Status of the ATLAS Luminosity Measurements

ATLAS Collaboration Week
CERN, February 21, 2024

F. Dattola, on behalf of the Luminosity Working Group



Counting collisions: luminosity in ATLAS

Overview

- Luminosity crucial input for any cross section measurement: $\sigma(pp \to X) = N(pp \to X)/L$
- Precise determination of luminosity is therefore essential for ATLAS physics analyses
 - Leading systematic uncertainty for some measurements

$$\sigma(Z \to ll) = 751 \pm (\text{stat.}) \pm 15(\text{syst.}) \pm 17(\text{lumi.}) \text{ pb at } 13.6 \text{ TeV [ATLAS-CONF-2023-006]}$$

Record 0.83% uncertainty achieved in Run 2 → Target: keep final Run-3 uncertainty within 1%



■ ATLAS Luminosity Highlights in 2024

- Preliminary calibration for 2023 pp data [ATL-DAPR-PUB-2024-001]
- Preliminary calibration for 2023 PbPb data [Preliminary HI Result (Talk)]
- Complete Run-3 tag with current best knowledge of luminosity [ATLGBLCONDTAGS-85]
- Final calibration for 2018 pp data at 900 GeV in EdBoard [Approval Talk]
- **3 separate vdM-calibration campaigns in 2024**: pp at 13.6 TeV (May), pp at 5.4 TeV (pp-ref, October), PbPb at 5.4 TeV NN (heavy ion, November)

Luminosity detectors and algorithms

ATLAS *luminometers* in Run 3

LUCID

- Reference luminometer: usable over full μ range
- Cherenkov light detector: 2×16 PMTs
- Several algorithms: event- and hit-counting
- Sensitive on single bunch resolution

Z-Counting

• Used for independent cross checks of baseline luminosity over time and μ

Tile calorimeter

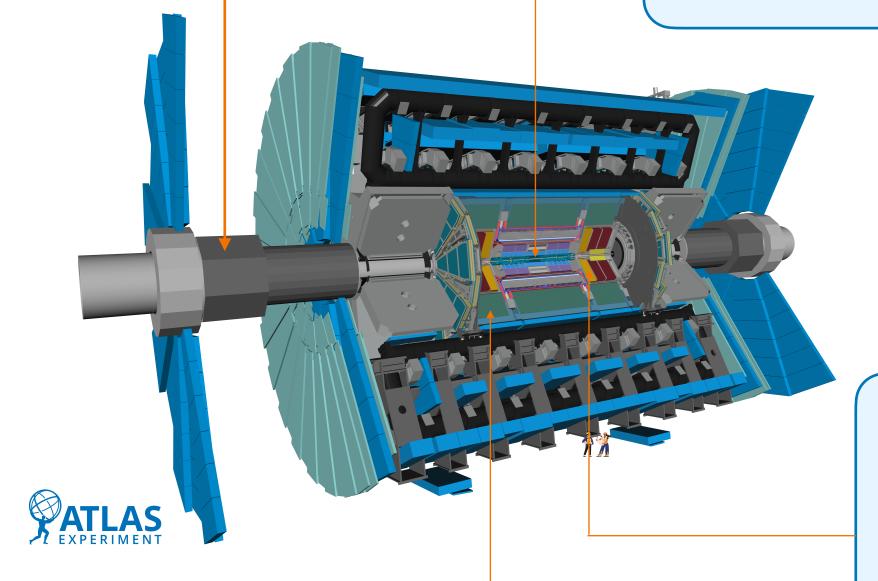
- Luminosity from PMT currents
- Usable over wide μ range
- Bunch-integrated sensitivity

Inner Detector

- **Track Counting**: count reconstructed tracks with Pixel + SCT in randomly triggered events
- Usable over wide μ range
- Very linear in μ for different track selections
- Sensitive on single bunch resolution
- Pixel Cluster Counting (in development): count reconstructed Pixel clusters in randomly triggered events

LAr calorimeters (EMEC and FCal)

- Luminosity measured with LAr gap HV currents over O(1) s integration times.
- Mostly usable in high- μ physics fills.
- Very stable over time.
- Bunch-integrated sensitivity.



Luminosity measurement strategy

Steps in luminosity determination

Basic idea:

• Measure luminosity through visible interaction rate in a luminosity-sensitive detector

$$\mathcal{L}_b = \frac{f_r n_1 n_2}{2\pi \Sigma_x \Sigma_y} = \frac{\mu \cdot f_r}{\sigma_{\text{inel}}} = \frac{\mu_{\text{vis}} \cdot f_r}{\sigma_{\text{vis}}} \leftarrow ---$$

LHC beam parameters —

- μ = number of inelastic collisions per bunch
- $\sigma_{\rm inel}$ = inelastic cross section
- \blacksquare μ_{vis} = visible interaction rate
- $\sigma_{\rm vis}$ = visible cross section

Step 1: vdM calibration

- Scan beams against each other in dedicated fills with speciallytailored LHC conditions
- Calibration of LUCID $\sigma_{
 m vis}$

Step 2: calibration transfer

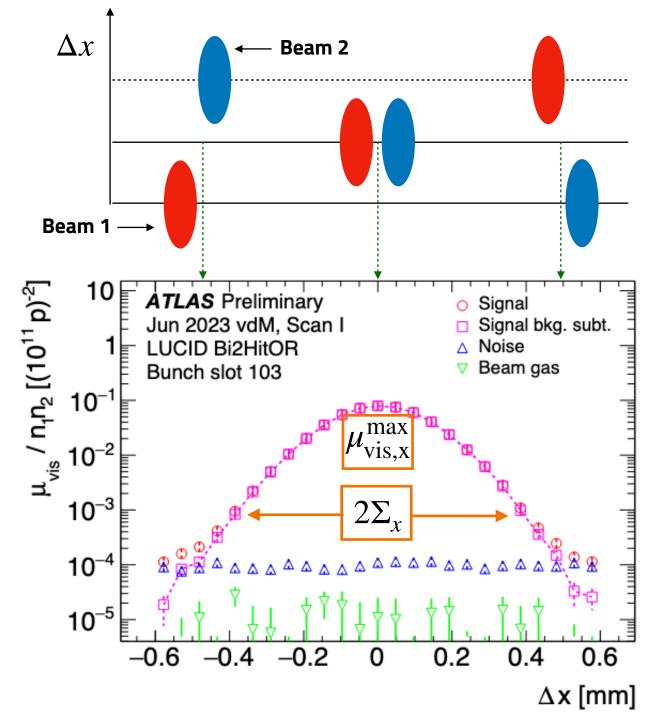
- Transfer LUCID measurement from vdM regime to physics regime
- Correct LUCID response with Track
 Counting measurement
- Cross-check with Tile measurement to assess uncertainties

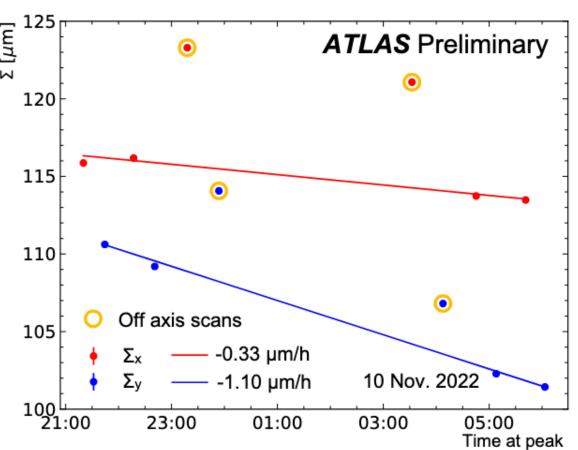
Step 3: long-term stability

- Verify stability of luminosity calibration from run to run over entire running period
- Compare run-integrated luminosities from LUCID, Tile, EMEC, FCAL

Analysis of van der Meer scans

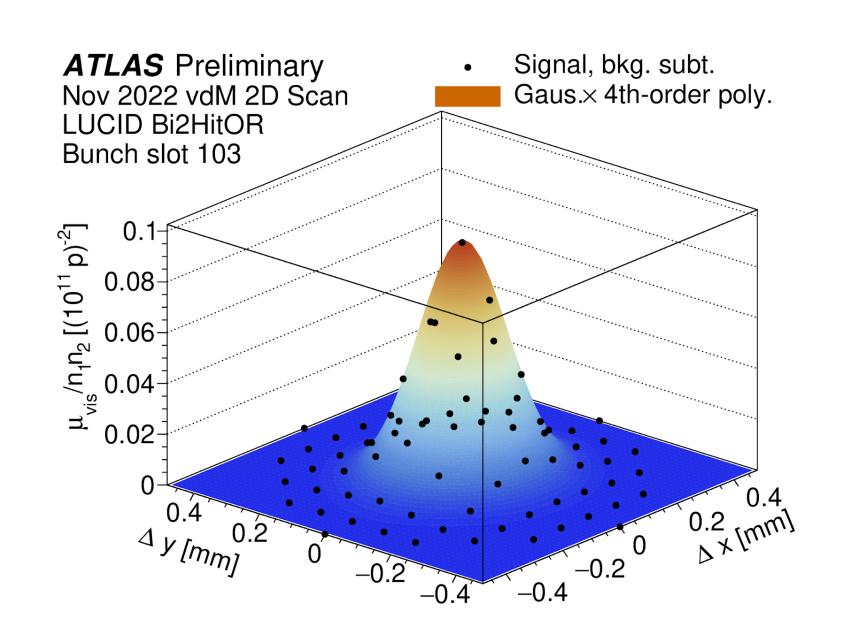
- Determine visible cross-section $\sigma_{\rm vis}=\mu_{\rm vis}^{\rm max}\cdot(2\pi\Sigma_{\rm x}\Sigma_{\rm y})/n_1n_2$ through fits of the measured interaction rate $\mu_{\rm vis}$ vs beam separation in x- and y-directions
- vdM calibration requires **specially-tailored LHC conditions**: low $\mu \sim 0.5$, isolated bunches, no crossing angle
- Many corrections needed: orbit drifts, emittance growth, non-factorisation, beam-beam effects, length scale, magnetic non-linearity
 - Non-factorisation main systematic in Run 3: 1.07% in 2022 and 1.39% in 2023 $(pp) \rightarrow$ combined 1.22% vs 0.24% in Run-2
 - Breaks factorisation of beam profiles assumed in the basic formalism
 - Transverse profile: $\rho(x,y) \neq \rho(x) \cdot \rho(y) \rightarrow [\Sigma_x \Sigma_y] \neq \Sigma_x \cdot \Sigma_y$
 - Evidence for strong non-factorisation in Run 3 from off-axis scans

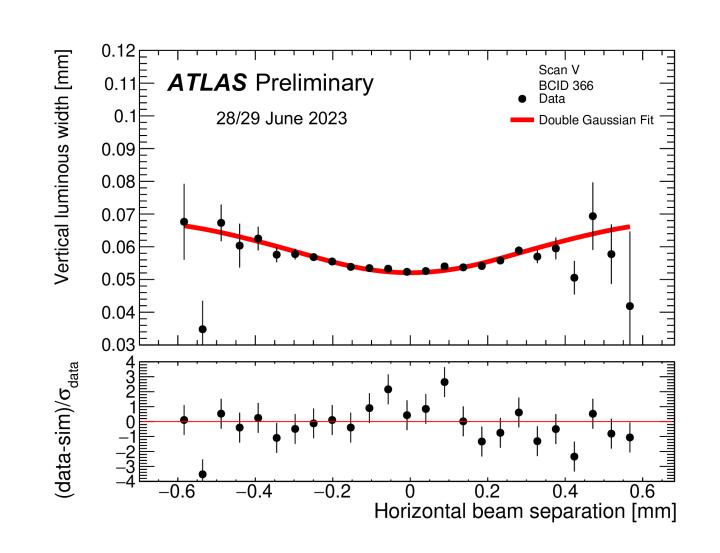




Analysis of van der Meer scans

- Two orthogonal approaches to improve control of non-factorisation
 - **LRE analysis**: combined fit of expected luminous region and luminosity to the data recorded at each scan step → model dependent
 - 2D grid scans: perform a scan over a grid in $(x, y) \rightarrow$ model independent





WE NEED YOU

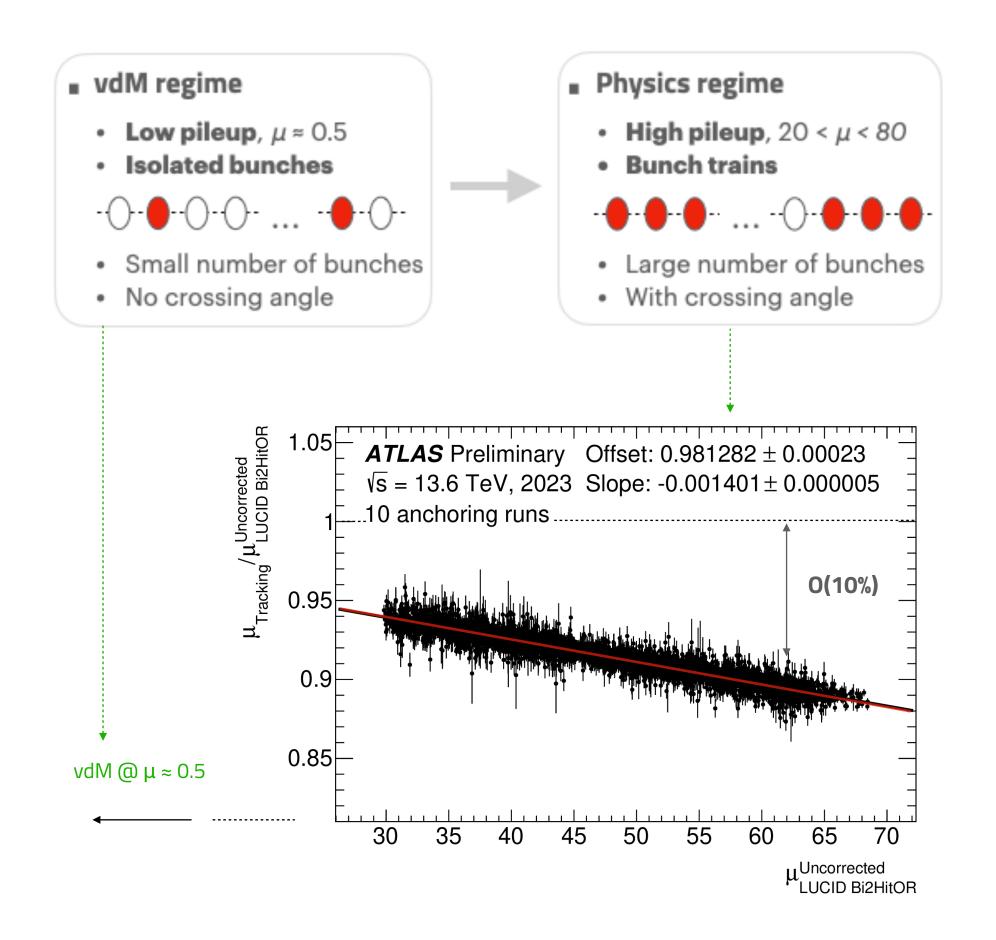
- Many vdM scans in 2024, including special scans to improve understanding of non factorisation (i.e. diagonal scans)
- Data to be analysed for the first time → not simple routine work, requires own intellectual input!

Calibration transfer

Going to high-pileup physics conditions

- Transfer LUCID vdM calibration to physics regime
 - Normalise Track Counting (TC) to LUCID in head-on parts of vdM fill and compare LUCID vs TC at high-luminosity in physics regime
 - \rightarrow LUCID response shows strong dependence on μ :
 - up to ~ 10% overestimation of luminosity at high- μ
 - \rightarrow correct using Track Counting (TC): $\langle \mu_{\rm corr} \rangle = p_0 \langle \mu_{\rm uncorr} \rangle + p_1 (\langle \mu_{\rm uncorr} \rangle)^2$
 - use a few long high-luminosity reference fills per year

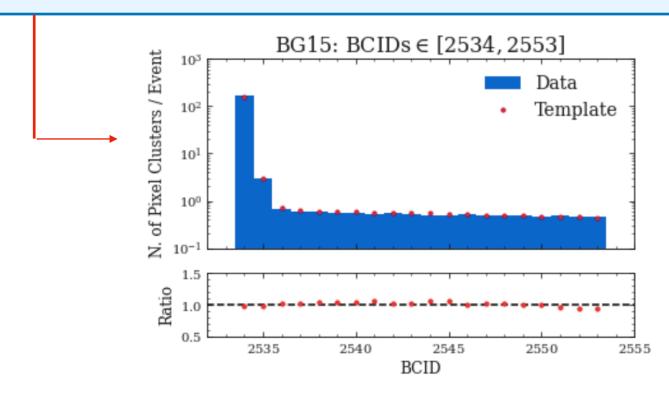
- TC stability constantly monitored with dedicated data stream
- Close collaboration with Tracking/Inner Detector experts



Intermezzo: luminosity with the Inner Detector

Not only corrections to LUCID...

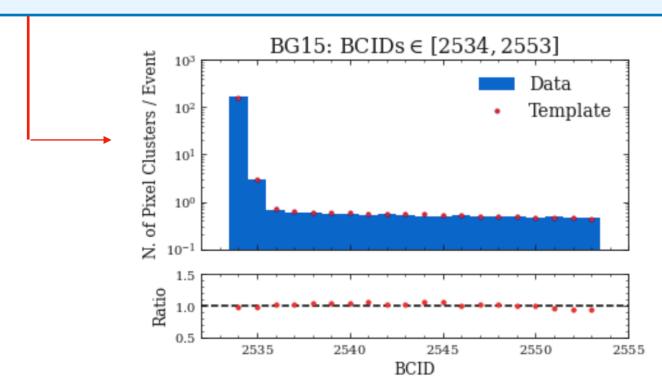
- Pixel Cluster Counting (PCC)
 - ATL-COM-DAPR-2023-028
 - **Assumption**: number of pixel clusters per bunch crossing proportional to μ
 - Already in use at CMS, could be key for Run 4
 - **Developing in Run 3**: implementation of PCC software and crucial studies for control of *afterglow* background



Intermezzo: luminosity with the Inner Detector

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 - ATL-COM-DAPR-2023-028
 - **Assumption**: number of pixel clusters per bunch crossing proportional to μ
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TC calibration in emittance scans

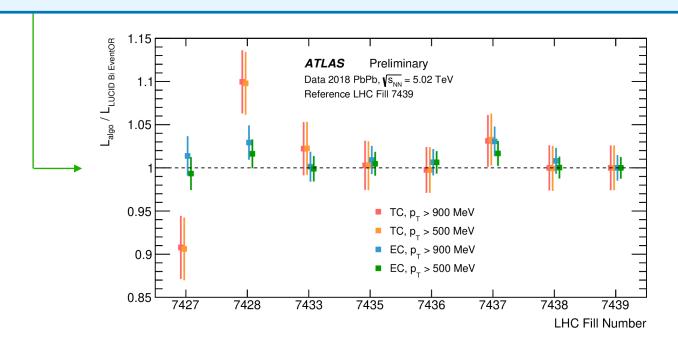
- Extraction TC $\sigma_{
 m vis}$ from emittance scans: vdM-like analysis
- Developing in Run 3: would provide TC calibration orthogonal to LUCID

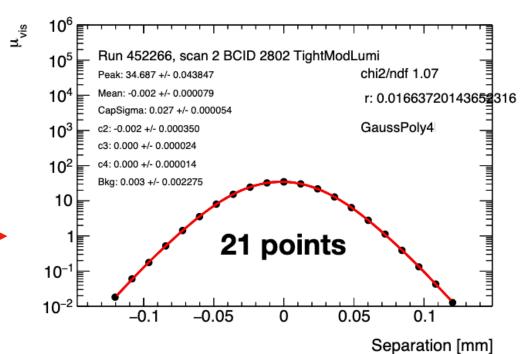
Track-based luminosity in PbPb

- ATL-COM-DAPR-2023-025
- Event Counting (EC): count number of events in a given time period with at least one track (Pixel + SCT) of a given working point

$$\mu_{\text{vis}}^{\text{EC}} = -\ln(1 - N_{\text{pass}}/N_{\text{tot}})$$

• EC is statistically precise and stable: sets longterm stability uncertainty in PbPb calibration





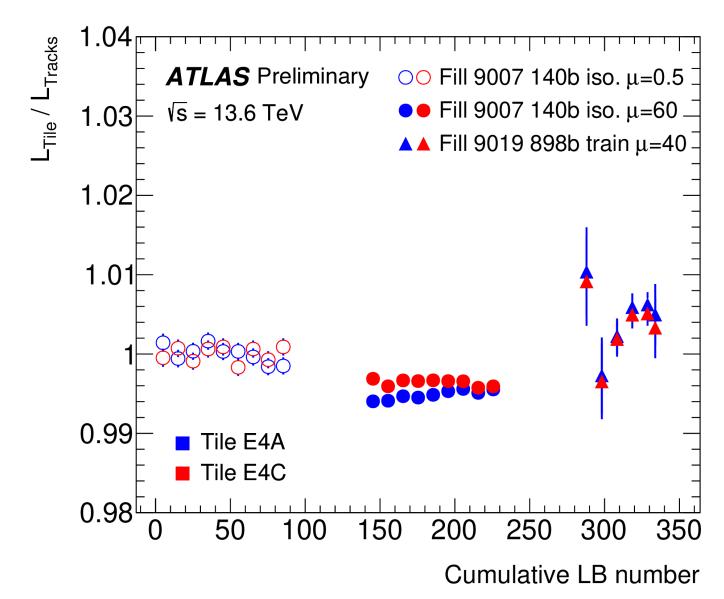
Calibration transfer

Uncertainty in track-counting luminosity

- Track-based LUCID correction assumes that TC is perfectly linear from vdM to physics regime
- **Probe TC linearity with Tile data** corrected for material activation, primarily using Tile E3/E4 gap scintillators
- Compare Tile/TC ratio in
 - vdM conditions
 - ▲ ▲ physics fill scheduled shortly after vdM
 - Combined 2022-2023 uncertainty = **1.23**%
 - → second largest systematic



Very limited personpower: new analysers with Tile experience needed!



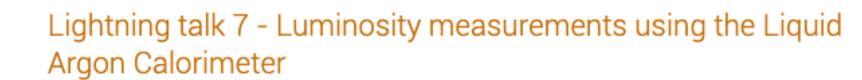
2022 → 1.5% tion	Used cell families	Range of shifts across used cell families
1-step extrap	oolation	
$(\mu \approx 0.5, 140b, isolated) \rightarrow (\mu \approx 40, 1154b, trains)$	A13, A14	[-0.1, 0.8]%
Alternative: 2-step	extrapolation	
$(\mu \approx 0.5, 140b, isolated) \rightarrow (\mu \approx 45, 144b, trains)$ $(\mu \approx 45, 144b, trains) \rightarrow (\mu \approx 40, 1154b, trains)$ (*)	A13, A14, E3, E4 A13, A14	[0.1, 0.7]% [0.0, 0.4]%
Combined 2-step extrapolation		[0.1, 1.1]%
Upper limit on extrapolation impact (rounded)		< 1%
Effect of missing laser corrections (linearly added)	A14	0.5%
Upper limit on total extrapolation impact		< 1.5%
2023 → 1.1% ion	Used cell families	Range of shifts acros

Long-term stability

Check of luminosity stability over time

- Use independent measurements from calorimeters normalised (anchored) to track-counting around vdM fill
 - LAr calorimeters: EMEC and FCal
 - Tile calorimeter: D6 Cells

 Long-term stability derived from luminosity-weighted difference between calorimeters and LUCID

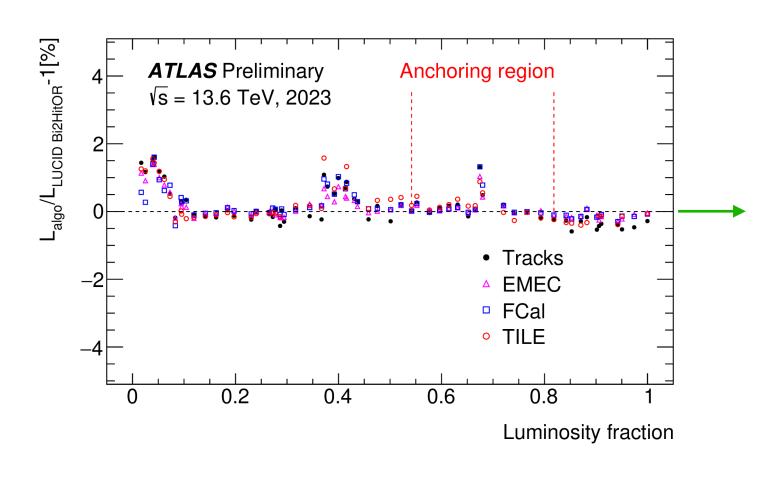


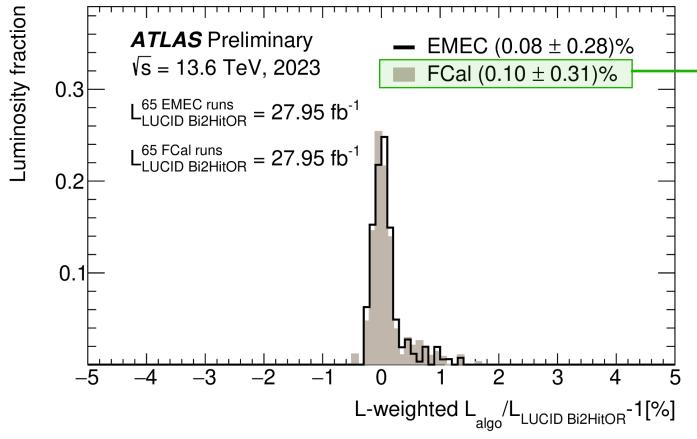
- Feb 21, 2025, 10:30 AM
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- 9 500/1-001 Main Auditorium (CERN)



Speaker

L Kyle Amirie (University of Toronto (CA))



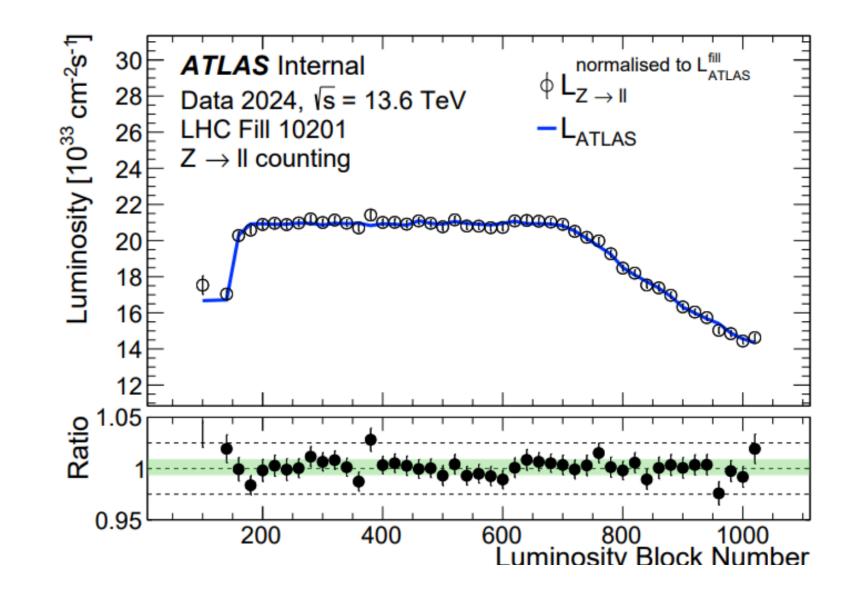


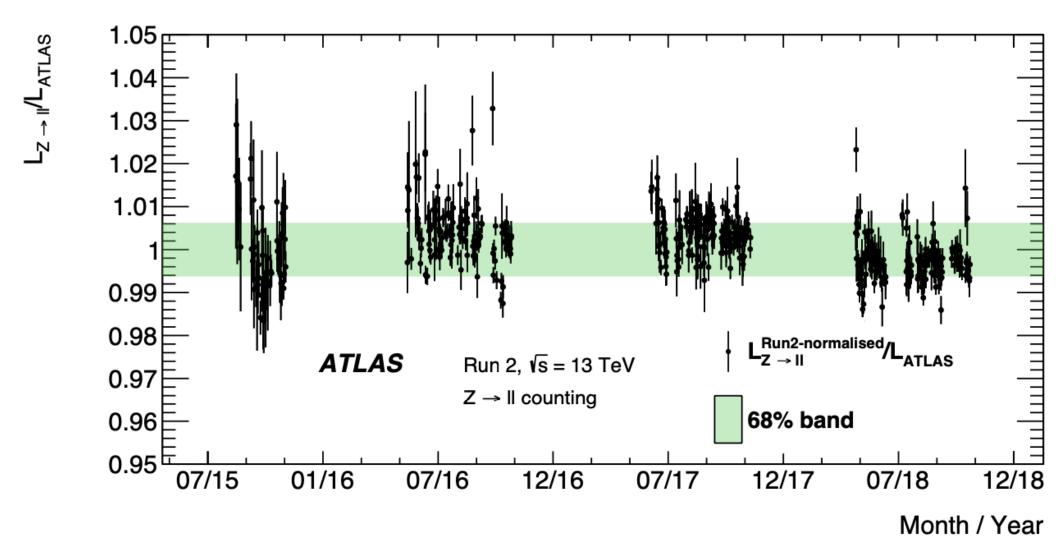
- Uncertainty \equiv largest mean $\Delta L/L$
 - \rightarrow equivalent to taking L_{calo} instead of L_{LUCID}
- 0.41% (2022), 0.10% (2023)
 - \rightarrow 0.22% combined

Z-Counting

Independent validation

- Counting $Z \rightarrow ee$ and $Z \rightarrow \mu\mu$ can be used for **relative luminosity measurement**
 - σ_Z only known to 3-4% (PDFs) cannot use for absolute luminosity scale :(
 - Comparison with baseline luminosity evaluated per data-taking period,
 with L(Z) normalised to L(ATLAS) in the period
 - Validates calibration stability with time and μ
 - Over full Run probes inter-year consistency of vdM calibration
 - Could backport future improved knowledge of non-factorisation to earlier years





Putting it all together

Status and plans

Calibration status:

- Good **accuracy in preliminary 2022-23** *pp* analysis (**~2% per year**) but dominated by correlated systematics: non-factorisation (1-1.5%), calibration transfer (1-1.5%)
- Preliminary uncertainty in 2023 PbPb analysis is 3.71%: expected to improve with better evaluation of systematics
- Calibration with 2024 pp data before the upcoming summer is the highest priority (less strict timelines for 2024 pp-ref and PbPb)



- Chronic lack of personpower and several critical areas understaffed:
 - 2024 luminosity measurement for EPS-HEP 2025 at risk
 - Urgent need for person-power in the Online Luminosity team

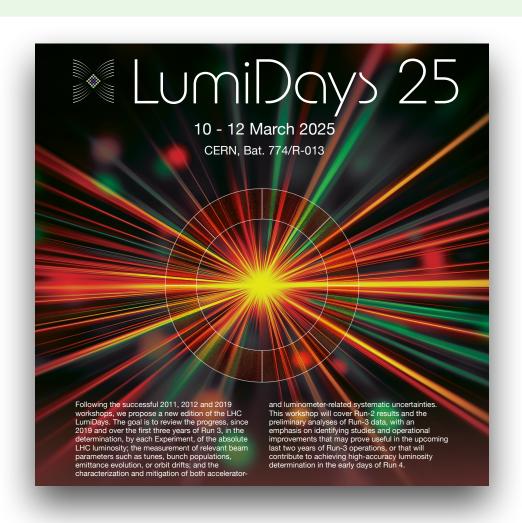
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Wk	25	26	27	28	
Мо	16	23	30	ZDCs out 7	
Tu			O ion setting up	VdM	
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- Preparing for a busy 2025
 - Intense scan program: pp, PbPb and Oxygen Run (pO, OO)
 - Lots of data to analyse!

[ATL-DAPR-PUB-2024-001]

Data sample	2022	2023	Comb.
Integrated luminosity [fb ⁻¹]	31.40	27.58	58.98
Total uncertainty [fb ⁻¹]	0.69	0.56	1.16
Uncertainty contributions [%]:			
Statistical uncertainty	0.01	0.01	0.01
Fit model*	0.24	0.15	0.20
Background subtraction*	0.06	0.30	0.17
FBCT bunch-by-bunch fractions*	0.01	0.01	0.01
Ghost-charge and satellite bunches [†]	0.17	0.04	0.11
DCCT calibration*	0.20	0.20	0.20
Orbit-drift correction	0.06	0.34	0.16
μ -dependence	0.00	0.30	0.14
Beam position jitter	0.00	0.01	0.01
Non-factorisation effects*	1.07	1.39	1.22
Beam-beam effects*	0.35	0.32	0.34
Emittance damping correction*	0.21	0.06	0.14
Length scale calibration	0.03	0.02	0.02
Inner detector length scale*	0.12	0.12	0.12
Magnetic non-linearity*	0.32	0.28	0.30
Bunch-by-bunch $\sigma_{ m vis}$ consistency	0.50	0.36	0.31
Scan-to-scan reproducibility	0.27	0.35	0.22
Reference specific luminosity*	0.43	0.44	0.43
Subtotal vdM calibration	1.44	1.71	1.49
Calibration transfer [†]	1.50	1.10	1.23
Calibration anchoring	0.53	0.16	0.29
Long-term stability	0.41	0.10	0.22
Total uncertainty [%]	2.19	2.04	1.97

■ Bonus: LHC LumDays 2025 (Mar 10-12)





Final Run 2 luminosity determination

Data sample	2015	2016	2017	2018	Comb.
Integrated luminosity [fb ⁻¹]	3.24	33.40	44.63	58.79	140.07
Total uncertainty [fb ⁻¹]	0.04	0.30	0.50	0.64	1.17
Uncertainty contributions [%]:					
Statistical uncertainty	0.07	0.02	0.02	0.03	0.01
Fit model*	0.14	0.08	0.09	0.17	0.12
Background subtraction*	0.06	0.11	0.19	0.11	0.13
FBCT bunch-by-bunch fractions*	0.07	0.09	0.07	0.07	0.07
Ghost-charge and satellite bunches*	0.04	0.04	0.02	0.09	0.05
DCCT calibration*	0.20	0.20	0.20	0.20	0.20
Orbit-drift correction	0.05	0.02	0.02	0.01	0.01
Beam position jitter	0.20	0.22	0.20	0.23	0.13
Non-factorisation effects*	0.60	0.30	0.10	0.30	0.24
Beam-beam effects*	0.27	0.25	0.26	0.26	0.26
Emittance growth correction*	0.04	0.02	0.09	0.02	0.04
Length scale calibration	0.03	0.06	0.04	0.04	0.03
Inner detector length scale*	0.12	0.12	0.12	0.12	0.12
Magnetic non-linearity	0.37	0.07	0.34	0.60	0.27
Bunch-by-bunch $\sigma_{\rm vis}$ consistency	0.44	0.28	0.19	0.00	0.09
Scan-to-scan reproducibility	0.09	0.18	0.71	0.30	0.26
Reference specific luminosity	0.13	0.29	0.30	0.31	0.18
Subtotal vdM calibration	0.96	0.70	0.99	0.93	0.65
Calibration transfer*	0.50	0.50	0.50	0.50	0.50
Calibration anchoring	0.22	0.18	0.14	0.26	0.13
Long-term stability	0.23	0.12	0.16	0.12	0.08
Total uncertainty [%]	1.13	0.89	1.13	1.10	0.83
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Personpower issues

Current scenario

[E. Torrence]

- Personnel are still (always) an issue. Currently missing people for
 - Online subgroup convener
 - vdM subgroup convener
 - Calibration transfer subgroup convener
 - Soon Lumi Convener
 - Non-factorization analysis losing main analyzer here
 - 0 ...
- Large number of scans in 2024 also means lots to analyze
 - Some new people interested, but still ramping up
 - Currently don't see that we will have 2024 pp calibrations before Summer (priority),
 ppref and PbPb timelines unclear
 - Situation will be similar (or worse) in 2025 pp, p-O, O-O, PbPb

Factorisation assumption

[T. Barklow]

$$L_b = \frac{f_r n_1 n_2}{2\pi \left[\sum_x \sum_y\right]} = \frac{\mu_{\text{vis}} f_r}{\sigma_{\text{vis}}} \qquad \qquad \left[\sum_x \sum_y\right] = \frac{1}{2\pi} \frac{\int R_{x,y}(\delta_x, \delta_y) d\delta_x d\delta_y}{R_{x,y}(0,0)} = \frac{1}{2\pi} \frac{1}{\int \hat{\rho}_1(x, y) \hat{\rho}_2(x, y) dx dy}$$

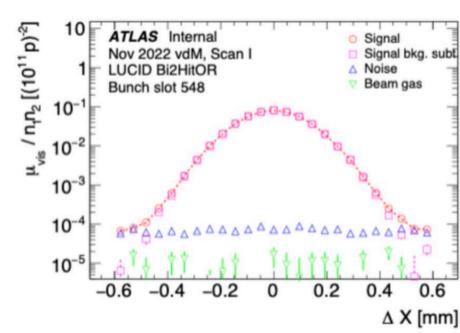
 f_r = bunch revolution frequency $n_1 n_2$ = bunch population product σ_{vis} = visible cross section

 $\mu_{\mathrm{vis}} = \mathrm{visible}$ interaction rate per bunch crossing as measured e.g. in LUCID

 δ_{x} , δ_{y} = horizontal, vertical beam separation

 $R_{x,y}(\delta_x,\delta_y)$ = any quantity proportional to the luminosity, such as LUCID μ_{vis}

 $\hat{\rho}_1(x,y)$, $\hat{\rho}_2(x,y)$ = normalized transverse particle density of beam 1, 2



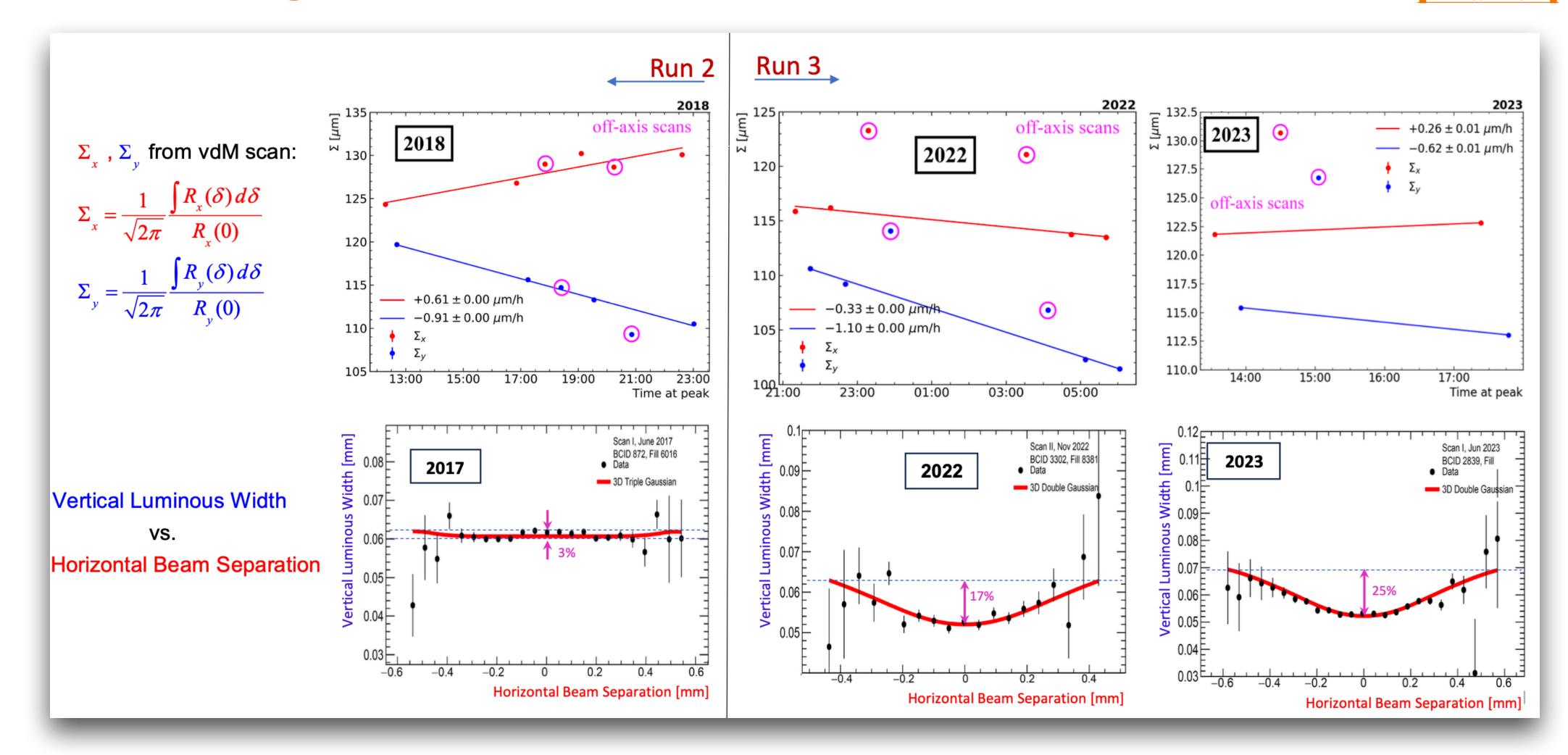
In practice, a grid scan covering the δ_x, δ_y 2-d plane is expensive in beam time, so the factorization assumption

 $R_{x,y}(\delta_x, \delta_y) = R_x(\delta_x)R_y(\delta_y) \text{ is often used } \Rightarrow [\Sigma_x \Sigma_y] = \Sigma_x \Sigma_y \text{ where 1-d scans provide } \Sigma_x = \frac{1}{\sqrt{2\pi}} \frac{\int R_x(\delta)d\delta}{R_x(0)} \text{ & } \Sigma_y = \frac{1}{\sqrt{2\pi}} \frac{\int R_y(\delta)d\delta}{R_y(0)}$

 $\Sigma_x(\Sigma_y) = \text{horizontal (vertical) convolved beam size:}$ $\Sigma_x = \sqrt{\sigma_{1x}^2 + \sigma_{2x}^2}$ for single Gaussian beams, where $\sigma_{jx} = \text{horiz beam size}$

Evidence for Large Run 3 Non-Factorisation from vdM Scans & LRE

[T. Barklow]



LRE analysis

[T. Barklow]

Single-beam parameters are obtained from a combined fit to the beam separation dependence of:

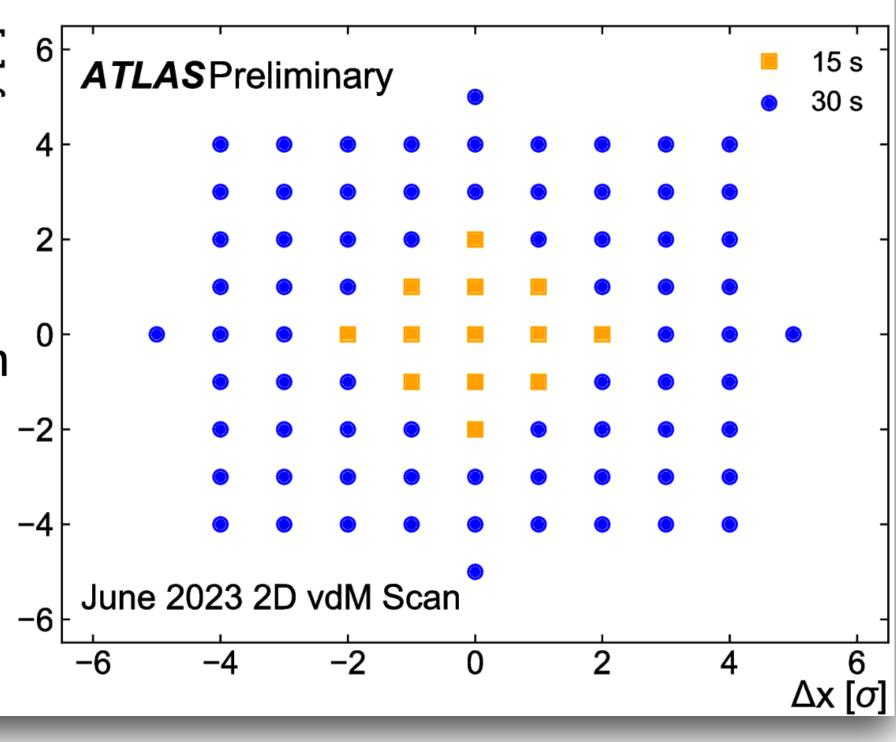
- $R(\delta_x, \delta_y)$ (\mathcal{L})
- luminous-region observables: <x,y,z>_L , σ_{x,y,z} , ...
- Single beam parameters provide normalized transverse beam particle densities $\hat{\rho}_1(x,y)$, $\hat{\rho}_2(x,y)$
- Directly calculate $\left[\sum_{x}\sum_{y}\right] = \frac{1}{2\pi \int \hat{\rho}_{1}(x,y)\hat{\rho}_{2}(x,y)dx\,dy}$.
- Combine $[\Sigma_x \Sigma_y]$ with bunch population product $n_1 n_2$ to get luminosity $L_b = \frac{f_r n_1 n_2}{2\pi [\Sigma_x \Sigma_y]}$ free of factorization assumptions.
- $\hat{\rho}_1(x,y)$, $\hat{\rho}_2(x,y)$ also used to generate simulated 1D vdM scans
- These scans are analyzed in the same fashion as real vdM scan data (i.e., with the factorization assumption)
- Non-factorization bias $R_{NF} \equiv \frac{L_b}{L_b^0}$ where $L_b^0 = \frac{f_r n_1 n_2}{2\pi \sum_x \sum_y}$ is from the simulated 1D vdM scan analysis. Historically $\frac{1}{R_{NF}}$ has been applied directly to $\sigma_{\text{vis,1DvdM}}$ to correct for non-factorization: $\sigma_{\text{vis,corrected}} = \sigma_{\text{vis,1DvdM}} / R_{NF}$
- Only 12 of the 136 lumi data beam crossings have beamspot data, so the average $\langle R_{NF} \rangle$ over the 12 BCID's is taken as the correction for each scan.

Non-factorisation analysis

Two roads to understand non factorisation

[K. Mönig]

- LRE analysis: parametrise single beam profile using scan data plus beamspot
 - can measure non-factorisation during scans
 - only possible for few bunches due to trigger limitations
 - needs some assumptions on shape
- Grid scan:
 - perform a scan over a 2d grid
 - can run on all bunches
 - in principle can measure non-factorisation in a model independent way
 - however due to time limitations can only be done once
 - Since non-factorisation changes with time cannot use it to correct 1d scans



Non-factorisation analysis

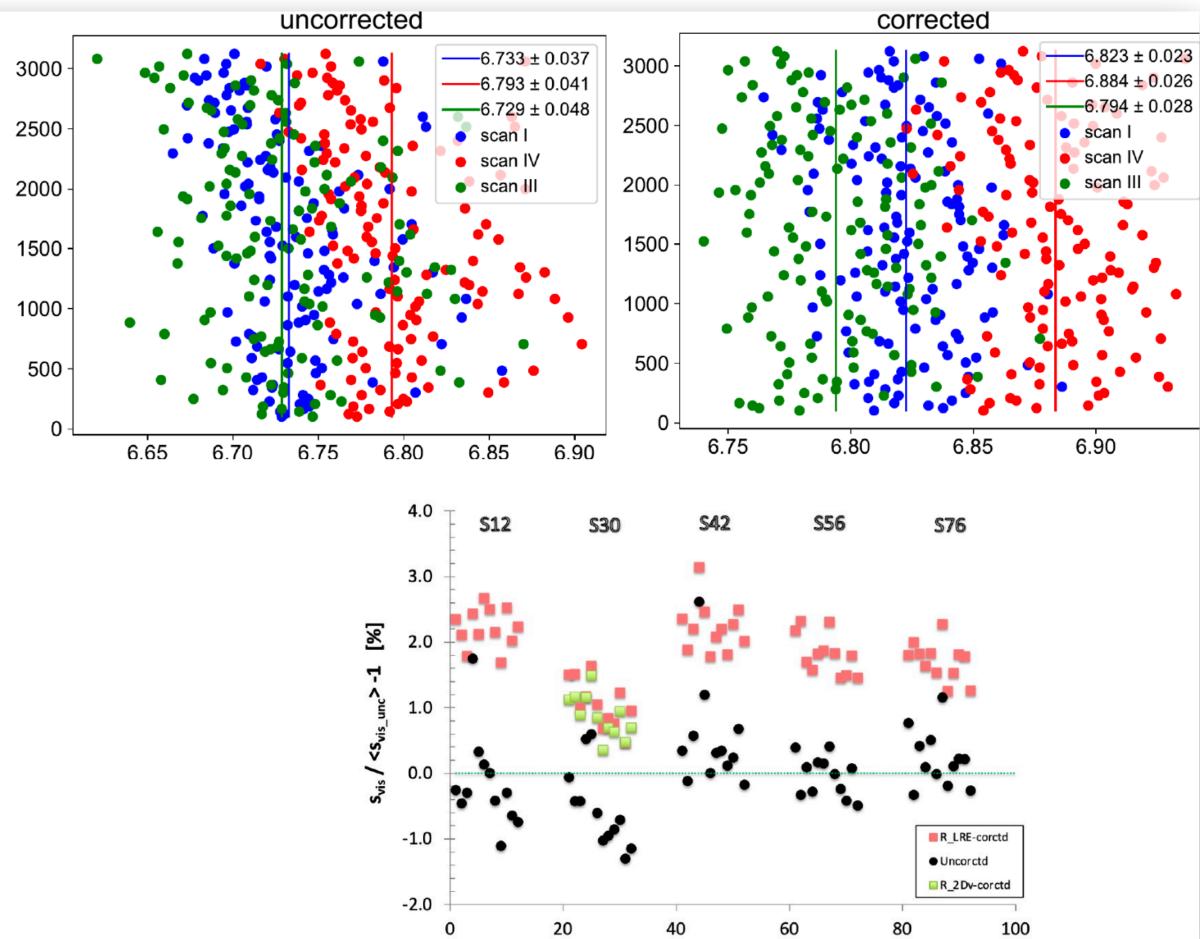
Current issues

[K. Mönig]

The 2d scan seems to catch the bunch-by-bunch pattern but cannot be used to correct the 1d scans due to time dependence

2D grid scan (Scan III)

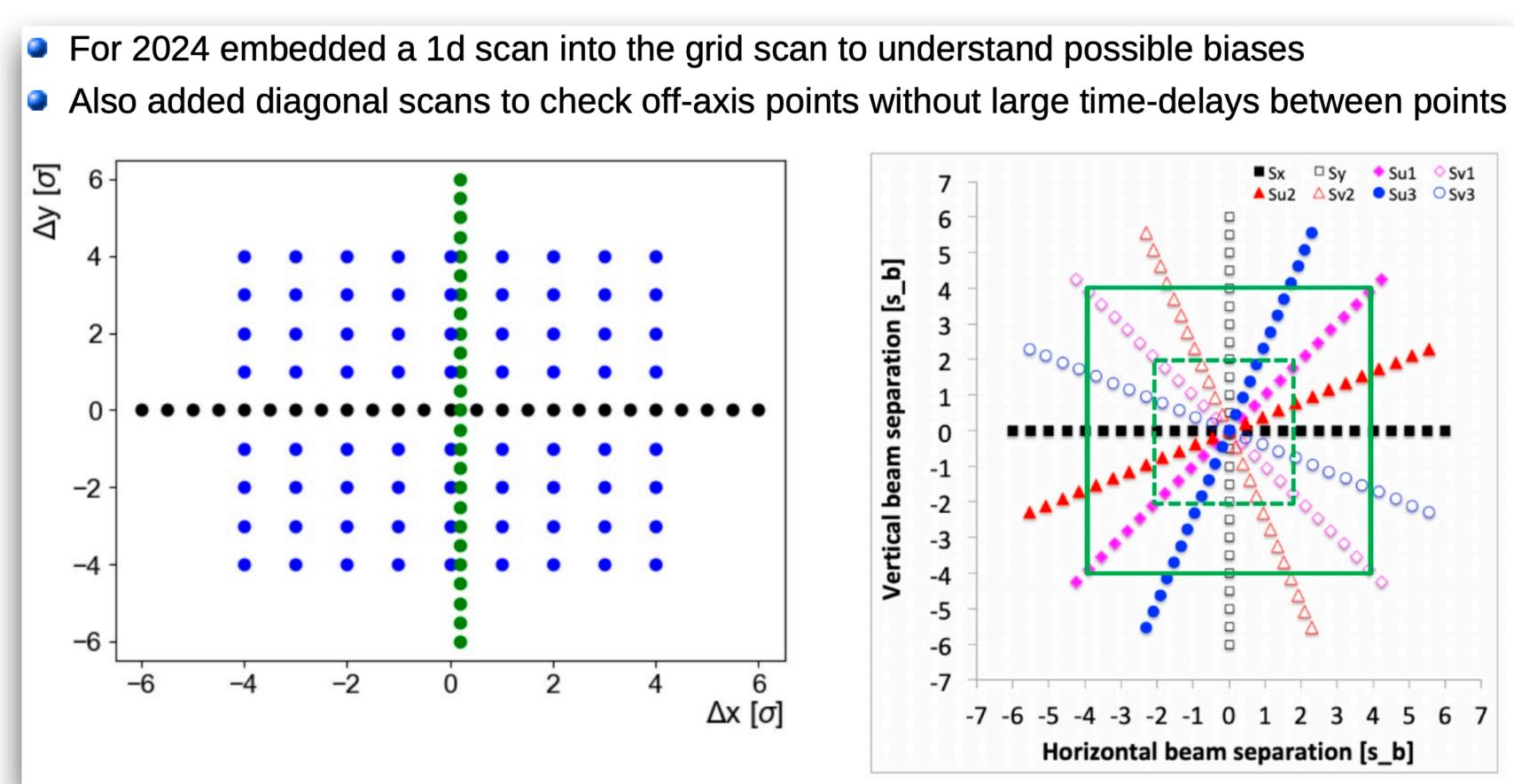
The LRE fit agrees well with the 2d fit but there seems a bias between the 1d scans and the 2d scan



Non-factorisation analysis

Mitigation strategies

[K. Mönig]

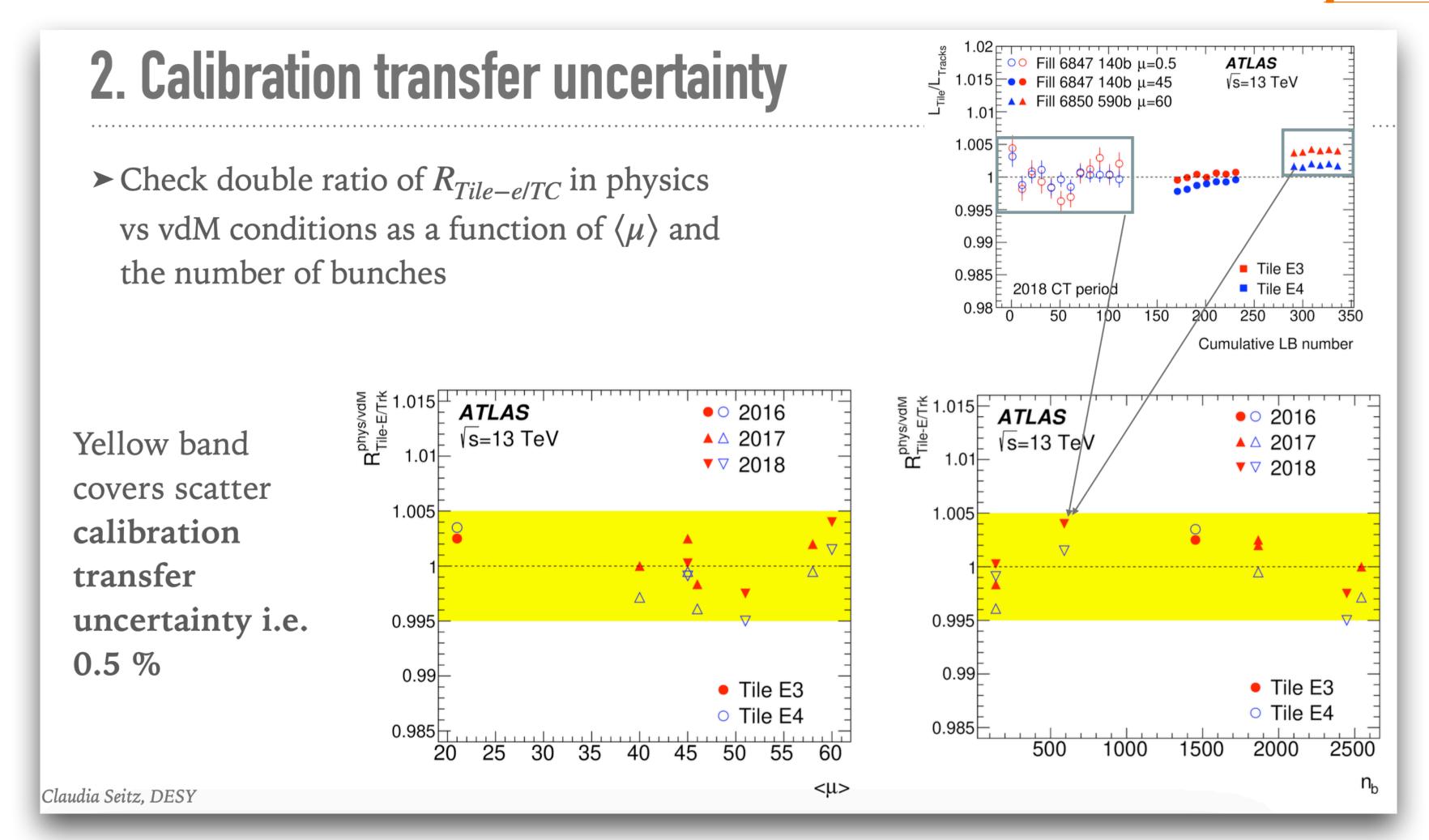


• Dedicated 'beam tailoring' in injectors to produce Gaussian-like beams reduces non-factorisation

Calibration Transfer

Uncertainty in track-counting luminosity

[C. Seitz]



Long-term stability

Methodology

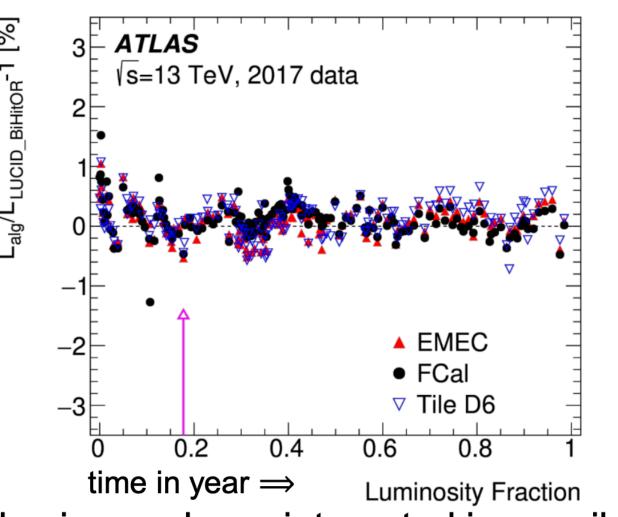
[R. Hawkings]

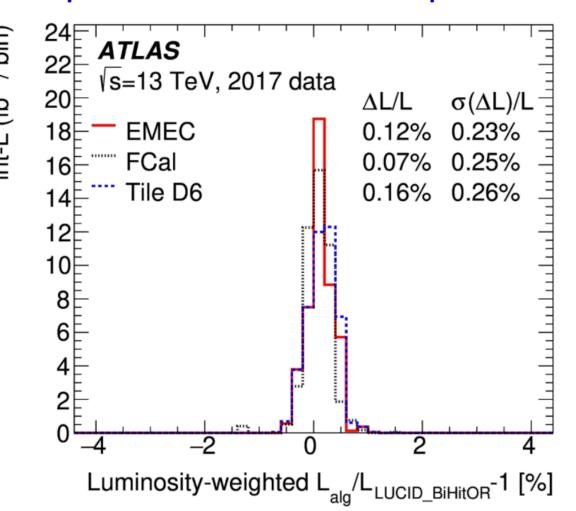


Long-term stability – methodology



- Now study calorimeter/LUCID differences for each run, for EMEC, FCal, Tile
 - ΔL/L typically within 0.5% for long runs, apart from start-up period with low n_b
 - Outliers tend to be short runs with low L_{int}
 - Integrated-luminosity-weighted histogram of per-run ΔL/L is more representative





- Physics analyses interested in possible deviations of L_{int} for whole sample
 - Mean of ΔL/L histogram captures this take largest mean from EMEC, FCal, Tile to define long-term stability uncertainty (equiv. to taking L_{calo} instead of L_{LUCID})

31st January 2023

Richard Hawkings

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Z-Counting

The Method

[J. Newell]

 Use the decays of Z bosons to electrons or muons to determine luminosity:

$$\mathcal{L}_Z = rac{N_Z}{\sigma_Z}$$



Monte Carlo correction factors and data-driven efficiencies are applied

$$\mathcal{L}_{Z}(\Delta t) = rac{N_{Z
ightarrow l^+ l^-}(\Delta t) imes (1 - f_{bkg})}{F^{ ext{MC}}(\mu) imes A^{ ext{MC}} imes rac{\epsilon_{Z
ightarrow l^+ l^-}^{T\&P}(\Delta t) imes \sigma_{theory} imes \Delta t}{}$$

• Theoretical cross-section does not affect plots shown (Cancels in $\mathcal{L}_{e^+e^-}/\mathcal{L}_{\mu^+\mu^-}$ and $\mathcal{L}_{Z\to l^+l^-}/\mathcal{L}_{ATLAS}$ ratios)

In-situ Z-event-level Efficiency per time (LB):

$$arepsilon_{Z
ightarrow l^+l^-}^{T\&P} \,=\, \left(1-\left(1-arepsilon_{ ext{trig, }1l}
ight)^2
ight) imes\left(arepsilon_{ ext{reco, }1l}
ight)^2$$

 Broken down into single-lepton reconstruction (tag-and-probe) and trigger efficiencies

Pileup-dependent Monte Carlo correction factor:

$$F^{ ext{MC}} = rac{N_{Z o l^+l^-}^{ ext{reco,fiducial,MC}}(\mu)}{N_{Z o l^+l^-}^{ ext{generated,MC}}(\mu) imes A^{ ext{MC}}} imes rac{1}{arepsilon_{Z o l^+l^-}^{ ext{T\&}P,MC}(\mu)}$$

Accounts for non-closure between data-driven reconstruction and trigger efficiencies and the true Z-event level efficiency given by Monte Carlo simulation

3

Oxygen Run

p-O Details

[E. Torrence]

- LHCf is one of the physics priorities for the p-O data, driving IP1 parameters
- Target 1.5 nb⁻¹ at $\mu \sim 0.01$ (0.03 for 10%) (~36h of stable beams)
- Max IP1/5 μ not clear, but $\mu \lesssim 0.05$ -0.1 was mentioned by Roderik in `21 talk
- LHCf prefers $\beta^* \sim 10-20$ m, probably will get more like 1m (unsqueezed)
- Likely fill pattern with ~24 colliding bunches, > 200 ns spacing
- Vertical (negative) crossing angle (bad for ZDC)

Physics motivation is measuring cosmic-ray cross-sections, needed by astrophysics community, need ~3% precision on lumi uncertainty to match other expected uncertainties

Oxygen Run

0-0 Details

[E. Torrence]

- Parameters here are expected to be more HI-like
- Expect something like 0.5 nb⁻¹ in one day (one long fill) or 2 days
- Energy requested to match PbPb nucleon-nucleon energy
- μ_{max} ~ 0.6 has been mentioned, may need to be levelled for physics
- May keep $\beta^* \sim 1m$ to save setup time decided to keep at $\beta^* \sim 0.5m$
- Vertical (positive) crossing for ZDC, Roderik wants to go to 0 xangle
- Probably 8 bunches colliding (12b_8_8_8) optimal for all IPs

ATLAS luminosity target probably around 5% to match expected statistics

Oxygen Run

Tentative plan

[E. Torrence]

- For p-O we really want a 5-point LSC non-linearity significant issue in '24
 PbPb, large risk if we go to 3 points
- For O-O optics identical to 2024 PbPb don't need additional LSC, can argue we roughly know this without too much of a systematic penalty
- Can achieve similar results with 1 head-on + grid as 2 head-on + diagonal, similar time estimates, choose between these internally
- For p-O, believe we can achieve something like 3% in 4 hours (Silver)
- For O-O, believe we can achieve something like 5% in 2.5 hours (Bronze), but with a bit more risk, 1.5 hours would be in the range of 5-7%. Depends on whether LRE analysis alone can constrain NF at all in this data.

2018 ALFA runs at 900 GeV

Dataset composed of 14 runs at $\sqrt{s} = 900$ GeV + dedicated vdM scans

[V. Maksimovic]

ALFA runs: notable characteristics

- $> <\mu> = 0.01 0.06, \beta* = 50/100 m$
- No crossing angle
 - 2 runs with with ID ON
 - 3 runs with crystal collimation (CC)
 - 2 runs with $\beta^* = 11$ m
 - Only 5 runs with unpaired bunches

vdM scans: notable characteristics

- Oct 2018: 363514, 363516
 - Fills 7299-7300
 - On-axis scans: \$1, \$3, \$4 (\$2 off-axis)
- Nov 2018: 365218, 365219
 - Fills 7406-7407
 - On-axis scans: S5, S6 (S7 off-axis)
- 11 m beam optics
- 150 colliding bunch-pairs
- 2 unpaired bunches per beam
- lucBi2HitOR algorithm used as nominal (using 5 available PMTs)

2018 ALFA runs at 900 GeV

Results

[V. Maksimovic]

LUCID data is used to to estimate the absolute luminosity value for the \sqrt{s} =900 GeV 2018 pp runs

- Calibrated visible cross-section from vdM: 0.85326 mb ± 0.07%(stat) ± 1.85%(syst)
 - o SigVis value calculated using November session scans only and the lucBi2HitOR algorithm
 - o Major contributors to total systematic uncertainty: Reference Specific Luminosity and Non-factorization
- LUCID and TC data are analysed to estimate the luminosity value and validate it
- The <u>main backgrounds</u> are <u>either subtracted</u> (afterglow + Bismuth) <u>or</u> their contribution is <u>evaluated</u> and accounted for as a systematic uncertainty (single-beam)
- A <u>calibration method for the LUCID AND</u> algorithm has been developed, allowing to improve the reliability of the stability uncertainty
- Given that the β *=11m dataset is composed only by 2 runs (only one with unpaired bunches, non with TC available), it was decided to quote the same systematic uncertainty as for the β *=90/100 m dataset

$$\beta^* = 11 \,\mathrm{m}$$
 $\mathcal{L} = 504.9 \pm 0.4_{stat} \pm 12.5_{sys} \,\mu\mathrm{b}^{-1}$

$$\beta^* = 50/100 \,\mathrm{m}$$
 $\mathcal{L} = 932.6 \pm 0.5_{stat} \pm 23.0_{sys} \,\mu\mathrm{b}^{-1}$

Uncertainty [%]	$\beta^* = 11 \mathrm{m}$	$\beta^* = 50/100\mathrm{m}$	
Stability	1.46 (sys)		
vdM calibration	1.85 (sys)		
Single-beam background	0.74	(sys)	
Total systematic	2.47		
Statistical uncertainty	0.079 (stat)	0.053 (stat)	
Total uncertainty	2.47	2.47	

2018 ALFA runs at 900 GeV

Track counting data

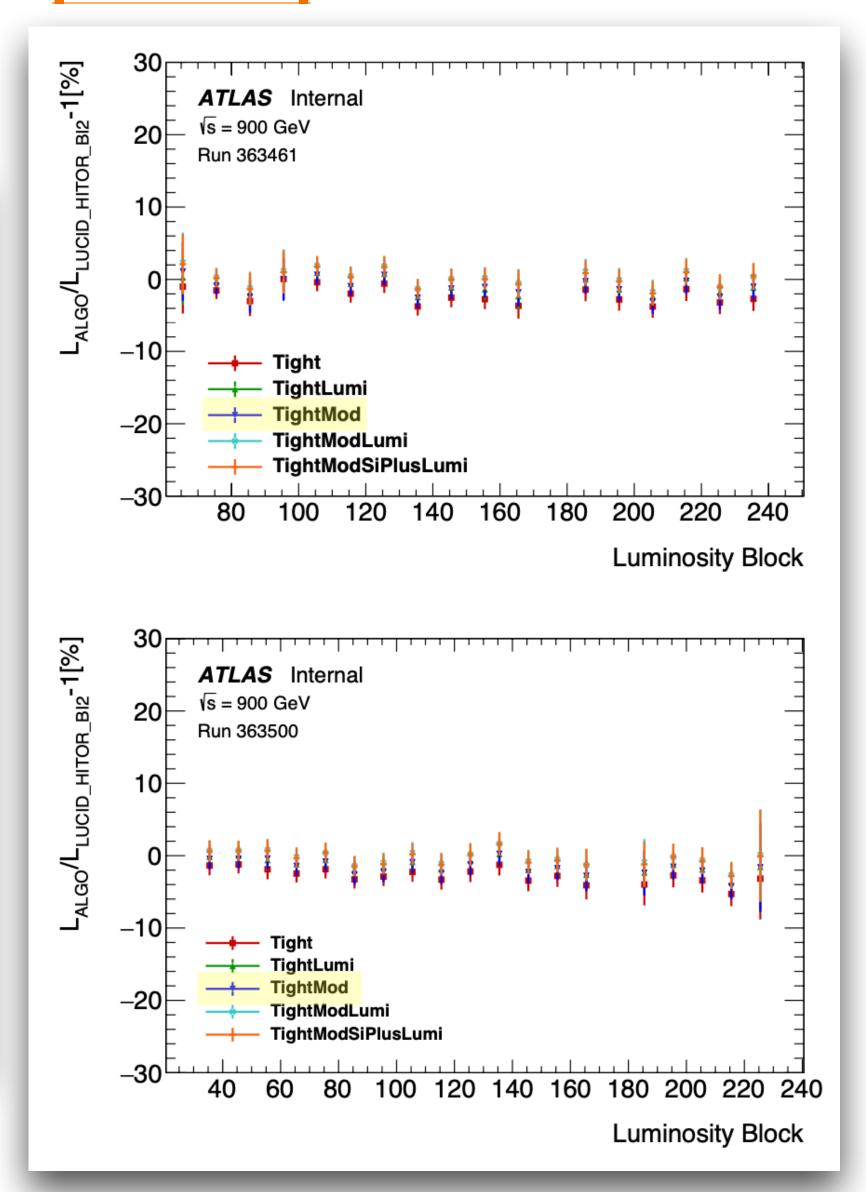
Track counting (TC) data from the Inner Detector is compared to LUCID's measurement **for stability evaluation**

Note: this is the only other detector that could provide a measurement in this regime (Calo has no sensitivity in this low- μ regime)

- ➤ Inner Detector data available in 2 runs, with a reduced LB selection
- TC data anchored to reference algo (LUCID HITOR BI2) in non-scan periods of the November vdM runs
- > TC selection optimised for low-luminosity: **TightMod**

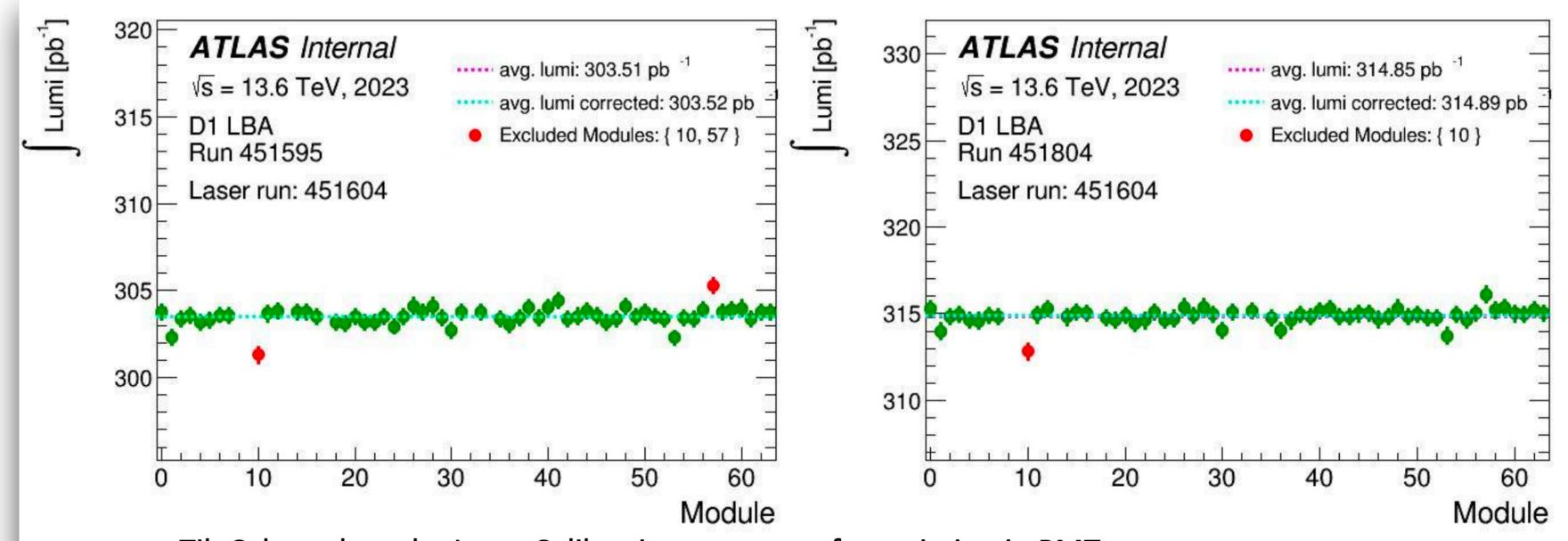
Algorithm	$\mathscr{L}_{tot} [\mu b^{-1}]$	$\Delta \mathscr{L}/\mathscr{L}_{ref}$ [%]
LUCID_HITOR_BI2	154.76 ± 0.22	//
TightMod	152.79 ± 0.20	-1.27 ± 0.19

[V. Maksimovic]



Laser corrections

[P. Rapheeha]



- TileCal employs the Laser Calibration to correct for variation in PMT responses
- See Rute's <u>presentation</u> on Laser correction derivation
- After Application of Laser corrections, modules that have luminosities at are 3 std. deviations from the mean are excluded
 - > Such modules are typically less than 2 for most runs
 - > The Laser run from which the correction are derived is given in the legend
- This exclusion is done on a per-run basis

0

Beam-beam effects

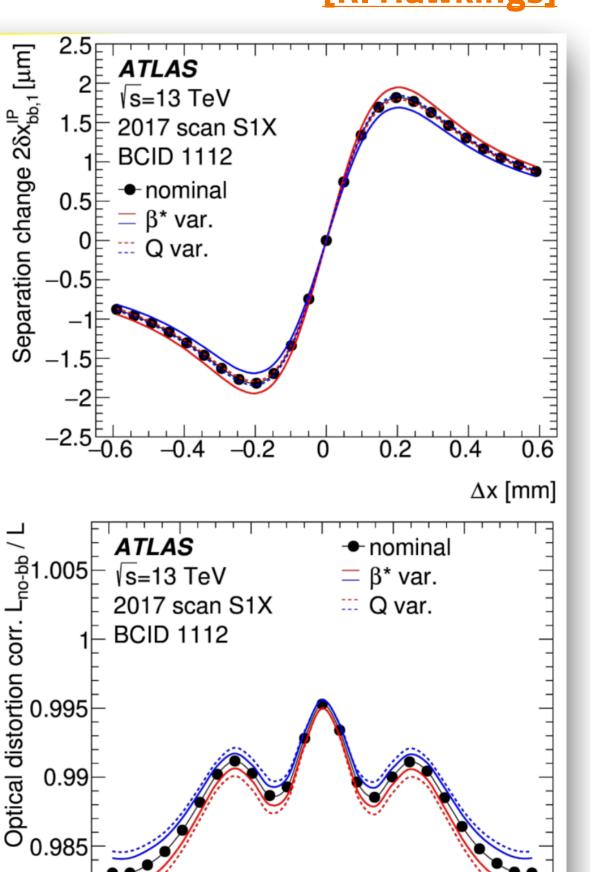
[R. Hawkings]

- Two effects from EM repulsion of colliding bunches
- 1) Separation-dependent beam-beam deflection
 - Beams give each other an angular kick θ_x resulting in an additional separation δ_x above that requested
 - Well-understood, depends on Δx , n_2 , Σ_x
 - Corrections to Δx of up to 2 μ m at Δx =200 μ m
- 2) Optical distortion, aka dynamic β
 - Opposing bunch defocuses the beam, changing its size and shape as a function of Δx , changing L_{inst}
 - Previously handled with a linear approximation
 - Beam-beam force is non-linear; proton in centre of the bunch feels a different force to one at the edge
 - Intensive common study in LHC luminosity WG, with multiparticle tracking codes <u>COMBI</u> and <u>B*B</u>
 - Developed parameterised correction as fn of ξ_R

$$\frac{L_{\text{no-bb}}}{L} (\Delta x) = f(\Delta x, Q_x, Q_y, \xi_R) \qquad \xi_R = \frac{r_p \bar{n} \beta^*}{2\pi \gamma \sum_x \sum_y}$$

31st January 2023

Richard Hawkings



0.2

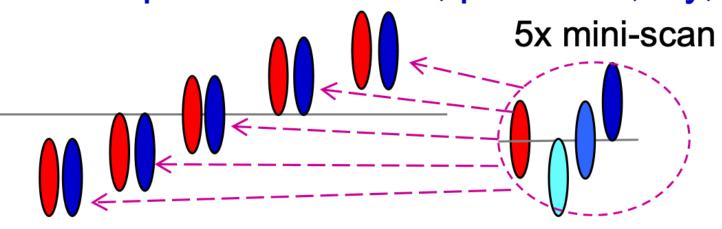
16_{Δx [mm]}

Length scale calibration

[R. Hawkings]

Relation between requested/real beam displacement

Calibrated in 5-point LSC scans, per beam, x/y, year



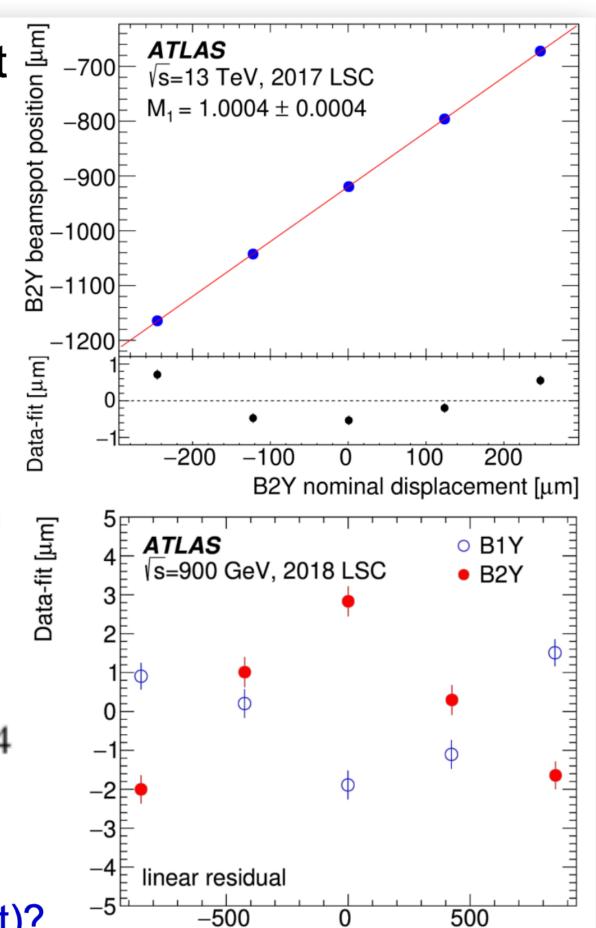
- True beam displacement measured from beamspot positions reconstructed from tracks in ATLAS ID
 - 3-point mini-scan of 'witness' beam around each position to interpolate to head-on collisions

Length scale factors M₁ all within ±0.4% of unity

- With random scatter from year to year
- σ_{vis} scales by L_{xy} : $L_{xy} = (M_1^{x,1} + M_1^{x,2})(M_1^{y,1} + M_1^{y,2})/4$

Fit residuals hint at potential non-linearities

- Much clearer in a LSC scan at 900 GeV in late 2018
- Reproducible? Magnet hysteresis (direction-dependent)?



Y nominal displacement [μm]

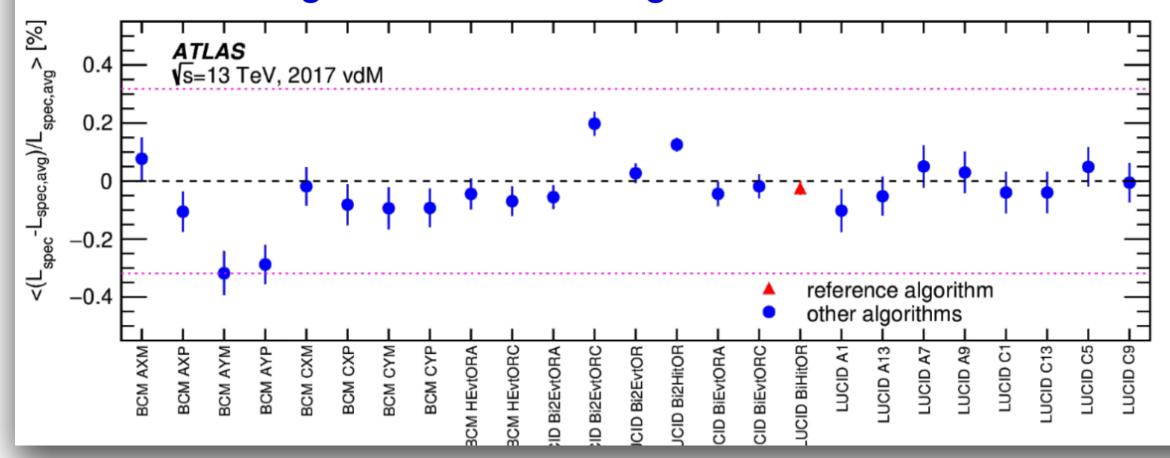
Consistency checks

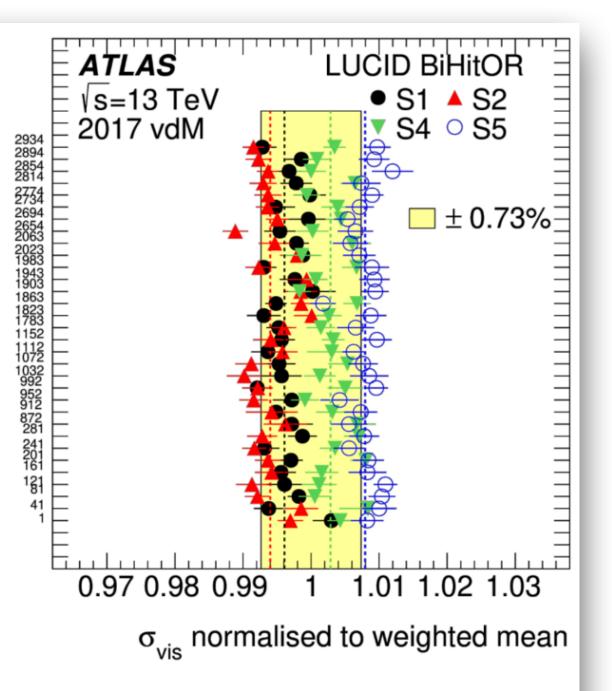
[R. Hawkings]

- Bunch-by-bunch consistency, up to 0.4%
 - Calculate RMS of σ_{vis} over all bunches in each scan, subtract expected spread due to stat. errors
- Scan-to-scan consistency, up to 0.7%
 - Sampling-corrected RMS of all scans in one year
- Reference specific luminosity L_{spec}:

$$\mathcal{L}_{\mathrm{spec}} = \frac{\mathcal{L}_{\mathrm{b}}}{n_1 n_2} = \frac{f_{\mathrm{r}}}{2\pi \sum_{x} \sum_{y}}$$

Per-algo. deviation averaged over all bunches/scans



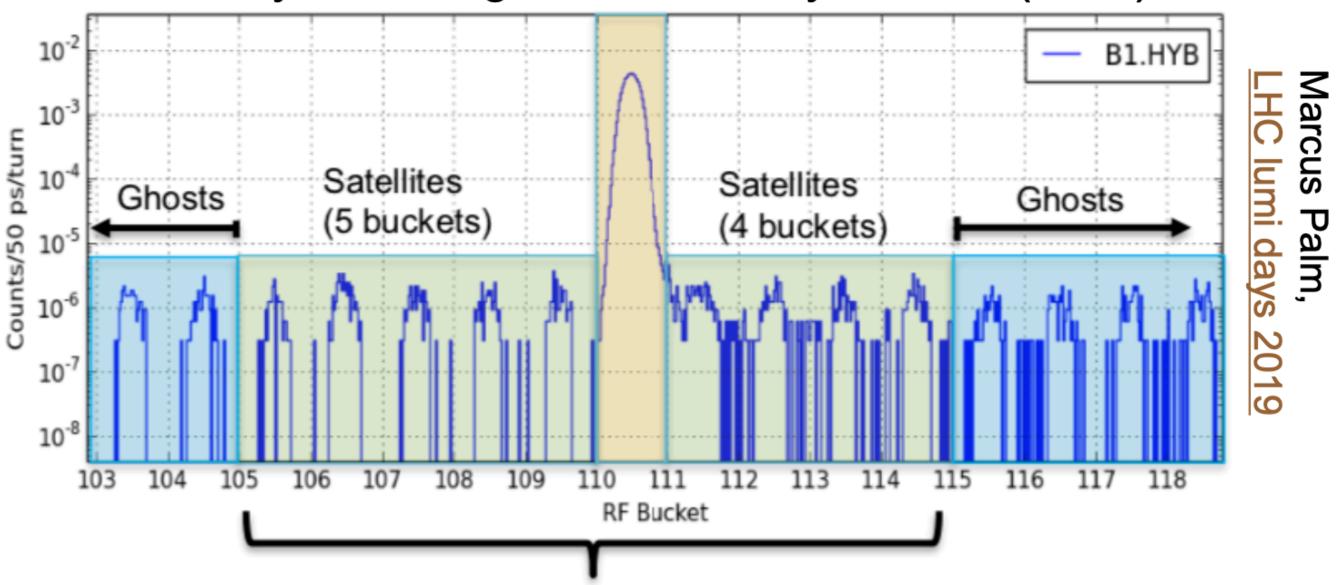


- **Total** per-year uncertainty on σ_{vis} :
 - **0.7-1.0%**

Bunch population measurements

[R. Hawkings]

- Measuring n₁ and n₂ relies on LHC instrumentation and LHCb SMOG
 - LHC ring has 3564 25ns bunch slots, each divided into 10 RF buckets (400 MHz)
 - Particles circulate stably in all these RF buckets
- Structure as seen by LHC longitudinal density monitor (LDM)



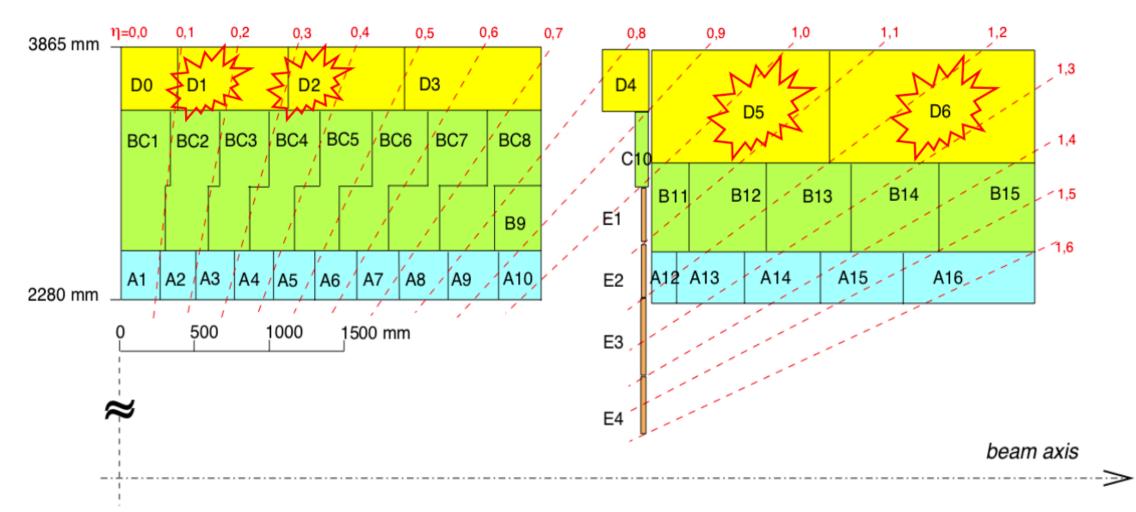
10 buckets = 1 slot \approx 25 ns

- Satellites: charge in filled slot, but in the wrong bucket
- Ghosts: charge in nominally-empty bunch slots (unfilled BCIDs)

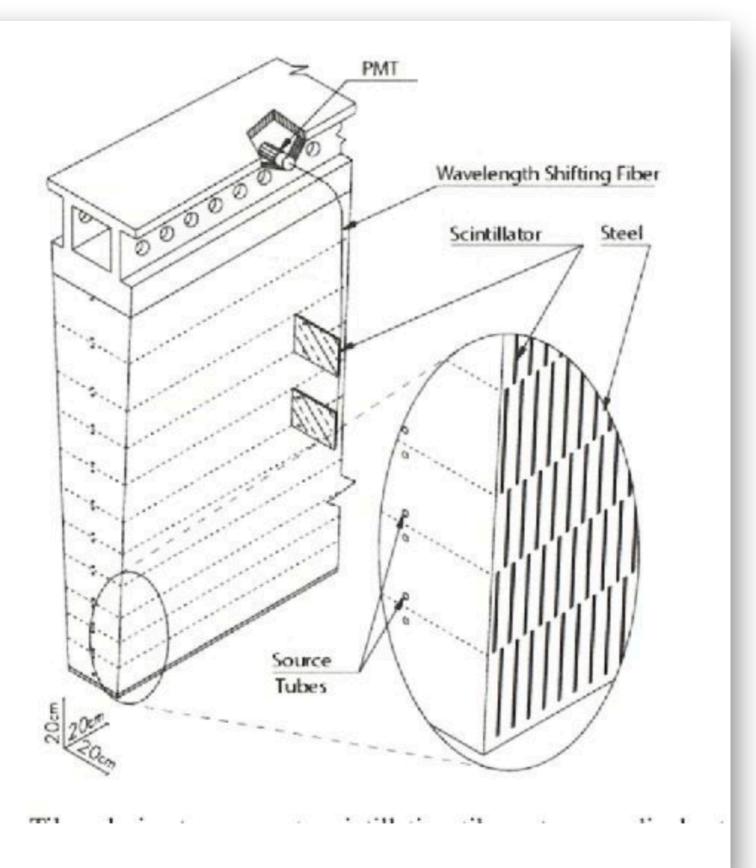
TileCal as Luminometer

[P. Rapheeha]

• The Tile Calorimeter has 64 wedge shaped modules around the beam axis and segmented longitudinal into three sections.



- Primary offline TileCal luminosity measurements for the Long-Term stability studies are obtained from D5/D6 cells
- D1 and cells will be used for systematic comparisons



TileCal as Luminometer

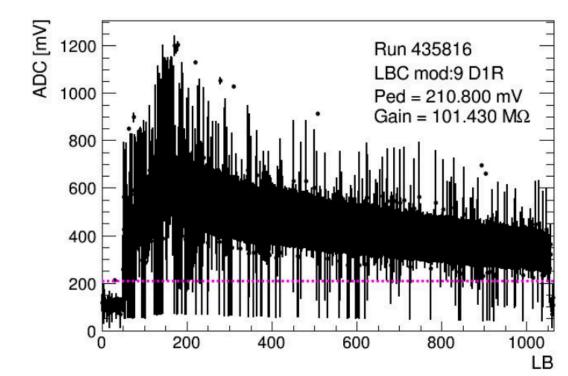
[P. Rapheeha]

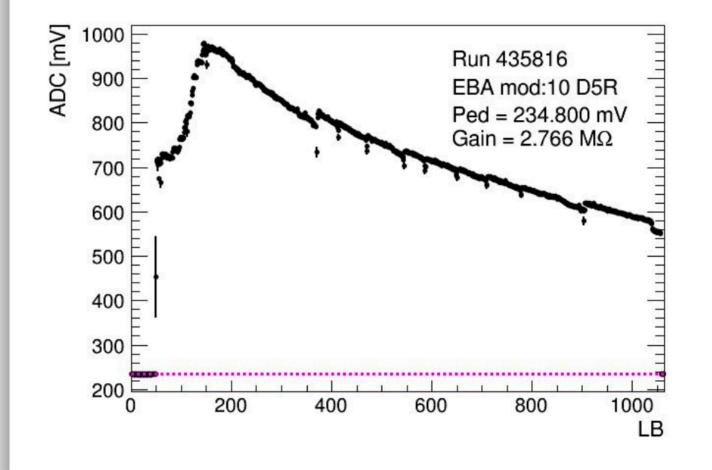
- Each cell is read by two PMTs
- The collision induced PMT current is given by:

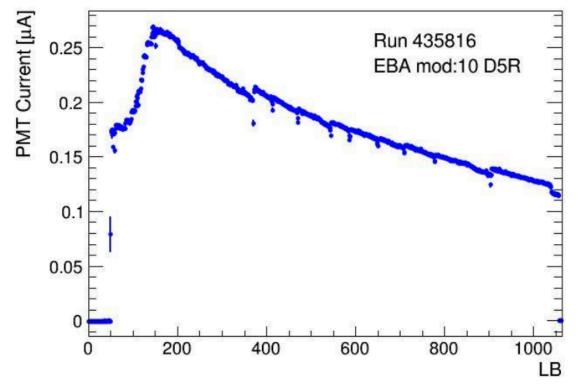
$$I_{PMT} = \frac{ADCs - pedestal}{Gain_{PMT}}$$

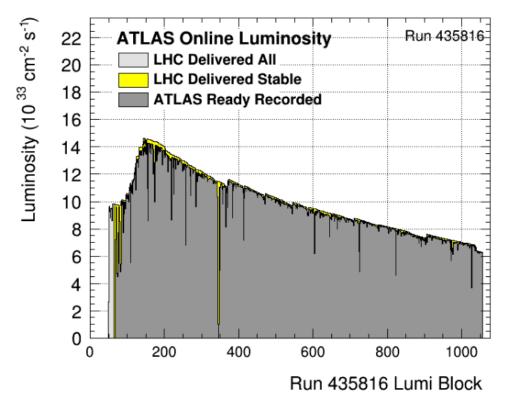
- The pedestal is taken as the average ADCs before collisions
- The PMT current of a given cell is proportional to the number of particles passing through the cell

- Not all modules are used in Luminosity studies
 - Modules that get power-cycled during runs, saturated or noisy modules are excluded









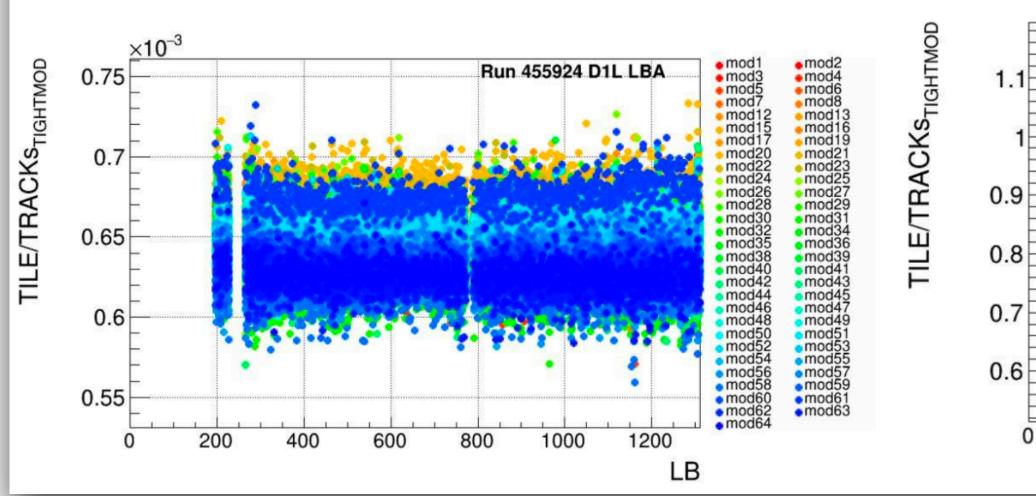
TileCal as Luminometer

[P. Rapheeha]

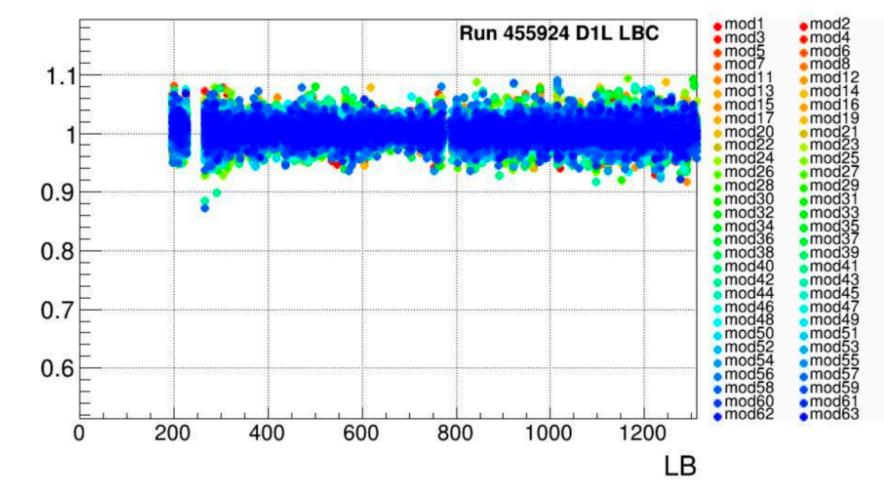
 PMT currents are cross-calibrated to trackcounting luminosity to PMT luminosities

$$\alpha_{module} = \frac{L_{Track}}{\langle I_{PMT} \rangle_{module}}$$

- The calibration constants are determined in an "anchoring run"
- Run 455924 is chosen as the anchoring run for the 2023 dataset
- LB 675 1350 Anchoring range

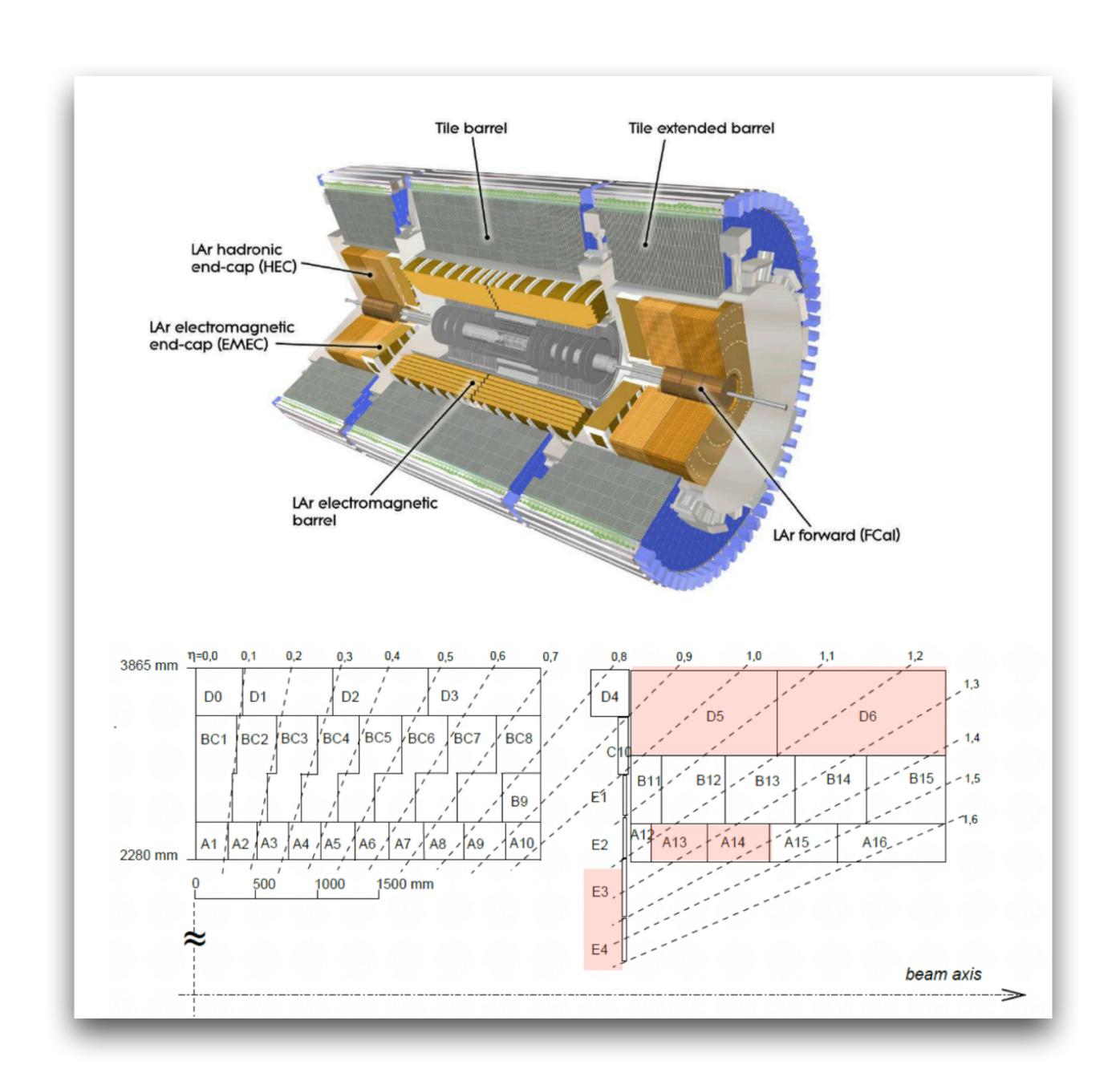


- The TileCal luminosity of a given cell in a given module obtained by taking the average current in the left and right PMTs
- The luminosity of a cell is given by the average luminosity across all good modules



Tile calorimeter

TILE cell division scheme



Combination across years

Standard method

and the absolute uncertainty in the total luminosity, $\sigma_{\mathcal{L}_{tot}}$, is given by standard error propagation as

$$\sigma_{\mathcal{L}_{tot}}^2 = \mathbf{e^T V_L e}$$
.

Here, V_L is the covariance matrix of the absolute luminosity uncertainties for the different years, and e is a column vector with unit entries. The covariance matrix is made up of the sum of terms corresponding to each uncertainty source in Table 8; uncorrelated uncertainties give rise to terms on the diagonal, whilst correlated sources are represented by terms with non-zero off-diagonal entries.

as $V_L = \sigma_L C \sigma_L^T$, where the vector σ_L of total absolute uncertainties on \mathcal{L}_i and (symmetric) correlation matrix C are given by

$$\boldsymbol{\sigma_{L}} = \begin{pmatrix} 0.0367 \\ 0.296 \\ 0.504 \\ 0.644 \end{pmatrix} \text{ fb}^{-1}, \; \mathbf{C} = \begin{pmatrix} 1.000 \\ 0.579 & 1.000 \\ 0.368 & 0.437 & 1.000 \\ 0.480 & 0.510 & 0.362 & 1.000 \end{pmatrix}.$$