Neutrino Physics: Part 1

The (non)Standard Model Particle

DESY Summer School 2024 - Dr. S. Blot

Overview

Part 1:

- **Introduction to neutrinos**
- **Neutrino cross sections**
- **Sources of neutrinos**
- **Massive neutrinos and oscillations**

Part 2:

- Overview of neutrino detection techniques
- Review current landscape and key measurements
- Open questions and future prospects

Let's go back in time…. 1920s TREATY OF PEACE THE ALLIED AND ASSOCIATED POWERS

Meanwhile…

The Protoco

FRANCE

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Offener Brief an die Gruppe der Radioaktiven bei der Gauversins-Tagung zu Nibingen.

Abschrift

Physikalisches Institut der Eidg. Technischen Hochschule Zurich

Zürich. L. Dez. 1930 **Cloriastrasse**

Liebe Radioaktive Daman und Harren.

Wie der Ueberbringer dieser Zeilen, dan ich huldvollst ansuhören bitte, Ihnen des näheren auseinandersetsen wird, bin ich angesichts der "falschen" Statistik der N- und Li-6 Kerne, sowie des kontinuierlichen beta-Spektrums auf einen vermweifelten Ausweg verfallen um den "Wechselsats" (1) der Statistik und den Energiesats zu retten. Månlich die Möglichkeit, es könnten elektrisch neutrale Teilchen, die ich Neutronen nammen will, in den Kernen existieren, welche den Spin 1/2 haben und das Ausschliessungsprinsip befolgen und aid von Lichtquanten unseerden noch dadurch unterscheiden, dass sie misk mit Lichtgeschwindigkeit laufen. Die Masse dar Neutronen mante von derselben Grossenordnung wie die Elektronenmasse sein und indenfalls nicht grösser als O,00. Protonemasse.- Das kontinuierliche hote- Spektrum wäre dann verständlich unter der Aunahme, dass beim betm-Zerfall mit den klektron jeweils noch ein Neutron emittiert wird, derart, dass die Summe der Energien von Neutron und klektron konstant ist.

- Proposal of a new particle that carries away **E, p**
	- Electrically neutral
	- \circ Spin- $\frac{1}{2}$
	- Very hard to detect

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- Proposal of a new particle that carries away **E, p**
	- Electrically neutral
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- Community initially skeptical
- Incorporated into theory of weak interactions by Enrico Fermi and renamed - *neutrino*

$$
n \rightarrow p + e^- + \overline{V}_e
$$

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Abschrift

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Crazy ideas inspire more crazy ideas…

Los Alamos proposal, 1940s

- Use nuclear bomb as an intense source of neutrinos
- Suspend a neutrino detector in a deep hole with vacuum
- Release detector when bomb goes off
- Hope it lands softly...

Discovery of the neutrino - 1956

Reines and Cowan - Project Poltergeist

- Use nuclear reactor as a neutrino source
- Capture neutrinos through inverse beta decay

 \overline{V}_e + p \rightarrow e⁺ + n

- Tank filled with water $+$ cadmium chloride, monitored by light sensors
	- Prompt signal: **e+ + e- → γ + γ**
	- Delayed signal: **113Cd + n → 114Cd + γ**

Inverse β decay

- 3 known types
- Leptons with spin $\frac{1}{2}$
- No electric charge
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Key properties:

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Neutrinos interactions

- Particle energy and momentum
- Type of scattering process
- Phase space available for final state

$$
d\sigma=\frac{1}{4|\mathbf{p}_{\text{initial}}|\sqrt{s}}\left(\frac{1}{(2\pi)^2}\right)|\mathcal{M}|^2d\Phi_2
$$

Using Feynman rules for scattering amplitude:

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\mathcal{M}=\boxed{\left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_3)\gamma^{\mu}(1-\gamma^5)u(p_4)\hspace{0.1cm}\right]\hspace{0.1cm}-ig_{\mu\nu}}\hspace{0.1cm}\cdot\hspace{0.1cm}\boxed{-ig_{\mu\nu}\over q^2-M_W^2}.
$$

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$$

$$
\mathcal{M} = \boxed{ \left(- \frac{g}{2 \sqrt{2}} \right) \bar{u}(p_3) \gamma^{\mu} (1 - \gamma^5) u(p_4) \cdot \boxed{-ig_{\mu\nu}}{q^2 - M_W^2}} \cdot \boxed{ \left(- \frac{g}{2 \sqrt{2}} \right) \bar{u}(p_1) \gamma^{\nu} (1 - \gamma^5) u(p_2) }
$$

Using Feynman rules for scattering amplitude:

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$$

When
$$
q^2 \ll M_{\gamma V}^2
$$
, then $\mathcal{M} = -\frac{g^2}{8M_W^2}$ (matrix stuff]

Using Feynman rules for scattering amplitude:

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$$
d\sigma = \frac{1}{4|\mathbf{p}_{\text{initial}}|\sqrt{s}}\left(\frac{1}{(2\pi)^2}\right)|\mathcal{M}|^2 d\Phi_2
$$

$$
\mathcal{M} = \underbrace{\left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_3)\gamma^{\mu}(1-\gamma^5)u(p_4)}_{\text{When }q^2<< M_{\text{W}}^2} \cdot \underbrace{\left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_1)\gamma^{\nu}(1-\gamma^5)u(p_2)}_{\text{Since...}}\right]
$$
\n
$$
\text{When }q^2<< M_{\text{W}}^2 \text{ , then } \mathcal{M} = -\frac{g^2}{8M_W^2} \cdot \text{[matrix stuff]} \quad \text{Such that, } \mathcal{M} \propto -\frac{G_F}{\sqrt{2}} \quad |\mathcal{M}|^2 \sim 10^{-10} \text{ GeV}^{-4}
$$
\n
$$
\text{Since...} \quad g^2 = 4\sqrt{2}G_F M_W^2 \quad \text{31}
$$

Using Feynman rules for scattering amplitude:

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$$
d\sigma = \frac{1}{4|\mathbf{p}_{\text{initial}}|\sqrt{s}}\left(\frac{1}{(2\pi)^2}\right)|\mathcal{M}|^2 d\Phi_2
$$

$$
M = \underbrace{\left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_3)\gamma^{\mu}(1-\gamma^5)u(p_4)}_{\text{W}}\cdot \underbrace{\left(-\frac{g}{2\sqrt{2}}\right)\bar{u}(p_1)\gamma^{\nu}\left(1-\gamma^5\right)u(p_2)}_{\text{fileds}}\quad \text{picks out} \text{with} \quad \text{fields} \quad \text{fields} \quad \text{We have} \quad \text{the second} \quad \text{fields} \quad \text{where} \quad \text{the second} \quad \text{fields} \quad \text{where} \quad \text{the second} \quad \text{folds} \quad \text{folds} \quad \text{where} \quad \text{g}^2 = 4\sqrt{2}G_F M_W^2 \quad \text{where} \quad \text{g}^2 = 4\sqrt{2}G_F M_W^2 \quad \text{where} \quad \text{g}^2 = 4\sqrt{2}G_F M_W^2 \quad \text{and} \quad \text{g}^2 = 4\sqrt{2}G_F M_W^2 \quad \text{and} \quad \text{g}^2 = 4\sqrt{2}G_F M_W^2 \quad \text{and} \quad \text{higgs} \quad \text{where} \quad \text{the second} \quad \text{folds} \quad \text{g}^2 = 4\sqrt{2}G_F M_W^2 \quad \text{g}^2 =
$$

Figure adapted from: [Formaggio, Zeller](https://arxiv.org/abs/1305.7513) ● For sake of time, have to skip a few Cross-Section (mb) $10¹$ steps in the calculations \rightarrow Excellent review with details can be found here: 10^{-1} <https://arxiv.org/abs/1305.7513> 10^{-1} Key points: 10^{-19} 10^{-22} 10^{-25} 10^{-28} 10^{-3} 10^{16} $10⁴$ 10^{-2} $10²$ 10^{10} 10^{14} 10^{18} $10⁴$ $10⁶$ $10⁸$ 10^{12}

For reference, σ (γ) ~ 1-10 $^{\circ}$ barn

Neutrino Energy (eV)

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Figure adapted from: [Formaggio, Zeller](https://arxiv.org/abs/1305.7513) For sake of time, have to skip a few 2 ross-Section (mb)
 $\frac{1}{2}$ a a a a a steps in the calculations Excellent review with details can be found here: <https://arxiv.org/abs/1305.7513> Key points: 10^{-19} → at low energies, **σ** ~ **G**_F²**s** / **π** (linear) 10^{-22} and σ (**v**-bar) / σ (**v**) \sim 1/2 (due to spin) 10^{-25} \rightarrow As energy grows, propagator matters 10^{-28} $\sigma_{\nu N}^{CC}=5.53\times 10^{-36} \ {\rm cm}^2 (\frac{E_{\nu}}{1 \ {\rm GeV}})^{\alpha},$ 10^{-31} 10^{-2} 10^{16} $10⁻⁴$ 10^{18} $10²$ $10⁴$ $10⁶$ 10^{10} 10^{12} $\sigma_{\nu N}^{NC} = 2.31 \times 10^{-36} \text{ cm}^2 \left(\frac{E_{\nu}}{1 \text{ GeV}}\right)^{\alpha}, \ \ \alpha \simeq 0.363.$ **Neutrino Energy (eV)**

For reference, σ (γ) ~ 1-10 $^{\circ}$ barn

- For sake of time, have to skip a few steps in the calculations
	- Excellent review with details can be found here: <https://arxiv.org/abs/1305.7513>
	- Key points:
		- → at low energies, **σ** ~ **G**_F²**s** / **π** (linear) and σ (**v**-bar) / σ (**v**) \sim 1/2 (due to spin)
		- \rightarrow As energy grows, propagator matters

$$
\sigma_{\nu N}^{CC} = 5.53 \times 10^{-36} \text{ cm}^2 \left(\frac{E_{\nu}}{1 \text{ GeV}}\right)^{\alpha},
$$

$$
\sigma_{\nu N}^{NC} = 2.31 \times 10^{-36} \text{ cm}^2 \left(\frac{E_{\nu}}{1 \text{ GeV}}\right)^{\alpha}, \quad \alpha \simeq 0.363
$$

Glashow resonance \sim 6 PeV

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For reference, σ (γ) ~ 1-10⁶ barn

Solar Neutrinos

- Fusion reactions in the sun generate anti-v_e through several mechanisms
- Dominant reaction is through *proton-proton (pp)* **chain** (99%)

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- **● Carbon-Nitrogen-Oxygen (CNO)** cycle is a subdominant process
	- More important for heavier, hotter stars

Solar Neutrinos

- Fusion reactions in the sun generate anti-v_e through several mechanisms
- Dominant reaction is through *proton-proton (pp)* **chain** (99%)
- **● Carbon-Nitrogen-Oxygen (CNO)** cycle is a subdominant process More important for heavier, hotter stars
- Standard Solar Model provides estimates of neutrino fluxes

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Once again, neutrinos are causing trouble…

• Neutrinos flavour states are a superposition of mass states

$$
\ket{\nu_{\alpha}} = \sum_i U_{\alpha i} \ket{\nu_i}
$$

 α = e, μ , τ (weak eigenstates) $i = 1, 2, 3$ (mass eigenstates)

● Neutrinos flavour states are a superposition of mass states

$$
|\nu_{\alpha}\rangle = \sum_{i} U_{\alpha i} | \nu_{i} \rangle \qquad \qquad \alpha = e, \mu, \tau \text{ (weak eigenstates)}
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 $i = 1, 2, 3 \text{ (mass eigenstates)}$

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$$
\n $\alpha = e, \mu, \tau \text{ (weak eigenstates)}$ \n $i = 1, 2, 3 \text{ (mass eigenstates)}$

• Connected by 3x3 matrix - up to $9 \times 2 = 18$ parameters (real + imaginary)

$$
U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}
$$

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\n
$$
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\n
$$
i = 1, 2, 3 \text{ (mass eigenstates)}
$$

- Connected by 3x3 matrix up to 9 x 2 = 18 parameters (real + imaginary)
- Unitarity constraints reduce dimensionality:

$$
U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix} \qquad \text{UtU} = 1 \rightarrow U_{\text{i}} \cdot U_{\text{j}} = \delta_{\text{ij}} \text{ and } U_{\text{i}}^{\text{T}} \cdot U_{\text{j}}^{\text{T}} = \delta_{\text{ij}}
$$

The Pontecorvo-Maki-Nakagawa-Sakata matrix

● PMNS matrix is most widely used parameterization

$$
U = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{bmatrix}
$$
 Free parameters: $\theta_{12}, \theta_{13}, \theta_{23}, \delta_{CP}$ (maybe $\alpha_1 \& \alpha_2$)
\n
$$
= \begin{bmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{bmatrix} \begin{bmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta} & 0 & c_{13} \end{bmatrix} \begin{bmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}
$$

\n"atmospheric" "reactor" "solar" "Majorana"

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$$
\n
$$
= \begin{bmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{bmatrix} \begin{bmatrix} e^{i\alpha_1/2} & 0 & 0 \\ 0 & e^{i\alpha_2/2} & 0 \\ 0 & 0 & 1 \end{bmatrix}
$$

● To understand how this helps answer the Solar neutrino problem, we apply Schrödinger's equation and get a plane wave solution

 $\mid \nu_j(t)\,\rangle = e^{-i\,\left(\,E_j t-\vec{p}_j\cdot\vec{x}\,\right)}\,\left|\, \nu_j(0)\,\right\rangle$

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\mid \nu_j(t) \,\rangle = e^{-i \, \left(\,E_j t \,-\, \vec{p}_j \cdot \vec{x} \,\right)} \, \mid \nu_j(0) \, \rangle
$$

\n- Ultra-relativistic limit:
$$
|p_j| = p_j \gg m_j
$$
\n- $E_j = \sqrt{p_j^2 + m_j^2} \simeq p_j + \frac{J}{2\,p_j} \approx E + \frac{J}{2\,E}$
\n

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$$
|\,\nu_j(t)\,\rangle=e^{-i\,\left(\,E_jt-\,\vec{p}_j\cdot\vec{x}\,\right)}\,|\,\nu_j(0)\,\rangle
$$

 \bullet Ultra-relativistic limit: $|\vec{p}_j| = p_j \gg m_j$

$$
E_j=\sqrt{p_j^2+m_j^2}\simeq p_j+\frac{m_j^2}{2\,p_j}\approx\boxed{E+\frac{m_j^2}{2\,E}}
$$

• Substituting this, with $t \sim L$ (distance)

$$
|\,\nu_j(L)\,\rangle=e^{-i\,\left(\frac{m_j^2\,L}{2\,E}\right)}\,\left|\,\nu_j(0)\,\right>
$$

● To understand how this helps answer the Solar neutrino problem, we apply Schrödinger's equation and get a plane wave solution

$$
|\,\nu_j(t)\,\rangle=e^{-i\,\left(\,E_jt-\,\vec{p}_j\cdot\vec{x}\,\right)}\,|\,\nu_j(0)\,\rangle
$$

- Ultra-relativistic limit: $\left|\vec{p}_j\right|=p_j\gg m_j$ $E_j=\sqrt{p_j^2+m_j^2}\simeq p_j+\dfrac{m_j^2}{2\,p_i}\approx\left|E+\dfrac{m_j^2}{2\,E}\right|$
- Substituting this, with $t \sim L$ (distance)

$$
\displaystyle \left| \; \nu_j (L) \, \right\rangle = e^{-i \left(\frac{m_j^2 \, L}{2 \, E} \right)} \, \left| \; \nu_j (0) \, \right\rangle \qquad \qquad P_{\alpha \to \beta} = \, \left| \; \left< \, \nu_\beta \, \right| \, \nu_\alpha (L) \, \right\rangle \, \left|^2 \\[0.4em] \qquad \qquad = \, \left| \; \sum_j \, U_{\alpha j}^* \, U_{\beta j} \, e^{-i \frac{m_j^2 \, L}{2 E}} \, \right|^2 \\[0.4em] \qquad \qquad = \, \left| \; \sum_j \, U_{\alpha j}^* \, U_{\beta j} \, e^{-i \frac{m_j^2 \, L}{2 E}} \, \right|^2
$$

Vacuum oscillation probability

$$
\begin{aligned} P_{\alpha\to\beta} &= \delta_{\alpha\beta} - 4 \, \sum_{j>k} \, \mathcal{R}_e \Big\{ \, U_{\alpha j}^* \, U_{\beta j} \, U_{\alpha k} \, U_{\beta k}^* \, \Big\} \, \sin^2 \! \left(\frac{\Delta_{jk} m^2 \, L}{4E} \right) \\ &+ 2 \, \sum_{j>k} \, \mathcal{I}_m \Big\{ \, U_{\alpha j}^* \, U_{\beta j} \, U_{\alpha k} \, U_{\beta k}^* \, \Big\} \, \sin \! \left(\frac{\Delta_{jk} m^2 \, L}{2E} \right), \end{aligned}
$$
 where $\Delta_{jk} m^2 \ \equiv m_j^2 - m_k^2$.

Vacuum oscillation probability

 $\alpha = \beta$ "appearance" or "survival" α ≠ β "disappearance"

$$
P_{\alpha \to \beta} = \delta_{\alpha \beta} - 4 \sum_{j > k} \mathcal{R}_e \Big\{ U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \Big\} \sin^2 \left(\frac{\Delta_{jk} m^2 L}{4E} \right)
$$

+
$$
2 \sum_{j > k} \mathcal{I}_m \Big\{ U_{\alpha j}^* U_{\beta j} U_{\alpha k} U_{\beta k}^* \Big\} \sin \left(\frac{\Delta_{jk} m^2 L}{2E} \right),
$$

where
$$
\Delta_{jk} m^2 \equiv m_j^2 - m_k^2.
$$
 Phases determined by squared mass splittings

$$
\Delta m_{12}^2 \sim 10^{-5} \text{ eV}^2
$$

$$
\Delta m_{32}^2 \sim 10^{-3} \text{ eV}^2
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Vacuum oscillation probability - electron neutrino

Vacuum oscillation probability - muon neutrino

Solving the solar neutrino problem

- Solar neutrino detectors were only sensitive to antiνe
- Low anti-ve survival probability when reaching Earth
- The SNO experiment was designed to measure both anti-νe and NC (all)
	- **→ NC rates match expectation!**

Theory

 $0.48 + 0.02$

SuperK

 $7.6 + 1.3$

 $2.56 + 0.23$

 $C1$

2015

Neutrino mixing vs Quark mixing

- Quarks also mix in weak interactions, governed by the CKM matrix
- How is the situation different for neutrinos?

Neutrino mixing vs Quark mixing

- Quarks also mix in weak interactions, governed by the CKM matrix
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Key takeaways from today

- **Neutrinos are Standard Model odd-balls**
- **Extremely low interaction rates**
- **Naturally produced in abundance by many sources**
- **Unexpectedly massive**

