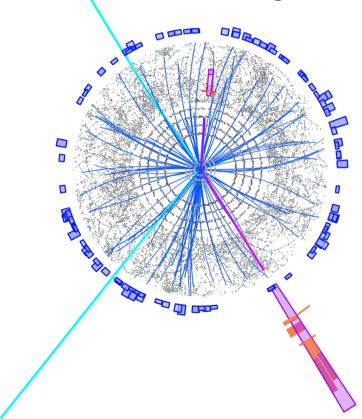
Detector Physics I

Event-Display:

Particle Detection in High Energy Physics



Simon Spannagel, DESY Introduction to the Terascale 17 March 2025

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Prologue: Units



Elementary charge: e = 1.602 176 6208 x 10⁻¹⁹ C

• Energy unit eV: $1 \text{ eV} = 1.602 \ 176 \ 6208 \ \text{x} \ 10^{-19} \ \text{J} \ (\text{CV})$

Energy of an electron, accelerated through 1V

• Electron mass: $m_p = 9.1 \times 10^{-31} \text{ kg}$

 $= 511 \text{ keV} / c^2$

• Factor c^2 often dropped, units defined with c = 1 (and $\overline{h} = 1$):

MeV, GeV, TeV...

for energy, momentum, (rest-) mass

Scientific notation: Mantissa plus exponent to Base 10: a x 10^b



What do we measure?

Determination of particle properties

Which particles can we measure?



- Particles have to be long-lived enough to reach the detectors...
 - Many elementary particles have very short life time (Higgs, W, Z...)
 - Measuring their decay products

$$d = c \tau \gamma$$
, $\gamma = \frac{1}{\sqrt{(1-\beta^2)}}$, $\beta = \frac{v}{c}$

- Particles have to interact with detector material!
 - Not every particle participates in every force
 - Possibilities of detection depends strongly on forces
- Elementary particles:

$$e^{\pm}, \mu^{\pm}, \quad V^{e}, \overline{V}^{e}, V^{\mu}, \overline{V}^{\mu}, V^{\tau}, \overline{V}^{\tau}, \quad \gamma$$

Baryons:

$$p^{\pm}$$
, n , Σ^{\pm} , Ξ_0 , Ξ^- , Ω^-

Mesons:

$$\pi^{\pm}$$
, K^{\pm} , K_{0} (K_{0}^{S} , K_{0}^{L})

$$\tau_n \approx 15 \, min$$
 $\tau_u \approx 2 \cdot 10^{-6} \, s$

Observables



• Velocity **v** Time of flight, RICH, etc.

Charge Q Bending radius in magnetic field

Lifetime τ Measurement of decay length

• Energy **E** Absorption in calorimeters

• Rest mass **m** Indirect measurement e.g. from momentum and energy or velocity

$$E^{2} = m^{2}c^{4} + p^{2}c^{2}$$
 $p = \gamma m v = \frac{mv}{\sqrt{(1-v^{2}/c^{2})}}$



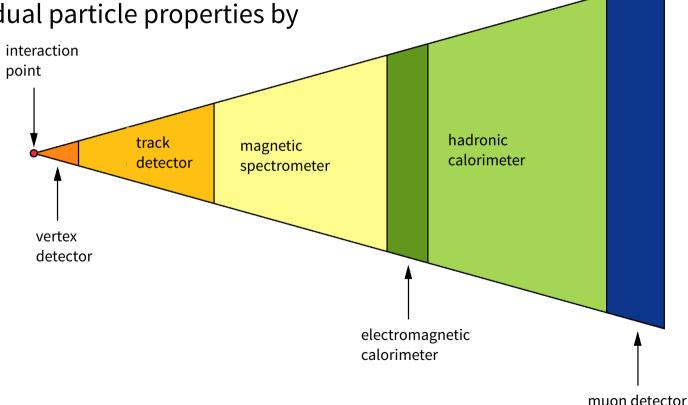
Typical Design of an Experiment



Combination of complementary detection methods

Measurement of individual particle properties by separate detectors

- Order is important!
 Some measurements are "destructive"
- Design of many experiments is very similar



Detector Terminology



Dead time:

Time period immediately after the detection of a particle during which the detector is not yet ready again to detect another particle.

- Non-paralyzable detector: newly occurring event does nothing
- Paralyzable detector: newly occurring event extends dead time

Resolution:

Achievable uncertainty on the observable

Efficiency:

Number of recorded/detected events divided by number of events that occurred



Interaction of Radiation with Matter

Energy loss and interaction processes

Interaction with Matter



- High energy particles interact via different processes with matter, depending on
 - Particle type
 - Energy
 - Material
- Energy loss of the particles via interaction
 - Energy transfer to material or other (free) particles
 - In Detectors: Energy loss = signal!

Charged Particles

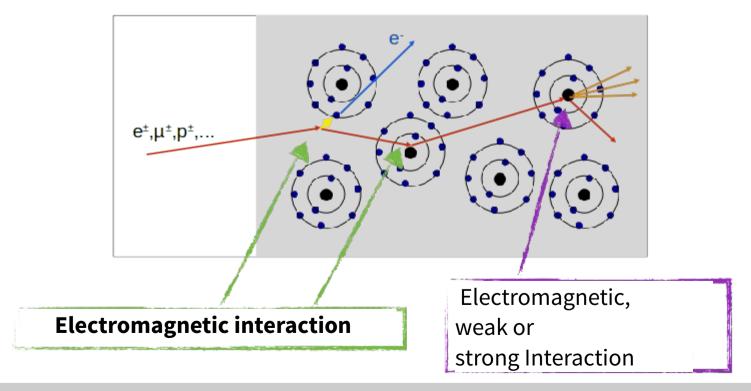


Ionization

Elastic scattering

Recoil from atom lattice → photons

Inelastic scattering



Mean Energy Loss

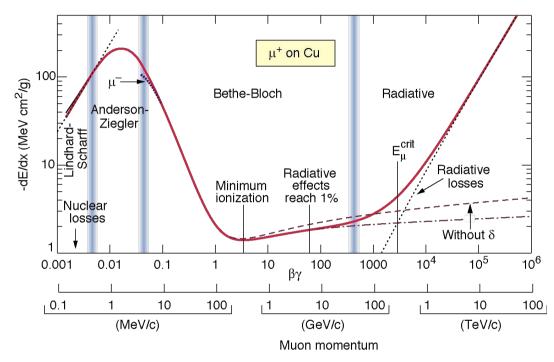


- Charged particles interact with electrons in matter
- For heavy charged particles
 Bethe-Formula

$$-\left\langle \frac{dE}{dx}\right\rangle \approx K q^2 \frac{1}{\beta^2} \frac{Z}{A} \left[\ln\left(\frac{2 m_e c^2}{I} \beta^2 \gamma^2\right) - \beta^2 \right]$$



- Properties of projectile: charge, energy
- Target properties: atomic number, ionization energy, density



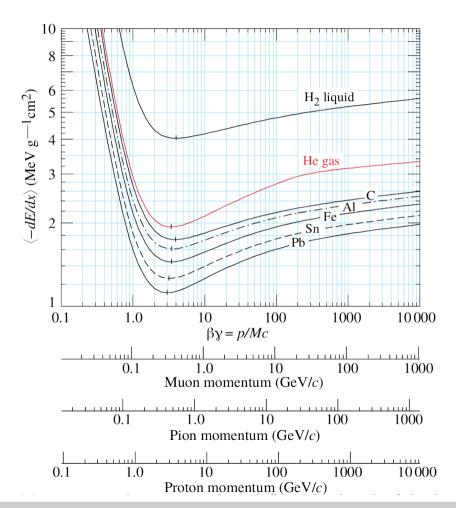
Mean Energy Loss



- Different components are dominating:
 - At low energies: $\sim 1/\beta^2$
 - At high energies: ~ ln γ

$$-\left\langle \frac{dE}{dx}\right\rangle \approx K q^2 \frac{1}{\beta^2} \frac{Z}{A} \left[\ln\left(\frac{2 m_e c^2}{I} \beta^2 \gamma^2\right) - \beta^2 \right]$$

- Minimal energy loss for $\sim \beta \gamma = 3$ MIP: Minimum Ionizing Particle
- $[dx] = g / cm^2 = cm x g / cm^3$

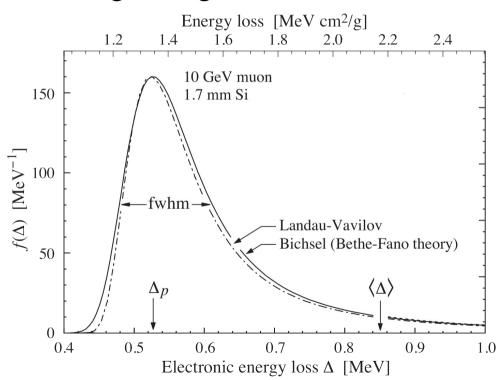


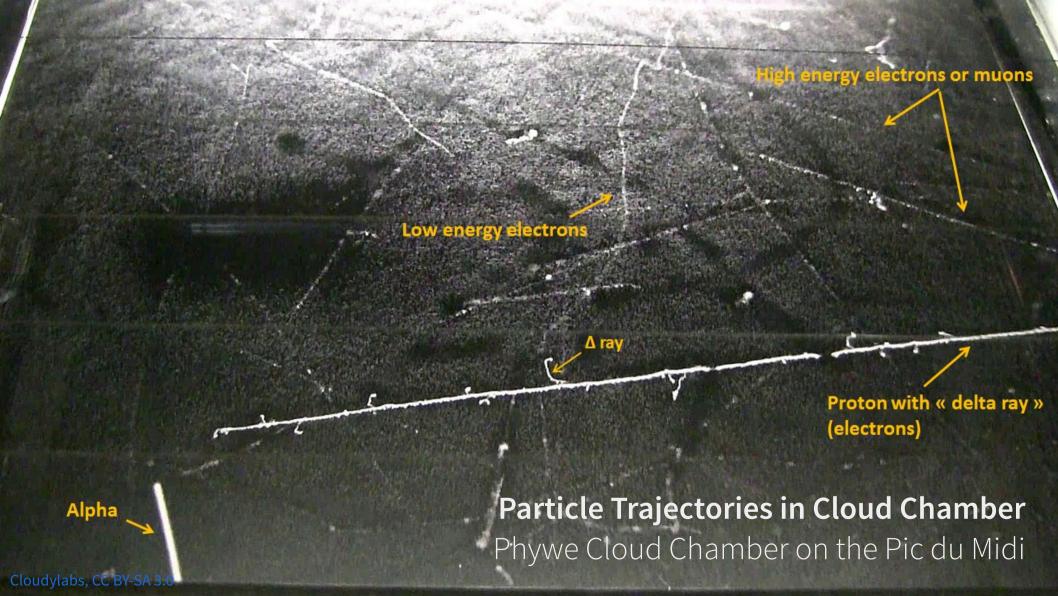
Fluctuations in Energy Loss



- Actual energy loss fluctuates around mean value
- Landau-Vavilov distribution with long tails to high energies
 - Most probable value << mean value

- Orbital electrons can receive very large energy transfers
 - Creation of delta electrons
 - Delta electrons have enough energy for further ionization

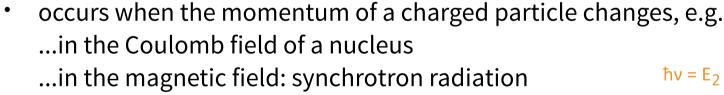




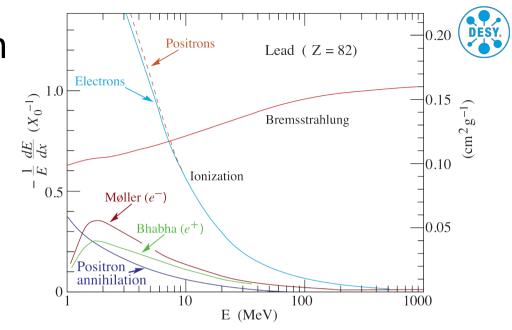
Exception: Electron & Positron

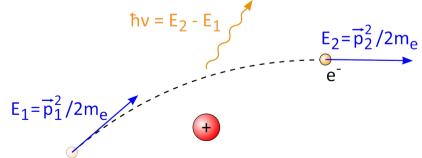
 Special case: low mass m_e = 0.511 MeV / c² m_u = 106 MeV / c² ≈ 200 m_e

At high energies: Bremsstrahlung



- Here: particles slowed down in matter
- Relevant for electrons: $-\frac{dE}{dx} \sim E \cdot \frac{1}{m^2}$





Cherenkov Radiation



Is emitted, when particle velocity > speed of light

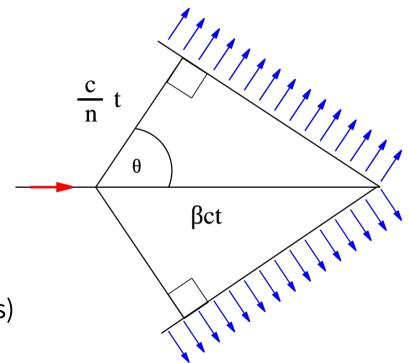
$$c_{Medium} = \frac{c_{Vacuum}}{n} < v_{Particle} < c_{Vacuum}$$

n: refractive index of medium

 Electromagnetic shockwave with conical shape is emitted under angle θ:

$$\cos(\theta) = \frac{1}{\beta n}, \quad \beta = \frac{v}{c}$$

Very low energy loss (ca. 1% of total energy loss)

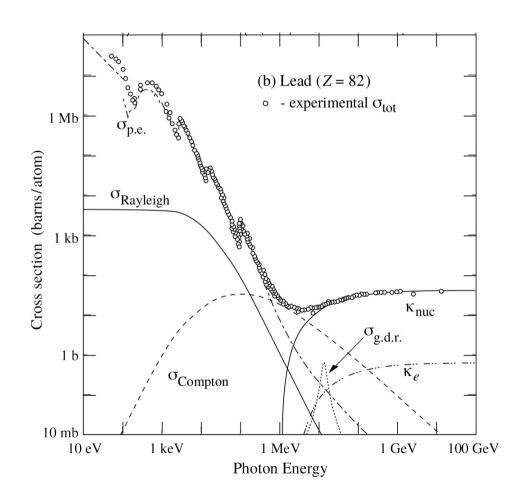




Photons



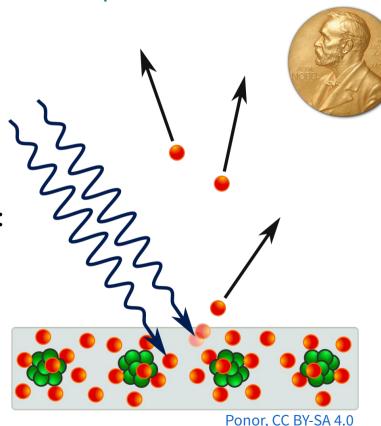
- Electromagnetic interaction
- Different processes dominate, depending on the photon energy:
 - Photo(electric) effect
 - Rayleigh scattering
 - Compton effect
 - Pair production



Photoelectric Effect



- Theoretical description of photo effect: Einstein
- Photon is absorbed by electron in atomic shell
- Transferred energy releases/frees electron γ + Atom = e^- + Ion
- Process only possible in the field of the nucleus:
 - Momentum conservation
 - Nucleus absorbs recoil
- Cross section of the photoelectric effect shows shell structure of the atom



Nobel price 1921 for Albert

Compton Effect

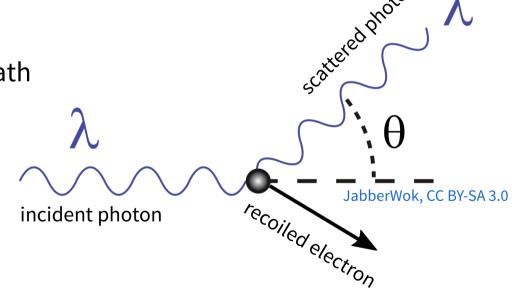


• Describes scattering of photon on "quasi-free" electron

$$\gamma$$
 + Atom = γ + e^- + Ion

Photon is deflected from its original path

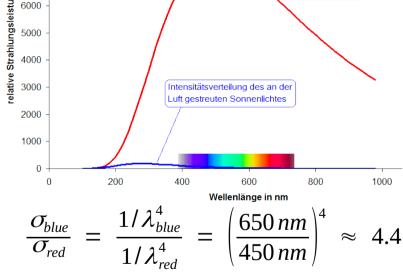
 Wavelength of photon changes through energy transfer to electron



Photons: Thomson/Rayleigh Scattering



- Elastic scattering: almost no energy transfer to material
- Thomson scattering:
 Photon scattering on free electron
 - Low-energy limit of Compton scattering
- Rayleigh scattering:
 Photon scattering off an entire atom
 - Scattering cross section $\sigma_{Rayleigh} \sim f^4$
- Reason for blue / red coloration of the sky depending on the zenith angle



7000

- Noon: short path through atmosphere, hardly any blue light scattered
- Morning/evening: long path through atmosphere, much blue light scattered

Photons: Pair Production



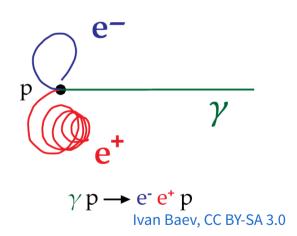
- Pair production is the generation of an electron-positron pair by the photon
- Pair production occurs in the field of a partner, which absorbs the recoil (atomic nucleus, but also shell electron)

$$\gamma + p = e^+ + e^- + p$$

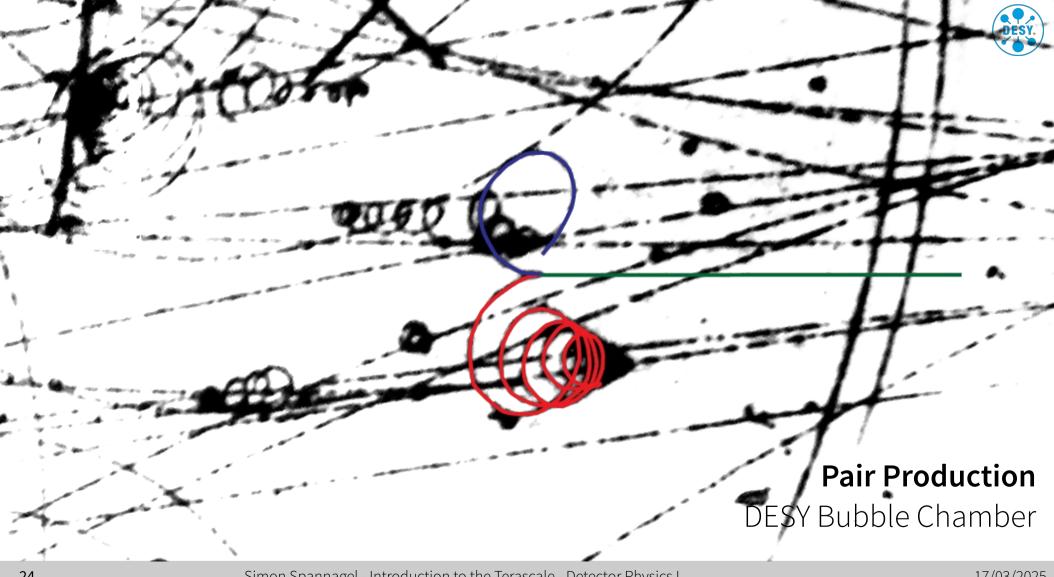
 Photon must provide at least rest mass of e+e- pair plus recoil energy:

$$E_{\gamma} \geq 2 m_e c^2 \left(1 + \frac{m_e}{M}\right)$$

Recoil energy can often be neglected, e.g. Ge detector:



$$\frac{m_e}{M} \approx 7.6 \cdot 10^{-6}$$

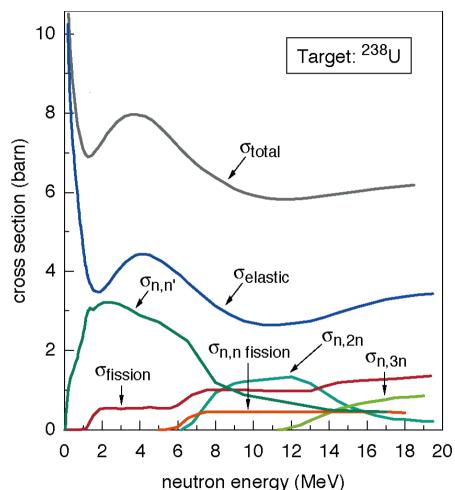


Hadronic Interaction



- Interactions of hadrons with nuclei
- Based on the strong interaction, QCD
 - Low range
 - Low probability of hadronic reactions
 - Neutrons can only interact strongly: very penetrating
- Many possible processes (energy dependent)

Elastic, inelastic scattering; neutron capture; reactions with radiation of charged particles; nuclear fission.



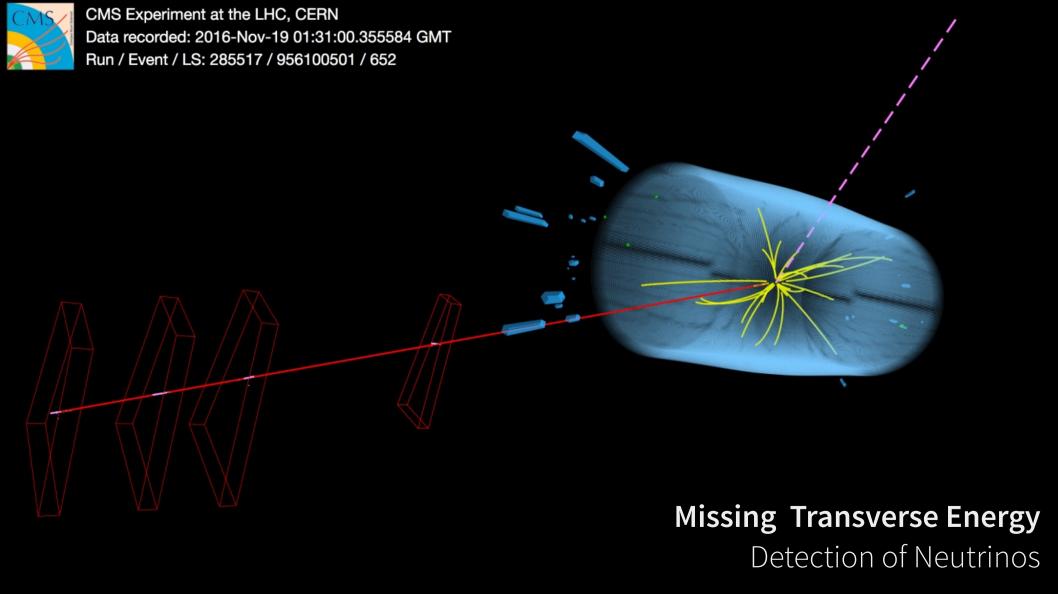
Interaction of Neutrinos



- Neutrinos are exclusively subject to the weak interaction
- Possible Interactions

$$\overline{v}_l + p \rightarrow n + l^+$$
 $v_l + n \rightarrow p + l^-$

- Neutrino interactions have very (very!) low cross sections
- Detection of neutrinos requires
 - Very large detector and high neutrino fluxes (direct detection) or
 - Hermetic detector for measurement of missing energy (indirect evidence)





Particle Detectors

Historic Overview

Geiger-Müller Counter

"Click"

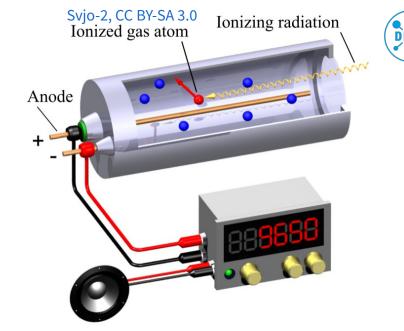
= Particle passing

"Many Click"

= Many particles

"A Great Many Click" = Run…!

- Detector filled with noble gas
 - Charged particles ionize noble gas atoms
 - High voltage applied between electrodes amplifies
- Signal: Current pulse to loudspeaker: Click!

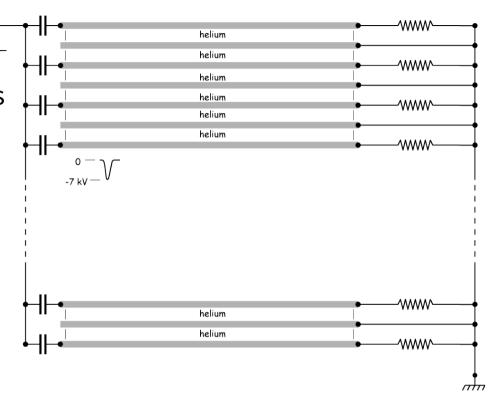




Spark Chamber



- "Many Geiger counters"
- Transparent chamber filled with inert gas and many parallel plates
 - Voltage (~ kV) between plates
 - Particles ionize noble gas atoms
 - Small sparks along the particle track
- Analysis by photos or microphones
- Relatively large dead time: quenching time of the avalanche
- Used during the 1930s-1960s





Cloud Chamber



Optical detection of charged particles

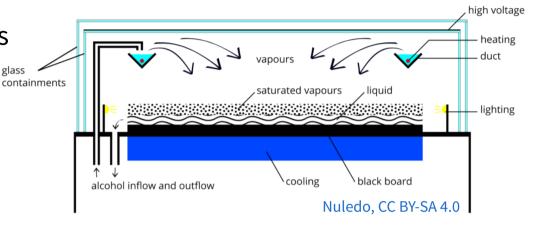
 Transparent chamber with supersaturated air-alcohol mixture

 High-energy charged particle generates ions by impact ionization

 Ions act as condensation nuclei, droplet formation in the gas mixture

Type of traces can (sometimes) be identified & assigned to particles

Nobel Prize 1927 for Charles Thomson Rees Wilson



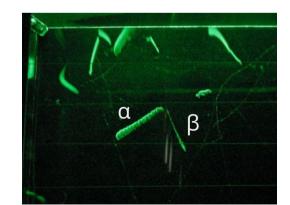






Photo emulsion / Nuclear emulsion

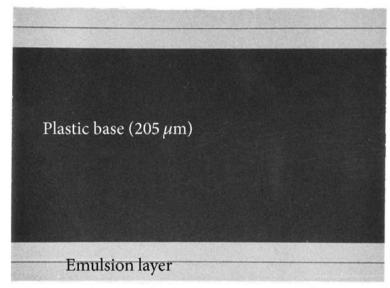


- Photographic plate with thick sensitive layer and very uniform grain size.
 - Ionizing radiation leaves traces
 - Development of the plate
 - Trajectories of particles (blackened by silver) visible with microscope

- Nobel Prize 1950 for Cecil Powell
- **OPERA** experiment in Gran Sasso

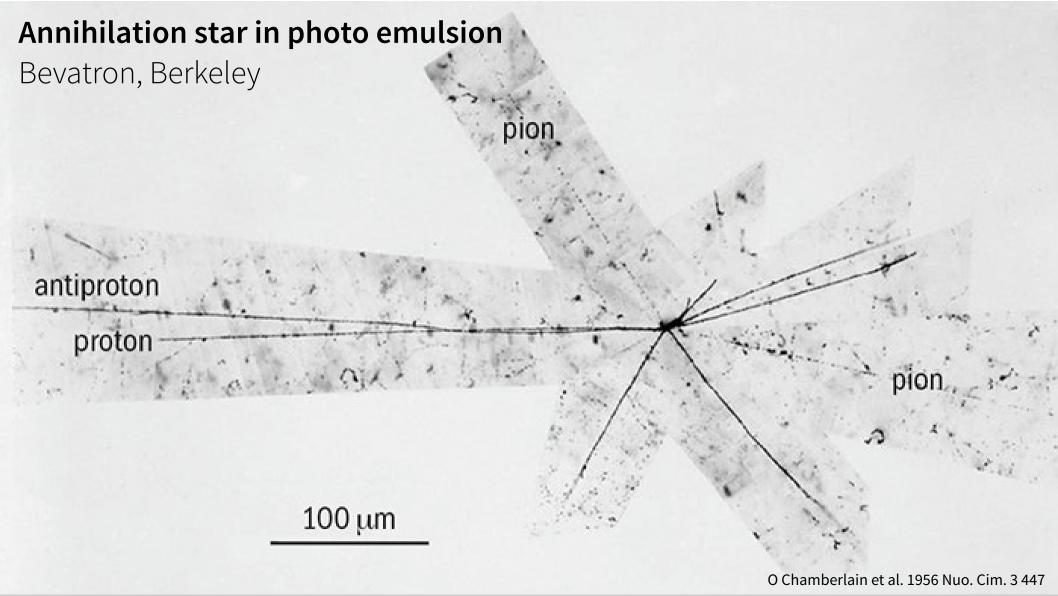






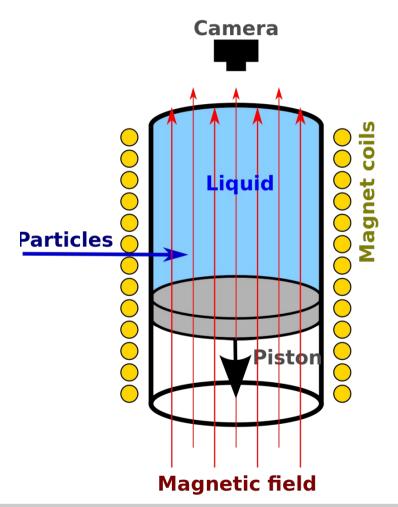
OPERA

- Photoemulsion plates for particle reconstruction
- AgBr emulsion, cooperation with Fuji Film: 9 million films, each about 10 x 12 cm



Bubble Chamber





- Optically transparent chamber filled e.g. with liquid hydrogen
 - Temperature of the liquid close to the boiling point
 - Reduction of chamber pressure with piston
 - Temperature of liquid is now above boiling point
 - Charged particles generate ions along track
 - Ions serve as nuclei for gas bubbles
- Analysis of photos of the traces



Nobel Prize 1960 for Donald A. Glaser

Prominent Bubble Chambers







- Most important particle detector type in the 1970s
- Particle source: proton synchrotron

- **Gargamelle** 1970-1978
 - 4.8 m x 1.88 m, 12 000 l
 - First detection of the Z boson
- **BEBC** 1971-1984
 - Big European Bubble Chamber
 - 3.5 T magnetic field of superconducting coils
 - Discovery of the D meson

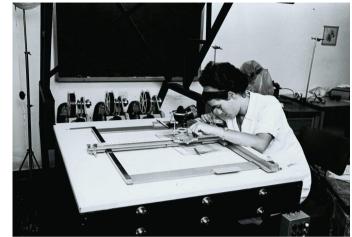


Problems of Early Particle Detectors

- Many based on photographic images
 - Data evaluation complex and only possible by hand
 - Limitation of the amount of data

- Long detector dead times
 - No further measurement possible until gas avalanche has been quenched/bubbles have disappeared
 - Only low particle rates possible
- Only little information in measurement
 - Information about particle location (& momentum), but no energy/time measurement





Weizmann Institute Archives





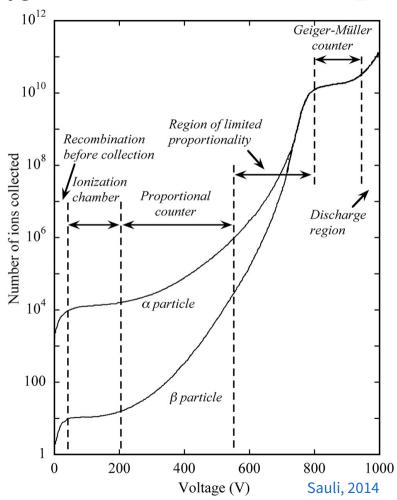
Gaseous Detectors

Primary & secondary ionization of gas atoms

Gaseous Detectors: Operating Principle

DESY.

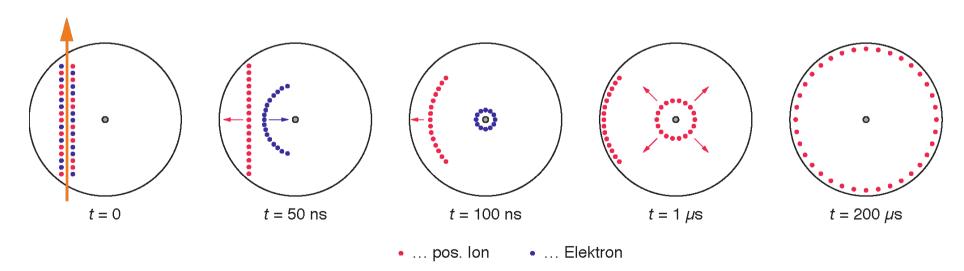
- Primary signal: charged particle generates electron-ion pairs by ionization
 - Noble gases: relatively low ionization energy
 - Average energy to generate a pair ~30 eV
 - Number proportional to deposited energy
- Amplification: different working ranges depending on applied voltage
 - Medium voltages: proportional amplification
 - High voltages: Avalanche formation due to secondary ionizations



Proportional Counter

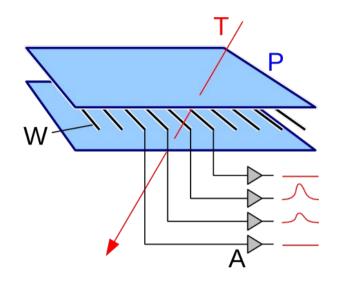


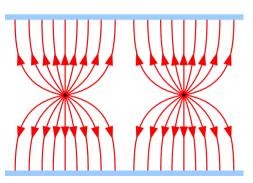
- Very similar to the Geiger counter: anode in the form of thin wire
- High field near wire leads to electron multiplication / signal amplification
- Choice of voltage: proportional range
 - Output signal proportional to original number of ionizations



Multi Wire Proportional Chamber (MWPC)







- Essentially many proportional counters next to each other, without separating walls
- Wires spaced a few millimeters apart
 - Good spatial resolution of a traversing particle
 - Large areas possible
 - Electronic selection
- High rates possible: 1000 particles/s for comparison, bubble chamber: 1-2 particles/s

Nobel Prize 1992 for Georges Charpak

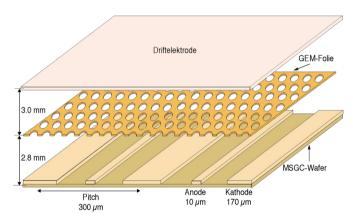


Micropattern Gas Detectors



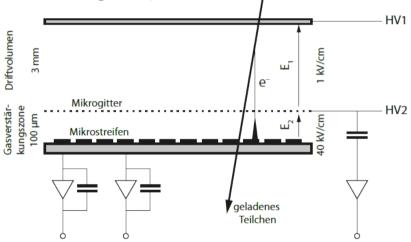
Replacement of fragile wires by micro structures

Potentially better spatial resolution and applicable for higher particle rate



Gas Electron Multiplier (GEM)

- Perforated, metalized Kapton foil, High voltage between electrodes
- Strong dipole field in perforation holes: Gas amplification

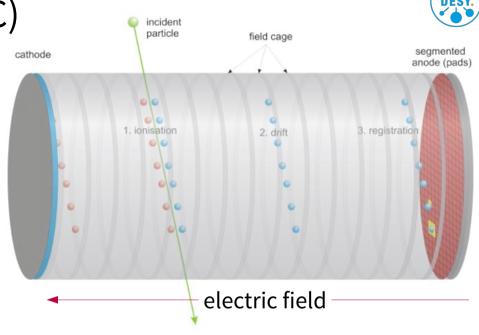


Micro-Mesh Gas Detectors (Micromegas)

- Metallic micro-grid
- Electron avalanche evolution near the lattice: Gas amplification

Time Projection Chambers (TPC)

- Large gas detector system
- Ionization along the particle track
 - Electrons and ions drifting in the E-field
 - Segmented anode: 2D information
 - Measurement of drift time:3D information
- Readout at anode side e.g. via multi-wire proportional chamber, GEMs, ...





Summary



- Particle detection
 - Measurement is performed by interaction of the particle to be measured with detector material
 - There is no single detector concept that can detect all particle types / properties
 - One needs several detector technologies and concepts
- Interaction of radiation and matter
 - Ionization and excitation of detector atoms, Bremsstrahlung, Cherenkov radiation
 - Photoelectric effect, Compton effect, pair production,
 - Hadronic interaction, missing energy, ...

- Historical detectors
 - Cloud chamber, photo emulsion, bubble chamber, spark chamber
- Gas detectors
 - Operating principle
 - Proportional counters, multi-wire proportional chamber
 - Micropattern gas detectors
 - Time Projection Chamber

