## HERMES: HERA MEASUREMENT OF SPIN The Spin of the Nucleon from HERMES point of view

Polina Kravchenko **Universität Erlangen-Nürnberg** 



**GPDs** 

Introduction

HERMES results

Conclusions

Experimental overview

longitudinal nucleon structure

$$S_{z} = \frac{1}{2} = \frac{1}{2}\Delta\Sigma$$
Naïve Parton Model  
SU<sub>spin</sub>(2) × SU<sub>flavor</sub>(3)
$$p\uparrow = \frac{1}{\sqrt{18}} \{u\uparrow u\uparrow d\downarrow \rangle - |u\uparrow u\downarrow d\uparrow \rangle - |u\downarrow u\uparrow d\uparrow \rangle + (u \Leftrightarrow d)\}$$

$$\Delta\Sigma = \Delta u + \Delta d = 1$$

$$\Delta d = \langle v\uparrow | N_{u\uparrow} - N_{u\downarrow} | v\uparrow \rangle = \frac{3}{18}(10-2) = \frac{4}{3}$$

$$\Delta d = \langle v\uparrow | N_{u\uparrow} - N_{u\downarrow} | v\uparrow \rangle = \frac{3}{18}(2-4) = -\frac{1}{3}$$

$$S_z = \frac{1}{2} = \frac{1}{2}\Delta\Sigma$$

#### In 1988 EMC measured:

$$\Gamma_{1}^{p} = \int_{0}^{1} g_{1}(x) dx = \frac{1}{2} \left\{ \frac{4}{9} \left( \Delta u + \Delta \overline{u} \right) + \frac{1}{9} \left( \Delta d + \Delta \overline{d} \right) + \frac{1}{9} \left( \Delta s + \Delta \overline{s} \right) \right\} = 0.114 \pm 0.012$$







## Measurements of $F_2$ from HERA $\longrightarrow$ Gluons are important!

## **Axial anomaly contribution:**

The contribution of the quark spins  $\Delta \Sigma$  is NOT an <u>observable</u>. The observable is  $\mathbf{a}_0$ , the flavour-singlet axial vector.

$$\Delta q + \Delta \overline{q}$$

$$\Delta q + \Delta \overline{q} - \frac{\alpha_s}{2\pi} \Delta G$$

 $\gamma_{\mu}\gamma_{5}$ 



# For full description the knowledge of orbital angular momentum is needed





Clean processes – no proton remnants
Separation light/heavy flavors
No flavor/charge separation
Access to flavor singlet D<sub>Σ</sub>
Not precise at large z
Three-jet events qqgs
Gluon FF



Flavor/charge separation
Larger z, smaller Q<sup>2</sup>
Unpolarized PDFs very well
constrained from DIS
Gluon in DIS - a small NLO effect (PGF)
Dependence on PDFs



Very sensitive to D<sub>g</sub>
 Large z
 Charge separation
 Several subprocesses
 Different p<sub>⊥</sub> scales

## **Principle of measurements**



- Virtual photon can only couple to quarks of opposite helicity
- Select q<sup>-</sup>(x) or q<sup>+</sup>(x) by changing the orientation of target nucleon spin or helicity of incident lepton beam

#### **Spin Independent Structure Function F**<sub>1</sub>

$$\sigma_{1/2} + \sigma_{3/2} \propto F_1(x) = \frac{1}{2} \sum_{i} e^2 (q^+(x) + q^-(x))$$

**Spin Dependent Structure Function** g<sub>1</sub>

$$\sigma_{1/2} - \sigma_{3/2} \propto g_1(x) = \frac{1}{2} \sum_{f} e^2 (q^+(x) - q^-(x))$$

#### Virtual photon asymmetry:

$$A_{1} = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} \approx \frac{g_{1}(x, Q^{2})}{F_{1}(x, Q^{2})}$$

# **Experimental essentials**

## asymmetry measurement

#### •Longitudinally polarized beam: electron, positron

•**Polarized target:** hydrogen, deuterium, helium3

## •<u>Measurement with high accuracy:</u> incoming lepton energy

direction scattered lepton energy direction good identification in semi-inclusive case good particle identification

### •<u>Control of false asymmetries:</u> beam flux, target size, detector size, detector efficiency



N - number of DIS events Pb, Pt - beam and target polarizations f - target dilution factor=polarizable N/total N D – depolarization factor (polarization transfer from polarized lepton to photon)



## **HERMES:** longitudinally polarized lepton beam





## **HERA positron beam properties:** • E<sub>e</sub>=27.6GeV, I<sub>e</sub><50mA, P<sub>b</sub>=0.55

- Ifetime=12-14h
- transversely polarized e± in storage ring
- oplarization build-up by emission
- of synchrotron radiation (Sokolov-Ternov effect)
- Spin rotators around HERMES IP

## **HERMES:** gas target



## **Typical HERMES target properties :**

- $P_t \sim 0.85$ ; polarized  $H^{\uparrow \rightarrow}$ ,  $D^{\rightarrow}$
- dilution factor=1
- Thickness = $10^{14}$ - $10^{15}$  nucl/cm<sup>2</sup>
- Temperature=100K



## **HERMES** spectrometer



Tracking: Drift Vertex Chambers, Front Chambers, Magnet Chambers, Back Chambers

## Particle Identification:

Čerenkov (RICH) Detector, Transition Radiation Detector, Preshower, Calorimeter Luminosity Monitor (Bhabha/Møller scattering)

- Deep-inelastic scattering (DIS) plays major role in understanding of nucleon structure
- Lepton-nucleon scattering **cleanest way** to probe substructure of nucleon
- Exchange of virtual boson, breakup and hadronization in DIS regime



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## cross section

$$\sigma(l+N \rightarrow l'+h+X)$$

## Factorization:

$$\sigma^{eN \to ehX} = \sum_{q} DF^{N \to q}(x) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z)$$



## cross section

 $\sigma(l+N \rightarrow l' + h + X)$ 

## Factorization:



Selection of constituents according to their distribution within the nucleon P DF

Photon hard scattering off the nucleon's constituents

Hadronization of the struck parton into the final hadron

## cross section

 $\sigma(l+N \rightarrow l' + h + X)$ 

## Factorization:

$$\sigma^{eN \to ehX} = \sum_{q} DF^{N \to q}(x) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z)$$

DF(x,Q<sup>2</sup>): Parton Distribution Functions q(x,Q<sup>2</sup>), Δq(x,Q<sup>2</sup>), δq(x,Q<sup>2</sup>)
σ: the hard-scattering cross section la→l/b (perturbative theory)
FF(z,Q<sup>2</sup>): Fragmentation Functions D<sub>1</sub>(z,Q<sup>2</sup>), H<sub>1</sub><sup>⊥</sup>(z,Q<sup>2</sup>),...





$$f_1^q(x) = q(x) \quad g_1^q(x) = \Delta q(x)$$

## Inclusive DIS: polarized structure function g1

$$g_1(x,Q^2) = \frac{1}{1-\frac{y}{2}-\frac{1}{4}y^2\gamma} \left[\frac{Q^4}{8\pi\alpha^2 y} \frac{\partial^2\sigma_{unpol}}{\partial x\partial Q^2} A_{||}(x,Q^2) + \frac{y}{2}\gamma^2 g_2(x,Q^2)\right]^{\text{formula}}$$



Observable: inclusive double-spin asymmetry

$$g_1(x) = \frac{1}{2} \sum_q e_q^2 (\Delta q + \Delta \bar{q})$$



## **Inclusive DIS:** first moment $\Gamma_1$ and $\Delta\Sigma$

 $\Gamma_1^d = \int dx g_1^d(x)$ 

Use only deuteron data! Assuming saturation of  $\Gamma_1^{d:}$ 



 $a_0 \stackrel{MS}{=} \Delta \Sigma = 0.330 \pm 0.025(exp) \pm 0.011(theory) \pm 0.028(evol)$ 

#### Method (LO):

Use correlation between detected hadron and struck quark  $\rightarrow$  'Flavor separation'

## flavour separation

$$\sigma(l+N \rightarrow l'+h+X)$$

Observable: semi-inclusive double-spin asymmetry

$$\begin{split} A_1^h(x,Q^2) &\stackrel{LO}{\sim} \frac{\sum_q e_q^2 \Delta q(x,Q^2) \int dz D_q^h(z,Q^2)}{\sum_q e_q^2 q(x,Q^2) \int dz D_q^h(z,Q^2)} \\ & \sim \sum_q \frac{e_q^2 q(x) \int dz D_q^h(z)}{\sum_{q'} e_{q'}^2 q' \int dz D_{q'}^h(z)} \cdot \frac{\Delta q(x)}{q(x)} \\ & \mathbf{P_q h} \qquad \mathbf{\Delta}/\mathbf{Q} \end{split}$$
Purity  $\mathbf{P_q}^h$  (x,z) is conditional probability

that hadron of type h in the final state originated from a struck quark of flavour q



e'(E')

## flavour separation

$$(A_{1p}^{\pi^+}, A_{1p}^{\pi^-}, \dots A_{1d}^{K^-}) \stackrel{\rightarrow}{\overrightarrow{A}} = P \cdot \overrightarrow{Q} \quad \left(\frac{\Delta u}{u}(x), \frac{\Delta d}{d}(x), \dots \frac{\Delta s}{s}(x)\right)$$

#### SU(3)<sub>f</sub> symmetry implicitly assumed



**u** quark large positive polarization

• d quark negative polarization

sea quarks compatible with zero in measured x-range (0.023-0.6):

 $\Delta \bar{u} = -0.002 \pm 0.043$ 

$$\int \Delta \bar{d} = -0.054 \pm 0.035$$

$$\Delta s = +0.028 \pm 0.034$$



## Direct measurement of **AG**

Mechanism: photon-gluon fusion (PGF). Observable: asymmetry in the hadron production.

 $\sigma(l + N \rightarrow h + X)$ 

- golden channel: charm production
- theoretically very clean
- experimentally very challenging



● at HERMES (√s=7 GeV hadron production at high P<sub>T</sub> experimentally very clean

highly model dependent due to variety of background processes



## transverse structure

Transverse-momentum-dependent (TMD) DF

$$\sigma(l+N \rightarrow l'+h+X)$$

$$\sigma^{eN \to ehX} \propto \sum_{q} DF^{N \to q}(x, p_T) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z, p_T)$$

Observable: azimuthal asymmetries in SIDIS.



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Transverse-momentum-dependent (TMD) DF

$$\sigma(l+N{\rightarrow}l^{'}+h+X)$$

$$\sigma^{eN \to ehX} \propto \sum_{q} DF^{N \to q}(x, p_T) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z, p_T)$$

## The SIDIS cross section up to twist-3

$$\begin{aligned} d\sigma &= \ d\sigma_{UU}^{0} + \cos(2\phi)d\sigma_{UU}^{1} + \frac{1}{Q}\cos(\phi)d\sigma_{UU}^{2} + P_{l}\frac{1}{Q}\sin(\phi)d\sigma_{LU}^{3} \\ &+ S_{L}\Big[\sin(2\phi)d\sigma_{UL}^{4} + \frac{1}{Q}\sin(\phi)d\sigma_{UL}^{5} + P_{L}(d\sigma_{LL}^{6} + \frac{1}{Q}\sin(\phi)d\sigma_{LL}^{7})\Big] \\ &+ S_{T}\Big[\sin(\phi - \phi_{s})d\sigma_{UT}^{8} + \sin(\phi + \phi_{s})d\sigma_{UT}^{9} + \sin(3\phi - \phi_{s})d\sigma_{UT}^{10} + \frac{1}{Q}\sin(2\phi - \phi_{s})d\sigma_{UT}^{11} + \frac{1}{Q}\sin(\phi_{s})d\sigma_{UT}^{12} + \frac{1}{Q}\cos(2\phi - \phi_{s})d\sigma_{LT}^{15})\Big] \\ &+ P_{l}\Big(\cos(\phi - \phi_{s})d\sigma_{LT}^{1} 3 + \frac{1}{Q}\cos(\phi_{s})d\sigma_{LT}^{14} + \frac{1}{Q}\cos(2\phi - \phi_{s})d\sigma_{LT}^{15})\Big] \end{aligned}$$



Distribution Functions (DF)						
N/q	U	L		Т		
U	$f_1$			$\boldsymbol{h}_{1}^{\perp}$		
L		<b>g</b> <sub>1</sub>		$h_{1L}^{\perp}$		
Т	$f_{IT}^{\perp}$	$g_{1T}^{\perp}$		$\boldsymbol{h}_{l}, \boldsymbol{h}_{lT}^{\perp}$		
Fragmentation Functions (FF)						
N/q	U	U		Т		
U	D	<b>D</b> <sub>1</sub>		$H_{I}^{\perp}$		

## transverse structure

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$$\sigma(l+N \rightarrow l'+h+X)$$

$$\sigma^{eN \to ehX} \propto \sum_{q} DF^{N \to q}(x, p_T) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z, p_T)$$

$$\sum_{x \in h_1(x, p_T^2)} \frac{\text{Collins effect}}{\otimes H_1^{\perp}(z, k_T^2)}$$

correlation between parton transverse polarization in a transversely polarized nucleon and transverse momentum of the produced hadron

$$\begin{bmatrix} \phi \end{bmatrix} d\sigma_{UU}^2 + P_l \frac{1}{Q} sin(\phi) d\sigma_{LU}^3 \\ + \left( d\sigma_{LL}^6 + \frac{1}{Q} sin(\phi) d\sigma_{LL}^7 \right) \end{bmatrix}$$

$$+ S_T \left[ sin(\phi - \phi_s) d\sigma_{UT}^8 + sin(\phi + \phi_s) d\sigma_{UT}^9 + sin(3\phi - \phi_s) d\sigma_{UT}^{10} + \frac{1}{Q} sin(2\phi - \phi_s) d\sigma_{UT}^{11} + \frac{1}{Q} sin(\phi_s) d\sigma_{UT}^{12} +$$

$$P_l \Big( \cos(\phi - \phi_s) d\sigma_{LT}^1 3 + \frac{1}{Q} \cos(\phi_s) d\sigma_{LT}^{14} + \frac{1}{Q} \cos(2\phi - \phi_s) d\sigma_{LT}^{15} \Big)$$





+ S

## transverse structure

Transverse-momentum-dependent (TMD) DF

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$$\sigma^{eN \to ehX} \propto \sum_{q} DF^{N \to q}(x, p_T) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z, p_T)$$

## Sivers effect $\propto f_{1T}^{\perp}(x, p_T^2) \otimes \overline{D_1(z, k_T^2)}$

correlation between parton transverse momentum and nucleon transverse polarization

 $d\sigma_{LU}^3$ 

requires orbital angular momentum

$$d\sigma = d\sigma_{UU}^{0} + \cos(2\phi) d\sigma_{U}^{1}$$

$$requires orbital angular momentum
$$J\sigma_{LU}^{3}$$

$$S_{L}\left[sin(2\phi) d\sigma_{UL}^{4} + \frac{1}{q}sin(e^{-\frac{1}{q}sin(e^{-\frac{1}{q}sin(\phi + \phi_{s})}d\sigma_{UT}^{9} + sin(3\phi - \phi_{s})}d\sigma_{UT}^{10} + \frac{1}{Q}sin(2\phi - \phi_{s}) d\sigma_{UT}^{11} + \frac{1}{Q}sin(\phi_{s}) d\sigma_{UT}^{12} + \frac{1}{Q}cos(\phi - \phi_{s}) d\sigma_{LT}^{13} + \frac{1}{Q}cos(\phi_{s}) d\sigma_{LT}^{14} + \frac{1}{Q}cos(2\phi - \phi_{s}) d\sigma_{LT}^{15}\right)$$$$



Distribution Functions (DF)						
N/q	U	L	Т			
U	$f_{I}$		$h_1^{\perp}$			
L		<b>g</b> <sub>1</sub>	$h_{1L}^{\perp}$			
Т	$f_{IT}^{\perp}$	$g_{1T}^{\perp}$	$h_{l}, h_{lT}^{\perp}$			
Fragmentation Functions (FF)						
N/q	U	J	Т			
U		1	$H_I^{\perp}$			

+



# **Collins** moments

$$\sigma(l+N \rightarrow l'+h+X)$$

# TMDs can be studied by measuring azimuthal asymmetries in SIDIS

- $\pi^+ > 0, \pi^- < 0$
- π<sup>-</sup> unexpectedly large!
- large *unfavoured* Collins fragmentation function!



K+ consistent with π+
K- (all sea object) opposite sign from π-



#### 

# Sivers moments

$$\sigma(l+N \rightarrow l'+h+X)$$

- significantly positive π<sup>+</sup>
   moment
- ➡first evidence of a non zero naïve T-odd DF in DIS
- ➡requires non-zero orbital angular momentum
- $\pi$  consistent with zero



K<sup>+</sup> amplitude larger
 than for π<sup>+</sup>
 (Sivers for sea quarks)
 K<sup>-</sup> consistent with zero



## What do we learn from Collins moment?



$$\sigma(l+N \rightarrow l'+h+X)$$

# First extraction of transversity distribution!!









## Probing the orbital angular momentum

Accessing Generalized Parton Distributions (GPDs) via Deeply Virtual Compton Scattering (DVCS) and exclusive meson production



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## Sensitive to $J_q$

### A. Airapetian et al., JHEP05 (2011)126

A. Airapetian et al., JHEP08 (2010) 130



#### A. Airapetian et al., PLB 666, 446 (2008)]

A.Airapetian et al., JHEP06 (2008) 066



# Backup slides

- Deep-inelastic scattering (DIS) plays major role in understanding of nucleon structure
- Lepton-nucleon scattering cleanest way to probe substructure of nucleon
- Exchange of virtual boson, breakup and hadronization in DIS regime



## cross section

 $\sigma(l+N{\rightarrow}l^{'}+X)$ 

Assuming one photon exchange

$$\frac{d^2\sigma}{d\Omega dE'} = \frac{\alpha^2}{MQ^4} \frac{E}{E'} L_{\mu\nu} W_{\mu\nu} W_{\mu\nu}$$
hadronic tensor  
contains information  
about hadron structure  
leptonic tensor  
from QED

## cross section

Assuming one photon exchange

 $\frac{d^2\sigma}{d\Omega dE'} = \frac{\alpha^2}{MQ^4} \frac{E}{E'} \frac{L_{\mu\nu}}{W\mu\nu}$ 

$$\sigma(l+N{\rightarrow}l'+X)$$

hadronic tensor contains information about hadron structure

leptonic tensor from QED

unpolarized structure functions momentum distribution of quarks

$$W^{\mu\nu} = -\left(g^{\mu\nu} - \frac{q^{\mu}q^{\nu}}{q^{2}}\right)F_{1}\left(x,Q^{2}\right) + \left(P^{\mu} - \frac{P \times q}{q^{2}}q^{\mu}\right)\left(P^{\nu} - \frac{P \times q}{q^{2}}q^{\nu}\right)F_{2}\left(x,Q^{2}\right) + iM\epsilon^{\mu\nu\rho\sigma}q_{\rho}\left[\frac{S_{\sigma}}{P \times q}g_{1}\left(x,Q^{2}\right) + \frac{S_{\sigma}\left(P \cdot q\right) - P_{\sigma}\left(S \cdot q\right)}{\left(P \times q\right)^{2}}g_{2}\left(x,Q^{2}\right)}\right]$$

polarized structure functions spin distribution of quarks

## Inclusive DIS: unpolarized structure function F<sub>2</sub>



 $\frac{\text{measured}}{dxdQ^2} = \frac{4\pi\alpha_{em}^2}{Q^4} \frac{F_2}{x} \times \left[1 - y - \frac{Q^2}{4E^2} + \frac{y^2 + Q^2/E^2}{2(1 + R(x, Q^2))}\right]$ 



#### in the 1-photon exchange approximation

$$F_1(x) = \frac{1}{2} \sum_q e_q^2 q(x)$$
$$F_2(x) = x \sum_q e_q^2 q(x)$$



## Inclusive DIS: unpolarized structure function F<sub>2</sub>







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## Inclusive DIS: $\Delta q$ and $\Delta G$

 $g_1^{NLO}(x,Q^2) = g_1^{LO} + \frac{1}{2}e^2 \sum_{q} \left[ \Delta q(x,Q^2) \otimes S_q + \Delta g(x,Q^2) \otimes C_g \right]^{he}$ 

SU(3)<sub>f</sub> symmetry implicitly assumed





•valence quarks are well determined  $\Delta u_v > 0$  and  $\Delta d_v < 0$ 

•**gluons** and **sea** quarks are poorly constraint by data



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#### Method (LO):

Use correlation between detected hadron and struck quark → **'Flavor separation'** 

# flavour separation

$$\sigma(l+N \rightarrow l'+h+X)$$

$$\sigma^{eN \to ehX} = \sum_{q} DF^{N \to q}(x) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z)$$

Probability that struck quark of flavour q fragments into hadron of type h with energy fraction  $z=E_h/v$ 

$$\begin{split} A_1^h(x,Q^2) &\sim \frac{\sum_q e_q^2 \Delta q(x,Q^2) \int dz D_q^h(z,Q^2)}{\sum_q e_q^2 q(x,Q^2) \int dz D_q^h(z,Q^2)} \\ &\sim \sum_q \frac{e_q^2 q(x) \int dz D_q^h(z)}{\sum_{q'} e_{q'}^2 q' \int dz D_{q'}^h(z)} \cdot \frac{\Delta q(x)}{q(x)} \end{split}$$

**Purity**  $P_q^h(x,z)$  is conditional probability that hadron of type h in the final state originated from a struck quark of flavour q





## transverse structure

 $\sigma(l+N \rightarrow l'+h+X)$ Transverse-momentum-dependent (TMD) DF

$$\sigma^{eN \to ehX} \propto \sum DF^{N \to q}(x, p_T) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z, p_T)$$

$$\sum_{q} DF^{N \to q}(x, p_T) \otimes \sigma^{eq \to eq} \otimes FF^{q \to h}(z, p_T)$$



$$d\sigma = d\sigma_{UU}^{0} + \cos(2\phi)d\sigma_{UU}^{1} + \frac{1}{Q}\cos(\phi)d\sigma_{UU}^{2} + P_{l}\frac{1}{Q}\sin(\phi)d\sigma_{LU}^{3}$$

$$+ S_{L}\left[sin(2\phi)d\sigma_{UL}^{4} + \frac{1}{O}sin(\phi)d\sigma_{UL}^{5} + P_{L}(d\sigma_{LL}^{6} + \frac{1}{O}sin(\phi)d\sigma_{LL}^{7})\right]$$

$$+ S_{T}\left[sin(\phi - \phi_{s})d\sigma_{l}^{8}\right] \frac{Boer-Mulders Effect}{\int Q} \sum_{l=1}^{2} \frac{1}{Q}sin(2\phi - \phi_{s})}{P_{l}\left(cos(\phi - \phi_{s})\right)} \sum_{l=1}^{2} \frac{Boer-Mulders Effect}{\int Q} \sum_{l=1}^{2} \frac{1}{Q}sin(2\phi - \phi_{s})}{\int Q} \sum_{l=1}^{2} \frac{1}{Q}sin(2\phi - \phi_{s})}$$

