

Gravity from Entanglement: Observable Consequences

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Dutch Institute for Emergent
Phenomena

Gravity from Entanglement: Observable Consequences

- 1. Gravity, entanglement and the holographic principle.**
- 2. Observable quantum gravity effects in interferometers.**
- 3. Gravity from entanglement: observed consequences.**

Is quantum gravity observable?

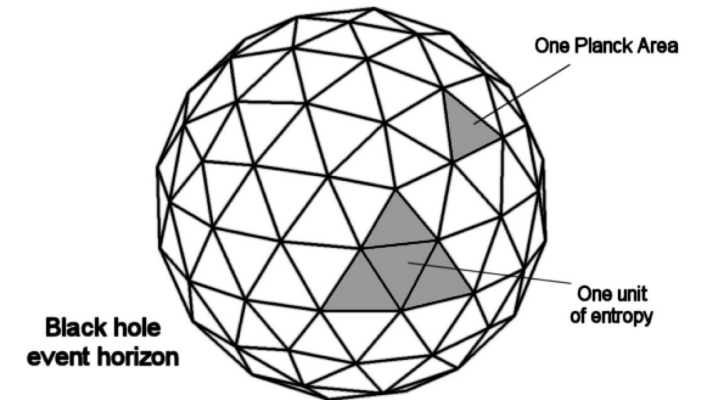
Standard lore:

$$\ell_p = \sqrt{\frac{\hbar G}{c^3}} \approx 10^{-35} m.$$

QG is only observable at the Planck scale.

Why this may not be true:

- *QG leads to divergences*
- *Black hole information paradox*
- *Holographic principle*
- *Infrared effects*

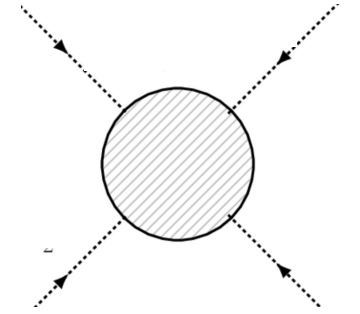


$$S = \frac{A c^3}{4G\hbar}$$

Is quantum gravity observable?

Effective field theory approach to General Relativity suggest

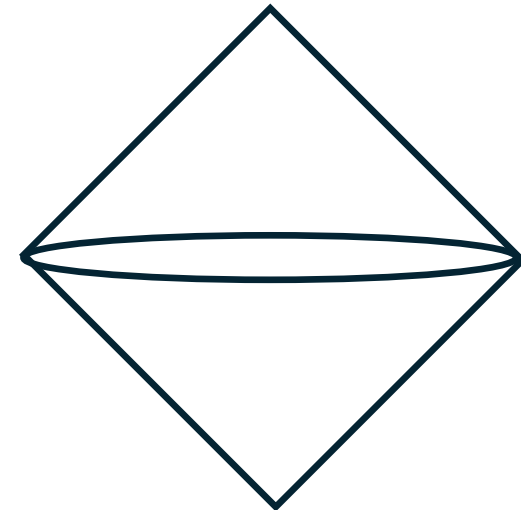
quantum gravity effects are suppressed by powers of : $\frac{E^2}{M_p^2} \sim \frac{L^2}{\ell_p^2}$

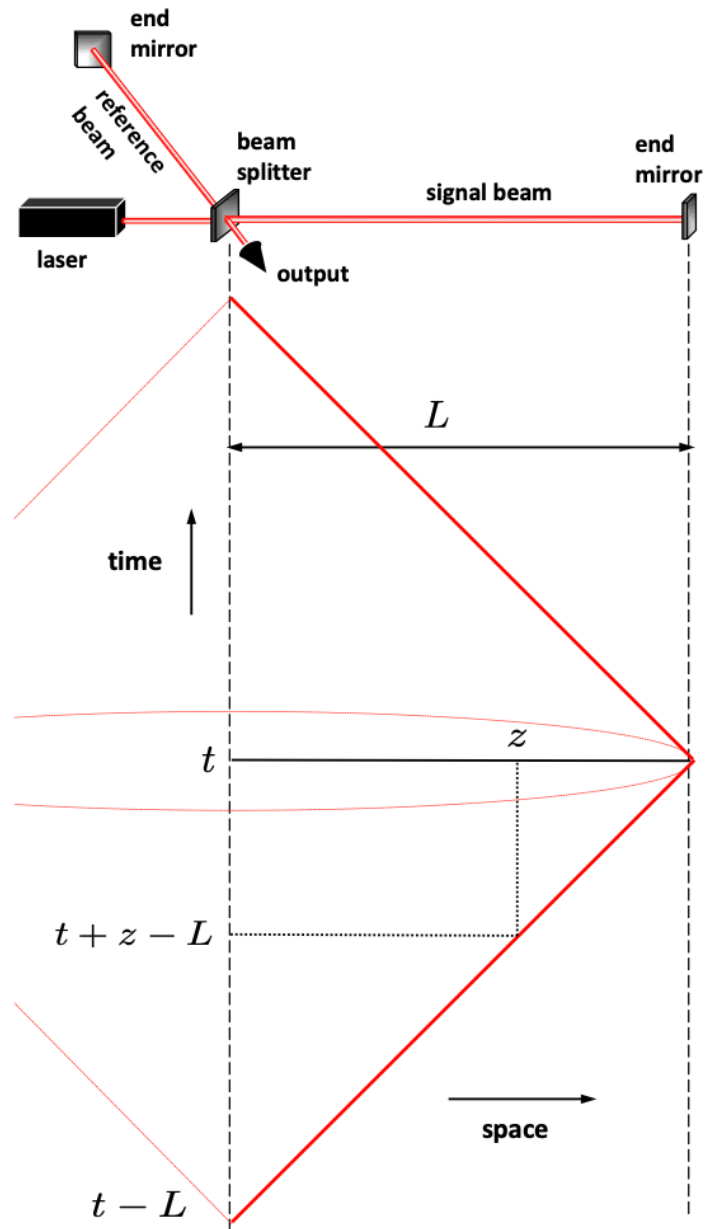


Observable signatures of QG require enhancement due to infrared effects

Theoretical arguments why QG involves situations where (naïve) EFT intuition breaks down.

- string theory: UV/IR correspondence
- black hole information paradox
- holography
- soft physics
- Inherent non-locality



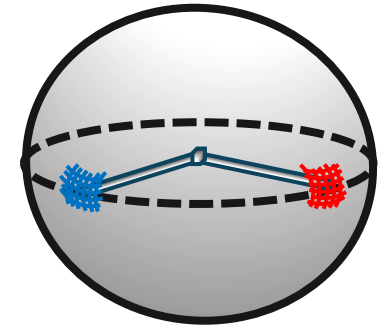


Observational signatures in interferometer experiments

with Kathryn Zurek (2019)

A diagram showing a random walk path with arrows between points x_i and x_f . The path is enclosed in a dashed box labeled D . Below the diagram is the equation:

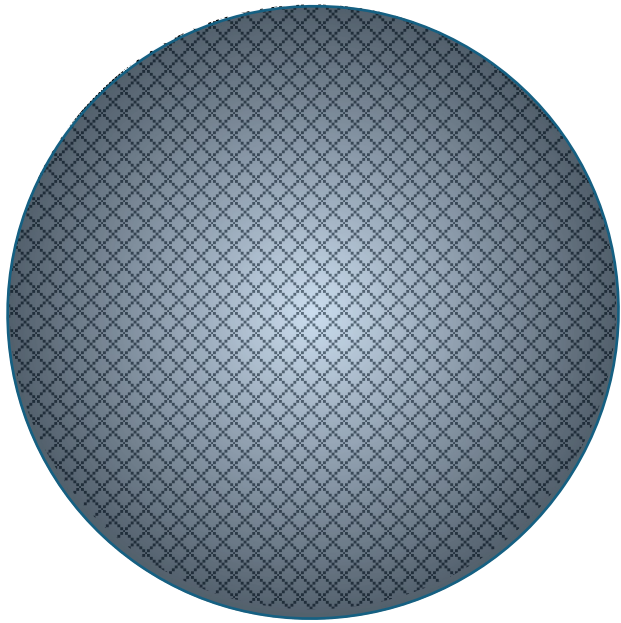
$$\langle \Delta x^2 \rangle = 2DT$$



Holographic Noise in Interferometers

Craig J. Hogan
University of Chicago and Fermilab

Bekenstein-Hawking Entropy



horizon

$$S = \frac{A}{4G\hbar}$$

Hawking temperature

$$T = \frac{\hbar\kappa}{2\pi}$$

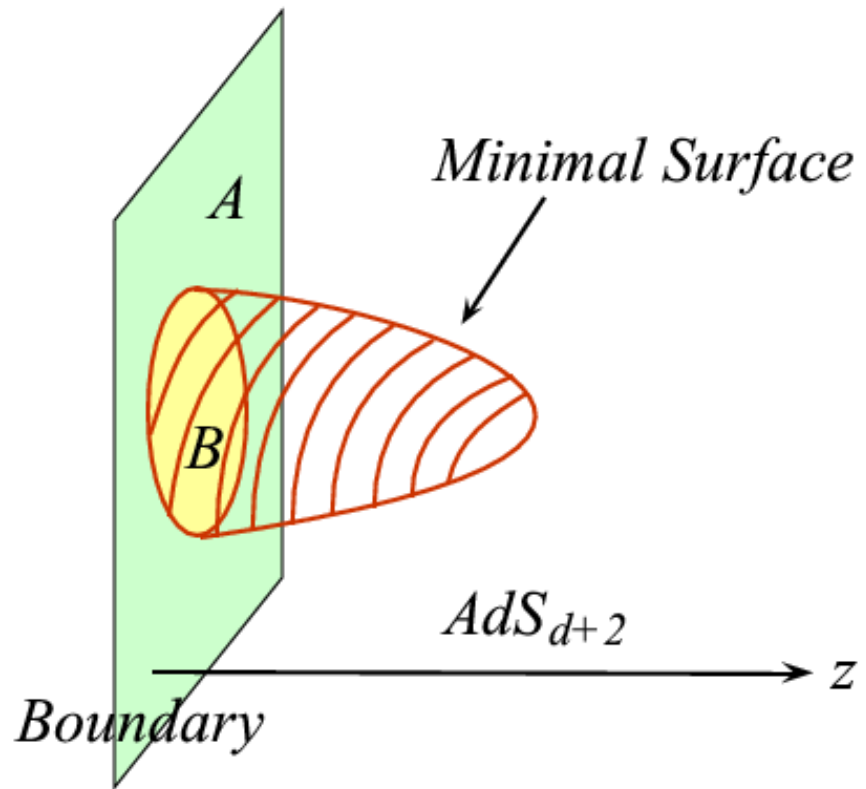
κ = surface gravity

Two possible interpretations

S = $\log(\# \text{ black hole microstates})$

S = entanglement entropy of spacetime vacuum

Ryu-Takanayagi Formula in AdS/CFT

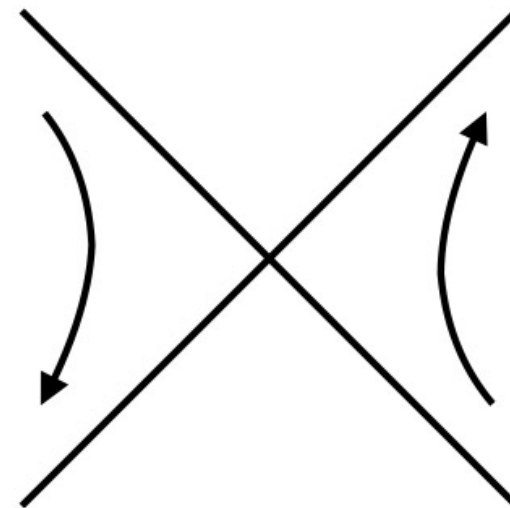


$$S_{ent} = \frac{Area}{4G\hbar}$$

The entanglement entropy of a boundary region B is given by the area of a minimal surface that is homologous to this region

Quantum Entanglement of the Vacuum

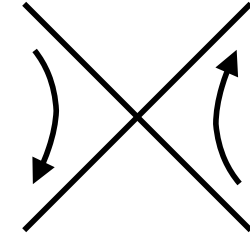
$$|\text{vac}\rangle = \frac{1}{\sqrt{Z}} \sum_i |K_i\rangle_L |K_i\rangle_R e^{-\pi K_i / \hbar}$$



The Minkowski vacuum, when viewed from the perspective of an accelerated observer, looks like a purified thermal state.

This is nowadays called the Thermo-field double state.

1st law of entanglement entropy



$$S = -\text{tr}(\rho \log \rho)$$

$$\rho = \text{tr}_{\mathcal{H}'} (|\Psi\rangle\langle\Psi|)$$

$$\rho = \frac{e^{-K}}{Z}$$

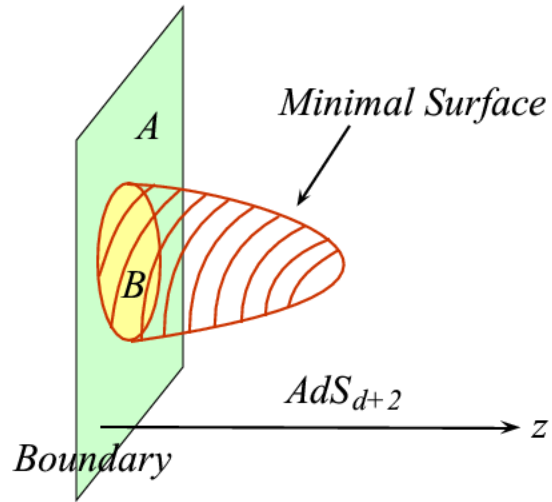
$$Z = \text{tr} (e^{-K})$$

Modular Hamiltonian K : defined by entanglement =>
becomes boost generator (= emergent time)

$$\delta S = \delta\langle K\rangle$$

$$\delta\langle K\rangle = \text{tr} (\delta\rho K)$$

Gravity from Quantum Entanglement



Gravitational equations are derived from 1st law of **entanglement entropy**

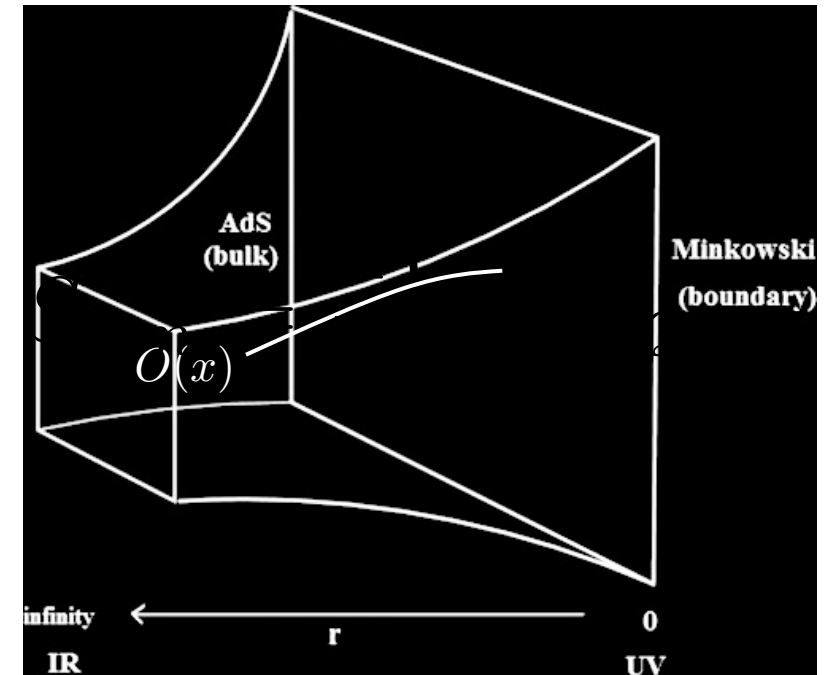
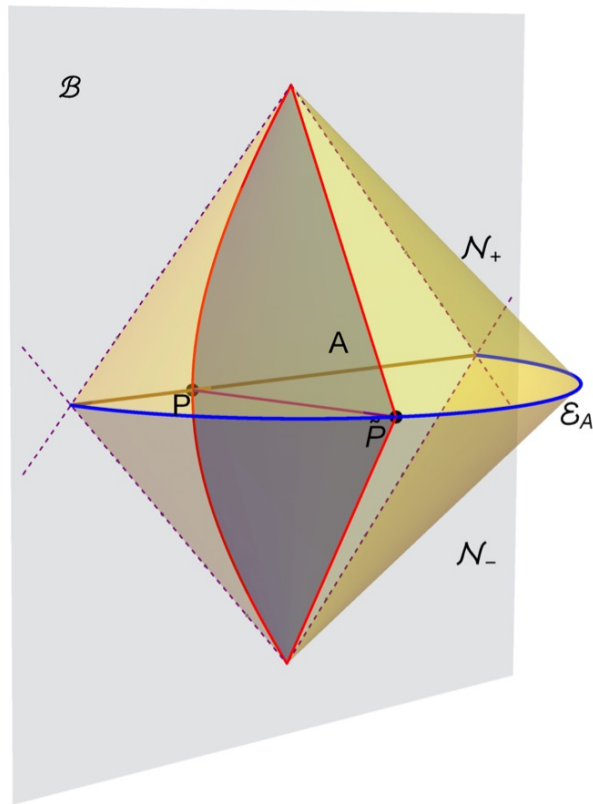
$$S_{ent} = \frac{Area}{4G\hbar}$$

Ryu, Takanayagi
van Raamsdonk
Myers, Casini et al.

Observables in quantum gravity are inherently non-local

General covariance => observables are defined relationally.

In AdS/CFT observables are defined in relation to the boundary.



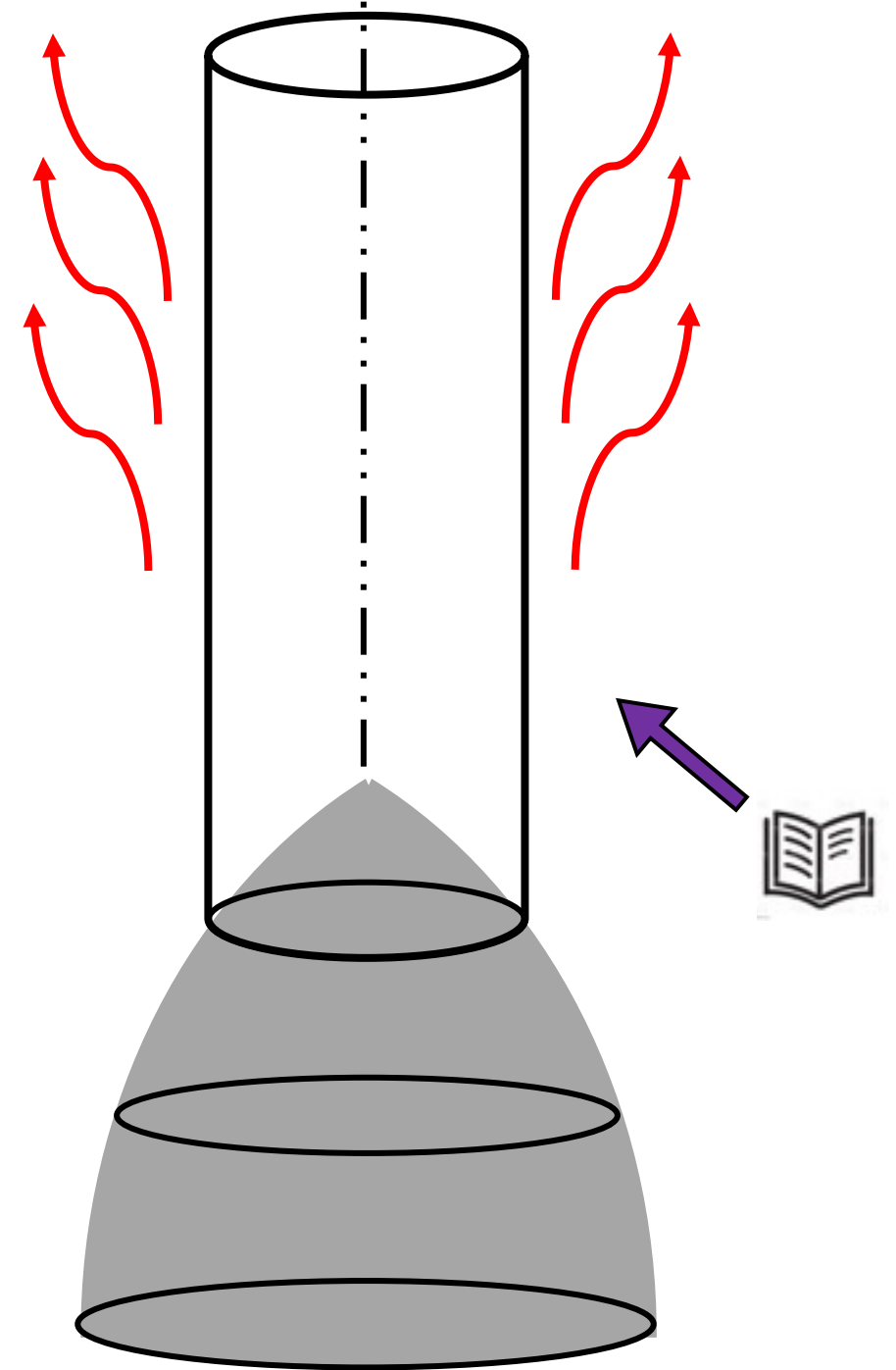
We do not live in Anti-de Sitter space...

but we can learn from AdS/CFT

Black hole information paradox

Semiclassical GR + local EFT:
BH formation + evaporation
destroys (quantum) information.

Unitarity in QM: can only be restored
via non-local quantum gravity effects.

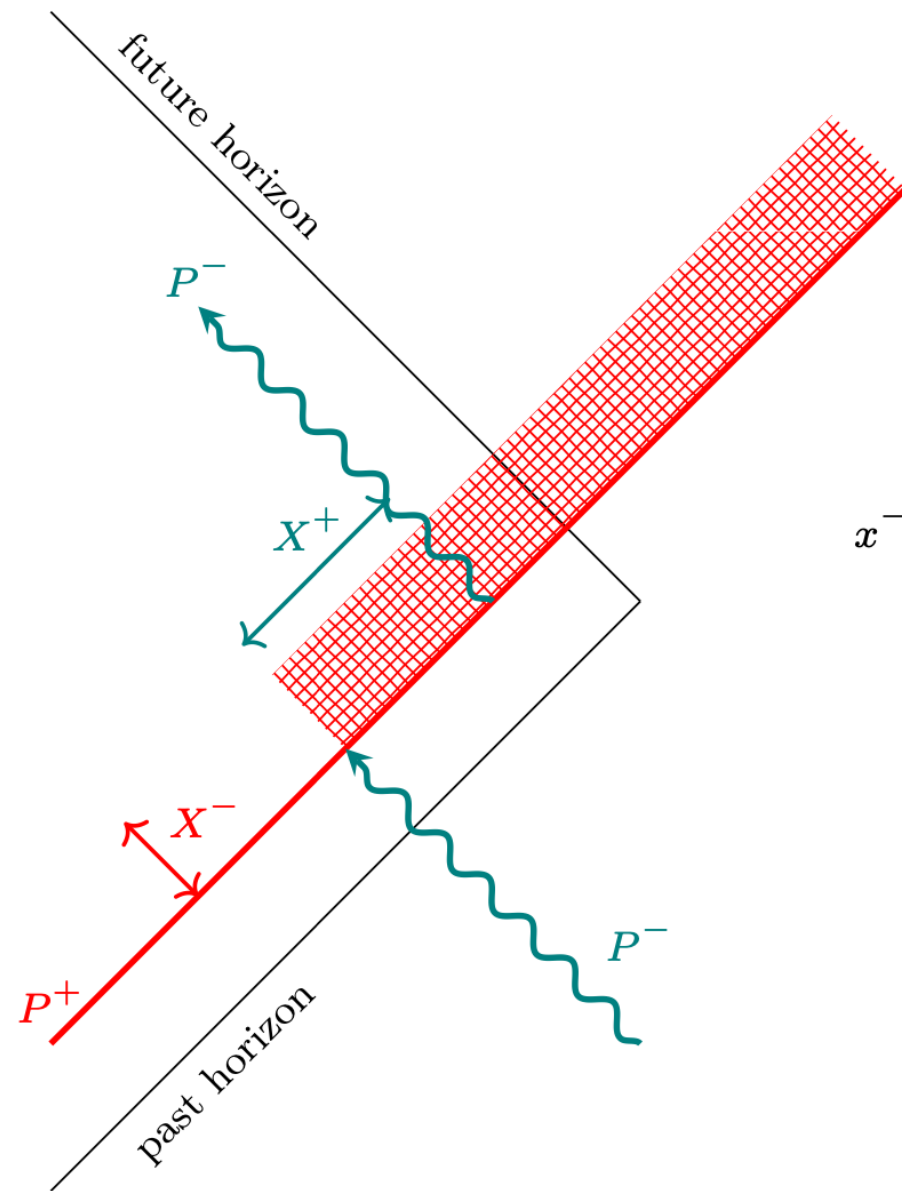


Lessons learned from the Black hole information paradox

- Horizon scale physics
- Quantum entanglement.
- Quantum chaos
- Shockwaves

't Hooft commutation relations

$$[X^+(z, \bar{z}), X^-(z', \bar{z}')] = 8\pi i G_N G(z - z')$$



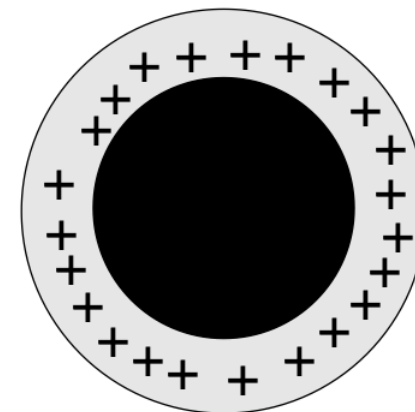
On the quantum width of a black hole horizon

Donald Marolf

Physics Department, UCSB, Santa Barbara, CA 93106. marolf@physics.ucsb.edu

$$\delta L^2 \sim \frac{L^2}{\sqrt{S_{BH}}},$$

$$\delta L \sim \sqrt{l_p L}.$$

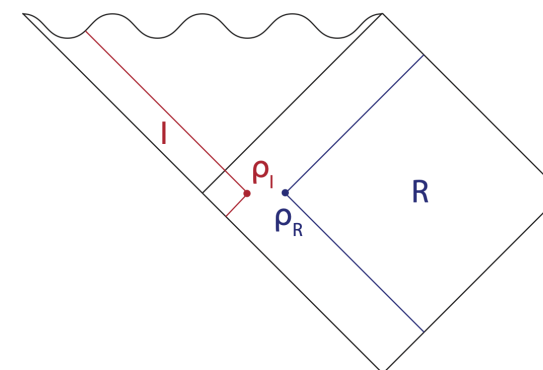


	+	-	+	-	0
	0	+	+	+	+
	0	+	+	+	0
	0	+	+	+	-
	+	-	-	0	-

Islands Far Outside the Horizon

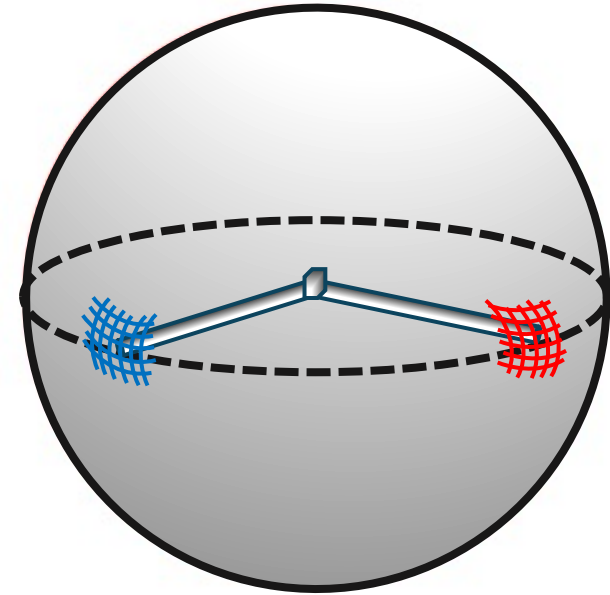
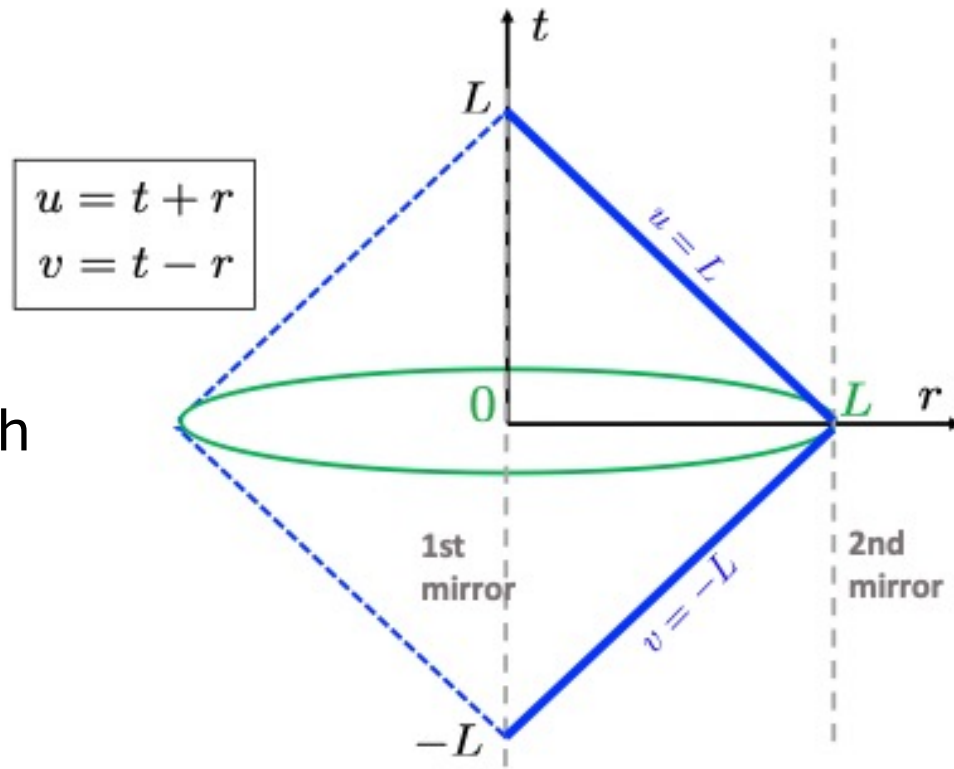
Raphael Bousso and Geoff Penington

*Center for Theoretical Physics and Department of Physics,
University of California, Berkeley, California 94720, U.S.A.*



Causal Diamond

contains single light path



Length fluctuations

$$\delta L = (\delta v(L) + \delta u(L)) / 2.$$

$$\delta L^2 = \delta u(L) \delta v(L)$$

$$\delta v(L) = \int_{-L}^L du \, h_{uu}(u, L)$$

$$\delta u(L) = \int_{-L}^L dv \, h_{vv}(L, v).$$

Calculation follows Marolf's argument

Via the coordinate transformation

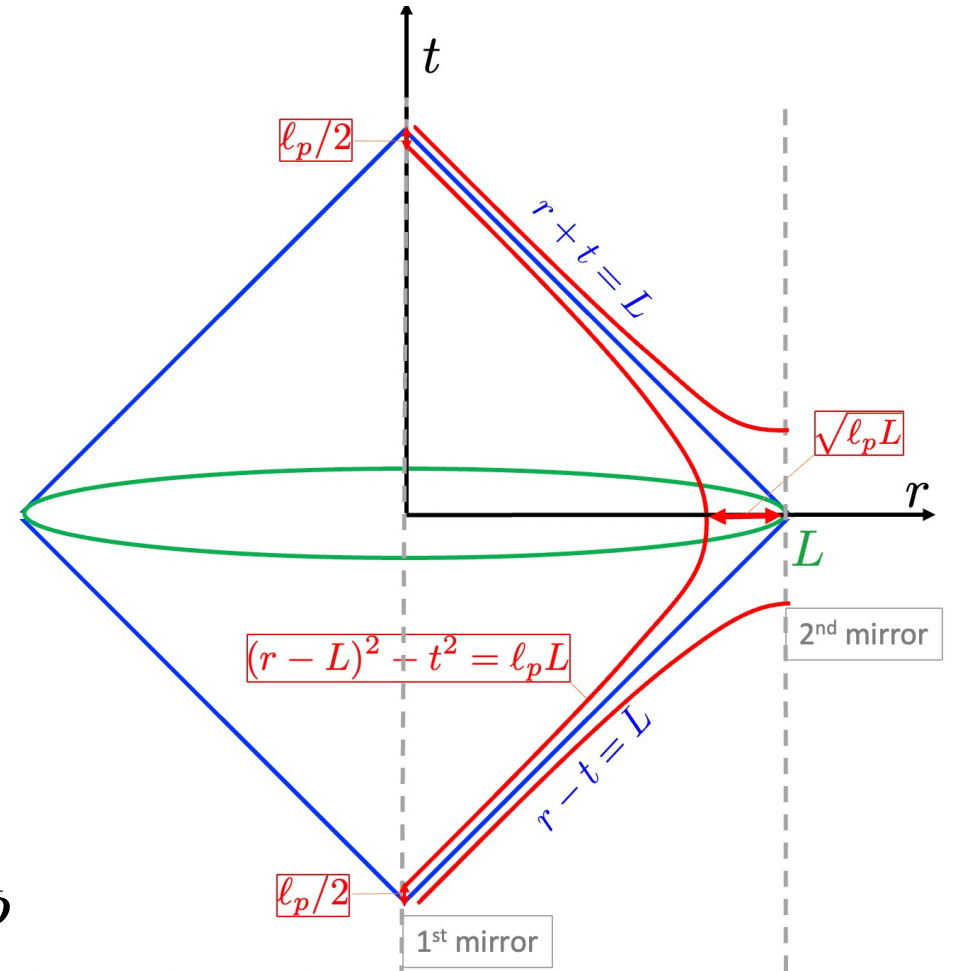
$$(u - L)(v - L) = 4L^2 f(R), \quad \log \frac{u - L}{v - L} = \frac{T}{L}$$

The metric becomes spherical Rindler space

$$ds^2 = -f(R)dT^2 + \frac{dR^2}{f(R)} + r^2(d\theta^2 + \sin^2 \theta d\phi^2),$$

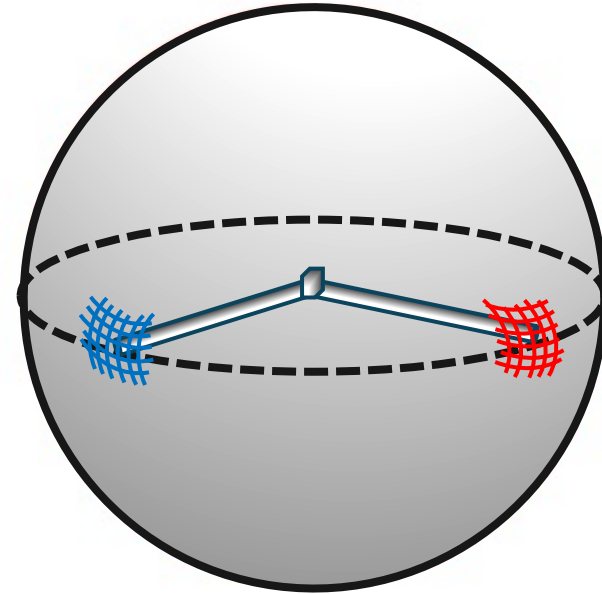
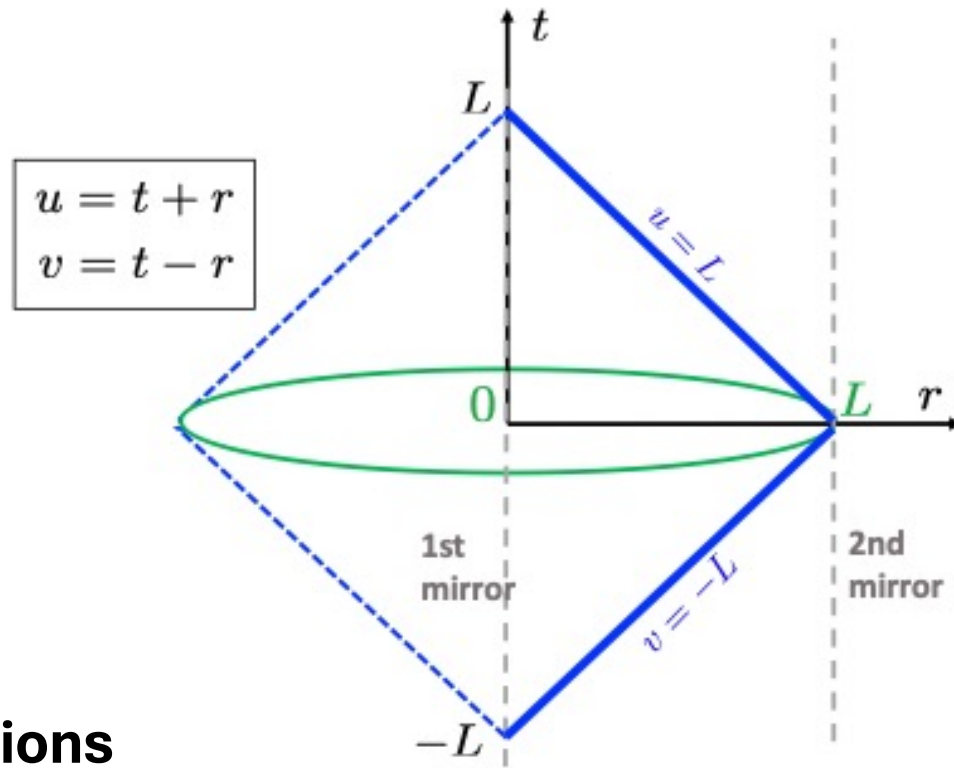
Backreaction effects are encoded in the potential Φ

$$f(R) = 1 - \frac{R}{L} + 2\Phi.$$



Causal Diamond

boundary behaves like
Rindler horizon



Modular energy fluctuations
lead to variations of the Newtonian potential

$$\frac{\delta v(L) \delta u(L)}{L^2} = 2\Phi(L)$$

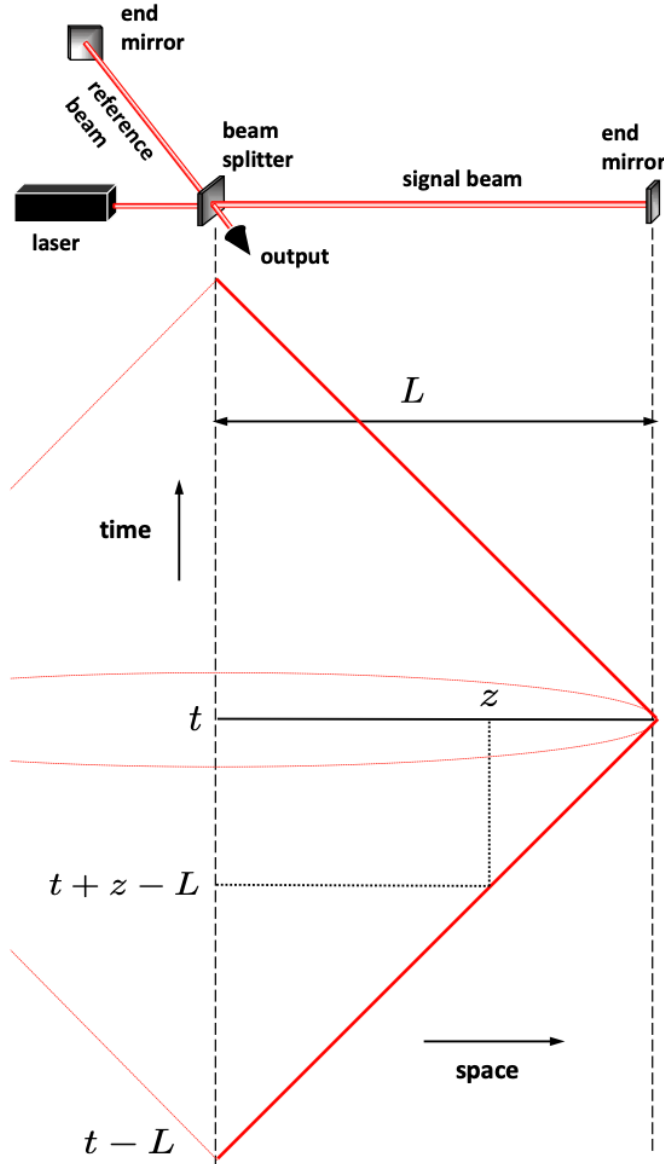
$$\delta L^2 = \delta u(L) \delta v(L)$$

$$\langle \Delta M^2 \rangle = -\frac{\partial^2}{\partial \beta^2} (\beta F) = \frac{1}{l_p^2}.$$

$$\left\langle \frac{\delta L^2}{L^2} \right\rangle = \frac{l_p^2 \Delta M}{4\pi L} = \frac{l_p}{4\pi L},$$

Observational signatures in interferometer experiments

with Kathryn Zurek (2019)



Vacuum energy fluctuations

$$\langle \Delta M^2 \rangle = -\frac{\partial^2}{\partial \beta^2} (\beta F) = \frac{1}{l_p^2}.$$

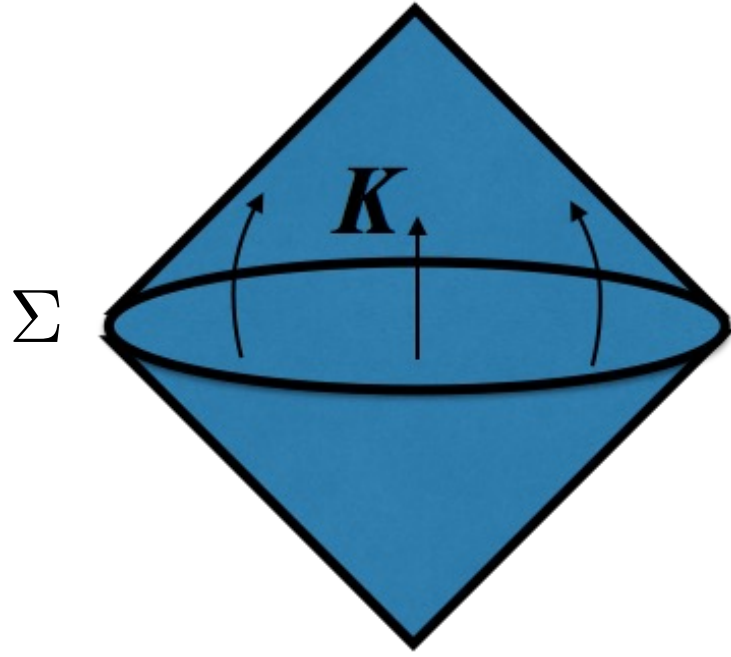
lead to length fluctuations

$$\left\langle \frac{\delta L^2}{L^2} \right\rangle = \frac{l_p^2 \Delta M}{4\pi L} = \frac{l_p}{4\pi L},$$

Actual observable is the power spectral density

$$S(\omega, t) = \int_{-\infty}^{\infty} d\tau \left\langle \frac{\delta L(t)}{L} \frac{\delta L(t - \tau)}{L} \right\rangle e^{-i\omega\tau}$$

Causal Diamonds, Entanglement and Modular Energy



Entanglement entropy:

$$\langle K \rangle = \frac{A(\Sigma)}{4G}.$$

Region bounded by past and future lightcones of two timelike separated points.

Described by a density matrix:

$$\rho = \text{tr}_{\mathcal{H}'} (|vac\rangle\langle vac|) = e^{-K}$$

Modular Energy fluctuations

$$\langle \Delta K^2 \rangle = \frac{A(\Sigma)}{4G}.$$

Spacetime Fluctuations in AdS/CFT

(with Kathryn Zurek)

The AdS-metric

$$ds^2 = L^2 \frac{dz^2 + dx_i^2 - dx_0^2}{z^2}$$

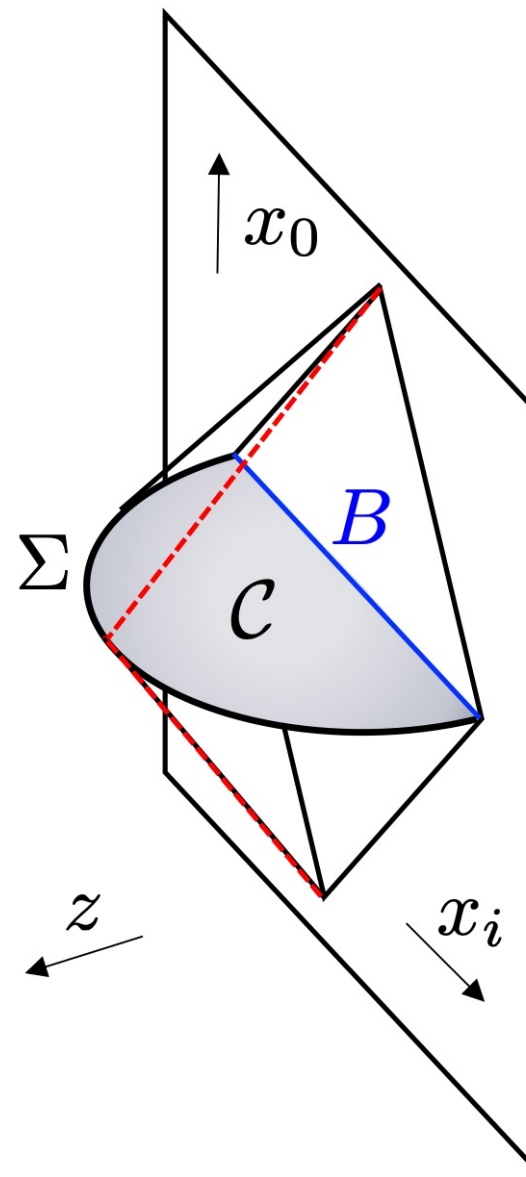
can be transformed to the AdS-Rindler metric

$$ds^2 = - \left(\frac{r^2}{L^2} - 1 \right) dt^2 + \left(\frac{r^2}{L^2} - 1 \right)^{-1} dr^2 + r^2 d\Sigma_{d-1}^2,$$

The fluctuations in the modular Hamiltonian

$$\Delta K = \int_{\mathcal{C}} \xi_K^\mu T_{\mu\nu}^{\text{bulk}} d\mathcal{C}^\nu,$$

Can be computed using (a gravitational version of) the replica method



Spacetime Fluctuations in AdS/CFT

From the free energy on the n -fold cover

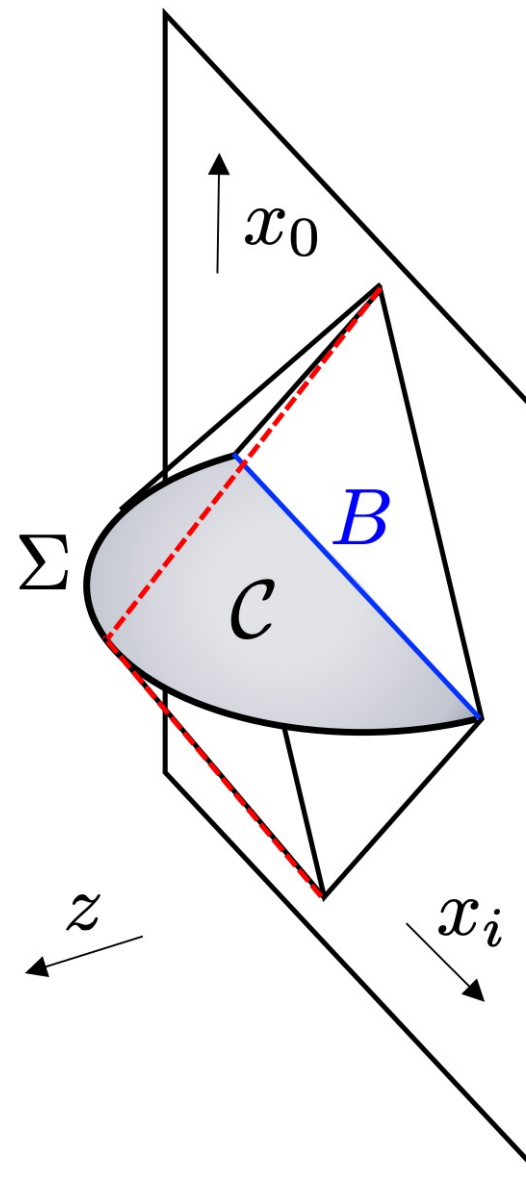
$$Z_n = \text{tr}(e^{-nK}) = e^{-nF_n}$$

One derives the expectation value and fluctuations of the modular Hamiltonian

$$\langle K \rangle = \frac{d}{dn} (nF_n) \Big|_{n=1} \quad \langle \Delta K^2 \rangle = -\frac{d^2}{dn^2} (nF_n) \Big|_{n=1}$$

One finds

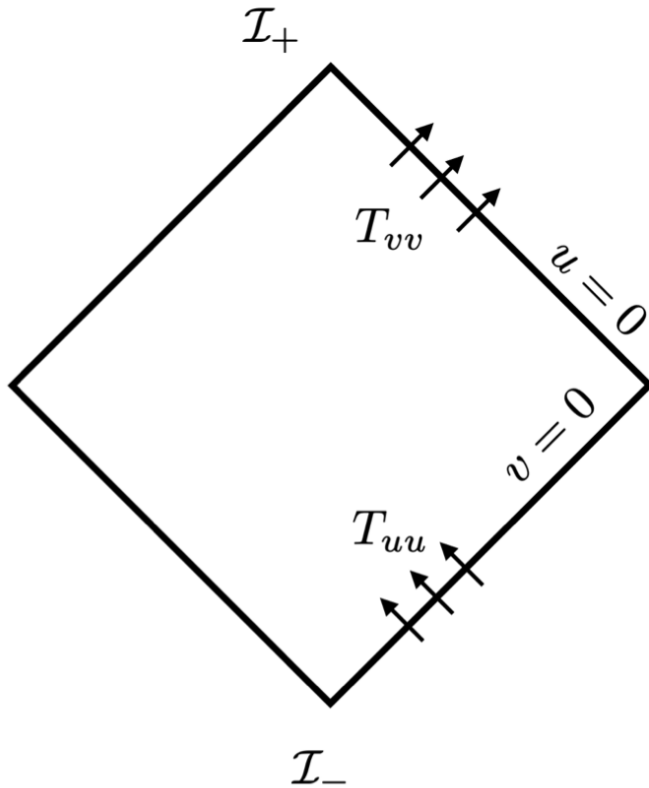
$$\langle K \rangle = \frac{A(\Sigma)}{4G} \quad \langle \Delta K^2 \rangle = \frac{A(\Sigma)}{4G}$$



Modular Energy Fluctuations from 't Hooft commutation relations

Metric fluctuations

$$ds^2 = -dudv + \nabla_y X^v dudy + \nabla_y X^u dvdy + dy^2.$$



The modular energy equals

$$\begin{aligned} K &= \frac{1}{\ell_p^{d-2}} \int d^{d-2}y \nabla_y X^u(y) \nabla_y X^v(y) \\ &= \int d^{d-2}y \left[\int_{-\infty}^0 du X^u T_{uu} + \int_0^{\infty} dv X^v T_{vv} \right] \end{aligned}$$

As operators the coordinate shifts obey

$$\langle X^u(y) X^v(y') \rangle = \ell_p^{d-2} f(y, y').$$

Modular Energy Fluctuations from 't Hooft commutation relations

Metric fluctuations

$$ds^2 = -dudv + \nabla_y X^v dudy + \nabla_y X^u dvdy + dy^2.$$

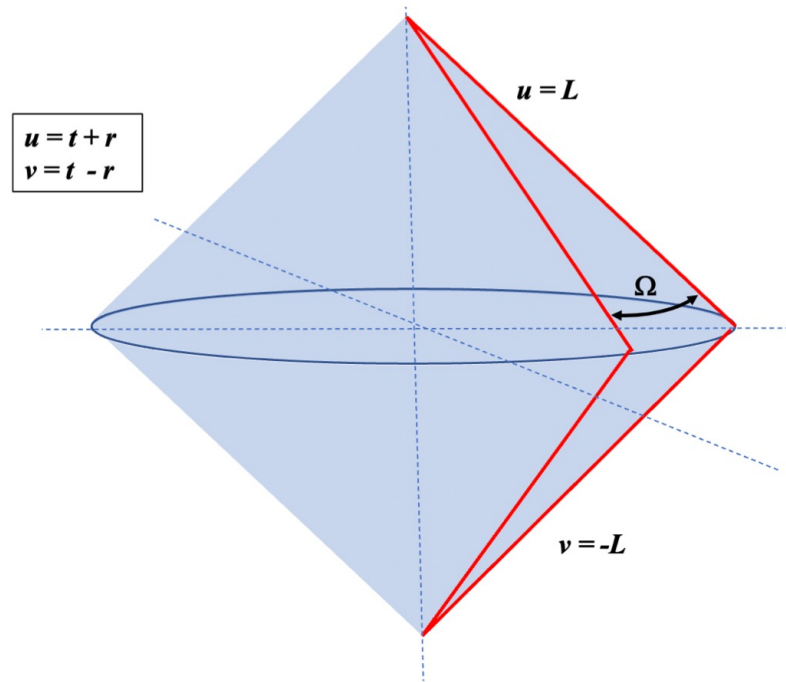
The modular energy equals

$$K = \frac{1}{\ell_p^{d-2}} \int d^{d-2}y \nabla_y X^u(y) \nabla_y X^v(y)$$

We derive using $\langle X^u(y) X^v(y') \rangle = \ell_p^{d-2} f(y, y').$

the following size of the metric fluctuations

$$\langle (\delta g_y^u \delta g_y^v)^2 \rangle \sim \langle [\nabla_y X^u \nabla_y X^v]_{avg}^2 \rangle \sim \left(\frac{\ell_p}{L} \right)^{d-2}$$



Heising-Simons Foundation (HSF) Collaboration on Quantum Gravity and Its Observational Signatures (QuRIOS)

Seeking to bridge the divide between theory and observability



Kathryn Zurek (Caltech)

Particle theory, Effective field theory and models



Erik Verlinde (University of Amsterdam)

String theory, emergent gravity



Maulik Parikh (Arizona State University)

Quantum gravity, classical and quantum black holes



Cynthia Keeler (Arizona State University)

String theory, fluid gravity



Steve Giddings (UC Santa Barbara)

Quantum gravity, black holes

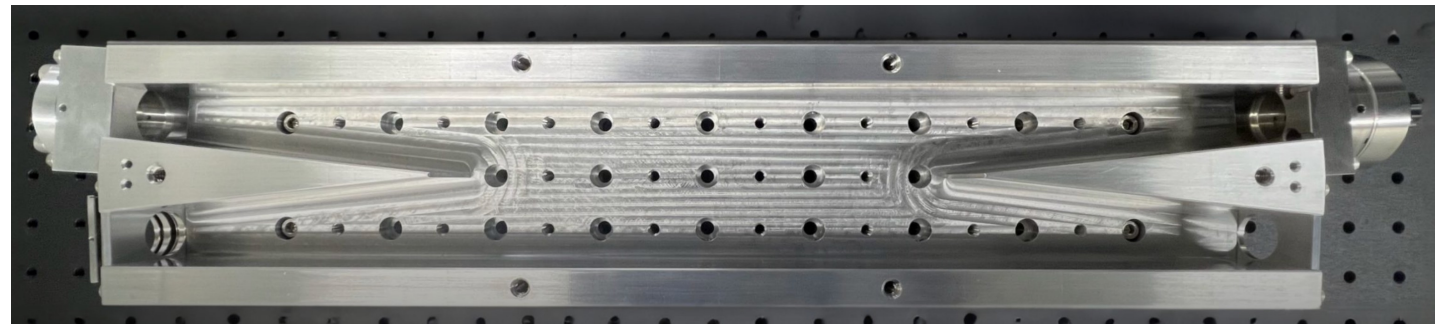
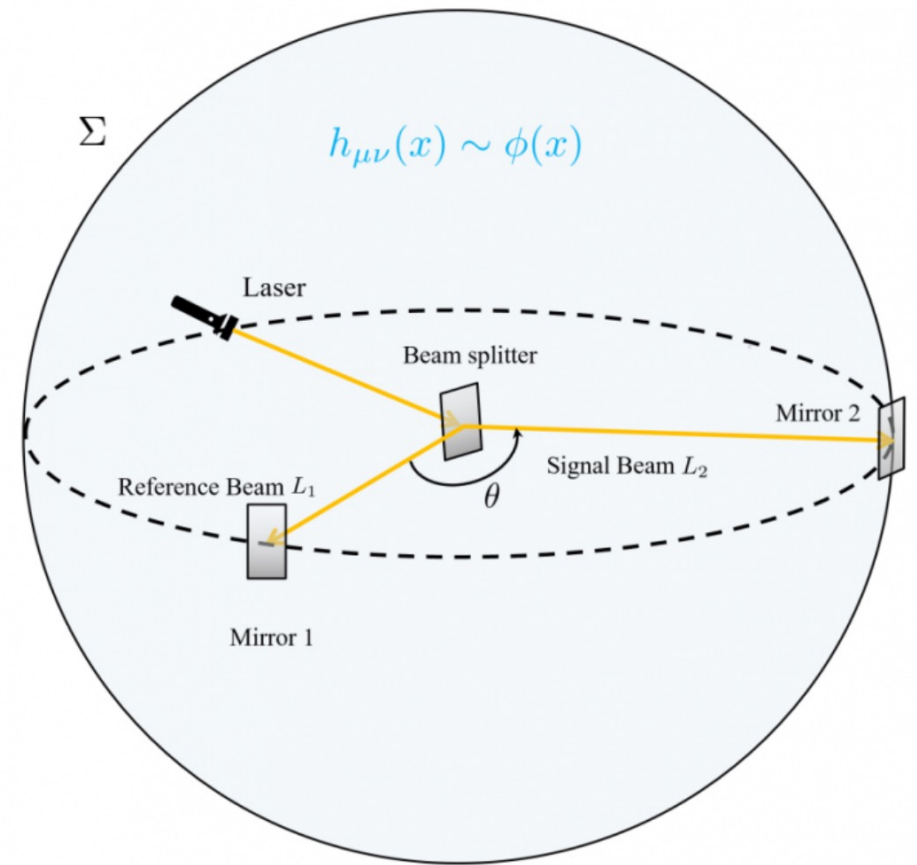


Ben Freivogel (University of Amsterdam)

String theory, cosmology, early universe

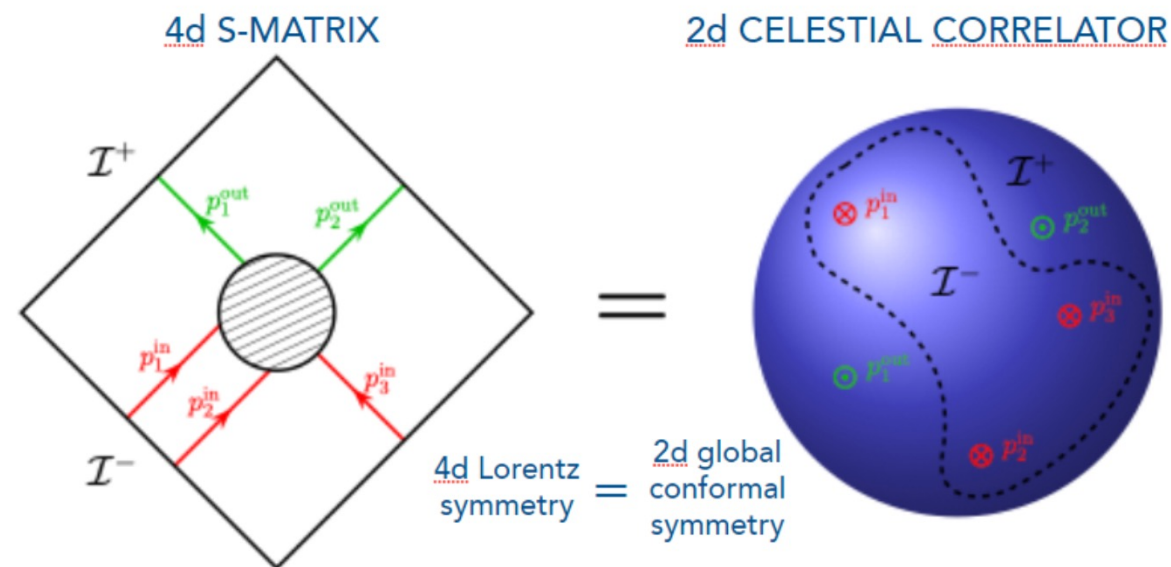
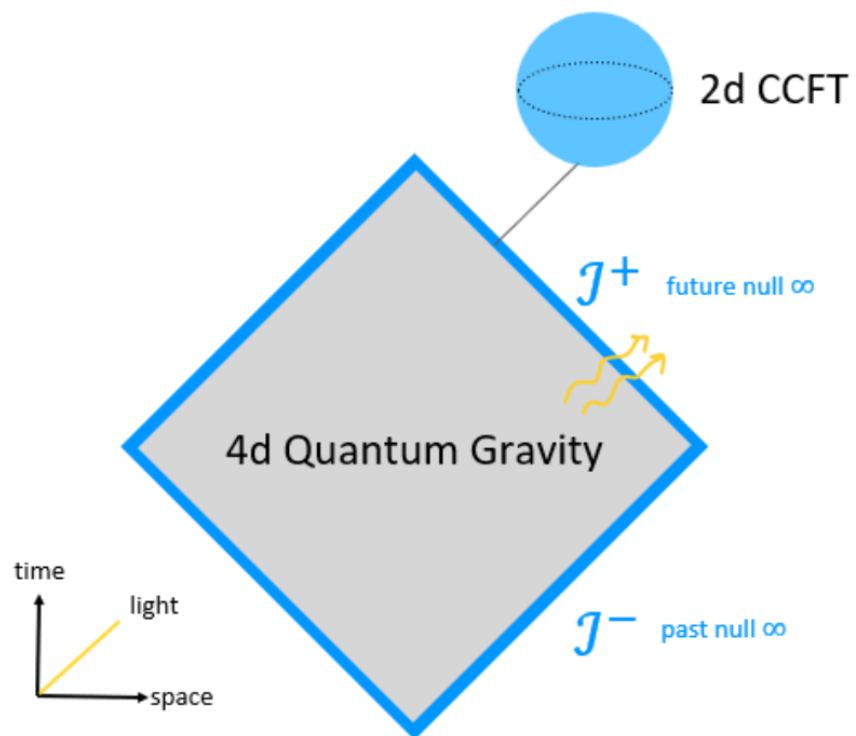
GQuEST

Gravity from the Quantum Entanglement of Space Time



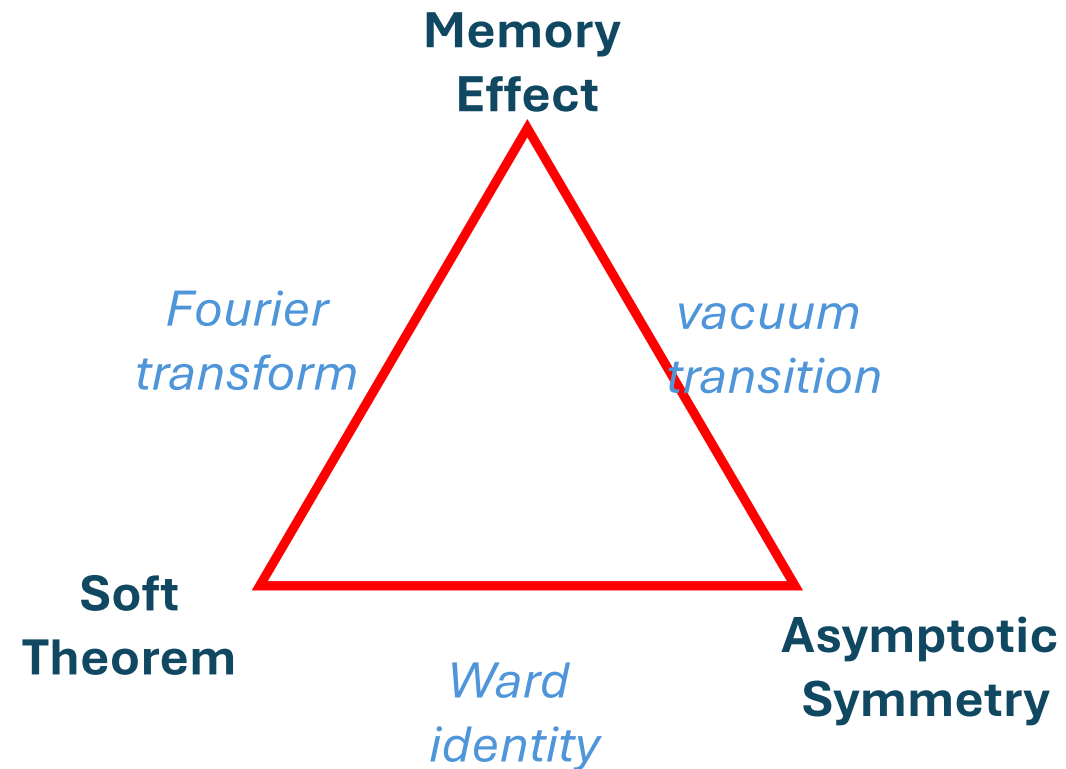
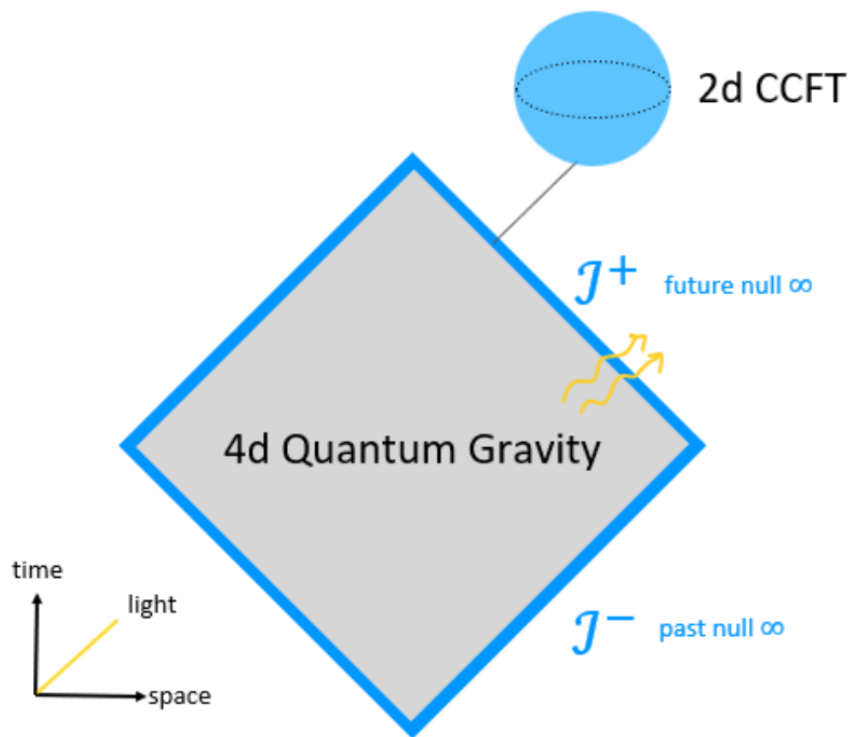
Relation with Celestial Holography

*talks by Ana, Andrea, Temple
and Kathryn*

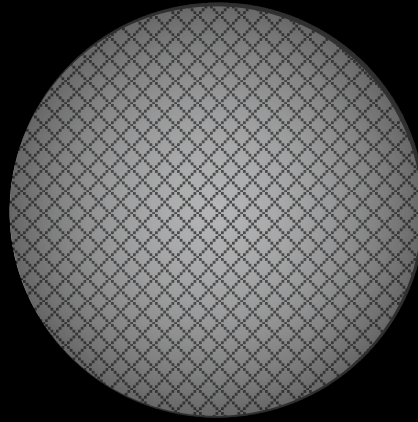


Relation with Celestial Holography

*talks by Ana, Andrea, Temple
and Kathryn*



The shockwave geometries are related to the gravitational memory



$$S = k_B \frac{Ac^3}{4\pi G \hbar}$$

$$k_B T = \frac{\hbar \kappa}{2\pi c}$$

Black hole thermodynamics relates Einstein equations
to 1st law of thermodynamics = derivable from microscopic theory

$$dM = \frac{\kappa}{2\pi} \frac{dA}{4G}$$

κ = surface gravity

$$dE = T dS$$

Emergent gravity = derivable
from microscopic theory!

Gravity and spacetime from quantum entanglement.

A change in the modular energy for a region of size **R**

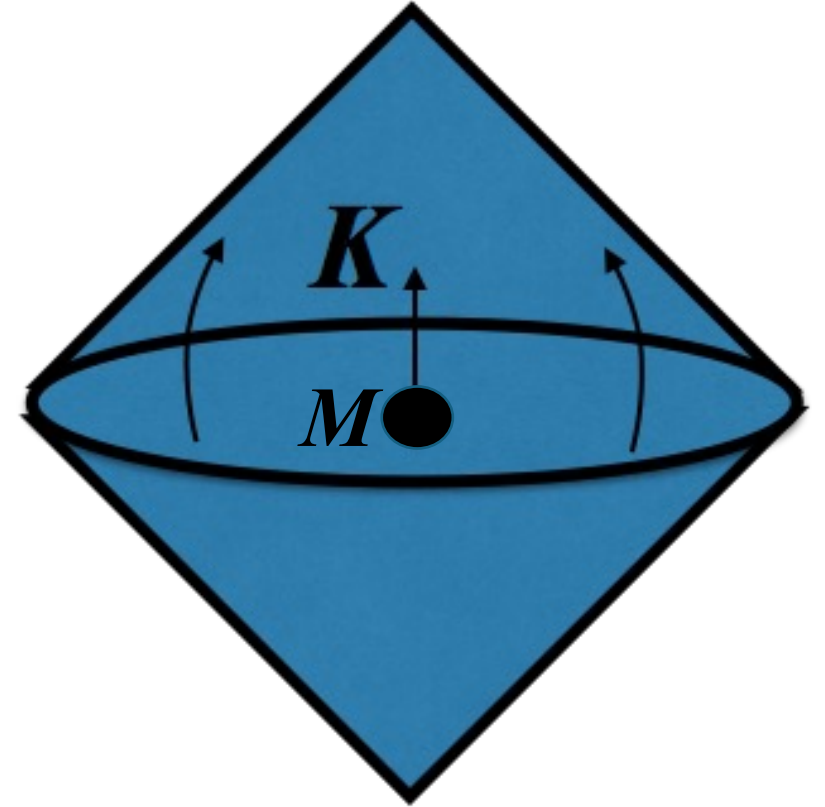
$$K = 2\pi \int d^n x \left(\frac{r^2 - |x|^2}{2r} \right) T_{00}(x)$$

gives a change in the entanglement entropy

$$\frac{\Delta A(R)}{4G} = -2\pi M R$$

=> the Einstein equations when the volume is kept fixed.

Can gravity and spacetime be derived purely from Quantum entanglement? What are the additional assumptions?



Ted Jacobson,
Manus Visser

What is the Microscopic Quantum description of de Sitter space

$$ds^2 = - \left(1 - \frac{R^2}{L^2} \right) dt^2 + \frac{dR^2}{1 - R^2/L^2} + R^2 d\Omega^2$$

What is the interpretation of the de Sitter entropy?

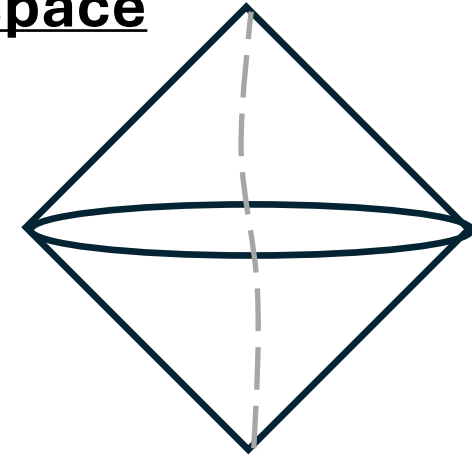
What is the microscopic explanation of

$$\Delta S = -2\pi M L$$

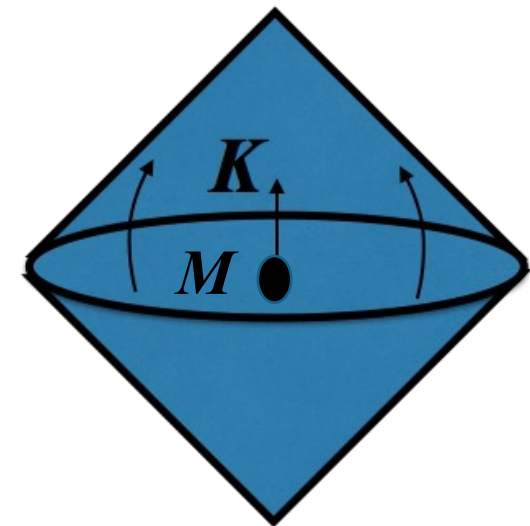
Can we again derive the gravitational equations?

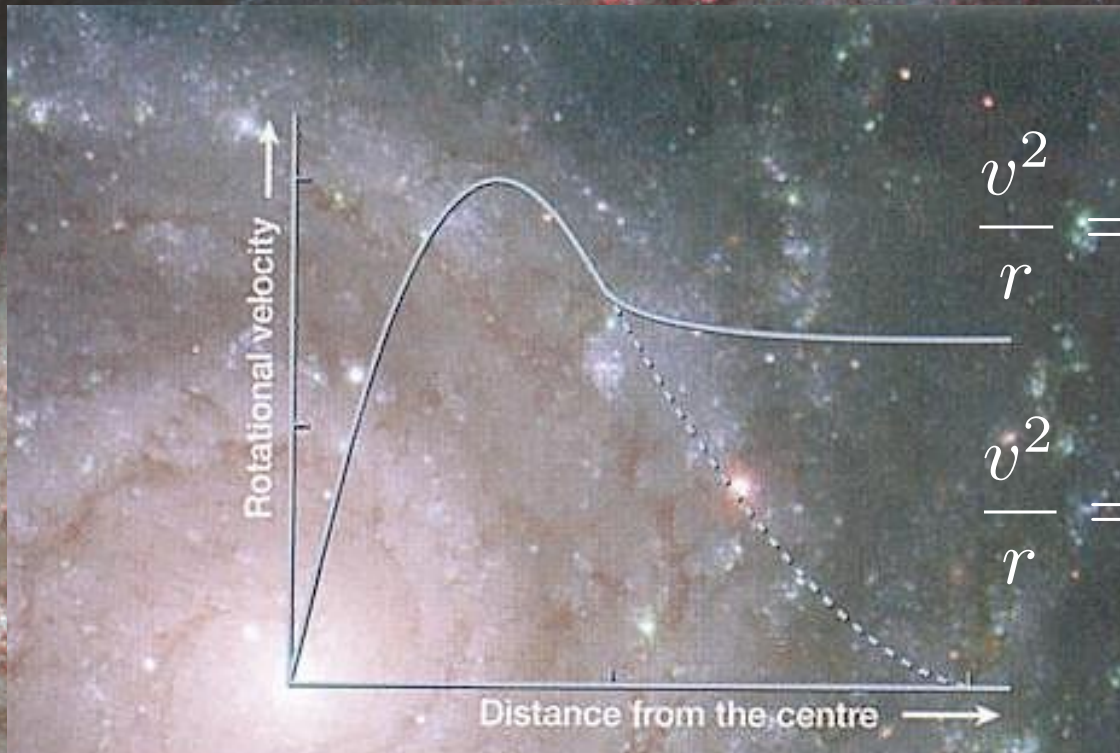
Why does gravity appear to behave differently when

$$2\pi M L \leq \frac{A(R)}{4G} \quad ?$$



$$S = \frac{A(L)}{4G}$$





$$\frac{v^2}{r} = \frac{GM_B}{r^2} + \frac{GM_D}{r^2}$$

$$\frac{v^2}{r} = \frac{GM_B}{r^2}$$

Empirically DM-effects appear when

$$\frac{GM}{r^2} < \frac{cH_0}{2}$$

$$\frac{Mc^2}{\hbar H_0 / 2\pi} < \frac{Ac^3}{4G\hbar}$$

THANK YOU

Matter entangles with Dark Energy

- The empirical fact

$$\frac{GM}{R^2} < \frac{cH_0}{2}$$

- implies that DM-effects appear when

$$2\pi \frac{McR}{\hbar} < \frac{A(R)c^3}{4G\hbar} \frac{R}{L}$$

- The left hand side is the entanglement entropy of matter.
- The right hand side represents the entropy contained in DE.

