

Muophobic forces at the muon collider

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Finanziato
dall'Unione europea
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Italiadomani
PIANO NAZIONALE
DI RIPRESA E RESILIENZA

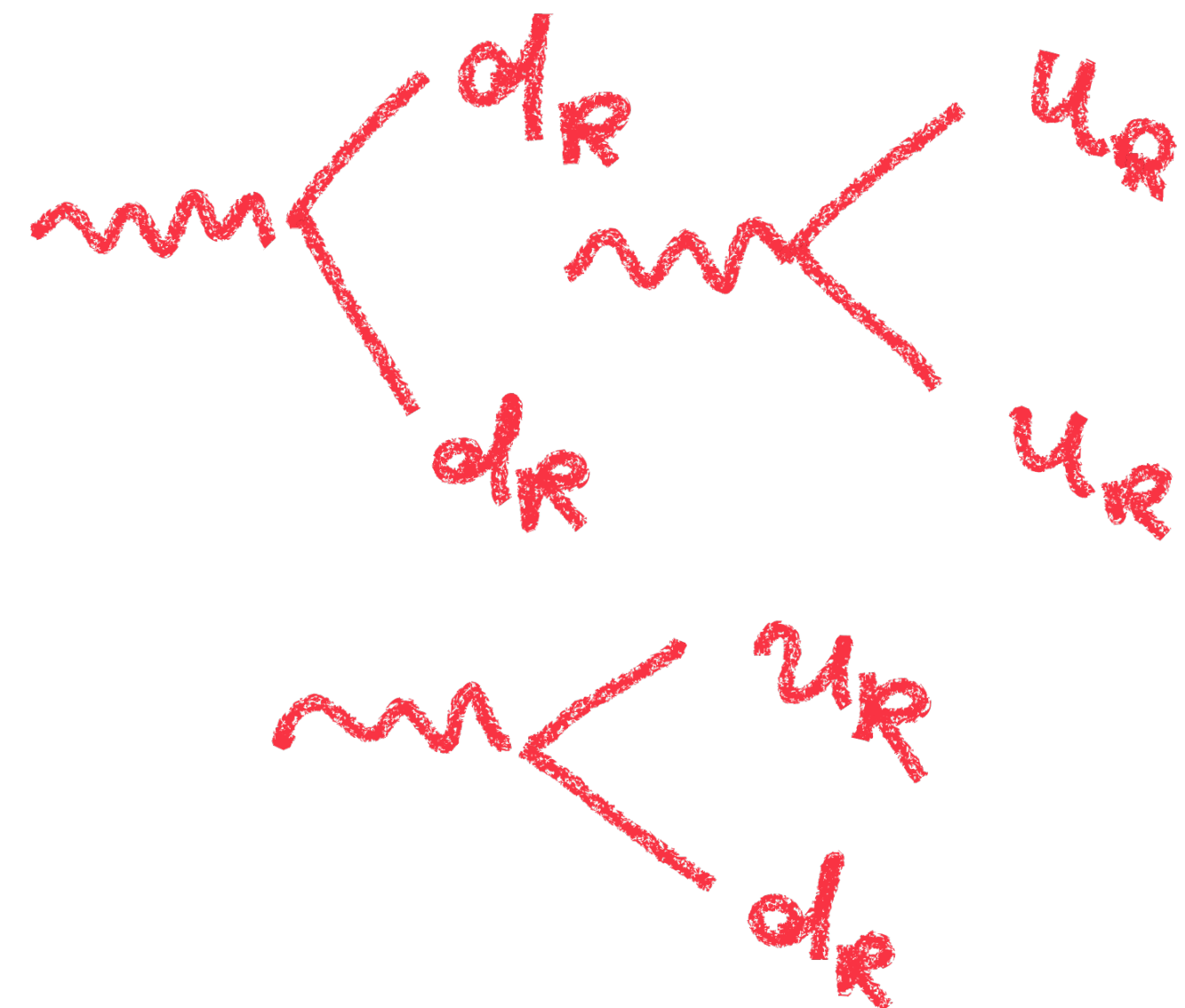


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New Gauge Forces in BSM

- New forces arise in motivated extensions of the SM
- They come with new gauge bosons (both electrically charged and neutral)
- Example is to group all the right-handed fields of the SM, that are singlets under the gauge group of the SM $SU(3)_C \times SU(2)_W \times U(1)_Y$, and make them n -plets of a new $SU(2)_R \supset U(1)_Y$

$$u_R \sim \mathbf{1}, d_R \sim \mathbf{1} \rightarrow \begin{pmatrix} u_R \\ d_R \end{pmatrix} \sim \mathbf{2}$$



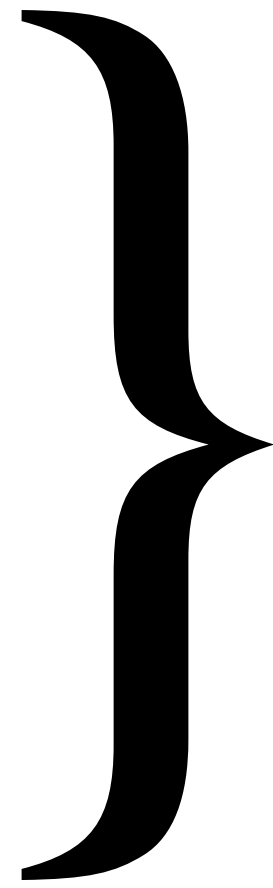
New Gauge Forces in BSM

connected to flavor puzzle (?)

$$B_3 - L_3$$

$$B_3 + L_3$$

$$\mathcal{U}_3$$



coupled only to bottom and top quark, and tau and its neutrino
(each scenario has distinct couplings pattern → dramatic consequences)

$$B_3$$

coupled only to bottom and top quark

$$L_3$$

coupled only to tau and its neutrino

New Gauge Forces in BSM

connected to flavor puzzle (?)

$$B_3 - L_3 \quad \begin{aligned} g_{Z'} &\equiv 3 \cdot g_{V,qL,3} = 3 \cdot g_{V,bR} = 3 \cdot g_{V,tR}, \\ -g_{Z'} &\equiv g_{V,\ell L,3} = g_{V,\tau R}. \end{aligned}$$

$$B_3 + L_3 \quad \begin{aligned} g_{Z'} &\equiv 3 \cdot g_{V,qL,3} = 3 \cdot g_{V,bR} = 3 \cdot g_{V,tR}, \\ g_{Z'} &\equiv g_{V,\ell L,3} = g_{V,\tau R}. \end{aligned}$$

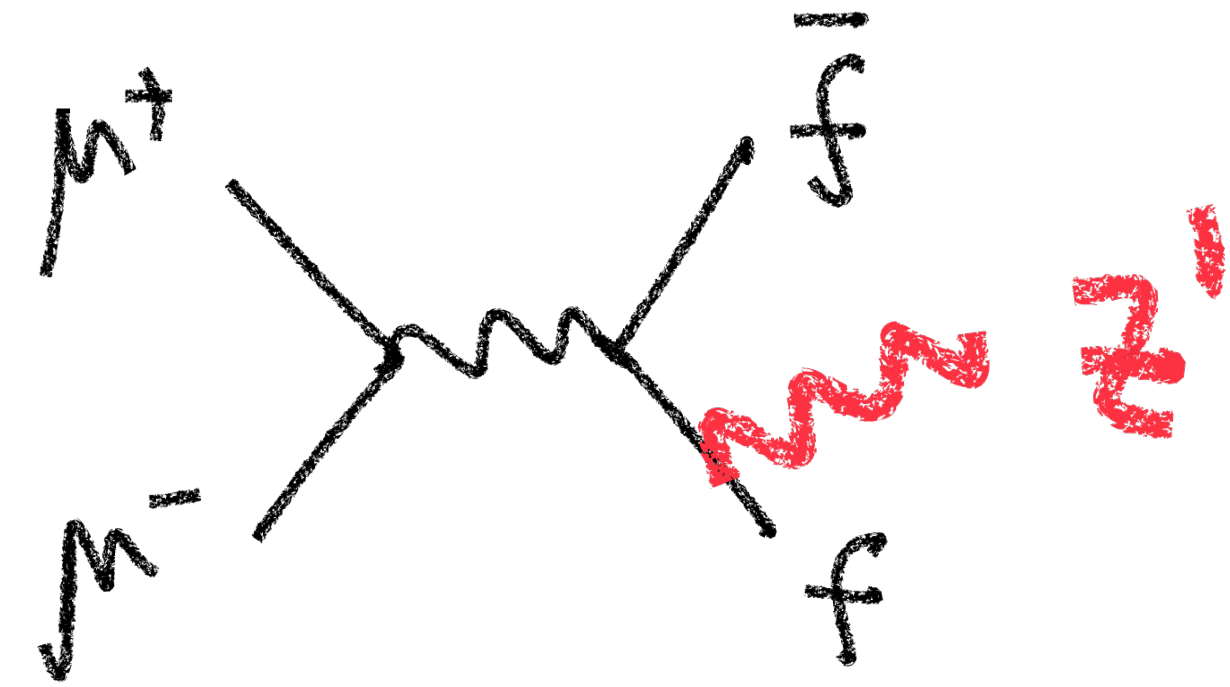
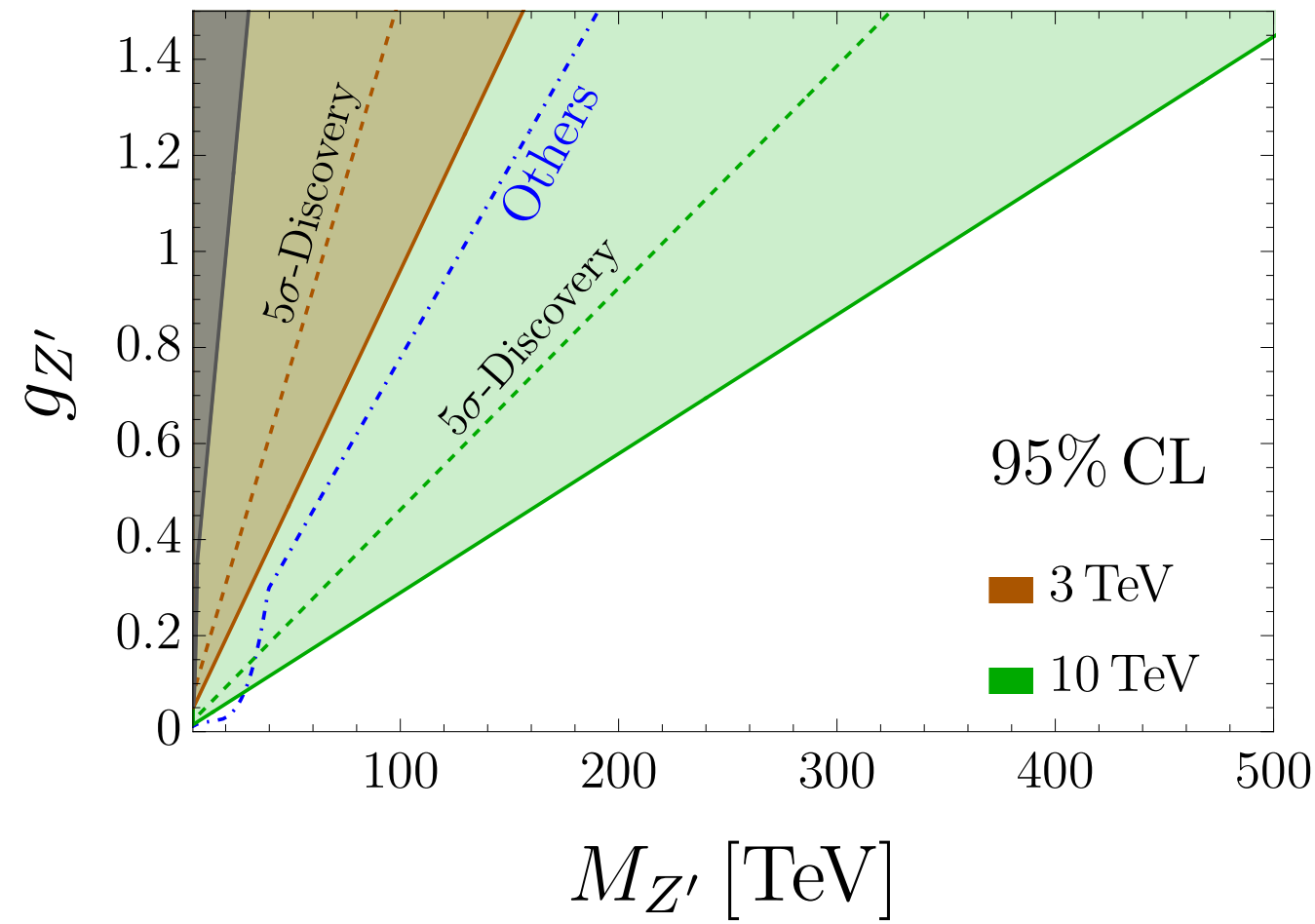
$$\mathcal{U}_3 \quad g_{Z'} \equiv g_{V,qL,3} = g_{V,bR} = g_{V,tR} = g_{V,\ell L,3} = g_{V,\tau R}.$$

$$B_3 \quad g_{Z'} \equiv 3 \cdot g_{V,qL,3} = 3 \cdot g_{V,bR} = 3 \cdot g_{V,tR}.$$

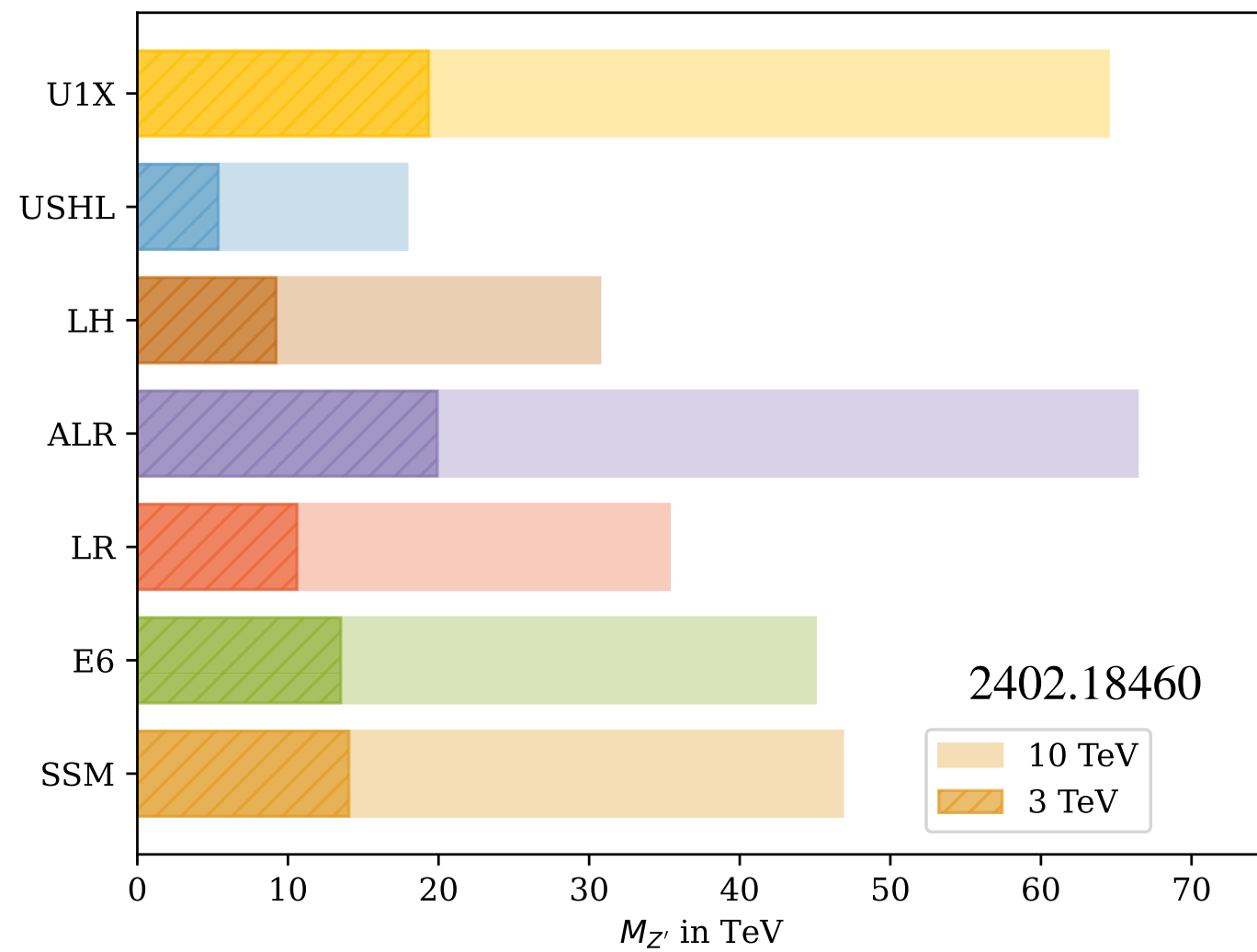
$$L_3 \quad g_{Z'} \equiv g_{V,\ell L,3} = g_{V,\tau R}.$$

New Gauge Forces in BSM

Indirect (Direct) search if heavier (lighter) than 10 TeV

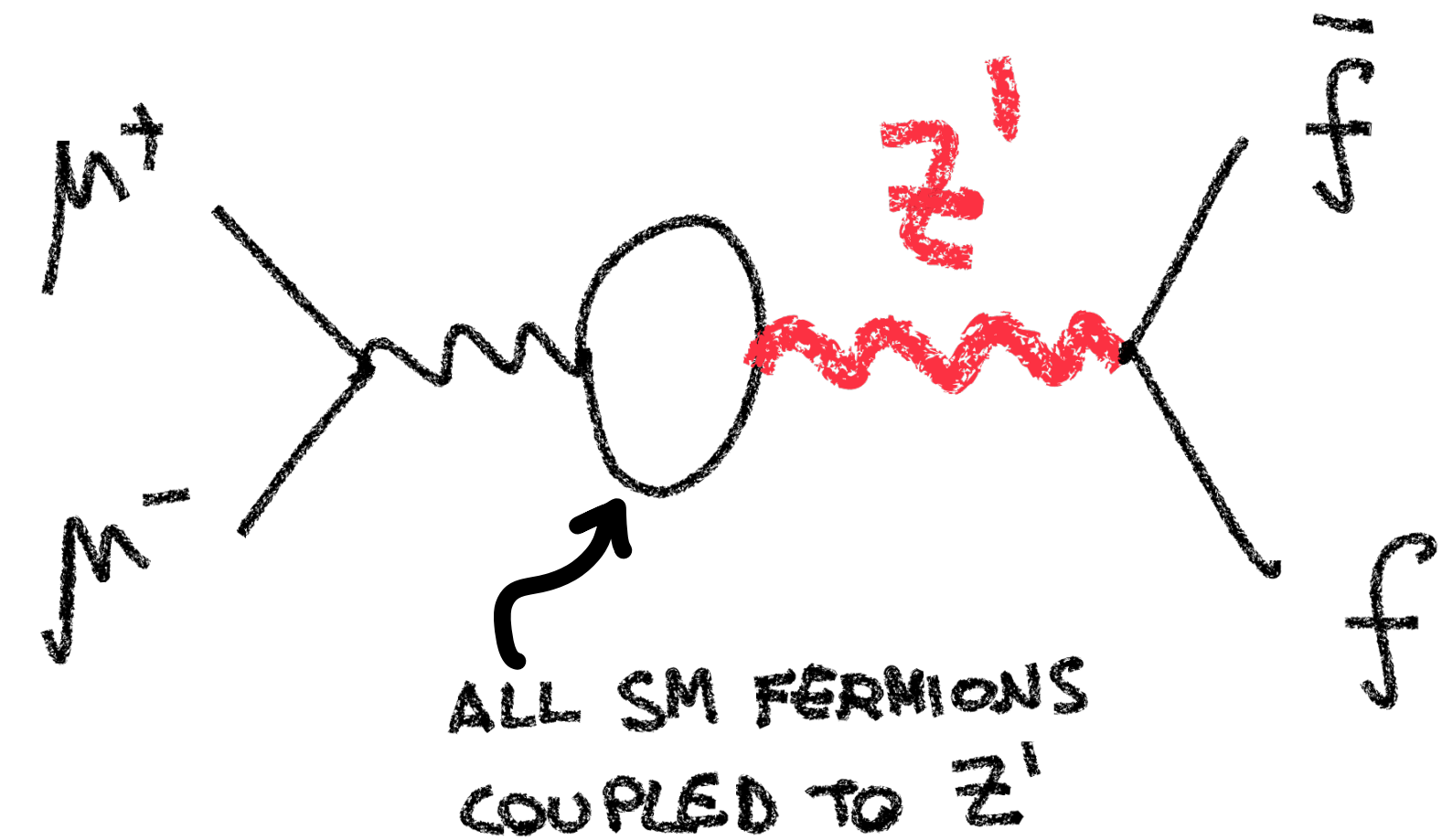


Do not need to assume muon coupling to Z'



Model	$g_{Z'}$	$2v_l$	$2a_l$
SSM	$\frac{e}{s_W c_W}$	$2s_W^2 - \frac{1}{2}$	$-\frac{1}{2}$
E_6	$\frac{e}{c_W}$	$\frac{2\cos\beta}{\sqrt{6}}$	$\frac{\cos\beta}{\sqrt{6}} + \frac{\sqrt{10}\sin\beta}{6}$
LR	$\frac{e}{c_W}$	$\frac{1}{\alpha} - \frac{\alpha}{2}$	$\frac{\alpha}{2}$
ALR	$\frac{e}{s_W c_W \sqrt{1-2s_W^2}}$	$\frac{5}{2}s_W^2 - 1$	$-\frac{1}{2}s_W^2$
LH	$\frac{e}{s_W}$	$-\frac{c}{4s}$	$-\frac{c}{4s}$
USLH	$\frac{e}{c_W \sqrt{3-4s_W^2}}$	$\frac{1}{2} - 2s_W^2$	$\frac{1}{2}$
$U(1)_X$	$\frac{e}{4c_W}$	-8	2

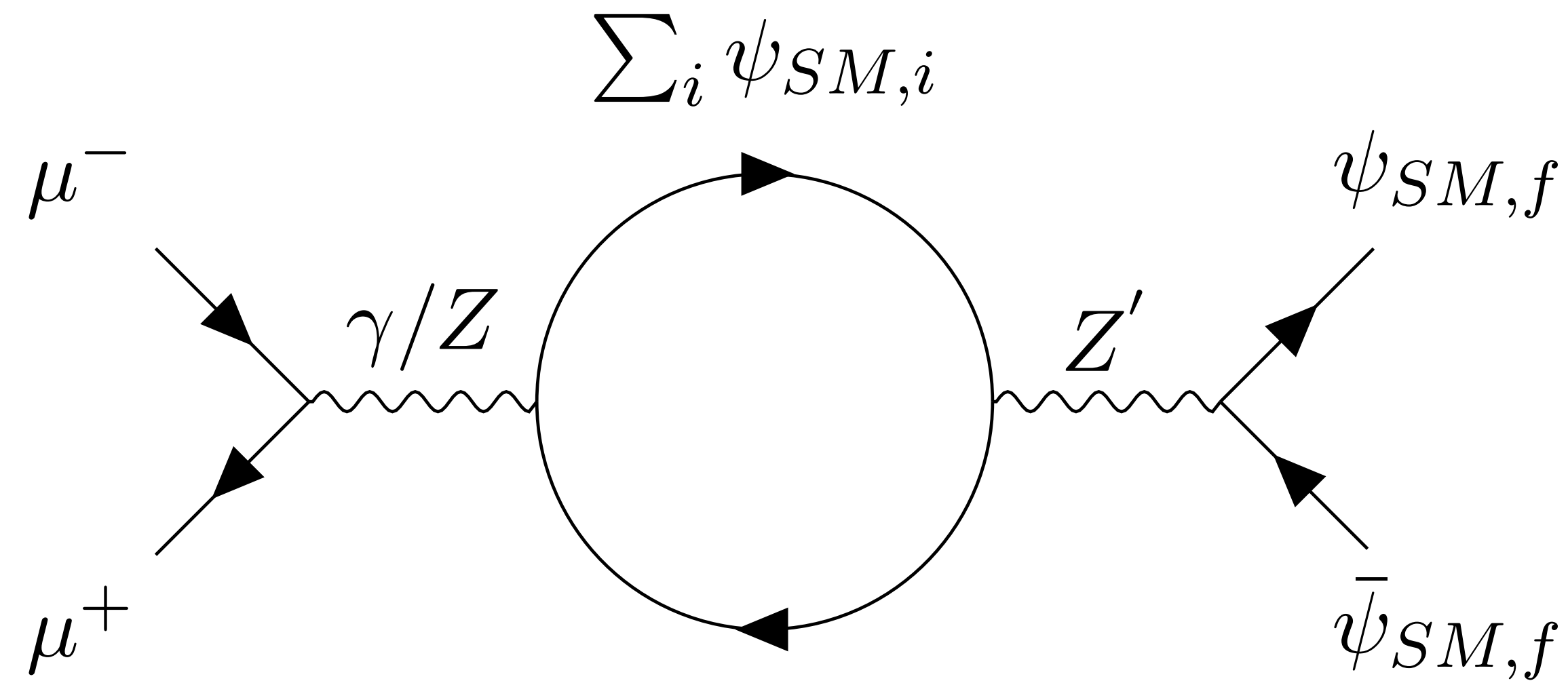
Table 1: The Z' couplings to leptons in the models considered in this paper. The sine of the Weinberg angle is denoted by s_W . For the description of the models and the explanation of the model-specific parameters see the main text. The table and notation are adapted from [22].



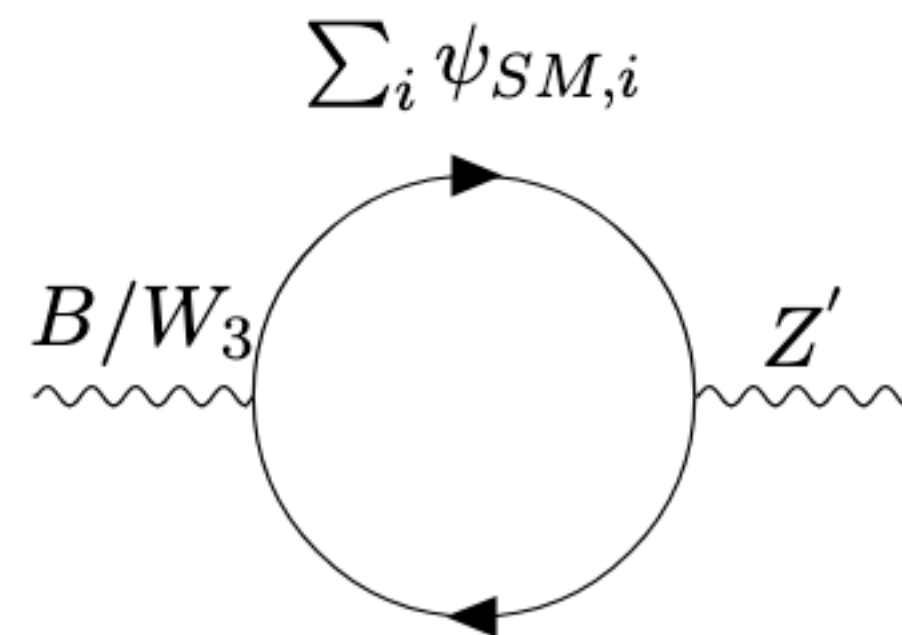
2312.17427 for 6 and 8 TeV results

New Gauge Forces in BSM

connected to flavor puzzle (?)



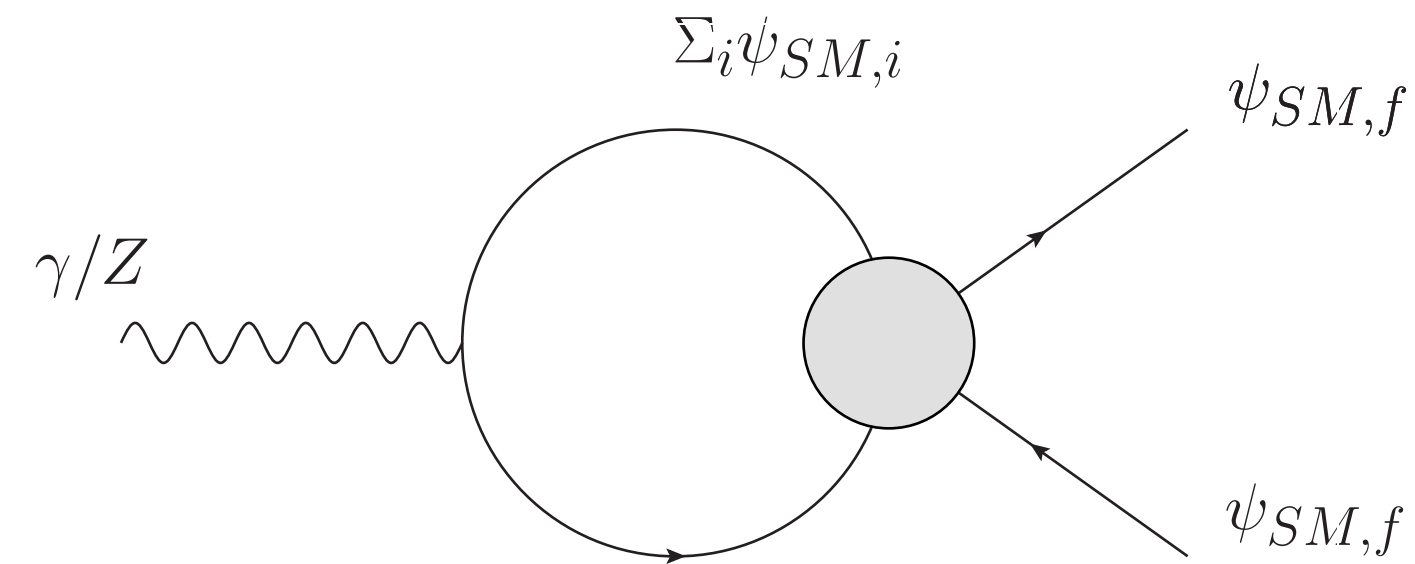
Even if Z' does not couple to muon directly, a coupling can be generated by loops



effectively we rely on a loop-level mixing between Z' and SM gauge bosons

New Gauge Forces in BSM

The crucial ingredient is the piece of the amplitude, which sums over intermediate fermions



Each intermediate state enters with its own SM charge and charge w.r.t. Z'

Cancellations of contributions are possible, and even expected in some cases

$$L_e - L_\tau$$

Handwritten Feynman diagrams illustrating the cancellation of contributions for the operator $L_e - L_\tau$. The first diagram shows a wavy line entering a loop with an electron (e) and two neutrinos (ν). The second diagram shows a wavy line entering a loop with a tau (τ) and two neutrinos (ν). The diagrams are subtracted, and the result is approximately $0 + \left(\frac{m_\tau}{\sqrt{s}}\right)^2$.

$$\mathcal{U}_3$$

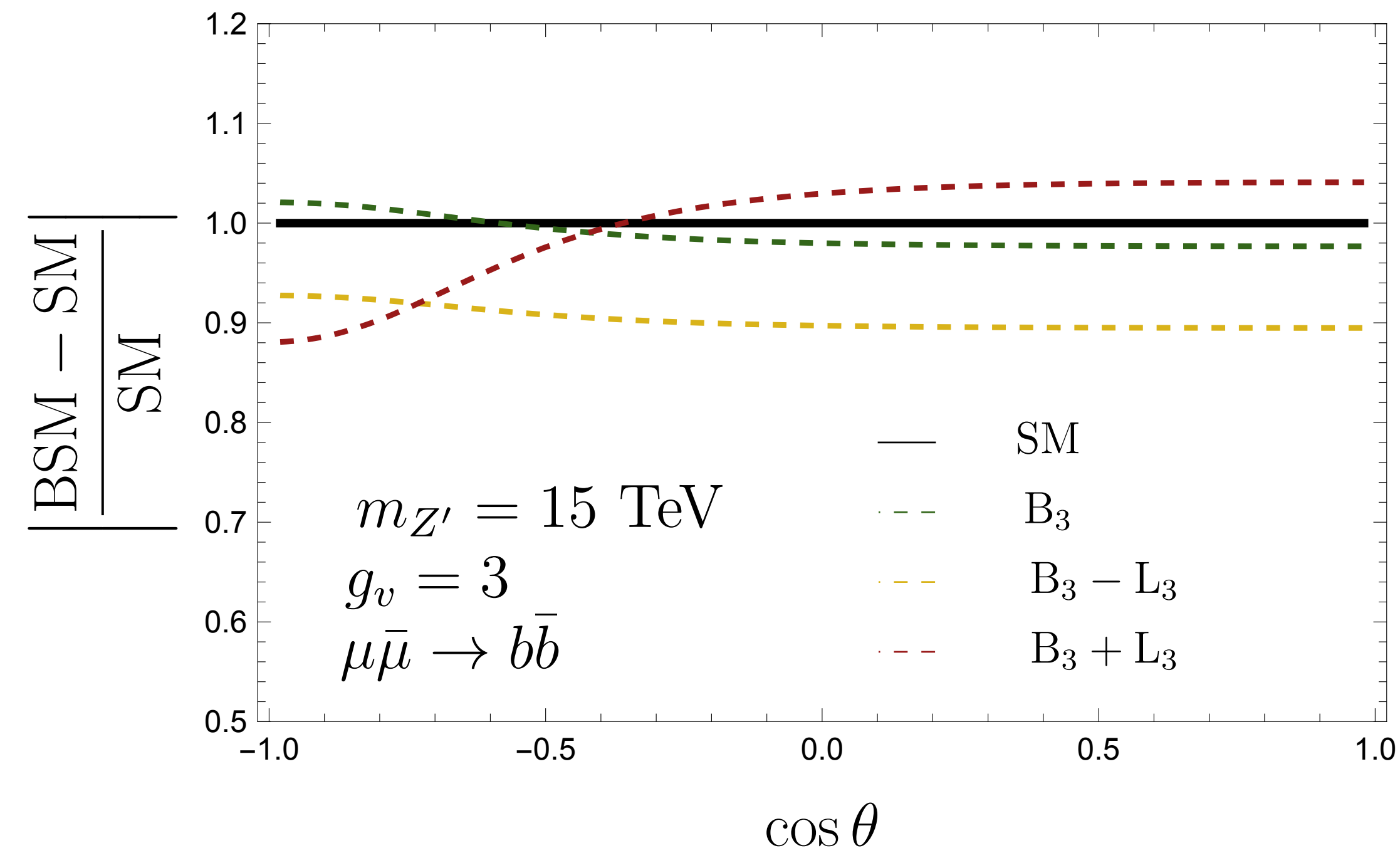
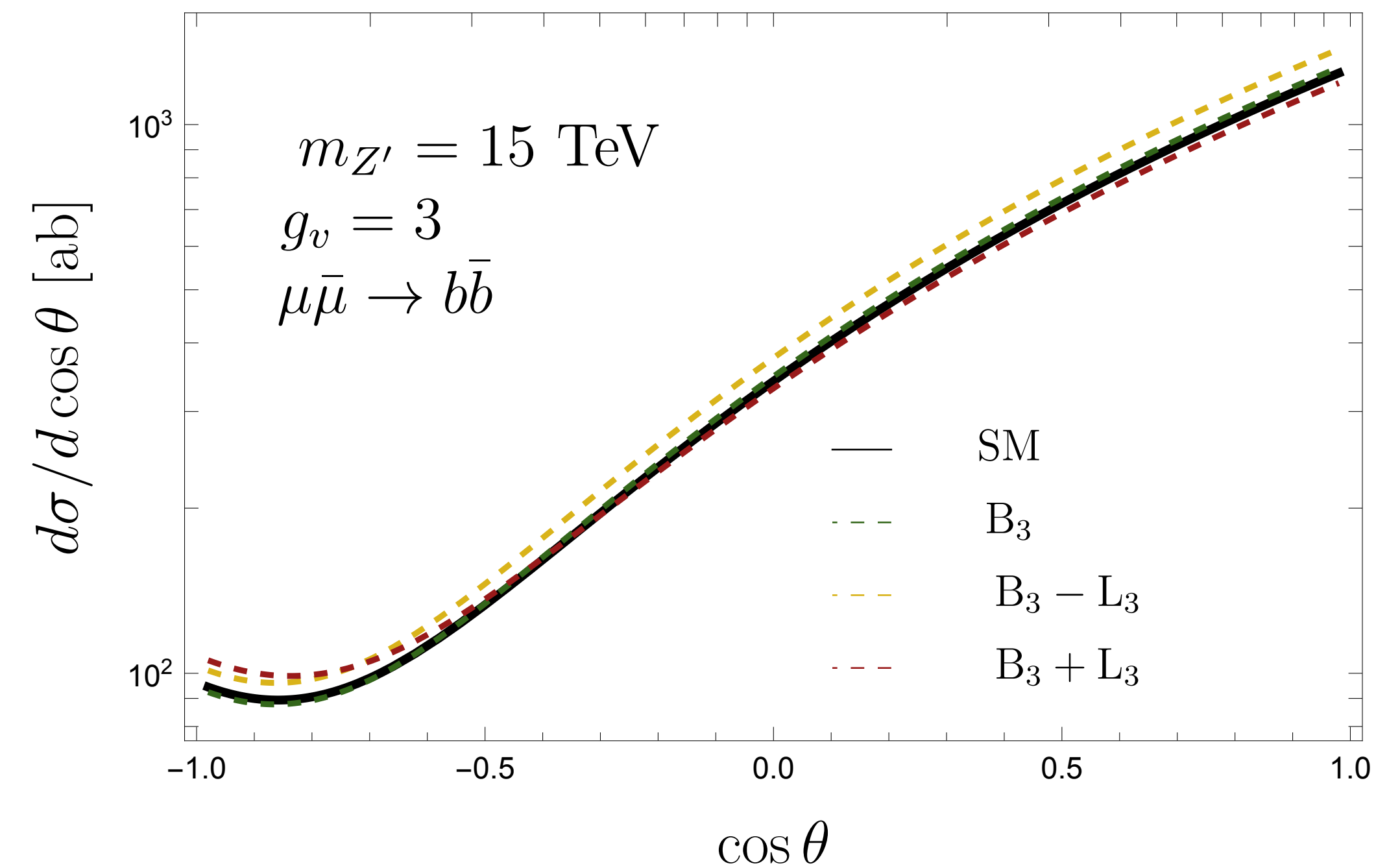
$$g_1 \sum_{SM} N_c Y_{SM} \equiv 0 + \left(\frac{m_t}{\sqrt{s}}\right)^2$$

$$SU(2)_R$$

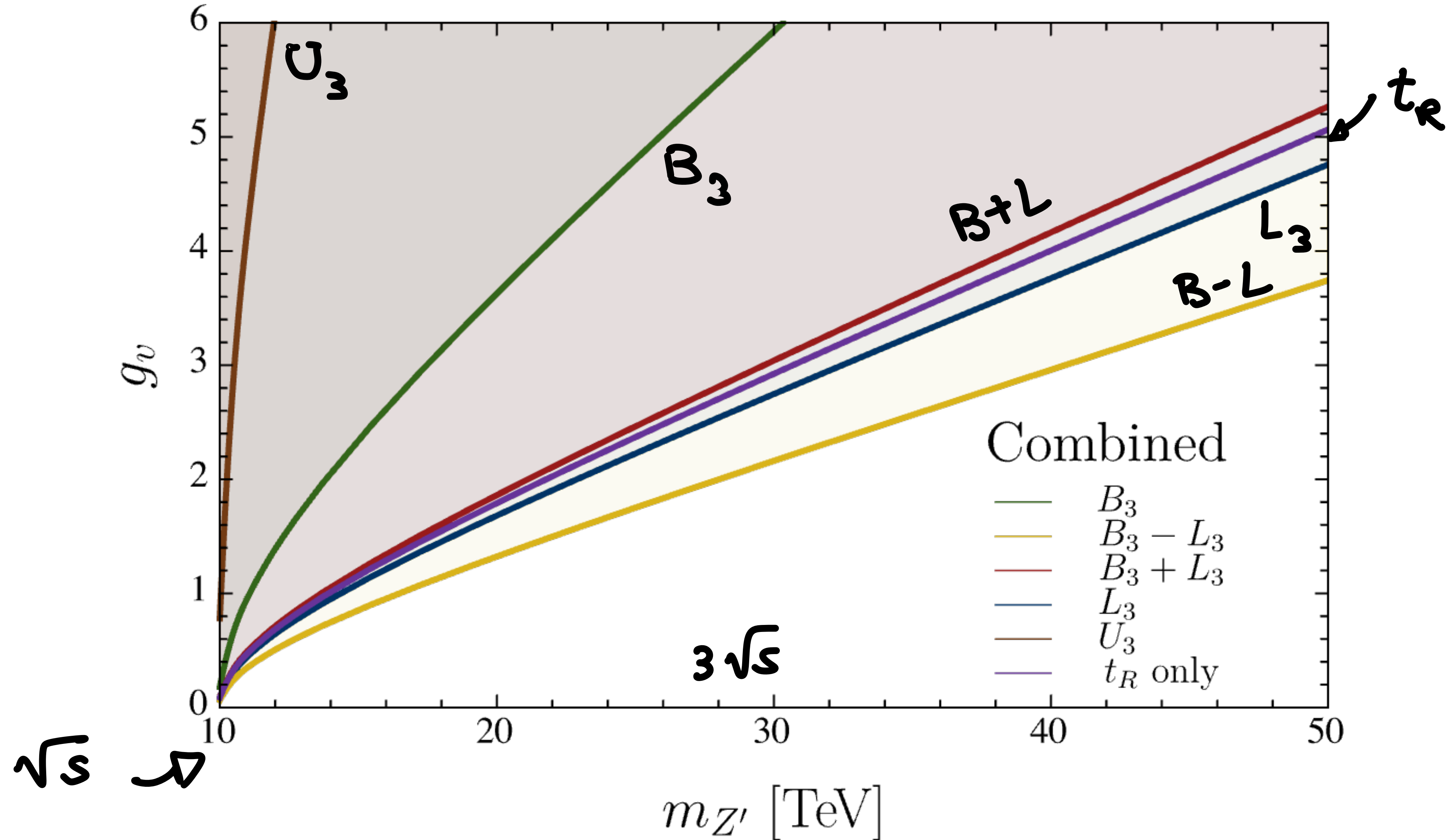
Handwritten Feynman diagrams illustrating the cancellation of contributions for the operator $SU(2)_R$. The first diagram shows a wavy line entering a loop with a right-handed bottom quark (b_R) and two right-handed top quarks (t_R). The second diagram shows a wavy line entering a loop with a right-handed top quark (t_R) and two right-handed bottom quarks (b_R). The diagrams are subtracted, and the result is approximately $0 + \left(\frac{m_t}{\sqrt{s}}\right)^2$.

Indirect sensitivity

- We do not deal with $m_{Z'} < 10$ TeV as this should be addressed by a family of direct searches
- We go for $m_{Z'} > 10$ TeV
- Consider $b\bar{b}$, $\tau\tau$, and $t\bar{t}$ final states and check the size of the deviation for a given Z' mass and coupling with respect to the SM.
- So for analysis limited to equal couplings to left- and right-handed fermions

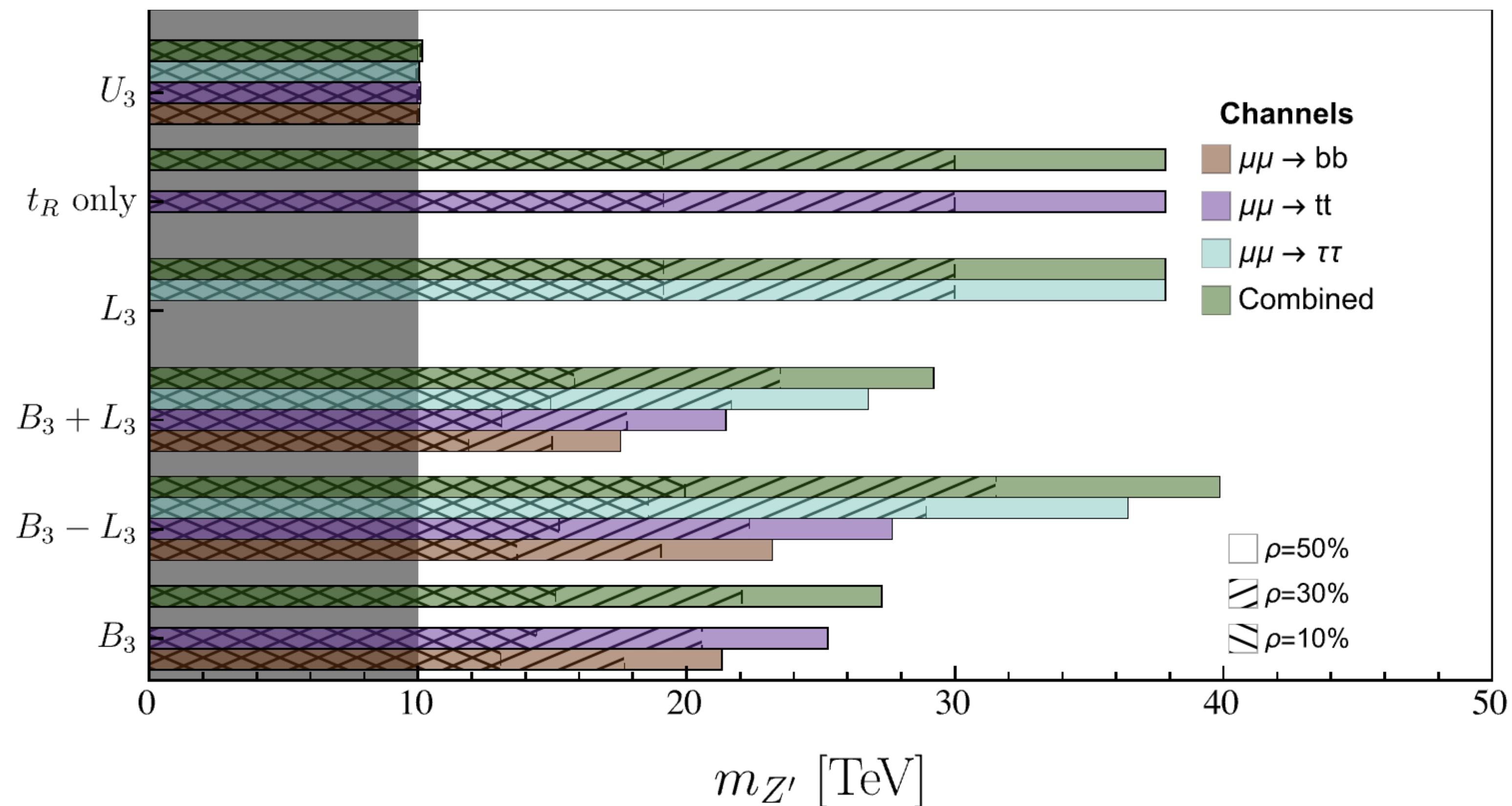


New Gauge Forces in BSM



Constraints from perturbativity

$$\Gamma(Z' \rightarrow ff) \propto \sum_f \frac{g_f^2}{8\pi} M_{Z'} < \rho M_{Z'} \quad \rho = [0.1, 0.5]$$



Outlook

- Physics case for the 10 TeV machine starts to be mature.
- Lots of refinement still needed (Detector Concept development, etc)
- Time to consider also “things you are not supposed to do” at the muon collider (muophobic and more!)

Conclusions

- New gauge forces are part of a wide set of extensions of the Standard Model
- They come in all size and shapes, hard to make a map and really focus on just one class of vector bosons. Can't really exclude any possibility a priori!
- If the new force carries is somehow not coupled to muons it might pose a problem for the muon collider! The typical example is a force mostly coupled to third generation fermions.
- Even in these “decoupled” scenarios muon collider is sensitive to loop effects
- Mass reach up to few times the center of mass energy for generic couplings structure!
- Some cases are problematic (cancellation of the interference, e.g. $L_e - L_\tau$, \mathcal{U}_3 , ...)
- Potential to connect with Dark Matter phenomenology (similar to $\mu^+ \mu^- \rightarrow \nu_\tau \nu_\tau + X$)

Thank you!

