

# Analytic bootstrap for defect CFTs

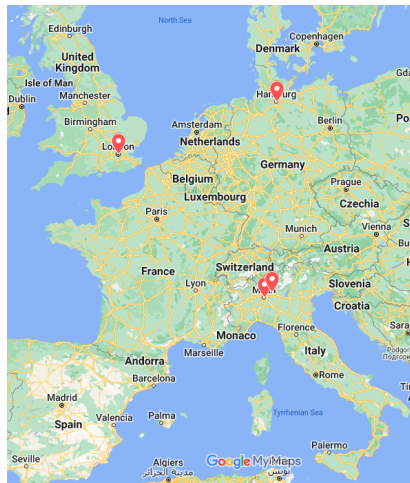
Davide Bonomi

DESY



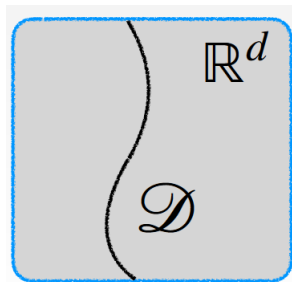
# About me

- Name: Davide Bonomi.
- Born in Bergamo, Italy.
- Bachelor's and Master's in Milan.
- PhD in London.
- Hobbies: reading, trekking,...



# My research

Main focus: Conformal Field Theories (CFTs) that contain extended objects ([defects](#)).

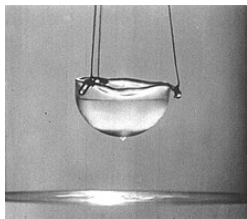


Main strategy: use conformal symmetry + properties of the Operator Product Expansion (OPE) to constrain or solve CFTs ([conformal bootstrap](#)).

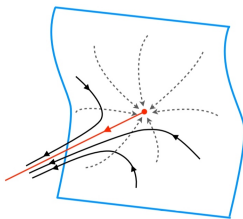
# Why CFT?

Conformal invariance  $\sim$  Poincaré + scale invariance.

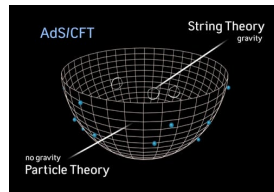
Conformal field theories are ubiquitous in physics:



Phase transitions



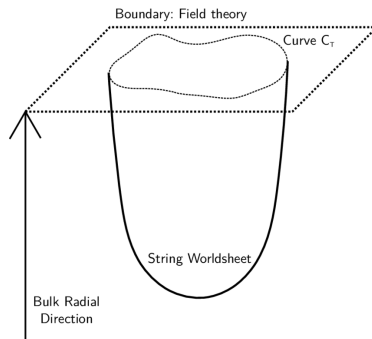
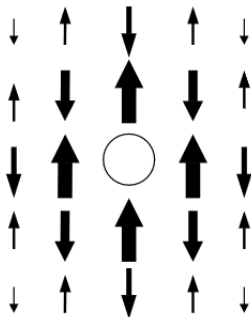
RG flows



AdS/CFT

# Why defects?

Many CFTs admit interesting defects:



Phase transitions with impurities  
(e.g. metal with external atom)

Wilson lines in CFTs  
 $\sim$   
Strings and branes in AdS

# Conformal Bootstrap

2pt and 3pt fixed up to **scaling dimensions** and **OPE coefficients**:

$$\langle \mathcal{O}_i(x) \mathcal{O}_j(0) \rangle = \frac{\delta_{ij}}{x^{2\Delta_{\mathcal{O}}}}, \quad \langle \mathcal{O}_i \mathcal{O}_j \mathcal{O}_k \rangle \sim \lambda_{ijk}$$

All other correlators can be computed using the OPE:

$$\mathcal{O}_i(x) \mathcal{O}_j(0) \stackrel{\text{OPE}}{=} \sum_k \frac{\lambda_{ijk} \mathcal{O}_k(0)}{|x|^{\Delta_i + \Delta_j - \Delta_k}} + \dots$$

Different OPEs  $\rightarrow$  crossing equation  $\rightarrow$  constraints on  $\{\Delta_i, \lambda_{ijk}\}$

$$\sum_i \lambda_{12\phi} \langle \phi(x_1) \phi(x_2) \phi(x_3) \rangle = \sum_j \lambda_{23\phi} \langle \phi(x_1) \phi(x_2) \phi(x_4) \rangle$$

The crossing eq is positive definite and can be studied numerically.

# Defect bootstrap

A  $p$ -dimensional defect breaks the bulk conformal group as

$$\underbrace{SO(d+1,1)}_{\mathcal{O}_{\Delta,\ell}} \rightarrow \underbrace{SO(p+1,1) \times SO(d-p)}_{\hat{\mathcal{O}}_{\hat{\Delta},\hat{\ell},s}}$$

Bulk and defect OPE  $\rightarrow$  crossing equation  $\rightarrow \left\{ \hat{\Delta}_i, \hat{\lambda}_{ijk}, a_i, b_{ij} \right\}$

$$\sum_{\mathcal{O}} \phi \begin{array}{c} \diagup \\ \diagdown \end{array} \begin{array}{c} \mathcal{O} \\ a \end{array} \Big| = \sum_{\hat{\mathcal{O}}} \phi \begin{array}{c} \text{---} b \\ \text{---} b \end{array} \Big| \hat{\mathcal{O}}$$

Problem: not positive definite  $\rightarrow$  no numerical bootstrap.

# Analytic bootstrap

The OPE controls the **singularities** of correlators. We can use this input to reconstruct correlators and extract CFT data. No need of a positive definite expansion.

E.g.: consider a function  $f(z)$  such that

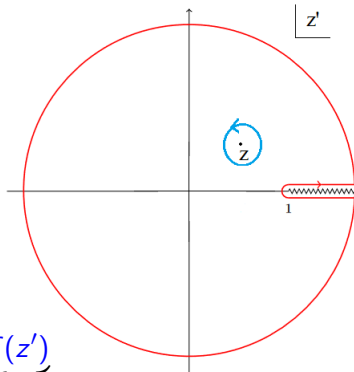
- $f(z)$  has a **branch cut** for  $z > 1$ .
- $|f(z)/z| \rightarrow 0$  as  $|z| \rightarrow \infty$ .

Cauchy's theorem implies:

$$f(z) = \frac{1}{2\pi i} \oint \frac{dz'}{z' - z} f(z')$$

By deforming the contour:

$$f(z) = \frac{1}{2\pi i} \int_1^\infty \frac{dz'}{z' - z} \underbrace{\text{Disc} f(z')}_{f(z' + i\epsilon) - f(z' - i\epsilon)}$$





# Example: defects in the $O(N)$ model

We studied the critical  $O(N)$  model

$$S = \int d^d x \left[ \frac{1}{2} (\partial_\mu \phi_i)^2 + \frac{\lambda}{4!} (\phi_i \phi_i)^2 \right]$$

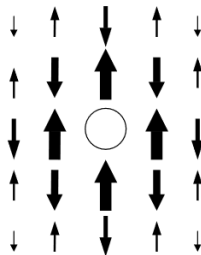
in presence of defects in  $d = 4 - \varepsilon$  with  $\varepsilon \ll 1$ .

- Localized magnetic field

$$D = e^{-h_0 \int d\tau \phi_1(\tau)}$$

- Spin impurity

$$D_j = \text{Tr}_{2j+1} \left( P e^{\frac{\zeta_0}{\sqrt{\kappa}} \int d\tau \phi^i T^i} \right), \quad T^i \in \mathfrak{su}(2)$$



Main result:  $\langle \phi \phi \rangle$  and  $\infty$  new defect CFT data from few bulk data.

# Ongoing work and future directions

- Combine analytic bootstrap with numerical bulk data.
- Correlators of defects (e.g. correlator of two impurities).
- Holographic defects and interplay with integrability, localization ect.

Thank you for your attention!