

Generalized Geometry

Ümit Ertem

August 5, 2024

Contents

1	Basic Concepts	2
2	Spinors and Complex Structures	6
3	Supergravity	10
4	T-duality	14

Lecture 1

Basic Concepts

Let us consider an n -dimensional manifold M . We denote the tangent and cotangent bundles whose sections correspond to vector fields and 1-forms as TM and T^*M , respectively. The basic idea of generalized geometry is to combine vectors and 1-forms into a single object. So, we consider the direct sum of tangent and cotangent bundles $E = TM \oplus T^*M$ which is called the generalized tangent bundle. Sections of E are called generalized vectors and a generalized vector $\mathcal{X} \in \Gamma E$ can be written in terms of a vector field $X \in \Gamma TM$ and a 1-form $\xi \in \Gamma T^*M$ as

$$\mathcal{X} = X + \xi$$

or we can also denote it in the matrix form as $\mathcal{X} = \begin{pmatrix} X \\ \xi \end{pmatrix}$. On E , one can define a bilinear form \langle, \rangle which is written for generalized vectors $\mathcal{X} = X + \xi$ and $\mathcal{Y} = Y + \eta$ as follows

$$\langle \mathcal{X}, \mathcal{Y} \rangle = \frac{1}{2}(i_X \eta + i_Y \xi) \quad (1.1)$$

where i_X denotes the interior derivative or contraction with respect to the vector field X . So, for a single generalized vector, we have $\langle \mathcal{X}, \mathcal{X} \rangle = i_X \xi$. The symmetry group of this bilinear form is the orthogonal group

$$O(TM \oplus T^*M) = \{\mathbb{A} \in GL(TM \oplus T^*M) \mid \langle \mathbb{A} \cdot, \mathbb{A} \cdot \rangle = \langle \cdot, \cdot \rangle\}$$

and this bilinear form has signature (n, n) . We can see this as follows. For the basis vectors $\{X_a\} \in \Gamma TM$ and basis 1-forms $\{e^a\} \in \Gamma T^*M$ for $a = 1, \dots, n$, we can define a basis for generalized vectors as

$$\mathcal{X}_A = X_a + \epsilon e_a$$

where $A = 1, \dots, 2n$ and $\epsilon = \pm 1$. For $\epsilon = 1$, \langle, \rangle gives a + sign and for $\epsilon = -1$, \langle, \rangle gives a - sign. The elements of $O(n, n)$ are in the form of

$$\mathbb{Q} = \begin{pmatrix} \mathcal{A} & \beta \\ B & -\mathcal{A}^T \end{pmatrix}$$

where

$$\begin{aligned} \mathcal{A} : \Gamma TM &\rightarrow \Gamma TM & , & & \mathcal{A}^T : \Gamma T^*M &\rightarrow \Gamma T^*M \\ \beta : \Gamma T^*M &\rightarrow \Gamma T^*M & , & & B : \Gamma TM &\rightarrow \Gamma T^*M \end{aligned}$$

and we can write $B \in \Lambda^2 T^*M$ as a 2-form and $\beta \in \Lambda^2 TM$ as a 2-vector. The invariance of \langle, \rangle under $O(n, n)$ gives shear transformations for generalized vectors. We have B - and β -transforms of \mathcal{X} as follows

$$e^B(X + \xi) = X + \xi + i_X B \quad (1.2)$$

$$e^\beta(X + \xi) = X + i_\xi \beta + \xi. \quad (1.3)$$

A metric g defined on M can be seen as a map $g : TM \rightarrow T^*M$ which is invertible. So, we can define a generalized metric on E induced by g as follows

$$\mathcal{G} = \begin{pmatrix} 0 & g^{-1} \\ g & 0 \end{pmatrix}. \quad (1.4)$$

So, the \mathcal{G} -dual of a generalized vector $\mathcal{X} = X + \xi$ can be written as

$$\tilde{\mathcal{X}} = \mathcal{G}(\mathcal{X}) = \begin{pmatrix} \tilde{X} \\ \tilde{\xi} \end{pmatrix} \in \Gamma E^*$$

where \tilde{X} denotes the g -dual 1-form and $\tilde{\xi}$ denotes the g -dual vector field and $E^* \simeq E$. In general, the B -transform of \mathcal{G} can be written as

$$\mathcal{G}_B = \begin{pmatrix} -g^{-1}B & g^{-1} \\ g - Bg^{-1}B & Bg^{-1} \end{pmatrix} \quad (1.5)$$

which is the generalized metric induced by g and B . Since $\mathcal{G}^2 = I$, ± 1 eigenspaces of \mathcal{G} which are denoted by V^\pm give a metric splitting of E

$$TM \oplus T^*M = V_+ \oplus V_-$$

corresponding to maximally positive/negative definite subbundles. The generalized metric \mathcal{G} can be written in terms of the bilinear form restricted to V_\pm as follows

$$\mathcal{G}(,) = \langle, \rangle_+ - \langle, \rangle_-. \quad (1.6)$$

We can also define an admissible metric for which the choices of V_\pm correspond to

$$V_\pm = \left\{ X \pm \tilde{X} : X \in \Gamma TM \right\}.$$

Similar to the Lie bracket $[\cdot, \cdot]$ of vector fields on ΓTM , we can also define a bracket operation on E called Courant bracket as $[\cdot, \cdot]_C : \Lambda^2 E \rightarrow \Gamma E$. For two generalized vectors $\mathcal{X} = X + \xi$ and $\mathcal{Y} = Y + \eta$, Courant bracket is defined as

$$[\mathcal{X}, \mathcal{Y}]_C = [X, Y] + \mathcal{L}_X \eta - \mathcal{L}_Y \xi - \frac{1}{2} d(i_X \eta - i_Y \xi) \quad (1.7)$$

where \mathcal{L}_X is the Lie derivative with respect to the vector field X and d is the exterior derivative. If we define the anchor map $\pi : \Gamma E \rightarrow \Gamma TM$ as $\pi(\mathcal{X}) = X$, then the Courant bracket satisfies

$$\pi([\mathcal{X}, \mathcal{Y}]_C) = [\pi(\mathcal{X}), \pi(\mathcal{Y})] \quad (1.8)$$

which can easily be seen from the definition. Although the Courant bracket is an antisymmetric bracket, it does not satisfy the Jacobi identity and hence does not correspond to a Lie bracket. However, with the definition of the Courant bracket, E admits a Courant algebroid structure. For $\mathcal{X}, \mathcal{Y}, \mathcal{Z} \in \Gamma E$, a Courant algebroid $(E, [,]_C, \langle, \rangle, \pi)$ satisfies the following properties

$$\begin{aligned} \pi([\mathcal{X}, \mathcal{Y}]_C) &= [\pi(\mathcal{X}), \pi(\mathcal{Y})] \\ [\mathcal{X}, f\mathcal{Y}]_C &= f[\mathcal{X}, \mathcal{Y}]_C + (\pi(X)f)\mathcal{Y} - \langle \mathcal{X}, \mathcal{Y} \rangle df \\ \langle df, dg \rangle &= 0 \\ [\mathcal{X}, [\mathcal{Y}, \mathcal{Z}]_C]_C + [\mathcal{Y}, [\mathcal{Z}, \mathcal{X}]_C]_C + [\mathcal{Z}, [\mathcal{X}, \mathcal{Y}]_C]_C &= dNij(\mathcal{X}, \mathcal{Y}, \mathcal{Z}) \end{aligned}$$

which is satisfied by the generalized tangent bundle $E = TM \oplus T^*M$. Here Nij denotes the Nijenhuis operator defined by

$$Nij(\mathcal{X}, \mathcal{Y}, \mathcal{Z}) = \frac{1}{3} (\langle [\mathcal{X}, \mathcal{Y}]_C, \mathcal{Z} \rangle + \langle [\mathcal{Y}, \mathcal{Z}]_C, \mathcal{X} \rangle + \langle [\mathcal{Z}, \mathcal{X}]_C, \mathcal{Y} \rangle).$$

On E , one can define an isotropic splitting $s : \Gamma TM \rightarrow \Gamma E$ and this determines a closed 3-form H on M given by

$$H(X, Y, Z) = (i_Y i_X H)(Z) = \langle [s(X), s(Y)]_C, s(Z) \rangle \quad (1.9)$$

for $X, Y, Z \in \Gamma TM$. The cohomology class of H classifies the Courant algebroids on M up to isomorphism. In the presence of non-zero H , one can modify the Courant bracket and define the twisted Courant bracket $[,]_H$ as follows

$$[\mathcal{X}, \mathcal{Y}]_H = [\mathcal{X}, \mathcal{Y}]_C - i_X i_Y H. \quad (1.10)$$

$[,]_H$ also defines a Courant algebroid structure for $dH = 0$. Another bracket operation on E which satisfies the Jacobi identity but not antisymmetric is the Dorfman bracket $[,]_D$ and it is defined as

$$[\mathcal{X}, \mathcal{Y}]_D = [X, Y] + \mathcal{L}_X \eta - i_Y d\xi. \quad (1.11)$$

Courant bracket correspond to the antisymmetrization of the Dorfman bracket and it can also be written as

$$[\mathcal{X}, \mathcal{Y}]_C = [\mathcal{X}, \mathcal{Y}]_D - d \langle \mathcal{X}, \mathcal{Y} \rangle. \quad (1.12)$$

Dorfman bracket originates from the Lie derivative on differential forms and it is also considered as the generalized Lie derivative

$$\mathbb{L}_X \mathcal{Y} = [\mathcal{X}, \mathcal{Y}]_D$$

and it satisfies

$$\begin{aligned}\mathbb{L}_{\mathcal{X}}(f\mathcal{Y}) &= (\pi(\mathcal{X})f)\mathcal{Y} + f\mathbb{L}_{\mathcal{X}}\mathcal{Y} \\ [\mathbb{L}_{\mathcal{X}}, \mathbb{L}_{\mathcal{Y}}] &= \mathbb{L}_{[\mathcal{X}, \mathcal{Y}]_D} \\ \mathbb{L}_{\mathcal{X}}(\mathcal{Y} \otimes \mathcal{Z}) &= \mathbb{L}_{\mathcal{X}}\mathcal{Y} \otimes \mathcal{Z} + \mathcal{Y} \otimes \mathbb{L}_{\mathcal{X}}\mathcal{Z}.\end{aligned}\tag{1.13}$$

We can define a generalized connection \mathbb{D} on E which is compatible with the bilinear form \langle, \rangle corresponding to the linear operator

$$\mathbb{D} : \Gamma E \rightarrow \Gamma(E^* \otimes E)$$

satisfying the following identities

$$\mathbb{D}_{\mathcal{X}}(f\mathcal{Y}) = \pi(\mathcal{X})(f)\mathcal{Y} + f\mathbb{D}_{\mathcal{X}}\mathcal{Y}\tag{1.14}$$

$$\pi(\mathcal{X})\langle \mathcal{Y}, \mathcal{Z} \rangle = \langle \mathbb{D}_{\mathcal{X}}\mathcal{Y}, \mathcal{Z} \rangle + \langle \mathcal{Y}, \mathbb{D}_{\mathcal{X}}\mathcal{Z} \rangle\tag{1.15}$$

where f is a function, $\mathcal{X}, \mathcal{Y}, \mathcal{Z} \in \Gamma E$ and we denote $\mathbb{D}_{\mathcal{X}}\mathcal{Y} = i_{\mathcal{X}}\mathbb{D}(\mathcal{Y})$. Here i is the contraction operator of elements in ΓE^* with elements in ΓE . Given a standard connection ∇ on TM , we can define the generalized connection \mathbb{D}^{∇} induced by ∇ as

$$\mathbb{D}^{\nabla} = \nabla \oplus \nabla$$

corresponding to

$$\mathbb{D}_{\mathcal{X}}^{\nabla}\mathcal{Y} = \nabla_{\pi(\mathcal{X})}\mathcal{Y}.\tag{1.16}$$

However, there is a freedom in choosing a connection as $\mathbb{D}' = \mathbb{D} + \alpha$ also defines a connection where $\alpha \in \Gamma(E^* \otimes o(E))$ and $o(E)$ denotes the bundle of skew-symmetric endomorphisms of E w.r.t. \langle, \rangle . However, by considering relevant torsion-free, metric compatible and divergence-fixing conditions, we can define a unique connection on E [ref].