

Modeling radio and X-ray afterglows of jetted tidal disruption events

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Deutsches Elektronen Synchrotron DESY

Hamburg “Matter and the Universe” Days
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TDE schematic, not to scale

Tidal disruption events

When a massive star passes close enough to a SMBH

- ~ half of the star's mass remains bounded by the SMBH gravitational force

Mass accretion -> relativistic jet -> months/year-long flare (optical transient)

- Energy to be reprocessed by accretion $\sim 10^{54} \text{ erg E}$

Fallback rate $\propto t^{-5/3}$ (Phinney 1989)

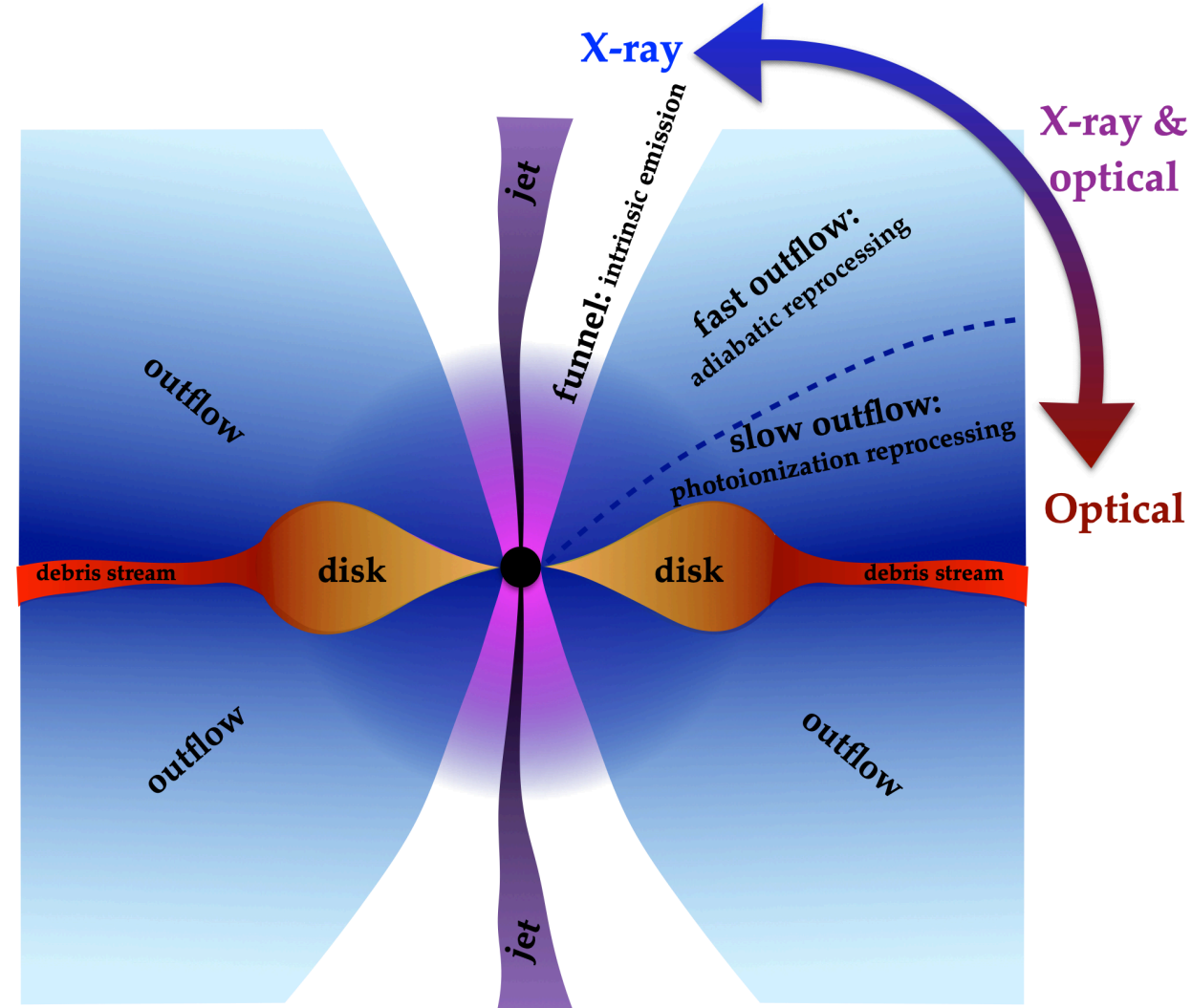
Thermal black body (bb) emissions in optical/UV (OUV) bands

Some (~1/4) TDEs are observed in (thermal) X-ray and infrared (IR)

TDE models

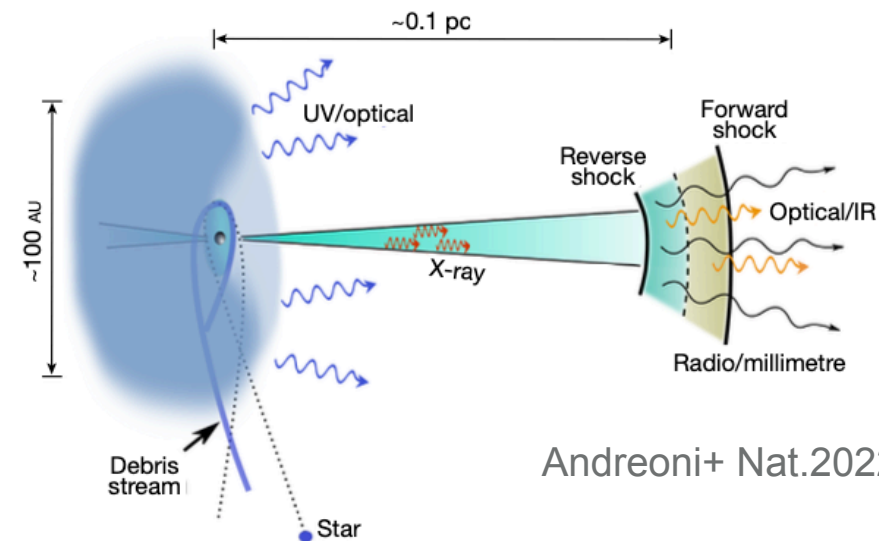
- **Radio, γ -rays, non-thermal X-rays:**
relativistic jet, sub relativistic wind
- **Thermal X-rays:** close to jet/funnel & hot disk corona
- **Optical/UV:** photosphere of hot disk corona (beyond which integrated optical depth < 1)
- **Infrared (IR):** dust-echo, corona...
- **4 TDEs/candidates with luminous X-ray jets among hundreds of TDEs**
 1. AT 2022cmc
 2. Sw 1112-82
 3. Sw1644+57
 4. Sw 2058+05

Disks - Hayashaki & Yamazaki 19 (HY19)
Wide angle winds - Fang 20, Murase+ 20
Stream-stream - Dai + 15,, HY19,
Jets - Wang + 11, Wang & Liu 16, Dai & Fang 17, Lunardini & Winter 17, Senno + 17



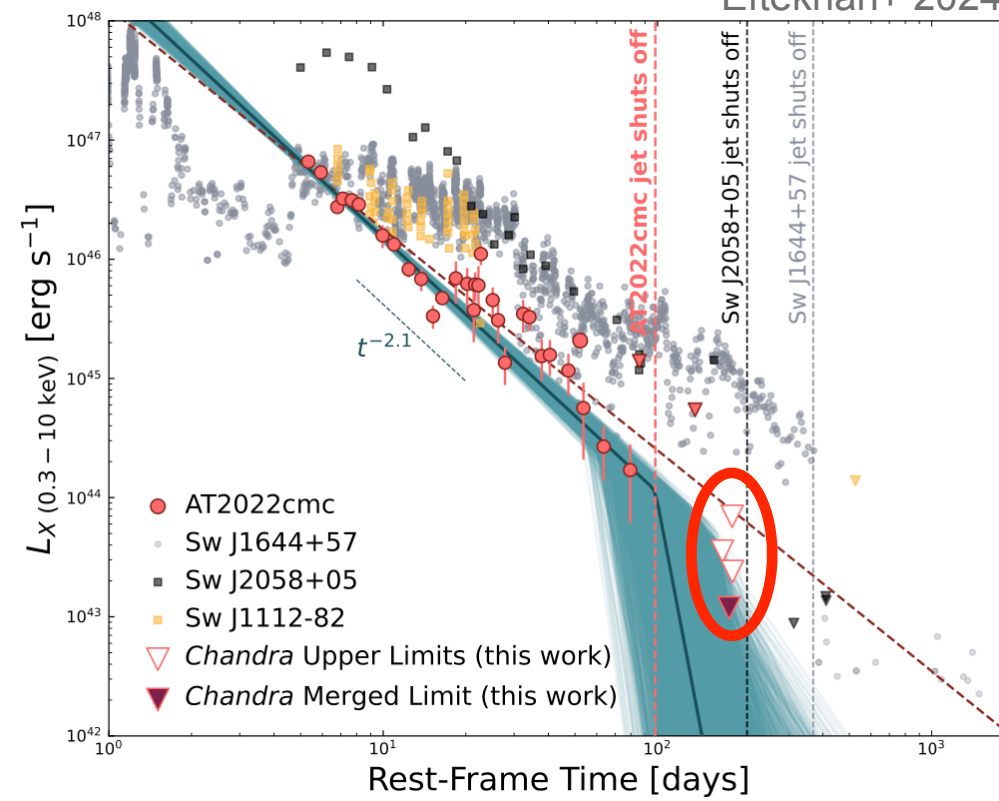
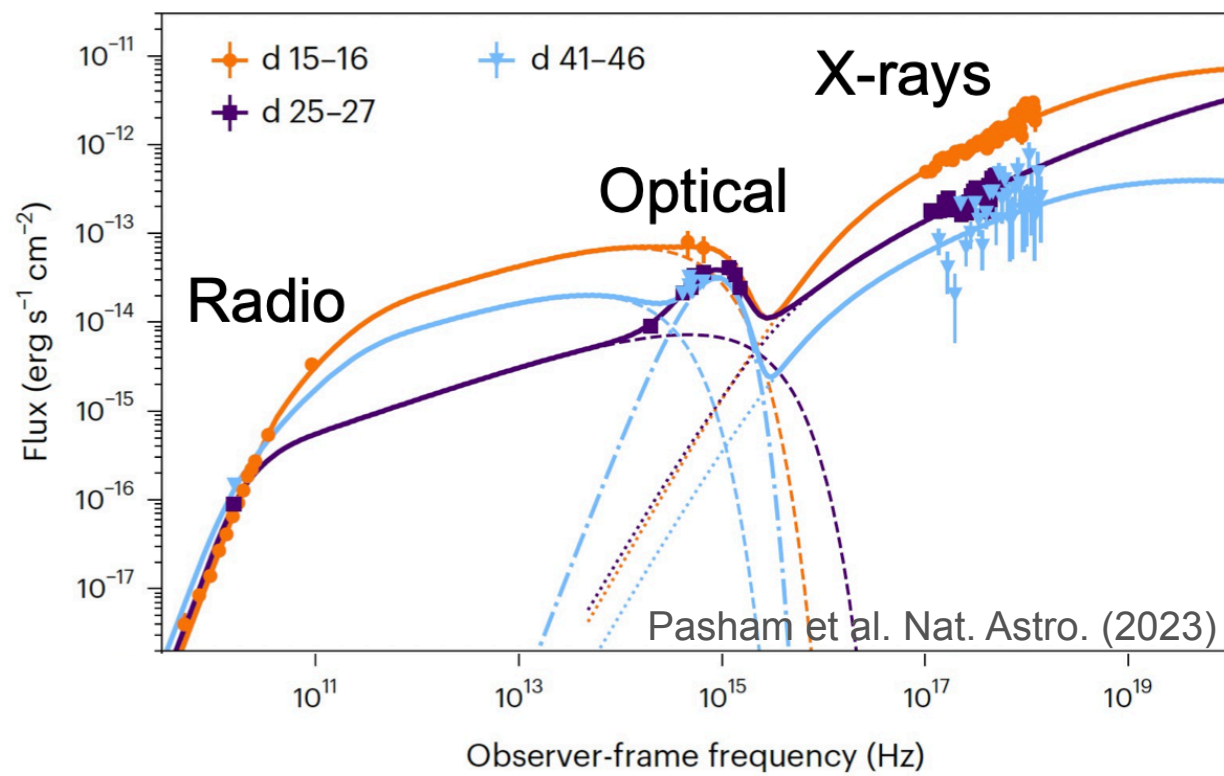
Jetted TDE: AT 2022cmc

- Recently documented jetted TDE ($z = 1.193$, Andreoni+ Nat.2022)
- **Bright non-thermal X-rays:** $L_{X,\text{iso}} \sim 3 \times 10^{47} \text{ erg/s } (T/5 \text{ d})^{-2}$
relativistic jets (Pasham+ 2023) + **later-time steepening** (Eftekhari+ 2024)
- **Optical:** black body emission from thermal envelope (Yao+ 2024)
- **Radio:** GRB-like jet forward shocks ($\Gamma \sim 2 - 5$) propagating in the circumnuclear medium (CNM) $n_{\text{cnm}} \propto R^{-k}$, $1.5 \lesssim k \lesssim 2.0$
(e.g., Matsumoto & Metzger 2023; Yao+ 2024; Zhou+ 2024)



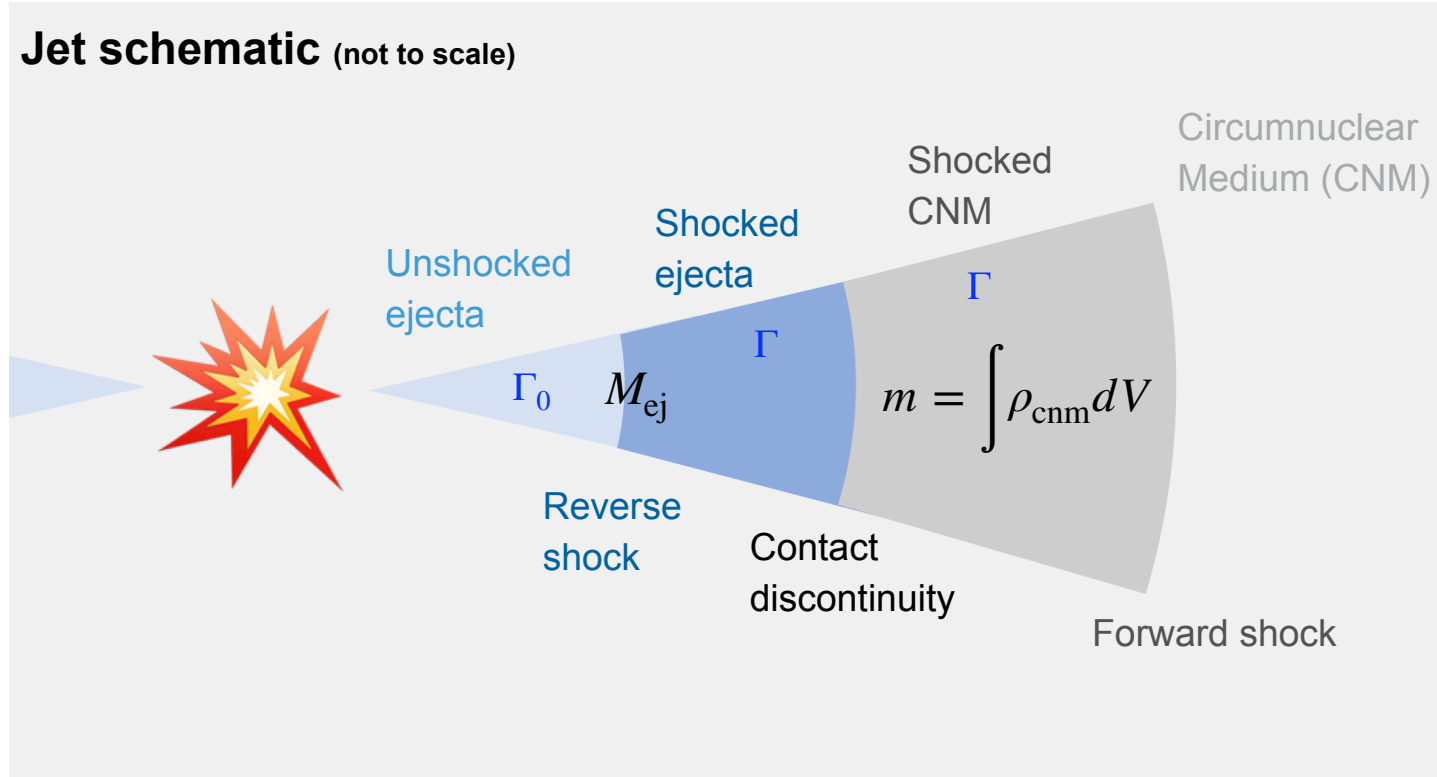
Andreoni+ Nat.2022

Eftekhari+ 2024



Jetted TDE: forward shock model

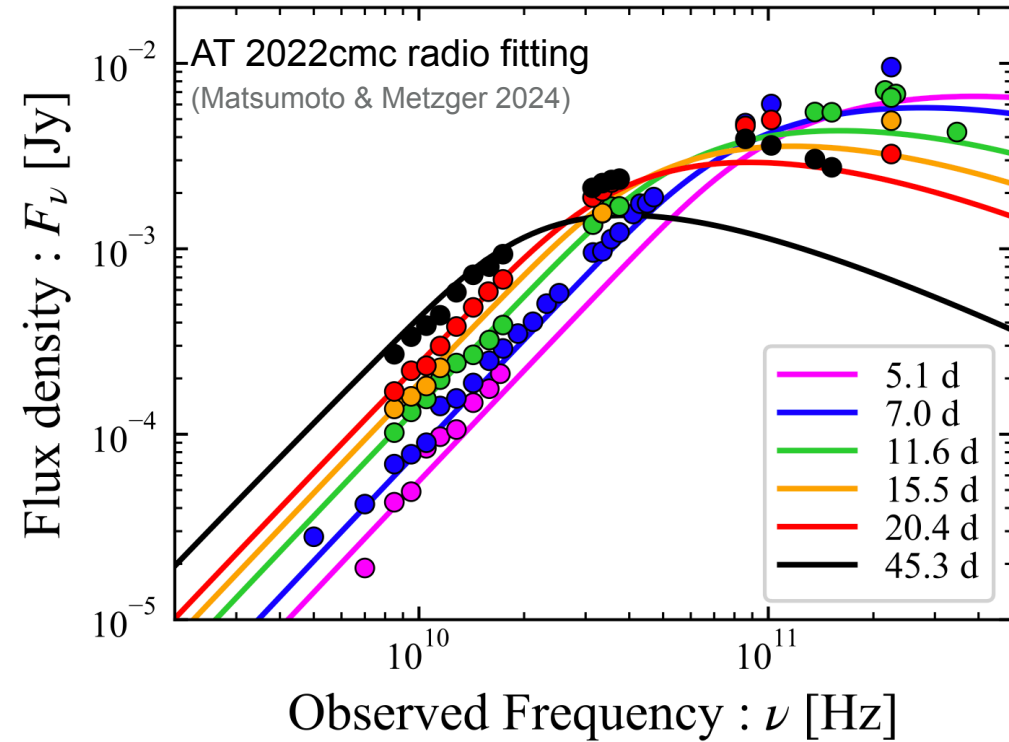
Jet schematic (not to scale)



FS dominated jet evolution with explosion energy $E_0 = \Gamma_0^2 M_{ej} c^2$:

Isotropic equivalent energy of the jet

$$\left. \begin{aligned} E_{iso} &= \Gamma M_{ej}^2 + \Gamma m c^2 + \frac{(\Gamma - 1)(\hat{\gamma} \Gamma^2 - \Gamma + 1)}{\Gamma} m c^2 \\ dE_{iso} &= c^2 dm \text{ (No internal energy injection)} \end{aligned} \right\} \frac{d\Gamma}{dm} = -\frac{\Gamma - 5\Gamma^3 + 4\Gamma^5}{3M_{ej}\Gamma^3 - 2m + 8\Gamma^4 m}.$$

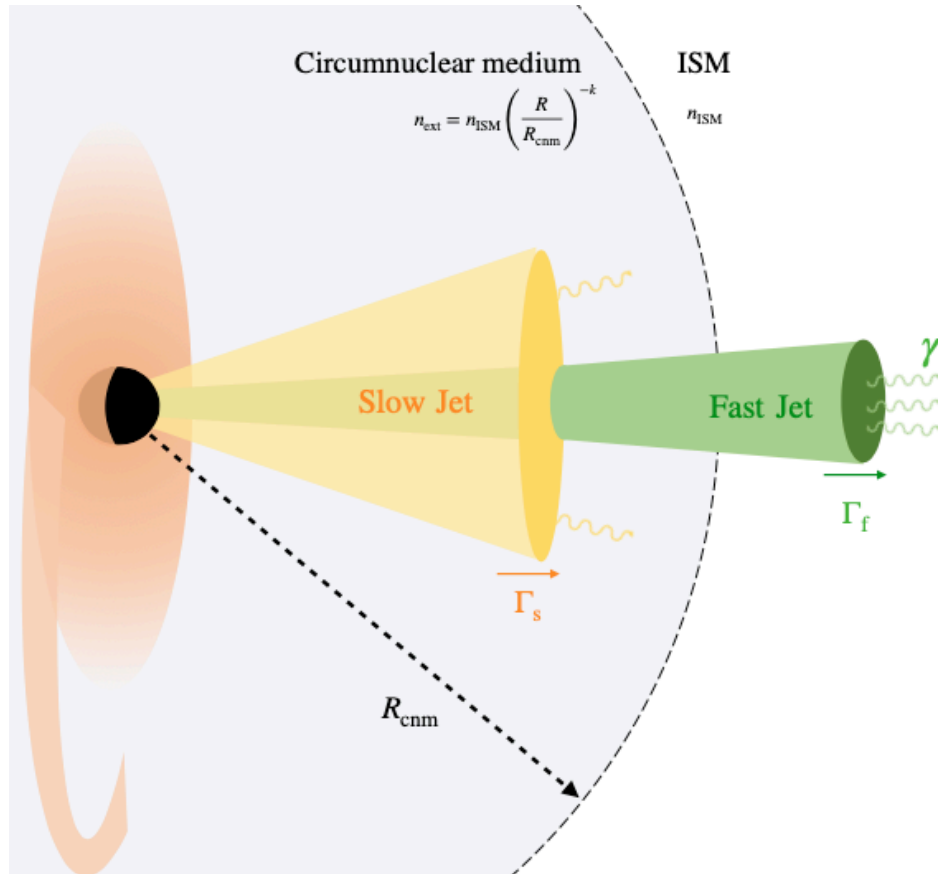


TDE FS model → radio spectra/lightcurves:

- Explosion, no continuous power injection
- $\Gamma_0 \sim 1 - 10$
- Synchrotron emission from accelerated e^-
- Wind-like CNM profile $\rho_{cnm} \propto R^{-k}$ ($1.5 < k < 2$)
- **Cannot explain X-ray observations → two component jet**

Jetted TDE: AT 2022cmc — structured jet

Narrow fast jet and wide slow jet



CY, Zhang, Winter & Murase, arXiv: 2406.11513 (ApJ)

Accretion rate ($\eta_{\text{acc}} \sim 0.01 - 0.1$)

$$\dot{M}_{\text{BH}} = \frac{\eta_{\text{acc}} M_{\star}}{C t_{\text{fb}}} \times \begin{cases} \left(\frac{t}{t_{\text{fb}}}\right)^{-\alpha}, & t < t_{\text{fb}} \\ \left(\frac{t}{t_{\text{fb}}}\right)^{-5/3}, & t > t_{\text{fb}}, \end{cases}$$

Fallback time

$$t_{\text{fb}} \simeq 3.3 \times 10^6 \text{ s } f_{T,-1.2}^{1/2} M_{\text{BH},7}^{1/2} M_{\star,0.7}^{-1/10}$$

Jet luminosities ($\eta_{\text{f/s}} \sim 0.1$)

$$L_{\text{f/s}} = \eta_{\text{f/s}} \dot{M}_{\text{BH}} c^2,$$

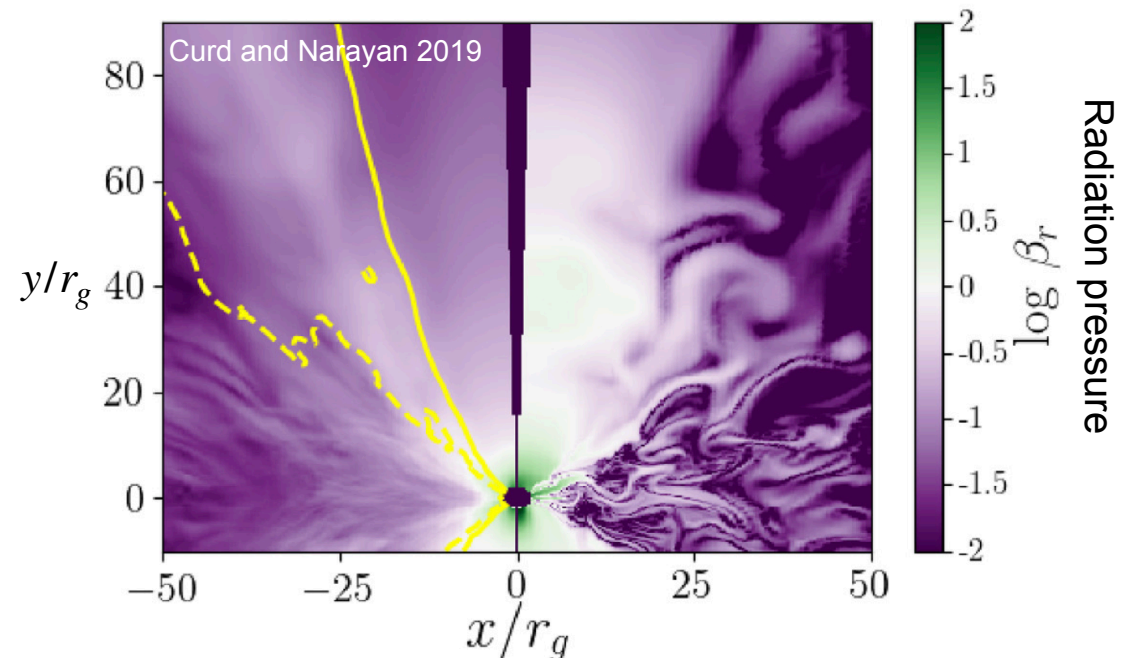
CNM density profile

$$k = 1.8, R_{\text{cnm}} = 10^{18} \text{ cm}$$

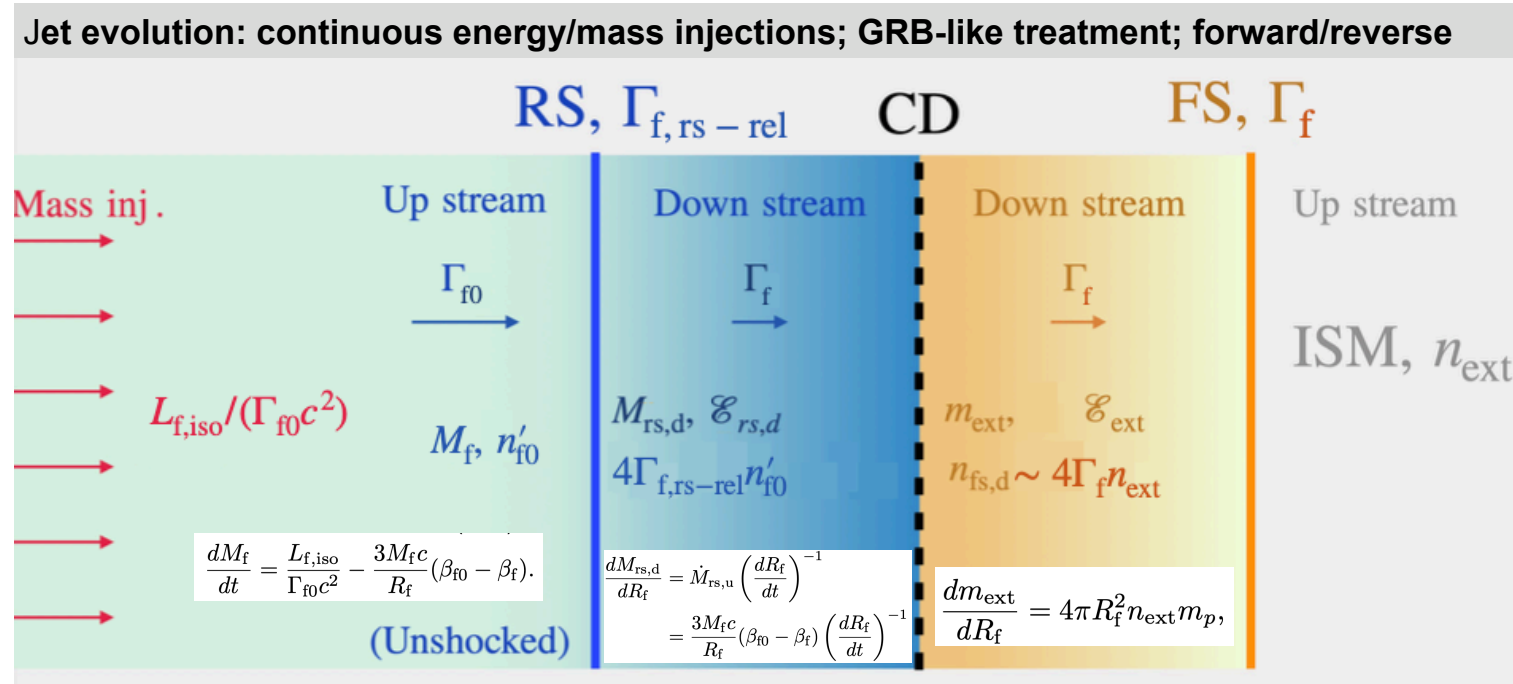
$$n_{\text{ext}}(R) = \begin{cases} n_{\text{ISM}} \left(\frac{R}{R_{\text{cnm}}}\right)^{-k}, & R < R_{\text{cnm}} \\ n_{\text{ISM}}, & R > R_{\text{cnm}} \end{cases}$$

Two components

- Slow outflow (quasi-isotropic)
- Fast relativistic jet (beamed)



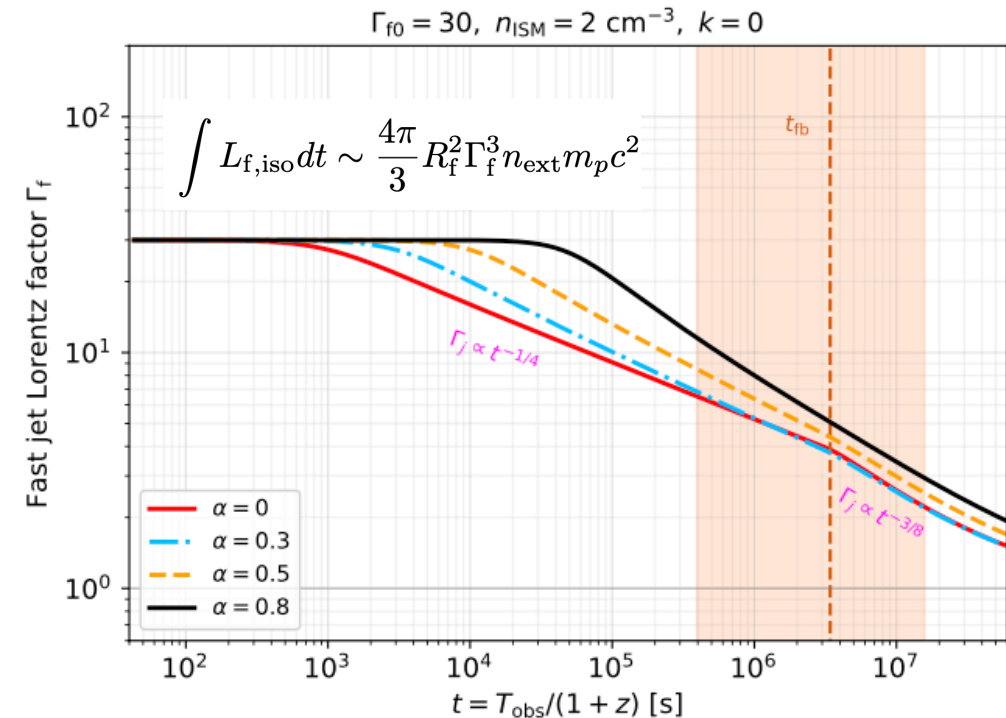
Jetted TDE: AT 2022cmc — jet dynamics



$$\frac{d\Gamma_f}{dR_f} = \frac{\overbrace{(\Gamma_{fs,eff} + 1)(\Gamma_f - 1)c^2 \frac{dm_{ext}}{dR_f}}^{\text{FS term}} + \overbrace{(\Gamma_f - \Gamma_{f0} - \Gamma_{rs,eff} + \Gamma_{rs,eff}\Gamma_{f,rs-rel})c^2 \frac{dM_{rs,d}}{dR_f}}^{\text{RS term}}}{(M_{rs,d} + m_{ext})c^2 + \mathcal{E}_{ext,in} \frac{d\Gamma_{fs,eff}}{d\Gamma_f} + \mathcal{E}_{rs,in} \frac{d\Gamma_{rs,eff}}{d\Gamma_f}}$$

A more comprehensive treatment including:

- Continuous power injection from central engine
- Reverse shock deceleration of unshocked ejecta



CY, Zhang, Winter & Murase, arXiv: 2406.11513 (ApJ)

Numerical method: AM³ (Astrophysical Multi-Messenger Modeling)

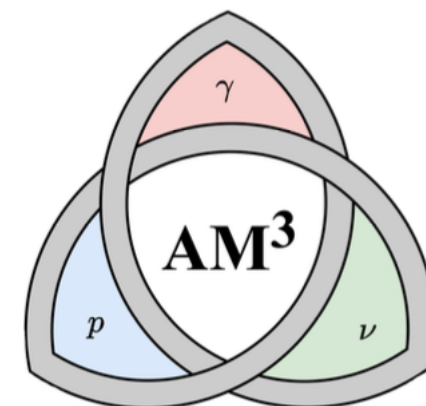
Numerically solving the coupled PDEs for **electron, proton, neutrons, neutrino and photon** distributions.

$$\partial_t n_i = \underbrace{Q_{i,ext}}_{\text{Injection}} + \sum_k Q_{int,k \rightarrow i} - \underbrace{\partial_E(\dot{E} \cdot n_i)}_{\text{Cooling}} - \underbrace{(\alpha_{i,esc} + \alpha_{i,adv})}_{\text{Escape/Sinking}} n_i$$

- An **Open-Source** Tool for **Time-Dependent** Leptoic-Hadronic Radiation Modeling
- Applied to active galactic nuclei, gamma-ray bursts, TDEs

[Klinger, Rudolph, Rodrigues, CY, Clairfontaine, Fedynitch, Winter, Pohl, Gao, arXiv: 2312.13371 \(ApJS\)](#)

- Developed at DESY; Public to the community (<https://gitlab.desy.de/am3/am3>)
- Dedicated for single-zone isotropic multi-messenger (EM and neutrinos) emissions
- Compiled in C++ with a Python interface: fast and easy to use
- Reliable and trackable
- Well-documented with detailed instructions and examples
- More than 14 papers on GRBs, AGNs, and TDEs are published based on AM3. The number keeps expanding!
- Check the doc website to learn more: <https://am3.readthedocs.io>



Jetted TDE: AT 2022cmc — spectra

CY, Zhang, Winter & Murase, arXiv: 2406.11513 (ApJ)

Fast jet reverse shock: X-ray (fast cooling)
Slow jet forward shock: radio (SSA)

Radiation modeling

Powerlaw injection $Q_e \propto \gamma_e^{-s}$

$$\text{Norm.} \cdot (4\pi R_f^2 t'_{f,\text{dyn}}) \int Q_e d\gamma_e = f_e N_e$$

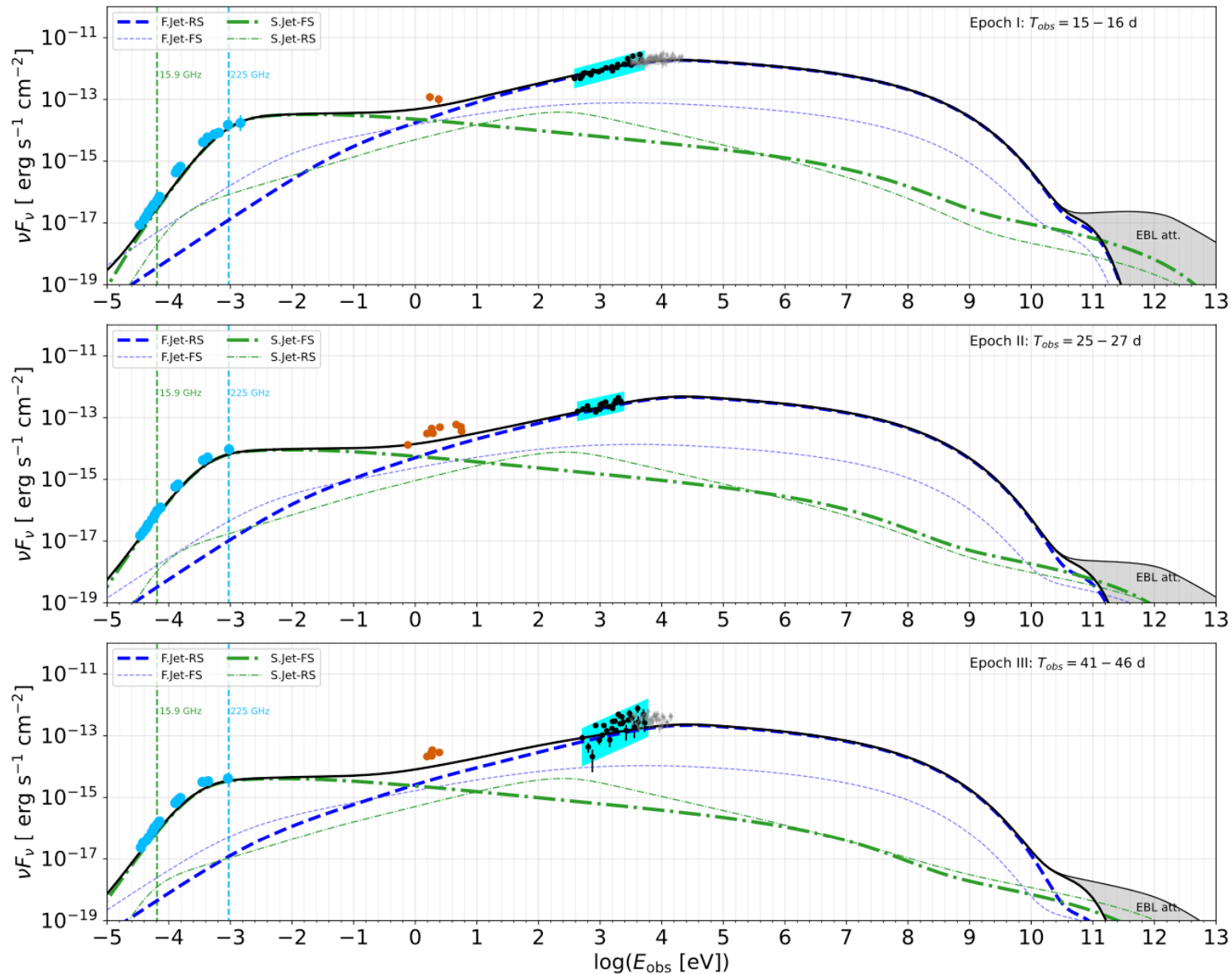
$$\gamma_{e,\text{min}} = (\Gamma - 1) \frac{s - 2}{s - 1} \frac{\epsilon_e m_p}{f_e m_e}$$

$$B_d = \sqrt{32\pi\epsilon_B \Gamma(\Gamma - 1) n_{p,d} m_p c^2}$$

$\Gamma_{f/s}$ for FS, $\Gamma_{\text{rel}} = (\Gamma_{f/s}/\Gamma_{f/s,0} + \Gamma_{f/s,0}/\Gamma_{f/s})/2$ for RS

Fitting parameters

Universal	α	0.8
	n_{ISM}	2.0 cm^{-3}
	s	2.3
Fast, slow jets	$\eta_{f,s}$	0.12, 0.04
	$\theta_{f,s}$	0.15, 0.3
	$\Gamma_{f0,s0}$	30, 4.0
	$\epsilon_e^{\text{fs,rs}}$	0.1, 0.2
FS, RS	$\epsilon_B^{\text{fs,rs}}$	$3.0 \times 10^{-3}, 0.1$
	$f_e^{\text{fs,rs}}$	$1.0, 1.5 \times 10^{-3}$



AT 2022cmc: structured jet light curves

Optical:

- Originated from a thermal envelope (Yao+, 2024)
- ULs for structured jets

X-rays:

- Continuously powered reverse shock model: **works good!**
- Lightcurve steepening after ~ 100 d and the late time ULs after ~ 200 d (**red points**, Eftekhari+ 2024): *Jet break correction*

$$f_{\text{br}} = \frac{1}{1 + (\Gamma_f \theta_f)^{-2}} \rightarrow (\Gamma_f \theta_f)^2, T_{\text{obs}} > T_{\text{br}} (\Gamma_f < \theta_f^{-1})$$

- Analytically consistent,

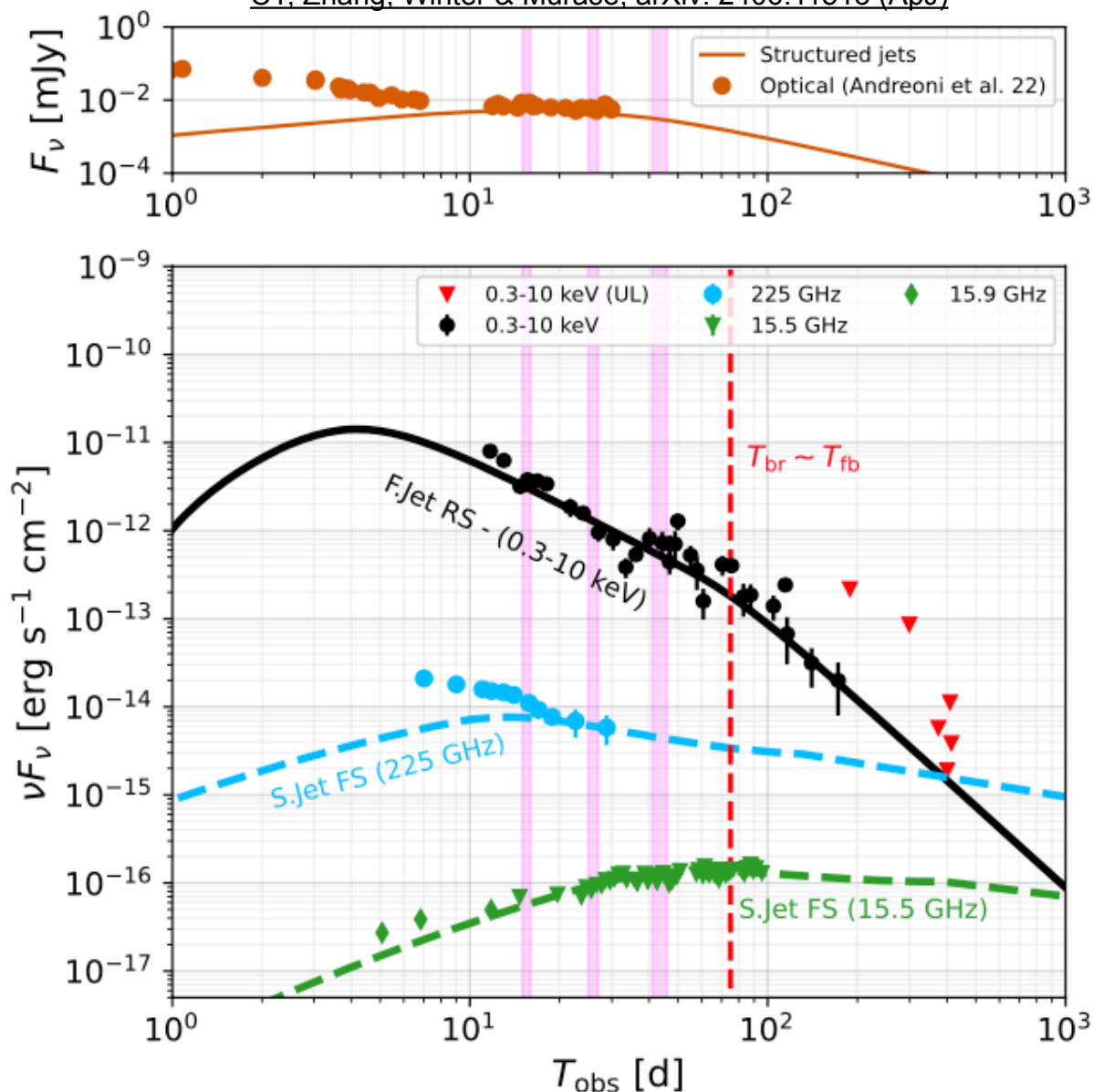
$$\nu F_{\nu}^{(\text{rs})} \propto \begin{cases} T_{\text{obs}}^{-[5\alpha + \alpha(s-1)]/4}, & T_{\text{obs}} < T_{\text{br}} \simeq T_{\text{fb}} \\ T_{\text{obs}}^{-(2s+25)/12}, & T_{\text{obs}} > T_{\text{br}} \simeq T_{\text{fb}}. \end{cases}$$

- Variability timescale: *active engine* ($\sim R_{\text{Sch}}/c$, short term) and *reverse shock* ($\sim R_f/(\Gamma_f^2 c)$, long term)

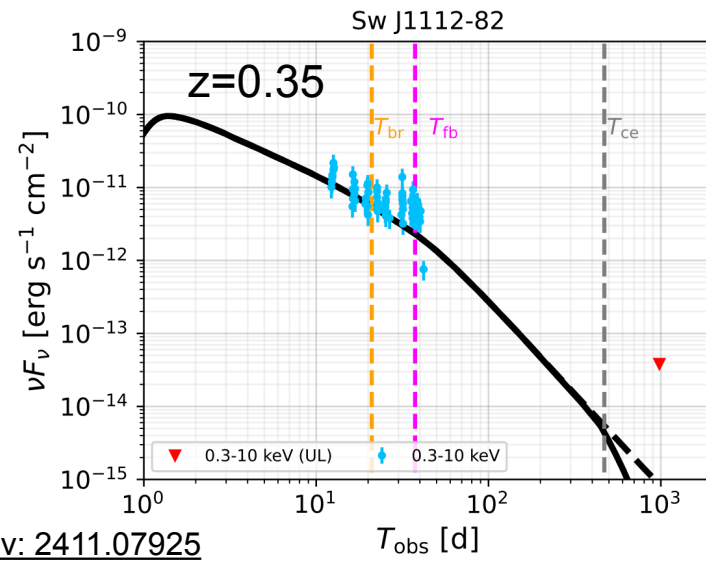
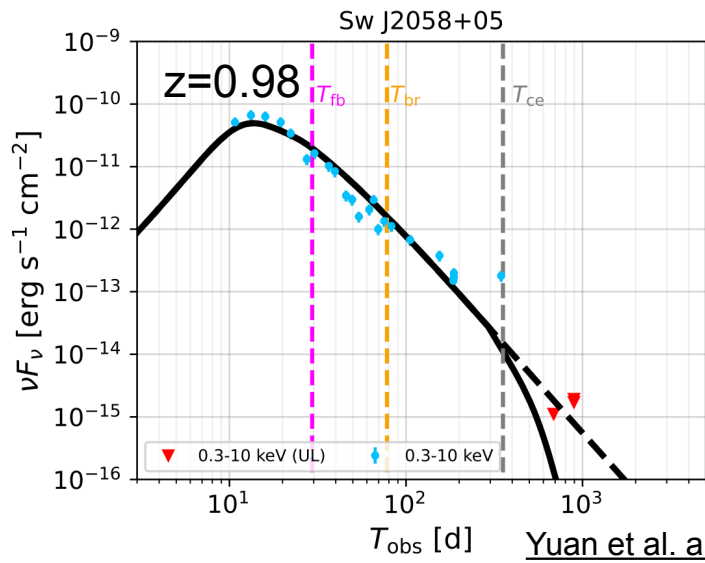
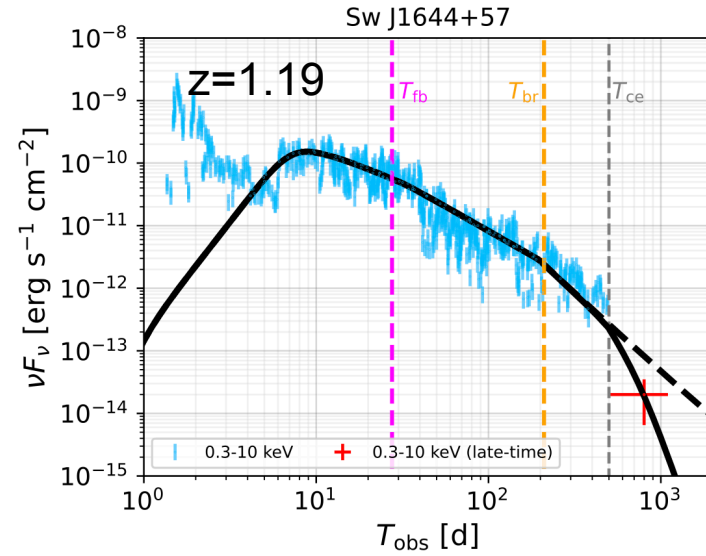
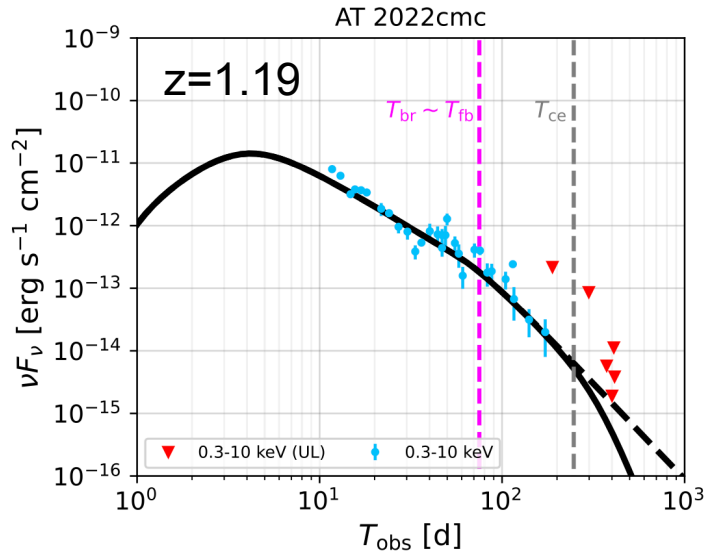
Radio:

Slow jet forward shock: 16 GHz and later-time 225 GHz light curves

CY, Zhang, Winter & Murase, arXiv: 2406.11513 (ApJ)



Revisiting X-ray afterglows of 4 jetted TDEs



Fast jet reverse shock model

Late-time X-ray limits:

- Ceased central engine + jet break: SW J1644 (sharp X-ray break)
- Transition to sub-Eddington rate → jet power injection stops

$$\dot{M}_{\text{BH}}(t_{\text{ce}}) = \frac{L_{\text{Edd}}}{\eta_{\text{rad}} c^2}$$

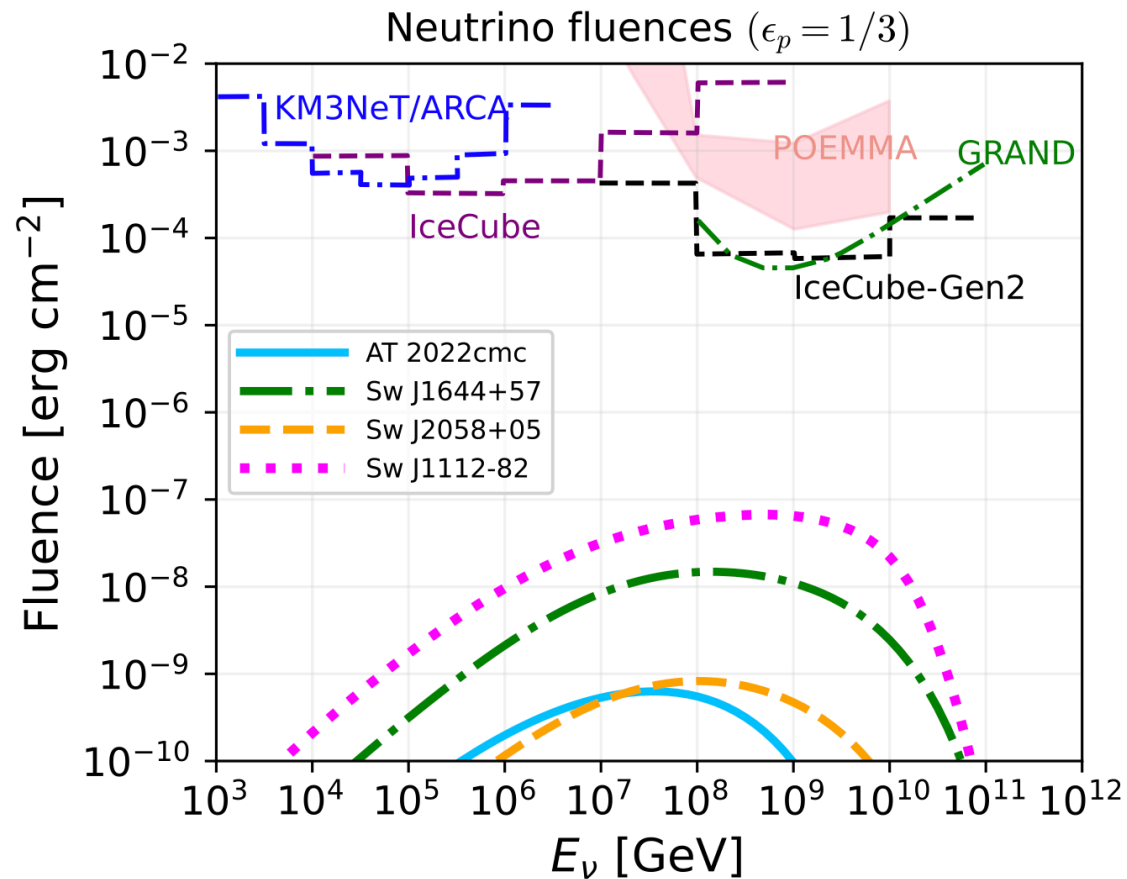
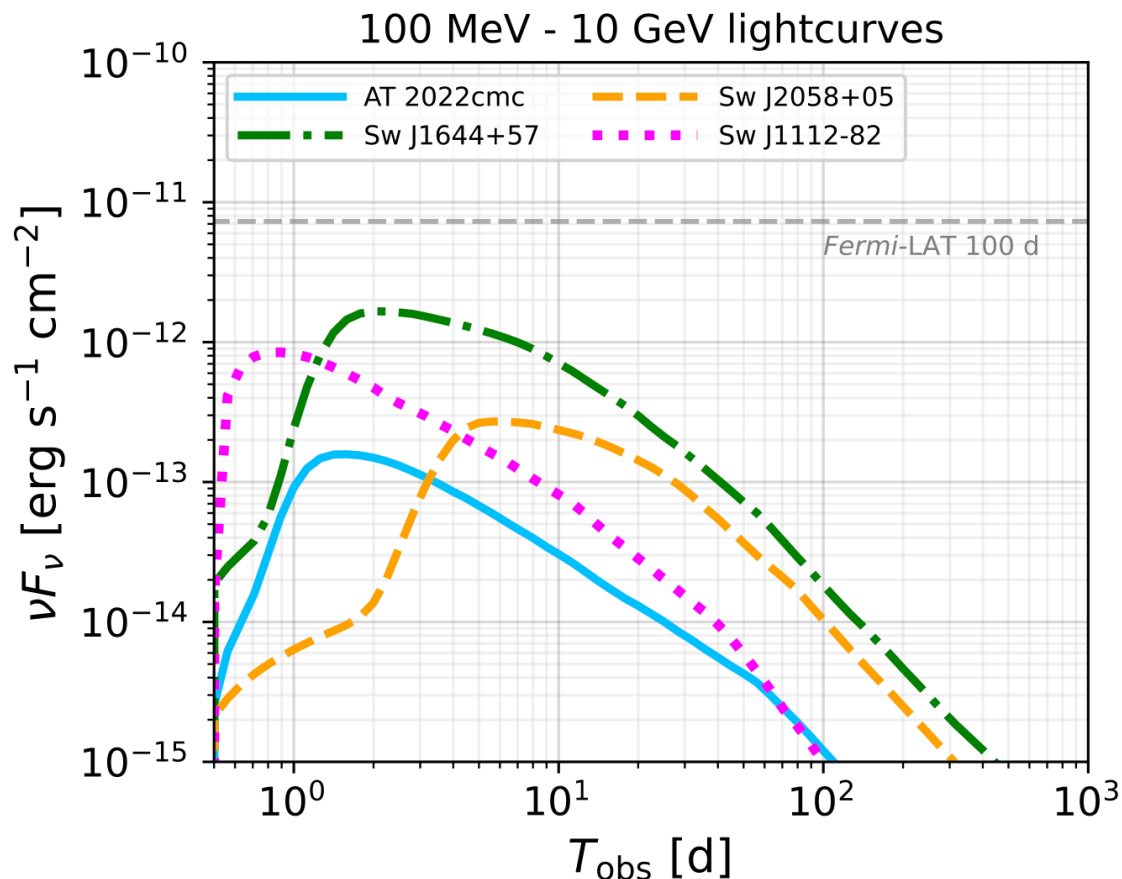
TDEs ^(a)	AT 2022cmc	J1644	J2058	J1112
z	1.19	0.35	1.19	0.89
$M_{\text{BH}} [M_{\odot}]$	10^7	10^6	10^6	2×10^6
Model parameters				
α	0.80	0.65	0.85	0.70
$\mathcal{E}_j [10^{52} \text{ erg}]$	5.4	3.5	2.9	6.3
$n_{\text{ISM}} [\text{cm}^{-3}]$	10	6.0	1.0	10
θ_j	0.15	0.1	0.1	0.1
Γ_0	30	25	42	35
ϵ_B	0.10	0.15	0.20	0.15
Results				
$M_{\star}^{(b)} [M_{\odot}]$	3.0	1.9	1.6	3.5
$T_{\text{fb}} [\text{d}]$	77	28	27	21
$T_{\text{br}} [\text{d}]$	79	212	76	37
$T_{\text{ce}}^{(c)} [\text{d}]$	227	352	331	470

Yuan et al. arXiv: 2411.07925

Jetted TDE X-ray afterglows

Gamma-ray and neutrino detectability

- 4 jetted TDEs: lower than Fermi-LAT 100d; Detection horizon for AT2022cmc-like TDEs: $z = 0.17$, rate ~ 0.02 - 0.1 per year
- Neutrino fluence from jetted TDEs: two orders lower than IceCube-Gen2 sensitivity \leftarrow low target photon density for $p\gamma$



Summary and conclusions

- The continuous power injection effects the early time light curves
- A persistently powered structured (two-component) jet model could explain the **radio (e.g., slow jet, forward shock) and X-ray (fast jet, reverse shock)** spectra/lightcurves of jetted AT 2022cmc
- Joint spectral and lightcurve fitting could reduce the parameter degeneracy
- The jet break and central engine cessation may lead to the late-time steepening in X-ray lightcurves in all 4 jetted TDEs
- TDE relativistic jets are challenging to be detected by current gamma-ray and neutrino detectors.
- Additional contributions, such as hidden winds and external photon fields, are needed to be efficient gamma-ray/neutrino emitters. *The neutrino and electromagnetic cascade signals from neutrino-emitting TDEs will be presented in the poster session.*

Thanks for your attention!

Tidal disruption events

When a massive star passes close enough to a SMBH

- ~ half of the star's mass remains bounded by the SMBH gravitational force
 - Mass accretion -> relativistic jet -> months/year-long flare (optical transient)
 - Energy to be reprocessed by accretion $\sim 10^{54}$ erg
- Fallback rate $\propto t^{-5/3}$ (Phinney 1989)
 - Thermal black body (bb) emissions in **optical/UV (OUV) bands**.
 - Some (~1/4) TDEs are observed in (thermal) **X-ray and infrared (IR)** ranges

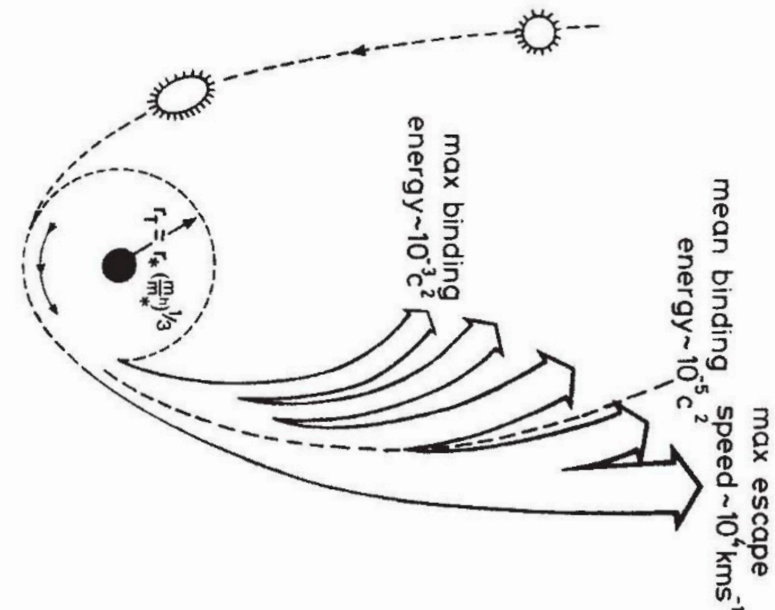
Tidal disruption of stars by black holes of 10^6 – 10^8 solar masses in nearby galaxies

Martin J. Rees

Institute of Astronomy, Madingley Road, Cambridge CB3 0HA, UK

Stars in galactic nuclei can be captured or tidally disrupted by a central black hole. Some debris would be ejected at high speed; the remainder would be swallowed by the hole, causing a bright flare lasting at most a few years. Such phenomena are compatible with the presence of 10^6 – $10^8 M_{\odot}$ holes in the nuclei of many nearby galaxies. Stellar disruption may have interesting consequences in our own Galactic Centre if a $\sim 10^6 M_{\odot}$ hole lurks there.

Martin J. Rees, Nature 1988



Numerical Method: AM³ (Astrophysical Multi-Messenger Modeling)

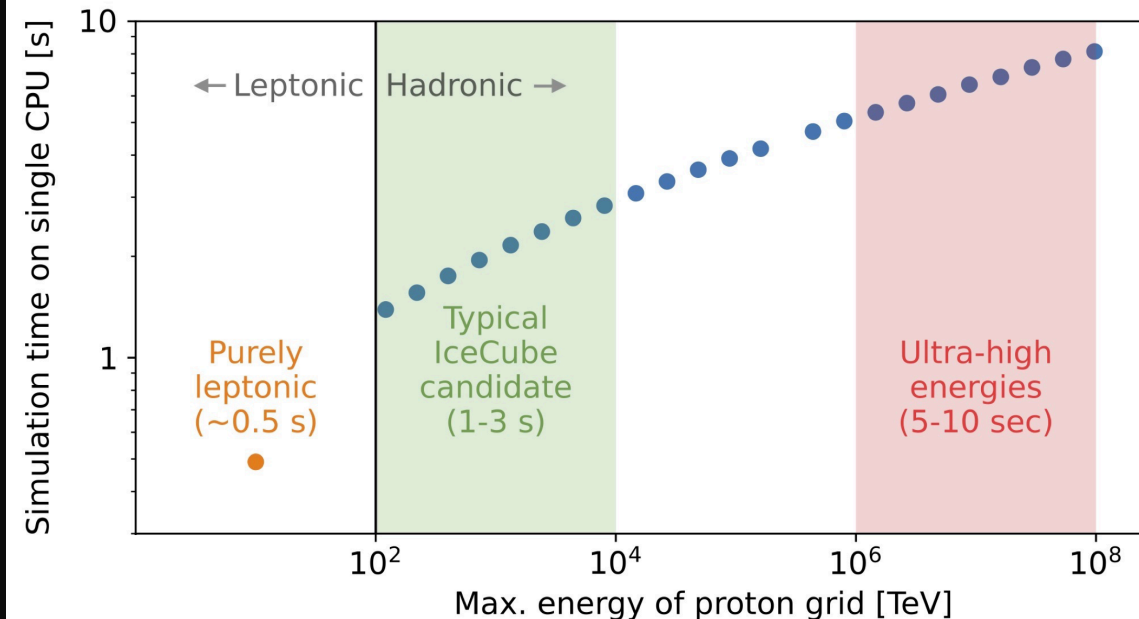
Numerically solving the coupled PDEs for **electron, proton, neutrons, neutrino and photon** distributions.

$$\partial_t n_i = \underbrace{Q_{i,ext}}_{\text{Injection}} + \sum_k Q_{int,k \rightarrow i} - \underbrace{\partial_E(\dot{E} \cdot n_i)}_{\text{Cooling}} - \underbrace{(\alpha_{i,esc} + \alpha_{i,adv})}_{\text{Escape/Sink}} n_i$$

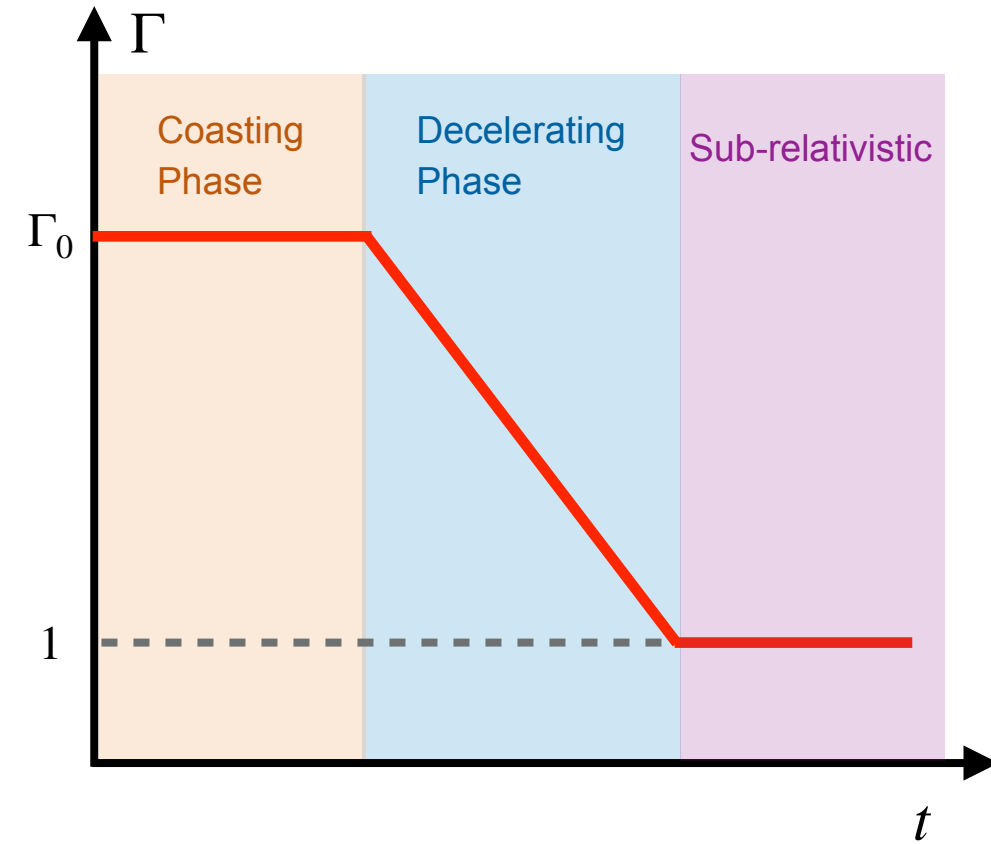
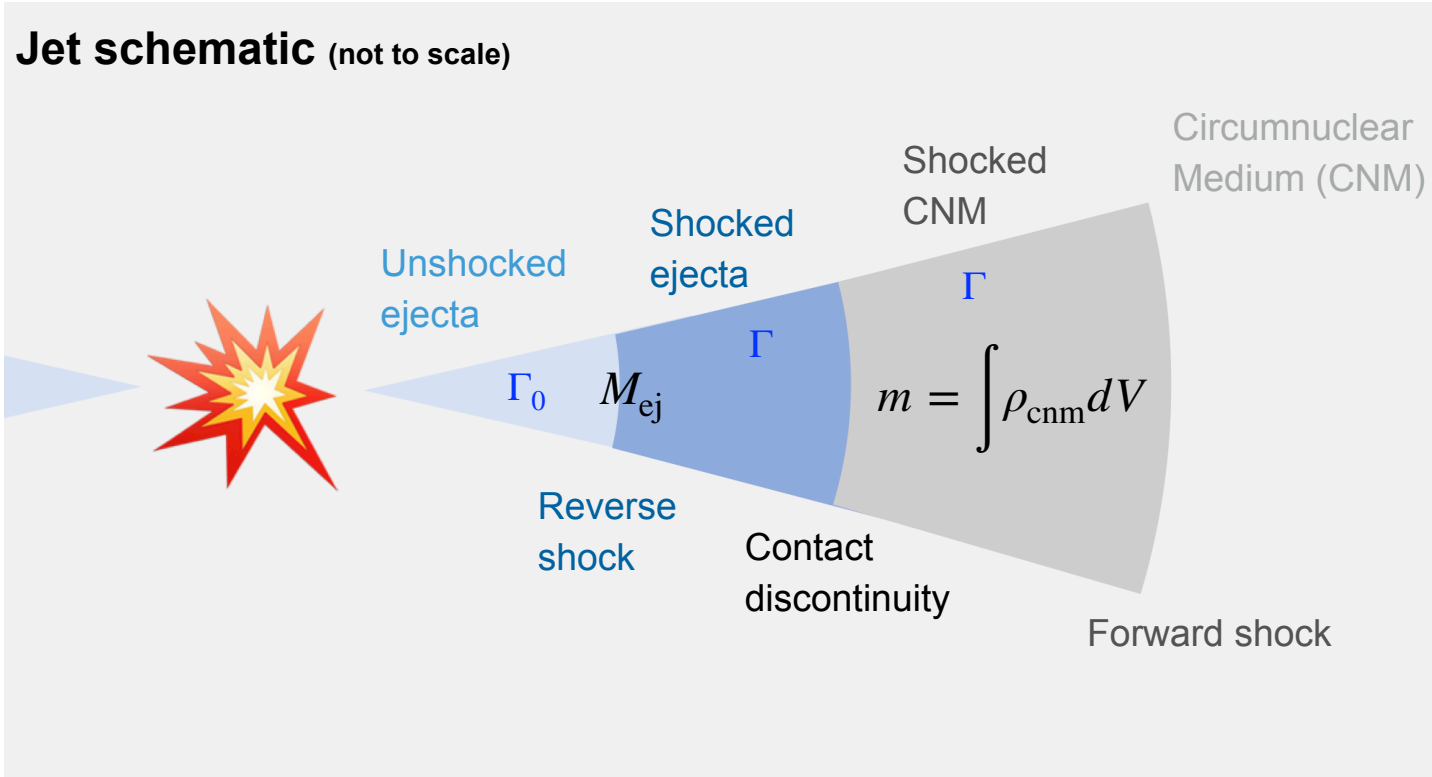
- An **Open-Source** Tool for **Time-Dependent** Lepto-Hadronic Modeling of Astrophysical Sources
- Blazars, GRBs, TDEs, etc
(Klinger, Rudolph, Rodrigues, CY +, arXiv: 2312.13371, ApJS)

Performance: C++ source code, python interface

- Simulation: kernel initialization, particle injection, ~30 steps to steady state
- Tested on a single CPU on Apple M2 chip
- Very fast for lepto-hadronic simulations (< 0.5 s/step)



Jetted TDE: forward shock model



FS dominated jet evolution with explosion energy $E_0 = \Gamma_0^2 M_{ej} c^2$:

Isotropic equivalent energy of the jet

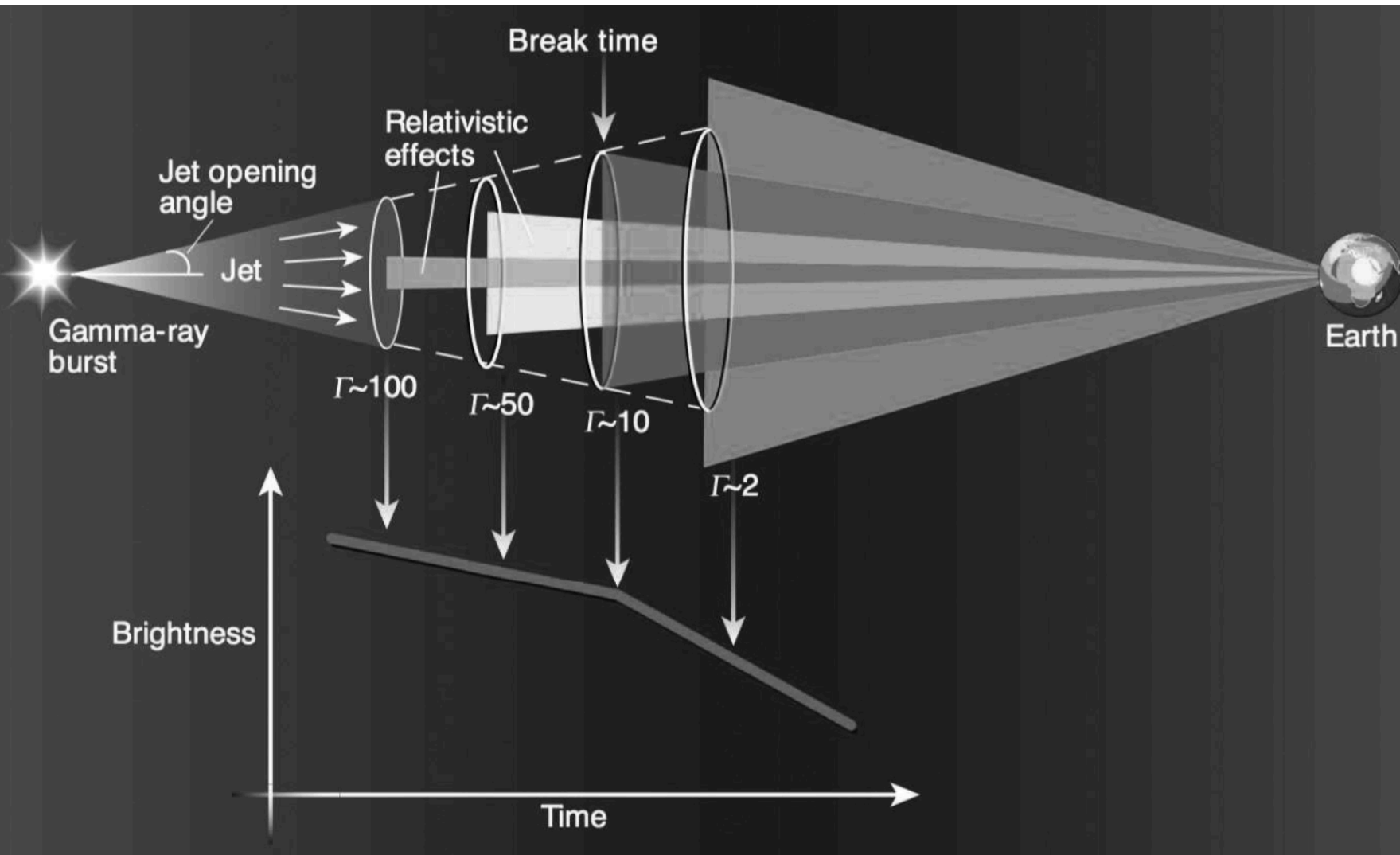
$$E_{iso} = \Gamma M_{ej}^2 + \Gamma m c^2 + \frac{(\Gamma - 1)(\hat{\gamma} \Gamma^2 - \Gamma + 1)}{\Gamma} m c^2$$

$$dE_{iso} = c^2 dm \text{ (No internal energy injection)}$$

$$\left. \vphantom{\begin{matrix} E_{iso} \\ dE_{iso} \end{matrix}} \right\} \frac{d\Gamma}{dm} = -\frac{\Gamma - 5\Gamma^3 + 4\Gamma^5}{3M_{ej}\Gamma^3 - 2m + 8\Gamma^4 m}$$

Jet break

Geometric effect of decelerating jet

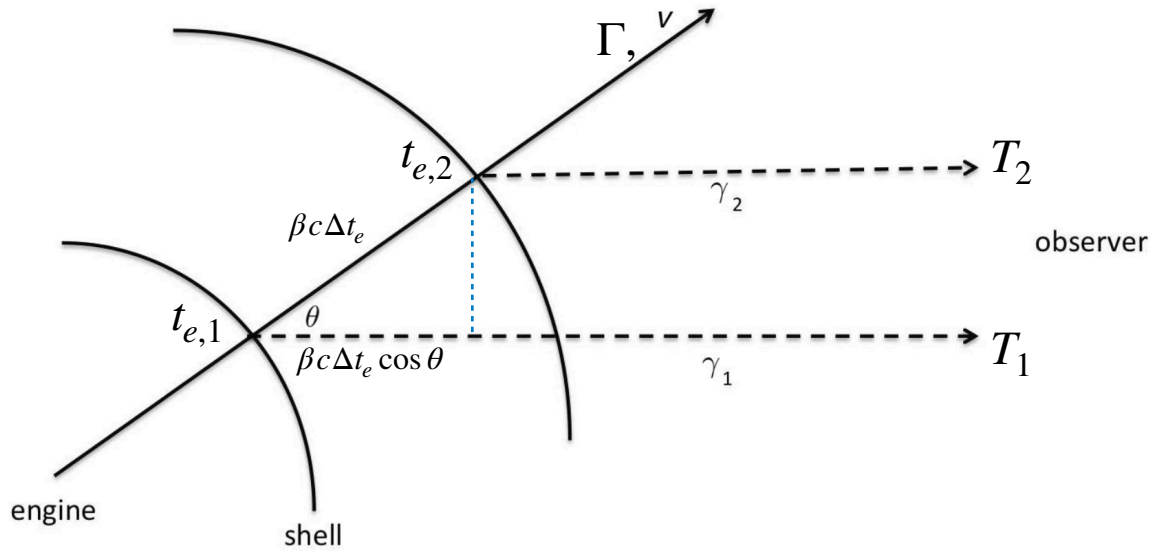


For jets with opening angle θ_j

- $\theta_j > 1/\Gamma$: observer has no knowledge outside Γ^{-1} cone
- $\theta_j < 1/\Gamma$: observer feels the progressive deficit of energy with in Γ^{-1} since no emission outside the jet cone is available
- $\theta_j = 1/\Gamma$ defines break time t_{br} , correction factor $(\theta_j \Gamma)^2$ applies after t_{br}

Equal-arrival-time surface (EATS)

Relate the time interval ΔT in the observer's frame to the emission time interval in the jet Δt_e (engine frame)



$$T_1 = t_{e,1} + d/c$$

$$T_2 = t_{e,2} + d/c - \beta \Delta t_e \cos \theta$$

$$\Delta T = T_2 - T_1 = \Delta t_e (1 - \beta \cos \theta)$$

Set $t_{e,1} = 0$, radius $R = \int \beta c dt_e \sim \frac{\beta c T}{1 - \beta \cos \theta}$

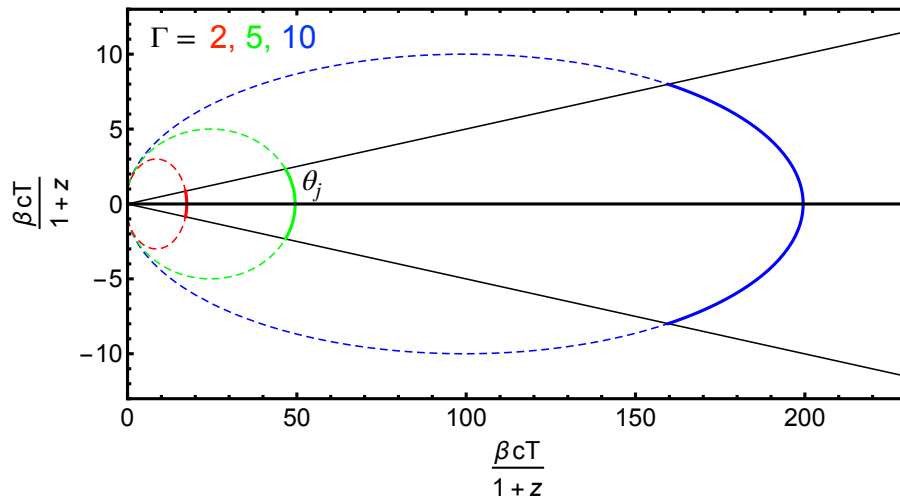
For $\theta = 0$, $T = (1 + z)t_e(1 - \beta) \simeq (1 + z)t_e/(2\Gamma^2)$

Photons from the EATS arrives simultaneously

$$R = \frac{\beta c T}{(1 - \beta \cos \theta)(1 + z)}$$

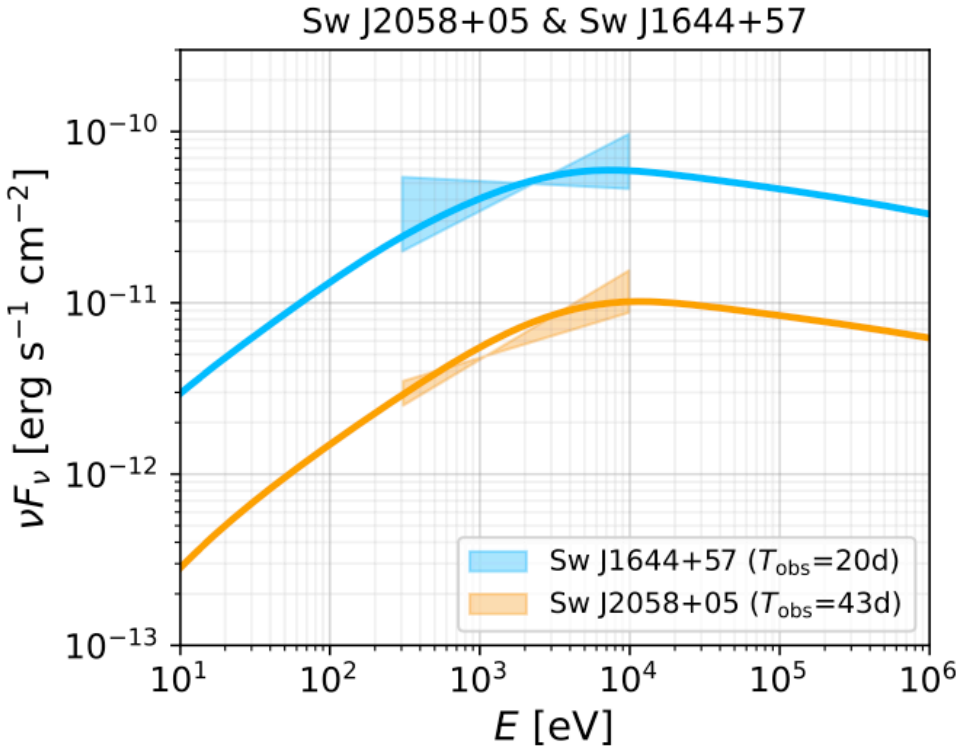
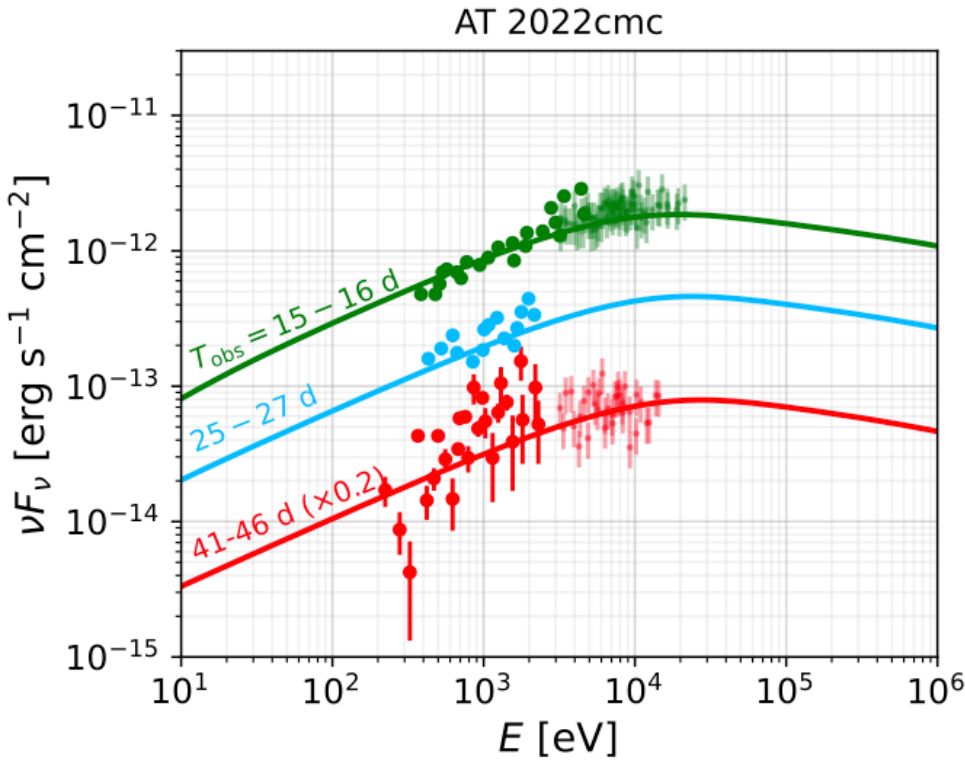
Integrate the emissivity over the EATS:

$$F(E, T_{\text{obs}}) = \frac{2\pi(1+z)}{d_L^2} \int_0^{\theta_j} d\theta \sin(\theta) R_j^2 \times \Delta R_j f_{\text{br}} \frac{j'(\hat{t}, R_j, \varepsilon')}{\Gamma^2(1 - \beta\mu)^2} \Big|_{\hat{t} = \frac{T_{\text{obs}}}{1+z} + \frac{\mu R_j}{c}}$$



Revisiting X-ray afterglows of 4 jetted TDEs

Spectral fitting



Yuan et al. arXiv: 2411.07925