Neutrino Physics: Part 1

The (non)Standard Model Particle

DESY Summer School 2024 - Dr. S. Blot

Overview

Part 1:

- Introduction to neutrinos
- Neutrino cross sections
- Sources of neutrinos
- Massive neutrinos and oscillations

Part 2:

- Overview of neutrino detection techniques
- Review current landscape and key measurements
- Open questions and future prospects













Total Energy Released in Decay





Offener Brief an die Gruppe der Radicaktiven bei der Gauvereins-Tagung zu Tubingen.

Abschrift

Physikalisches Institut der Eidg. Technischen Hochschule Zürich

Zirich, 4. Des. 1930 Cloriastrasse

Liebe Radioaktive Damen und Herren;

Wie der Ueberbringer dieser Zeilen, den ich huldvollst anguhören bitte, Ihnen des näheren auseinandersetsen wird, bin ich angesichts der "falschen" Statistik der No und Lie6 Kerne, sowie des kontinuierlichen beta-Spektrums auf einen versweifelten Ausweg verfallen um den "Wechselasts" (1) der Statistik und den Energiesats zu retten. Nämlich die Möglichkeit, es könnten elektrisch neutrale Teilchen, die ich Neutronen namen will, in den Kernen acistieran, welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und alma von Lichtquanten musserden noch dadurch unterscheiden, dass sie mant mit Lichtgeschwindigkeit laufen. Die Masse der Meutronen timate van derselben Grossenordnung wie die Elektronenmasse sein und interfalls night grösser als 0.01. Protonemasses- Das kontinuisrliche hatas Spektrum wäre dann verständlich unter der Annahme, dass beim bein-Zerfall mit den blektron jeweils noch ein Neutron emittiert wird, derart, dass die Summe der Energien von Neutron und Elektron konstant ist.

- Proposal of a new particle that carries away **E**, **p**
 - Electrically neutral
 - Spin-½
 - Very hard to detect

Offener Brief an die Gruppe der Radicaktiven bei der Gauvereins-Tagung zu Tübingen.

Abschrift

Physikalisches Institut der Eidg. Technischen Hochschuls Zürich

Zirich, 4. Des. 1930 Cloriastrasse

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich huldvollst anguhören bitte, Ihnen des näheren auseinandersetsen wird, bin ich angesichts der "falschen" Statistik der No und Lie6 Kerne, sowie des kontinuierlichen beta-Spektrums auf einen versweifelten Ausweg verfellen um den "Wechselsats" (1) der Statistik und den Energiesats zu retten. Nämlich die Möglichkeit, es könnten elektrisch neutrale Teilchen, die ich Neutronen namen will, in den Kernen existioren, welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und alma von Lichtquanten musserden noch dadurch unterscheiden, dass sie manht mit Lichtgeschwindigkeit laufen. Die Masse der Neutronen finate von derselben Grossenordnung wie die Elektronenmagse sein und tempfalls nicht grösser als 0.01 Protonemasse.- Das kontinuierliche hans Spertrum ware dann verständlich unter der Annahme, dass beim hate-Zerfall mit den blektron jeweils noch ein Neutron emittiert wirde derart, dass die Sume der Energien von Meutron und Elektron konstant ist.

- Proposal of a new particle that carries away **E**, **p**
 - Electrically neutral
 - Spin-1/2
 - Very hard to detect
- Community initially skeptical

Offener Brief an die Gruppe der Radicaktiven bei der Gauvereins-Tagung zu Tübingen.

Abschrift

Physikalisches Institut der Eidg. Technischen Hochschuls Zürich

Zirich, L. Des. 1930 Cloriastrasse

Liebe Radioaktive Damen und Herren,

Wie der Ueberbringer dieser Zeilen, den ich huldvollst anguhören bitte, Ihnen des näheren auseinandersetsen wird, bin ich angesichts der "falschen" Statistik der No und Lie6 Kerne, sowie des kontimuierlichen beta-Spektrums auf einen versweifelten Ausweg verfellen um den "Wechselsats" (1) der Statistik und den Energiesats zu retten. Nämlich die Möglichkeit, es könnten elektrisch neutrale Tailchen, die ich Neutronen namen will, in den Kornen atistierang welche den Spin 1/2 haben und das Ausschliessungsprinzip befolgen und she von Lichtquanten musserden noch dadurch unterscheiden, dass sie manht mit Lichtgeschwindigkeit laufen. Die Masse der Neutronen finate von derselben Grossenordnung wie die Elektronenmagse sein und inimfalls nicht grösser als 0,00 Protonemasse.- Das kontinuierliche hans Spertrum ware dann verständlich unter der Annahme, dass beim beth-Zerfall mit dem blektron jeweils noch ein Neutron emittiert wirde derart, dass die Sume der Energien von Meutron und Elektron konstant ist.

- Proposal of a new particle that carries away **E**, **p**
 - Electrically neutral
 - Spin-½
 - Very hard to detect
- Community initially skeptical
- Incorporated into theory of weak interactions by Enrico Fermi and renamed - *neutrino*

$$n \rightarrow p + e^{-} + \overline{V}_{e}$$

Offener Brief an die Grunpe der Radicaktiven bei der Gauvereins-Tagung zu Tübingen.





Mant nit Lichtgeschwindigkeit laufen. Die Masse der Meutronen Susste von derselben Grossenordnung wie die Elektronenmasse sein und Jesenfalls nicht grösser als 0,01 Protonenmasse.- Das kontinuierliche Beta- Spektrum wäre dann verständlich unter der Annahme, dass beim Deta-Zerfall mit dem blektron jeweils noch ein Neutron emittiert wärste derart, dass die Summe der Energien von Meutron und blektron konstent ist.

Crazy ideas inspire more crazy ideas...

Los Alamos proposal, 1940s

- Use nuclear bomb as an intense source of neutrinos
- Suspend a neutrino detector in a deep hole with vacuum
- Release detector when bomb goes off
- Hope it lands softly...



Discovery of the neutrino - 1956

Reines and Cowan - Project Poltergeist

- Use nuclear reactor as a neutrino source
- Capture neutrinos through inverse beta decay

 \overline{v}_{e}^{+} + p \rightarrow e⁺ + n

- Tank filled with water + cadmium chloride, monitored by light sensors
 - Prompt signal: $\mathbf{e}^+ + \mathbf{e}^- \rightarrow \mathbf{\gamma} + \mathbf{\gamma}$
 - $\circ \quad \text{Delayed signal: } ^{113}\text{Cd} + n \rightarrow {}^{114}\text{Cd} + \gamma$



Inverse β decay





- 3 known types
- Leptons with spin $\frac{1}{2}$
- No electric charge
- Only experience weak interactions



- 3 known types
- Leptons with spin $\frac{1}{2}$
- No electric charge
- Only experience weak interactions
 - \rightarrow W, Z bosons are heavy
 - Short-range ~10⁻¹⁸m
 - Suppresses interactions



- 3 known types
- Leptons with spin $\frac{1}{2}$
- No electric charge
- Only experience weak interactions
 - \rightarrow W, Z bosons are heavy
 - Short-range ~10⁻¹⁸m
 - Suppresses interactions
 - → Recall from SM lectures: weak interaction is maximally parity-violating!



Key properties:

- 3 known types
- Leptons with spin $\frac{1}{2}$
- No electric charge
- Only experience weak interactions •
 - W, Z bosons are heavy \rightarrow
 - Short-range $\sim 10^{-18}$ m
 - Suppresses interactions
 - Recall from SM lectures: weak interaction is \rightarrow maximally parity-violating!



DESY Summer School 2024: Neutrino Physics



- 3 known types
- Leptons with spin $\frac{1}{2}$
- No electric charge
- Only experience weak interactions
 - \rightarrow W, Z bosons are heavy
 - Short-range ~10⁻¹⁸m
 - Suppresses interactions
 - → Recall from SM lectures: weak interaction is maximally parity-violating!





Neutrinos interactions









DESY Summer School 2024: Neutrino Physics

Cross-section for CC elastic scattering of anti- v_+e



- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}}\left(rac{1}{(2\pi)^2}
ight)|{\cal M}|^2 d\Phi_2$$

Cross-section for CC elastic scattering of anti- v_+e



Using Feynman rules for scattering amplitude:

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$



Using Feynman rules for scattering amplitude:

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$



Using Feynman rules for scattering amplitude:

$$\mathcal{M}=\left(-rac{g}{2\sqrt{2}}
ight)ar{u}(p_3)\gamma^\mu(1-\gamma^5)u(p_4)$$
 \cdot

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$



Using Feynman rules for scattering amplitude:

$$\mathcal{M}=\left(-rac{g}{2\sqrt{2}}
ight)ar{u}(p_3)\gamma^\mu(1-\gamma^5)u(p_4)iggred{\cdot}iggred{-ig_{\mu
u}}{q^2-M_W^2}iggred{\cdot}$$

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$



Using Feynman rules for scattering amplitude:

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}}\left(rac{1}{(2\pi)^2}
ight)|{\cal M}|^2 d\Phi_2$$

$$\mathcal{M}=\left(-rac{g}{2\sqrt{2}}
ight)ar{u}(p_3)\gamma^\mu(1-\gamma^5)u(p_4)ig\cdot rac{-ig_{\mu
u}}{q^2-M_W^2}ig\cdot \left(-rac{g}{2\sqrt{2}}
ight)ar{u}(p_1)\gamma^
u(1-\gamma^5)u(p_2)$$



Using Feynman rules for scattering amplitude:

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$

$$\mathcal{M} = \left(-rac{g}{2\sqrt{2}}
ight) ar{u}(p_3) \gamma^\mu (1-\gamma^5) u(p_4) \cdot rac{-ig_{\mu
u}}{q^2 - M_W^2} \cdot \left(-rac{g}{2\sqrt{2}}
ight) ar{u}(p_1) \gamma^
u (1-\gamma^5) u(p_2)$$

When `q²` <<
$${\sf M}_{\sf W}^{-2}$$
 , then ${\cal M}=-rac{g^2}{8M_W^2}\cdot$ [matrix stuff]

Cross-section for CC elastic scattering of anti- v_+e



Using Feynman rules for scattering amplitude:

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$

$$\mathcal{M} = \left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_3)\gamma^{\mu}(1-\gamma^5)u(p_4) \cdot \left[\frac{-ig_{\mu\nu}}{q^2 - M_W^2}\right] \cdot \left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_1)\gamma^{\nu}(1-\gamma^5)u(p_2)$$

$$When q^2 << M_W^2 \text{, then } \mathcal{M} = -\frac{g^2}{8M_W^2} \cdot [\text{matrix stuff}] \qquad \text{Such that, } \mathcal{M} \propto -\frac{G_F}{\sqrt{2}} \qquad |\mathcal{M}|^2 \sim 10^{-10} \,\text{GeV}^{-4}$$

$$\sup_{g^2 = 4\sqrt{2}G_F M_W^2} \qquad 31$$



Using Feynman rules for scattering amplitude:

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$d\sigma = rac{1}{4|{f p}_{
m initial}|\sqrt{s}} \left(rac{1}{(2\pi)^2}
ight) |{\cal M}|^2 d\Phi_2$$

$$\mathcal{M} = \left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_3)\gamma^{\mu}(1-\gamma^5)u(p_4) \cdot \frac{-ig_{\mu\nu}}{q^2 - M_W^2} \cdot \left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_1)\gamma^{\nu}(1-\gamma^5)u(p_2) \quad \begin{array}{l} \text{Important!} \\ \text{picks out} \\ \text{left-handed} \\ \text{fields} \end{array}$$

$$\text{When } q^2 << M_W^2 \text{ , then } \mathcal{M} = -\frac{g^2}{8M_W^2} \cdot [\text{matrix stuff}] \quad \begin{array}{l} \text{Such that, } \mathcal{M} \propto -\frac{G_F}{\sqrt{2}} \\ \text{since...} \\ g^2 = 4\sqrt{2}G_F M_W^2 \end{array} \quad |\mathcal{M}|^2 \sim 10^{-10} \text{ GeV}^{-4}$$

Cross-section for CC elastic scattering of anti- v_e +e

Cross-Section (mb) For sake of time, have to skip a few steps in the calculations Excellent review with details can be \rightarrow found here: https://arxiv.org/abs/1305.7513 101 Key points: 10-19 10-22 10-25 10-28 10-3 10-2 10¹⁰ 1016 10-4 10² 1014 1018 10⁴ 10⁶ 108 10¹²

For reference, $\sigma(\gamma) \sim 1-10^6$ barn

Figure adapted from: Formaggio, Zeller

Neutrino Energy (eV)

Cross-section for CC elastic scattering of anti- v_{a} +e

Figure adapted from: Formaggio, Zeller Cross-Section (mb) For sake of time, have to skip a few steps in the calculations Excellent review with details can be found here: https://arxiv.org/abs/1305.7513 Key points: 10-19 \rightarrow at low energies, $\sigma \sim G_{F}^{2} s / \pi$ (linear) 10-22 and $\sigma(v-bar) / \sigma(v) \sim 1/2$ (due to spin) 10-25 10-28 10-31 10-2 1016 10-4 10² 10¹⁰ 104 10⁶ 1012 1014 1018 Neutrino Energy (eV)

For reference, $\sigma(\gamma) \sim 1-10^6$ barn

Cross-section for CC elastic scattering of anti- v_e +e

For sake of time, have to skip a few steps in the calculations Excellent review with details can be found here: https://arxiv.org/abs/1305.7513 Key points: 10-19 \rightarrow at low energies, $\sigma \sim G_{F}^{2} s / \pi$ (linear) 10-22 and $\sigma(v-bar) / \sigma(v) \sim 1/2$ (due to spin) 10-25 As energy grows, propagator matters \rightarrow 10-28 $\sigma_{\nu N}^{CC} = 5.53 \times 10^{-36} \text{ cm}^2 (\frac{E_{\nu}}{1 \text{ GeV}})^{\alpha},$ 10-31 10-2 104 10² 10⁶ $\sigma_{\nu N}^{NC} = 2.31 \times 10^{-36} \text{ cm}^2 (\frac{E_{\nu}}{1 \text{ GeV}})^{\alpha}, \quad \alpha \simeq 0.363.$ Neutrino Energy (eV) For reference, $\sigma(y) \sim 1-10^6$ barn

DESY Summer School 2024: Neutrino Physics

Figure adapted from: Formaggio, Zeller

Cross-section for CC elastic scattering of anti- v_e +e

- For sake of time, have to skip a few steps in the calculations
 - → Excellent review with details can be found here: https://arxiv.org/abs/1305.7513
 - Key points:
 - → at low energies, $\sigma \sim G_F^2 s / \pi$ (linear) and $\sigma(v-bar) / \sigma(v) \sim 1/2$ (due to spin)
 - \rightarrow As energy grows, propagator matters

$$\begin{split} \sigma_{\nu N}^{CC} &= 5.53 \times 10^{-36} \text{ cm}^2 (\frac{E_{\nu}}{1 \text{ GeV}})^{\alpha}, \\ \sigma_{\nu N}^{NC} &= 2.31 \times 10^{-36} \text{ cm}^2 (\frac{E_{\nu}}{1 \text{ GeV}})^{\alpha}, \quad \alpha \simeq 0.36 \end{split}$$

 \rightarrow Glashow resonance ~ 6 PeV

DESY Summer School 2024: Neutrino Physics



For reference, $\sigma(\gamma) \sim 1-10^6$ barn













Solar Neutrinos

- Fusion reactions in the sun generate anti-v_e through several mechanisms
- Dominant reaction is through proton-proton (pp) chain (99%)



Solar Neutrinos

- Fusion reactions in the sun generate anti-v_e through several mechanisms
- Dominant reaction is through proton-proton (pp) chain (99%)
- Carbon-Nitrogen-Oxygen (CNO) cycle is a subdominant process
 - \rightarrow More important for heavier, hotter stars



Solar Neutrinos

- Fusion reactions in the sun generate anti-v_e through several mechanisms
- Dominant reaction is through proton-proton (pp) chain (99%)
- Carbon-Nitrogen-Oxygen (CNO)
 cycle is a subdominant process
 → More important for heavier, hotter stars
- Standard Solar Model provides estimates of neutrino fluxes



• First solar model developed in 1960s by John Bahcall

- First solar model developed in 1960s by John Bahcall
- Prediction of anti-v_e flux higher than experimental measurements



- First solar model developed in 1960s by John Bahcall
- Prediction of anti-v_e flux higher than experimental measurements
- Confirmed by several targets
 - \rightarrow Unlikely to be a cross-section error



- First solar model developed in 1960s by John Bahcall
- Prediction of anti-v_e flux higher than experimental measurements
- Confirmed by several targets
 - \rightarrow Unlikely to be a cross-section error

Once again, neutrinos are causing trouble...



• Neutrinos flavour states are a superposition of mass states

$$\ket{
u_lpha} = \sum_i U_{lpha i} \ket{
u_i}$$

 $\alpha = e, \mu, \tau$ (weak eigenstates) i = 1, 2, 3 (mass eigenstates)

• Neutrinos flavour states are a superposition of mass states

$$\ket{
u_lpha} = \sum_i U_{lpha i} \ket{
u_i}$$

 $\alpha = e, \mu, \tau$ (weak eigenstates) i = 1, 2, 3 (mass eigenstates)



• Neutrinos flavour states are a superposition of mass states

$$|
u_{lpha}\rangle = \sum_{i} U_{lpha i} \, |
u_i\rangle$$
 $lpha$ = e, µ, T (weak eigenstates)
i = 1, 2, 3 (mass eigenstates)

• Connected by 3x3 matrix - up to 9 x 2 = 18 parameters (real + imaginary)

$$U = egin{bmatrix} U_{e1} & U_{e2} & U_{e3} \ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \ U_{ au 1} & U_{ au 2} & U_{ au 3} \end{bmatrix}$$

• Neutrinos flavour states are a superposition of mass states

$$|
u_{lpha}
angle = \sum_{i} U_{lpha i} \, |
u_i
angle$$
 $lpha$ = e, µ, т (weak eigenstates)
i = 1, 2, 3 (mass eigenstates)

- Connected by 3x3 matrix up to 9 x 2 = 18 parameters (real + imaginary)
- Unitarity constraints reduce dimensionality:

$$U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix} \qquad U^{\dagger} U = 1 \rightarrow U_{i} \cdot U_{j} = \delta_{ij} \text{ and } U^{\top}_{i} \cdot U^{\top}_{j} = \delta_{ij}$$

i.e. $|U_{e1}|^{2} + |U_{\mu 1}|^{2} + |U_{\tau 1}|^{2} = 1$

The Pontecorvo-Maki-Nakagawa-Sakata matrix

• PMNS matrix is most widely used parameterization

$$U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}$$
 Where cij = cos θ ij and sij = cos θ ij
Free parameters: θ_{12} , θ_{13} , θ_{23} , δ_{CP} (maybe $\alpha_1 \& \alpha_2$)

$$= \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

"atmospheric" "reactor" "solar" "Majorana"

The Pontecorvo-Maki-Nakagawa-Sakata matrix

• PMNS matrix is most widely used parameterization

$$U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}$$
 Where cij = cos θ ij and sij = cos θ ij

$$= \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$= \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & c_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

• To understand how this helps answer the Solar neutrino problem, we apply Schrödinger's equation and get a plane wave solution

 $|\,
u_j(t)\,
angle = e^{-i\,igl(\,E_jt\,-\,ec p_j\cdotec x\,igr)}\,|\,
u_j(0)\,
angle$

• To understand how this helps answer the Solar neutrino problem, we apply Schrödinger's equation and get a plane wave solution

• Ultra-relativistic limit: $|ec{p}_j| = p_j \gg m_j$ $E_j = \sqrt{p_j^2 + m_j^2} \simeq p_j + rac{m_j}{2\,p_j} pprox E + rac{m_j}{2\,E}$

• To understand how this helps answer the Solar neutrino problem, we apply Schrödinger's equation and get a plane wave solution

$$|\,
u_j(t)\,
angle = e^{-i\,igl(\,E_jt\,-\,ec p_j\cdotec x\,igr)}\,|\,
u_j(0)\,
angle$$

ullet Ultra-relativistic limit: , $\left|ec{p}_{j}
ight|=p_{j}\gg m_{j}$

$$E_j = \sqrt{p_j^2 + m_j^2} \simeq p_j + rac{m_j^2}{2\,p_j} pprox E + rac{m_j^2}{2\,E}$$

9

• Substituting this, with t ~ L (distance)

$$|\,
u_j(L)\,
angle = e^{-i\left(rac{m_j^2\,L}{2\,E}
ight)}\,|\,
u_j(0)\,
angle$$

9

• To understand how this helps answer the Solar neutrino problem, we apply Schrödinger's equation and get a plane wave solution

$$|\,
u_j(t)\,
angle = e^{-i\,igl(\,E_jt\,-\,ec{p}_j\cdotec{x}\,igr)}\,|\,
u_j(0)\,
angle$$

- Ultra-relativistic limit: , $\left| ec{p}_{j}
 ight| = p_{j} \gg m_{j}$ $E_{j} = \sqrt{p_{j}^{2} + m_{j}^{2}} \simeq p_{j} + rac{m_{j}^{2}}{2 \, p_{i}} pprox E + rac{m_{j}^{2}}{2 \, E}$
- Substituting this, with t ~ L (distance)

$$\begin{array}{l} |\nu_{j}(L)\rangle = e^{-i\left(\frac{m_{j}^{2}L}{2E}\right)} & \text{Transition probability for } \alpha \to \beta \\ |\nu_{j}(L)\rangle = e^{-i\left(\frac{m_{j}^{2}L}{2E}\right)} |\nu_{j}(0)\rangle & P_{\alpha \to \beta} = \left|\left\langle \nu_{\beta} \mid \nu_{\alpha}(L)\right\rangle\right|^{2} \\ & = \left|\sum_{j} U_{\alpha j}^{*} U_{\beta j} e^{-i\frac{m_{j}^{2}L}{2E}}\right|^{2} \end{array}$$

DESY Summer School 2024: Neutrino Physics

Vacuum oscillation probability

$$egin{aligned} P_{lpha o eta} &= \delta_{lphaeta} - 4\,\sum_{j>k}\,\mathcal{R}_e\Big\{\,U^*_{lpha j}\,U_{eta j}\,U_{lpha k}\,U^*_{eta k}\,\Big\}\,\sin^2\!\left(rac{\Delta_{jk}m^2\,L}{4E}
ight) \ &+ 2\,\sum_{j>k}\,\mathcal{I}_m\Big\{\,U^*_{lpha j}\,U_{eta j}\,U_{lpha k}\,U^*_{eta k}\,\Big\}\,\sin\!\left(rac{\Delta_{jk}m^2\,L}{2E}
ight), \end{aligned}$$
 where $\Delta_{jk}m^2\,\equiv m_j^2 - m_k^2$.

Vacuum oscillation probability

 $\alpha = \beta$ "appearance" or "survival" $\alpha \neq \beta$ "disappearance"

$$egin{aligned} P_{lpha o eta} &= \delta_{lphaeta} - 4\,\sum_{j>k}\,\mathcal{R}_e\Big\{\,U^*_{lpha j}\,U_{eta j}\,U_{lpha k}\,U^*_{eta k}\,\Big\}\,\sin^2\left(\!\left|rac{\Delta_{jk}m^2\,L}{4E}
ight) \ &+ 2\,\sum_{j>k}\,\mathcal{I}_m\Big\{\,U^*_{lpha j}\,U_{eta j}\,U_{lpha k}\,U^*_{eta k}\,\Big\}\,\sin\left(\!\left|rac{\Delta_{jk}m^2\,L}{2E}
ight), \end{aligned}$$
 where $\Delta_{jk}m^2\,\equiv m_j^2 - m_k^2$. Phases determined by squared mass splittings

$$\Delta m_{12}^2 \sim 10^{-5} \text{ eV}^2$$
 $\Delta m_{32}^2 \sim 10^{-3} \text{ eV}^2$

Vacuum oscillation probability

 $\alpha = \beta$ "appearance" or "survival" $\alpha \neq \beta$ "disappearance"

$$P_{\alpha \to \beta} = \delta_{\alpha\beta} - 4 \sum_{j > k} \mathcal{R}_e \left\{ U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \right\} \sin^2 \left(\frac{\Delta_{jk} m^2 L}{4E} \right) \\ + 2 \sum_{j > k} \mathcal{I}_m \left\{ U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \right\} \sin \left(\frac{\Delta_{jk} m^2 L}{2E} \right),$$
where $\Delta_{jk} m^2 \equiv m_j^2 - m_k^2$. Phases determined by squared mass splittings
$$\Delta m_{12}^2 \sim 10^{-5} \text{ eV}^2 \qquad \Delta m_{32}^2 \sim 10^{-3}$$

1

Vacuum oscillation probability - electron neutrino



Vacuum oscillation probability - muon neutrino



Solving the solar neutrino problem

- Solar neutrino detectors were only sensitive to antive
- Low anti-ve survival probability when reaching Earth
- The SNO experiment was designed to measure both anti-ve and NC (all)
 - $\rightarrow \ \ NC \ rates \ match \\ expectation!$



2015



Neutrino mixing vs Quark mixing

- Quarks also mix in weak interactions, governed by the CKM matrix
- How is the situation different for neutrinos?

Neutrino mixing vs Quark mixing

- Quarks also mix in weak interactions, governed by the CKM matrix
- How is the situation different for neutrinos?



Key takeaways from today

- Neutrinos are Standard Model odd-balls
- Extremely low interaction rates
- Naturally produced in abundance by many sources
- Unexpectedly massive



DESY Summer School 2024: Neutrino Physics