Report for Physics Research Committee from



Open session DESY PRC 72

October 25-26, 2011 DESY Hamburg

spin and hadronization

HERMES main research topics:

- r origin of nucleon spin
 - Iongitudinal spin/momentum structure
 - ransverse spin/momentum structure
- ▶ hadronization/fragmentation

spin and hadronization



HERMES main research topics:

- r origin of nucleon spin
 - Iongitudinal spin/momentum structure
 - ransverse spin/momentum structure
- ➡ hadronization/fragmentation

spin-1/2 particles (quarks, leptons, proton, neutron...)

- reference are the fundamental constituents of matter
- are responsible for stability of matter (Pauli-principle)
 - no two spin-1/2 particles can occupy a state with identical quantum numbers

spin and hadronization



HERMES main research topics:

- r origin of nucleon spin
 - Iongitudinal spin/momentum structure
 - ransverse spin/momentum structure
- ➡ hadronization/fragmentation

spin-1/2 particles (quarks, leptons, proton, neutron...)
are the fundamental constituents of matter
are responsible for stability of matter (Pauli-principle)
no two spin-1/2 particles can occupy a state with identical quantum numbers

nucleon properties (charge, momentum, magnetic moment, spin...) should be explained by its constituents

- \blacktriangleright momentum: quarks carry ~ 50 % of the proton momentum
- ➡ spin: total quark spin contribution only ~30%

hac

David Gross' presentation on HEP 2011

NEW PREDICTIONS (10 years)

1. QCD tests & applications will greatly improve, incorporating NLO, NNLO, ... and a theory of fragmentation and hadronization. 2. Atlas and CMS will discover a candidate Higgs particle. 3. There will be convincing evidence for Susy particles. 4. Plans will be underway to build a LC (at Cern) to explore the superworld and the US will join CERN. 5. There will be direct detection of the Dark Matter wind. 6. Alice will see a crossover to the perturbative quark-gluon plasma. 7. Some new Z mesons will be discovered. 8. Gravitational waves and B modes will be observed. 9. String theory will start to be a **theory** with predictions. 10. We will have a plausible explanation of why Λ is so small.

ructure cture

nuc exp

Wigner functions:
$$W^q(\mathbf{k}, \mathbf{b})$$

probability to find a quark in a nucleon with a certain polarization in a position **b** and momentum **k**



 $oldsymbol{s}_q$ $oldsymbol{k}_\perp$

b

 $\int dp W(x,p) = |\psi(x)|^2$ **PRC 72,** Harto(x,p) $\frac{20}{b}(p)|^2$

Wigner functions: $W^q(\mathbf{k}, \mathbf{b})$

probability to find a quark in a nucleon with a certain polarization in a position **b** and momentum **k**



Ami Rostomyan for HERMES

 $m{k}_\perp$

 $\int dp W(x,p) = |\psi(x)|^2$ **PRC 72,** Harto(x.g., 20|b(p)|^2

b

Wigner functions: $W^q(\mathbf{k}, \mathbf{b})$

probability to find a quark in a nucleon with a certain polarization in a position **b** and momentum **k**



Wigner functions: $W^q(\mathbf{k}, \mathbf{b})$

probability to find a quark in a nucleon with a certain polarization in a position **b** and momentum **k**



Wigner functions: $W^q(\mathbf{k}, \mathbf{b})$

probability to find a quark in a nucleon with a certain polarization in a position **b** and momentum **k**



Wigner functions: $W^q(\mathbf{k}, \mathbf{b})$

probability to find a quark in a nucleon with a certain polarization in a position **b** and momentum **k**



The HERMES experiment, located at HERA, with its pure gas targets and advanced particle identification (π , K, p) is well suited for TMD and GPD measurements.

The HERMES experiment, located at HERA, with its pure gas targets and advanced particle identification (π , K, p) is well suited for TMD and GPD measurements.



The HERMES experiment, located at HERA, with its pure gas targets and advanced particle identification (π , K, p) is well suited for TMD and GPD measurements.



- **longitudinal** target polarization (H, D, ³He)
- transverse target polarization (H)
- unpolarized targets: H, D, ⁴He, ¹⁴N, ²⁰Ne, ⁸⁴Kr, ¹³¹Xe
- unpolarized H, D targets with recoil detector

Ami Rostomyan for HERMES

PRC 72, *Hamburg*, 2011

The HERMES experiment, located at HERA, with its pure gas targets and advanced particle identification (π , K, p) is well suited for TMD and GPD measurements.



Ami Rostomyan for HERMES

PRC 72, *Hamburg*, 2011

inclusive measurements

$$\frac{d^2\sigma}{dxdQ^2} = \frac{4\pi\alpha_{em}^2}{Q^4} \left[F_2(x,Q^2) x \left[1 - y - \frac{Q^2}{4E^2} + \frac{y^2 + Q^2/E^2}{2\left[1 + R(x,Q^2)\right]} \right] \quad F_2(x) = x \sum_q e_q^2 f_1^q(x)$$

- HERMES collaboration - JHEP 05 (2011) 126



Ami Rostomyan for HERMES

PRC 72, Hamburg, 2011

- HERMES collaboration - JHEP 05 (2011) 126



6

- HERMES collaboration - JHEP 05 (2011) 126



- HERMES collaboration - JHEP 05 (2011) 126

PRC 72, *Hamburg*, 2011

$$\frac{d^{2}\sigma}{dxdQ^{2}} = \frac{4\pi\alpha_{em}^{2}(F_{2}(x,Q^{2}))}{Q^{4}} \left[1 - y - \frac{Q^{2}}{4E^{2}} + \frac{y^{2} + Q^{2}/E^{2}}{2[1 + R(x,Q^{2})]}\right] F_{2}(x) = x \sum_{q} e_{q}^{2}f_{1}^{q}(x)$$
HERMES data

• cover the transition region between the perturbative and non-perturbative regimes of QCD

• in good agreement with existing data

• provides data in previously uncovered kinematic region

• end to region

• e

Ami Rostomyan for HERMES

- HERMES collaboration - JHEP 05 (2011) 126

$$\frac{d^2\sigma}{dxdQ^2} = \frac{4\pi\alpha_{em}^2}{Q^4} \left[\frac{F_2(x,Q^2)}{x} \left[1 - y - \frac{Q^2}{4E^2} + \frac{y^2 + Q^2/E^2}{2[1 + R(x,Q^2)]} \right] F_2(x) = x \sum_{q} e_q^2 f_1^q(x)$$

$$\frac{HERMES data}{q} F_2 \text{ proton} F_2 \text{ proton} F_2 \text{ proton} F_2 \text{ determine} F_2 \text{ proton} F_2 \text{ determine} F_2 \text{ proton} F_2 \text{ determine} F_2 \text{ determin$$

Ami Rostomyan for HERMES



semi-inclusive measurements

semi-inclusive DIS cross section and TMDs

$$\frac{d^6\sigma}{dx \, dy \, dz \, dP_{h\perp}^2 d\phi \, d\phi_s} \overset{\text{leading}}{\propto} \left\{ F_{UU} + \epsilon \, F_{UU}^{\cos 2\phi} \cos 2\phi \right\} + S_{||} \lambda_e \left\{ \sqrt{1 - \epsilon^2} F_{LL} \right\} + S_{\perp} \left\{ \dots \right\} + \dots$$







$$+ |S_{\perp}| \lambda_{e} \left\{ \sqrt{1 - \varepsilon^{2}} \cos(\phi - \phi_{S}) F_{LT}^{\cos(\phi - \phi_{S})} \right\} +$$

 $g_{1T}^{\perp} \otimes D_1$







HERMES: access to all TMDs thanks to the polarized beam and target

Ami Rostomyan for HERMES

PRC 72, Hamburg, 2011



HERMES: access to all TMDs thanks to the polarized beam and target

Ami Rostomyan for HERMES

PRC 72, Hamburg, 2011



$\sigma_{UU} \propto f_1 \otimes D_1$ $f_1 = \bigcirc$

LO interpretation of multiplicity results (integrated over $\mathbf{P}_{h\perp}$):

$$M^{h} \propto \frac{\sum_{q} e_{q}^{2} \int dx f_{1q}(x, Q^{2}) D_{1q}^{h}(z, Q^{2})}{\sum_{q} e_{q}^{2} \int dx f_{1q}(x, Q^{2})}$$

 \checkmark charge-separated multiplicities of pions and kaons sensitive to the individual quark and antiquark flavors in the fragmentation process



$\sigma_{UU} \propto f_1 \otimes D_1$

LO interpretation of multiplicity results (integrated over $\mathbf{P}_{h\perp}$):



$$M^{h} \propto \frac{\sum_{q} e_{q}^{2} \int dx f_{1q}(x, Q^{2}) D_{1q}^{h}(z, Q^{2})}{\sum_{q} e_{q}^{2} \int dx f_{1q}(x, Q^{2})}$$

 \checkmark charge-separated multiplicities of pions and kaons sensitive to the individual quark and antiquark flavors in the fragmentation process

➡ calculations using DSS and Kretzer
FF fits together with CTEQ6L PDFs

➡ fair agreement for positive hadrons

disagreement for negative hadrons



$\sigma_{UU} \propto f_1 \otimes D_1$

LO interpretation of multiplicity results (integrated over $\mathbf{P}_{h\perp}$):



$$M^{h} \propto \frac{\sum_{q} e_{q}^{2} \int dx f_{1q}(x, Q^{2}) D_{1q}^{h}(z, Q^{2})}{\sum_{q} e_{q}^{2} \int dx f_{1q}(x, Q^{2})}$$

 \checkmark charge-separated multiplicities of pions and kaons sensitive to the individual quark and antiquark flavors in the fragmentation process

➡ calculations using DSS and Kretzer
FF fits together with CTEQ6L PDFs

➡ fair agreement for positive hadrons

disagreement for negative hadrons

✓ inclusion of the data in the future global analyses can give an improved knowledge on FF

Ami Rostomyan for HERMES



PRC 72, *Hamburg*, 2011

fragmentation in nuclear matter



Tuesday, September 27, 11

Tuesday, September 27, 11

20

Ami Rostomyan for HERMES

20

fragmentation in nuclear matter



r typical hadronization length $\propto (1 - z)\nu$ is of the order of nucleus size (1-10 fm)

the time development of the hadronization can be studied using nuclei of increasing size

the struck quark or the $q\bar{q}$ pair propagate through a "cold" nuclear medium

Tuesday, September 27, 11

Ami Rostomyan for HERMES

20

Tuesday, September 27, 11

20

fragmentation in nuclear matter



For typical hadronization length $\propto (1-z)\nu$ is of the order of nucleus size (1-10 fm)

the time development of the hadronization can be studied using nuclei of increasing size

the struck quark or the $q\bar{q}$ pair propagate through a "cold" nuclear medium

➡ interaction signature: reduction of the number of hadrons per DIS event and per nucleon

$$R_{A}^{h}(\nu, Q^{2}, z, p_{t}^{2}) = \frac{\left(\frac{N^{h}(\nu, Q^{2}, z, p_{t}^{2})}{N^{e}(\nu, Q^{2})}\right)_{A}}{\left(\frac{N^{h}(\nu, Q^{2}, z, p_{t}^{2})}{N^{e}(\nu, Q^{2})}\right)_{D}}$$

Tuesday, September 27, 11
fragmentation in nuclear matter



Tuesday, September 27, 11

✓ leptonic probes: well determined^{teday, September 27, 11} energy and momentum transferred to the ²⁰ quark

 \checkmark useful for understanding the fundamental aspects of hadronization

✓ input for calculation of nuclear parton distributions

For typical hadronization length $\propto (1-z)\nu$ is of the order of nucleus size (1-10 fm)

➡ the time development of the hadronization can be studied using nuclei of increasing size

the struck quark or the $q\bar{q}$ pair propagate through a "cold" nuclear medium

➡ interaction signature: reduction of the number of hadrons per DIS eyent and per nucleon

$$R_{A}^{h}(\nu, Q^{2}, z, p_{t}^{2}) = \frac{\left(\frac{N^{h}(\nu, Q^{2}, z, p_{t}^{2})}{N^{e}(\nu, Q^{2})}\right)_{A}}{\left(\frac{N^{h}(\nu, Q^{2}, z, p_{t}^{2})}{N^{e}(\nu, Q^{2})}\right)_{D}}$$

fragmentation in nuclear medium

- Eur. Phys. J. A 47 (2011) 113



fragmentation in nuclear medium

- Eur. Phys. J. A 47 (2011) 113



► Ne data show less pronounced trends compared to Kr and Xe

fragmentation in nuclear medium

- Eur. Phys. J. A 47 (2011) 113



► Ne data show less pronounced trends compared to Kr and Xe

- $\checkmark \pi^+, \pi^-, K^-$: increase of R_A with virtual photon energy v
- \mathbf{k} K⁺: clear increase of R_A with v for the lowest z-slice and flatter behavior for higher z
- $rac{p}{r}$: weak v-dependence

 $rac{r}{r}$ p: R_A exceeding unity at higher values of v and low z (apart from hadronization, different production mechanisms contribute) Ami Rostomyan for HERMES *PRC* 72, *Hamburg*, 2011

12





$$H \qquad D$$

$$A_{1p}^{h^+ - h^-} = \frac{4\Delta u_v(x) - \Delta d_v(x)}{4u_v(x) - d_v(x)} \qquad A_{1d}^{h^+ - h^-} = \frac{\Delta u_v(x) + \Delta d_v(x)}{u_v(x) + d_v(x)}$$





assumption: charge conjugation symmetry in fragmentation: $D_q^h(z) = D_{\bar{q}}^h(z)$

➡ cancellation of fragmentation functions in the charge difference asymmetry



0.02

0.1

0.2 0.3

0.01

$\sigma_{LL} \propto g_{1L} \otimes D_1$



► <u>assumption</u>: charge conjugation symmetry in fragmentation: $D_q^h(z) = D_{\bar{q}}^h(z)$

cancellation of fragmentation functions in the charge difference asymmetry

Ami Rostomyan for HERMES

0.2 0.3

X

0.1

0.02

0.01







$\sigma_{LL} \propto g_{1L} \otimes D_1$



► <u>assumption</u>: charge conjugation symmetry in fragmentation: $D_q^h(z) = D_{\bar{q}}^h(z)$

 cancellation of fragmentation functions in the charge difference asymmetry

these equations can be solved for the flavor separated polarized valence quark densities



Ami Rostomyan for HERMES





✓ good agreement between results from current analysis and purity method (- HERMES collaboration - *Phys. Rev. D* 71 (2005) 012003) involving different assumptions

$\sigma_{LL} \propto g_{1L} \otimes D_1$



► <u>assumption</u>: charge conjugation symmetry in fragmentation: $D_q^h(z) = D_{\bar{q}}^h(z)$

 cancellation of fragmentation functions in the charge difference asymmetry

these equations can be solved for the flavor separated polarized valence quark densities



PRC 72, Hamburg, 2011

 $\sigma_{UU} \propto h_1^\perp \otimes H_1^\perp$





$$\sigma_{UU} \propto h_1^\perp \otimes H_1^\perp$$



Ami Rostomyan for HERMES



$$\sigma_{UU} \propto h_1^\perp \otimes H_1^\perp$$



14



$$\sigma_{UU} \propto h_1^\perp \otimes H_1^\perp$$



close to submission

exclusive measurements



theoretically the cleanest probe of GPDs $\gamma^* N \to \gamma N : H, E, \widetilde{H}, \widetilde{E}$





theoretically the cleanest probe of GPDs $\gamma^* N \to \gamma N : H, E, \widetilde{H}, \widetilde{E}$









• theoretically the cleanest probe of GPDs $\gamma^* N \to \gamma N : H, E, \widetilde{H}, \widetilde{E}$





✓ HERMES measured complete set of beam helicity, beam charge and target polarization asymmetries



• theoretically the cleanest probe of GPDs $\gamma^* N \to \gamma N : H, E, \widetilde{H}, \widetilde{E}$





✓ HERMES measured complete set of beam helicity, beam charge and target polarization asymmetries

unpolarized target

$$F\mathcal{H} + \frac{x_B}{2 - x_B} (F_1 + F_2) \widetilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$

$$= \text{ longitudinally polarized target} \frac{x_B}{2 - x_B} (F_1 + F_2) \left(\mathcal{H} + \frac{x_B}{2} \mathcal{E} \right) + F_1 \widetilde{\mathcal{H}} - \frac{x_B}{2 - x_B} \left(\frac{x_B}{2} F_1 + \frac{t}{4M^2} F_2 \right) \widetilde{\mathcal{E}}$$

transversely polarized target

$$\frac{t}{4M^2} \left[(2-x_B)F \mathcal{E} - 4\frac{1-x_B}{2-x_B}F_2 \mathcal{H} \right]$$



• theoretically the cleanest probe of GPDs $\gamma^* N \to \gamma N : H, E, \widetilde{H}, \widetilde{E}$



 \checkmark need information about GPDs E and H

$$J_q = \frac{1}{2} \lim_{t \to 0} \int_{-1}^{1} dx \, x \left[H_q(x,\xi,t) + E_q(x,\xi,t) \right]$$



✓HERMES measured complete set of beam helicity, beam charge and target polarization asymmetries

unpolarized target

$$F\mathcal{H} + \frac{x_B}{2 - x_B} (F_1 + F_2) \widetilde{\mathcal{H}} - \frac{t}{4M^2} F_2 \mathcal{E}$$

$$= \text{ longitudinally polarized target}$$

$$= \frac{x_B}{2 - x_B} (F_1 + F_2) \left(\mathcal{H} + \frac{x_B}{2} \mathcal{E} \right)$$

$$+ F_1 \widetilde{\mathcal{H}} - \frac{x_B}{2 - x_B} \left(\frac{x_B}{2} F_1 + \frac{t}{4M^2} F_2 \right) \widetilde{\mathcal{E}}$$

transversely polarized target

$$\frac{t}{4M^2} \left[(2 - x_B)F\mathcal{E} - 4\frac{1 - x_B}{2 - x_B}F_2\mathcal{H} \right]$$

PRC 72, *Hamburg*, 2011



unique and complete set of DVCS asymmetries

$$\gamma^* N \to \gamma N : F(H, E, \widetilde{H}, \widetilde{E})$$
$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx \, C_q(\xi, x) \, F^q(x, \xi, t)$$



unique and complete set of DVCS asymmetries

(pre-recoil publications) - JHEP 11 (2009) 083 -- Nucl. Phys. B 829 (2010) 1-27 -

- JHEP 11 (2009) 083 -- Nucl. Phys. B 829 (2010) 1-27- $\gamma^* N \to \gamma N : F(H, E, \widetilde{H}, \widetilde{E})$ $\mathcal{F}(\xi,t) = \sum_{\alpha} \int_{-1}^{1} dx \, C_q(\xi,x) \, F^q(x,\xi,t)$

- JHEP 06 (2008) 066 -



-JHEP 06 (2010) 019 -- Nucl. Phys. B 842 (2011) 265-298 -

- JHEP 06 (2010) 019 -- Nucl. Phys. B 842 (2011) 265-298 -

Ami Rostomyan for HERMES



(pre-recoil data)

double -spin asymmetries

$$\sigma(\phi, \phi_s, e_{\ell}, S_{\perp}, \lambda_l) = \sigma_{\mathrm{UU}}(\phi) \left\{ 1 + e_{\ell} \mathcal{A}_C(\phi) + \lambda_l \mathcal{A}_{LU}^{DVCS}(\phi) + e_{\ell} \lambda_l \mathcal{A}_{LU}^{I}(\phi) + S_{\perp} \mathcal{A}_{UT}^{DVCS}(\phi, \phi_S) + e_{\ell} S_{\perp} \mathcal{A}_{UT}^{I}(\phi, \phi_S) + \lambda_l S_{\perp} \mathcal{A}_{LT}^{BH+DVCS}(\phi, \phi_S) + e_{\ell} \lambda_l S_{\perp} \mathcal{A}_{LT}^{I}(\phi, \phi_S) \right\}$$









the due to different kinematic pre-factors, this amplitude is suppressed

Ami Rostomyan for HERMES



18

• $A_{LT,I}^{\sin(\phi-\phi_s)\sin\phi}$ could provide a similar constraint to the real part of Re \mathcal{E} • due to different kinematic pre-factors, this amplitude is suppressed

 \checkmark nevertheless, may serve as additional constraints in global fits



before PRC 71

- Nucl. Phys. B842 (2011) 265-298
- Eur. Phys. J. C 71 (2011) 1609

after PRC 71

- JHEP 05 (2011) 126
- Phys. Lett. B 704 (2011) 15-23
- Eur. Phys. J. A 47 (2011) 113

before PRC 71

- Nucl. Phys. B842 (2011) 265-298
- Eur. Phys. J. C 71 (2011) 1609

after PRC 71

- JHEP 05 (2011) 126
- Phys. Lett. B 704 (2011) 15-23
- Eur. Phys. J. A 47 (2011) 113

- ▶ in total 5 publications so far in 2011
- ► 3 more publications are expected
- ✤ 8 more in circulation within collaboration

before PRC 71

- Nucl. Phys. B842 (2011) 265-298
- Eur. Phys. J. C 71 (2011) 1609

after PRC 71

- JHEP 05 (2011) 126
- Phys. Lett. B 704 (2011) 15-23
- Eur. Phys. J. A 47 (2011) 113

- ➡ in total 5 publications so far in 2011
- ► 3 more publications are expected
- ✤ 8 more in circulation within collaboration



before PRC 71

- Nucl. Phys. B842 (2011) 265-298
- Eur. Phys. J. C 71 (2011) 1609

after PRC 71

- *JHEP 05 (2011) 126*
- Phys. Lett. B 704 (2011) 15-23
- Eur. Phys. J. A 47 (2011) 113

- in total 5 publications so far in 2011
- ► 3 more publications are expected
 - ✤ 8 more in circulation within collaboration



average citations per paper: 69

HERMES clearly in publication phase

before PRC 71

- Nucl. Phys. B842 (2011) 265-298
- Eur. Phys. J. C 71 (2011) 1609

after PRC 71

- JHEP 05 (2011) 126
- Phys. Lett. B 704 (2011) 15-23
- Eur. Phys. J. A 47 (2011) 113

- ➡ in total 5 publications so far in 2011
- ► 3 more publications are expected
 - 8 more in circulation within collaboration



average citations per paper: 69

HERMES clearly in publication phase

- major conferences (PANIC, HEP, EINN'11 etc.): 17
- major workshops (DIS, EDS-Blois, etc.): 30+

data preservation
status of data preservation

✓ HERMES participates in DPHEP initiative





HERMES has been the pioneering collaboration in TMD and GPD fields
still very important player in the field of nucleon (spin) structure

- ▶ polarized e^{+/-} beams
- pure gas target

- good particle identification
- recoil detector

Ami Rostomyan for HERMES



backup slides



GPD H: unpolarized hydrogen target

(recoil data)

$$\sigma(\phi, P_{\ell}, e_{\ell}) = \sigma_{UU}(\phi) \times \left[1 + P_{\ell} \mathcal{A}_{LU}^{DVCS}(\phi) + e_{\ell} \mathcal{A}_{LU}^{I}(\phi) + e_{\ell} \mathcal{A}_{C}(\phi)\right]$$
$$\mathcal{A}_{LU}(\phi) \simeq \sum_{n=1}^{2} \mathcal{A}_{LU}^{\sin(n\phi)} \sin(n\phi)$$

➡ extraction of single-charge beam-helicity asymmetry amplitudes for elastic data sample (background < 0.1%)



► indication for slightly larger magnitude of the leading amplitude for elastic process compared the one in the recoil detector acceptance

PRC 72, *Hamburg*, 2011