

VBF-HH at NLO QCD in HEFT

DESY theory seminar

Marius Höfer | 27 January 2025

with Jens Braun, Pia Bredt, Gudrun Heinrich

Outline

1. Introduction

2. EFTs for VBF-HH

3. Framework: Whizard+GoSam

4. Process setup and validation

5. Phenomenology

6. Summary and Outlook

Introduction
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EFTs for VBF-HH
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Framework: Whizard+GoSam
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Phenomenology
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Summary and Outlook
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We found the Higgs ...

Introduction
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EFTs for VBF-HH
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Summary and Outlook
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We found the Higgs ... Now we have to understand it!

Many parameters of the Higgs sector still poorly constrained

↳ relates to many open questions: origin of the Higgs mass, hierarchy problem, vacuum stability, ...

HL-LHC: exploring parameters of the Higgs sector, theory has to keep up

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HL-LHC: exploring parameters of the Higgs sector, theory has to keep up

Multi-Higgs couplings

Current bounds provide ample room for New Physics! How to measure?

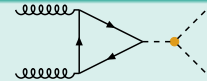
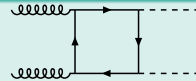
- Single Higgs production at (at least) NLO EW
- Double and triple Higgs production at LO → cross section is extremely small

Higgs Pair production

Most relevant channels: gluon fusion (ggF) and vector boson fusion (VBF)

ggF

- Main production channel
- Loop induced



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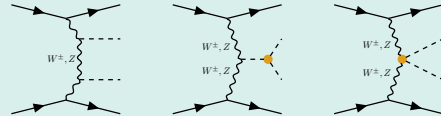
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VBF

- Relatively clean wrt. QCD background
- Sensitive to $g_{hhvv} - g_{hvv}^2$: unitarity, HEFT vs SMEFT
- In SM: $\sigma_{VBF} \approx 1.7\text{fb} \approx 5\%$ of σ_{ggF} (at 13 TeV)



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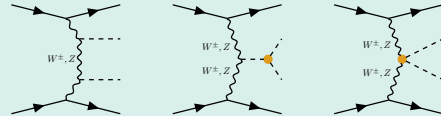
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- “new” couplings in HH production

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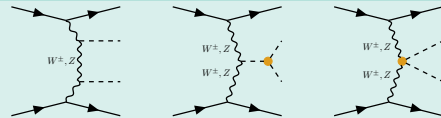
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Status of VBF-HH

Experiment

- Bounds on prod. cross sec.: $\sigma_{VBF-HH}^{exp.} \lesssim \sigma_{VBF-HH}^{SM} \times \begin{cases} 50 & [\text{ATLAS @ LHC-HWG'24}] \\ 91 & [\text{CMS PAS '24}] \end{cases}$
- Rather loose constraints on g_{hvv} , g_{hhvv} and g_{hhh} [ATLAS 2024][CMS 2024]

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Theory

- NNLO QCD + NLO EW [Dreyer et al. 2020]
 - N³LO QCD [Dreyer et al. 2018]
 - NNLO QCD [Dreyer et al. 2019][Ling et al. 2014]
 - NLO QCD [Fidjy 2008][Baglio et al. 2013][Frederix et al. 2014]
- } Standard Model

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 - NLO QCD [Fidgy 2008][Baglio et al. 2013][Frederix et al. 2014]
 - In EFT context, at tree-level: g_{hhvv} only [Bishara et al. 2017], g_{hhvv} and g_{hhh} [Kilian et al. 2020]
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New Physics hiding in VBF-HH?

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graph TD; A[New Physics hiding in VBF-HH?] --> B[Model building]; A --> C[Effective Field Theory];
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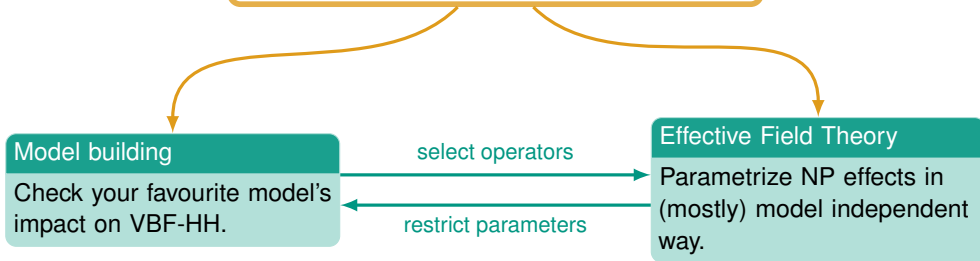
Model building

Check your favourite model's impact on VBF-HH.

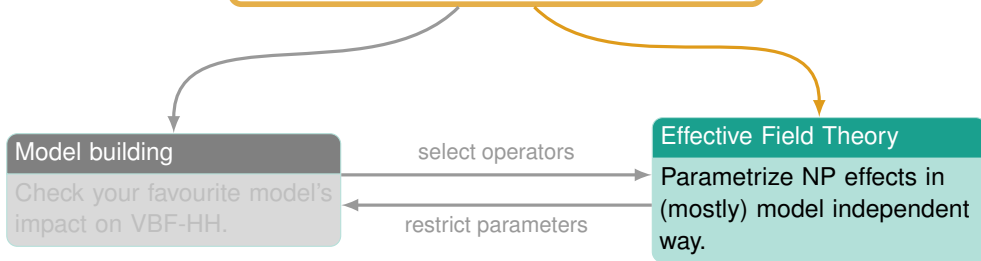
Effective Field Theory

Parametrize NP effects in (mostly) model independent way.

New Physics hiding in VBF-HH?



New Physics hiding in VBF-HH?



The EFT approach to New Physics

The SM as an EFT at the EW scale

General requirements:

- Obey (gauge and global) symmetries of the SM
- Field content = SM fields

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Construction

$$\mathcal{L}_{eff} = \mathcal{L}_{SM} + \sum_i C_i \mathcal{O}_i$$

- Coefficients C_i absorb unknown UV dynamics
- Systematic power counting organizes terms according to relevance (e.g. dim-6, dim-8, ... in SMEFT)

Two common EFT realizations

SMEFT

- EWSB realized as in SM, have ordinary Higgs doublet Φ
- Must have weakly coupled NP, NP decouples from EW scale physics
- Organization of power counting according to canonical dimension

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HEFT

- EWSB not necessarily SM-like, split Goldstone modes U and Higgs scalar h
- Can have strongly coupled NP in Higgs sector, does not decouple
- Organization of power counting in chiral dimension (loop counting)

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- Organization of power counting in chiral dimension (loop counting)
- Note: sometimes called the “Electroweak Chiral Lagrangian” (EWChL)

The HEFT-Lagrangian at LO

$$\begin{aligned}
\mathcal{L}_{SM} = & -\frac{1}{2} \langle G_{\mu\nu} G^{\mu\nu} \rangle - \frac{1}{2} \langle W_{\mu\nu} W^{\mu\nu} \rangle - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} + \bar{\psi} i \not{D} \psi \\
& + \frac{1}{2} \partial_\mu h \partial^\mu h - v^2 \left[\frac{m_h^2}{2} \left(\frac{h}{v} \right)^2 + \frac{m_h^2}{2} \left(\frac{h}{v} \right)^3 + \frac{m_h^2}{8} \left(\frac{h}{v} \right)^4 \right] \\
& + \frac{v^2}{4} \langle D_\mu U^\dagger D^\mu U \rangle \left[1 + 2 \frac{h}{v} + \left(\frac{h}{v} \right)^2 \right] \\
& - v \left[\bar{q}_L \left(Y_u + Y_u \frac{h}{v} \right) U \begin{pmatrix} u_R \\ 0 \end{pmatrix} + \text{h.c.} + \dots \right]
\end{aligned}$$

- Goldstone bosons of EWSB:

$$U = \exp \left(2i \frac{\varphi_\alpha t^\alpha}{v} \right)$$

- Higgs doublet:

$$\Phi = U \begin{pmatrix} 0 \\ \frac{v+h}{\sqrt{2}} \end{pmatrix}$$

$\Rightarrow$ correlations between
 h -couplings

- Renormalizable in the
classical sense

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& + \frac{1}{2} \partial_\mu h \partial^\mu h - v^2 \left[\frac{m_h^2}{2} \left(\frac{h}{v} \right)^2 + \sum_{i=3}^{\infty} v^{(i)} \left(\frac{h}{v} \right)^i \right] \\
& + \frac{v^2}{4} \langle D_\mu U^\dagger D^\mu U \rangle \left[1 + \sum_{i=1}^{\infty} f_U^{(i)} \left(\frac{h}{v} \right)^i \right] \\
& - v \left[\bar{q} \left(Y_u + \sum_{i=1}^{\infty} Y_u^{(i)} \left(\frac{h}{v} \right)^i \right) U \begin{pmatrix} u_R \\ 0 \end{pmatrix} + \text{h.c.} + \dots \right]
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- Goldstone bosons of EWSB:

$$U = \exp \left(2i \frac{\varphi_\alpha t^\alpha}{v} \right)$$

- h not part of a doublet
 $\Rightarrow$ no correlations
- h couples strongly to Goldstone sector $\Rightarrow$ arbitrary powers of h
- Non-renormalizable in the classical sense

HEFT power counting

- Characteristic scale of HEFT: v
 - ↳ LO Lagrangian contains all terms of order $\sim v^4$
- Strong dynamics in scalar sector implies EFT cutoff: $\Lambda \sim 4\pi v$
- HEFT is constructed as an expansion in $\frac{v^2}{\Lambda^2} \sim \frac{1}{16\pi^2} \Rightarrow$ loop expansion

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Chiral dimension d_χ

$$[A_\mu, \varphi, h]_\chi = 0, \quad [\partial, \bar{\psi}\psi, g, y]_\chi = 1 \quad \Rightarrow \quad [\mathcal{O}]_\chi = 2L + 2 \quad (L: \text{loop order})$$

↳ LO Lagrangian $\mathcal{L}_2$ has $d_\chi = 2$ (0-loop order / tree level)

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Note: custodial symmetry breaking operator at $d_\chi = 2$: $\mathcal{O}_{\beta_1} \sim v^2 \langle U^\dagger D_\mu U t^3 \rangle^2 F_{\beta_1}(h)$

- Highly constrained from EW T -parameter $\Rightarrow$ loop-suppressed $\Rightarrow$ NLO term

The HEFT Lagrangian at NLO

Construction of the NLO Lagrangian

- $\mathcal{L}_2$ renormalizable in the modern sense: introduction of counterterms order by order in loop expansion
- At NLO ($d_\chi = 4$): terms needed as counterterms for 1-loop diagrams constructed from $\mathcal{L}_2$
- From (momentum) power counting arguments: finite number of divergent diagrams at each order

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Operator classes

UhD^4 , $g^2 X^2 Uh$, $gXUhD^2$, $y^2 \psi^2 UhD$, $y\psi^2 UhD^2$, $y^2 \psi^4 Uh$

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Full set of 1-loop RGEs available since 2020 [\[Buchalla et al. 2020\]](#)

Parametrization of anomalous Higgs couplings

HEFT for the practitioner

Not all Operators relevant for given processes: e.g. (double) Higgs production, Higgs decays

- Assumptions:
- CP conservation
 - Custodial symmetry conserved at LO (Wh couplings = Zh couplings)

$$\mathcal{L}_{eff} \supset \left(2c_V \frac{h}{v} + c_{2V} \frac{h^2}{v^2} \right) \left(m_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) - \sum_f m_f \left(c_f \frac{h}{v} + c_{2f} \frac{h^2}{v^2} \right) \bar{f} f + c_\lambda \frac{m_h^2}{2v} h^3$$

$$+ \frac{\alpha}{8\pi} \left(c_{\gamma\gamma} \frac{h}{v} + c_{2\gamma\gamma} \frac{h^2}{v^2} \right) F_{\mu\nu} F^{\mu\nu} + \frac{\alpha_s}{8\pi} \left(c_{gg} \frac{h}{v} + c_{2gg} \frac{h^2}{v^2} \right) G_{\mu\nu}^a G^{a\mu\nu}$$

SM: $c_V = c_{2V} = c_f = c_\lambda = 1$, $c_{2f} = c_{\gamma\gamma} = c_{2\gamma\gamma} = c_{gg} = c_{2gg} = 0$

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 $\Rightarrow$ operators $\sim h \langle W_{\mu\nu} W^{\mu\nu} \rangle$ from $\mathcal{L}_4$ subleading (all other couplings = Zh couplings)

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SM: $c_V = c$

h -photon and h -gluon interactions loop-induced
 $\Rightarrow \mathcal{L}_4$ terms contribute to leading NP effects

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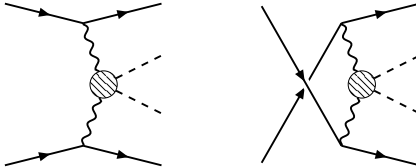
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$$+ \frac{\alpha}{8\pi} \left(c_{\gamma\gamma} \frac{h}{v} + c_{2\gamma\gamma} \frac{h^2}{v^2} \right) F_{\mu\nu} F^{\mu\nu} + \left[\begin{array}{l} \text{■ At NLO QCD: No additional vertex structures!} \\ \Rightarrow \text{corresponds to } \kappa\text{-framework} \\ \text{■ } \mathcal{L}_4 \text{ becomes relevant at NLO EW} \end{array} \right] \leftarrow \mathcal{L}_4$$

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VBF-HH in HEFT: LO

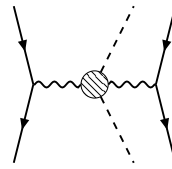
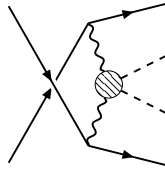
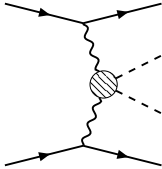
$$\mathcal{L}_{\text{eff}} \supset \left(2c_V \frac{h}{v} + c_2 v \frac{h^2}{v^2} \right) \left(m_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} m_Z^2 Z_\mu Z^\mu \right) + c_\lambda \frac{m_h^2}{2v} h^3$$



$$\text{Shaded Circle} \in \left\{ \begin{array}{c} \text{Diagram 1: } q \bar{q} \rightarrow q \bar{q} h \text{ via } t\text{-channel } h \text{ exchange} \\ \text{Diagram 2: } q \bar{q} \rightarrow q \bar{q} h \text{ via } s\text{-channel } h \text{ exchange} \\ \text{Diagram 3: } q \bar{q} \rightarrow q \bar{q} h \text{ via } s\text{-channel } Z \text{ exchange} \end{array} \right\}$$

VBF-HH in HEFT: LO

$$\mathcal{L}_{\text{eff}} \supset \left(2c_V \frac{h}{v} + c_2 v \frac{h^2}{v^2} \right) (m_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} m_Z^2 Z_\mu Z^\mu) + c_\lambda \frac{m_h^2}{2v} h^3$$

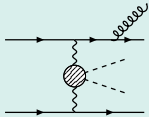


Same coupling order $\mathcal{O}(g_{ew}^4)$
 $\Rightarrow$ included in our calculation,
suppressed by VBF-cuts:
large rapidity separation and
invariant mass in jet-system

$$\text{Higgs} \in \left\{ \begin{array}{c} \text{W} \text{ exchange} \\ \text{Z} \text{ exchange} \\ \text{Box} \end{array} \right\}$$

VBF-HH in HEFT: NLO QCD

Real corrections:



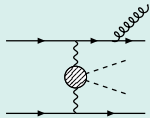
- + u - and s -channel crossings
- + process with gluon crossed into initial state

$$\text{shaded circle} \in \left\{ \begin{array}{c} \text{t-channel gluon exchange} \\ \text{s-channel gluon exchange} \\ \text{u-channel gluon exchange} \end{array} \right\}$$

No QCD loop corrections to 
 $\Rightarrow$ renormalization as in SM

VBF-HH in HEFT: NLO QCD

Real corrections:

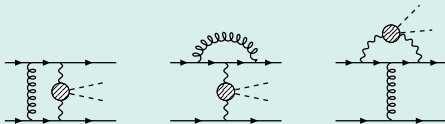


- + u - and s -channel crossings
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$$\text{shaded circle} \in \left\{ \begin{array}{c} \text{gluon exchange} \\ \text{gluon exchange} \\ \text{gluon exchange} \end{array} \right\}$$

No QCD loop corrections to shaded circle
 $\Rightarrow$ renormalization as in SM

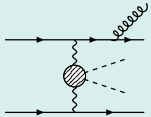
Virtual corrections:



+ u - and s -channel crossings

VBF-HH in HEFT: NLO QCD

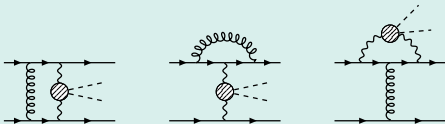
Real corrections:



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Same coupling order $\mathcal{O}(g_s^2 g_{ew}^4)$
 $\Rightarrow$ included in our calculation,
suppressed by VBF-Cuts.

Virtual corrections:



+ u - and s -channel crossings

Framework: Whizard + GoSam

- Use GoSam to generate amplitudes based on custom HEFT UFO model
- Whizard as Monte Carlo generator
- Implement flexible interface, usable for other processes

Whizard + GoSam

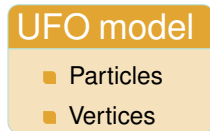
sindarin runcard:
define process, specify parameters



Whizard (Monte Carlo event generator)

GoSam (amplitude provider)

Whizard + GoSam



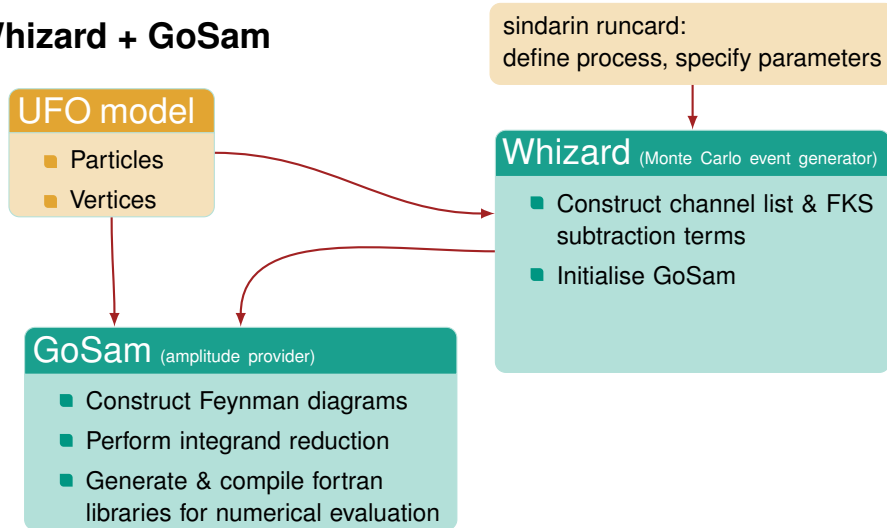
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Whizard (Monte Carlo event generator)

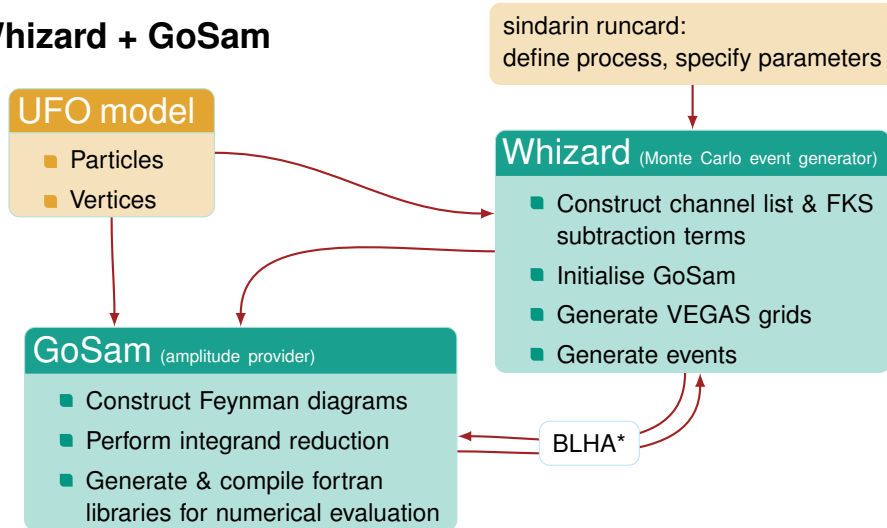
- Construct channel list & FKS subtraction terms

GoSam (amplitude provider)

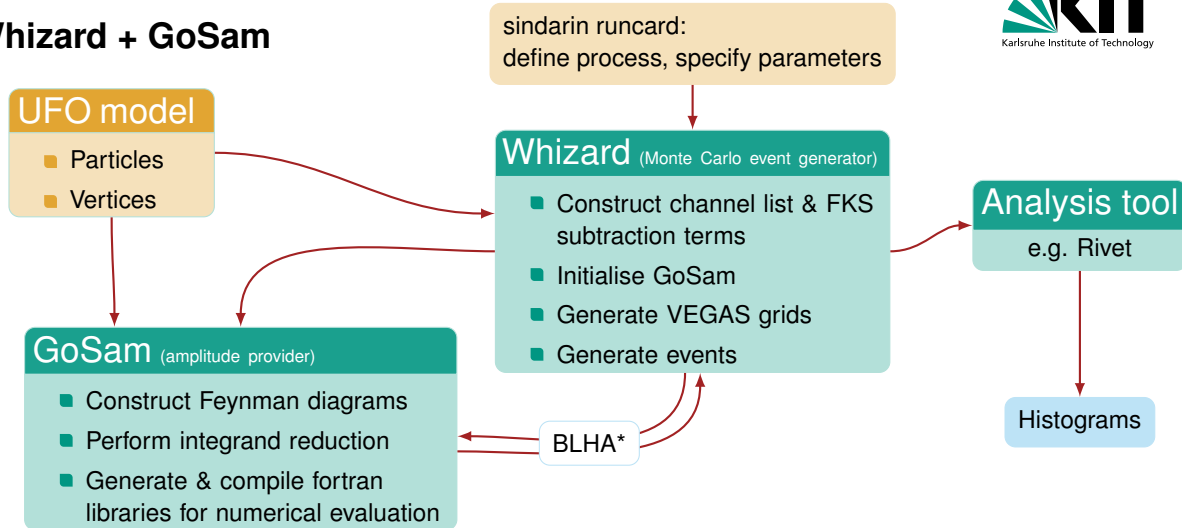
Whizard + GoSam



Whizard + GoSam



Whizard + GoSam



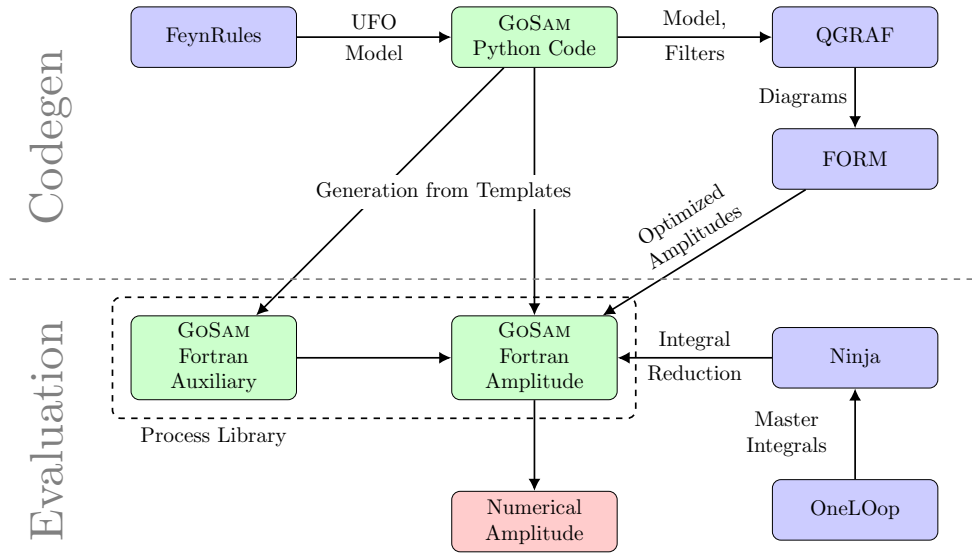


figure by Jens Braun

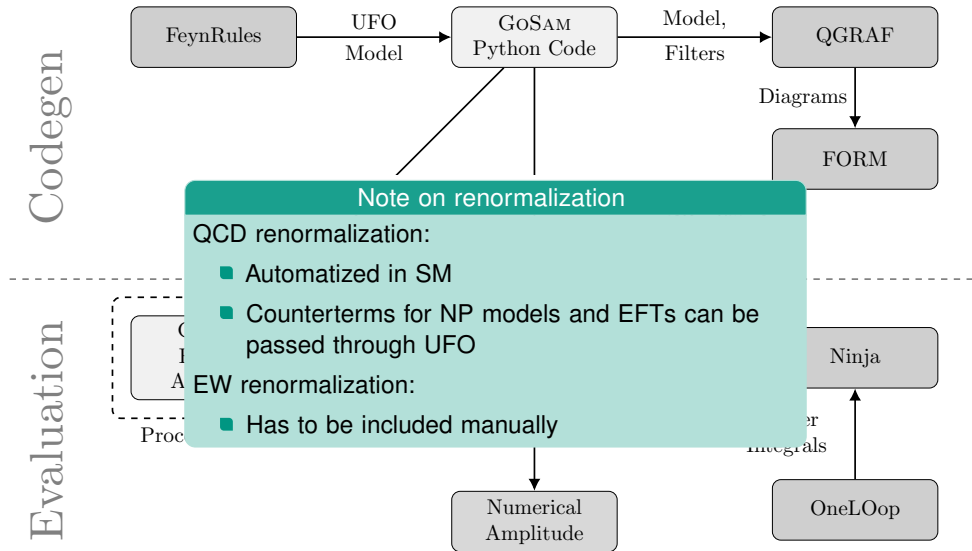


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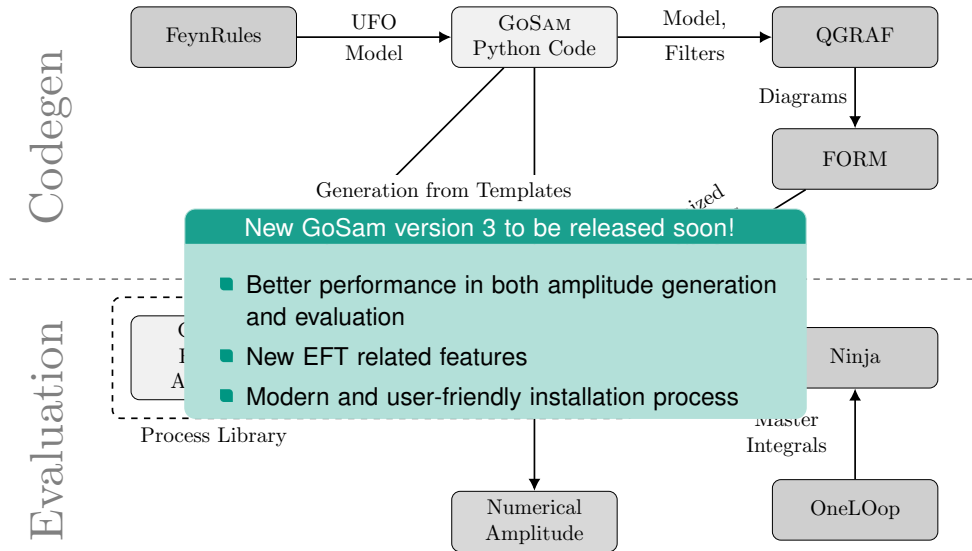
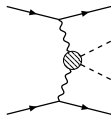


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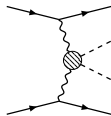
VBF-HH: Process setup



$pp \rightarrow hh + 2\text{jets}$ at $\mathcal{O}(\alpha_{ew}^4, \alpha_{ew}^4 \alpha_s) + c_V, c_{2V}, c_\lambda$

- $\sqrt{s} = 13.6\text{TeV}$
- PDF set: PDF4LHC21_mc
- Central renormalization and factorization scale: $\mu_0 = \left(\frac{m_h^2}{4} \left(\frac{m_h^2}{4} + p_{T,hh}^2 \right) \right)^{\frac{1}{4}}$
- 3-point scale variation $\mu_R = \mu_F = a\mu_0$ with $a = 0.5, 1, 2$
- Jets: anti- k_t with $R = 0.4$, at least two jets with $p_{T,j} > 20\text{GeV}$ and $|y_j| < 4.5$
- VBF cuts: $m_{j_1 j_2} > 600\text{GeV}$ and $\Delta\eta(j_1, j_2) > 4.0$

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For Validation:

$c_V = c_{2V} = c_\lambda = 1$ (SM)

$\mu_R = \mu_F = 2m_h$

VBF-HH: Validation

SM case ($c_V = c_{2V} = c_\lambda = 1$) can be validated against other tools.

- Matrix elements: check against OpenLoops2 for 10^6 random phase space points
↳ avg. 13 (8) digit agreement for tree (1-loop)

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- Distributions:

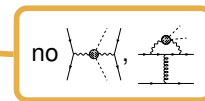
	LO	NLO	
Whizard+OpenLoops2	✓	✓	Checks OL provider
VBFNLO	✓	—	VBF-approximation
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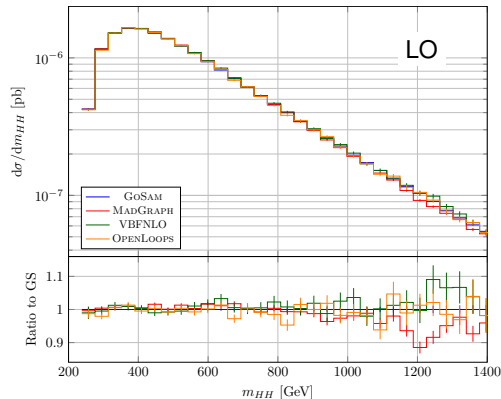
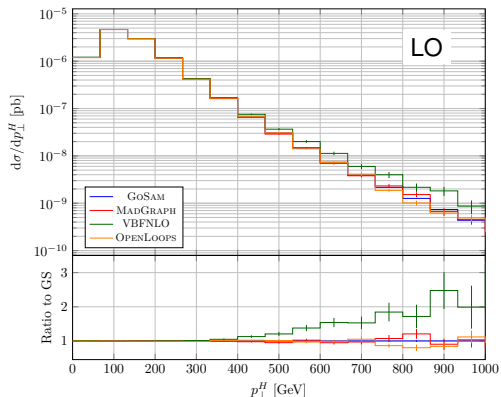
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Phenomenology

Introduction
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EFTs for VBF-HH
oooooooo

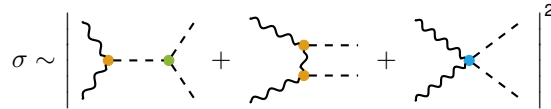
Framework: Whizard+GoSam
ooo

Process setup and validation
ooo

Phenomenology
●oooooooo

Summary and Outlook
o

Parametrization of the total cross section

$$\sigma \sim \left| \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} \right|^2$$


$$\Rightarrow \frac{\sigma}{\sigma_{\text{SM}}} = A_0 c_\lambda^2 c_V^2 + A_1 c_V^4 + A_2 c_{2V}^2 + A_3 c_\lambda c_V^3 + A_4 c_\lambda c_V c_{2V} + A_5 c_V^2 c_{2V}$$

- Coefficients A_i can be determined by fit after calculating σ for different values of c_i
- No new A_i at NLO QCD, but different numerical values
- Note: $\sum_i A_i = 1$ by construction

Parametrization of the total cross section

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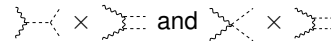
NLO coeff.	$\mu_F = \mu_r = \mu_0/2$	$\mu_F = \mu_r = \mu_0$	$\mu_F = \mu_r = 2\mu_0$
A_0	0.7011(46)	0.6889(35)	0.6830(30)
A_1	22.15(12)	21.71(9)	21.55(8)
A_2	11.86(7)	11.59(5)	11.55(4)
A_3	-6.139(42)	-6.025(33)	-5.984(27)
A_4	3.865(29)	3.786(23)	3.773(19)
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Destructive interferences:

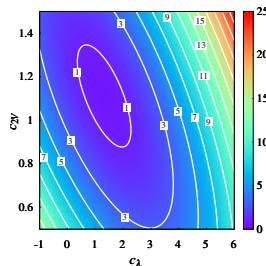
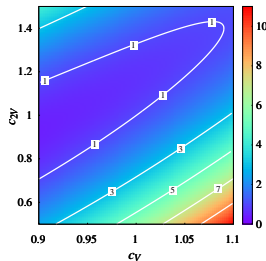
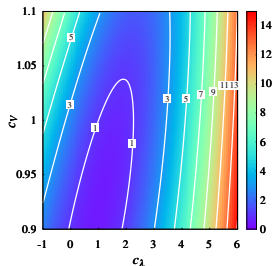


Note: depends on signs of c_i

Total cross section

Parameter ranges are chosen roughly like current experimental bounds [ATLAS 2024][CMS 2024]

$$c_V \in [0.9, 1.1], \quad c_{2V} \in [0.5, 1.5], \quad c_\lambda \in [-1, 6]$$



$\sigma/\sigma_{SM} \text{ (NLO),}$
one c_i fixed to SM

- Can have increase of cross section by more than a factor of 20.
- Current best exp. bound: $\sigma_{VBF-HH}^{exp.} \lesssim 50\sigma_{VBF-HH}^{SM}$ [ATLAS @ LHC-HWG'24]

Introduction
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○○○

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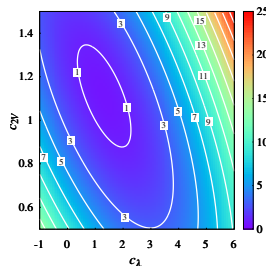
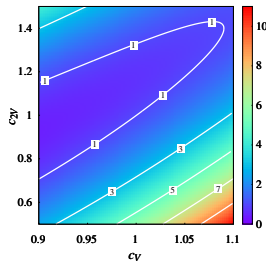
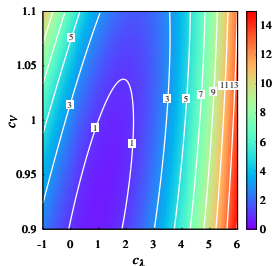
Phenomenology
○○○●○○○○○○

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○

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σ/σ_{SM} (NLO),
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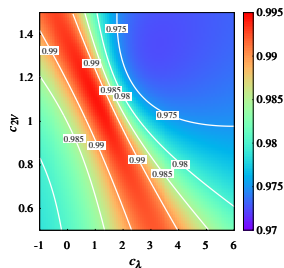
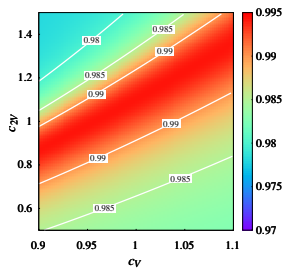
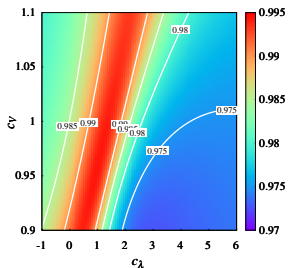
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When all three c_i allowed to float:
up to factor ~ 130 possible

Total cross section: K-factors

Parameter ranges are chosen roughly like current experimental bounds [ATLAS 2024][CMS 2024]

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$\sigma_{\text{NLO}}/\sigma_{\text{LO}}$,
one c_i fixed to SM

■ NLO QCD K-factor mostly flat: structure of EFT contributions unaffected by QCD

Distributions

Parametrization with A_i in principle possible with distributions, but:

- Requires very high statistics in each bin
- Detailed analysis of correlations for uncertainty estimation

Distributions

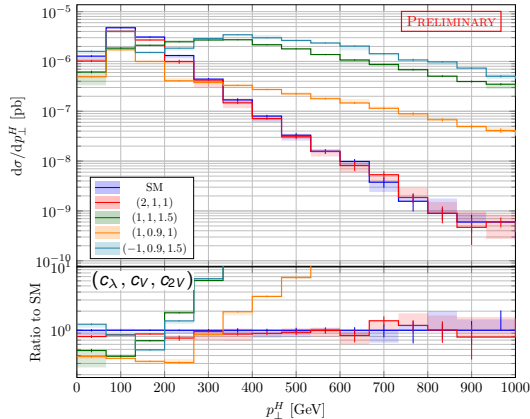
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Alternative: distributions for 12 benchmark points (c_λ, c_V, c_{2V}) incl. SM

c_λ	1	0	1	1	1	2	-1	-1	2	3	4	6
c_V	1	1	0.9	1	1	1	0.9	1.05	0.9	1.1	0.95	1.1
c_{2V}	1	1	1	0.5	1.5	1	1.5	1.3	1.4	0.5	0.5	1

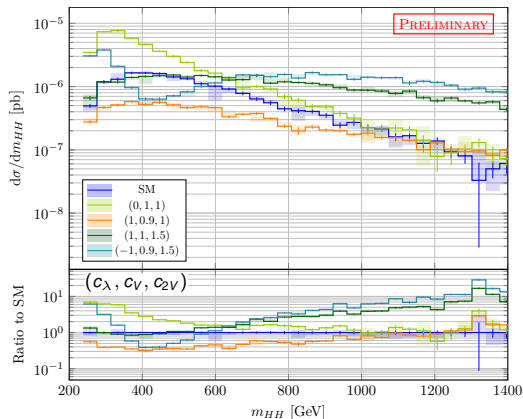
Distributions: Higgs- p_T



Note: p_T of any Higgs $\Rightarrow$ two histogram entries per event

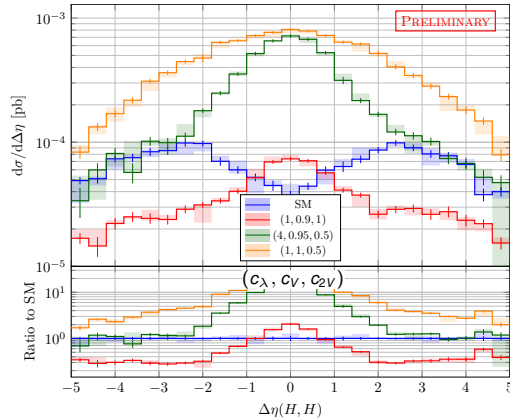
- SM: Max. around 100 GeV, exponentially falling tail.
- Doubling c_λ has no significant effect, slight reduction for low p_T .
- $c_{2V} = 1.5$ (upper exp. bound) enhances tail drastically. Reduction in low p_T . Shifts maximum.
- $c_V = 0.9$ (lower exp. bound) similar as previous, does not shift maximum.
- All c_i at exp. boundaries: enhancements in tail and in first bin, reduction in second and third bin.

Distributions: Higgs pair invariant mass



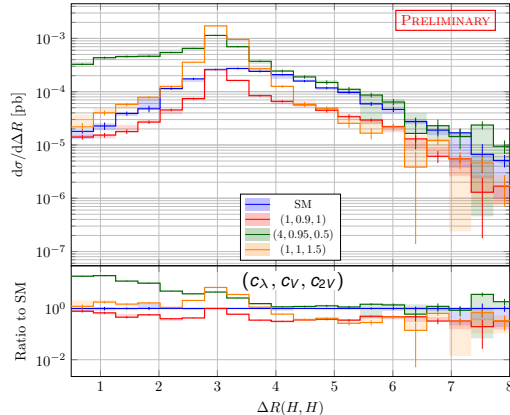
- SM: broad max. around 400 GeV, exponentially falling tail.
- Switching of triple h coupling has almost no effect on tail. Increases low m_{hh} significantly.
- $c_{2V} = 1.5$ (upper exp. bound) flattens distribution. Large enhancements in the tail.
- $c_V = 0.9$ (lower exp. bound) similar as previous, but overall reduction of the cross-section, except for high m_{hh}
- All c_i at exp. boundaries: enhancements for low and high m_{hh} . Dip around 400 GeV.

Distributions: Higgs pair rapidity separation



- SM: low rap.-separation suppressed, peaks around ± 2.5 .
- $c_V = 0.9$ (lower exp. bound) very different from SM. Low rapidity-separation favoured.
- Larger c_λ , smaller c_V , c_{2V} : similar as previous, much larger peak.
- $c_{2V} = 0.5$ (lower exp. bound) even larger shape distortions, large overall enhancement.

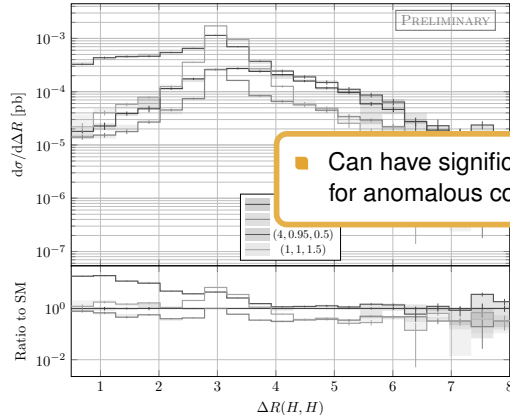
Distributions: Higgs pair separation



$$\Delta R = \sqrt{\Delta\eta^2 + \Delta\Phi^2}$$

- SM: form of $\Delta\eta$ distr. decisive: low ΔR suppressed, peak at $\Delta R \approx 3.5$.
- $c_V = 0.9$ (lower exp. bound): clear peak at $\Delta R \approx 3$, overall suppression.
- $c_{2V} = 1.5$ (upper exp. bound) clear peak at $\Delta R \approx 3$, enhancement for low ΔR , suppression for high ΔR .
- Larger c_λ , smaller c_V , c_{2V} : high ΔR comparable to SM, but low ΔR region enhanced.

Distributions: Higgs pair separation



Can have significant and interesting shape distortions for anomalous couplings within experimental bounds.

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Conclusions

- Identified leading HEFT operators for VBF-HH at NLO QCD
 - ↳ Three anomalous couplings c_λ , c_V , c_{2V} (corresponds to κ -framework)
 - ↳ Sub-leading operators will be relevant at NLO EW
- Updated Whizard-GoSam interface to investigate HEFT effects
- Determined set of coefficients A_i : total cross section for arbitrary values of the anom. couplings
 - ↳ Large enhancements wrt. SM cross section possible
- Observe significant shape distortions in distributions for anomalous couplings within allowed range

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- Consider possible constraints from unitarity
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- Add NLO EW (a lot of work...)

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Thank you!

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