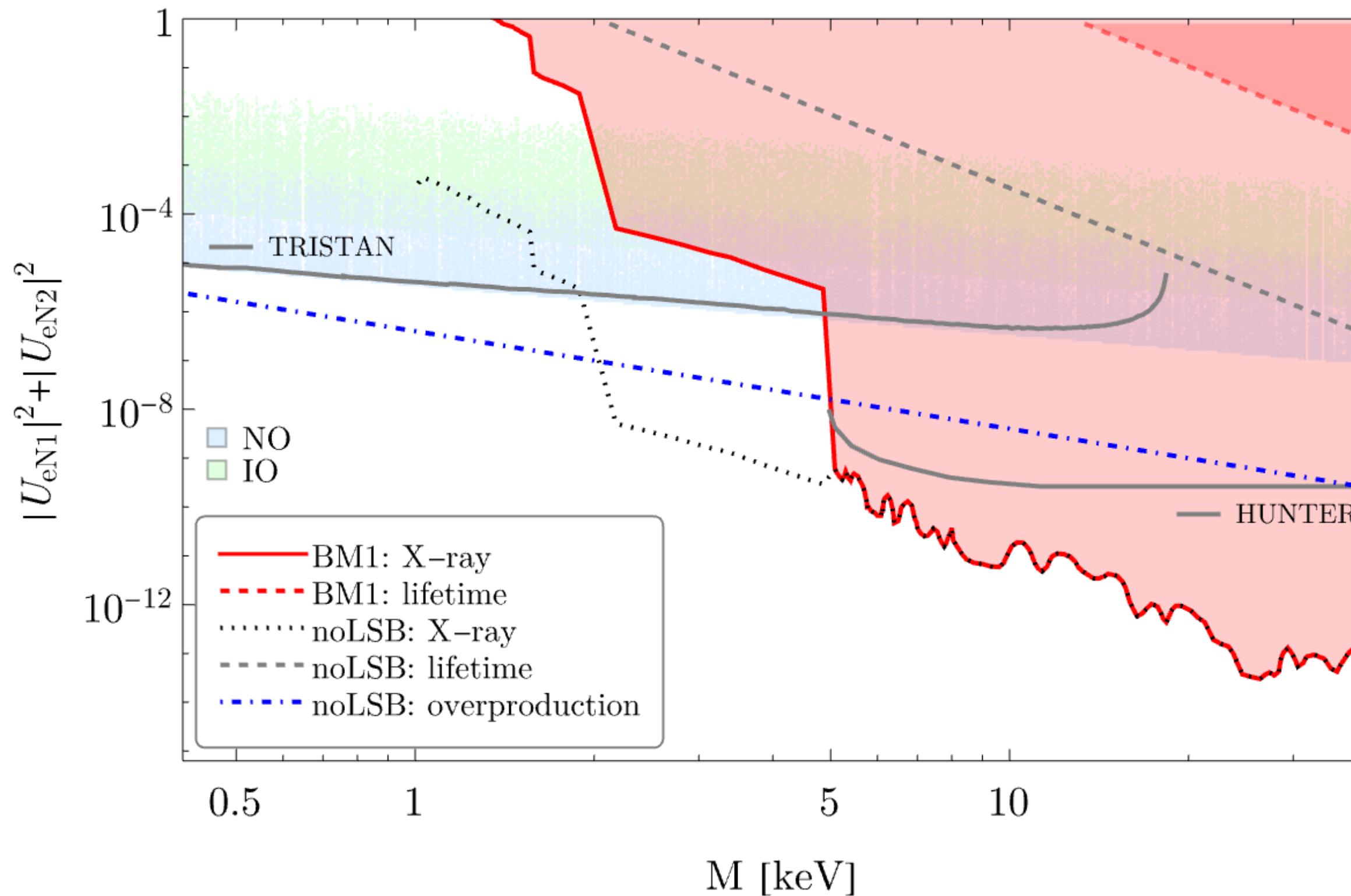
- in the recent universe, the potential of the spectator scalar field becomes tachyonic \Rightarrow field starts rolling \Rightarrow expectation value $\bar{S}(t > t_{SB}) \neq 0$
- neutrino masses are induced!



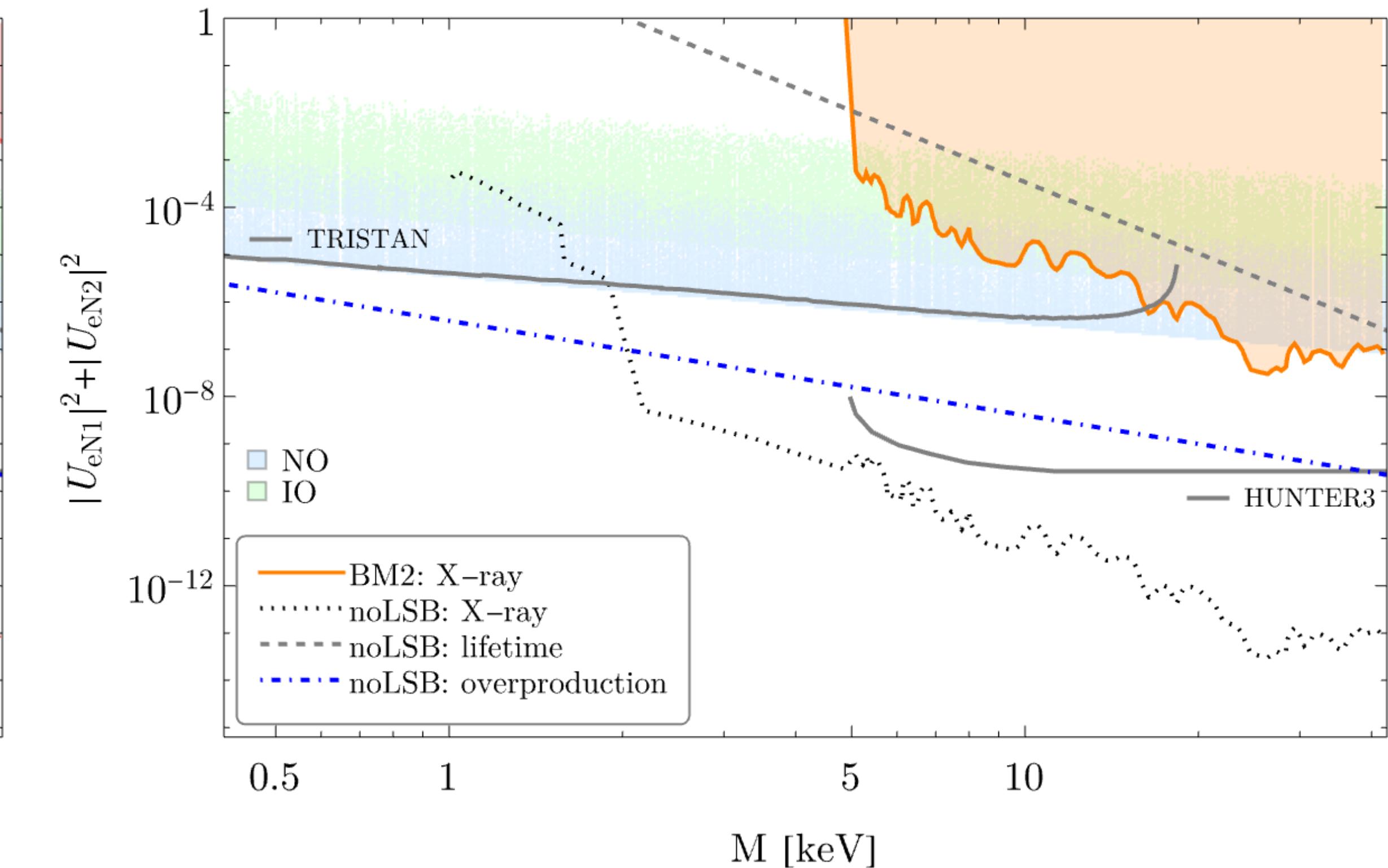
→aligns with
DESI results!

X-ray Constraints

BM 1



BM 2



BM1: 0.4 - 5 keV DM

TRISTAN

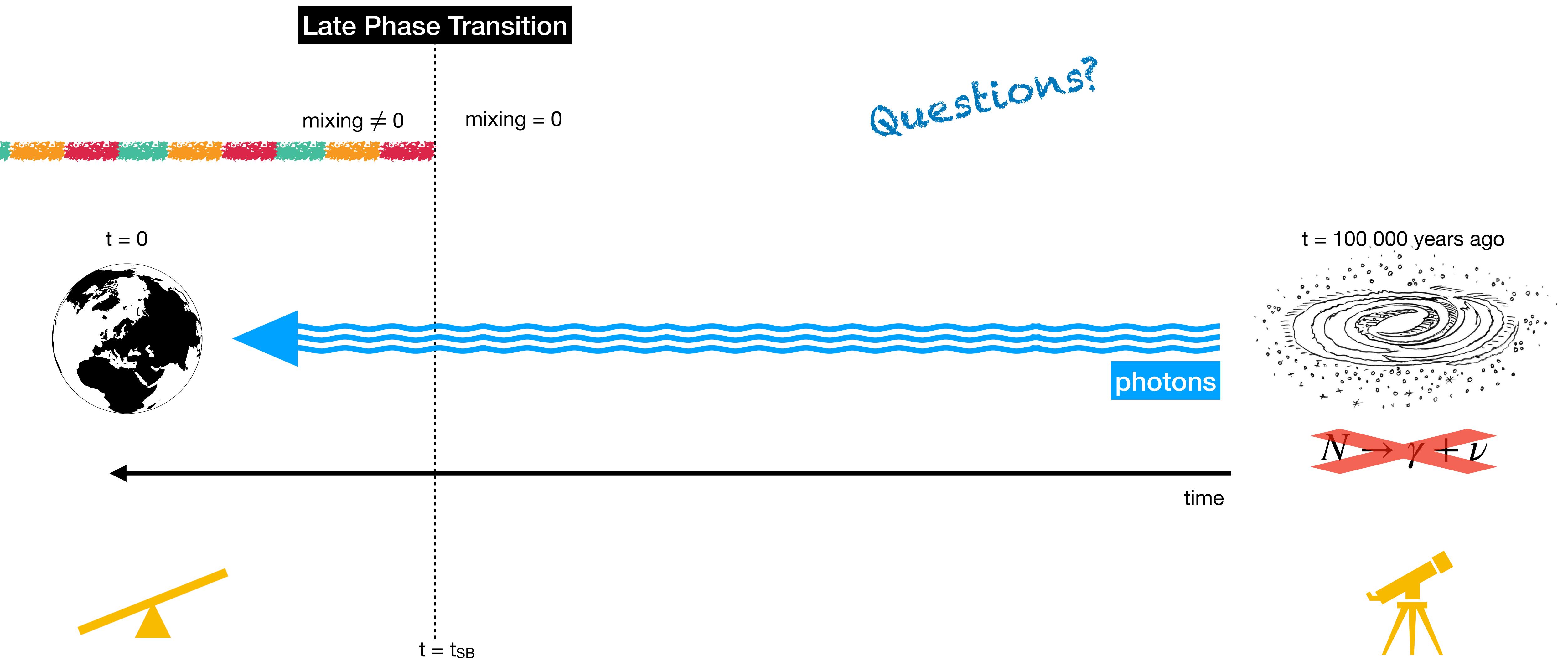


BM2: 0.4 - 22 keV DM

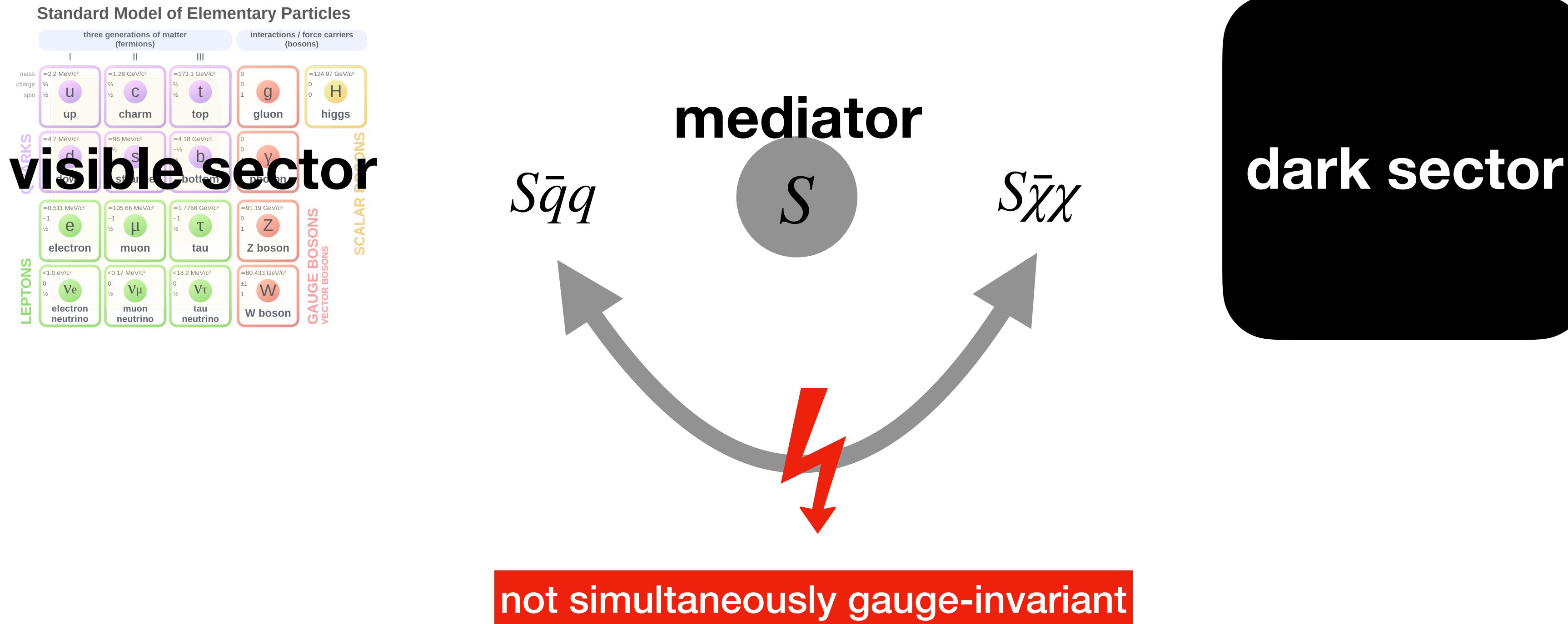
TRISTAN + HUNTER3



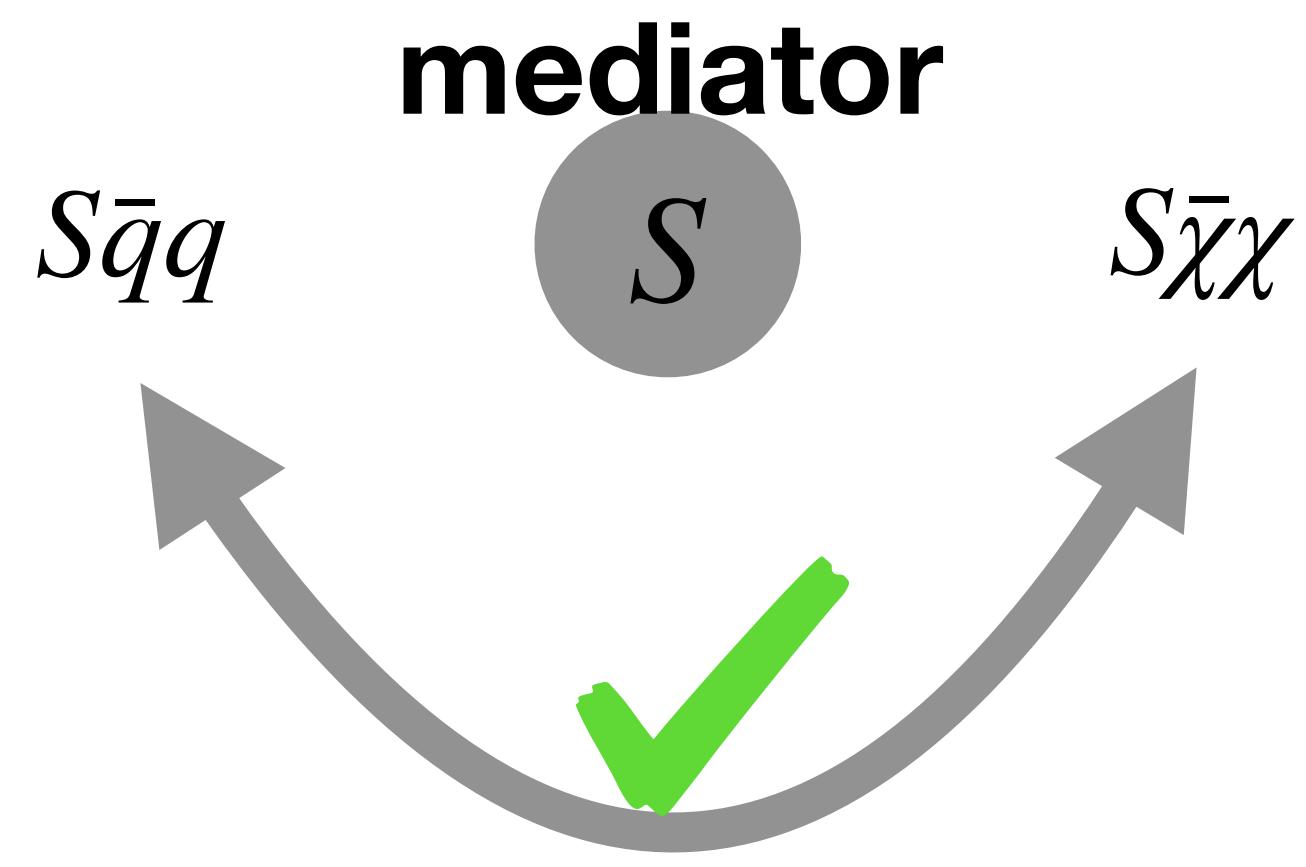
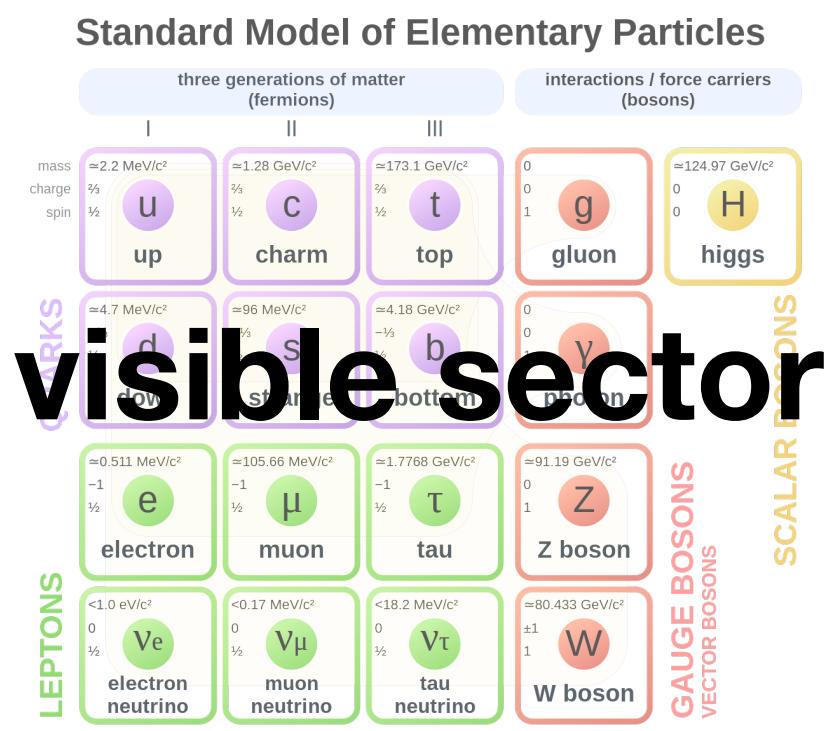
How to reconcile the bounds



Simplified Models



extended DMEFT



up to dim-5:

$$\begin{aligned} \mathcal{L} \supset & \frac{1}{2} \sum_{\phi=S_1, S_2, h} \partial_\mu \phi \partial^\mu \phi - \mathcal{O}_5^\lambda + \frac{v^2}{4} \text{Tr} \left[(D_\mu \Sigma)^\dagger (D^\mu \Sigma) \right] \mathcal{O}_3^\kappa \\ & + i \frac{v^2}{4} \text{Tr} \left[\Sigma^\dagger (D^\mu \Sigma) \sigma^3 \right] (\partial_\mu h \mathcal{O}_2^h + \partial_\mu S_1 \mathcal{O}_2^{s1} + \partial_\mu S_2 \mathcal{O}_2^{s2}) \\ & - \frac{v}{\sqrt{2}} \left((\overline{u}_{i,L} \ \overline{d}_{i,L}) \Sigma \begin{pmatrix} Y_{ij}^u u_{j,R} \\ Y_{ij}^d d_{j,R} \end{pmatrix} \mathcal{O}_2^{cq} + (\overline{\nu}_{i,L} \ \overline{\ell}_{i,L}) \Sigma \frac{1-\sigma_3}{2} Y_{ij}^\ell \begin{pmatrix} \nu_{j,R} \\ \ell_{j,R} \end{pmatrix} \mathcal{O}_2^{c\ell} + \text{h.c.} \right) \\ & - \sum_{\phi} \frac{\phi}{16\pi^2} \left[g'^2 c_B^\phi B^{\mu\nu} B_{\mu\nu} + g^2 c_W^\phi W^{I\mu\nu} W_{\mu\nu}^I + g_s^2 c_G^\phi G^{a\mu\nu} G_{\mu\nu}^a \right] \\ & - \sum_{\phi} \frac{\phi}{16\pi^2} \left[g'^2 \tilde{c}_B^\phi B^{\mu\nu} \tilde{B}_{\mu\nu} + g^2 \tilde{c}_W^\phi W^{I\mu\nu} \widetilde{W}_{\mu\nu}^I + g_s^2 \tilde{c}_G^\phi G^{a\mu\nu} \widetilde{G}_{\mu\nu}^a \right], \end{aligned} \quad (2.1)$$

2411.05914 Arcadi, Cabo-Almeida,
Fabian, Goertz



- ✓ gauge-invariant, consistent
- ✓ richer than simplified models
(interplay between different operators can become apparent)
- ✓ more hope for LHC signals compared to standard DM EFT
(light mediator)

preliminary

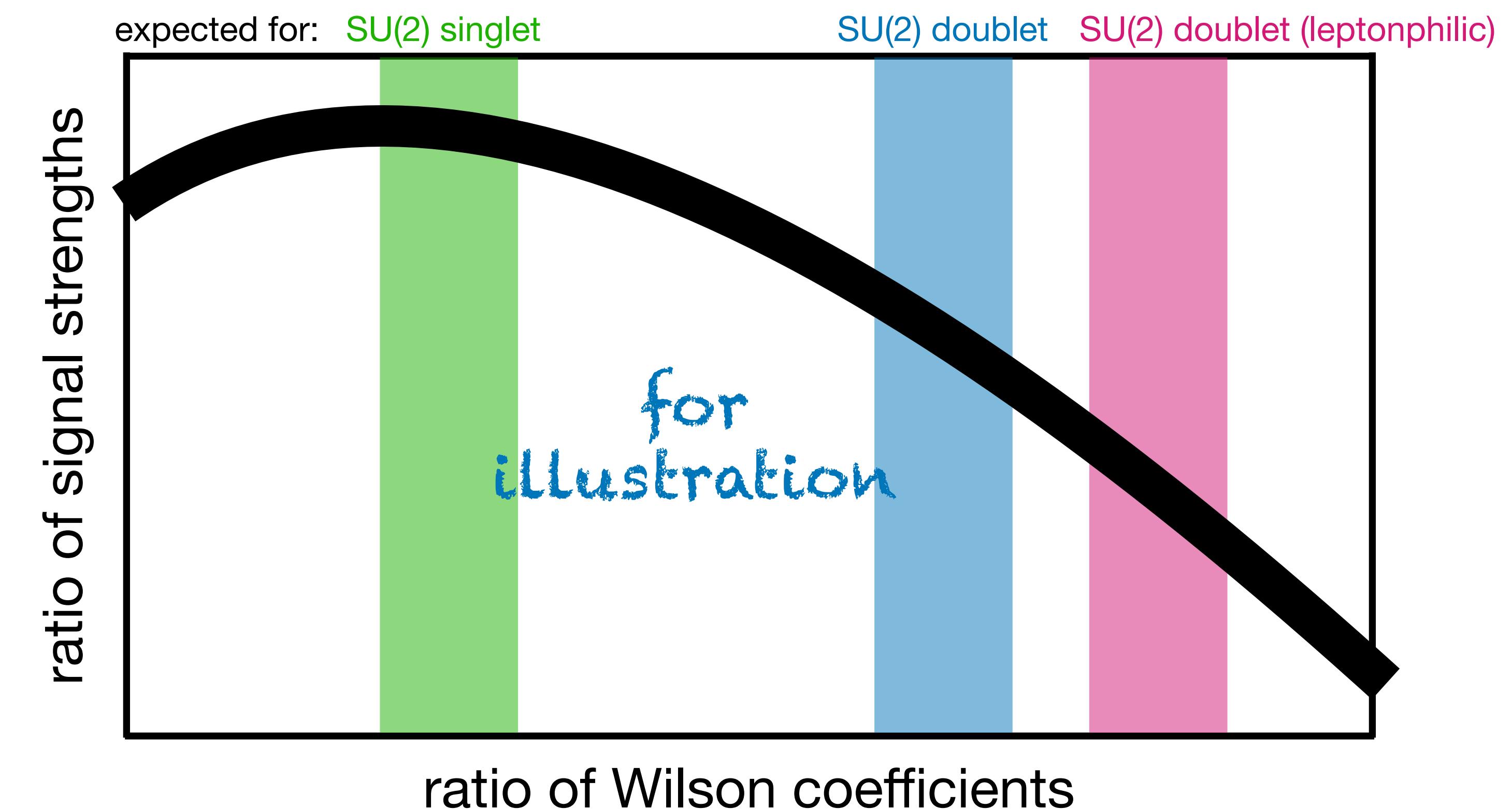
Characterising New Resonances At LHC in EFT

- focus on **one new scalar degree of freedom** in a mass range **accessible at the LHC**
- input: generalised HEFT + dimensional analysis
- output sth like:
(in ideal case)

example: $S\bar{q}q$ operator

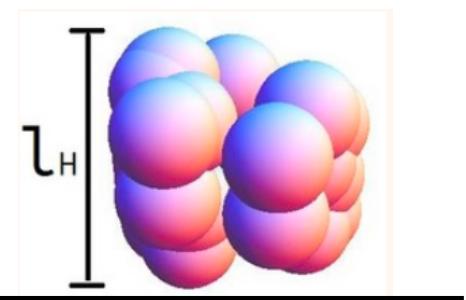
$SU(2)$ doublet:
allowed at dim-4

$SU(2)$ singlet:
not gauge-inv, needs Higgs insertion
→ additional suppression

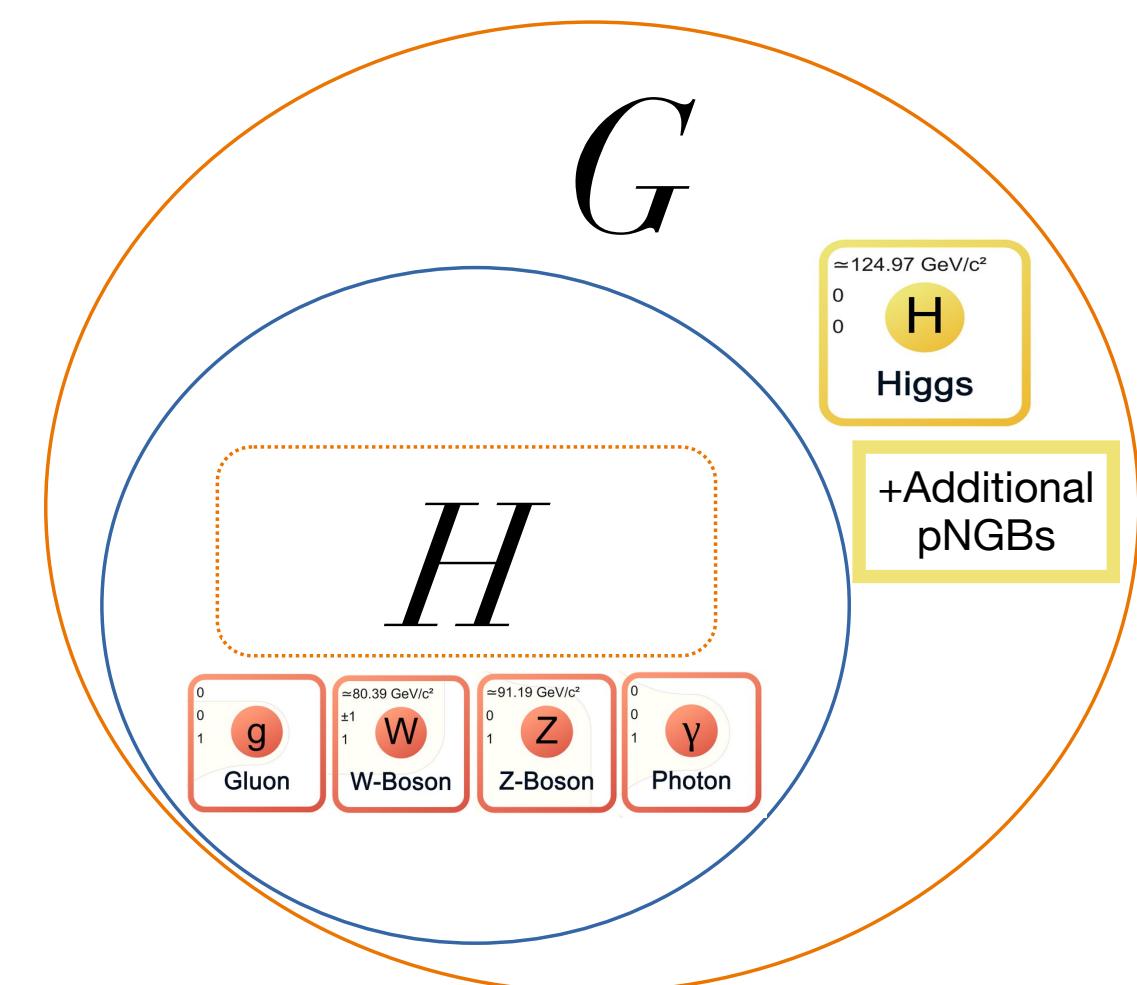


Composite Higgs

Corrections @ $\mathcal{O}(\text{TeV})$



Analogue:
Pions in QCD



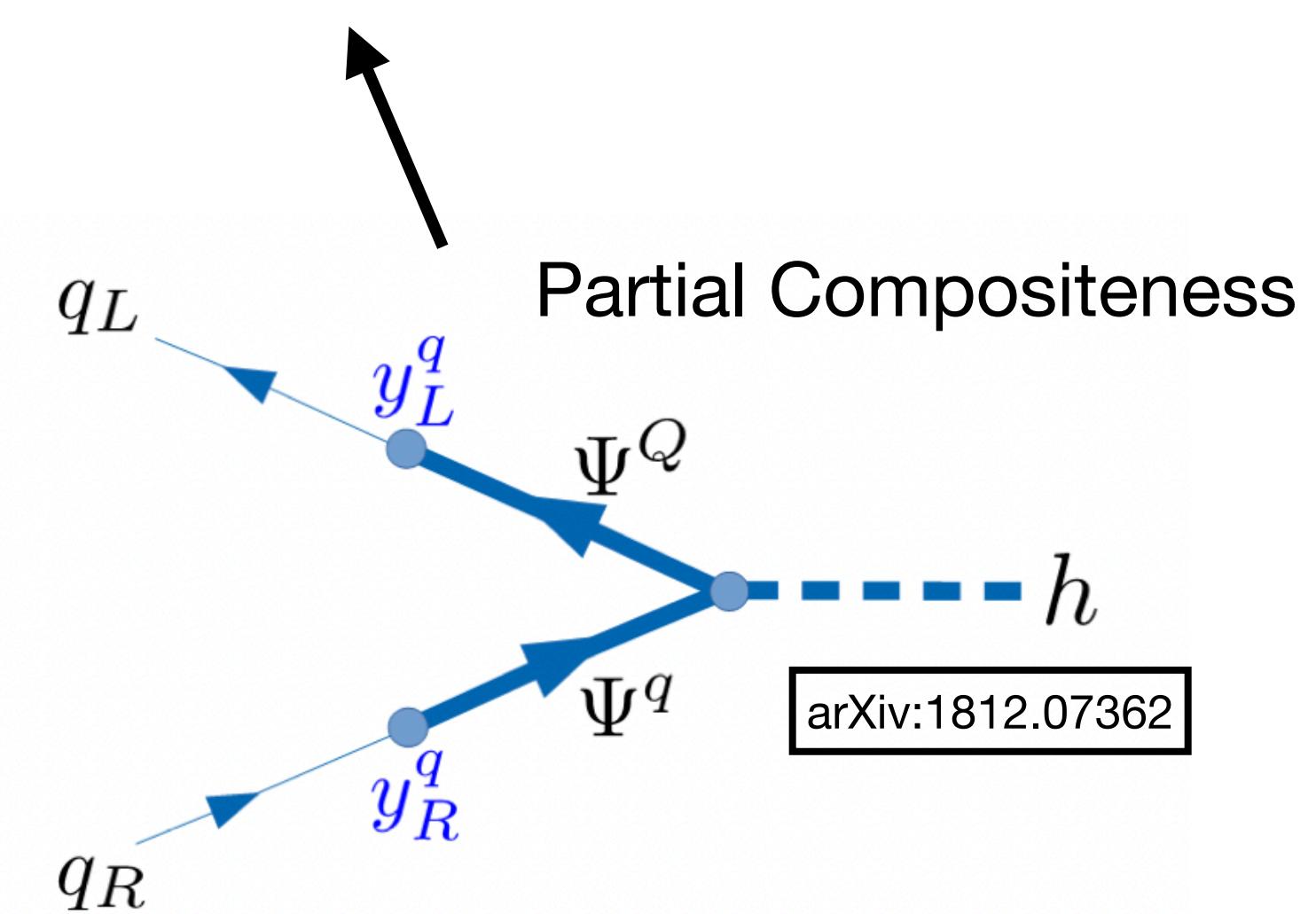
Higgs as a pseudo-Nambu Goldstone Boson

No tree-level potential \Rightarrow naturally light Higgs
(protected by shift symmetry & compositeness)

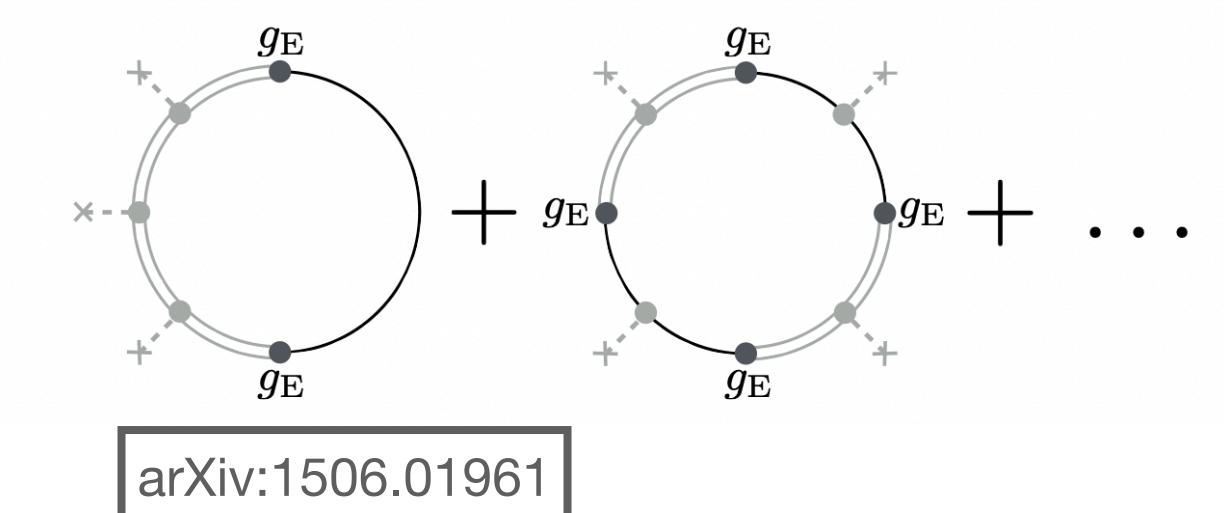
$$\mathcal{L} \supset \mathcal{L}_{\text{el}} + \mathcal{L}_{\text{comp}} + \mathcal{L}_{\text{mix}}$$

Standard Model

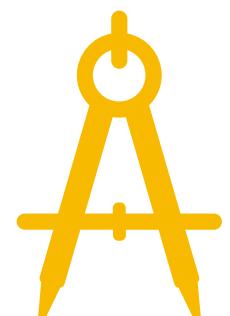
Higgs



Explicit Breaking of global symmetry
by SM fields which transform under
 G_{SM} but not G
main source: top quark



Common CH Problems

-  **(double-)tuning** the Higgs potential


$$V_h = -\alpha \sin^2 \frac{h}{f} + \beta \sin^4 \frac{h}{f}$$

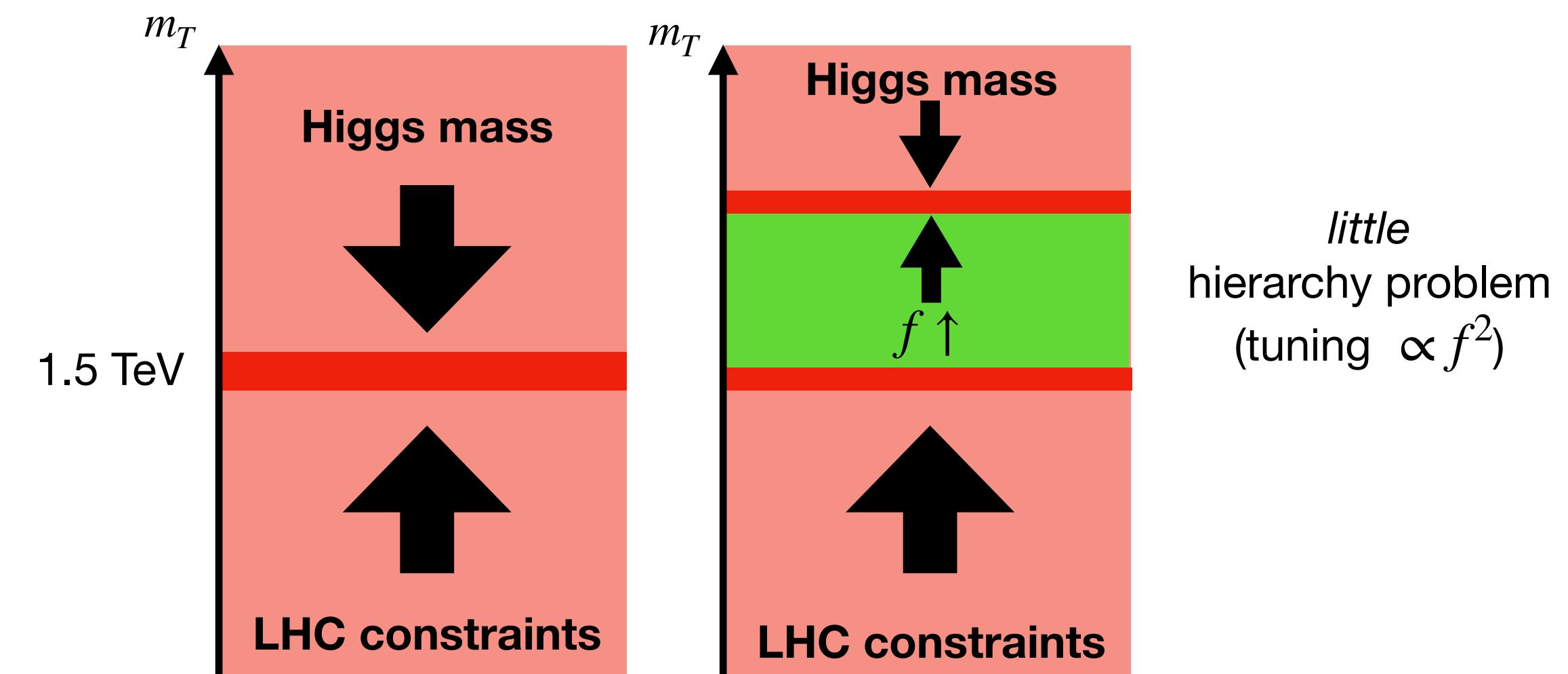
$$\xi = \sin^2(\langle h \rangle / f) = \frac{v_{SM}^2}{f^2} \propto \frac{\alpha}{\beta}$$

vacuum misalignment angle

-  **LHC constraints** on light composite resonances

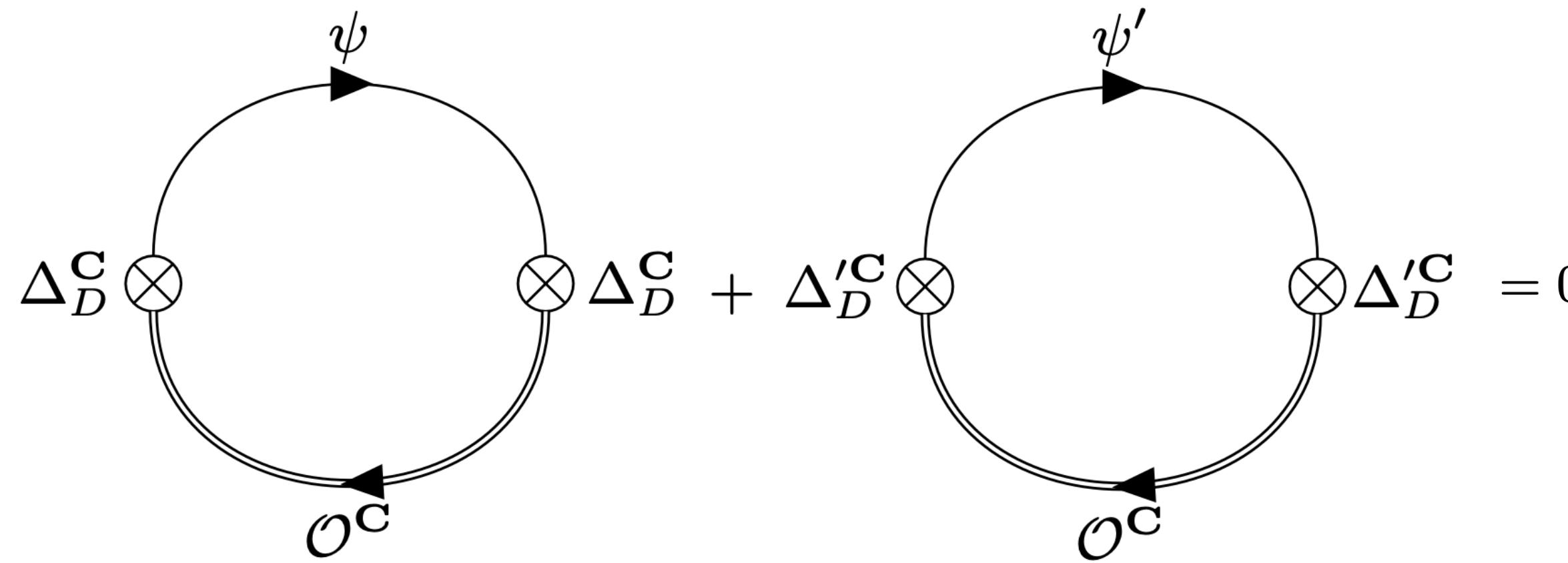
$$m_H \propto \frac{\min(m_T)}{f} m_t$$

analytical estimate
(from 2-site model)



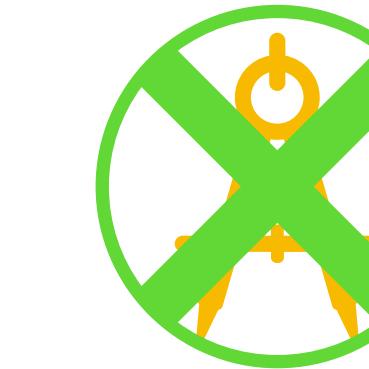
Conjugate Fermion

Mechanism



explicitly:

The quadratic contribution of the top quark to the Higgs potential is cancelled!



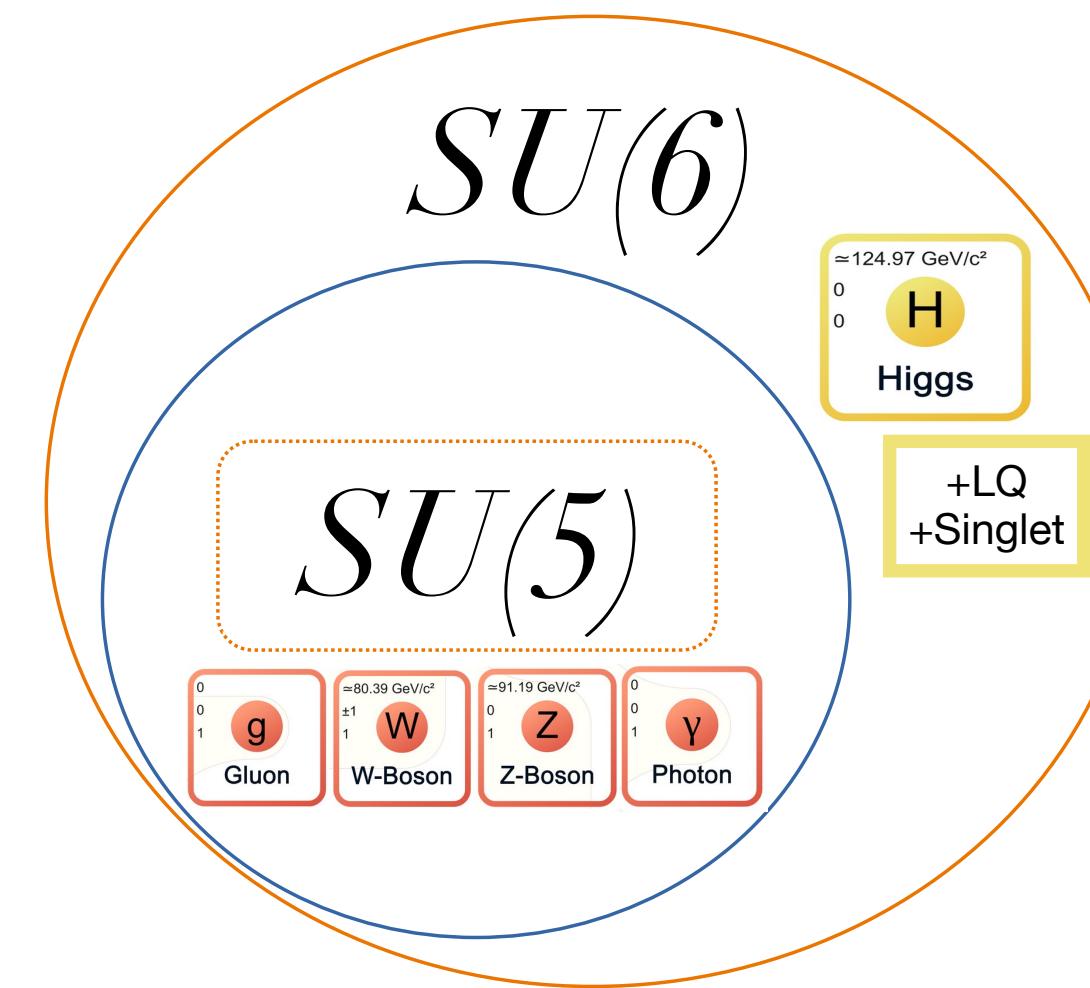
The quadratic contribution of a chiral fermion ψ to the pNGB potential of a coset G/H is cancelled when a new chiral fermion ψ' with conjugated gauge quantum numbers is added, called mirror fermion, if the fermions talk to the same composite operator in a real representation \mathbf{R} of the group G which decomposes as $\mathbf{R} \rightarrow \mathbf{C} \oplus \bar{\mathbf{C}}$ under H , with \mathbf{C} a complex representation and $\bar{\mathbf{C}}$ its complex conjugate.

mirror fermion
has Dirac mass

Conjugate Fermion

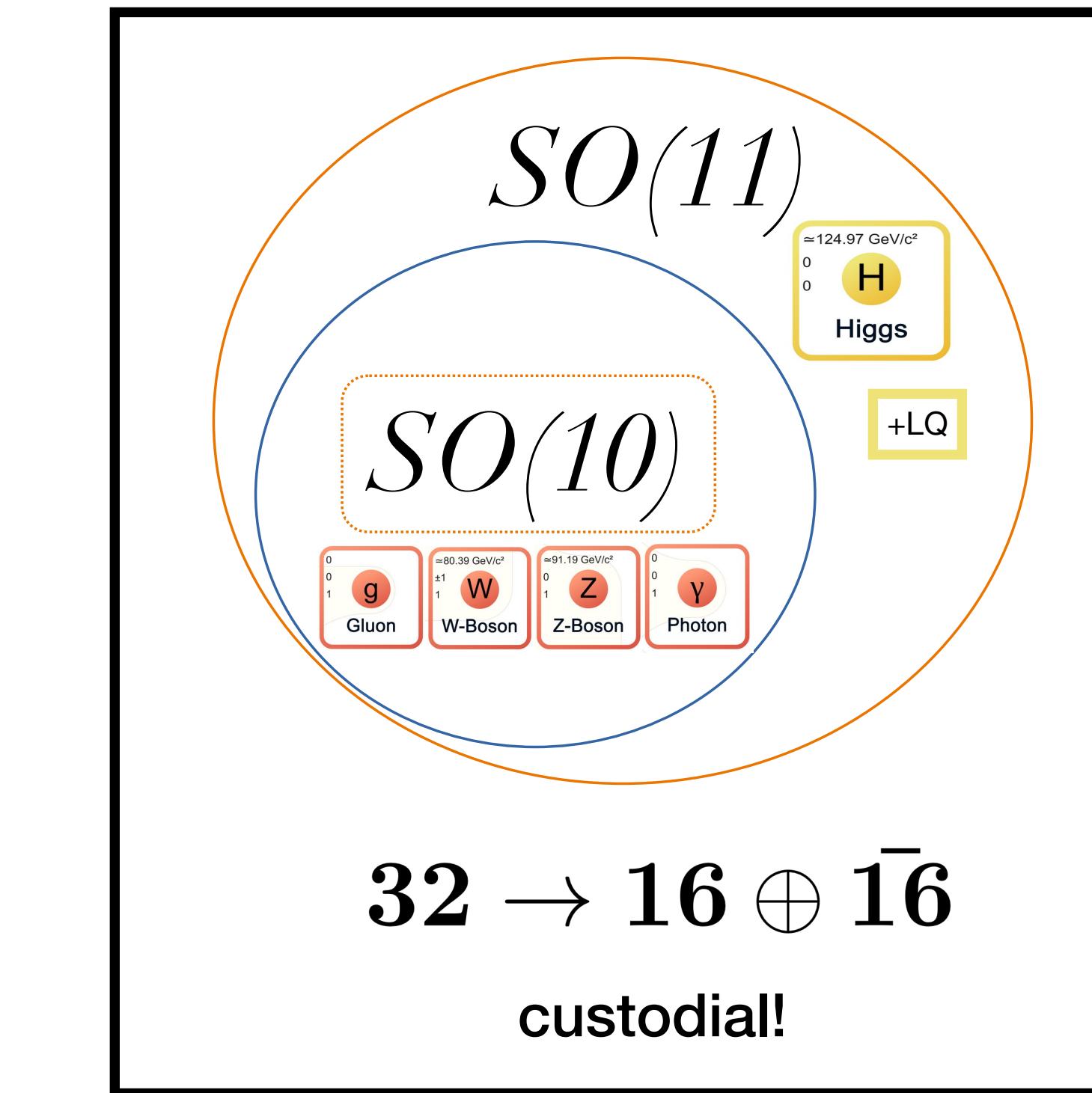
Explicit Model - Composite Grand Unified Theories

minimal:



$$20 \rightarrow 10 \oplus \bar{10}$$

non-custodial \Rightarrow tree-level corrections to T parameter



$$SO(10) \rightarrow SU(5)
16 \rightarrow 10 \oplus 5^* \oplus 1$$

q_L

ω_R

$$10 \rightarrow (3, 2)_{1/6} \oplus (3^*, 1)_{-2/3} \oplus (1, 1)_1$$

$$10^* \rightarrow (3^*, 2)_{-1/6} \oplus (3, 1)_{2/3} \oplus (1, 1)_{-1}$$

θ_L

t_R

Scan:
f = 1600 GeV
 $\lambda_L = \lambda_R$
 $m_t(f) \sim 150$ GeV
parameter range [-5 f , 5 f]
 b_R included
SU(6)/SU(5) or SO(11)/SO(10)

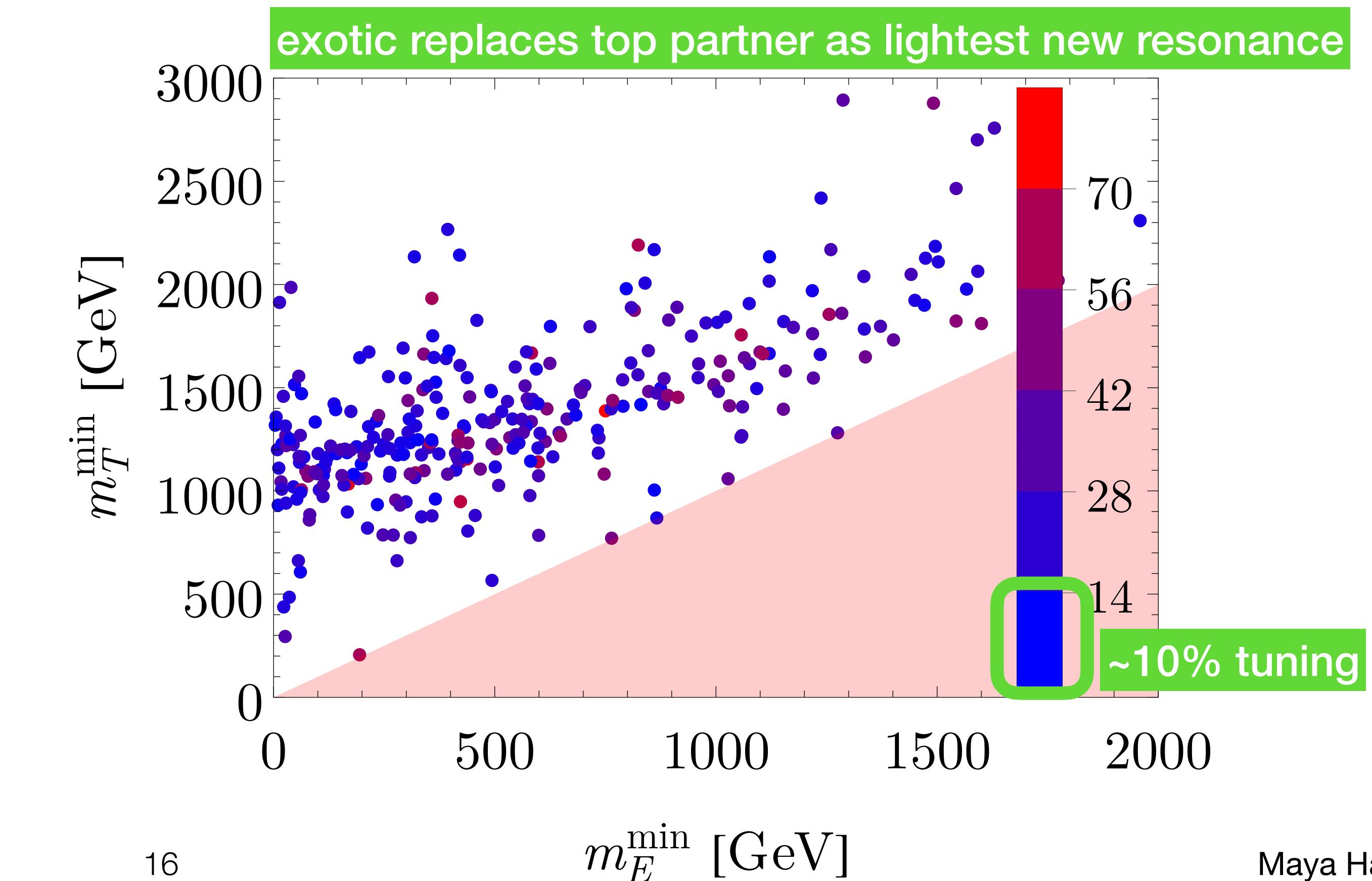
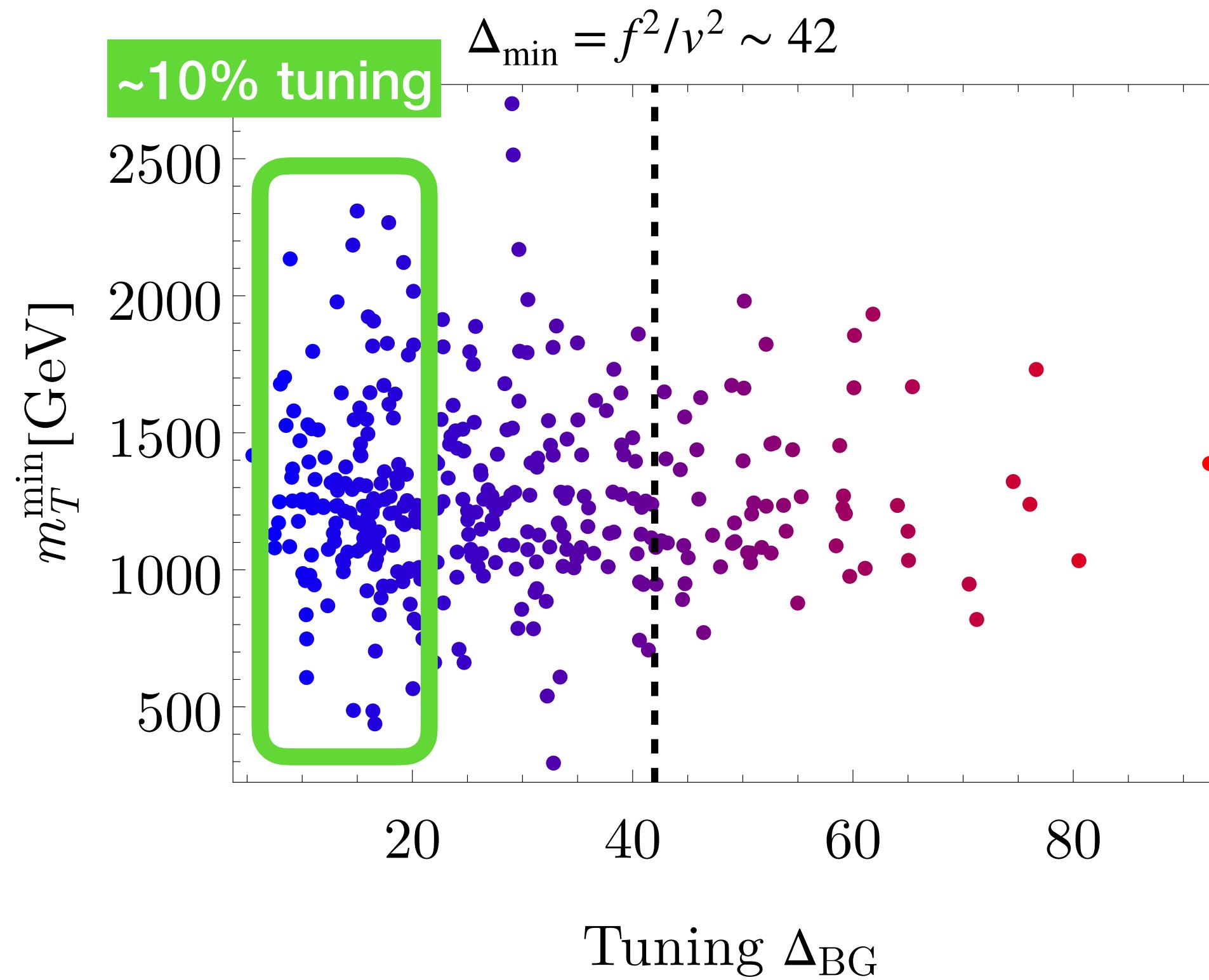
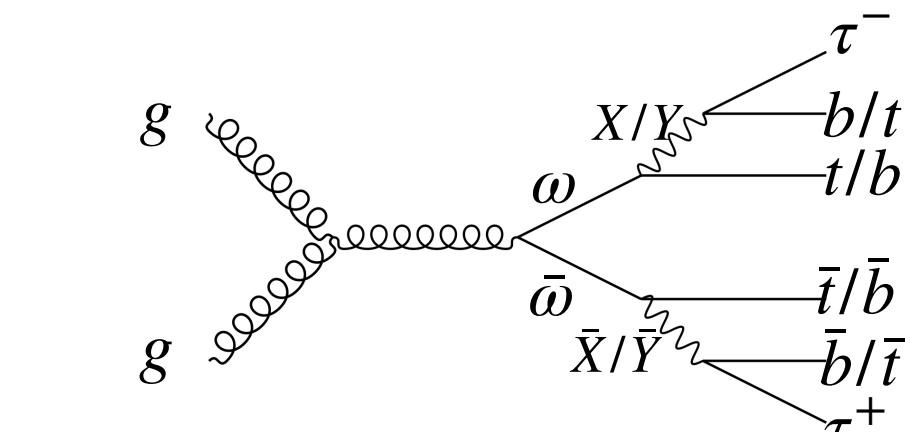
Anglescu, Bally, Goertz, MH
arXiv:2309.05698, PRD

CMS, arXiv:2209.0737; ATLAS, arXiv:2210.15413;...
+many more!

Top Partner ≥ 1500 GeV

Conjugate Fermions

exotics ω carry $B = 2/3$
 \Rightarrow 6 particle final state
 \Rightarrow **unexplored signature!**



$$m_* = g_* f$$


Conjugate Fermion

Proof

Three ingredients why the cancellation works!

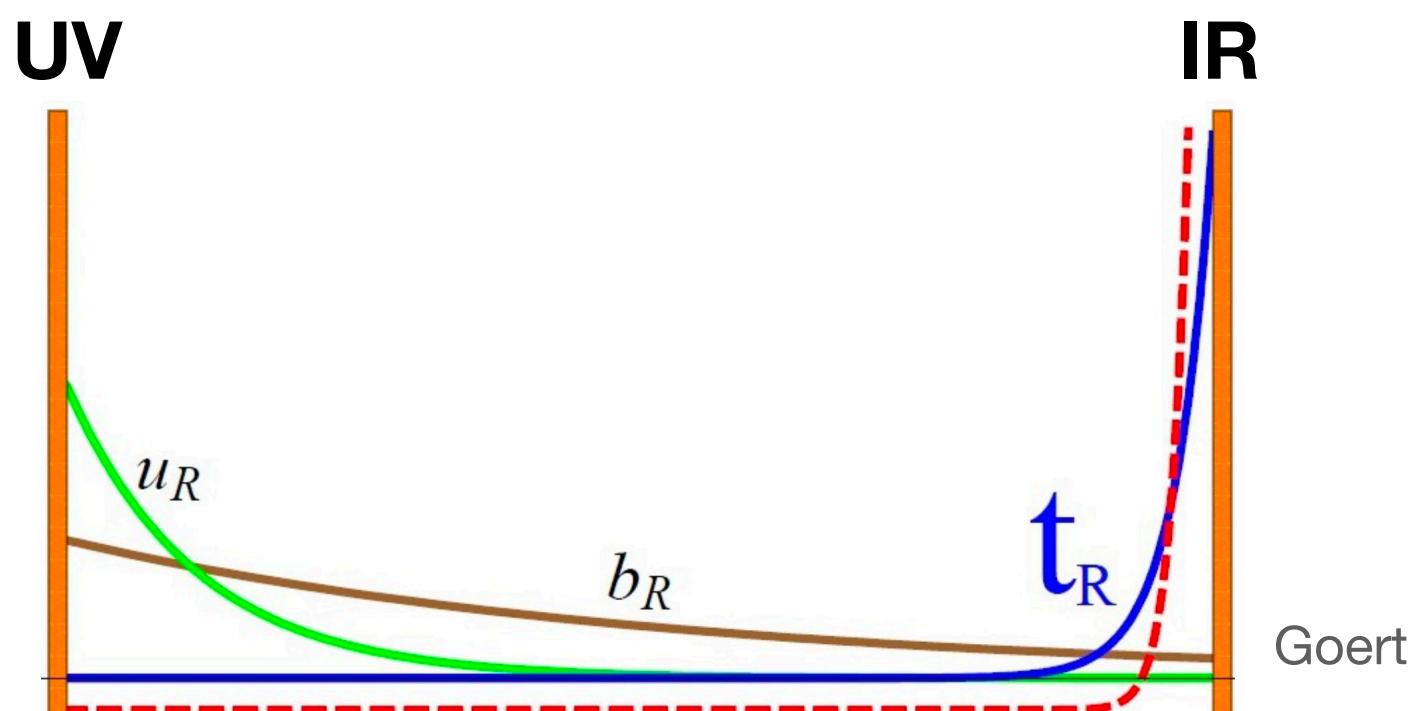
$$(\Delta_D^{\bar{C}})^\dagger \Delta_D^{\bar{C}} = (\Delta'^C_D)^\dagger \Delta'^C_D$$

embedding &
quantum numbers

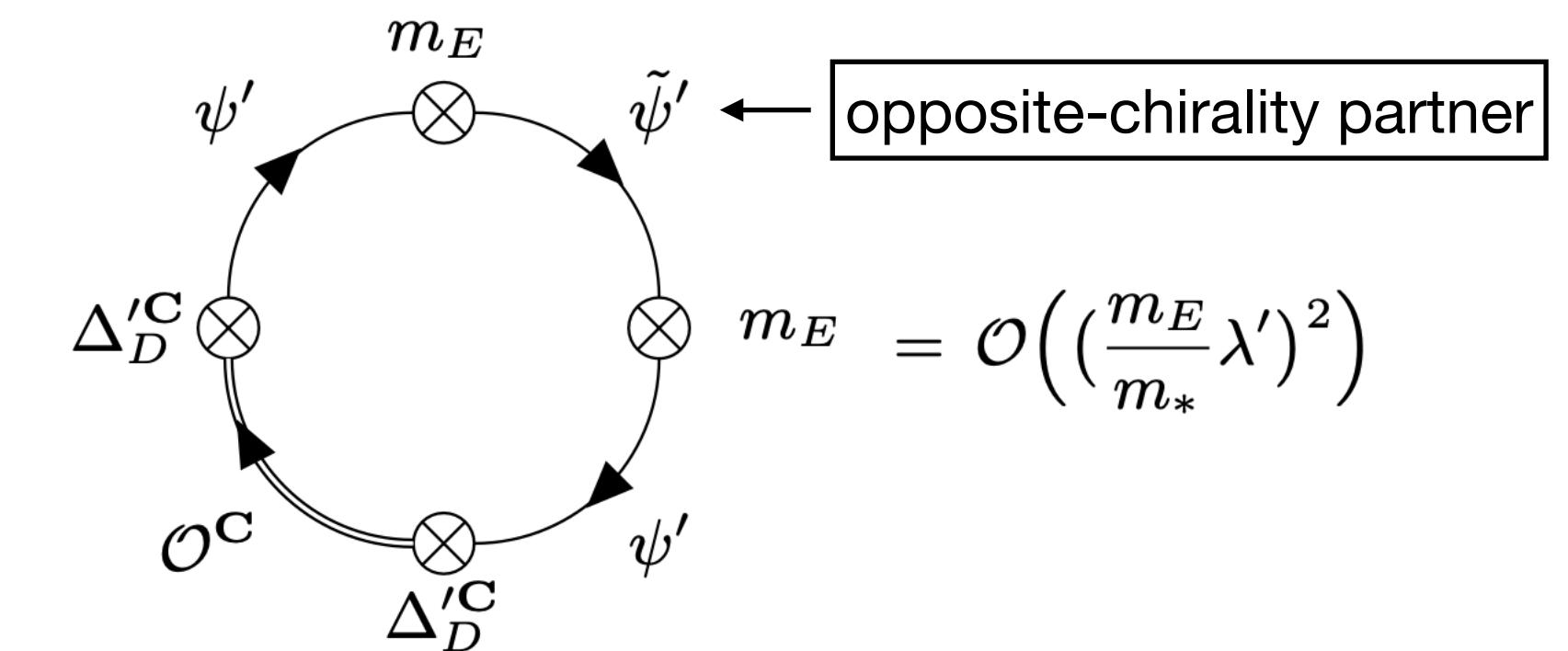
$$R \rightarrow C \oplus \bar{C} , \quad \lambda = \lambda' , \quad m_E \ll m_*$$

5D dual:
same bulk field
same localisation

coincidence?



Dirac mass for mirror fermion
⇒ additional contribution to potential



5D dual:
opposite-chirality partner: **UV brane-localised field**
mirror-fermion: **IR-localised field**
→ **exponential suppression of m_E**