

Theoretical aspects of early SUSY searches at the LHC

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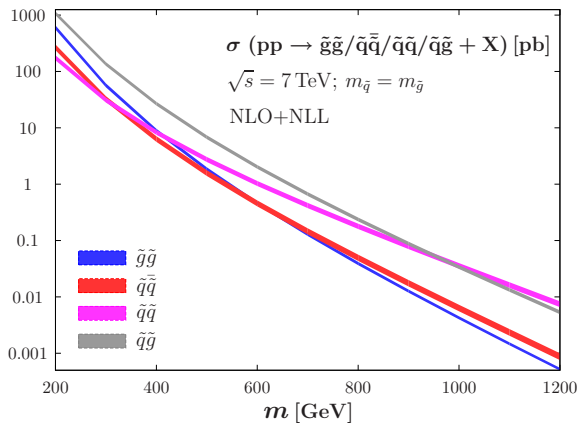
- ▶ SUSY particle production at the LHC:
MSSM cross section calculations
- ▶ Exploring SUSY at the LHC I:
constraining SUSY models from current and future exclusions
- ▶ Exploring SUSY at the LHC II:
parameter determination from cross sections and kinematic edges

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Squark and gluino production at the LHC

NLO+NLL QCD cross section

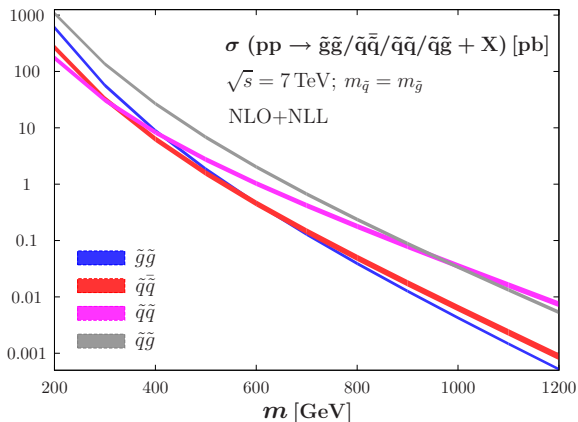
[Beenakker, Höpker, Spira, Zerwas '96; Beenakker, Brensing, MK, Laenen, Kulesza, Niessen '09]



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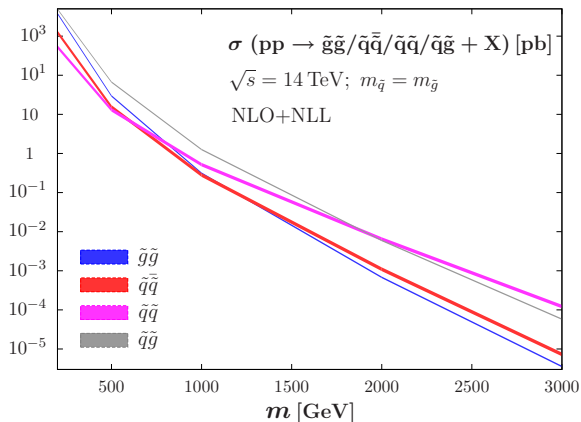


$\rightarrow \sigma \approx 100$ fb for $m \approx 1000$ GeV at the LHC with 7 TeV

Squark and gluino production at the LHC

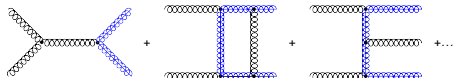
NLO+NLL QCD cross section

[Beenakker, Höpker, Spira, Zerwas '96; Beenakker, Brensing, MK, Laenen, Kulesza, Niessen '09]

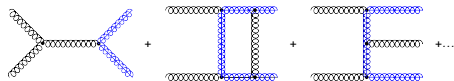


$\rightarrow \sigma \approx 2.5$ pb for $m \approx 1000$ GeV at the LHC with 14 TeV

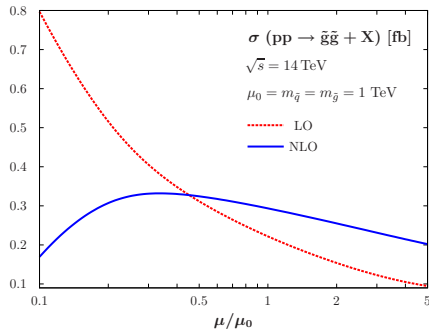
Take gluino-pair production as an example



Take gluino-pair production as an example

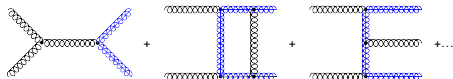


Scale dependence

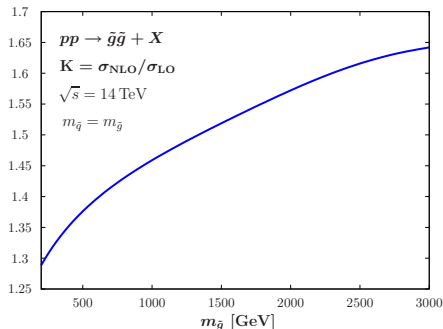


→ $\Delta\sigma \approx \pm 20\%$ at NLO [Beenakker, Höpker, Spira, Zerwas '96]

Take gluino-pair production as an example



K-factor



→ large NLO corrections for $m_{\tilde{g}} \gtrsim 1 \text{ TeV}$

Threshold summation

NLO cross section near threshold $\beta = \sqrt{1 - 4m^2/s} \ll 1$:

$$\begin{aligned} \sigma^{\text{NLO}}[gg \rightarrow \tilde{g}\tilde{g} + X] \approx & \frac{\alpha_s^2(\mu^2)}{m^2} \frac{27\pi}{64} \beta \left(1 + 4\pi\alpha_s(\mu^2) \left\{ \frac{1}{16\beta} \right. \right. \\ & \left. \left. + \frac{3}{2\pi^2} \ln^2(8\beta^2) - \frac{29}{4\pi^2} \ln(8\beta^2) - \frac{3}{2\pi^2} \ln(8\beta^2) \ln\left(\frac{\mu^2}{m^2}\right) \right\} \right) \end{aligned}$$

Threshold summation

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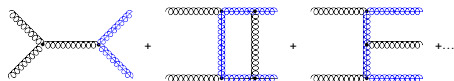
$$\sigma^{\text{NLO}}[gg \rightarrow \tilde{g}\tilde{g} + X] \approx \frac{\alpha_s^2(\mu^2)}{m^2} \frac{27\pi}{64} \beta \left(1 + 4\pi\alpha_s(\mu^2) \left\{ \frac{1}{16\beta} + \frac{3}{2\pi^2} \ln^2(8\beta^2) - \frac{29}{4\pi^2} \ln(8\beta^2) - \frac{3}{2\pi^2} \ln(8\beta^2) \ln\left(\frac{\mu^2}{m^2}\right) \right\} \right)$$

→ requires all-order summation of large logarithmic corrections:

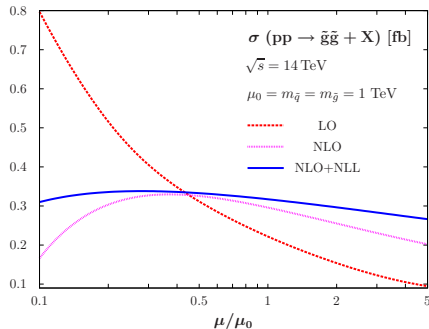
$$d\sigma^{\text{res}} \propto d\sigma_{\text{LO}} \times \exp \left(\underbrace{\log \beta^2 g_1(\alpha_s \log \beta^2)}_{\text{LL}} + \underbrace{g_2(\alpha_s \log \beta^2)}_{\text{NLL}} + \dots \right) + \dots$$

[Kulesza, Motyka; Beenakker, Brensing, MK, Kulesza, Laenen, Niessen '09
cf. Kidonakis, Sterman; Bonciani, Catani, Mangano, Nason; ...]

Take gluino-pair production as an example

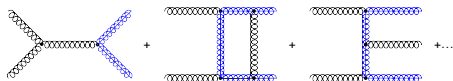


Scale dependence

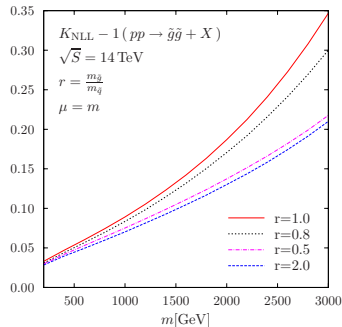


$\rightarrow \Delta\sigma \lesssim \pm 10\%$ at NLO+NLL [Kulesza, Motyka; Beenakker et al.]

Take gluino-pair production as an example



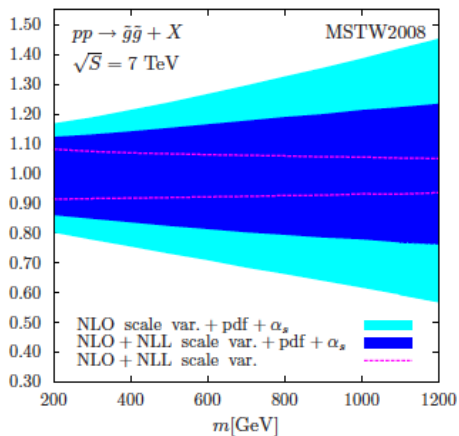
NLL K-factor



→ up to 30% increase from NLL summation [Kulesza, Motyka; Beenakker et al.]

Theory uncertainty

Theory uncertainty: $\Delta\sigma = \Delta\mu \pm \sqrt{\Delta(\text{pdf})^2 + \Delta(\alpha_s)^2}$



$\rightarrow \Delta\sigma \lesssim \pm 20\%$ at NLO+NLL [Beenakker et al.]

- ▶ **NLO QCD** → public computer code PROSPINO

[Beenakker, Höpker, MK, Plehn, Spira, Zerwas]

- ▶ **Threshold summation**

[Kulesza, Motyka; Langenfeld, Moch; Beenakker, Brensing, MK, Kulesza, Laenen, Niessen; Beneke, Falgari, Schwinn]

- ▶ **Electroweak corrections**

[Hollik, Kollar, Mirabella, Trenkel; Bornhauser, Drees, Dreiner, Kim; Bozzi, Fuks, Klasen; Beccaria, Macorini, Panizzi, Renard, Verzegnassi]

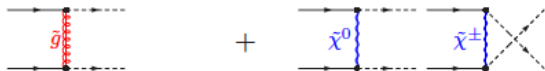
- ▶ **Bound state effects**

[Hagiwara, Yokoya; Kauth, Kühn, Marquard, Steinhauser]

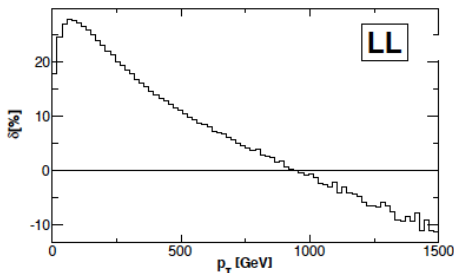
and more for SUSY models beyond the MSSM. . .

Squark pair production with EWK corrections

EWK effects include $\mathcal{O}(\alpha)$ loop corrections and QCD/EWK interference



$p_t(\tilde{q})$ distribution:



[Germer et al.]

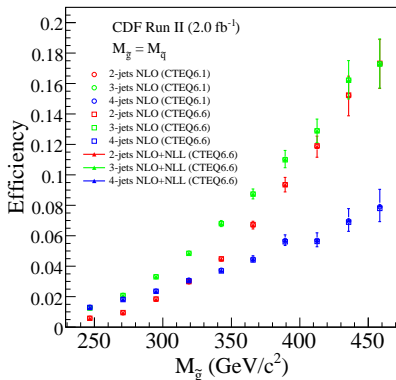
→ potentially significant effects, depending in detail on the process and the SUSY scenario

Squark and gluino searches at the Tevatron

Impact of precision calculations on CDF SUSY limits

[Beenakker et al. with Martinez and D'Onofrio]

- ▶ relative weight of $\tilde{g}\tilde{g}/\tilde{q}\tilde{q}/\tilde{q}\tilde{g}$ enters efficiency and thus cross section limit

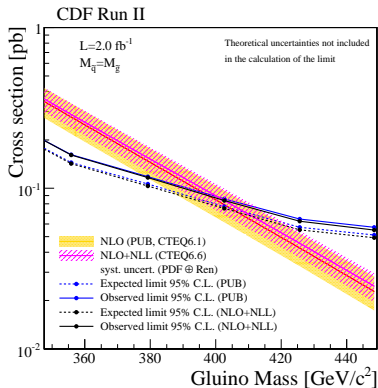


Squark and gluino searches at the Tevatron

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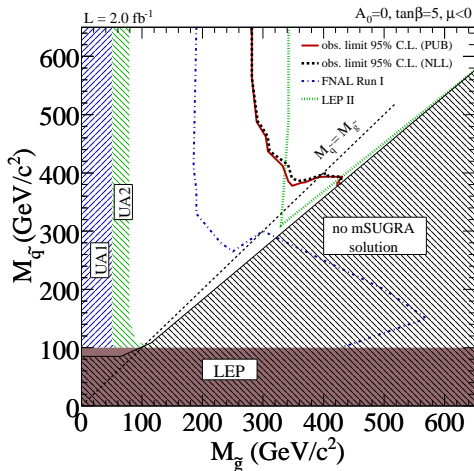
- ▶ relative weight of $\tilde{g}\tilde{g}/\tilde{q}\tilde{q}/\tilde{q}\tilde{g}$ enters efficiency and thus cross section limit
- ▶ mass limit from comparison of cross section limit and theory prediction



Squark and gluino searches at the Tevatron

Impact of precision calculations on CDF SUSY limits

[Beenakker et al. with Mario Martinez and Monica D'Onofrio]



Squark and gluino searches at the LHC

SUSY searches are often interpreted in the **constrained MSSM**,
where the breaking is universal at the GUT scale

- ▶ universal scalar masses: $M_{\tilde{Q}}^2, M_{\tilde{U}}^2, M_{\tilde{D}}^2, M_{\tilde{L}}^2, M_{\tilde{E}}^2 \rightarrow M_0^2$ at M_{GUT}
- ▶ universal gaugino masses: $M_1, M_2, M_3 \rightarrow M_{1/2}$ at M_{GUT}
- ▶ universal trilinear couplings $A_{ij}^e, A_{ij}^d, A_{ij}^u \rightarrow A \cdot h_{ij}^e, A \cdot h_{ij}^d, A \cdot h_{ij}^u$ at M_{GUT}

The masses at the weak scale are determined through the RGE.

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The masses at the weak scale are determined through the RGE.

In the CMSSM the sparticle masses are strongly correlated, e.g.

$$M_{\tilde{g}} \simeq 3M_{\tilde{\chi}^\pm} \simeq 3M_{\tilde{\chi}_2^0} \simeq 6M_{\tilde{\chi}_1^0}$$

and

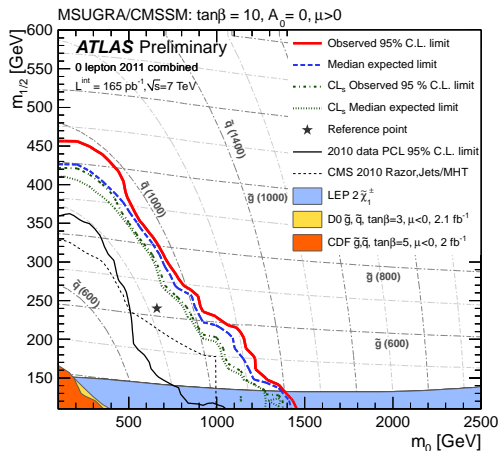
$$m_{\tilde{u}_L}^2 \simeq M_0^2 + 6.3 M_{1/2}^2 + D_{\tilde{u}_L}$$

$$m_{\tilde{e}_L}^2 \simeq M_0^2 + 0.5 M_{1/2}^2 + D_{\tilde{e}_L}$$

where $D_{\tilde{f}_L} = M_Z^2 \cos(2\beta)(T_{3f} - Q_f \sin(2\theta_W))$

Squark and gluino searches at the LHC

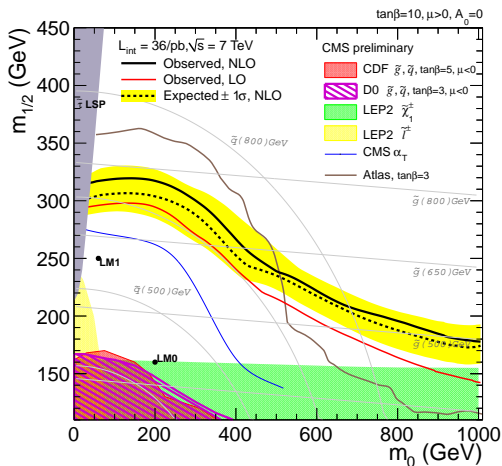
Atlas limits (165 pb^{-1})



$$\rightarrow m_{\tilde{q}} \approx m_{\tilde{g}} \gtrsim 950 \text{ GeV}$$

Squark and gluino searches at the LHC

CMS limits (36 pb⁻¹)



- ▶ SUSY particle production at the LHC:
MSSM cross section calculations
- ▶ Exploring SUSY at the LHC I:
constraining SUSY models from current and future exclusions
[P. Bechtle, K. Desch, H.K. Dreiner, M.K., B. O'Leary, C. Robens, B. Sarrazin, and P. Wienemann]
- ▶ Exploring SUSY at the LHC II:
parameter determination from cross sections and kinematic edges

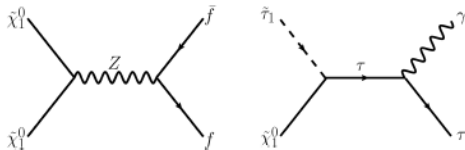
Indirect SUSY searches

There is a wealth of precision measurements from B/K physics, $(g - 2)$, astrophysics (DM) and collider limits which show sensitivity to supersymmetry, in particular

▶ $(g - 2)_\mu$

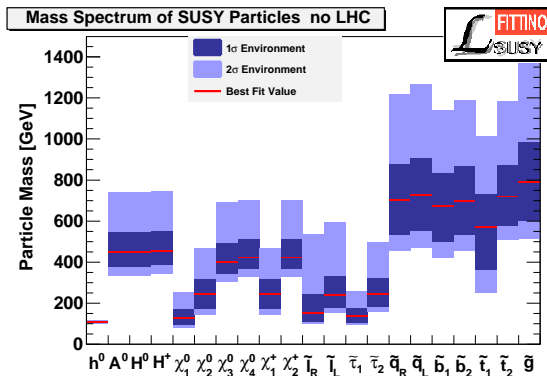


▶ DM relic abundance

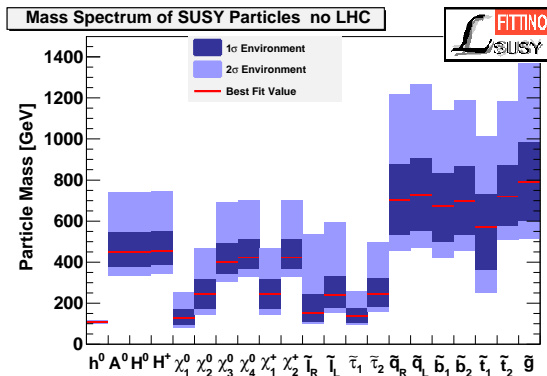


Indirect SUSY searches

→ global fit of constrained SUSY models



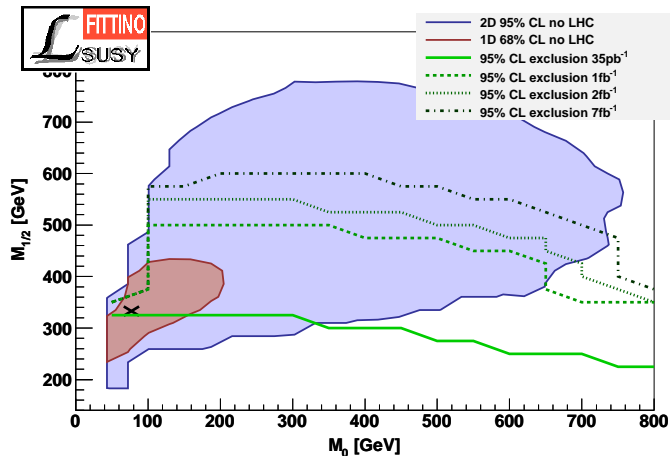
→ global fit of constrained SUSY models



- ▶ global fits point to light particle spectrum with $\tilde{m} < 1$ TeV
- ▶ current data cannot constrain more general SUSY models

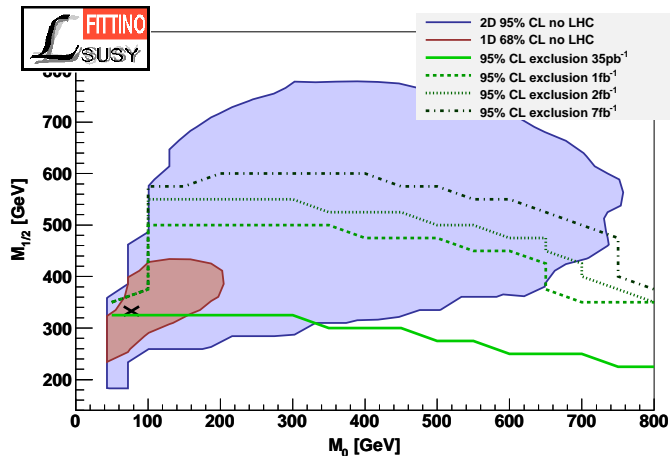
Direct SUSY searches at the LHC: expected limits

The LHC is probing the preferred region of SUSY parameter space



Direct SUSY searches at the LHC: expected limits

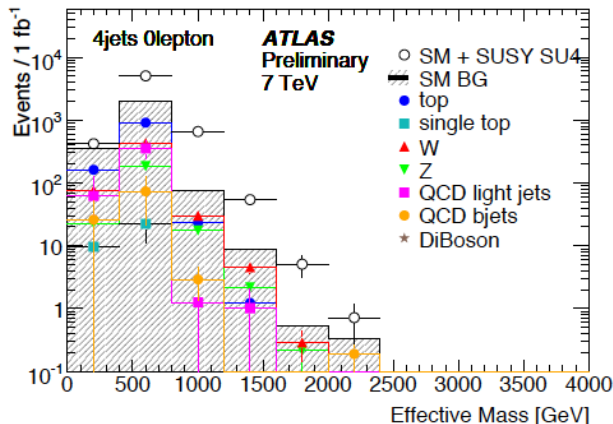
But what if we do not see any SUSY signal at the LHC?



Direct SUSY searches at the LHC: expected limits

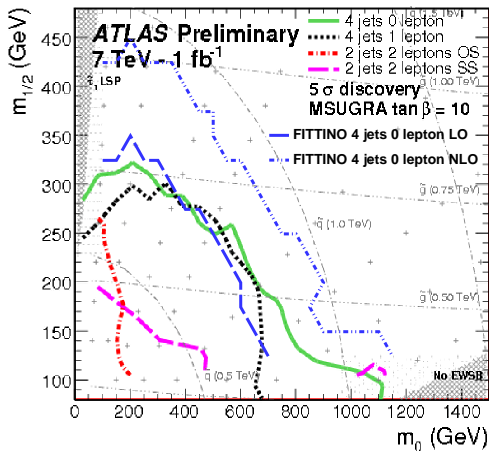
We have considered the SUSY search in the 4 jets + $E_{T,\text{miss}}$ signature

$$\text{with } M_{\text{eff}} = \sum_i p_{T,i} + E_{T,\text{miss}}$$



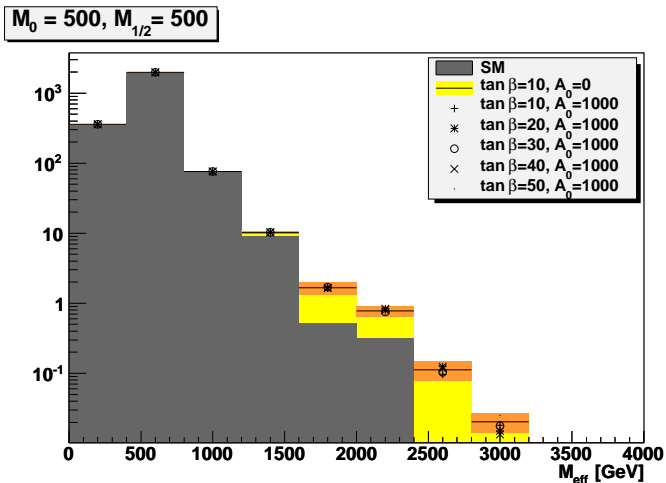
Direct SUSY searches at the LHC: expected limits

- ▶ The simulation of M_{eff} is based on Herwig++, Delphes and NLO+NLL K-factors.



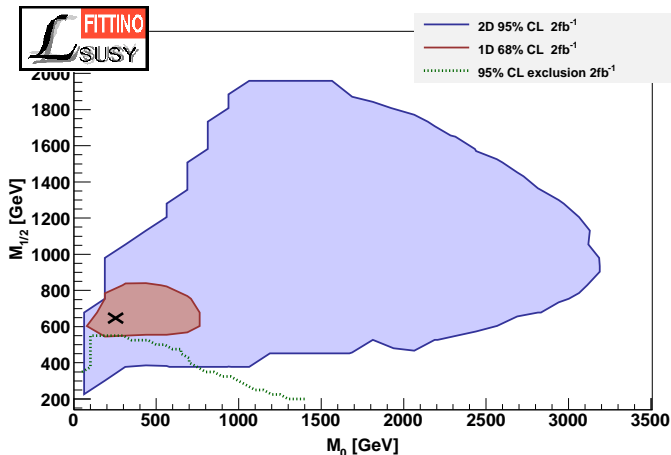
Direct SUSY searches at the LHC: expected limits

- ▶ The 4 jets + $E_{T,miss}$ signature is rather independent of $\tan\beta$ and A_0



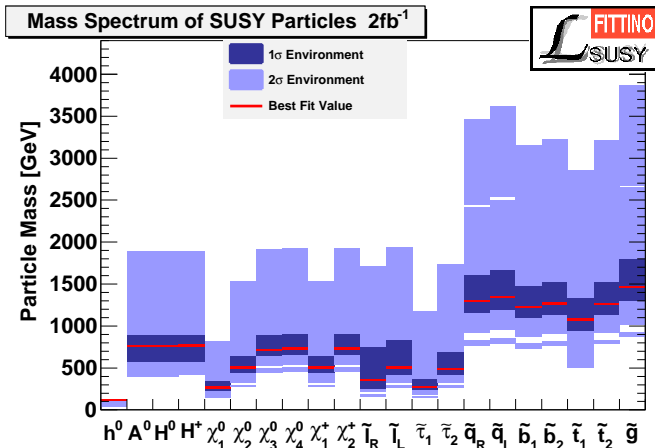
Global SUSY fits with projected LHC exclusions

- ▶ Low-energy observables, DM and LHC exclusions with 2 fb^{-1}



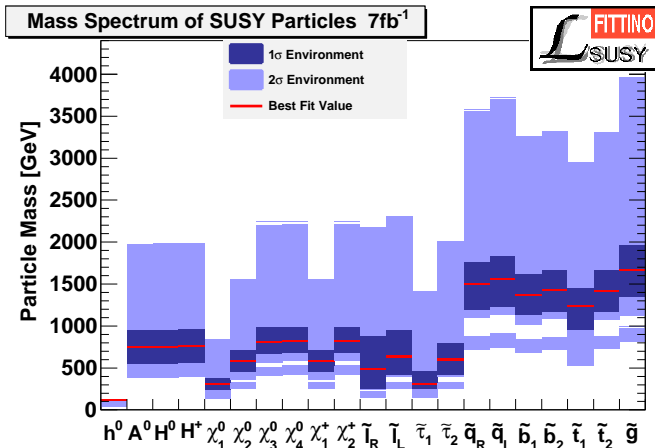
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Global SUSY fits with projected LHC exclusions

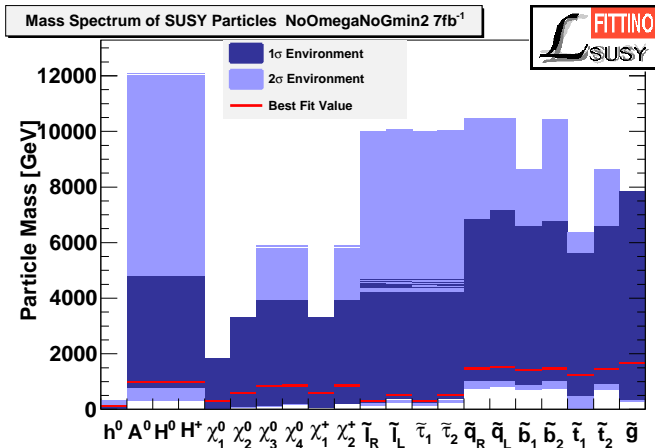
- ▶ Low-energy observables, DM and LHC exclusions with 7 fb^{-1}



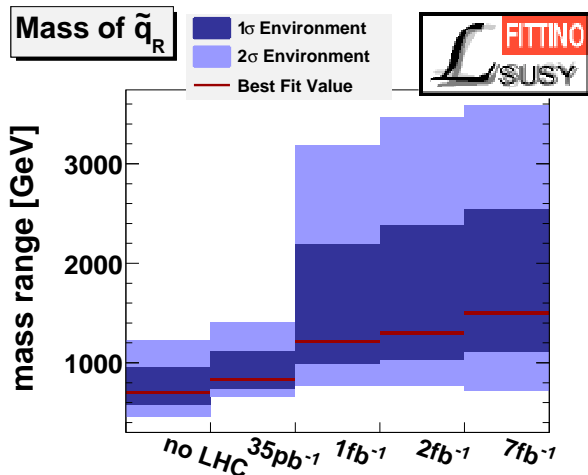
- ▶ what happens if we take out $(g - 2)_\mu$ and Ω_{DM} ?

Global SUSY fits with projected LHC exclusions

- ▶ what happens if we take out $(g - 2)_\mu$ and Ω_{DM} ?



- ▶ LHC mass limits on squarks are rather model independent



Global SUSY fits with projected LHC exclusions: is there a tension?

- LEOs prefer low mass scales (for non-coloured sector)
- LHC prefers high mass scales (for coloured sector)

Is there a tension building up?

Global SUSY fits with projected LHC exclusions: is there a tension?

- LEOs prefer low mass scales (for non-coloured sector)
- LHC prefers high mass scales (for coloured sector)

Is there a tension building up?

Let us look at the best fit points:

	M_0	$M_{1/2}$	A_0	$\tan \beta$	χ^2/ndf
no LHC	77^{+114}_{-31}	333^{+89}_{-87}	426^{+70}_{-735}	13^{+10}_{-8}	19/20
35 pb ⁻¹	126^{+189}_{-54}	400^{+109}_{-40}	724^{+722}_{-780}	17^{+14}_{-9}	20/21
1 fb ⁻¹	235^{+389}_{-103}	601^{+148}_{-63}	627^{+1249}_{-717}	31^{+19}_{-18}	24/21
2 fb ⁻¹	254^{+456}_{-128}	647^{+157}_{-74}	771^{+1254}_{-879}	30^{+20}_{-19}	24/21
7 fb ⁻¹	403^{+436}_{-281}	744^{+142}_{-150}	781^{+1474}_{-918}	43^{+11}_{-33}	25/21

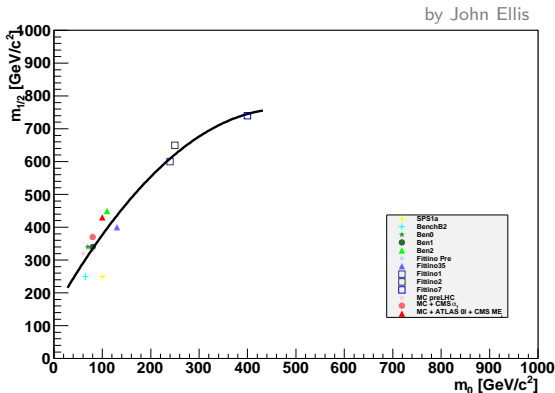
→ even the CMSSM would "survive" the 2011/2012 LHC run

[Note: $a_\mu^{\text{SUSY}} \sim \text{sgn}(\mu) \tan \beta M_{\text{SUSY}}^{-2}$ and Ω_{DM} require larger $\tan \beta$]

Comparison of global CMSSM fits with and without LHC exclusions

There has been a lot of activity recently, see e.g.

Allanach, arXiv:1102.3149 [hep-ph], Buchmueller et al., arXiv:1102.4585 [hep-ph], Bechtle et al., arXiv:1102.4693 [hep-ph], Allanach et al., arXiv:1103.0969 [hep-ph]

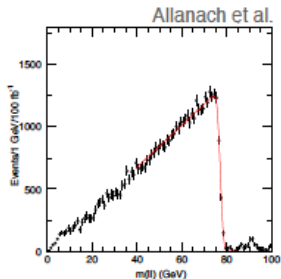
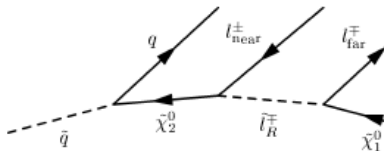


→ the analyses differ in detail, but there is good agreement overall

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SUSY parameter determination at the LHC

Mass measurements from cascade decays, e.g.



→ kinematic endpoints sensitive to masses:

$$\begin{aligned}(m_{ll}^2)^{\max} &= (m_{\tilde{\chi}_2^0}^2 - m_{l_R}^2)(m_{l_R}^2 - m_{\tilde{\chi}_1^0}^2)/m_{l_R}^2 \\(m_{qll}^2)^{\max} &= (m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{\tilde{\chi}_1^0}^2)/m_{\tilde{\chi}_2^0}^2 \\(m_{ql,\min}^2)^{\max} &= (m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{\tilde{\chi}_2^0}^2 - m_{l_R}^2)/m_{\tilde{\chi}_2^0}^2 \\(m_{ql,\max}^2)^{\max} &= (m_{\tilde{q}_L}^2 - m_{\tilde{\chi}_2^0}^2)(m_{l_R}^2 - m_{\tilde{\chi}_1^0}^2)/m_{l_R}^2\end{aligned}$$

SUSY parameter determination with cross sections

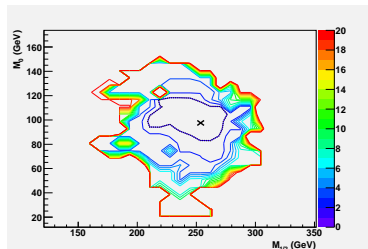
How well could we do at 7 TeV and 1 fb^{-1} ? [Dreiner, MK, Lindert, O'Leary]

[cf. Baer et al., Altunkaynak et al., ...]

SUSY parameter determination with cross sections

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SPS1a with 4 kinematic edges only

→ no stable fit

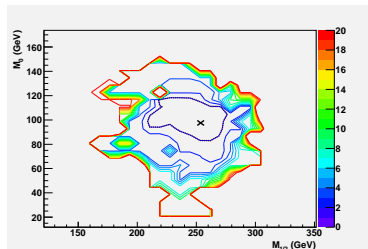
→ add **cross sections**:

sensitive to masses and spins

SUSY parameter determination with cross sections

How well could we do at 7 TeV and 1 fb^{-1} ? [Dreiner, MK, Lindert, O'Leary]

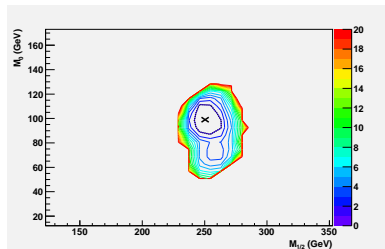
[cf. Baer et al., Altunkaynak et al., ...]



SPS1a with 4 kinematic edges only

→ no stable fit

→ add **cross sections**:
sensitive to masses and spins



kinematic edges \oplus **cross sections**

→ $m_0 = 99 \pm 9 \text{ GeV}$

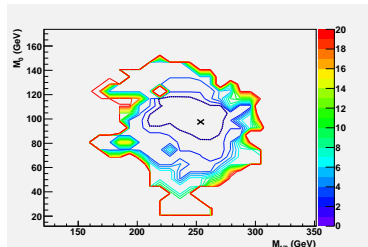
$m_{1/2} = 250 \pm 7 \text{ GeV}$

$\tan \beta = 11 \pm 6$

SUSY parameter determination with cross sections

How well could we do at 7 TeV and 1 fb⁻¹? [Dreiner, MK, Lindert, O'Leary]

[cf. Baer et al., Altunkaynak et al., ...]



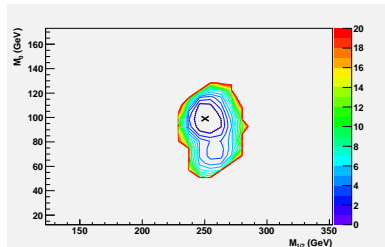
SPS1a with 4 kinematic edges only

→ no stable fit

→ add **cross sections**:

sensitive to masses and spins

→ **cross sections are crucial to determine BSM parameters**



kinematic edges ⊕ **cross sections**

→ $m_0 = 99 \pm 9$ GeV

$m_{1/2} = 250 \pm 7$ GeV

$\tan \beta = 11 \pm 6$

Beyond mSUGRA: explore mirage mediation models

[Choi, Nilles, Falkowski, Ratz, Loewen and many others]

→ characteristic pattern of soft SUSY-breaking terms

$$M_1 : M_2 : M_3 \simeq (1 + 0.66\alpha) : (2 + 0.2\alpha) : (6 - 1.8\alpha)$$

→ relative size of modulus and anomaly mediation controlled by α

How well can we determine α from future LHC data?

SUSY parameter determination with cross sections

Beyond mSUGRA: explore mirage mediation models

[Choi, Nilles, Falkowski, Ratz, Loewen and many others]

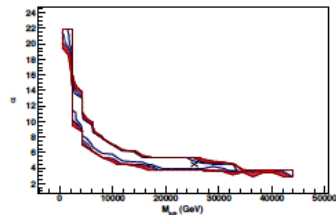
→ characteristic pattern of soft SUSY-breaking terms

$$M_1 : M_2 : M_3 \simeq (1 + 0.66\alpha) : (2 + 0.2\alpha) : (6 - 1.8\alpha)$$

→ relative size of modulus and anomaly mediation controlled by α

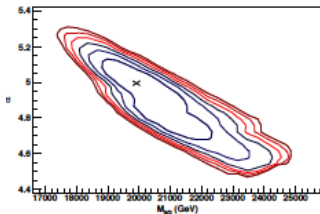
How well can we determine α from future LHC data?

Assume LHC @ 14 TeV with 1 fb^{-1}



kinematic edges only

[Conley, Dreiner, Glaser, MK, Tattersall]



kinematic edges \oplus cross sections

$$\rightarrow \alpha = 4.8 \pm 0.2$$

- ▶ Squark and gluino cross section predictions are in good shape
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 - ▶ EWK corrections are significant for specific channels and/or phase space regions
- ▶ The LHC is now probing the SUSY parameter space favoured by low-energy observables and DM
 - ▶ It is possible to reconcile LE measurements (mostly $(g-2)_\mu$ and Ω_{DM}) with a possible non-discovery of SUSY in the 7 TeV run, even in very constrained models like the CMSSM.
 - ▶ LHC searches mostly constrain the coloured sparticle sector and can push squark and gluino mass limits up to about 1.5 TeV in 2011/2012.
 - ▶ The LHC limits on squarks and gluinos are rather model independent.

Conclusion & outlook

- ▶ Cross sections are important input for LHC BSM parameter fits, in particular
 - ▶ for early LHC data
 - ▶ for non-universal models
 - ▶ and to disentangle models with different spin, eg. SUSY, UED, little Higgs, ...

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And if you want to know more...

SUSY/BSM cross section predictions are addressed within a LPCC working group [Mangano, MK et al.]

see <http://web.physik.rwth-aachen.de/service/wiki/bin/view/Main/BSMCrossSectionWorkingGroup>

BSM parameter determination now coordinated effort within the Terascale Alliance [Wienemann et al.]

see <http://www.terascale.de>