# Introduction to Accelerator Physics

Part 3

Pedro Castro / Accelerator Physics Group (MPY) Zeuthen, 9th August 2024





#### Accelerator lectures framework in Summer Student Prog.

<u>12th Aug.</u>: Accelerator R&D at DESY, Anne Oppelt

<u>Today:</u> focus on synchrotrons and synchrotron technology

#### synchrotrons: machines for discoveries

Facility	Particle(s) discovered	Year of discovery	Nobel Price
SPEAR	charm quark	1974	1976
SPEAR	tau lepton	1975	1995
PETRA	gluon	1979	
S $ar{p}$ pS	$W^{\pm}$ , Z bosons	1983	1984
SLC, LEP	$N_{\rm v}=3$		
Tevatron	top quark	1995	
LHC	Higgs	2012	2013

#### Main HEP discoveries at synchrotrons in the last 50 years



#### Main HEP discoveries at synchrotrons in the last 50 years

![](_page_3_Figure_1.jpeg)

#### Scope of this lecture:

- 1. Synchrotrons: key components and their <u>challenges to reach high energies</u>:
  - Dipole magnetic fields
  - Superconducting dipoles
  - Quadrupole magnets to focus beams
- 2. Synchrotrons and Linear Accelerators:
  - Acceleration using radio-frequency electomagnetic fields

![](_page_4_Picture_7.jpeg)

![](_page_5_Figure_0.jpeg)

## Motion in electric and magnetic fields

Equation of motion under Lorentz Force

![](_page_6_Figure_2.jpeg)

# Motion in magnetic fields

if the electric field is zero ( $\vec{E} = 0$ ), then

![](_page_7_Figure_2.jpeg)

Magnetic fields do not change the particles energy

# **Motion in magnetic fields**

if the electric field is zero (E=0), then

$$\vec{F} = \frac{d\vec{p}}{dt} = q \cdot \vec{v} \times \vec{B}$$

$$E^2 = \vec{p}^2 c^2 + E_0^2 \qquad \text{energy-momentum relation in special relativity}}$$

$$\text{total energy} \qquad \qquad \text{energy at rest}$$

$$\text{momentum}$$

# Motion in magnetic fields

if the electric field is zero (E=0), then

$$\vec{F} = \frac{d\vec{p}}{dt} = q \cdot \vec{v} \times \vec{B}$$

$$E^{2} = \vec{p}^{2}c^{2} + E_{0}^{2}$$

$$E\frac{dE}{dt} = c^{2}\vec{p}\frac{d\vec{p}}{dt} = c^{2}q\vec{p}(\vec{v}\times\vec{B}) = c^{2}q|\vec{p}| |\vec{v}\times\vec{B}| \cos \phi = 0$$
since  $\vec{v}\times\vec{B} \perp \vec{v} \Rightarrow \phi = 90^{\circ}$ 

Magnetic fields do not change the particles energy, only electric fields do !

#### acceleration with DC electric fields

![](_page_10_Figure_1.jpeg)

#### In general:

- Static magnetic fields  $\rightarrow$  to guide (bend + focus) particle beams
- Static electric fields  $\rightarrow$  accelerate particle beams (low energy)
- Radio-frequency EM fields  $\rightarrow$  accelerate particle beams (high E)

![](_page_11_Picture_4.jpeg)

#### acceleration with RF (radio-frequency) electric fields

![](_page_12_Figure_1.jpeg)

![](_page_12_Figure_2.jpeg)

### **RF cavity basics: a cylindrical cavity**

![](_page_13_Figure_1.jpeg)

# **RF cavity basics: a cylindrical cavity**

![](_page_14_Figure_1.jpeg)

LC circuit (or resonant circuit) analogy:

![](_page_14_Figure_3.jpeg)

### **RF cavity basics: a cylindrical cavity**

![](_page_15_Figure_1.jpeg)

LC circuit (or resonant circuit) analogy:

![](_page_15_Figure_3.jpeg)

![](_page_16_Figure_0.jpeg)

### Equations for the electric and magnetic fields in a pill box cavity

![](_page_17_Picture_1.jpeg)

![](_page_18_Figure_0.jpeg)

![](_page_19_Figure_0.jpeg)

set of solutions with  $B_z = 0$  (that is,  $\vec{B}$  is transverse)  $E_{z} = E_{0}J_{m}\left(x_{mn}\frac{r}{R}\right)\cos m\theta \cos\left(\frac{p\pi}{l}z\right)e^{j\omega t}$   $E_{r} = -\frac{p\pi}{l}\frac{R}{x_{mn}}E_{0}J'_{m}\left(x_{mn}\frac{r}{R}\right)\cos m\theta \sin\left(\frac{p\pi}{l}z\right)e^{j\omega t}$   $E_{\theta} = -\frac{p\pi}{l}\frac{mR^{2}}{x_{mn}^{2}r}E_{0}J_{m}\left(x_{mn}\frac{r}{R}\right)\sin m\theta \sin\left(\frac{p\pi}{l}z\right)e^{j\omega t}$  $B_z = 0$  $B_{r} = -j\omega \frac{mR^{2}}{x_{mn}^{2} rc^{2}} E_{0}J_{m}\left(x_{mn}\frac{r}{R}\right) \sin m\theta \cos\left(\frac{p\pi}{l}z\right)e^{j\omega t}$  $B_{\theta} = -j\omega \frac{R}{x_{mn}c^{2}} E_{0}J'_{m}\left(x_{mn}\frac{r}{R}\right) \cos m\theta \cos\left(\frac{p\pi}{l}z\right)e^{j\omega t}$ indices: m = 0,1,2,...: number of full period variations in  $\theta$  of the fields n = 1,2,...: number of zeros of the axial field component in  $\vec{r}$ 

p = 0,1,2,...: number of half period variations in z of the fields

 $J_m$  : Bessel's functions  $J'_m$  : derivative of the Bessel's

$$x_{mn}$$
: n-th root of  $J_m$  (that is,  $J_m(x_{mn}) = 0$ )

а

ngular frequency : 
$$\omega = c \sqrt{\left(\frac{x_{mn}}{R}\right)^2 + \left(\frac{p\pi}{l}\right)^2}$$

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![](_page_20_Figure_0.jpeg)

set of solutions with 
$$B_z = 0$$
 (that is,  $\vec{B}$  is transverse)  

$$\begin{cases}
E_z = E_0 J_m \left( x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t} \\
E_r = -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 J'_m \left( x_{mn} \frac{r}{R} \right) \cos m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t} \\
E_\theta = -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left( x_{mn} \frac{r}{R} \right) \sin m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t} \\
B_z = 0 \\
B_r = -j\omega \frac{mR^2}{x_{mn}^2 rc^2} E_0 J_m \left( x_{mn} \frac{r}{R} \right) \sin m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t} \\
B_\theta = -j\omega \frac{R}{x_{mn}c^2} E_0 J'_m \left( x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t} \\
\text{indices:}
\end{cases}$$

m = 0, 1, 2, ...: number of full period variations in  $\theta$  of the fields n = 1,2,...: number of zeros of the axial field component in  $\vec{r}$ p = 0,1,2,...: number of half period variations in z of the fields

 $x_{mn}$ : n-th root of  $J_m$  (that is,  $J_m(x_{mn}) = 0$ )  $J_m$  : Bessel's functions  $J'_m$ : derivative of the Bessel's functions angular frequency :  $\omega = c \left| \left( \frac{x_{mn}}{R} \right)^2 + \left( \frac{p\pi}{I} \right)^2 \right|$ DESY.

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![](_page_21_Figure_0.jpeg)

set of solutions with  $B_z = 0$  (that is,  $\vec{B}$  is transverse)  $E_{z} = E_{0}J_{m}\left(x_{mn}\frac{r}{P}\right)\cos m\theta \cos\left(\frac{p\pi}{I}z\right)e^{j\omega t}$  $E_r = -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 \underline{J'_m} \left( x_{mn} \frac{r}{R} \right) \cos m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t}$  $E_{\theta} = -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left( x_{mn} \frac{r}{R} \right) \sin m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t}$  $B_{\pi} = 0$  $B_r = -j\omega \frac{mR^2}{x_{mn}^2 rc^2} E_0 J_m \left( x_{mn} \frac{r}{R} \right) \sin m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t}$  $\int B_{\theta} = -j\omega \frac{R}{x_{mn}c^2} E_0 J'_m \left( x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left( \frac{p\pi}{I} z \right) e^{j\omega t}$ indices:

m = 0,1,2,...: number of full period variations in  $\theta$  of the fields n = 1,2,...: number of zeros of the axial field component in  $\vec{r}$ p = 0,1,2,...: number of half period variations in z of the fields

 $J_m$  : Bessel's functions

 $x_{mn}$ : n-th root of  $J_m$  (that is,  $J_m(x_{mn}) = 0$ )

$$J'_m$$
: derivative of the Bessel's functions  
angular frequency:  $\omega = c \sqrt{\left(\frac{x_{mn}}{R}\right)^2 + \left(\frac{p\pi}{l}\right)^2}$ 

DESY.

 $J_m$  : Bessel's functions

![](_page_22_Figure_1.jpeg)

![](_page_23_Figure_0.jpeg)

set of solutions with 
$$B_z = 0$$
 (that is,  $\vec{B}$  is transverse)  
 $E_z = E_0 J_m \left( \frac{x_{mn}}{R} \frac{r}{R} \right) \cos m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t}$   
 $E_r = -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 J'_m \left( \frac{x_{mn}}{R} \frac{r}{R} \right) \cos m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t}$   
 $E_\theta = -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left( \frac{x_{mn}}{R} \frac{r}{R} \right) \sin m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t}$   
 $B_z = 0$   
 $B_r = -j\omega \frac{mR^2}{x_{mn}^2 rc^2} E_0 J_m \left( \frac{x_{mn}}{R} \frac{r}{R} \right) \sin m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t}$   
 $E_\theta = -j\omega \frac{R}{x_{mn}^2 rc^2} E_0 J_m \left( \frac{x_{mn}}{R} \frac{r}{R} \right) \cos m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t}$   
hdices:  
 $= 0, 1, 2, ...:$  number of full period variations in  $\theta$  of the fields

n = 1,2,...: number of zeros of the axial field component in  $\vec{r}$ 

p = 0,1,2,...: number of half period variations in z of the fields

 $J_{m} : \text{Bessel's functions} \qquad x_{mn} : \text{n-th root of } J_{m} \text{ (that is, } J_{m}(x_{mn}) = 0)$   $J'_{m} : \text{derivative of the Bessel's functions} \qquad \text{angular frequency} : \omega = c \sqrt{\left(\frac{x_{mn}}{R}\right)^{2} + \left(\frac{p\pi}{l}\right)^{2}}$ DESY.
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![](_page_24_Figure_0.jpeg)

set of solutions with  $B_z = 0$  (that is,  $\vec{B}$  is transverse)  $E_{z} = E_{0}J_{m}\left(x_{mn}\frac{r}{P}\right)\cos m\theta \cos\left(\frac{p\pi}{I}z\right)e^{j\omega t}$  $E_r = -\frac{p\pi}{l} \frac{R}{x_{mn}} E_0 J'_m \left( x_{mn} \frac{r}{R} \right) \cos m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t}$  $E_{\theta} = -\frac{p\pi}{l} \frac{mR^2}{x_{mn}^2 r} E_0 J_m \left( x_{mn} \frac{r}{R} \right) \sin m\theta \sin \left( \frac{p\pi}{l} z \right) e^{j\omega t}$  $B_{7} = 0$  $B_r = -j\omega \frac{mR^2}{x_{mn}^2 rc^2} E_0 J_m \left( x_{mn} \frac{r}{R} \right) \sin m\theta \cos \left( \frac{p\pi}{l} z \right) e^{j\omega t}$  $\int B_{\theta} = -j\omega \frac{R}{x_{mn}c^2} E_0 J'_m \left( x_{mn} \frac{r}{R} \right) \cos m\theta \cos \left( \frac{p\pi}{I} z \right) e^{j\omega t}$ indices: m = 0,1,2,...: number of full period variations in  $\theta$  of the fields n = 1,2,...: number of zeros of the axial field component in  $\vec{r}$ p = 0,1,2,...: number of half period variations in z of the fields

 $J_m$  : Bessel's functions

$$x_{mn}$$
: n-th root of  $J_m$  (that is,  $J_m(x_{mn}) = 0$ )

$$J'_m$$
 : derivative of the Bessel's functions

angular frequency : 
$$\omega = c \sqrt{\left(\frac{x_{mn}}{R}\right)^2 + \left(\frac{p\pi}{l}\right)^2}$$

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![](_page_25_Figure_0.jpeg)

- m = 0 : rotation symmetry of the fields n = 1 : no zeros of the axial field component in  $\vec{r}$ p = 0 : no variation in z of the fields
- $J_m$  : Bessel's functions
- $J'_m$ : derivative of the Bessel's functions

angular frequency : 
$$\omega = c \frac{x_{01}}{R}$$
  $x_{01} = 2.405$ 

![](_page_26_Figure_0.jpeg)

- m = 0 : rotation symmetry of the fields n = 1 : no zeros of the axial field component in  $\vec{r}$ p = 0 : no variation in z of the fields
- $J_m$  : Bessel's functions

 $J'_m$ : derivative of the Bessel's functions

angular frequency : 
$$\omega = c \frac{x_{01}}{R}$$
  $x_{01} = 2.405$ 

#### Pill box cavity: 3D visualisation of E and B

![](_page_27_Figure_1.jpeg)

Equations for the electric and magnetic fields in a pill box cavity

#### **Examples of pill box cavities**

DESY cavity (pill box)

![](_page_29_Picture_2.jpeg)

#### ADONE cavity 51 MHz (pill box) Frascati lab, Italy

![](_page_29_Picture_4.jpeg)

# ADONE cavity 51 MHz (pill box) **Examples of pill box cavities** Frascati lab, Italy ADONE in 1963, Laboratori Nazionali di Frascati, Italy Page 31

# Superconducting cavity used at DESY

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)

![](_page_31_Picture_2.jpeg)

	Free-electron LASer in Hamburg	0.3 km	DESY	2004-	?	e-	1.2 GeV
European <u>X</u> -ray <u>F</u> ree- <u>E</u> lectron <u>L</u> aser		3 km	DESY	2016-	?	e-	17.5 GeV
	International Linear Collider	30 km	?	?		e-/e+	2x250 GeV

# Superconducting cavity used at DESY

![](_page_32_Picture_1.jpeg)

material: pure Niobium

operating temperature: 2 K

accelerating field gradient: up to 35 MV/m

### **Cavities inside a cryostat**

![](_page_33_Picture_1.jpeg)

![](_page_34_Figure_0.jpeg)

# **Cavities inside a cryostat**

![](_page_35_Picture_1.jpeg)
## **Cavities inside an accelerator module (cryostat)**



module installation in FLASH (2004)

### **<u>Free-electron LAS</u>er in <u>Hamburg</u> (FLASH)**



### **<u>Free-electron LAS</u>er in <u>Hamburg</u> (FLASH)**



### 100 accelerator modules (cryostats) in XFEL



### Superconducting cavities at HERA

16 cavities 500 MHz



#### **Superconducting Particle Accelerator**

From 1992 to 2007, eight of these superconducting accelerator components were used in the 6.3-kilometre long storage ring HERA to accelerate electrons and their antiparticles, positrons.

Two four-cell cavities are arranged in one thermal vessel (cryostat). The cavities are made of the metal niobium which becomes superconducting at a temperature of minus 269 degrees Celsius. At this temperature, particles are accelerated almost without electric resistance and thus very efficiently with a very high electric alternating voltage which is injected in the middle between the cavities. During HERA operation, this cavity reached an accelerating gradient of 5 million volts per metre.

### Superconducting cavities at LEP

272 cavities 352 MHz



### **Superconducting cavities at LHC**

16 cavities 400 MHz



### **Other accelerators using superconducting cavities**

- 5 de-commissioned
- 11 in operation
- 4 in construction
- 10 in design phase

Total = 30

full list: <u>http://tesla-new.desy.de/srf\_accelerators</u>

### Superconducting cavity used in FLASH and in XFEL

Superconducting cavity used in FLASH (0.3 km) and in XFEL (3 km)





# **Accelerating field map**



## Is there a net acceleration?

Simulation of the fundamental mode: electric field lines



























Is there a net acceleration? ..... timing is the key



#### Is there a net acceleration? ..... timing is the key

for protrons,  $\beta < 1$ 

example: ESS (European Spallation Source), Lund, Sweden



# Superconducting cavity used at DESY



### **Frequently Asked Questions**

- 1) Why this shape?
- 2)

#### 3)

#### 4)

### **Multipacting mitigation in superconducting cavities**



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1) Why this shape? ..... to reduce/avoid multipacting

2) How to feed  $\vec{E}$  in?

3)

4)

### Superconducting cavity used in FLASH and in XFEL



### Fundamental mode coupler (input coupler)



- 1) Why this shape? ...... to reduce/avoid multipacting
- 2) How to feed  $\vec{E}$  in? ..... with input couplers
- 3) How to measure  $\vec{E}$  ?

4)

### Superconducting cavity used in FLASH and in XFEL



1) Why this shape? ...... to reduce/avoid multipacting

2) How to feed  $\vec{E}$  in? ...... with input couplers

3) How to measure  $\vec{E}$  ? ...... with pick up antennas

4) What are HOM couplers for?


1)	Why this shape?		to reduce/avoid multipacting
2)	How to feed $\vec{E}$ in?		with input couplers
3)	How to measure $ec{E}$	?	with pick up antennas

4) What are HOM couplers for? ..... to reduce HOM (wakefields)

## **Summing-up of this part**

Particle acceleration using radio-frequency fields:



## **MEDIA DATABASE.** "Electron acceleration – a virtual simulation"



DESY→Press→Media database→European XFEL (with filter: media type=movies)

https://media.desy.de/DESYmediabank/?l=en#l=en&cid=3980&cname=European%20XFEL&f=2165&s=&p=&r=

YouTube: <u>https://www.youtube.com/watch?v=FJO\_DmM4q7M</u> search text: electron acceleration

## Contact

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www.desy.de