# Theory prediction for $t\bar{t}$ threshold enhancement at hadron colliders

Maria V. Garzelli mostly on the basis of M.V.G., G. Limatola, S.-O. Moch, M. Steinhauser, O. Zenaiev, [arXiv:2412.16685]

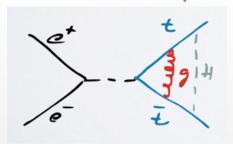
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FH Particle Physics Discussion: a spotlight on the  $t\bar{t}$  threshold at the LHC DESY, January 17th, 2024

## $e^+e^- ightarrow tar{t}$ close to threshold at lepton colliders



- \* Top quark non relativistic close to threshold
- $(v_t/c = \beta_t = \sqrt{1 4m_t^2/M_{t\bar{t}}^2} <<< 1)$  in  $t\bar{t}$  rest frame. Regime for NRQCD. Plenty of studies, due to sensitivity to  $m_t$ .
- \* Simultaneous sensitivity to  $\Gamma_t$ ,  $\alpha_s$ ,  $y_{t\bar{t}h}$ , to be considered when the process is used for  $m_t$  extraction via top-threshold scan in short-distance  $m_t$  renormalization scheme ( $\Delta m_t \sim$  some tens MeV).
- \* Even  $e^+e^- \rightarrow W^+bW^-\bar{b}$  predictions are available (non-resonant contributions to  $m_t$ ).

# $t\bar{t} + X$ production close to threshold at hadron colliders

- \* Far less studies at hadron coliders w.r.t. lepton colliders
- \* Interest suddenly enhanced by the threshold excesses seen in the results of recent CMS analyses searching for BSM  $H \to t\bar{t}$  decays (to be cross-checked by ATLAS).
- \* NRQCD predictions, including singlet and octet contributions for the  $m_{t\bar{t}}$  distribution  $\Rightarrow$  Revival of the theory chain in Kiyo, Kühn, Moch, Steinhauser and Uwer, [arXiv:0812.0919]
- \* bound-state effects related to resummation of terms depending on  $\alpha_s^n/\beta_t^m$ , associated to the exchanges of Coulomb-like virtual gluons between the t and  $\bar{t}$  quarks
- \* plus resummation of large logarithms due to real soft-gluon emissions close to threshold at NLL.

#### $M_{t\bar{t}}$ differential distribution

$$M_{t\bar{t}}\frac{\mathrm{d}\sigma_{P_1P_2\to T}}{\mathrm{d}M_{t\bar{t}}}(S,M_{t\bar{t}}^2) = \sum_{i,j} \int_{\rho}^1 \mathrm{d}\tau \left[\frac{\mathrm{d}\mathcal{L}_{ij}}{\mathrm{d}\tau}\right](\tau,\mu_f^2) \; M_{t\bar{t}}\frac{\mathrm{d}\hat{\sigma}_{ij\to T}}{\mathrm{d}M_{t\bar{t}}}(\hat{s},M_{t\bar{t}}^2,\mu_f^2)$$

with  $au=\hat{s}/S$ ,  $ho=M_{t\bar{t}}^2/S$ ,

$$\left[\frac{d\mathcal{L}_{ij}}{d\tau}\right](\tau,\mu_f^2) = \int_0^1 dx_1 \int_0^1 dx_2 f_{i/P_1}(x_1,\mu_f^2) f_{j/P_2}(x_2,\mu_f^2) \delta(\tau - x_1 x_2)$$

and

$$M_{t\bar{t}} \frac{{\rm d}\hat{\sigma}_{ij\to T}}{{\rm d}M_{t\bar{t}}} (\hat{s}, M_{t\bar{t}}^2, \mu_f^2) \ = \ F_{ij\to T} (\hat{s}, M_{t\bar{t}}^2, \mu_f^2) \ \frac{1}{m_t^2} \, {\rm Im} \, G^{[1,8]} (M_{t\bar{t}} + i\Gamma_t)$$

valid in NRQCD for the considered distribution at NLO:

- still to be generalized at higher-orders
- still to be generalized for other distributions

### Factorization of the partonic cross-section

$$M_{t\bar{t}} \frac{{\rm d}\hat{\sigma}_{ij\to T}}{{\rm d}M_{t\bar{t}}} (\hat{s}, M_{t\bar{t}}^2, \mu_f^2) \ = \ F_{ij\to T} (\hat{s}, M_{t\bar{t}}^2, \mu_f^2) \ \frac{1}{m_t^2} \, {\rm Im} \, G^{[1,8]} (M_{t\bar{t}} + i \Gamma_t)$$

valid in NRQCD for the considered distribution at NLO

- $F_{ij \to T}$ : hard-scattering function for producing the  $T = ^{2S+1} L_j^{[1,8]}$  state, containing threshold logarithms becoming large close to the partonic threshold (i.e. for  $\tau$  reaching its minimum value  $\rho$ , equivalent to  $z \to 1$ ).
- $G^{[1,8]}$ : attractive and repulsive non-relativistic Green's functions, solution of Schroedinger eq. accounting for exchange of potential gluons between t and  $\overline{t}$ . They depend on  $\Gamma_t$  and  $m_t$ .
  - $G^{[1]}$  from attractive Coulomb-like potential, responsible for "toponium" quasi-bound states  $\Gamma^t >> \Lambda_{QCD}$ : the top quark decays before a proper bound state can be formed.
  - $G^{[8]}$  from repulsive Coulomb-like potential.

### Resummation of threshold logarithms

- \* We resum threshold logarithms up to NLL for the most relevant T channels  $(gg \to^1 S_0^{[1]},^1 S_0^{[8]})$  and  $q\bar{q} \to^3 S_1^{[8]})$ , working in N-space (N conjugate to z).
- \* We consider the Mellin transform of the hard function

$$F_{ij\to T}^{N}(M_{t\bar{t}}^{2},\mu_{f}^{2}) = \int_{0}^{1} dz z^{N-1} F_{ij\to T}(\hat{s}, M_{t\bar{t}}^{2},\mu_{f}^{2})$$

and the Mellin transform  $\mathcal{L}^{\textit{N}}_{ij}$  of the luminosity function  $\mathcal{L}_{ij}( au)$ .

- \* We then recover predictions in z-space from the inverse Mellin transform of  $(\mathcal{L} \otimes \mathcal{F})^{\mathcal{N}} = \mathcal{L}^{\mathcal{N}} \mathcal{F}^{\mathcal{N}}$ , performed numerically, using the minimal prescription.
- $\ast$  This includes matching of the resummed results to the NLO ones, avoiding double-counting, leading to NLO+NLL predictions.
- \* Uncertainties related to different prescriptions for avoiding the Landau pole in the inverse Mellin transform and matching not yet included.

## Theory predictions: input

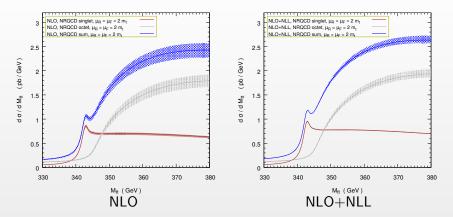
- \* Present setup:
  - NNPDF3.1 NNLO PDFs +  $\alpha_s(M_Z)$ ,
  - $m_t = 172.5 \text{ GeV}$ ,
  - $\Gamma_t$ =1.36 GeV,
  - $\mu_R = \mu_F = [m_t, 4m_t].$
- \* Ongoing study of parametric uncertainties.
- \* Possibility to provide predictions with different inputs: please feel free to ask.

### Predictions for the convolution $\mathcal{L} \otimes \mathcal{F}$

	NLO			resummed		
$gg \to {}^{1}S_0^{[1]}$	18.2	18.7	18.3	19.4	20.5	21.1
$gg\to{}^1S_0^{[8]}$	55.8	55.2	52.8	60.0	61.5	62.0
$q\bar{q}\to^3S_1^{[8]}$	21.7	22.3	22.0	22.4	22.4	22.0

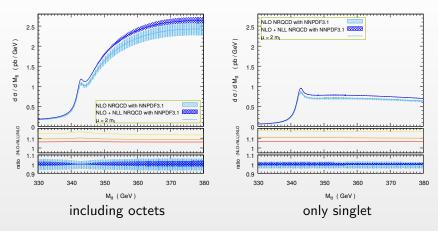
TABLE I. Comparison of the NLO and NLO+NLL resummed result of the convolution  $\mathcal{L}\otimes F$  (in units  $10^{-6}~\text{GeV}^{-2}$ ) for the LHC configuration  $\sqrt{S}=13~\text{TeV}$  with NNPDF3.1 PDFs at the reference point  $M_{t\bar{t}}=2m_t$ . The three columns correspond to the scale choices  $\mu_r=\mu_f\in\{m_t,2m_t,4m_t\}$ .

## $M_{t\bar{t}}$ distribution for $pp \to t\bar{t} + X$ close to threshold



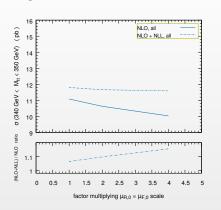
- \* Multiple singlet and octet contributions included at fixed-order: all S-wave channels, i.e. the  $gg, gq, q\bar{q} \rightarrow^1 S_0^{[1,8]}$  channels, the  $gg, gq, q\bar{q} \rightarrow^3 S_1^{[8]}$  channels, as well as the  $gg \rightarrow^3 S_1^{[1]}$  one.
- \* NLL effects enhance predictions and reduce considerably uncertainty bands w.r.t. NLO.

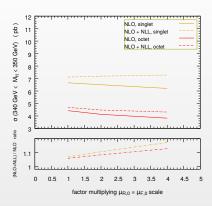
### Comparison of NLO and NLO+NLL predictions



- \* K-factors  $\sim 10\%$ , depending on the scale
- $\ast$  Uncertainties reduced in NLO+NLL w.r.t. NLO, especially in the peak region.

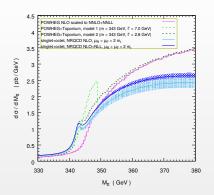
# Integrated cross-section in the region of validity of NRQCD vs. scale





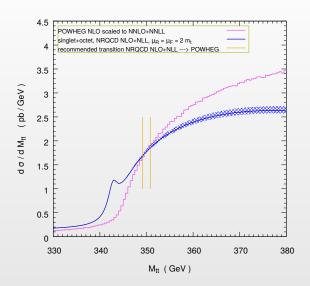
- \* K-factors  $\sim 10\%$  for both singlet and octet
- \* NLL effects increase with the value of the scale  $\mu_{R,0}=\mu_{F,0}$
- \* other scale choices under investigation

### Comparison with predictions of Fuks model 2021

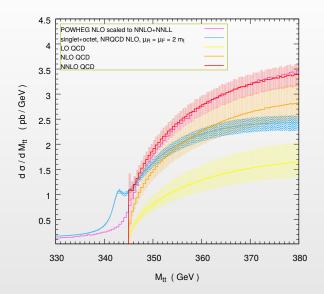


- \* POWHEG and Fuks et al. model input from Laurids et al.
- \* Height, position and width of Fuks-modelled singlet do not coincide with our NLO (+ NLL) QCD singlet + octet predictions.
- \* Big uncertainties on the Green's functions used in our NLO QCD predictions, not shown in our plots. Missing higher orders would enhance these functions.

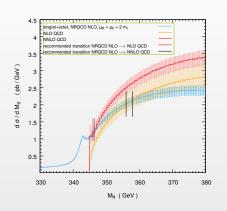
# Recommended transitions from NRQCD to QCD POWHEG predictions

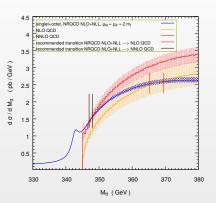


#### Reweighted Powheg vs. fixed-order predictions vs. NRQCD



# Recommended transitions from NRQCD to QCD fixed-order predictions





- \* fixed-order predictions with MATRIX at LO, NLO, NNLO, with their scale uncertainty bands.
- $\ast$  A proper matching between NRQCD and QCD predictions still to be devised

#### **Conclusions**

- \* Further studies ongoing, especially focused on uncertainties.
- \* So far, we are limited to inclusive predictions, cuts still to be applied.
- $\ast$  Investigation of how to best use our predictions: bin-by-bin reweighting ? Alternatives ?
- \* How safe/robust is extrapolating to other distributions (i.e. the distributions currently measured by the experimentalists)?
- \* Future interesting studies: compare shape of enhancements at threshold from models with BSM Higgs decaying into a  $t\bar{t}$  pair to our results, and understand up to which extent it is possible to discriminate between the effects of BSM Higgs and toponium.
- $\ast$  Please note that our threshold enhancement, corresponding to the color-singlet config., occurs mostly **below** threshold, in contrast with many BSM Higgs models.
- \* Our current predictions: lower limit to the true ones.