a program for automatic computation of multiparticle cross sections

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Higher luminosity of the colliders \Rightarrow higher precision of the measurements.

SM radiative corrections should be included, at least the leading ones.

Top quark pair production at ILC

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- As the top and antitop decay before they hadronize, predominantly into bW^+ and $\bar{b}W^-$
- and W's decay into $f\bar{f}'$ -pairs,
- \Rightarrow 6 fermion reactions of the form

$$e^+e^- \rightarrow bf_1\bar{f}_1'\bar{b}f_2\bar{f}_2',$$

where $f_1, f'_2 = v_e, v_\mu, v_\tau, u, c$ and $f'_1, f_2 = e^-, \mu^-, \tau^-, d, s$, should be studied with a complete set of the Feynman diagrams, possibly including radiative corrections.

Feynman diagrams of $e^+e^- \rightarrow bf_1\bar{f}'_1\bar{b}f_2\bar{f}'_2$



(a) 'signal', (b) and (c) 'background'.

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(a) 'signal', (b) and (c) 'background'.

- $e^+e^- \rightarrow b\nu_{\mu}\mu^+\bar{b}\mu^-\bar{\nu}_{\mu} \Rightarrow 452 \text{ diagrams}$
- $e^+e^- \rightarrow bud\bar{b}d\bar{u} \Rightarrow 1484$ diagrams

in the unitary gauge of SM, neglecting the Higgs boson couplings to fermions lighter than b.



 $e^+e^- \rightarrow t\bar{t} \rightarrow bf_1\bar{f}_1'\bar{b}f_2\bar{f}_2'$



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K. Kołodziej, A. Staroń, A. Lorca, T. Riemann, Eur. Phys. J. C 46 (2006) 357.

EW radiative corrections

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- (a) calculated with a subroutine based on topfit, by T. Riemann, et al.,
- (b)–(e) with subroutines written by A. Staroń.

$e^+e^- \rightarrow t\bar{t}H$

If the Higgs boson has mass below the $t\bar{t}$ threshold, $m_H < 2m_t$, then $t\bar{t}H$ Yukawa coupling can be best determined from this reaction [Djouadi, Kalinowski, Zerwas]. If the Higgs boson has mass below the $t\bar{t}$ threshold, $m_H < 2m_t$, then $t\bar{t}H$ Yukawa coupling can be best determined from this reaction [Djouadi, Kalinowski, Zerwas].

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Diagrams (a) i (b) dominate \Rightarrow the cross section $\sim g_{ttH}^2$.



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56550, or neglecting the Higgs boson coupling to light fermions, 26816 Feynman diagrams \Rightarrow A strong need for a complete automation of the cross section computation. Do it with carlomat! • a program written in Fortran 90/95

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- generates the matrix element for a user specified process

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- fermion masses are not neglected
- the maximum number of external particles is 12
- only the Standard Model is implemented at the moment

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How does carlomat work?

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They are copied to another directory where the numerical program is executed.

Generation of topologies

Start with 4 topologies of a 4 particle process.



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1 \\
2 \\
4
\end{array}$ $\begin{array}{c}
1 \\
3 \\
3
\end{array}$ $\begin{array}{c}
2 \\
3 \\
4
\end{array}$ $\begin{array}{c}
1 \\
2 \\
3
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Attach line 5 to each line, including the internal ones, and to each triple vertex \Rightarrow 25 topologies of a 5 particle process.

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1 & 2 \\
4 & 3
\end{array}$ $\begin{array}{c}
1 & 2 \\
3 & 4
\end{array}$ $\begin{array}{c}
1 & 2 \\
4 & 3
\end{array}$ $\begin{array}{c}
2 \\
3 & 4
\end{array}$ $\begin{array}{c}
2 \\
4 & 1
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Attach line 5 to each line, including the internal ones, and to each triple vertex \Rightarrow 25 topologies of a 5 particle process.

6 particles \Rightarrow 220 topologies

- 7 particles \Rightarrow 2485 topologies
- 8 particles \Rightarrow 34300 topologies
- 9 particles \Rightarrow 559405 topologies
- 10 particles \Rightarrow 10525900 topologies
- 11 particles \Rightarrow 224449225 topologies

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For a process with *n* external particles, topologies for n - 1 particles are generated recursively and then, while adding the *n*-th particle, the program checks whether a topology results in a Feynman diagram or not.

Topologies can be generated and stored on a disk prior to the program execution.

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Each topology is divided into two parts which are separately checked against the Feynman rules. Actual external particles are assigned to lines 1, 2, 3, ..., n in a strict order.



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Particles defined for one Feynman diagram are used as building blocks of other Feynman diagrams.



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Use is made of the routines developed for a Monte Carlo program eett6f v. 1.0, for calculating lowest order cross sections of $e^+e^- \rightarrow 6$ fermions relevant for a $t\bar{t}$ -pair production and decay K. Kołodziej, Comp. Phys. Commun. 151 (2003) 339, hep-ph/0210199; A. Biernacik, K. Kołodziej, Nucl. Phys. B (Proc. Suppl) 116 (2003) 33, hep-ph/0210405. which have been tailored to meet needs of the automatic generation of amplitudes.

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polarization vectors representing particles, both on- and off-shell ones, are computed only once, for all the helicities of the external particles they are made of, and stored in arrays – a novel feature with respect to e.g. MadGraph.

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Mappings of

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- $\sim 1/s$ behaviour of the photon and gluon propagators are performed.

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be used as MC generator of unweighted events as well.

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- a few reactions with 7 particles

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- reactions with 4 fermions and a photon in the final state have been checked against ee4fgamma
- several reactions with 6 particles in the final state have been checked against HELAC/PHEGAS and AMEGIC++ [T. Gleisberg, et al., hep-ph/0311273, EPJ C34 (2004) 173] and against results of eett6f.

Agreement, within one standard deviation.

Tests

Cross sections of $e^+e^- \rightarrow b\bar{b}b\bar{b}u\bar{d}\mu^-\bar{\nu}_{\mu}$ without QCD contributions have been checked against Whizard. Agreement, within one standard deviation.

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Cross sections of $e^+e^- \rightarrow b\bar{b}b\bar{b}u\bar{d}\mu^-\bar{\nu}_{\mu}$ without QCD contributions have been checked against Whizard. *Agreement, within one standard deviation.* Further thourogh comparisons with other MC generators as HELAC, MadEvent, Whizard, Sherpa, Alpgen, CompHEP should be done.

Cross sections of $e^+e^- \rightarrow b\bar{b}b\bar{b}u\bar{d}\mu^-\bar{\nu}_{\mu}$

\sqrt{s} [GeV]	σ _{all} [ab]	σ _{no QCD} [ab]	σ _{signal} [ab]	σno cuts signal	$\sigma_{NWA}^{no cuts}$ [ab]
500	26.8(4)	7.80(3)	3.095(3)	3.796(3)	3.920(1)
800	100.2(8)	66.8(1)	46.27(2)	58.36(2)	60.03(2)
1000	93.1(3)	61.4(1)	40.18(2)	51.74(2)	52.42(3)
2000	47.4(2)	28.5(1)	15.14(3)	22.14(4)	20.68(3)

 $5^{\circ} < \theta(q, \text{beam}), \ \theta(l, \text{beam}) < 175^{\circ}, \ \theta(q, q'), \ \theta(l, q) > 10^{\circ},$ $E_q, \ E_l, \not E^T > 15 \text{ GeV}$

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Large off resonance background!

[KK, S. Szczypiński]

Cross sections [ab] at $\sqrt{s} = 500 \text{ GeV of } e^+e^-ra$

$m_{bb}^{\rm cut}$	bbbbbudsc		$bar{b}bar{b}uar{d}\mu^-ar{ u}_\mu$		$bar{b}bar{b} au^+ u_ au \mu^- ar{ u}_\mu$	
[GeV]	σ_{all}	$\sigma_{sig.}$	σ_{all}	$\sigma_{sig.}$	σ_{all}	$\sigma_{sig.}$
20	13.88(6)	8.70(2)	3.50(2)	2.38	4.03(9)	0.86
5	10.17(4)	8.66(2)	2.62(1)	2.33	1.89(7)	0.86
1	9.07(4)	8.65(1)	2.37(1)	2.31	1.09(2)	0.86

 $60 \text{ GeV} < m(\sim b_1, \sim b_2) < 90 \text{ GeV}, \quad m(\mu, \not\!\!\!E^T) < 90 \text{ GeV}$ $|m(b, \sim b_1, \sim b_2) - m_t| < 30 \text{ GeV},$ $m_t - 30 \text{ GeV} < m_T(b, \mu, \not\!\!\!E^T) < m_t + 10 \text{ GeV},$ $|m(b_3, b_4) - m_H| < m_{bb}^{\text{cut}}$

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Background substantially reduced! [KK, S. Szczypiński]

Summary and Outlook

Automatic generation of matrix elements and phase space parametrizations in carlomat seems to work quite well.
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be worked on.

Extensions of SM can be implemented and the corresponding lowest order cross sections can be calculated in a fully automatic way.

Leading SM radiative corrections can be implemented manually.

Helicity amplitudes

The constant widths of unstable particles are introduced through the complex mass parameters:

$$M_V^2 = m_V^2 - im_V \Gamma_V, \qquad V = W, Z,$$

 $M_H^2 = m_H^2 - im_H \Gamma_H, \qquad M_t = \sqrt{m_t^2 - im_t \Gamma_t},$

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which replace masses in the corresponding propagators, both in the *s*- and *t*-channel Feynman diagrams,

$$\Delta_F^{\mu\nu}(q) = \frac{-g^{\mu\nu} + q^{\mu}q^{\nu}/M_V^2}{q^2 - M_V^2},$$

$$\Delta_F(q) = \frac{1}{q^2 - M_H^2}, \qquad S_F(q) = \frac{q + M_t}{q^2 - M_t^2}.$$

Propagators of a photon and a gluon are taken in the Feynman gauge.

carlomat - p.2

Helicity amplitudes

The EW mixing parameter may be defined either real (FWS) or complex (CMS),

$$\sin^2 \theta_W = 1 - \frac{m_W^2}{m_Z^2}$$
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 or $\sin^2 \theta_W = 1 - \frac{M_W^2}{M_Z^2}$.

Colour matrix is calculated only once at the beginning of execution of the numerical program after having reduced its size with the use of the SU(3) algebra properties.