LARGE NEUTRINO MASSES CONSISTENT WITH COSMOLOGY & keV STERILE NEUTRINO DARK MATTER **FROM A DARK SECTOR**

Cristina Benso, KIT



obel Kolbé Anna McCoy Farid Salazar, Yukari Yan



Karlsruher Institut für Technologie







• Standard Model is great







• Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature:





4

- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature:
 - active neutrino masses

* KATRIN Collaboration, [2406.13516 [nucl-ex]]

** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu







- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature:
 - active neutrino masses
 - dark matter

* KATRIN Collaboration, [2406.13516 [nucl-ex]]

** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu







- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature:
 - active neutrino masses ---- seesaw mechanism (3 heavy RH Majorana neutrinos N)
 - dark matter

* KATRIN Collaboration, [2406.13516 [nucl-ex]]

** I. Esteban et al, [2410.05380 [hep-ph]]

<u>cristina.benso@kit.edu</u>



east) two puzzles of Nature: B heavy RH Majorana neutrinos N)





- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature:
 - active neutrino masses ---- seesaw mechanism (3 heavy RH Majorana neutrinos N)
 - dark matter \longrightarrow sterile neutrino DM (ψ)

* KATRIN Collaboration, [2406.13516 [nucl-ex]]

** I. Esteban et al, [2410.05380 [hep-ph]]

<u>cristina.benso@kit.edu</u>



east) two puzzles of Nature: B heavy RH Majorana neutrinos N)





- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses; - KATRIN: current upper limit $m_{\nu_o} < 0.45$ eV *, expected final reach $m_{\nu_o} = 0.2$ eV;

 - Oscillation data: $\Sigma m_{\nu} > 0.058 \, (0.098)$ eV, for normal (inverted) neutrino mass ordering. **

* KATRIN Collaboration, [2406.13516 [nucl-ex]]

** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu



****DESI Collaboration, [2404.03002 [astro-ph.CO]]





- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses; - KATRIN: current upper limit $m_{\nu_a} < 0.45$ eV *, expected final reach $m_{\nu_a} = 0.2$ eV; - Oscillation data: $\Sigma m_{\nu} > 0.058 (0.098)$ eV, for normal (inverted) neutrino mass ordering. **

- Cosmological observations (assuming ΛCDM): stringent constraints on Σm_{ν} ; - DESI: $\Sigma m_{\nu} < 0.072 \text{ eV }^{***}$

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu



****DESI Collaboration, [2404.03002 [astro-ph.CO]]



- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses; - KATRIN: current upper limit $m_{\nu_o} < 0.45$ eV *, expected final reach $m_{\nu_o} = 0.2$ eV; - Oscillation data: $\Sigma m_{\nu} > 0.058 (0.098)$ eV, for normal (inverted) neutrino mass ordering. **
- Cosmological observations (assuming ΛCDM): stringent constraints on Σm_{μ} ; - DESI: $\Sigma m_{\nu} < 0.072 \text{ eV }^{***}$

• What if KATRIN measures something? How could laboratory results be reconciled with cosmological limits?

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu

MU Days 2024 - 12.12.2024, DESY Hamburg



****DESI Collaboration, [2404.03002 [astro-ph.CO]]



- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses; - KATRIN: current upper limit $m_{\nu_a} < 0.45$ eV *, expected final reach $m_{\nu_a} = 0.2$ eV; - Oscillation data: $\Sigma m_{\nu} > 0.058 (0.098)$ eV, for normal (inverted) neutrino mass ordering. **
- Cosmological observations (assuming Λ CDM): stringent constraints on Σm_{i} ; - DESI: $\Sigma m_{\nu} < 0.072 \text{ eV} ***$

- What if KATRIN measures something? How could laboratory results be reconciled with cosmological limits?
- Is it possible that the same dark sector that could make laboratory measurement compatible with cosmological limits can also provide a viable dark matter candidate?

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu



****DESI Collaboration, [2404.03002 [astro-ph.CO]]

MU Days 2024 - 12.12.2024, DESY Hamburg



Cosmological bounds on neutrino masses are established constraining ρ_v



Cosmological bounds on neutrino masses are establish

*DESI Collaboration, [2404.03002 [astro-ph.CO]]

cristina.benso@kit.edu



ned constraining
$$\rho_v \colon \Sigma m_v \times \left(\frac{n_v^0}{56 \text{ cm}^{-3}}\right) < 0.072 \text{ eV}^*$$

Cosmological bounds on neutrino masses are establish

 \longrightarrow if n_{ν}^0 changes, the upper bound on Σm_{ν} can be relaxed.

*DESI Collaboration, [2404.03002 [astro-ph.CO]]



ned constraining
$$\rho_v \colon \Sigma m_v \times \left(\frac{n_v^0}{56 \text{ cm}^{-3}}\right) < 0.072 \text{ eV}^*$$

Cosmological bounds on neutrino masses are establish

 \longrightarrow if n_{ν}^0 changes, the upper bound on Σm_{ν} can be relaxed.



*DESI Collaboration, [2404.03002 [astro-ph.CO]]



ned constraining
$$\rho_v \colon \Sigma m_v \times \left(\frac{n_v^0}{56 \text{ cm}^{-3}}\right) < 0.072 \text{ eV}^*$$

ontributes to
$$N_{eff} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \left(\frac{\rho_{rad} - \rho_{\gamma}}{\rho_{\gamma}}\right)$$

RELAXATION OF COSMOLOGICAL BOUND ON NEUTRINO MASSES

Cosmological bounds on neutrino masses are establish

 \longrightarrow if n_{ν}^0 changes, the upper bound on Σm_{ν} can be relaxed.



 \rightarrow the depletion of ν must be compensated by production of new light or massless dark species χ . Our hypothesis: ν are transformed into χ (fermionc singlets).

*DESI Collaboration, [2404.03002 [astro-ph.CO]]

cristina.benso@kit.edu

ned constraining
$$\rho_{v} \colon \Sigma m_{\nu} \times \left(\frac{n_{\nu}^{0}}{56 \text{ cm}^{-3}}\right) < 0.072 \text{ eV} *$$

ontributes to
$$N_{eff} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \left(\frac{\rho_{rad} - \rho_{\gamma}}{\rho_{\gamma}}\right)$$





RELAXATION OF COSMOLOGICAL BOUND ON NEUTRINO MASSES

Cosmological bounds on neutrino masses are establish

 \longrightarrow if n_{ν}^0 changes, the upper bound on Σm_{ν} can be relaxed.



 \rightarrow the depletion of ν must be compensated by production of new light or massless dark species χ . Our hypothesis: ν are transformed into χ (fermionc singlets).



*DESI Collaboration, [2404.03002 [astro-ph.CO]]

cristina.benso@kit.edu

ned constraining
$$\rho_{v} \colon \Sigma m_{\nu} \times \left(\frac{n_{\nu}^{0}}{56 \text{ cm}^{-3}}\right) < 0.072 \text{ eV} *$$

ontributes to
$$N_{eff} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \left(\frac{\rho_{rad} - \rho_{\gamma}}{\rho_{\gamma}}\right)$$





RELAXATION OF COSMOLOGICAL

Cosmological bounds on neutrino masses are established

 \longrightarrow if n_{ν}^0 changes, the upper bound on Σm_{ν} can be relaxed.



 \rightarrow the depletion of ν must be compensated by production of new light or massless dark species χ . Our hypothesis: ν are transformed into χ (fermionc singlets).



*DESI Collaboration, [2404.03002 [astro-ph.CO]]

cristina.benso@kit.edu



ned constraining
$$\rho_v \colon \Sigma m_v \times \left(\frac{n_v^0}{56 \text{ cm}^{-3}}\right) < 0.072 \text{ eV}^*$$

ontributes to
$$N_{eff} = \frac{8}{7} \left(\frac{11}{4}\right)^{4/3} \left(\frac{\rho_{rad} - \rho_{\gamma}}{\rho_{\gamma}}\right)$$

 \rightarrow the transformation of ν into χ must take place in the specific temperature range 100 keV $\gtrsim T \gtrsim 10$ eV.









Wishlist:





Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



MU Days 2024 - 12.12.2024, DESY Hamburg



Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



 $\tilde{\chi}$ N massless vector-like dark fields χ to deplete ν after their decoupling from the thermal bath







Wishlist:



3 heavy RH Majorana neutrinos N to give mass to ν via seesaw mechanism



 $\tilde{\chi}$ N massless vector-like dark fields χ to deplete ν after their decoupling from the thermal bath



 $one DM candidate \psi$







Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



 $\sum_{\chi} N_{\chi}$ massless vector-like dark fields χ to deplete ν after their decoupling from the thermal bath



 \bigvee one DM candidate ψ

They can all be accommodated in a dark sector that is mainly an extension of the SM neutrino sector,







Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



 N_{γ} massless vector-like dark fields χ to deplete u after their decoupling from the thermal bath



 \bigvee one DM candidate ψ

They can all be accommodated in a dark sector that is mainly an extension of the SM neutrino sector, complemented with

- a new gauge boson Z' that mediates the interactions between u, χ and ψ







Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



 N_{γ} massless vector-like dark fields χ to deplete u after their decoupling from the thermal bath



one DM candidate ψ

They can all be accommodated in a dark sector that is mainly an extension of the SM neutrino sector, complemented with

- a new gauge boson Z' that mediates the interactions between ν, χ and ψ
- a new singlet scalar ϕ , whose VEV breaks the new U(1) and gives mass to Z' and to dark neutrinos





Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



 N_{γ} massless vector-like dark fields χ to deplete u after their decoupling from the thermal bath



one DM candidate ψ

They can all be accommodated in a dark sector that is mainly an extension of the SM neutrino sector, complemented with

- a new gauge boson Z' that mediates the interactions between ν, χ and ψ
- a new singlet scalar ϕ , whose VEV breaks the new U(1) and gives mass to Z' and to dark neutrinos

Bonus:





Wishlist:



3 heavy RH Majorana neutrinos N to give mass to u via seesaw mechanism



 N_{γ} massless vector-like dark fields χ to deplete u after their decoupling from the thermal bath



 $one DM candidate \psi$

They can all be accommodated in a dark sector that is mainly an extension of the SM neutrino sector, complemented with

- a new gauge boson Z' that mediates the interactions between ν, χ and ψ
- a new singlet scalar ϕ , whose VEV breaks the new U(1) and gives mass to Z' and to dark neutrinos

Bonus:



one lighter copy N' of the heavy RH Majorana neutrinos N, participating in a second seesaw mechanism to give mass to ψ







New symmetries:







New symmetries:

• U(1)' gauge symmetry





New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms



MU Days 2024 - 12.12.2024, DESY Hamburg



New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms

New interactions:



MU Days 2024 - 12.12.2024, DESY Hamburg





New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms

New interactions:

• Yukawa interactions

$$-\mathscr{L}_{int} = Y_{\nu}\bar{N}l_{L}\tilde{H}^{\dagger} + Y_{\chi}\bar{N}\chi_{L}\phi + Y_{\psi}\bar{N}\psi_{L}\phi + Y_{\nu}\bar{N}'l_{L}\tilde{H}$$



$\tilde{H}^{\dagger} + Y'_{\chi}\bar{N}'\chi_L\phi + Y'_{\psi}\bar{N}'\psi_L\phi + \frac{1}{2}M\bar{N}N^c + \frac{1}{2}M'\bar{N}'N'^c + H.c.$



New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms

New interactions:

• Yukawa interactions

Interactions of N

$$-\mathscr{L}_{int} = Y_{\nu}\bar{N}l_{L}\tilde{H}^{\dagger} + Y_{\chi}\bar{N}\chi_{L}\phi + Y_{\psi}\bar{N}\psi_{L}\phi + Y_{\nu}\bar{N}'l_{L}\tilde{H}$$



Interactions of N'

$\tilde{H}^{\dagger} + Y'_{\chi}\bar{N}'_{\chi_L}\phi + Y'_{\psi}\bar{N}'\psi_L\phi + \frac{1}{2}M\bar{N}N^c + \frac{1}{2}M'\bar{N}'N'^c + H.c.$





New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms

New interactions:

• Yukawa interactions

Interactions of N

$$-\mathscr{L}_{int} = Y_{\nu}\bar{N}l_{L}\tilde{H}^{\dagger} + Y_{\chi}\bar{N}\chi_{L}\phi + Y_{\psi}\bar{N}\psi_{L}\phi + Y_{\nu}\bar{N}'l_{L}\tilde{H}$$









New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms

New interactions:

- Yukawa interactions Interactions of N
- Gauge interactions:

$$\mathscr{L} = \sum_{f} Q_{f} g Z'_{\mu} \bar{f} \gamma^{\mu} f \qquad \text{with} \quad f = \{\chi_{L}, \chi_{R}, \psi_{L}, \psi_{L}$$





 Ψ_R



New symmetries:

- U(1)' gauge symmetry
- \mathbb{Z}_2 symmetry, under which all fields but ψ_R and χ_R are even \longrightarrow forbids vector-like mass terms

New interactions:

• Yukawa interactions
Interactions of N
$$-\mathscr{L}_{int} = Y_{\nu}\bar{N}l_{L}\tilde{H}^{\dagger} + Y_{\nu}\bar{N}\chi_{L}\phi + Y_{\psi}\bar{N}\psi_{L}\phi + Y_{\nu}'\bar{N}'l_{L}\tilde{H}^{\dagger}$$

• Gauge interactions:

$$\mathscr{L} = \sum_{f} \mathcal{Q}_{f} g Z'_{\mu} \bar{f} \gamma^{\mu} f$$

with $f = \{\chi_I, \chi_R, \psi_I, \psi_R\}$

Parameters of interest:

$$\{m_{\psi}, m_{Z'}, v_{\phi}, \theta_{\nu\chi}, N_{\chi}\}$$











CB, T. Schwetz, D. Vatsyayan, [2410.23926 [hep-ph]]







• At early times, DM produced from decay of N and N' (little abundance)



CB, T. Schwetz, D. Vatsyayan, [2410.23926 [hep-ph]]







- At early times, DM produced from decay of N and N' (little abundance)
- (reaches equilibrium abundance thermalising in the dark sector)



CB, T. Schwetz, D. Vatsyayan, [2410.23926 [hep-ph]]

• Once the population of Z' becomes relevant, DM mainly produced from Z' decays, or $Z'Z' \leftrightarrow \psi \psi$, and $\chi \chi \leftrightarrow \psi \psi$





- At early times, DM produced from decay of N and N' (little abundance)
- (reaches equilibrium abundance thermalising in the dark sector)
- At late times, DM freezes-out via annihilations $\psi \psi \rightarrow \chi \chi$ (possibly avoiding DM overproduction)



CB, T. Schwetz, D. Vatsyayan, [2410.23926 [hep-ph]]

• Once the population of Z' becomes relevant, DM mainly produced from Z' decays, or $Z'Z' \leftrightarrow \psi \psi$, and $\chi \chi \leftrightarrow \psi \psi$









Suppression of cosmological bound on Σm_{ν} :

is the factor $\left(\frac{n_{\nu}}{n_{\nu}^{SM}}\right) < 1$ by which the limit on Σm_{ν} is relaxed in $\Sigma m_{\nu} \left(\frac{n_{\nu}}{n_{\nu}^{SM}}\right) \left(\frac{n_{\nu}^{0}}{56 \text{cm}^{-3}}\right) < 0.12 \text{ eV}$









Suppression of cosmological bound on Σm_{ν} :

is the factor $\left(\frac{n_{\nu}}{n_{\nu}^{SM}}\right) < 1$ by which the limit on Σm_{ν} is relaxed in $\Sigma m_{\nu} \left(\frac{n_{\nu}}{n_{\nu}^{SM}}\right) \left(\frac{n_{\nu}^{0}}{56 \text{cm}^{-3}}\right) < 0.12 \text{ eV}$

Deviation from standard value of effective number of neutrino species :

$$\frac{\frac{8}{7}\left(\frac{11}{4}\right)^{4/3}\frac{\rho_{dark}}{\rho_{\gamma}} = \frac{g_{\nu} + \tilde{g}}{2}\left(\frac{T_{dark}}{T_{\nu}^{SM}}\right)^{4} = \frac{g_{\nu} + \tilde{g}}{2}\left(\frac{g_{\nu} + \tilde{g} + g_{\psi} + \frac{8}{7}g_{Z'}}{(g_{\nu} + \tilde{g})^{1/3}}\right)^{4}$$

MU Days 2024 - 12.12.2024, DESY Hamburg



- We considered an extension of the SM neutrino sector, by addition of
 - 4 copies of heavy RH neutrinos, N and N', that participate in two separate seesaw mechanisms,
 - I sterile neutrino DM candidate ψ ,
 - N_{χ} families of massless dark fermions χ ,



- I gauge boson Z' relative to a new U(1) symmetry + I scalar singlet ϕ that breaks the new symmetry.





- We considered an extension of the SM neutrino sector, by addition of
 - 4 copies of heavy RH neutrinos, N and N', that participate in two separate seesaw mechanisms,
 - I sterile neutrino DM candidate ψ ,
 - N_{γ} families of massless dark fermions χ ,
- SM bath, depleting n_{ν}^0 and subsequently noticeably relaxing the cosmological bound on Σm_{ν} .



- I gauge boson Z' relative to a new U(1) symmetry + I scalar singlet ϕ that breaks the new symmetry.

• The N_{γ} species of χ fermions are produced at the expenses of active neutrinos after their decoupling from the





- We considered an extension of the SM neutrino sector, by addition of
 - 4 copies of heavy RH neutrinos, N and N', that participate in two separate seesaw mechanisms,
 - I sterile neutrino DM candidate ψ ,
 - N_{γ} families of massless dark fermions χ ,
- SM bath, depleting n_{ν}^0 and subsequently noticeably relaxing the cosmological bound on Σm_{ν} .
- that is efficiently populated via interactions with active neutrinos in the interval of time within BBN and recombination.



- I gauge boson Z' relative to a new U(1) symmetry + I scalar singlet ϕ that breaks the new symmetry.

• The N_{γ} species of χ fermions are produced at the expenses of active neutrinos after their decoupling from the

• The DM candidate ψ is produced in the correct abundance via freeze-out after thermalisation of the dark sector



- We considered an extension of the SM neutrino sector, by addition of
 - 4 copies of heavy RH neutrinos, N and N', that participate in two separate seesaw mechanisms,
 - I sterile neutrino DM candidate ψ ,
 - N_{γ} families of massless dark fermions χ ,
- SM bath, depleting n_{ν}^0 and subsequently noticeably relaxing the cosmological bound on Σm_{ν} .
- that is efficiently populated via interactions with active neutrinos in the interval of time within BBN and recombination.
- future CMB missions.



- I gauge boson Z' relative to a new U(1) symmetry + I scalar singlet ϕ that breaks the new symmetry.

• The N_{γ} species of χ fermions are produced at the expenses of active neutrinos after their decoupling from the

• The DM candidate ψ is produced in the correct abundance via freeze-out after thermalisation of the dark sector

• Our model predicts a sizable deviation of N_{eff} from the SM value at recombination, that may be observable by





<u>cristina.benso@kit.edu</u>

MU Days 2024 - 12.12.2024, DESY Hamburg



BACKUP SLIDES





CONSTRAINTS AND PREDICTIONS





Remember: Relevant parameters for phenomenology $\{m_{\psi}, m_{Z'}, v_{\phi}, \theta_{\nu\chi}, N_{\chi}\}$

here $N_{\chi} = 10$





CONSTRAINTS AND PREDICTIONS





Remember:

Relevant parameters for phenomenology $\{m_{\psi}, m_{Z'}, v_{\phi}, \theta_{\nu\chi}, N_{\chi}\}$

here $N_{\gamma} = 10$

Majority of constraints from requirement of equilibrium or non-equilibrium of various processes within the dark sector or involving also active neutrinos





CONSTRAINTS AND PREDICTIONS



cristina.benso@kit.edu



Thermalization:

• ν must thermalise with Z' in the interval 100 keV > T > 10 eV

 \longrightarrow condition: $\langle \Gamma(Z' \leftrightarrow \nu \nu) \rangle > H(T \sim m_{Z'}/3);$

• ν should not be in thermal equilibrium with Z' at T > 0.7 MeV

 \longrightarrow condition: $\langle \Gamma(Z' \leftrightarrow \nu \nu) \rangle < H(T = 0.7 \text{ MeV});$

• an existing abundance of χ must not grow exponentially before BBN \square condition: $\langle \Gamma(\nu\chi \leftrightarrow \chi\chi) \rangle < H(T = 0.7 \text{ MeV});$

• CMB must not be distorted by $\nu\nu \leftrightarrow Z'$ and $Z' \leftrightarrow \chi\chi$ at $z < 10^{\circ}$ $\longrightarrow \langle \Gamma(\nu\nu\leftrightarrow Z') \rangle < H(T = 23 \text{ eV}) \text{ and } \langle \Gamma(Z'\leftrightarrow \chi\chi) \rangle < H(T = 23 \text{ eV});$ • CMB must not be perturbed by χ free-streaming at $z < 10^5$ $\longrightarrow \langle \Gamma(\chi\chi \leftrightarrow \chi\chi) \rangle < H(T = 23 \text{ eV})$

$$\Delta N_{eff} \simeq 0.014 \sum_{\chi=1}^{N_{\chi}} \frac{|\theta_{e\chi}|^2 + 0.8(|\theta_{\mu\chi}|^2 + |\theta_{\tau\chi}|^2)}{10^{-6}} \left(\frac{m_{\nu}}{0.1 \text{ eV}}\right) < 0$$

MU Days 2024 - 12.12.2024, DESY Hamburg







- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature:
 - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)
 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses:

* KATRIN Collaboration, [2406.13516 [nucl-ex]]

** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu







- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses: - KATRIN aims to measure the effective electron antineutrino mass $m_{\nu_e} = \sqrt{\Sigma |U_{ei}|^2 m_{\nu_i}^2}$ current upper limit $m_{\nu_{e}} < 0.45$ eV *, expected final reach $m_{\nu_{e}} = 0.2$ eV;







- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses: - KATRIN aims to measure the effective electron antineutrino mass $m_{\nu_e} = \sqrt{\Sigma |U_{ei}|^2 m_{\nu_i}^2}$ current upper limit $m_{\nu_e} < 0.45$ eV *, expected final reach $m_{\nu_e} = 0.2$ eV; - Oscillation data put a lower limit on the sum of neutrino masses:

- $\Sigma m_{\nu} > 0.058 (0.098)$ eV for normal (inverted) neutrino mass ordering. **

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu







- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses: - KATRIN aims to measure the effective electron antineutrino mass $m_{\nu_e} = \sqrt{\Sigma |U_{ei}|^2 m_{\nu_i}^2}$ current upper limit $m_{\nu_e} < 0.45$ eV *, expected final reach $m_{\nu_e} = 0.2$ eV; - Oscillation data put a lower limit on the sum of neutrino masses:

- $\Sigma m_{\nu} > 0.058 (0.098)$ eV for normal (inverted) neutrino mass ordering. **
- Cosmological observations set stringent constraints on the sum of active neutrino masses, assuming ΛCDM : for example, DESI established an upper bound of $\Sigma m_{\nu} < 0.072$ eV ***

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu



****DESI Collaboration, [2404.03002 [astro-ph.CO]]

MU Days 2024 - 12.12.2024, DESY Hamburg





- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)
- - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses: - KATRIN aims to measure the effective electron antineutrino mass $m_{\nu_e} = \sqrt{\Sigma |U_{ei}|^2 m_{\nu_i}^2}$ current upper limit $m_{\nu_e} < 0.45$ eV *, expected final reach $m_{\nu_e} = 0.2$ eV; - Oscillation data put a lower limit on the sum of neutrino masses: $\Sigma m_{\nu} > 0.058 (0.098)$ eV for normal (inverted) neutrino mass ordering. **

- Cosmological observations set stringent constraints on the sum of active neutrino masses, assuming ΛCDM : for example, DESI established an upper bound of $\Sigma m_{\nu} < 0.072$ eV ***
- What if KATRIN measures something? How could laboratory results be reconciled with cosmological limits?

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu



****DESI Collaboration, [2404.03002 [astro-ph.CO]]

MU Days 2024 - 12.12.2024, DESY Hamburg



- Standard Model is great <u>but</u> it does not explain (at least) two puzzles of Nature: - active neutrino masses \longrightarrow seesaw mechanism (3 heavy RH Majorana neutrinos N)

 - dark matter \longrightarrow sterile neutrino DM (ψ)
- Lab. experiments aim to measure directly the small value of active neutrino masses: - KATRIN aims to measure the effective electron antineutrino mass $m_{\nu_e} = \sqrt{\Sigma |U_{ei}|^2 m_{\nu_i}^2}$ current upper limit $m_{\nu_e} < 0.45$ eV *, expected final reach $m_{\nu_e} = 0.2$ eV; - Oscillation data put a lower limit on the sum of neutrino masses:

- $\Sigma m_{\nu} > 0.058 (0.098)$ eV for normal (inverted) neutrino mass ordering. **
- Cosmological observations set stringent constraints on the sum of active neutrino masses, assuming $\Lambda {\sf CDM}$: for example, DESI established an upper bound of $\Sigma m_{\nu} < 0.072$ eV ***
- What if KATRIN measures something? How could laboratory results be reconciled with cosmological limits?
- Is it possible that the same dark sector that makes laboratory measurement compatible with cosmological limits provides also a viable dark matter candidate?

* KATRIN Collaboration, [2406.13516 [nucl-ex]] ** I. Esteban et al, [2410.05380 [hep-ph]]

cristina.benso@kit.edu



MU Days 2024 - 12.12.2024, DESY Hamburg



^{****}DESI Collaboration, [2404.03002 [astro-ph.CO]]