Formation of dark matter solar halos

DESY Summer Student programme 2024

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HELMHOLTZ



Content



Introduction

What is Dark Matter?

Evidence of dark matter in the Universe

- Rotation velocity of the stars in the galaxies;
- High-speed, gravitationally bound clusters of galaxies;
- Gravitational Lensing;
- Observations from cluster collisions;

A lot of different candidates! See Dark Matter Lecture

Why dark matter?

- Predicted production during Early Universe
- · Axions could explain other problems, such as in QCD





How much DM is in our Solar System?

DM density in our galaxy averaged on scales much bigger than the Solar System

Loose upper bounds from Solar System ephemerides from:

- planets (arXiv: 0903.4849)
- asteroids (arXiv: 2210.03749)

From arXiv: 2306.12477 [hep-ph]



Can the Sun gravitationally capture DM?



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Can the Sun gravitationally capture DM?

Non-relativistic approximation

$$\phi = \frac{1}{\sqrt{2m}} \left(\psi e^{-imt} + \text{c.c.} \right) \quad \text{with } \dot{\psi} \ll m\psi$$

 \rightarrow Schroedinger-Poisson equation

$$\left(i\partial_t + \frac{\nabla^2}{2m} + \frac{\alpha}{r}\right)\psi = g|\psi|^2\psi$$

From arXiv: 2306.12477 [hep-ph]

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+ gravitational potential

$$\Phi = -\frac{GM}{r}$$

$$\alpha = GMm$$
$$\alpha_{em} = \frac{e^2}{4\pi\epsilon_0}$$

Can the Sun gravitationally capture DM?

Analogy with Hydrogen atom $(g = 0) \rightarrow$ expansion in its eigenfunctions

$$\psi(\boldsymbol{x},t) = \sum_{nlm} c_{nlm}(t) e^{-i\omega_{nlm}t} \psi_{nlm}(\boldsymbol{x}) + \int [dk] c_{\boldsymbol{k}}(t) e^{-i\omega_{\boldsymbol{k}}t} \psi_{\boldsymbol{k}}(\boldsymbol{x})$$

and a lot of calculations (and sweat)...

ULDM: coherence time $\tau_{\rm dm} = \frac{2\pi}{mv_{\rm dm}^2}$ after which uncorrelated waves: $\langle a^*(\boldsymbol{k})a(\boldsymbol{k}')\rangle \equiv (2\pi)^3 f(\boldsymbol{k})\delta(\boldsymbol{k}-\boldsymbol{k}')$ use $\langle \cdot \rangle$ over times much larger than $\tau_{\rm dm}$

From arXiv: 2306.12477 [hep-ph]

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Analogy with Hydrogen atom $(g = 0) \rightarrow$ expansion in its eigenfunctions

$$\psi(\boldsymbol{x},t) = \sum_{nlm} c_{nlm}(t) e^{-i\omega_{nlm}t} \psi_{nlm}(\boldsymbol{x}) + \int [dk] c_{\boldsymbol{k}}(t) e^{-i\omega_{\boldsymbol{k}}t} \psi_{\boldsymbol{k}}(\boldsymbol{x})$$

and a lot of calculations (and sweat)...

$$i\dot{c}_{nlm}^{(1)}(t) + \sqrt{N_{nlm}^{(0)}}\omega_{nlm}^{(1)} = ge^{i\omega_n t} \int d^3x \left|\psi^{(0)}\right|^2 \psi^{(0)}\psi_{nlm}^*$$

 $2 \rightarrow 2$ scattering processes

From arXiv: 2306.12477 [hep-ph]

Can the Sun gravitationally capture DM?

Initial state: waves + ground state 100

$$\psi^{(0)} = \sqrt{N_{100}^{(0)}} e^{-i\omega_1 t} \psi_{100} + \int [dk] a(\mathbf{k}) e^{-i\omega_k t} \psi_{\mathbf{k}}$$

1st order: $2 \rightarrow 2$ scattering processes. 5 possible diagrams:



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Dark photons (spin-1 bosons)

Can the Sun gravitationally capture DP?



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Dark photons (spin-1 bosons)

Can the Sun gravitationally capture DP?

Non-relativistic limit + gravitational potential \rightarrow

$$\Rightarrow \left(i\partial_t + \frac{\nabla^2}{2m} + \frac{\alpha}{r}\right)\psi^i = g\sum_j \left(-2|\psi^j|^2\psi^i - (\psi^j)^2\psi^{i*}\right)$$

Same as for spin-0 bosons (but more sweat)... and finally

$$i\dot{c}_{100}^{i(1)}(t) = ge^{i\omega_{1}t} \sum_{j} \int d^{3}x \left(-2 \underbrace{|\psi^{j(0)}|^{2} \psi^{i(0)} \psi_{100}^{i*}}_{k_{3}^{j}} - \underbrace{(\psi^{j(0)})^{2} \psi^{i(0)*} \psi_{100}^{i*}}_{k_{3}^{j}}\right)$$

$$j = 1, 2, 3$$

$$k_{2}^{i} \underbrace{(3.1)}_{(3.1)} 100^{i}} \underbrace{k_{2}^{j} \underbrace{(3.2)}_{(3.2)} 100^{i}}_{k_{2}^{j}} \underbrace{k_{3}^{j}}_{(3.2)} \underbrace{(y^{j(0)})^{2} \psi^{i(0)*} \psi_{100}^{i*}}_{k_{3}^{j}}}_{(3.2)} \underbrace{(y^{j(0)})^{2} \psi^{i(0)*} \psi_{100}^{i*}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{(3.2)} \underbrace{(y^{j(0)})^{2} \psi^{i(0)*} \psi_{100}^{i*}}_{k_{3}^{j}}}_{(3.2)} \underbrace{(y^{j(0)})^{2} \psi^{i(0)*} \psi_{100}^{i*}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{(3.2)} \underbrace{(y^{j(0)})^{2} \psi^{i(0)*} \psi_{100}^{i*}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}_{k_{3}^{j}}}$$

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Dark photons (spin-1 bosons)

Can the Sun gravitationally capture DP?

But... problem! Non-zero interference (3.1)x(3.2)* complicates calculation Solution: interference from equal diagrams \rightarrow separate the culprit (j = i)

 \rightarrow 3 new diagrams \rightarrow Bose enhancement method \checkmark



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$$\begin{aligned} \begin{array}{l} \begin{array}{l} \begin{array}{l} \text{Dark photons (spin-1 bosons)} \\ \dot{N}_{100}^{i} = \dot{N}_{100}^{i(3.3)} + \dot{N}_{100}^{i(3.4)} + \dot{N}_{100}^{i(3.5)} \\ \dot{N}_{100}^{i(3.3)} = \frac{4}{27} \sum_{j \neq i} g^{2} \int [dk_{1}][dk_{2}][dk_{3}](2\pi)\delta(\Delta\omega)|\mathcal{M}|^{2} \times \\ & \times \left\{ f^{j}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2})f^{j}(\mathbf{k}_{3}) \\ & + N_{100}^{i}[f^{j}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{j}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{j}(\mathbf{k}_{3})] \right\} \\ \dot{N}_{100}^{i(3.4)} = \frac{2}{27} \sum_{j \neq i} g^{2} \int [dk_{1}][dk_{2}][dk_{3}](2\pi)\delta(\Delta\omega)|\mathcal{M}|^{2} \times \\ & \times \left\{ f^{j}(\mathbf{k}_{1})f^{j}(\mathbf{k}_{2})f^{i}(\mathbf{k}_{3}) \\ & + N_{100}^{i}[f^{j}(\mathbf{k}_{1})f^{j}(\mathbf{k}_{2}) - (f^{j}(\mathbf{k}_{1}) + f^{j}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3})] \right\} \\ \dot{N}_{100}^{i(3.5)} = \frac{18}{27} \quad g^{2} \int [dk_{1}][dk_{2}][dk_{3}](2\pi)\delta(\Delta\omega)|\mathcal{M}|^{2} \times \\ & \times \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2})f^{i}(\mathbf{k}_{3}) \\ & + N_{100}^{i}[f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{j}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3})] \right\} \\ \hline k_{100}^{i} = \frac{18}{27} \quad g^{2} \int [dk_{1}][dk_{2}][dk_{3}](2\pi)\delta(\Delta\omega)|\mathcal{M}|^{2} \times \\ & \times \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2})f^{i}(\mathbf{k}_{3}) \\ & + N_{100}^{i}[f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3})] \right\} \\ \hline k_{100}^{i} = \frac{18}{27} \quad g^{2} \int [dk_{1}][dk_{2}][dk_{3}](2\pi)\delta(\Delta\omega)|\mathcal{M}|^{2} \times \\ & \times \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3}) \right\} \\ \hline k_{100}^{i} = \frac{18}{100} \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3}) \right\} \right\} \\ \hline k_{100}^{i} = \frac{18}{100} \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3}) \right\} \\ \hline k_{100}^{i} = \frac{18}{100} \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3}) \right\} \\ \hline k_{100}^{i} = \frac{18}{100} \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3}) \right\} \\ \hline k_{100}^{i} = \frac{18}{100} \left\{ f^{i}(\mathbf{k}_{1})f^{i}(\mathbf{k}_{2}) - (f^{i}(\mathbf{k}_{1}) + f^{i}(\mathbf{k}_{2}))f^{i}(\mathbf{k}_{3}) \right\} \\ \hline k_{100}^{i} = \frac{18}{100} \left\{ f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^{i}(\mathbf{k})f^$$

What's next?

- Typical regimes for DP solar halo formation
 - 2 possibilities:
 - symmetric initial condition
 - asymmetric initial condition
- DM solar halo <u>inside</u> the Sun (or other massive object)
 - ightarrow gravitational potential $\Phi \,\,$ \propto

depending on mass distribution $\rho_M(r)$ inside the Sun (massive object)

 $\frac{1}{r}$



Thank you. Questions?





Thank you.

Contact

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Dark photons (spin-1 bosons)
$$\langle a^{i*}(\mathbf{k})a^j(\mathbf{k}')\rangle \equiv (2\pi)^3 f^i(\mathbf{k})\delta(\mathbf{k}-\mathbf{k}')\frac{\delta_{ij}}{3}$$

Can the Sun gravitationally capture DP?

$$\mathcal{L} = -\frac{1}{4} F_{\rho\sigma} F^{\rho\sigma} - \frac{1}{2} m^2 A_{\rho} A^{\rho} + \underbrace{\frac{1}{\Lambda^4} (F_{\rho\sigma} F^{\rho\sigma})^2}_{\text{self-interaction}} \qquad \begin{array}{l} F^{\mu\nu} = \partial_{\mu} A^{\nu} - \partial_{\nu} A^{\mu} \\ \partial_{\mu} A^{\mu} = 0 \quad \text{(Lorentz condition)} \end{array}$$
$$(-\Box + m^2) A^i + \frac{8}{\Lambda^4} (\partial_{\mu} F^2) F^{\mu i} + \frac{8}{\Lambda^4} F^2 \Box A^i = 0$$
$$A^i = \frac{1}{\sqrt{2m}} \left(\psi^i e^{-imt} + \text{c.c.} \right) \quad \text{with } \partial_{\nu} \psi^i \ll m \psi^i$$

$$\left(i\partial_t + \frac{\nabla^2}{2m} + \frac{\alpha}{r}\right)\psi^i = g\sum_j \left(-2|\psi^j|^2\psi^i - (\psi^j)^2\psi^{i*}\right)$$