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The Ternovsky-Shul'ga-Fomin effect of suppression of radiation in aligned crystals

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Spectral Radiation Density (SRD) and its suppression effect

$$l_c \sim 2c\gamma^2 / \omega$$

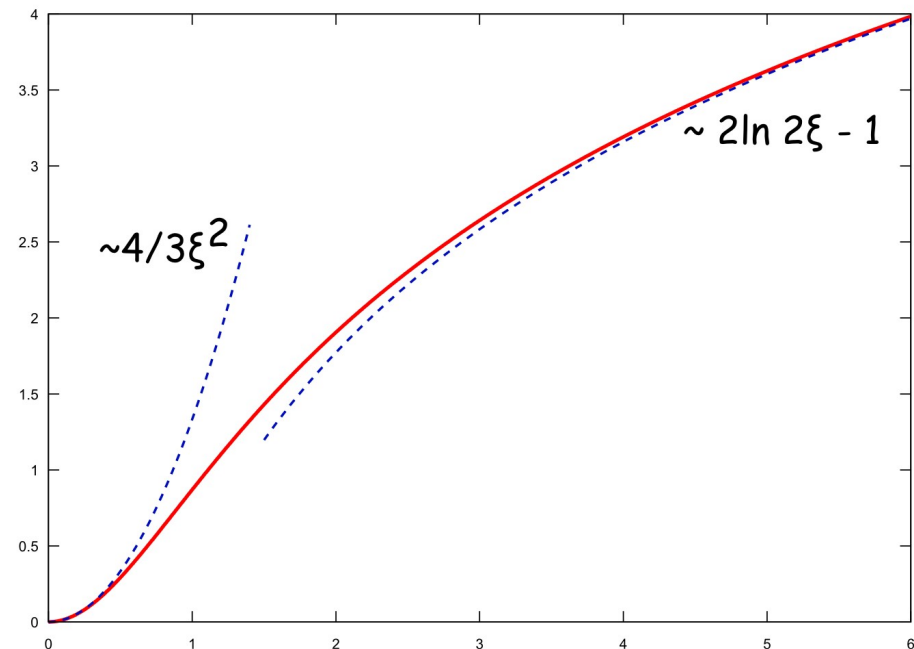
Radiation coherence length

Spectral radiation density

$$\frac{d\mathcal{E}}{d\omega} = \frac{2e^2}{\pi c} \left\{ \frac{2\xi^2 + 1}{\xi\sqrt{\xi^2 + 1}} \ln\left(\xi + \sqrt{\xi^2 + 1}\right) - 1 \right\}, \quad \xi = \gamma\vartheta/2$$

$$L \ll l_c$$

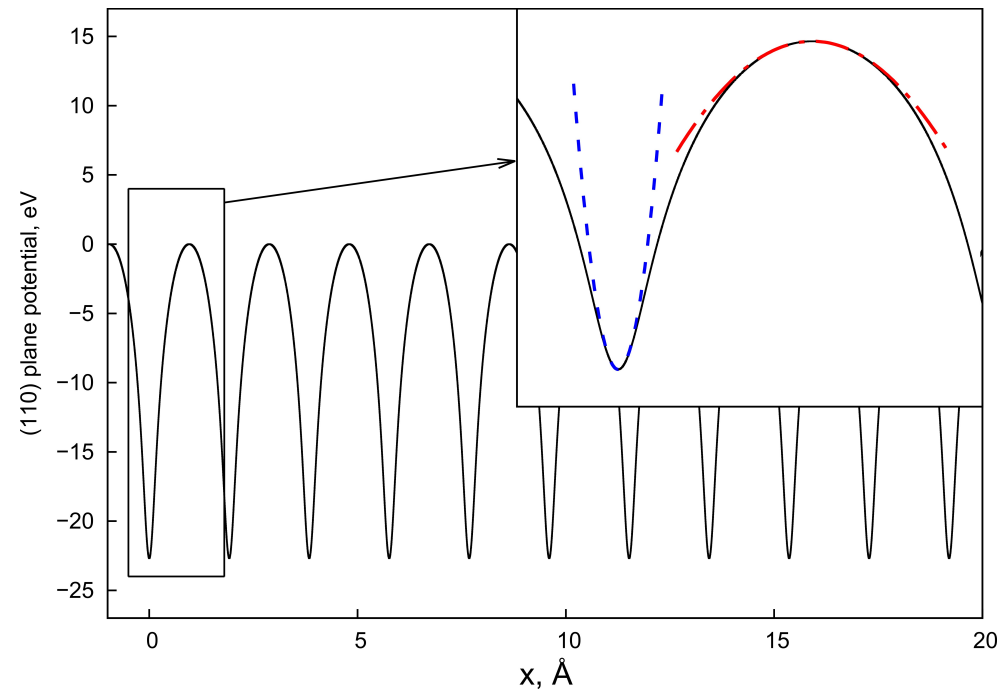
Radiation
suppression effect





Continuous planes potential

$$U_{pl}(x) = \frac{1}{a_{pl}} \int_0^{a_{pl}} dy U_{ax}(x, y)$$



$$U_{ax}(\vec{\rho}) = \frac{1}{a_1 \bar{v}^2} \int d^2 v e^{-v^2/2\bar{v}^2} \int_{-\infty}^{\infty} dz u(\vec{\rho} + \vec{v}, z) \quad - \quad \text{string potential}$$

v - heat oscillations of atom coordinates

$$u_m(r) = \sum_i \alpha_i \exp(-\beta_i r/R) \quad - \quad \text{Molière approximation for single atoms}$$

Angular trajectories of positrons in the Si (110) planes potential at different incidence angles

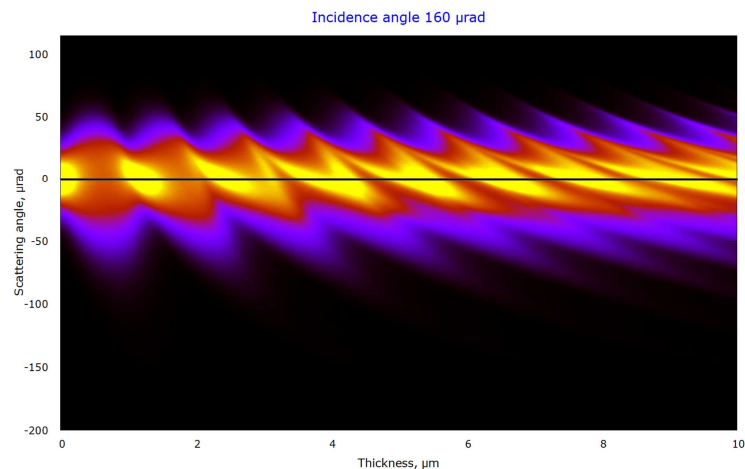
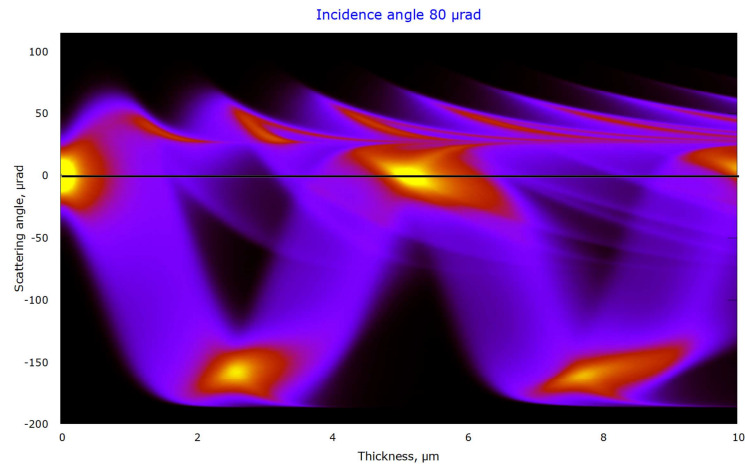
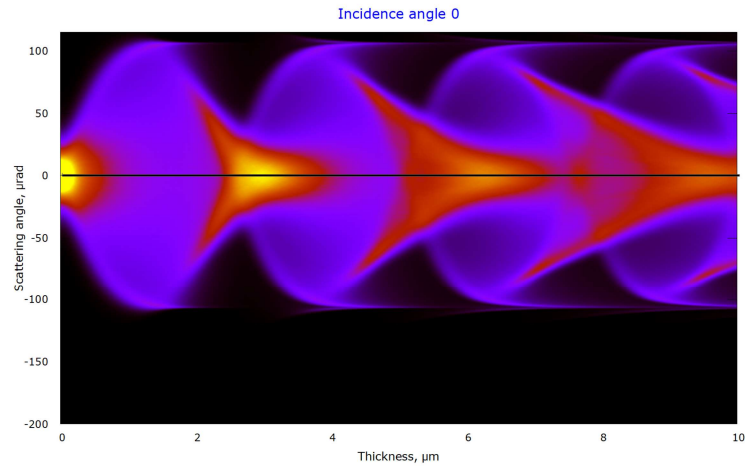
$$T_{1/2} [\text{\AA}] = b \cdot \beta \sqrt{\varepsilon [eV]}$$

$$\theta_c = \sqrt{\frac{2U_0}{pv}} \approx \sqrt{\frac{2U_0}{\varepsilon}}$$

$$T_{\text{ob}} = d/\psi$$

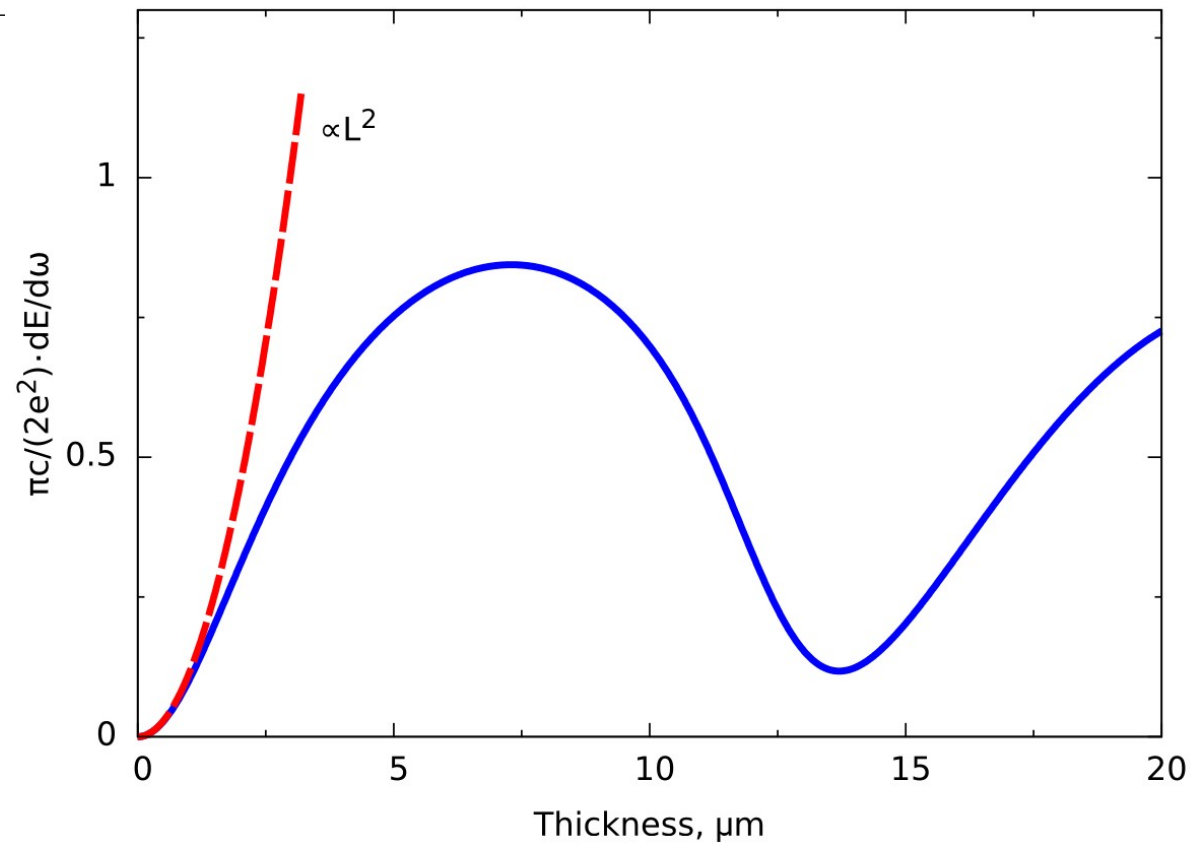
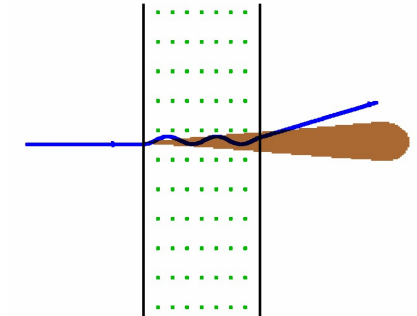
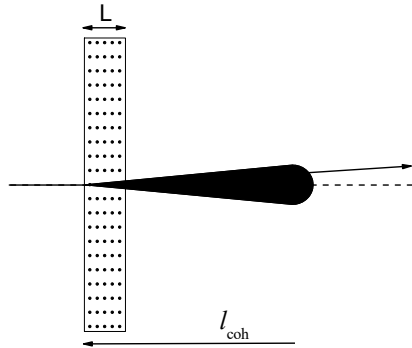
$$\Delta \mathcal{G} = \frac{d \cdot \overline{\nabla U}}{2} \cdot \frac{1}{\psi \varepsilon_{\parallel} \beta^2}$$

$$\left\langle \frac{d\mathcal{E}}{d\omega} \right\rangle = \frac{1}{N} \sum_n \left(\frac{d\mathcal{E}}{d\omega} \right)_n \quad \begin{array}{l} \text{Averaging over} \\ \text{impact parameters} \end{array}$$



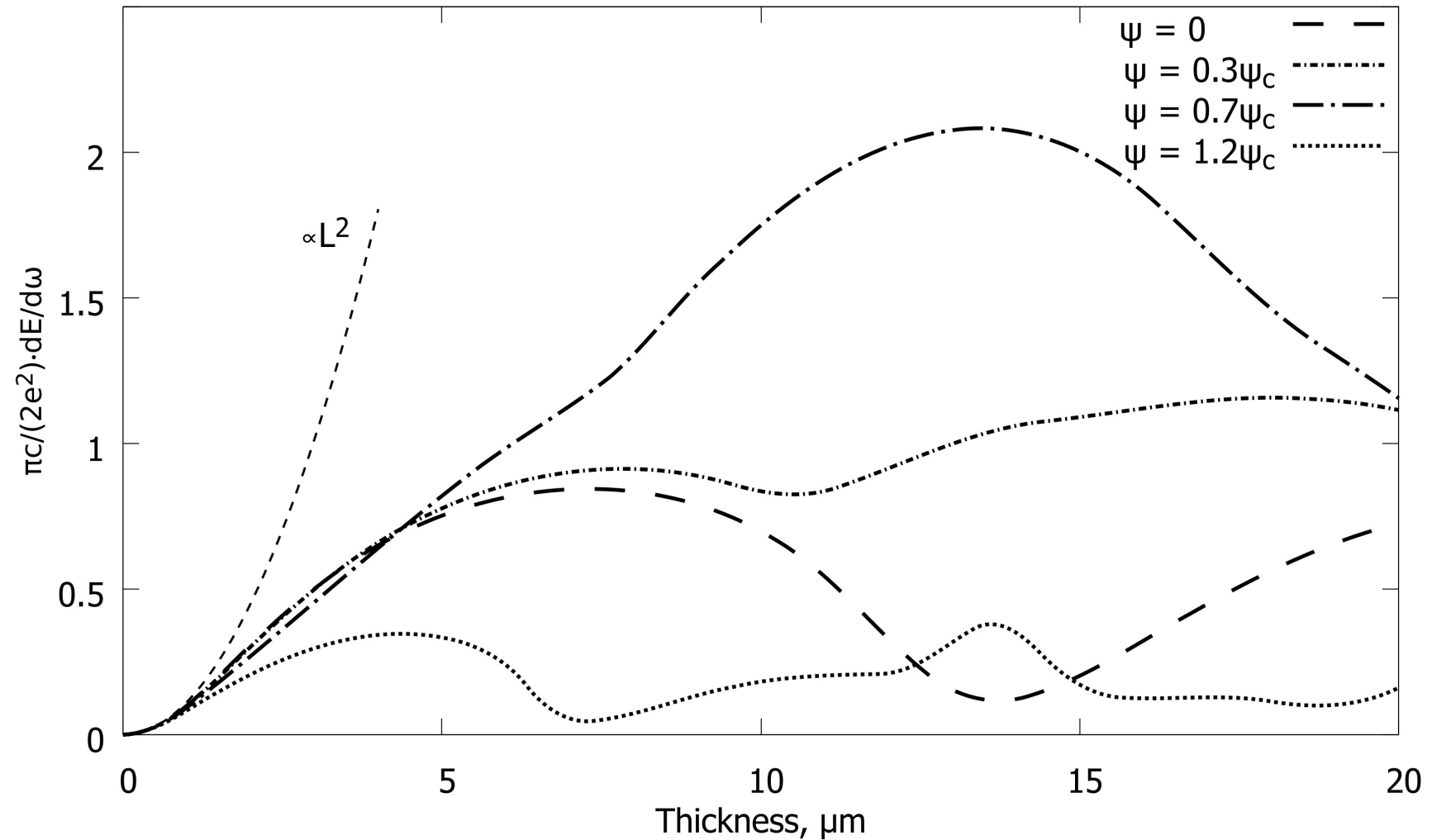


SRD of **100 Gev** positrons incident perpendicular to Si (110) plane



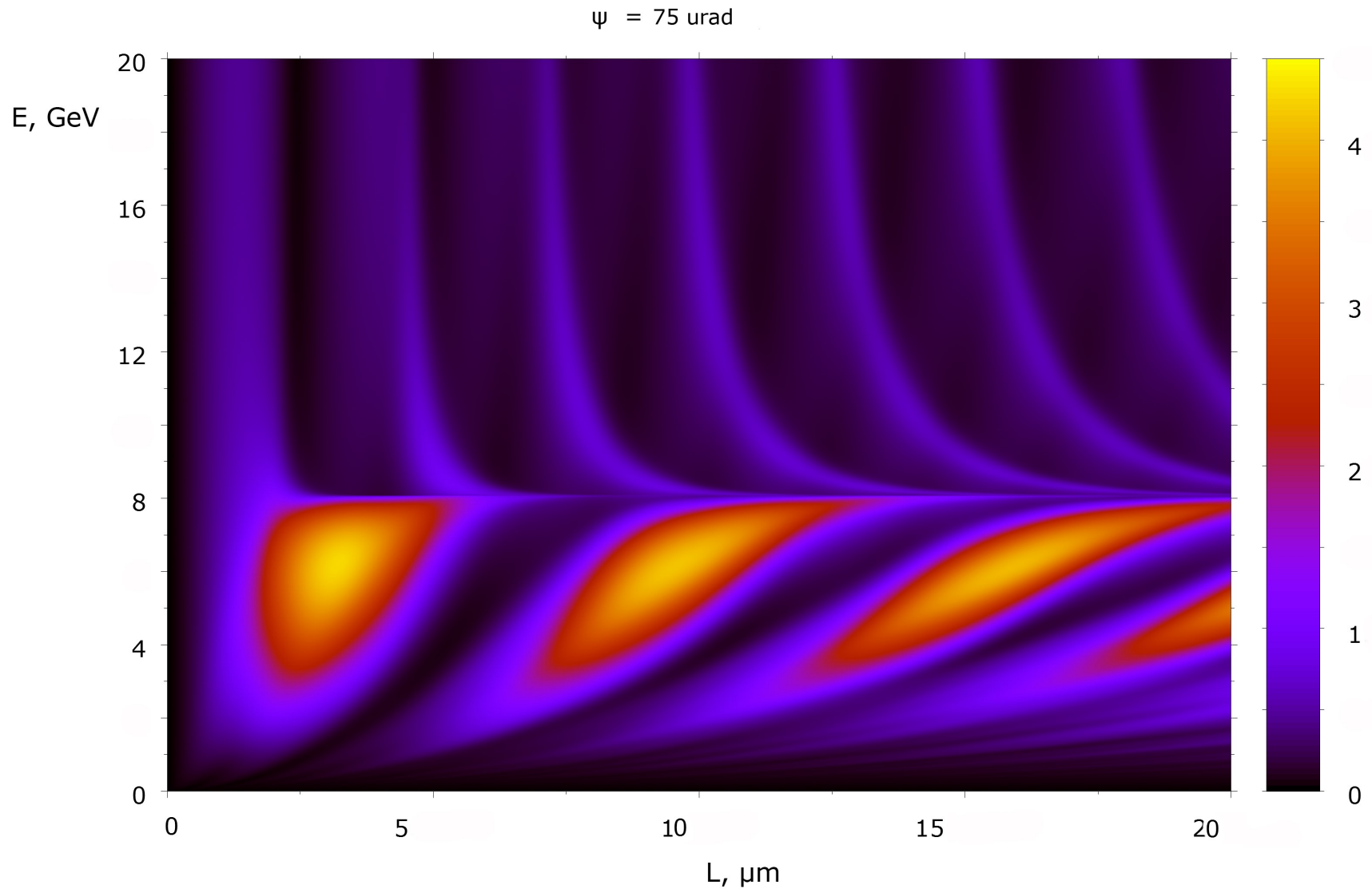


SRD of positrons at different incidence angles relatively Si (110) plane



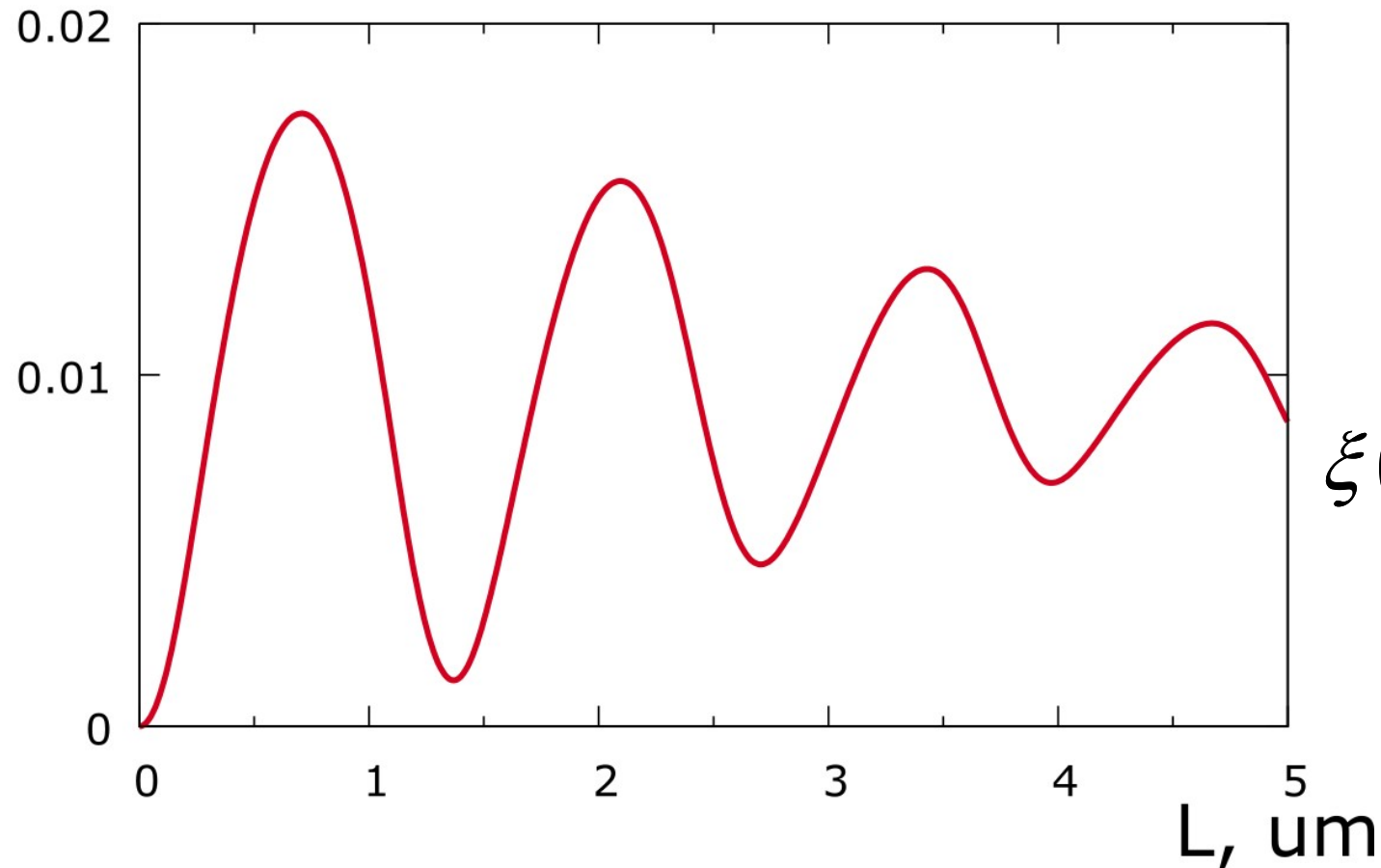


Dependence of SRD on thickness and particle energy at the incidence angle $\psi=75\mu\text{rad}$





SRD of 1GeV positrons at incidence parallel to Si (110) planes



$$T_{1/2} \sim 1.3 \mu\text{m}$$

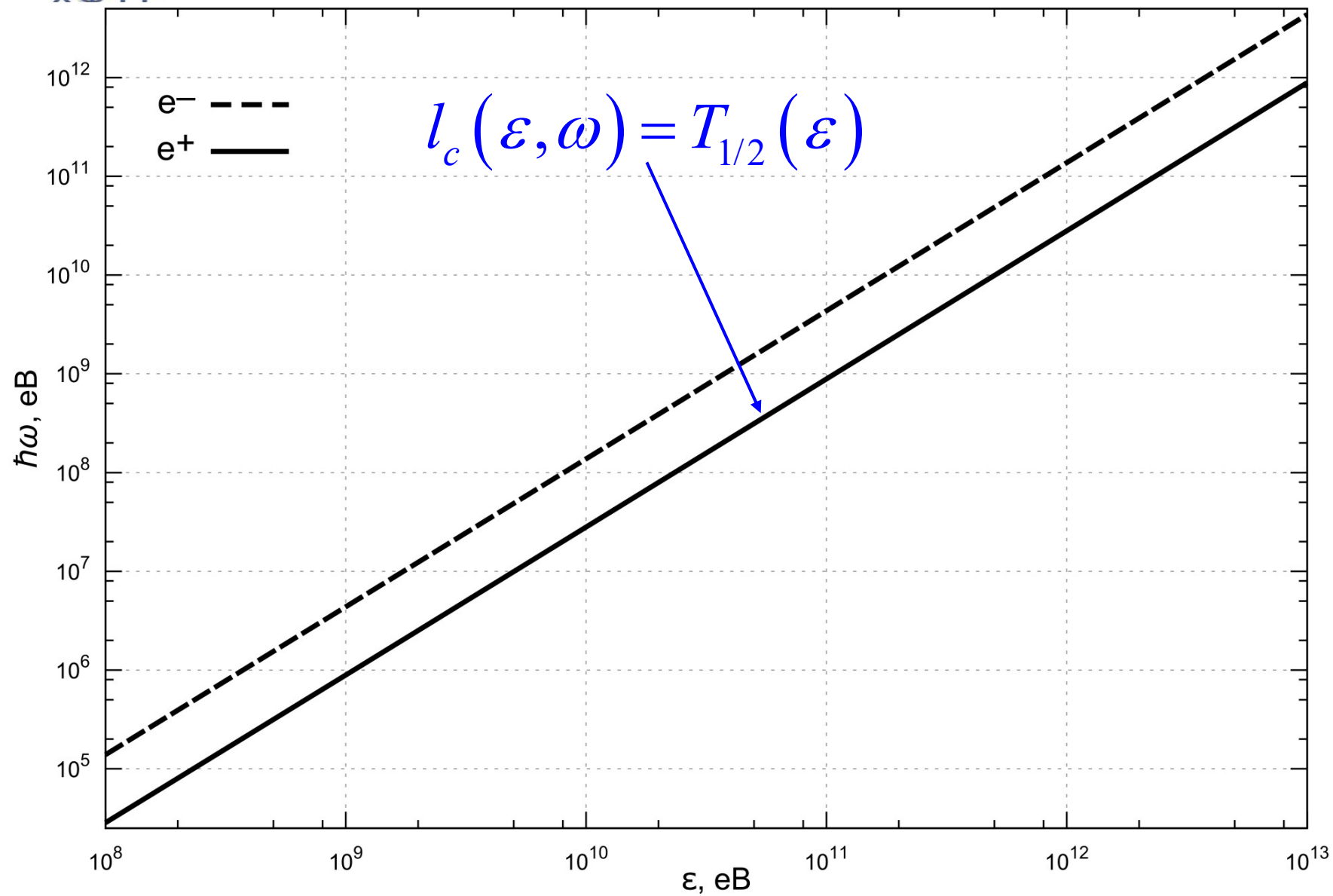
$$\xi = \gamma \vartheta / 2$$

$$\xi(\vartheta_c) \approx 0.21$$

$$T_{1/2} = l_c = 2c\gamma^2 / \omega \quad \Rightarrow \quad \hbar\omega < \sim 0.87 \text{ MeV}$$



Applicability range Si (110)





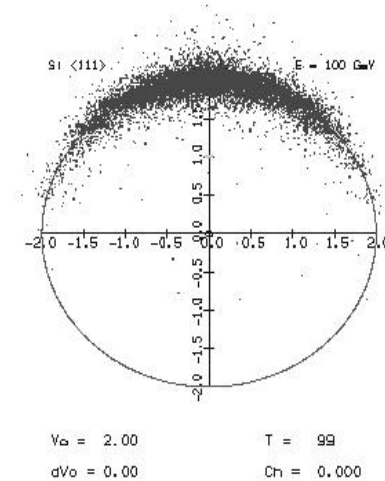
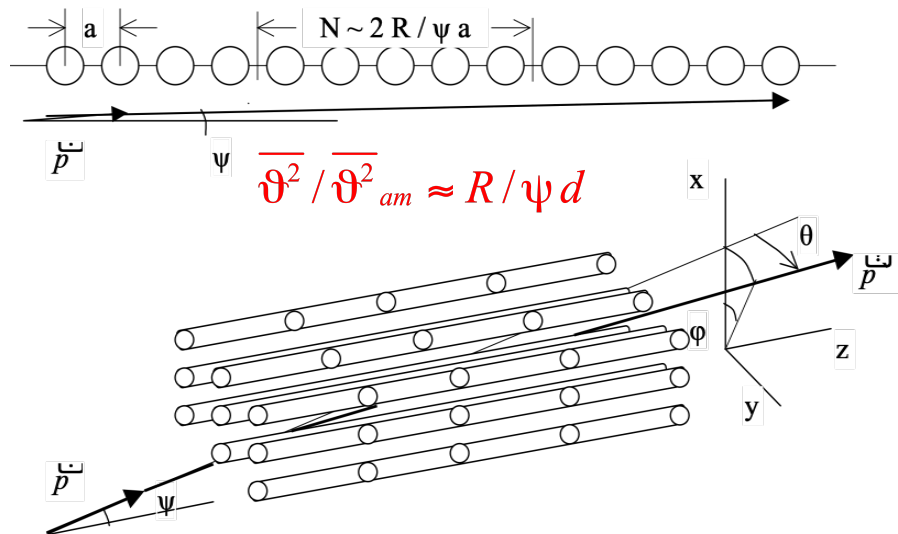
Conclusions

- For the motion of e^+ and, in some cases, e^- in the field of crystal planes, the effect of periodical spatial dependence of the coherent radiation density is possible besides with the effect of the radiation suppression effect
- At the angles ~ 0.7 of channeling angle the radiation maximum is reached
- At above-barrier motion energies neither spatial period of SRD oscillations, nor its amplitude does not depend on the particle energy if this energy is high enough
- The conditions for observing the considered effect can be reached at high energy experiments on existing experimental setups



TSF effect in a thin crystal: axial orientation

Multiple scattering in crystal: N. Shul'ga, V. Truten', S. Fomin, J. Techn. Phys. **52** (1982) 2279.

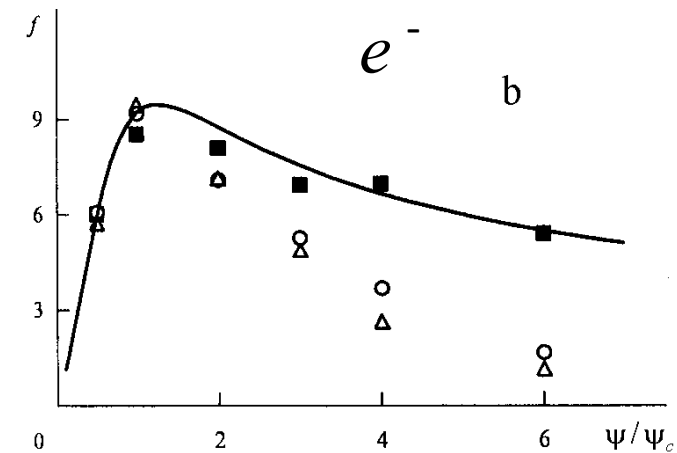
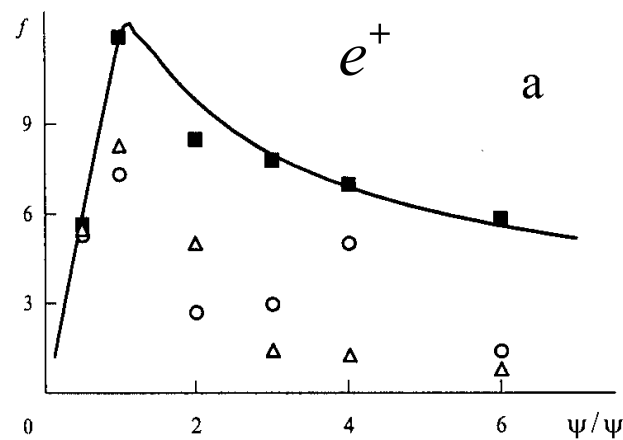


$$f = \sqrt{\langle \vartheta^2 \rangle / \langle \vartheta^2 \rangle_{am}}$$

Si <111>

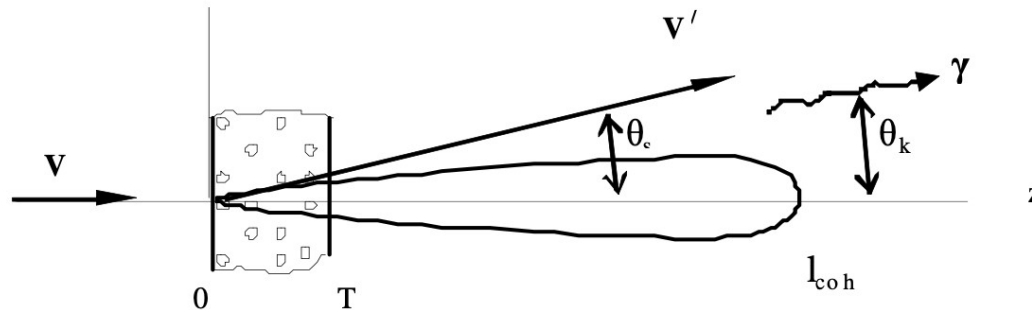
T = 50 μ m

E = 30 GeV





TSF effect in a thin crystal: axial orientation



$$l_c \gg T$$

$$l_c = \frac{2\varepsilon\varepsilon'}{m^2\omega}$$

$$\frac{dE}{d\omega} = \frac{2e^2}{\pi c} \left\{ \frac{2\xi^2 + 1}{\xi \sqrt{\xi^2 + 1}} \ln \left(\xi + \sqrt{\xi^2 + 1} \right) - 1 \right\}, \quad \xi = \gamma\vartheta/2$$

$$\left\langle \frac{dE}{d\omega} \right\rangle_c \approx \frac{R}{8\psi a} \frac{dE_{BH}}{d\omega}, \quad \gamma^2 \theta_{cr}^2 \ll 1,$$

$$\frac{dE_{BH}}{d\omega} = \frac{4}{3} \frac{T}{X_0}$$

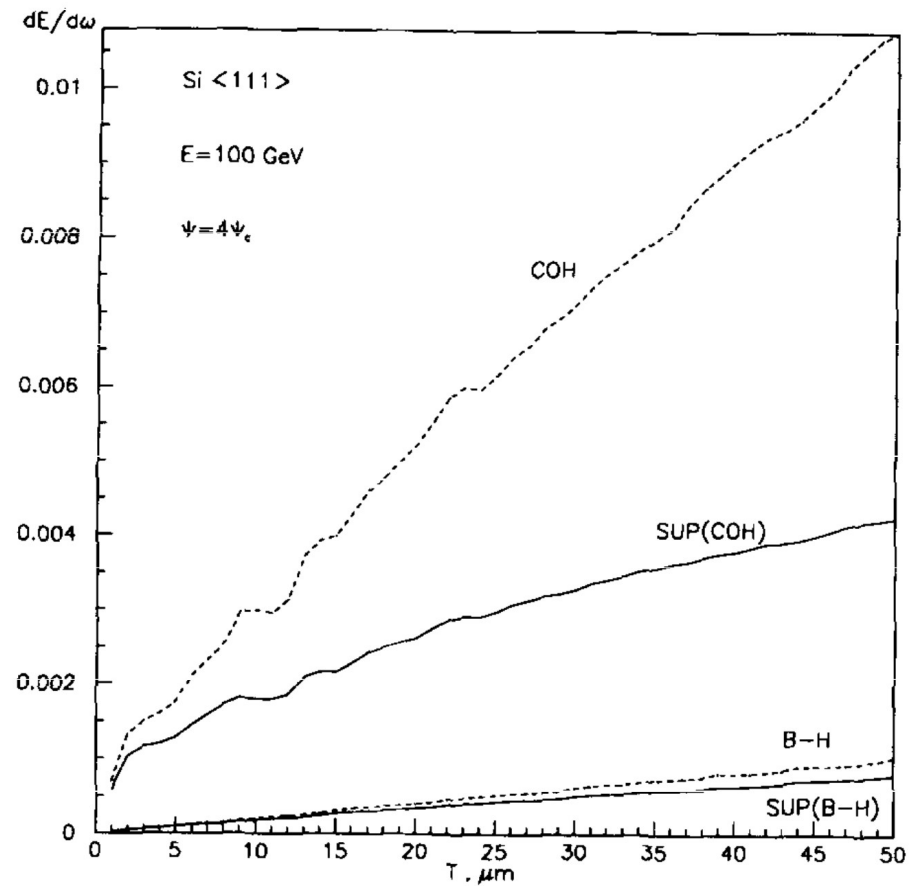
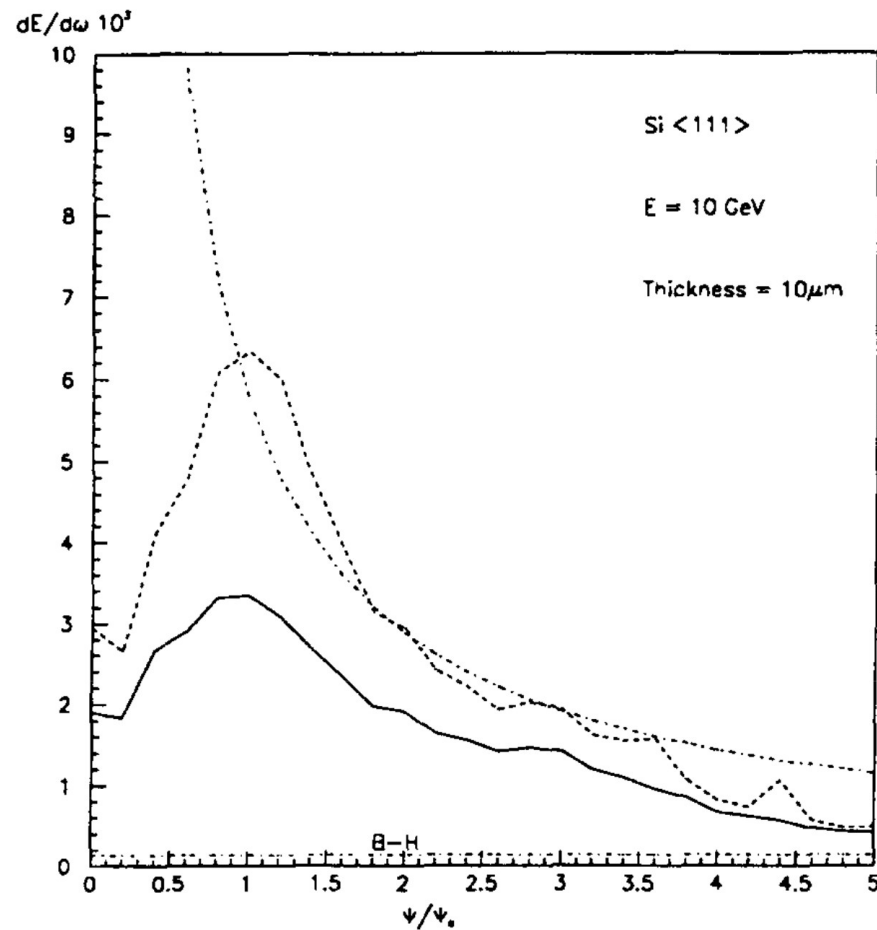
$$\left\langle \frac{dE}{d\omega} \right\rangle_c \approx \frac{2e^2}{\pi} \left[\ln \left(\frac{R}{8\psi a} \gamma^2 \theta_{am}^2 \right) - 1 \right], \quad \gamma^2 \theta_{cr}^2 \gg 1.$$

$$\omega \geq \gamma \omega_p$$



TSF effect in a thin crystal: axial orientation

S. Fomin, A. Jejcic, J. Maillard, N. Shul'ga, J. Silva, NIM B, 115 (1996) 375; 119 (1996) 59.





TSF effect in a thin crystal: **volume reflection**

A. Taratin, S. Vorobyov, Phys. Lett. A, 119 (1987) 425.

$$\theta_0 = 62 \text{ } \mu\text{rad}$$

$$\theta_B = \frac{L}{R_B}$$

Computer
simulations

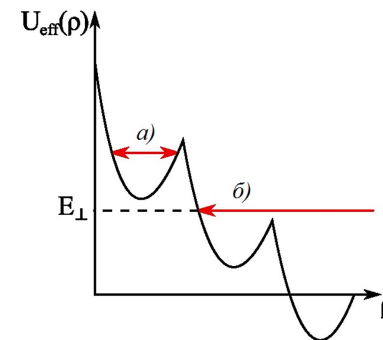
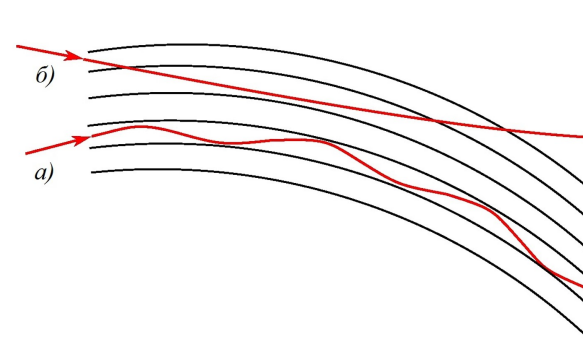
$$E = 300 \text{ GeV}$$

Si (110)

$$L = 0.62 \text{ mm}$$

$$R = 500 \text{ m}$$

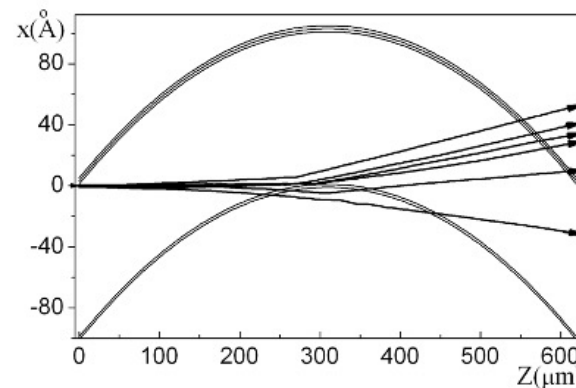
$$\theta_c \equiv \theta_L$$



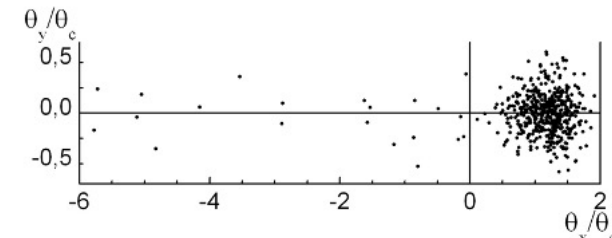
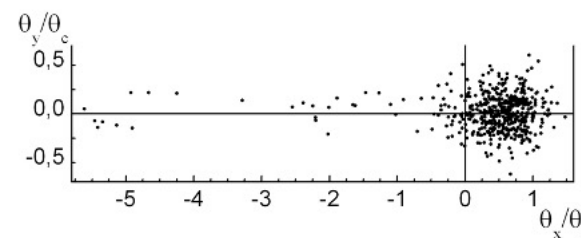
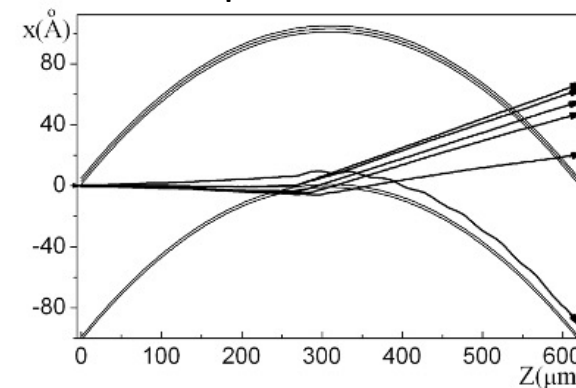
$$\theta_L = \sqrt{\frac{2U_0}{\varepsilon}},$$

$$U_0 = 2\pi n d a_{TF} Z e^2$$

electrons



positrons



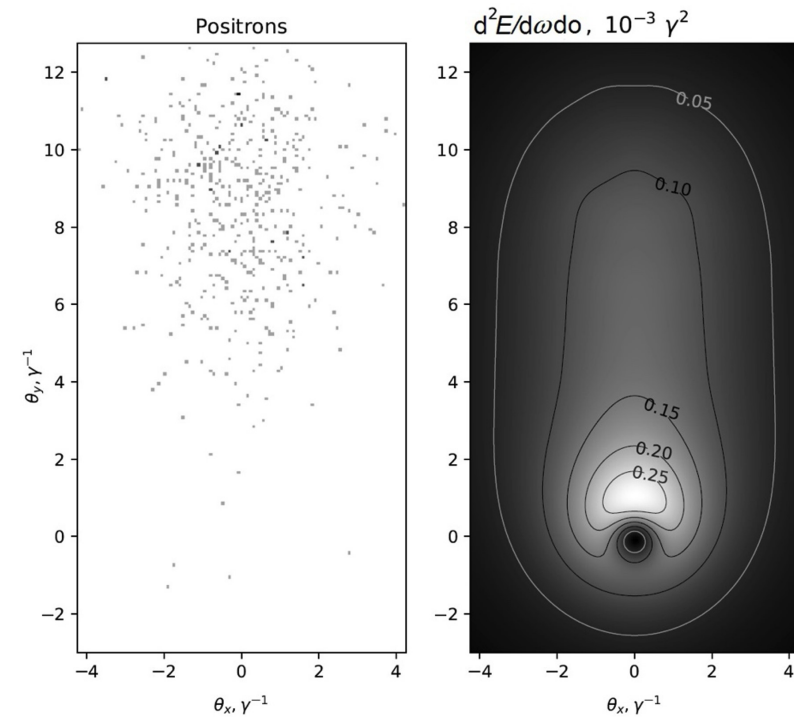
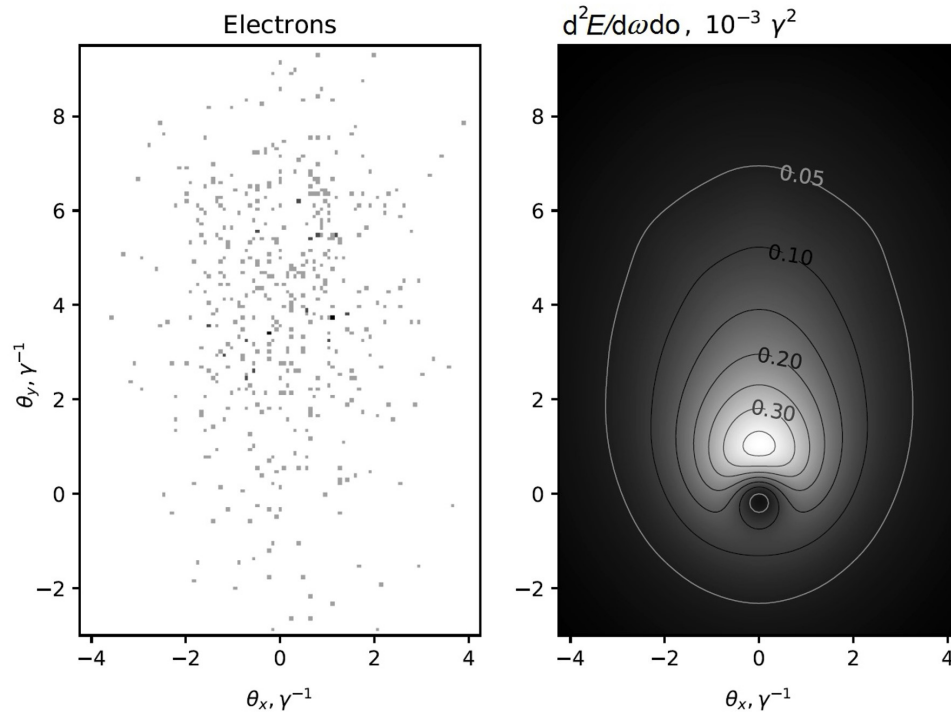


TSF effect in a thin crystal: **volume reflection**

S. Fomin, M. Malovytsia, V. Truten', N. Shul'ga, East European Physical Journal (2024) in print.

$$\gamma\theta_L \geq 1 \text{ \& } T \leq l_c$$

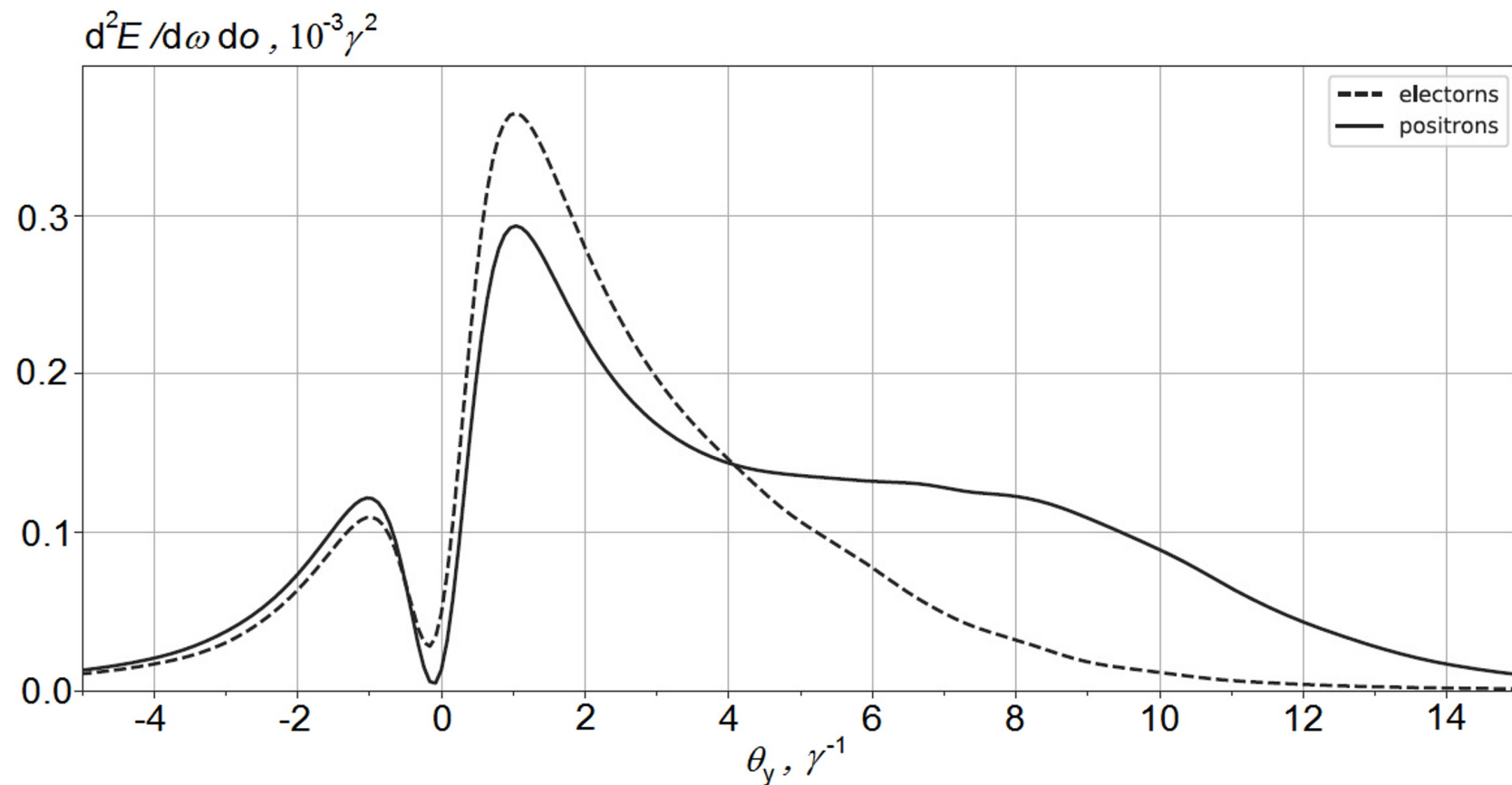
$$\theta_L \approx 16 \mu\text{rad} \quad \gamma\theta_L \approx 8$$





TSF effect in a thin crystal: **volume reflection**

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Thank you for attention !