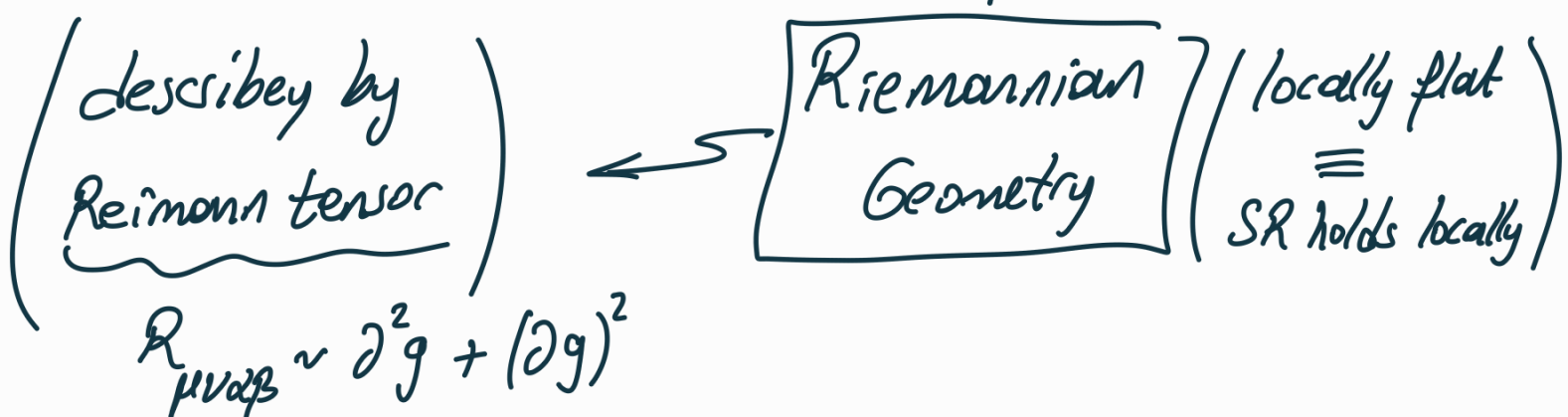
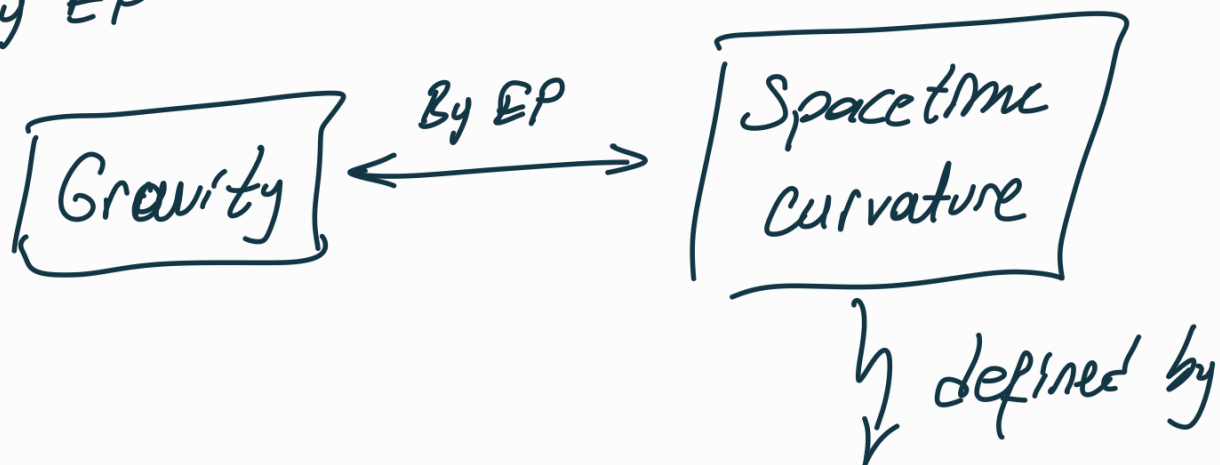
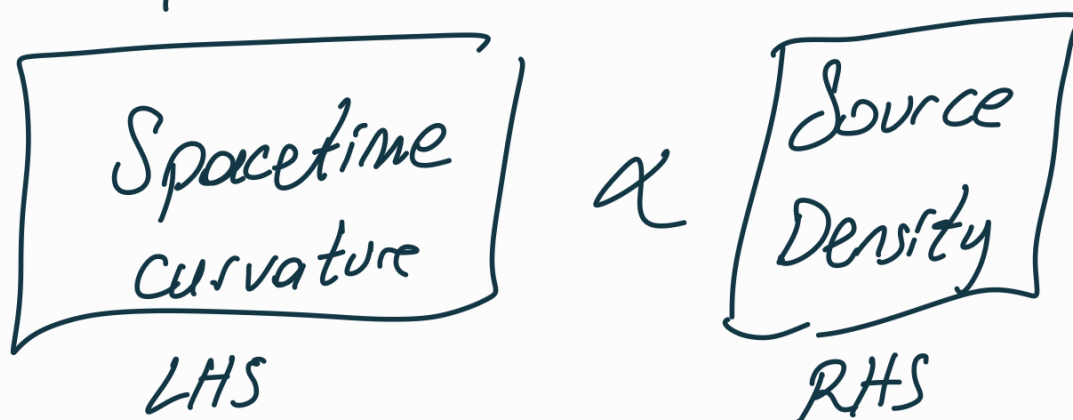


Heuristic Derivation of EFE:

* By EP



* Also from Poisson eqn.



$$\left. \begin{array}{l} E^2 - c^2 p^2 = m^2 c^4 \\ \text{In SR} \end{array} \right\} \begin{array}{l} \text{mass} \\ \text{energy} \\ \text{mom.} \end{array} \Bigg| \text{density}$$

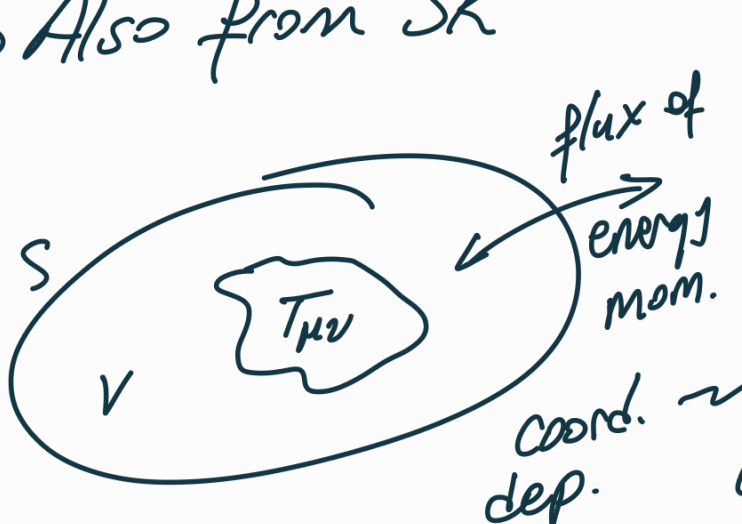
* So mass density ρ in Poisson eqn.
 should be generalized to some object
 including all these: mass/energy/mom.

$\Rightarrow$ we already have such a quantity
 in SR.

$$\left(\begin{array}{c} 2^{\text{nd}}\text{-rank} \\ \text{Tensor} \end{array} \right) T_{\mu\nu} = \left[\begin{array}{c|c} \text{energy density} & \text{energy flux density} \\ \hline \text{mom. density} & \text{mom. flux density (stress tensor)} \end{array} \right]$$

$\hookrightarrow$ EM tensor

$\Rightarrow$ Also from SR



flux of energy mom.

coord. dep. $\rightarrow$

(like $\partial_\mu j^\mu = 0$)

$$\boxed{\partial_\mu T^{\mu\nu} = 0}$$

Conservation of EM.

$\Rightarrow$ In GR, we should have general coords.

$$\left(\begin{array}{l} \eta_{\mu\nu} \rightarrow g_{\mu\nu} \\ \partial_\mu \rightarrow \nabla_\mu = \partial_\mu + \Gamma^\cdot_{\cdot\mu} \end{array} \right) \Rightarrow \boxed{\nabla_\mu T^{\mu\nu} = 0}$$

general covariance

$\Downarrow$

Covariantly const.

() principle of general covariance

* Then we have the properties of $T_{\mu\nu}$:

$$(T_{\mu\nu}) \Rightarrow \left\{ \begin{array}{l} - \text{vanishes in the absence of "matter"} \\ - \text{is a 2}^{\text{nd}}\text{-rank tensor} \\ - \text{is sym; } T_{\mu\nu} = T_{\nu\mu} \\ - \text{is cov. const.} \end{array} \right.$$

* All these suggest that

$$\underbrace{F[R_{\mu\nu\alpha\beta}(g_{\mu\nu})]}_{\text{spacetime curv. part.}} \propto \underbrace{F[T_{\mu\nu}]}_{\text{source density part}} \quad \left(\begin{array}{c} \text{Structure} \\ \text{of} \\ \text{GR} \\ \text{eqns.} \end{array} \right)$$

* Let's try some possibilities:

(1) The simplest possibility is

$$g_{\mu\nu} \propto T_{\mu\nu} ?$$

- satisfies $\nabla_\mu T^{\mu\nu} = 0$ since $\nabla_\mu g_{\mu\beta} = 0$

- gives $g_{\mu\nu} = 0$ in vacuum ($T_{\mu\nu} = 0$)

$\Rightarrow$ does not work! (trivial)

(ii) Another possibility is

$$R_{\mu\nu\alpha\beta} \propto T_{\mu\nu} T_{\alpha\beta} \quad ?$$

— does not satisfy $\nabla_\mu T^{\mu\nu} = 0$

— gives flat spacetime ($g_{\mu\nu} = \eta_{\mu\nu}$)
in vacuum ($T_{\mu\nu} = 0$)

$\Rightarrow$ So, does not work!

(iii) One last guess is

$$R_{\mu\nu} \propto T_{\mu\nu} \quad ? \quad R_{\mu\nu} \equiv R^\alpha{}_\alpha \mu_\nu$$

— does not satisfy $\nabla_\mu T^{\mu\nu} = 0$

$$\text{since } \nabla_\mu R^{\mu\nu} = \frac{1}{2} \nabla^\nu R \neq 0$$

$\Rightarrow$ So does not work! — Vacuum case is OK!
 $R_{\mu\nu} = 0$

* With all these possibilities, we should also note that

$$\left(\begin{array}{l} \text{Poisson} \\ \text{Eqn.} \end{array} \right) \quad \nabla^2 \phi = \underbrace{\partial_i \partial_i}_{\text{Divergence appears}} \phi = 4\pi G \rho \quad \sim (\text{in NG})$$

Divergence appears

$$\left(\begin{array}{l} \text{Since GR should} \\ \text{recover NG} \end{array} \right) \Rightarrow \frac{\text{GR}}{g_{\mu\nu}} \longleftrightarrow \frac{\text{NG}}{\phi}$$

$$\Rightarrow \text{LHS} \equiv F[\partial_i \partial_i g_{\mu\nu}]$$

in GR eqn.

$\Rightarrow$ Then we have

LHS

$$F[R_{\mu\nu}(g_{\mu\nu})] \propto F[T_{\mu\nu}]$$

* After these reasoning, there is one particular combination of $R_{\mu\nu\alpha\beta}$ which satisfies all the properties of RHS:

$$G_{\mu\nu} \equiv R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \kappa T_{\mu\nu} \quad (\text{EFE})$$

κ = proportionality const
(determining gravitational coupling)

— You can indeed show that

$$\nabla_{\mu} G^{\mu\nu} = 0 \rightarrow \text{Bianchi identity}$$

Exercise

* You can also add CC to LHS

$$G_{\mu\nu} + \underline{\underline{\Lambda g_{\mu\nu}}} = \kappa \bar{T}_{\mu\nu}$$

$\Downarrow$ can also be derived
from action principle

$$S = \int \sqrt{-g} R d^4x$$

Einstein-Hilbert action