

Physics beyond the standard model II

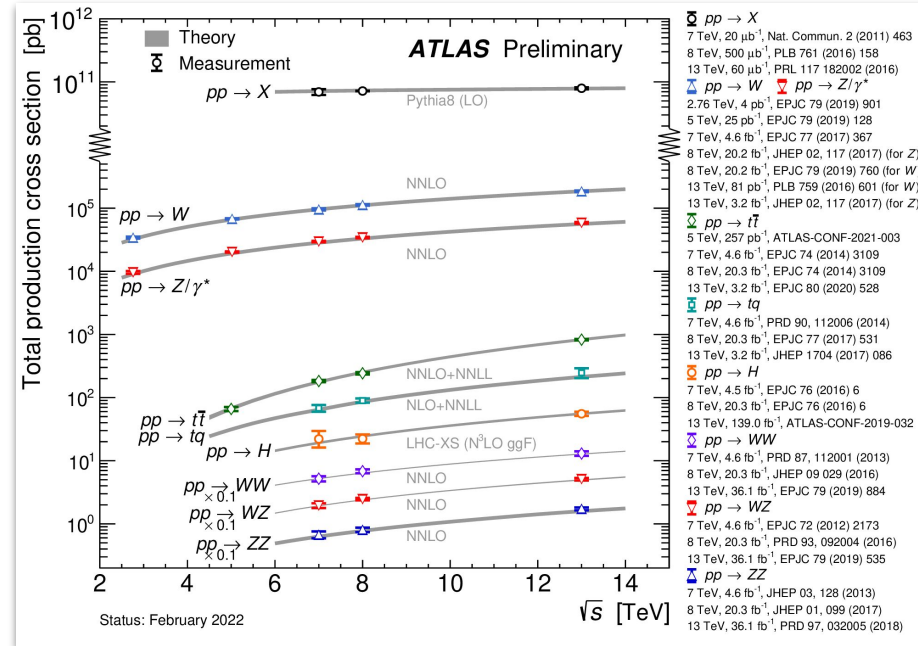
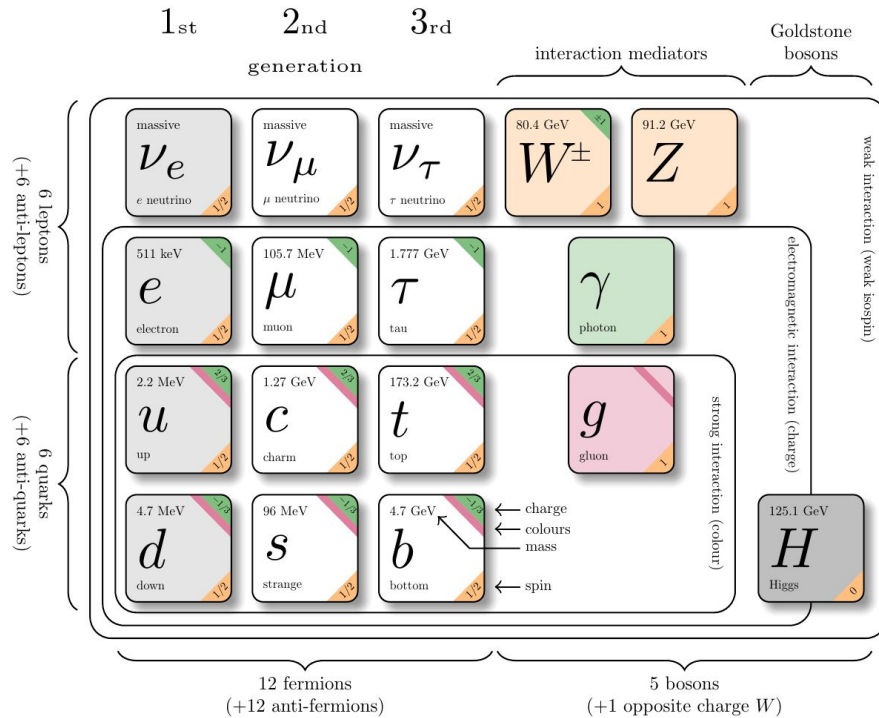
Ben Brüers

DESY Zeuthen, 05.09.2025

Based on [Elisa Pueschel's lecture](#) from 2019



Last time...



<http://cds.cern.ch/record/2804061>

The Standard model is very precise, but...

Last time (2)...

- ...there are a bunch of things it cannot explain:

9 fermion masses ($m_u, m_d, m_c, m_s, m_b, m_t, m_e, m_\mu, m_\tau$)

- + 2 Higgs boson parameters: the mass & VEV (m_H, v)
- + 3 coupling parameters (g_W, g', g_s)
- + 4 CKM parameters (3 mixing angles + 1 CP violating phase)
- + 1 CP violating phase in QCD (see later)

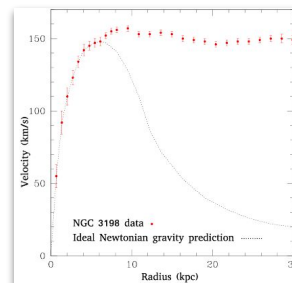
19 free parameters

$$\delta M_H^2 = \frac{G_F \Lambda^2}{4\pi^2 \sqrt{2}} (6M_W^2 + 3M_Z^2 + M_H^2 - 12m_t^2)$$

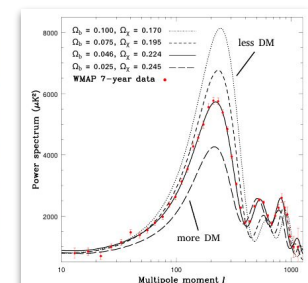
$$\delta M_H^2|_{t\text{-loop}} \approx -\frac{3G_F}{\pi^2 \sqrt{2}} m_t^2 \Lambda^2 \approx -0.075 \Lambda^2$$

“bare mass” tuned very finely, $O(10^{-4})$ - $O(10^{-34})$ GeV!

$$\mathcal{L}_{\text{QCD}} = \underbrace{\bar{\psi}_i (i\gamma^\mu (D_\mu)_{ij}) \psi_j}_{\text{quark dynamics}} - \underbrace{m \delta_{ij}}_{\text{quark mass}} \psi_j - \underbrace{\frac{1}{4} G_{\mu\nu}^a G_a^{\mu\nu}}_{\text{gluon dynamics}} - \underbrace{\bar{\Theta} \frac{\alpha_s}{8\pi} G^{\mu\nu a} \tilde{G}_{\mu\nu}^a}_{\text{CP-violating term}}$$



K. Begeman, Astron. Astrophys. 223 (1989), pp. 47–60

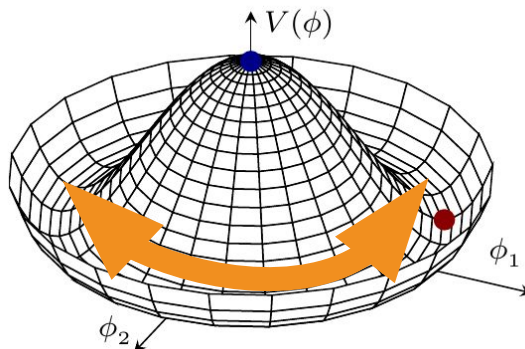


arxiv:1001.4635 arxiv:1006.2483

Last time (3)...

- Discussed multiple SM extension:

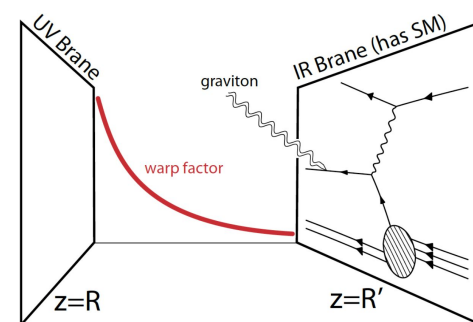
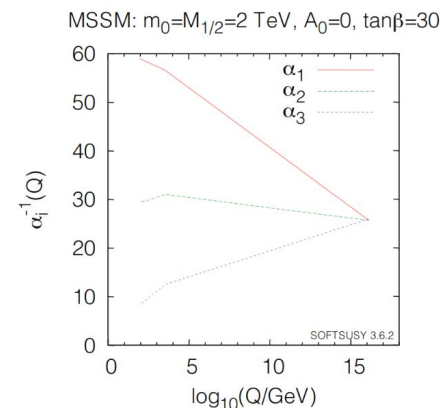
- Axions + ALPs
- 2HDM(+a)
- SUSY
- GUTs
- Extra dimensions



Drawing inspired by
https://indico.soc.kit.edu/event/477/contributions/4854/attachments/2575/3683/KSETA_Durbach_2019_pargner.pdf

$$Q |\text{fermion}\rangle = |\text{boson}\rangle$$

$$Q |\text{boson}\rangle = |\text{fermion}\rangle$$

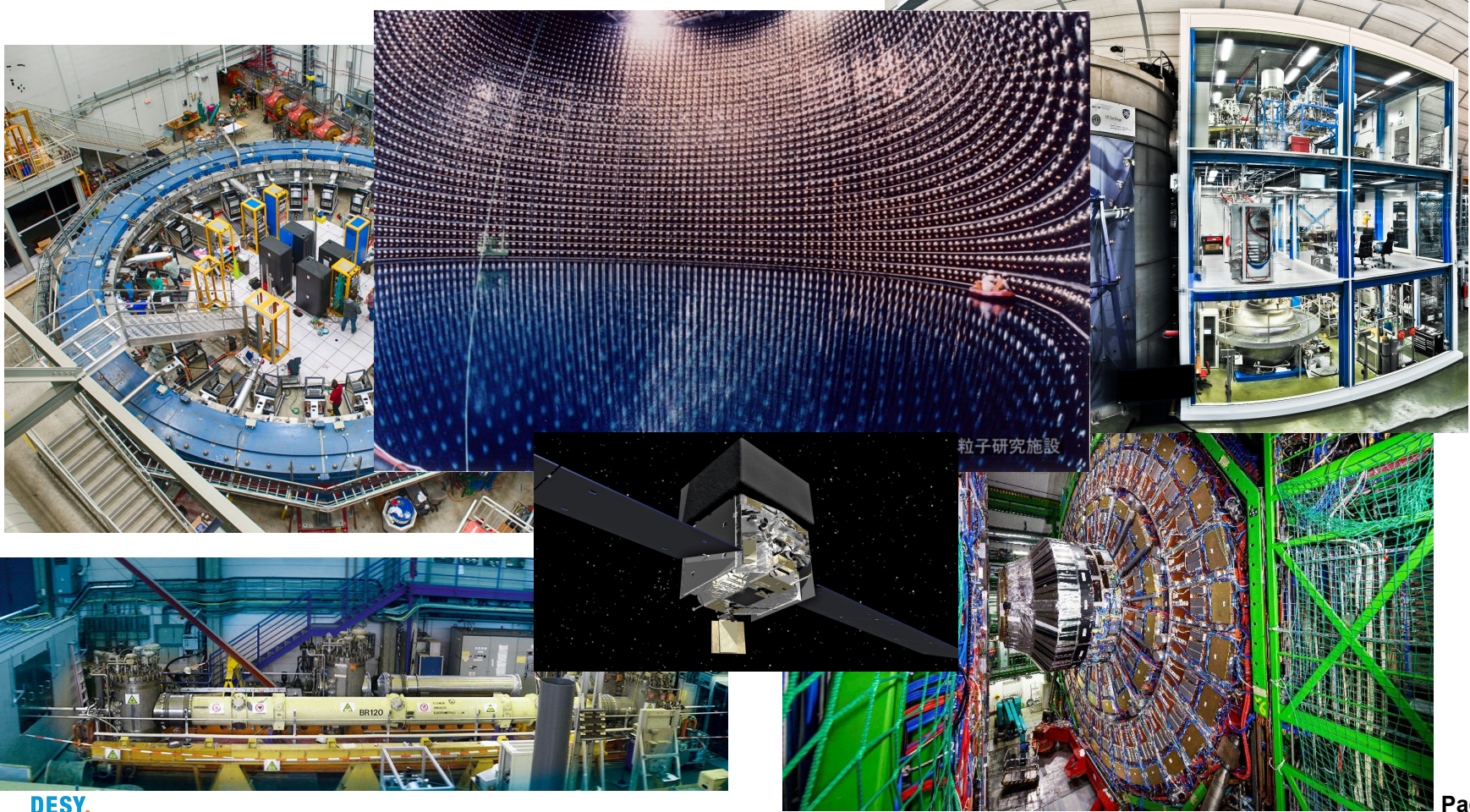


https://www.zeuthen.desy.de/students/2019/lectures/Pleschel_Lecture1_BeyondSM_2019.pdf

Model	Dark Matter	Hierarchy problem	Strong CP problem	Unification	Gravity
Axions	✓	-	✓	-	-
2HDM	✓	-	-	-	-
SUSY	✓	✓	possible	✓	e.g. mSUGRA
GUTs	-	-	-	✓	-
Extra dims.	✓	✓	-	possible	✓

Today:

- Will discuss experiments searching for physics beyond the Standard Model
- Will cover collider-based experiments and others



The LHC

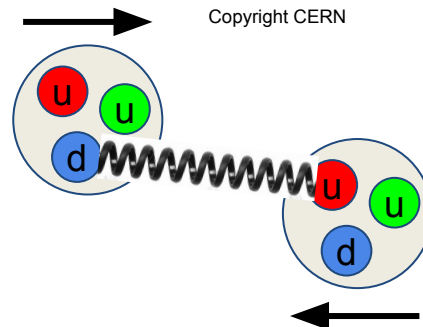
LHC

→ circumference
27 km

→ proton-proton
collisions with a
CME of
 $\sqrt{s} = 13 \text{ TeV}$

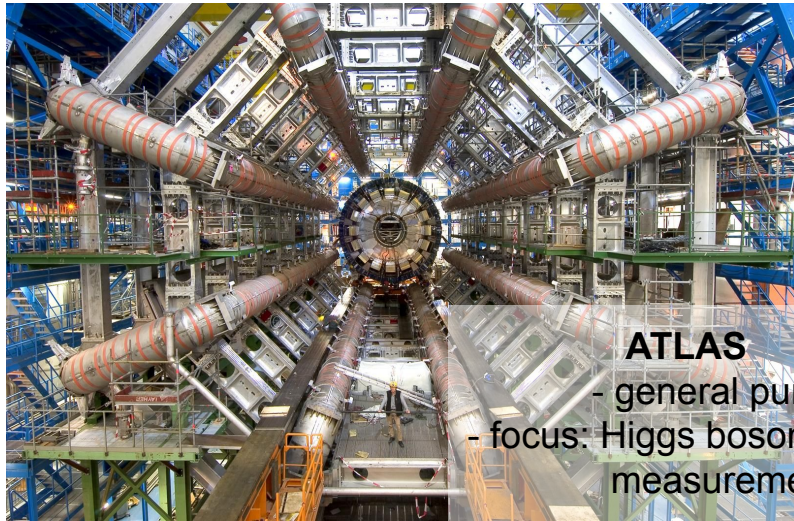
→ interacting
particles are
quarks, so their
CME often
smaller

→ detectors
surround points
of collision



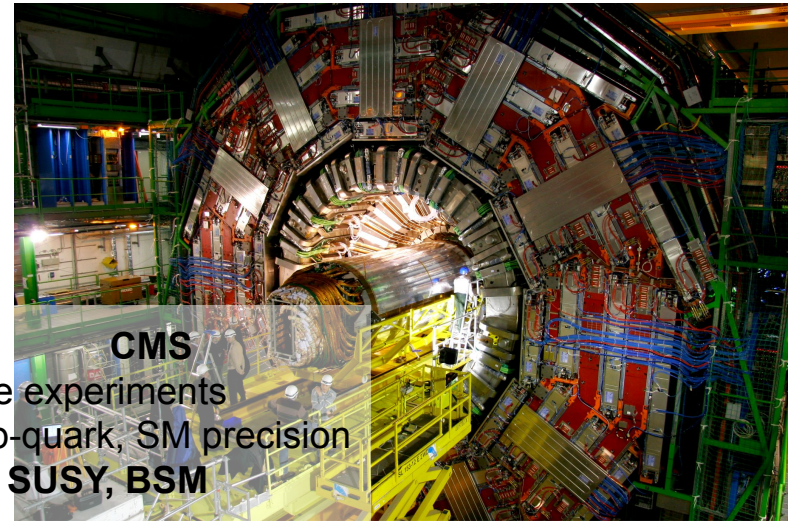
<https://previews.123rf.com/images/denisnata/denisnata1003/denisnata100300048/6555910-black-spiral-telescope-cable-isolated-on-white-background.jpg>

Experiments at the LHC

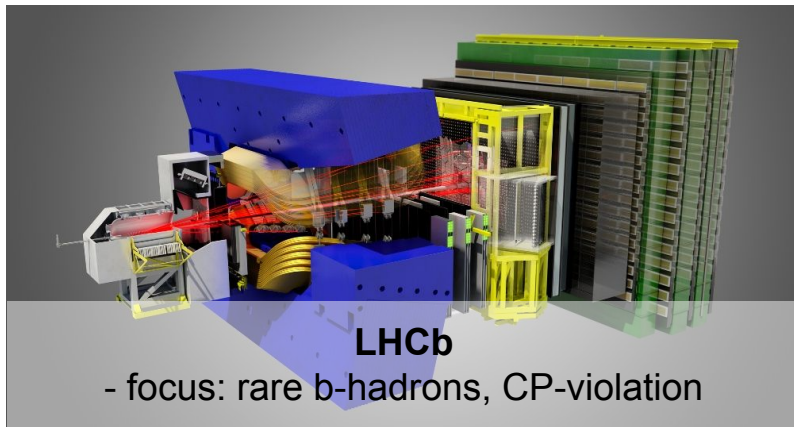


ATLAS & CMS
- general purpose experiments
- focus: Higgs boson, top-quark, SM precision measurements, **SUSY, BSM**

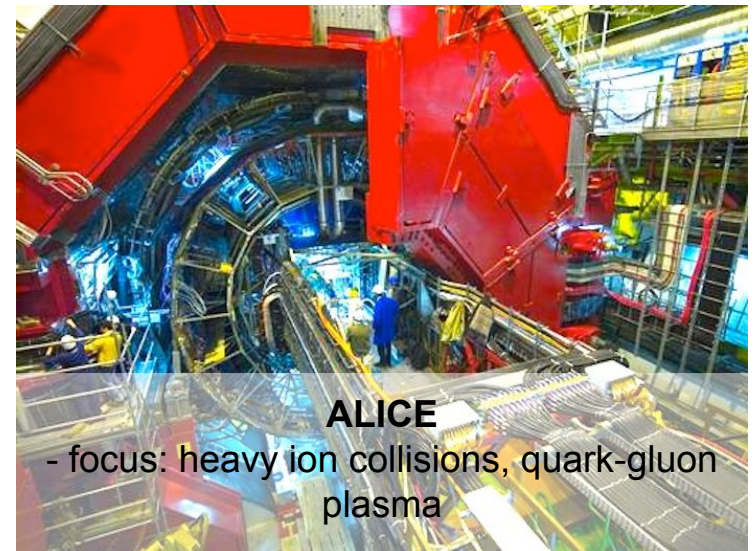
https://www.weltmaschine.de/sites/sites_custom/site_weltmaschine/content/e28861/e36564/e36588/e36608/0511013_01-A4-at-144-dpi.jpg



https://cms.cern/sites/default/files/field/image/cds-record-1275108-hoch-20071215_721-nice.jpg



LHCb
- focus: rare b-hadrons, CP-violation

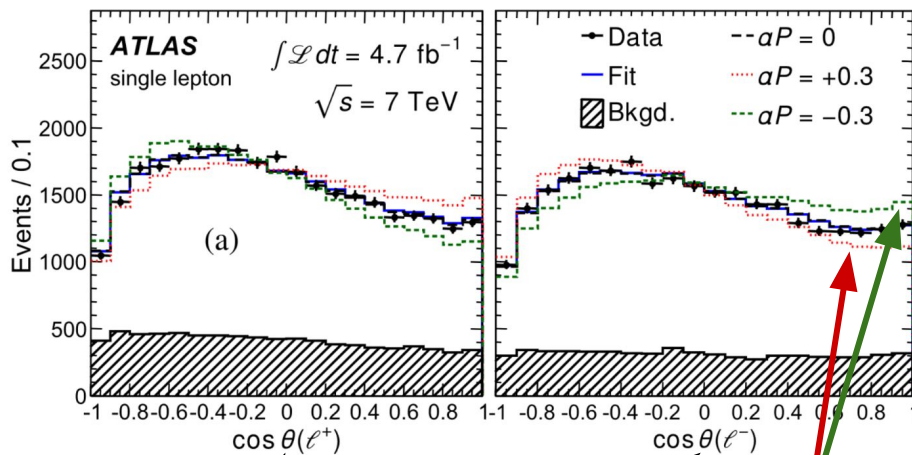


ALICE
- focus: heavy ion collisions, quark-gluon plasma

How to search for BSM physics at colliders

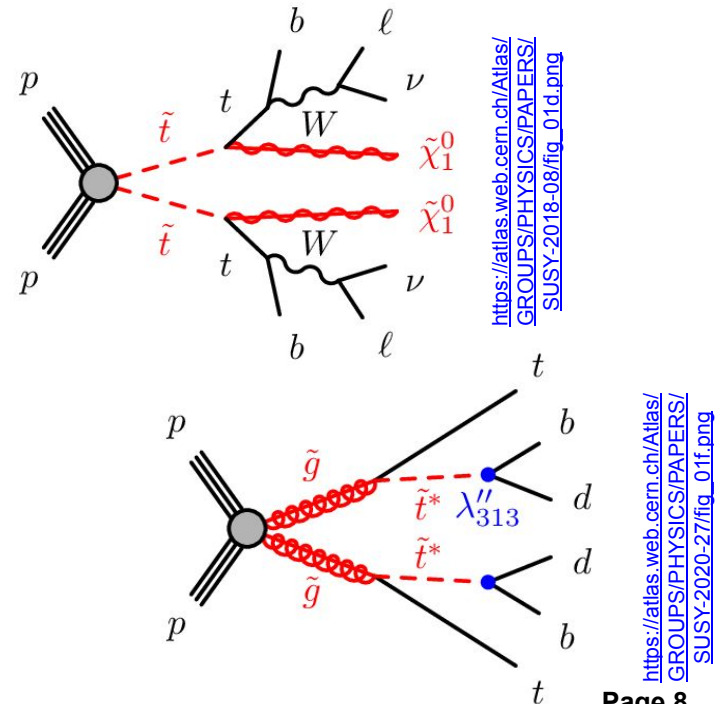
- BSM physics must be rare and/or involve heavy particles (else discovered)
- Two search principles:
 - BSM physics slightly modifies masses, couplings, etc.
→ investigate these effects in precision measurements
 - **BSM physics particles produced in proton collisions**
→ **search for the particles / reconstruct their decays, etc.**
→ **ideally investigate signatures that are rare in the SM**

[doi:10.1103/PhysRevLett.111.232002](https://doi.org/10.1103/PhysRevLett.111.232002)



polar angle of lepton w.r.t. beam axis

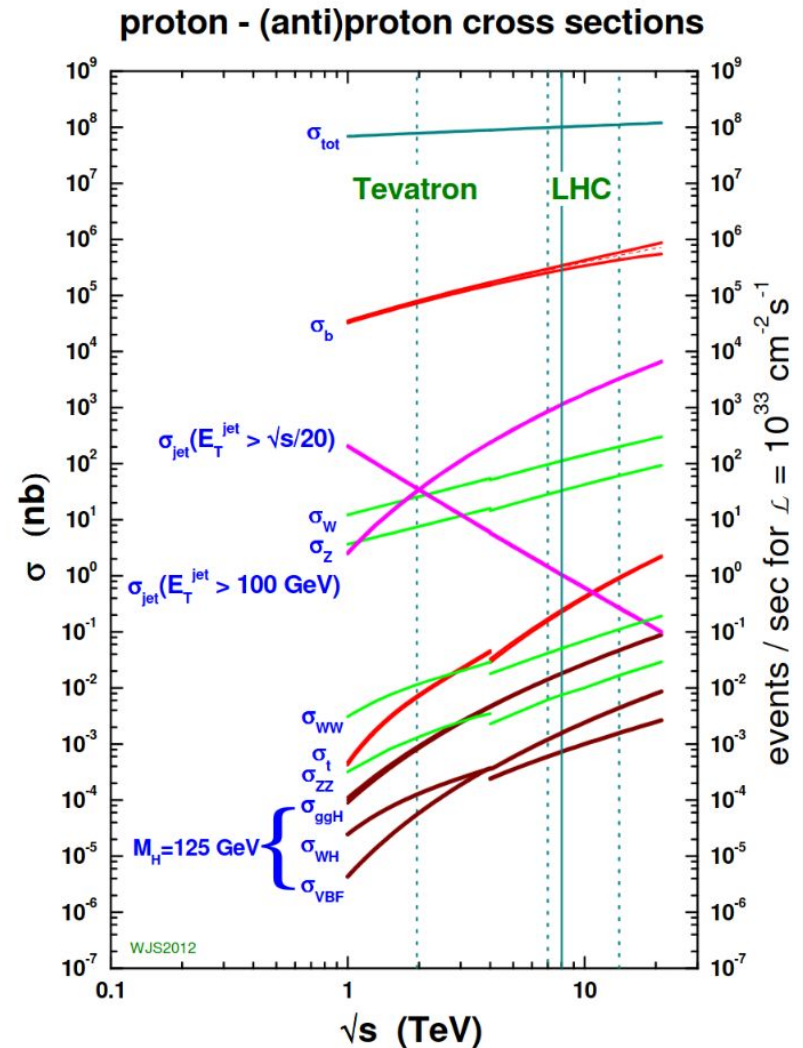
with BSM physics



Metrics to keep in mind

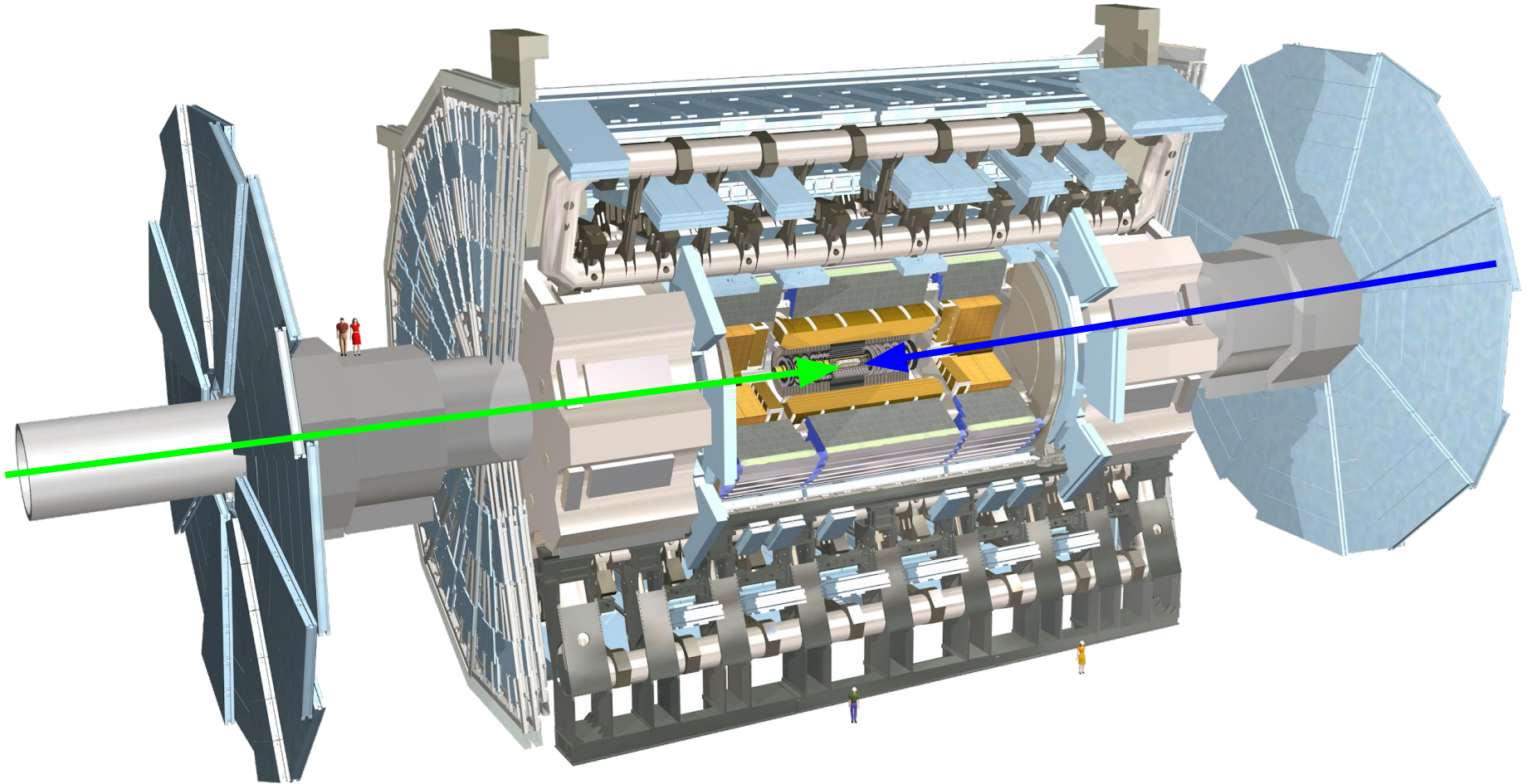
- BSM physics must be rare and/or involve heavy particles (else discovered)
- Cross-section must not be too low, else will not be able to discover anything
→ **important model metric**
- Too investigate very rare signatures need many collision events
- The mass of the particles must not be heavier than CME, else not produced
- Can only detect particles produced in area of detector (detector not infinite) and if detector sufficiently efficient
→ **important collider metrics: CME, luminosity, detector efficiency**

$$N_{\text{obs}} = \int \mathcal{L}(t) \sigma \varepsilon A dt$$



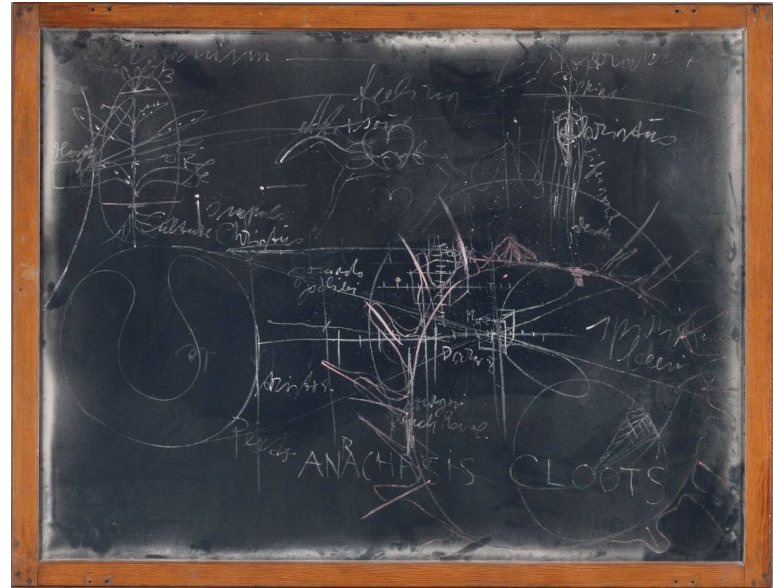
http://www.hep.ph.ic.ac.uk/~wstirling/plots/crosssections2012_v5.pdf

The ATLAS detector



How do we measure particles with the ATLAS detector?

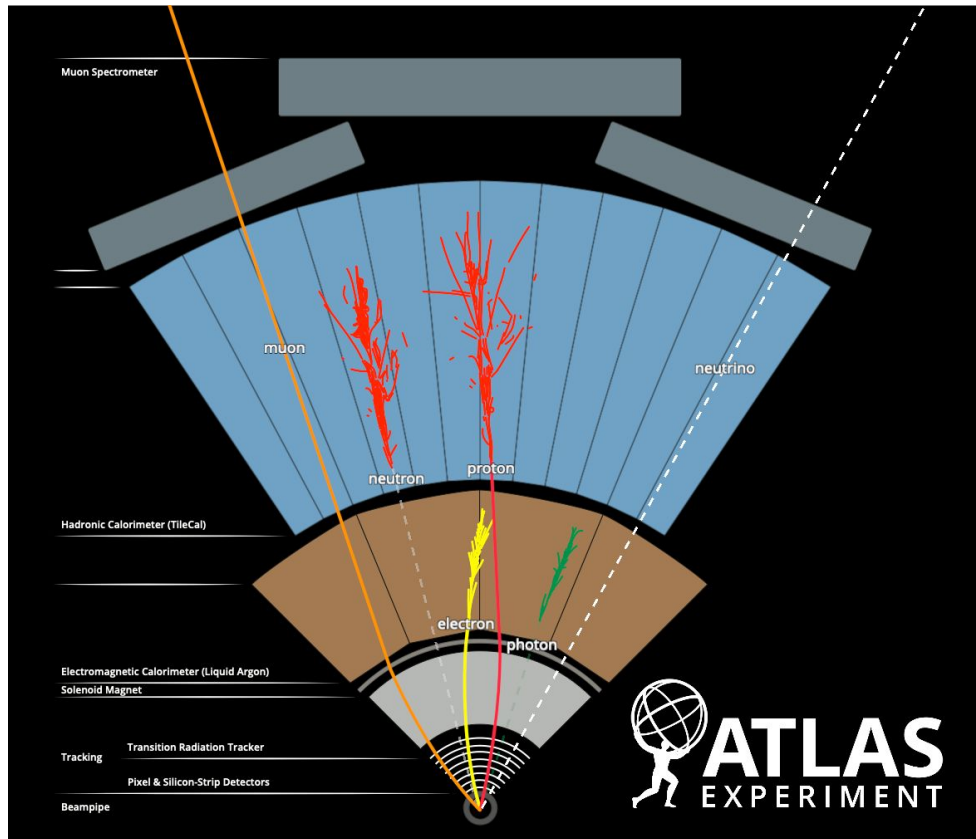
It's blackboard time!



<https://cameo.mfa.org/images/b/ba/2000.979-CR9834-d1.jpg>

(summary on the next slide)

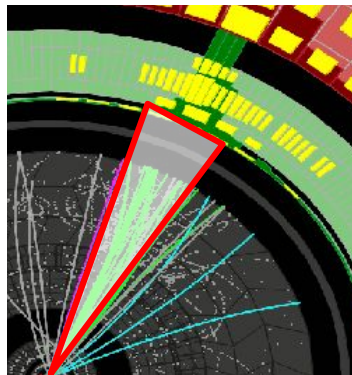
Particle reconstruction: ATLAS detector



- Multi-layer detector, measure particles via their interaction with the detector
- Inner-most detector: measure tracks of charged particles (electron, muon, charged hadron)
- ECAL / HCAL: measure energy of EM-interacting / hadronically interacting particle
→ concept: make particles lose all their energy and measure the loss
- Muon chambers for high precision muon momentum measurements
- NB: due to quark confinement, quarks cannot exist alone
→ form collimated hadron sprays, we call these “jets”

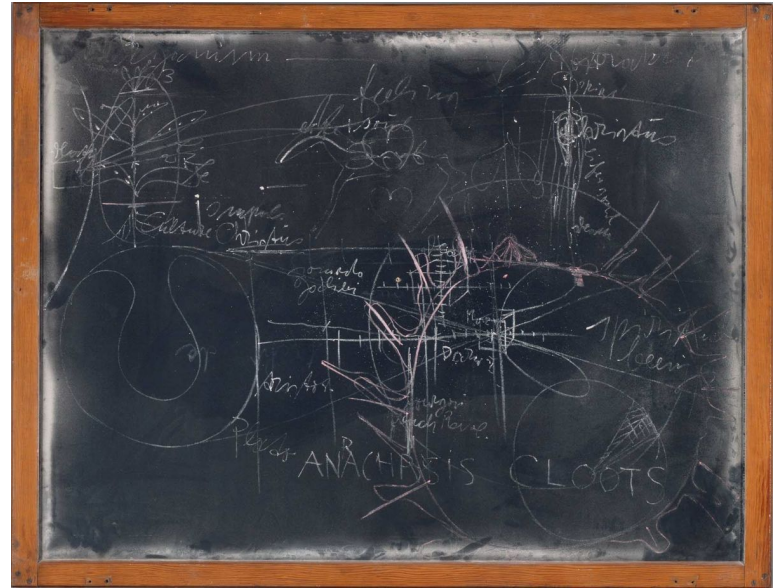
q → had. 1
had. 2
had. 3
...
had. n

confinement



What about invisible particles?

It's blackboard time!



<https://cameo.mfa.org/images/b/ba/2000.979-CR9834-d1.jpg>

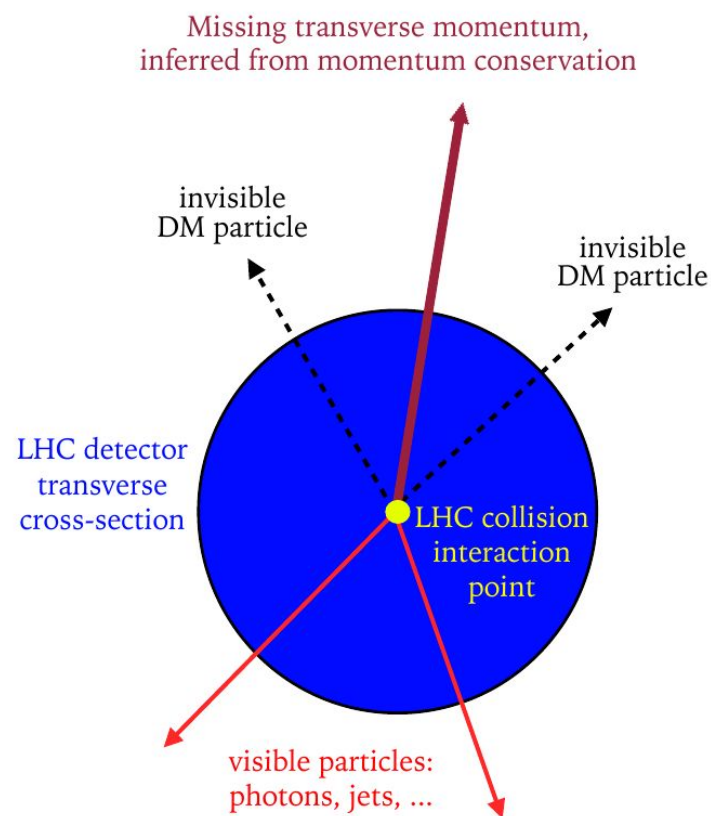
(summary on the next slide)

How do we measure invisible particles?

- Neutrinos and potential BSM particles interact weakly with detector→no signal
- But: if produced with particles that produce a signal: can use **momentum conservation** to **infer** on them
- Protons collide heads on: momentum of interacting particles in direction orthogonal to the beam axis is ~ 0
- Momentum is conserved → momenta of all particles in the plane transverse to the beam axis must sum to zero
→ can infer on the total momentum of the particles escaping detection and the direction of it in the transverse plane
→ **this is called Missing Transverse Momentum** and the magnitude **Missing Transverse Energy (MET)**

$$0 = \sum_{\text{all}} \vec{p}_{T,i} = \sum_{\text{visible}} \vec{p}_{T,i} + \sum_{\text{in-visible}} \vec{p}_{T,j}$$

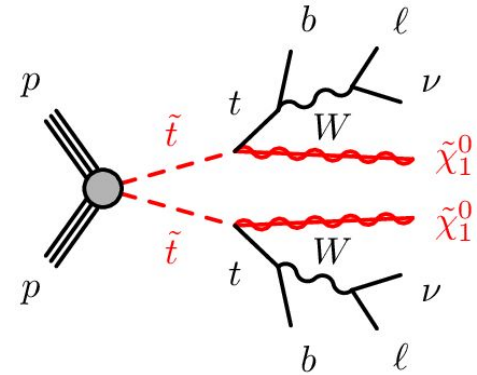
Missing transverse momentum



How do we search for the BSM particles?

- In the data, select those collision events with:

- 2 jets compatible with a bottom-quark
- 2 leptons
- MET compatible with $2 \times \tilde{\chi}_1^0$ and $2 \times \nu$
- Maybe an invariant mass requirement on the system of b - W - $\tilde{\chi}_1^0$



https://atlas.web.cern.ch/Atlas/GROUPS/PHYSICS/PAPERS/SUSY-2018-08/fig_01d.png

- **Caveat: we might not have saved these events in ATLAS!!**

- Protons collide every 25 ns in ATLAS
→ 1 Mb of data / event or 40 Tb of data per second!!
→ cannot store this!! & most not of interest (low energy)
- Solution: coarsely analyse all collision event as they happen
- Only store events fulfilling certain criteria, e.g. 6 jets, 2 leps → “triggering”
- **So: the BSM physics you can find depends on your trigger criteria!**



Triggering criteria in ATLAS (in 2018)

- Lots of triggers there! But sometimes need to add one...

[ATL-DAQ-PUB-2019-001/](https://atlas.cern/DAQ-PUB-2019-001/)

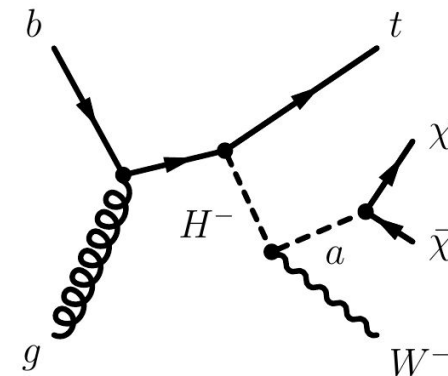
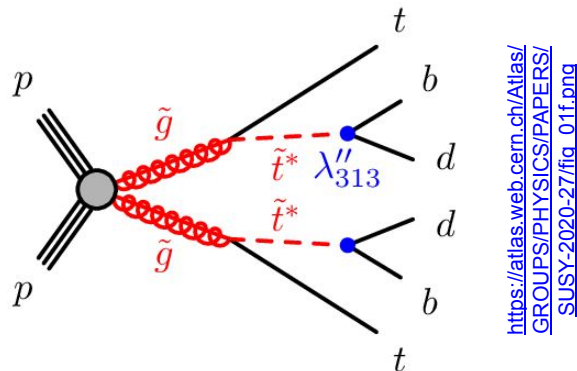
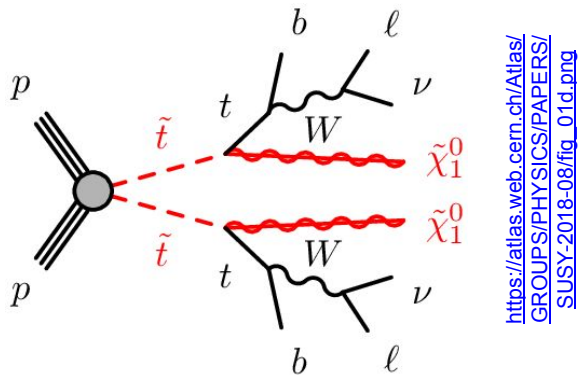
Trigger	Typical offline selection	Trigger Selection		L1 Peak Rate [kHz]	HLT Peak Rate [Hz]
		L1 [GeV]	HLT [GeV]	$L=2.0 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$	
Single leptons	Single isolated μ , $p_T > 27$ GeV	20	26 (i)	16	218
	Single isolated tight e , $p_T > 27$ GeV	22 (i)	26 (i)	31	195
	Single μ , $p_T > 52$ GeV	20	50	16	70
	Single e , $p_T > 61$ GeV	22 (i)	60	28	20
	Single τ , $p_T > 170$ GeV	100	160	1.4	42
Two leptons	Two μ , each $p_T > 15$ GeV	2×10	2×14	2.2	30
	Two μ , $p_T > 23, 9$ GeV	20	22, 8	16	47
	Two very loose e , each $p_T > 18$ GeV	2×15 (i)	2×17	2.0	13
	One e & one μ , $p_T > 8, 25$ GeV	20 (μ)	7, 24	16	6
	One loose e & one μ , $p_T > 18, 15$ GeV	15, 10	17, 14	2.6	5
	One e & one μ , $p_T > 27, 9$ GeV	22 (e, i)	26, 8	21	4
	Two τ , $p_T > 40, 30$ GeV	20 (i), 12 (i) (+jets, topo)	35, 25	5.7	93
	One τ & one isolated μ , $p_T > 30, 15$ GeV	12 (i), 10 (+jets)	25, 14 (i)	2.4	17
Three leptons	One τ & one isolated e , $p_T > 30, 18$ GeV	12 (i), 15 (i) (+jets)	25, 17 (i)	4.6	19
	Three very loose e , $p_T > 25, 13, 13$ GeV	$20, 2 \times 10$	$24, 2 \times 12$	1.6	0.1
	Three μ , each $p_T > 7$ GeV	3×6	3×6	0.2	7
	Three μ , $p_T > 21, 2 \times 5$ GeV	20	$20, 2 \times 4$	16	9
	Two μ & one loose e , $p_T > 2 \times 11, 13$ GeV	2×10 (μ)	$2 \times 10, 12$	2.2	0.5
Signle photon	Two loose e & one μ , $p_T > 2 \times 13, 11$ GeV	$2 \times 8, 10$	$2 \times 12, 10$	2.3	0.1
	One loose γ , $p_T > 145$ GeV	24 (i)	140	24	47
Two photons	Two loose γ , each $p_T > 55$ GeV	2×20	2×50	3.0	7
	Two γ , $p_T > 40, 30$ GeV	2×20	35, 25	3.0	21
	Two isolated tight γ , each $p_T > 25$ GeV	2×15 (i)	2×20 (i)	2.0	15
Single jet	Jet ($R = 0.4$), $p_T > 435$ GeV	100	420	3.7	35
	Jet ($R = 1.0$), $p_T > 480$ GeV	111 (topo: $R = 1.0$)	460	2.6	42
	Jet ($R = 1.0$), $p_T > 450$ GeV, $m_{\text{jet}} > 45$ GeV	111 (topo: $R = 1.0$)	$420, m_{\text{jet}} > 35$	2.6	36
b -jets	One b ($\epsilon = 60\%$), $p_T > 285$ GeV	100	275	3.6	15
	Two b ($\epsilon = 60\%$), $p_T > 185, 70$ GeV	100	175, 60	3.6	11
	One b ($\epsilon = 40\%$) & three jets, each $p_T > 85$ GeV	4×15	4×75	1.5	14
	Two b ($\epsilon = 70\%$) & one jet, $p_T > 65, 65, 160$ GeV	$2 \times 30, 85$	$2 \times 55, 150$	1.3	17
	Two b ($\epsilon = 60\%$) & two jets, each $p_T > 65$ GeV	$4 \times 15, \eta < 2.5$	4×55	3.2	15
Multijets	Four jets, each $p_T > 125$ GeV	3×50	4×115	0.5	16
	Five jets, each $p_T > 95$ GeV	4×15	5×85	4.8	10
	Six jets, each $p_T > 80$ GeV	4×15	6×70	4.8	4

Model dependence of collider searches

- Disadvantage of triggering: need to know what you are looking for...
- Need to select a model to design your search!
→ BSM searches at collider depend on the model you are investigating!!!!

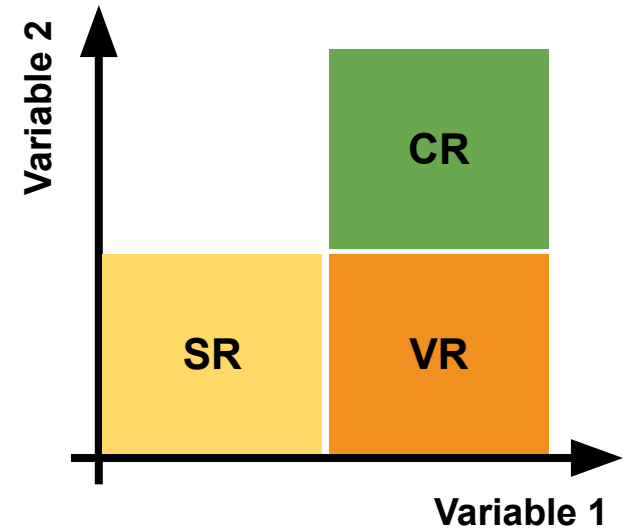
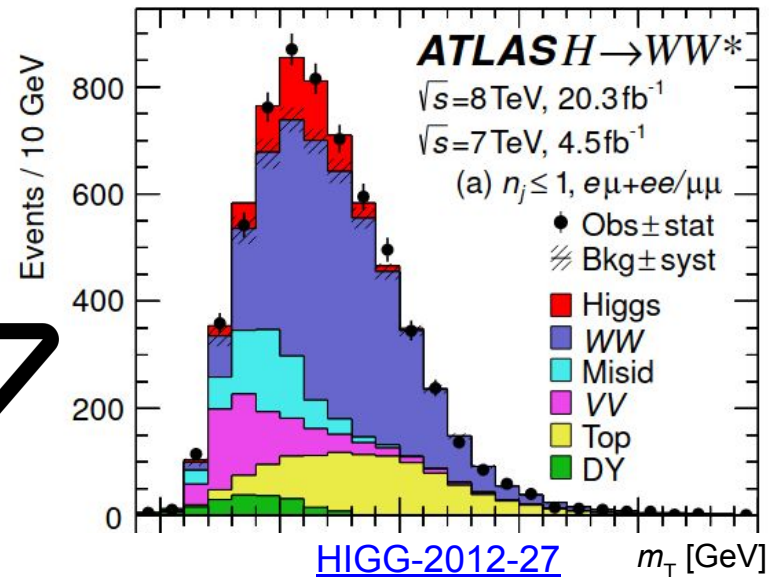
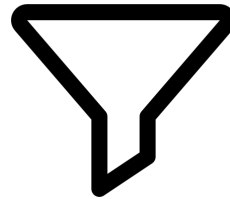


https://www.ataxia.org/wp-content/uploads/2022/07/needle-in-a-haystack-1752846_960_720.jpg



Ok fine, we triggered, we reconstructed our particles, what's next?

- The entire dataset still contains very many collision events
- BSM physics is rare → need to filter data to “see” our BSM physics
- Use e.g. number of reconstructed electrons, jet momenta, etc.
- Filtered data comprised of our signal & Standard Model background
- Define multiple “filters” = “regions”
 - **Signal regions:** much signal, little bkg.
 - **Control regions:** estimate bkg.
 - **Validation regions:** verify bkg. estimate

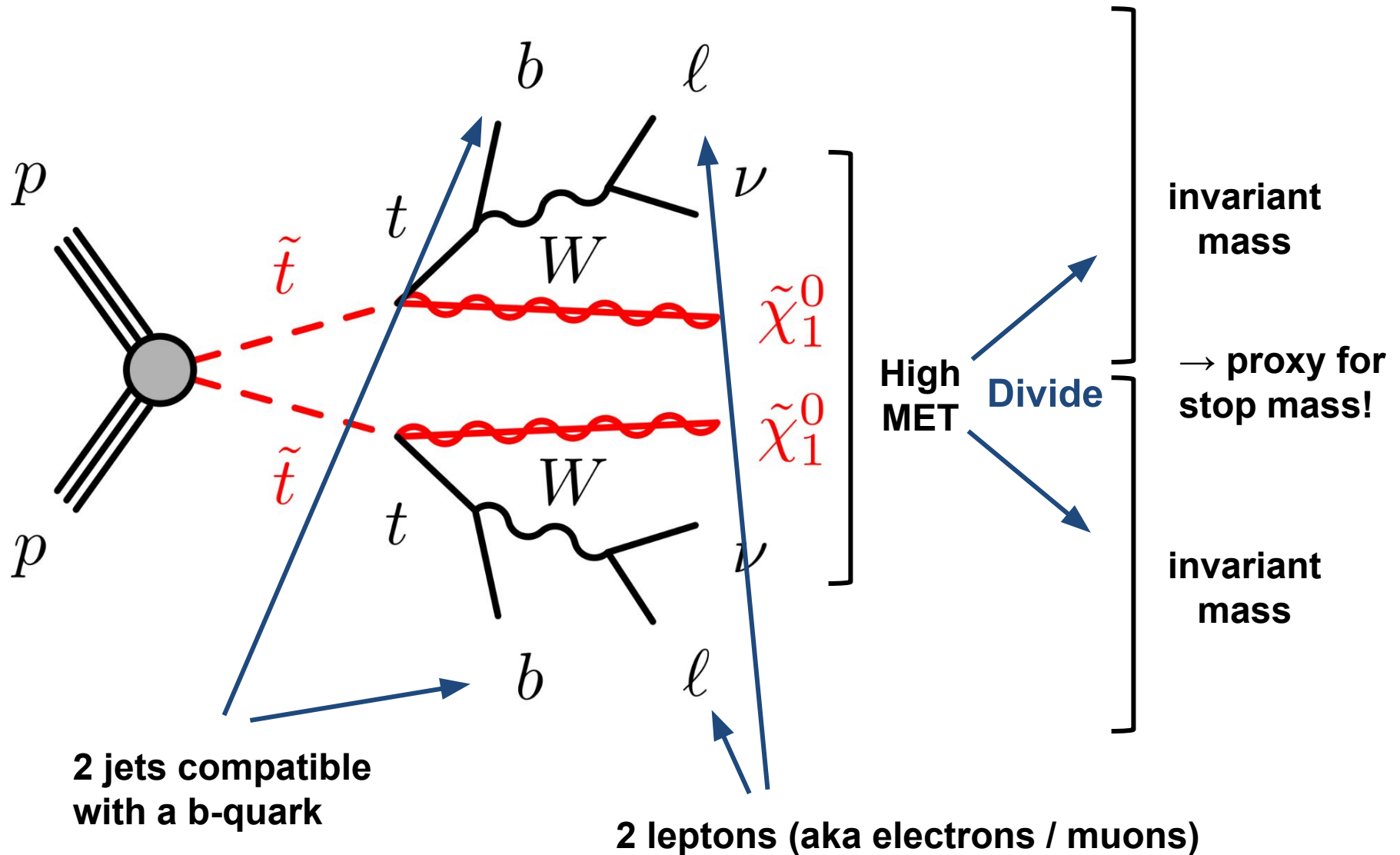


Searching for SUSY

- Recall: there are many SUSY models, e.g. MSSM, pMSSM, CMSSM, etc.
- In practise: study simplified versions of these models to reduce number of parameters to study (MSSM has 124 free parameters)
- Often also reduce number of production channels / decay channels
- Result give an **indication** for what would happen in the full model

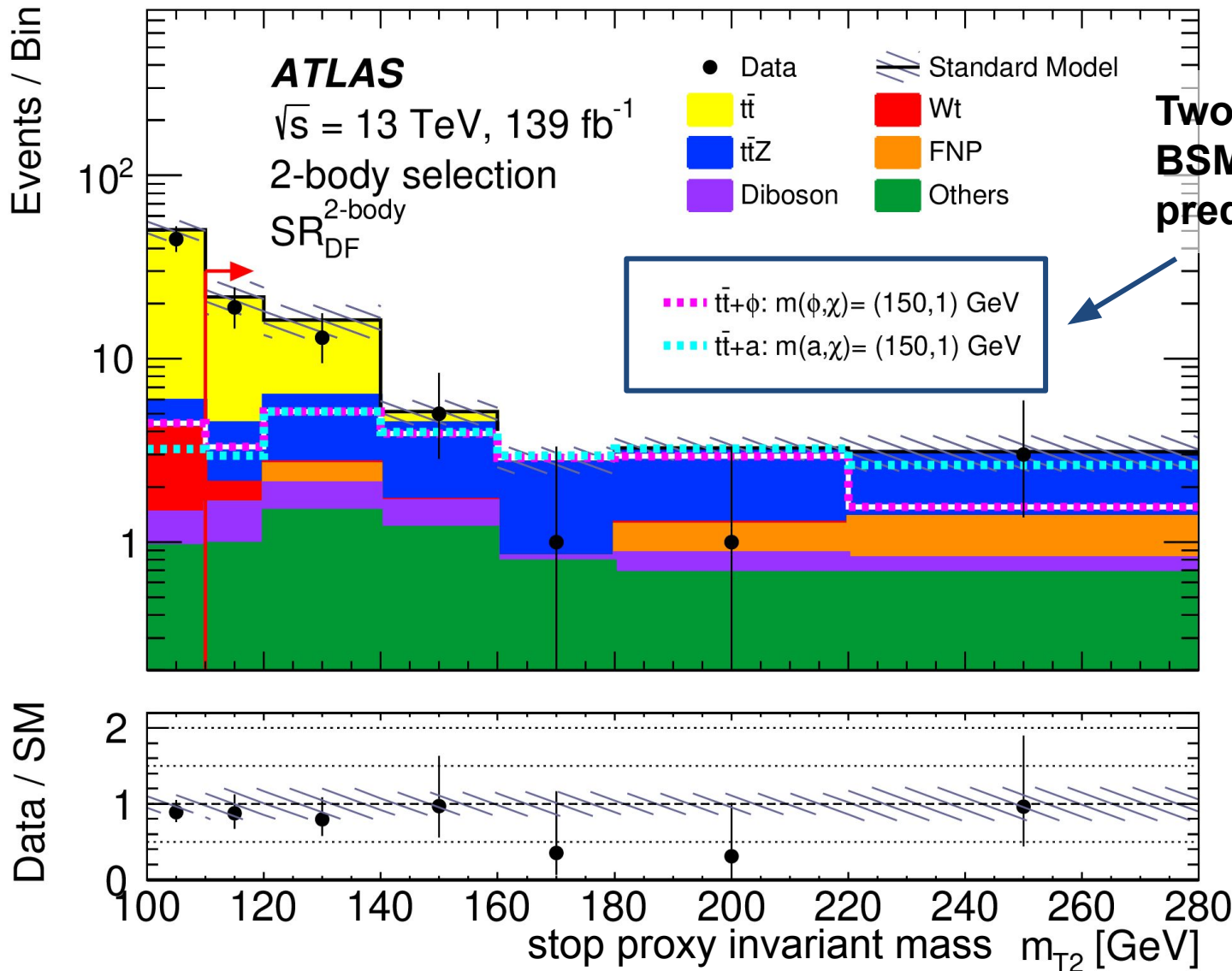
SUSY example: stop-2L

[SUSY-2018-08](#)



Stop-2L: final distribution

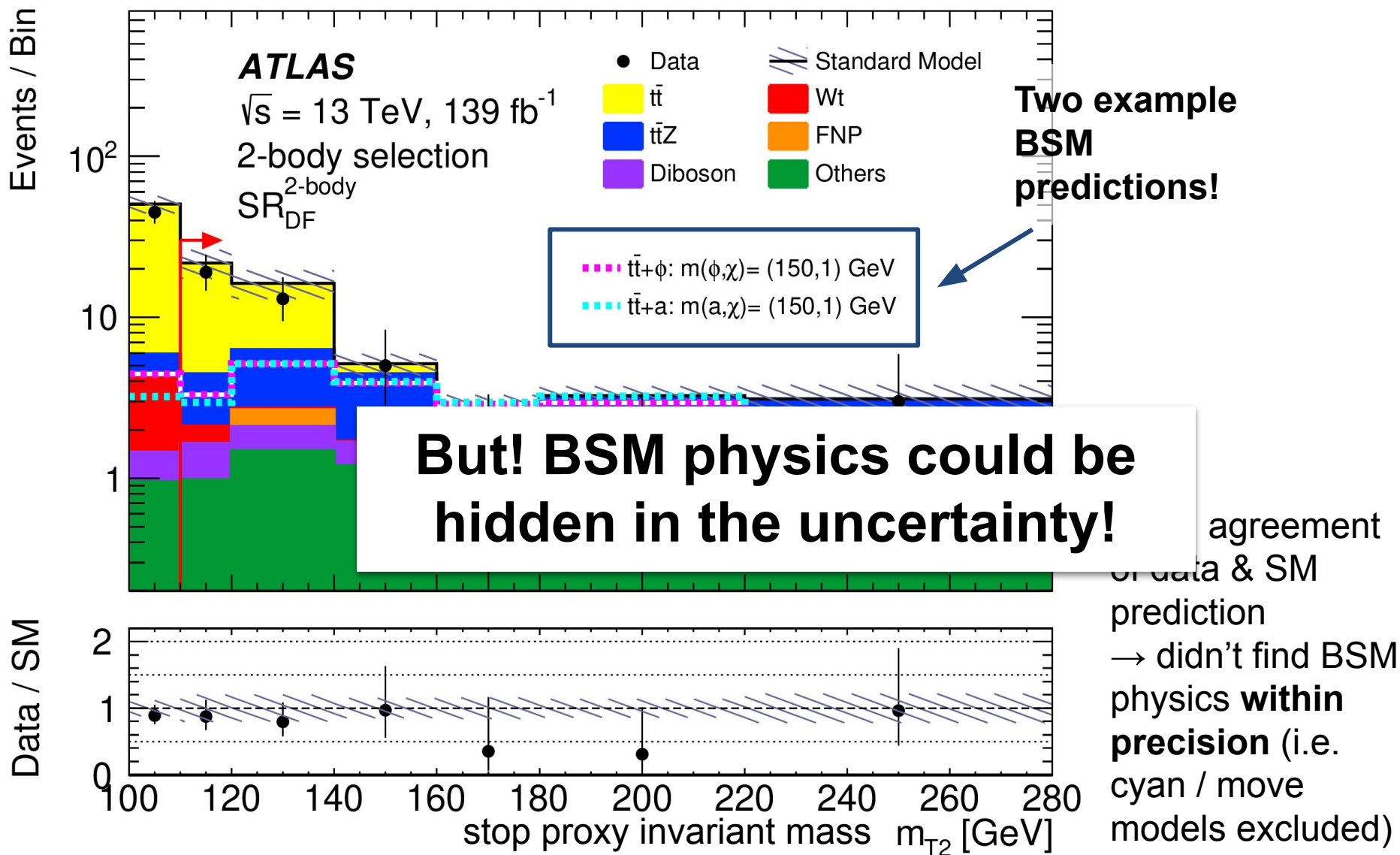
SUSY-2018-08



Good agreement of data & SM prediction
 → didn't find BSM physics **within precision** (i.e. cyan / pink models excluded)

Stop-2L: final distribution

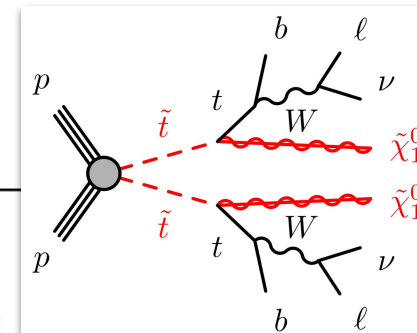
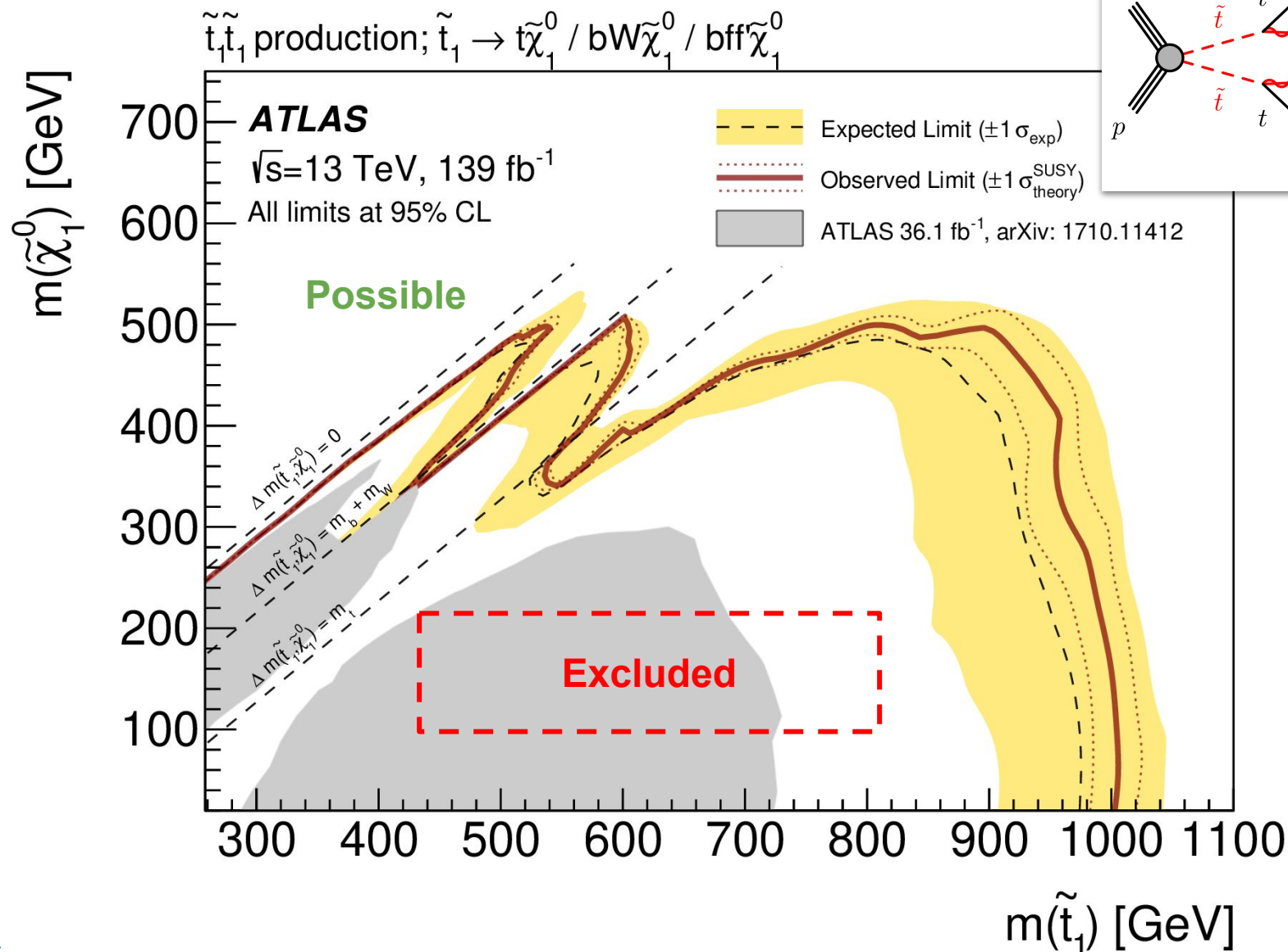
[SUSY-2018-08](#)



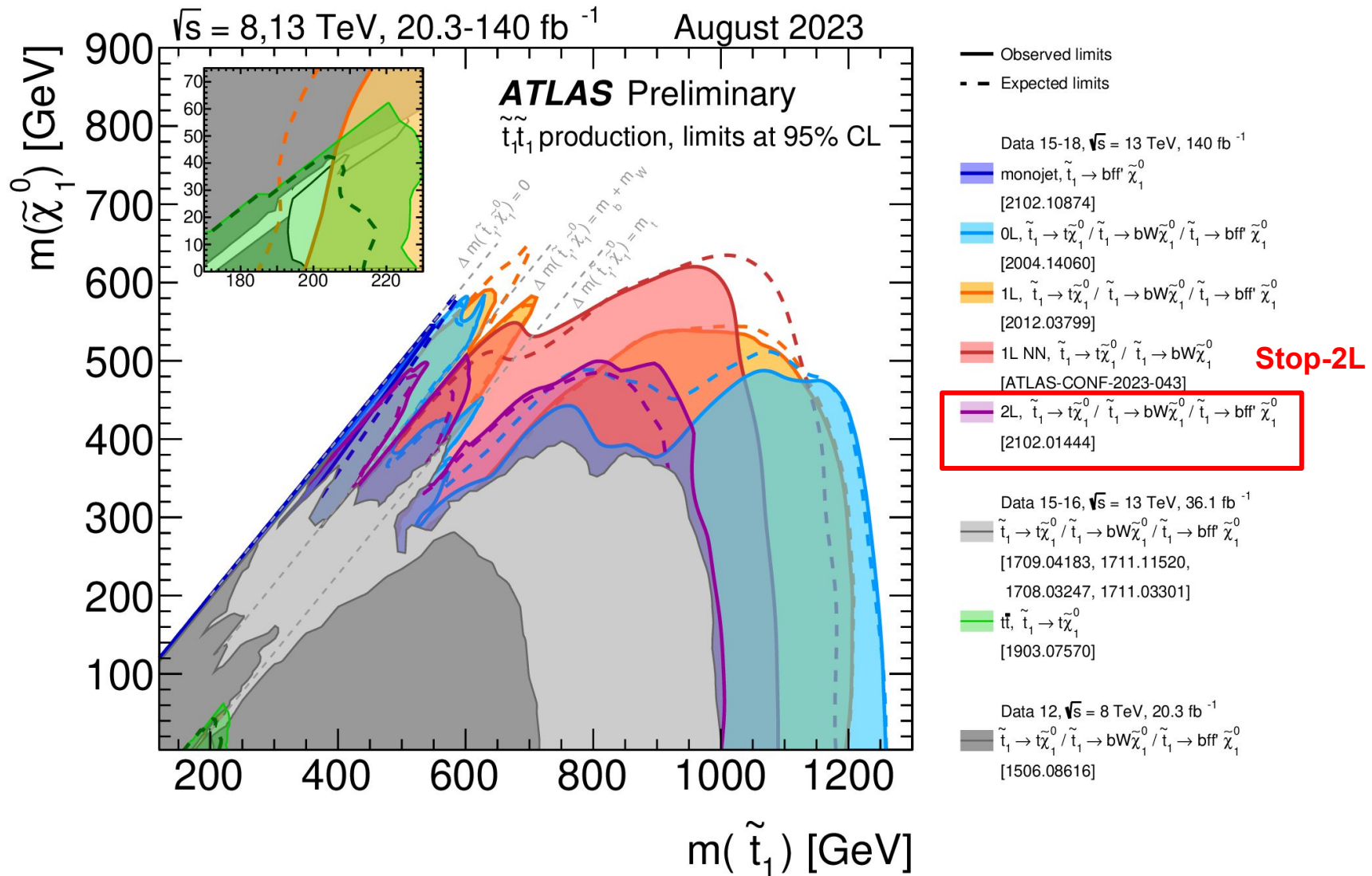
Stop-2L: results

Cannot exclude a model if it is hidden in the uncertainty* → exclusion limits

[SUSY-2018-08](#)



There are many SUSY searches out there



There are many SUSY searches out there

- SUSY particle masses between 100 GeV and 1 TeV largely excluded
- Some models under pressure; hierarchy problem only solved if masses “light”

July 2024

$\sqrt{s} = 13$ TeV

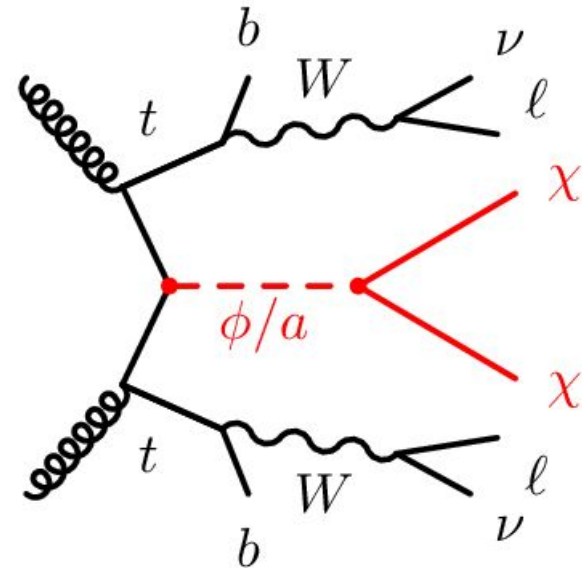
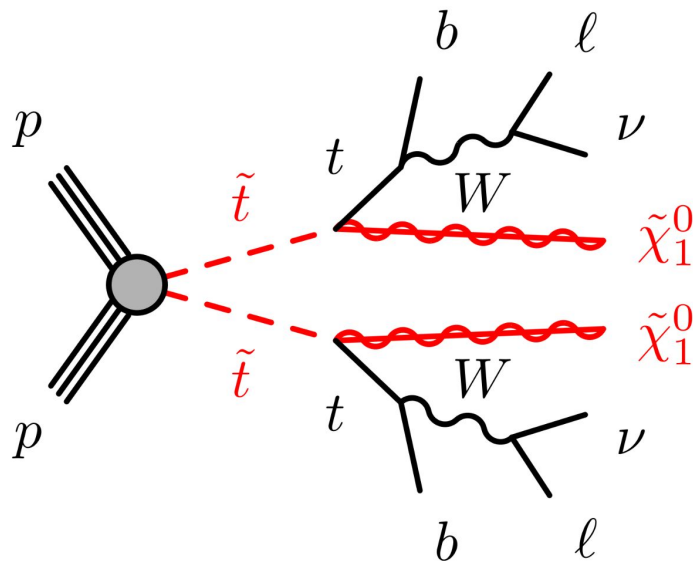
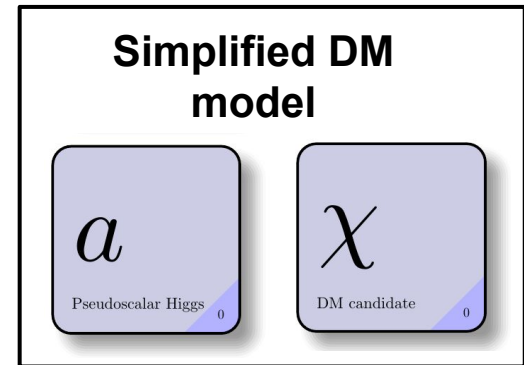
Model	Signature	$\int \mathcal{L} dt$ [fb ⁻¹]	Mass limit	Reference				
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	2-6 jets 1-3 jets	E_T^{miss} 140 E_T^{miss} 140	$\tilde{q}$ [1x, 8x Degen.] 1.0 $\tilde{q}$ [8x Degen.] 0.9	$m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5$ GeV	2010.14293 2102.10874	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	0 e, μ	2-6 jets	E_T^{miss} 140	$\tilde{g}$ $\tilde{g}$ Forbiden	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{\chi}_1^0) = 1000$ GeV	2010.14293 2010.14293	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}W\tilde{\chi}_1^0$	1 e, μ	2-6 jets	E_T^{miss} 140	$\tilde{g}$	$m(\tilde{\chi}_1^0) < 600$ GeV	2101.01629	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}(\ell\ell)\tilde{\chi}_1^0$	$ee, \mu\mu$	2 jets	E_T^{miss} 140	$\tilde{g}$	$m(\tilde{\chi}_1^0) < 700$ GeV	2204.13072	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\bar{q}WZ\tilde{\chi}_1^0$	0 e, μ SS e, μ	7-11 jets 6 jets	E_T^{miss} 140 E_T^{miss} 140	$\tilde{g}$ $\tilde{g}$	$m(\tilde{\chi}_1^0) < 600$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV	2008.06032 2307.01094	
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	0-1 e, μ SS e, μ	3 b 6 jets	E_T^{miss} 140 E_T^{miss} 140	$\tilde{g}$ $\tilde{g}$	$m(\tilde{\chi}_1^0) < 500$ GeV $m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV	2211.08028 1909.08457	
	3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1$	0 e, μ	2 b	E_T^{miss} 140	$\tilde{b}_1$ $\tilde{b}_1$	$m(\tilde{\chi}_1^0) < 400$ GeV 10 GeV $< \Delta m(\tilde{b}_1, \tilde{\chi}_1^0) < 20$ GeV	2101.12527 2101.12527
		$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow b\tilde{h}\tilde{\chi}_1^0$	0 e, μ 2 τ	6 b 2 b	E_T^{miss} 140 E_T^{miss} 140	$\tilde{b}_1$ Forbiden $\tilde{b}_1$	$\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 100$ GeV $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 0$ GeV	1908.03122 2103.08189
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$		0-1 e, μ	≥ 1 jet	E_T^{miss} 140	$\tilde{t}_1$	$m(\tilde{\chi}_1^0) = 1$ GeV	2004.14060, 2012.03799	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$		1 e, μ	3 jets/1 b	E_T^{miss} 140	$\tilde{t}_1$ Forbiden	$m(\tilde{\chi}_1^0) = 500$ GeV	2012.03799, 2401.13430	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$		1-2 τ	2 jets/1 b	E_T^{miss} 140	$\tilde{t}_1$ Forbiden	$m(\tilde{\tau}_1) = 800$ GeV	2108.07665	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{\chi}_1^0, \tilde{c} \rightarrow c\tilde{\chi}_1^0$		0 e, μ mono-jet	2 c mono-jet	E_T^{miss} 36.1 E_T^{miss} 140	$\tilde{t}_1$ $\tilde{t}_1$	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV	1805.01649 2102.10874	
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$		1-2 e, μ	1-4 b	E_T^{miss} 140	$\tilde{t}_1$	$m(\tilde{\chi}_2^0) = 500$ GeV	2006.05880	
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$		3 e, μ	1 b	E_T^{miss} 140	$\tilde{t}_2$ Forbiden	$m(\tilde{\chi}_1^0) = 360$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 40$ GeV	2006.05880	
EW direct	$\tilde{\chi}_1^+ \tilde{\chi}_2^0$ via WZ	Multiple ℓ/jets $ee, \mu\mu$	≥ 1 jet	E_T^{miss} 140 E_T^{miss} 140	$\tilde{\chi}_1^+ / \tilde{\chi}_2^0$ $\tilde{\chi}_1^+ / \tilde{\chi}_2^0$	$m(\tilde{\chi}_1^0) = 0$, wino-bino $m(\tilde{\chi}_1^+) - m(\tilde{\chi}_1^0) = 5$ GeV, wino-bino	2106.01676, 2108.07586 1911.12606	
	$\tilde{\chi}_1^+ \tilde{\chi}_1^+$ via WW	2 e, μ		E_T^{miss} 140	$\tilde{\chi}_1^+$	$m(\tilde{\chi}_1^0) = 0$, wino-bino	1908.08215	
	$\tilde{\chi}_1^+ \tilde{\chi}_2^0$ via Wh	Multiple ℓ/jets		E_T^{miss} 140	$\tilde{\chi}_1^+ / \tilde{\chi}_2^0$ Forbiden	$m(\tilde{\chi}_1^0) = 70$ GeV, wino-bino	2004.10894, 2108.07586	
	$\tilde{\chi}_1^+ \tilde{\chi}_1^+$ via $\tilde{\ell}_L/\bar{\nu}$	2 e, μ		E_T^{miss} 140	$\tilde{\chi}_1^+$	$m(\tilde{\ell}, \bar{\nu}) = 0.5(m(\tilde{\ell}_1^+) + m(\tilde{\chi}_1^0))$	1908.08215	
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$	2 τ		E_T^{miss} 140	$\tilde{\tau}$ [$\tilde{\tau}_R \tilde{\tau}_R$]	$m(\tilde{\chi}_1^0) = 0$	2402.06003	
	$\tilde{L}_{L,R}\tilde{L}_{L,R}, \tilde{L} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ $ee, \mu\mu$	0 jets ≥ 1 jet	E_T^{miss} 140 E_T^{miss} 140	$\tilde{L}$ $\tilde{L}$	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{L}) - m(\tilde{\chi}_1^0) = 10$ GeV	1908.08215 1911.12606	
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ	≥ 3 b 0 jets	E_T^{miss} 140 E_T^{miss} 140	$\tilde{H}$ $\tilde{H}$	$\text{BR}(\tilde{H}^0 \rightarrow h\tilde{G}) = 1$ $\text{BR}(\tilde{H}_\pm^1 \rightarrow Z\tilde{G}) = 1$	2401.14922 2103.11684	
		0 e, μ ≥ 2 large jets		E_T^{miss} 140	$\tilde{H}$	$\text{BR}(\tilde{H}_\pm^1 \rightarrow Z\tilde{G}) = 1$	2108.07586	
		2 e, μ ≥ 2 jets		E_T^{miss} 140	$\tilde{H}$	$\text{BR}(\tilde{H}_\pm^1 \rightarrow Z\tilde{G}) = \text{BR}(\tilde{H}_\pm^1 \rightarrow h\tilde{G}) = 0.5$	2204.13072	
	Long-lived particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	E_T^{miss} 140	$\tilde{\chi}_1^\pm$ $\tilde{\chi}_1^\pm$	Pure Wino Pure higgsino	2201.02472 2201.02472
Stable $\tilde{g}$ R-hadron		pixel dE/dx		E_T^{miss} 140	$\tilde{g}$		2205.06013	
Metastable $\tilde{g}$ R-hadron, $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$		pixel dE/dx		E_T^{miss} 140	$\tilde{g}$ [$\tau(\tilde{g}) = 10$ ns]	$m(\tilde{\chi}_1^0) = 100$ GeV	2205.06013	
$\tilde{\ell}\tilde{\ell}, \tilde{\ell} \rightarrow \ell\tilde{G}$		Displ. lep		E_T^{miss} 140	$\tilde{\ell}, \tilde{\mu}$	$\tau(\tilde{\ell}) = 0.1$ ns $\tau(\tilde{\ell}) = 0.1$ ns	ATLAS-CONF-2024-011 ATLAS-CONF-2024-011	
		pixel dE/dx		E_T^{miss} 140	$\tilde{\ell}$	$\tau(\tilde{\ell}) = 10$ ns	2205.06013	

100 GeV

1 TeV

Searching for Dark Matter with Stop-2L

- In a simplified Dark Matter model, can produce Dark Matter from a mediator radiating off an inner top-quark line
- This mediator decays into 2x DM particle
- Final state the same as in Stop-2L!!!
- Simultaneously search for SUSY and DM!!!



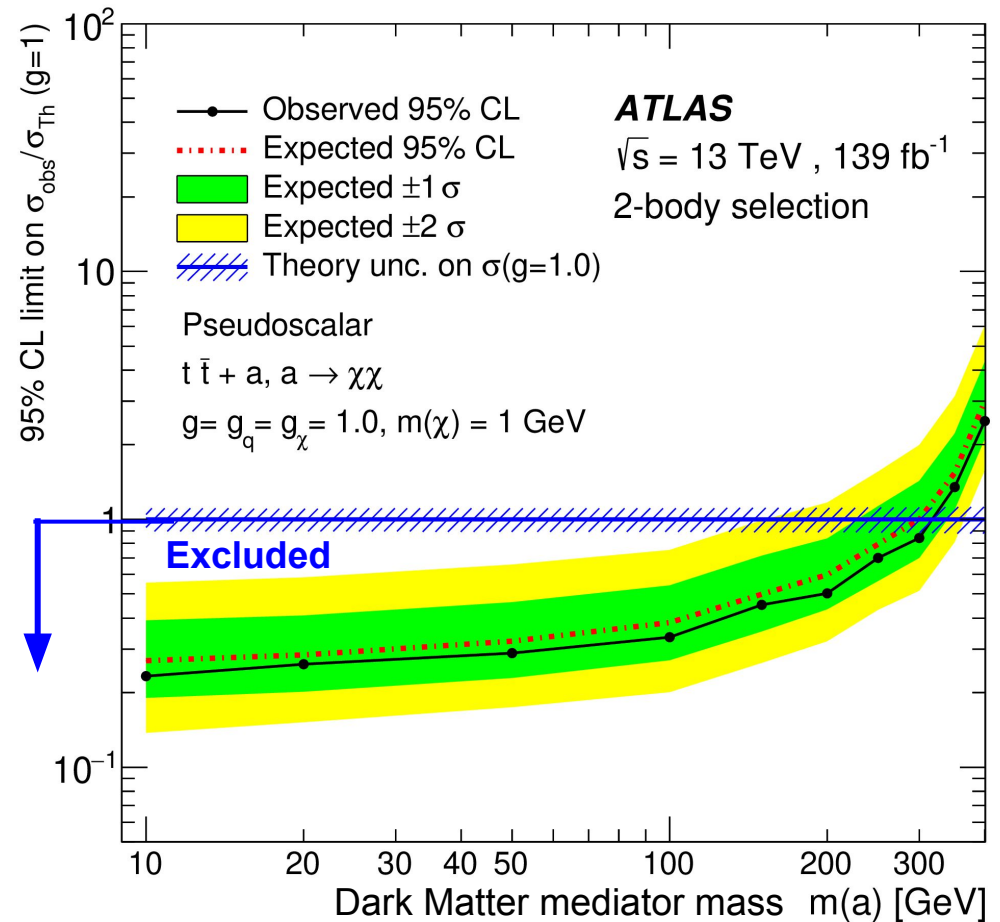
[SUSY-2018-08](#)

Stop-2L Dark Matter results

- Exclude DM mediator masses up to 300 GeV

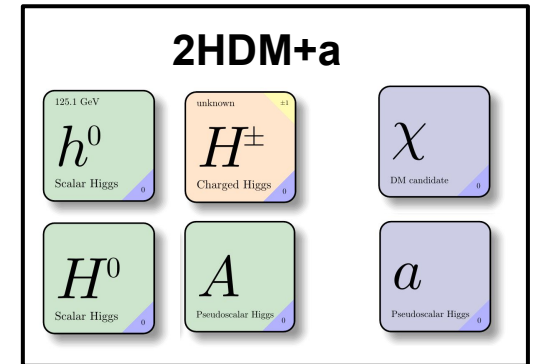
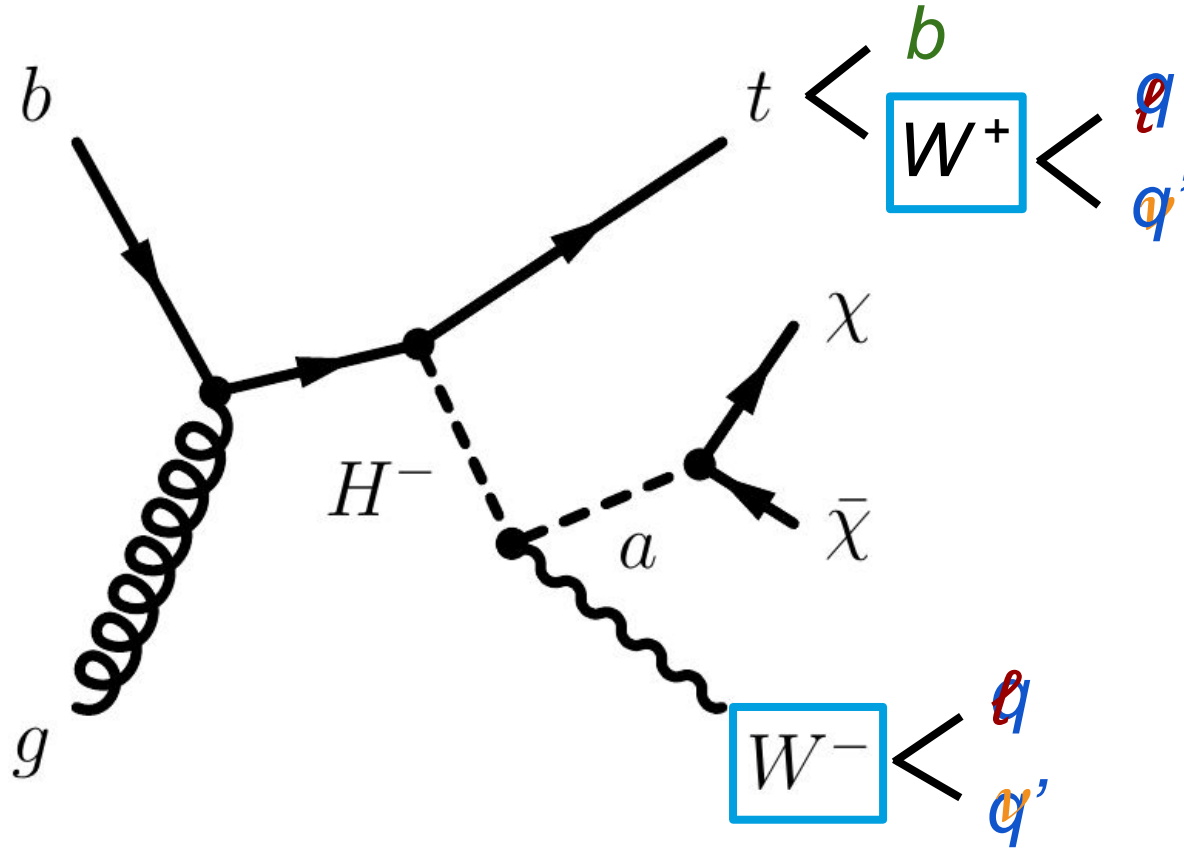
= “maximum factor with which I can multiply the BSM prediction so it is hidden in the uncertainty”*

If below 1: not hidden in uncertainty!
So excluded!



SUSY-2018-08

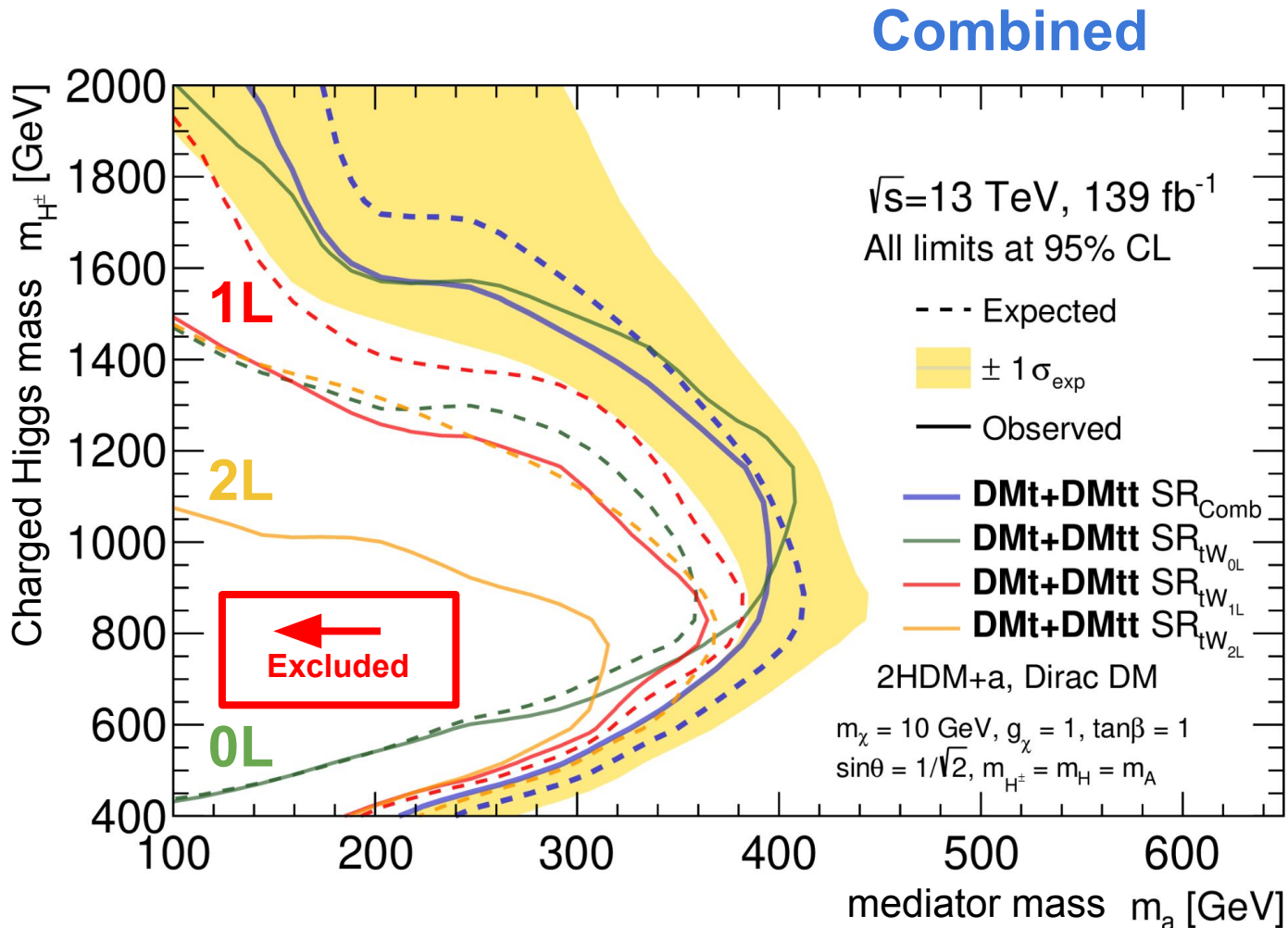
Searching for Dark Matter in tW+DM



- tW+DM high x-sec
- Two **W-bosons**:
zero, one or two visible leptons (e,μ)
- **0L, 1L, 2L channel**
→ basically 3 separate searches
- To get maximum sensitivity
→ combine the three searches

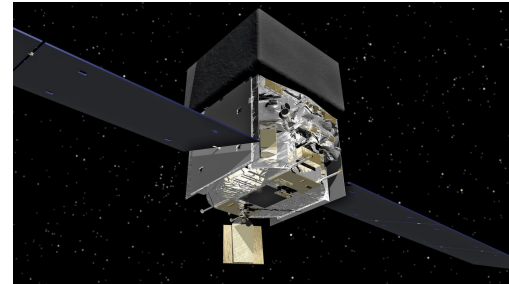
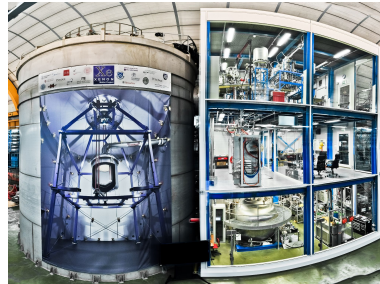
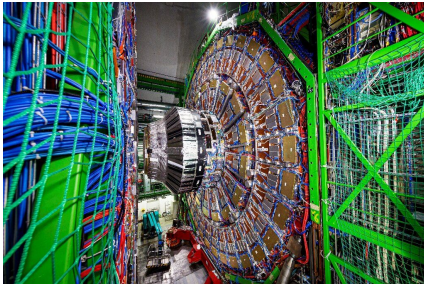
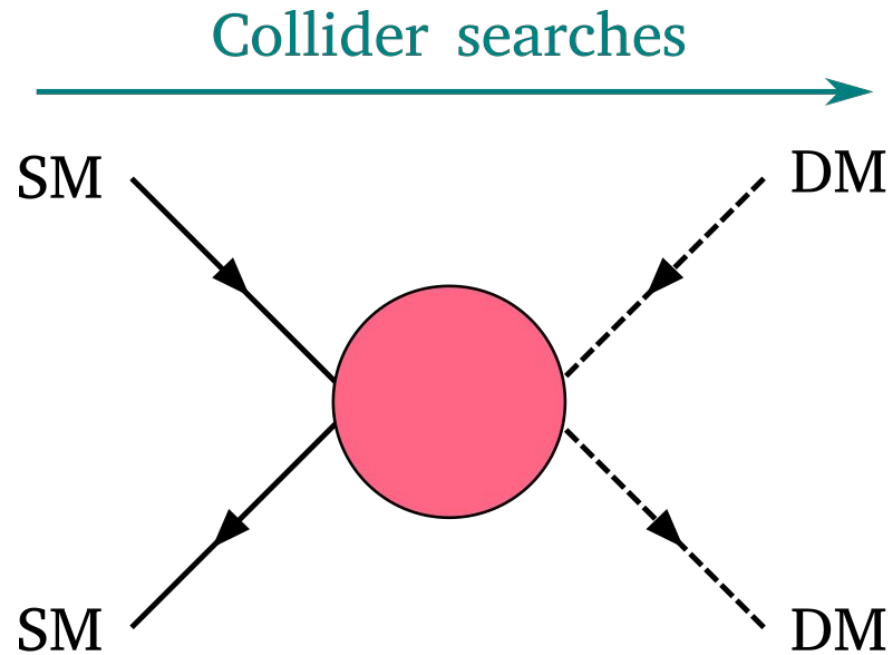
tW+DM: results

- Exclude DM mediator masses up to 300 GeV



BREAK
(5-10 mins)

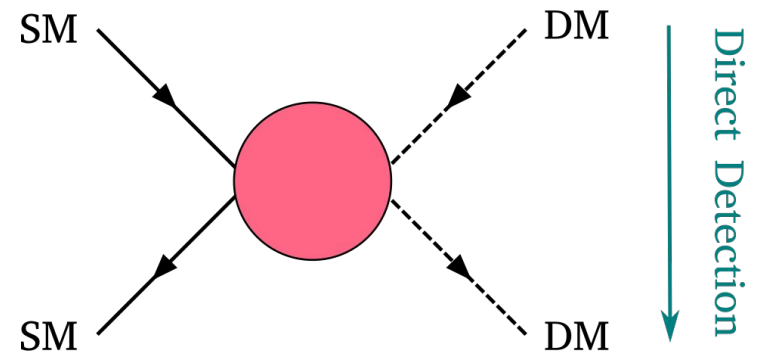
Other ways to search for Dark Matter



Direct detection Dark Matter searches

[arxiv:1903.03026](https://arxiv.org/abs/1903.03026)

- Assume: permanent flow of DM particles through planet earth
- **DM scatters with nuclei** (as no EM interaction) → **measurable**
- Number of scattered DM particles depends on:
 - Density of DM, $\rho_0 = 0.3 \text{ GeV}/c^2/\text{cm}^3$
 - DM velocity w.r.t. earth $\sim 220 \text{ km/s}$
 - DM-nuclei cross-section
 - Mass of DM & interacting nucleus
 - **Mass of detector → higher better**
 - **Spin of DM particle and spin sensitivity of target material**
 - **Minimum recoil energy sensitivity of detector → lower better**

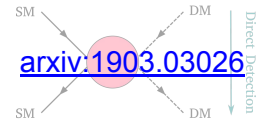


$$\frac{dR}{dE_{\text{nr}}} = \frac{\rho_0 M}{m_N m_\chi} \int_{v_{\text{min}}}^{v_{\text{esc}}} v f(v) \frac{d\sigma}{dE_{\text{nr}}} dv$$

Diagram illustrating the components of the direct detection rate equation:

- total detector mass** points to M
- DM density** points to ρ_0
- Events per recoil energy** points to $\frac{dR}{dE_{\text{nr}}}$
- mass of nucleus/DM** points to m_N and m_χ
- velocity distribution** points to $f(v)$
- differential interaction cross-section, $\propto m_N$** points to $\frac{d\sigma}{dE_{\text{nr}}}$

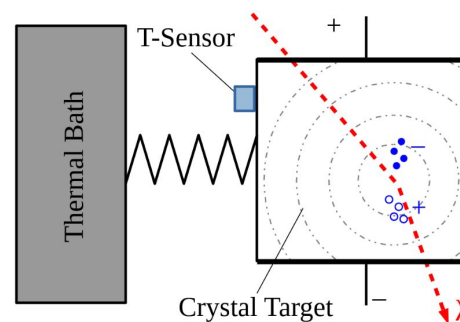
Experimental setups of DD experiments



- Measure recoil of nuclei by DM
→ phonons/heat
- Backgrounds:
 - cosmic muons, neutrons (from nuclear decays, cosmic μ)
 - γ , $\beta^\pm$ from radioactive decays
→ often interact via EM
- Use very pure, non-activated, well-selected materials & shielding
- Measure two signals: one from EM recoil (e.g. charge; higher for bkg.); one from nuclear recoil (e.g. heat)

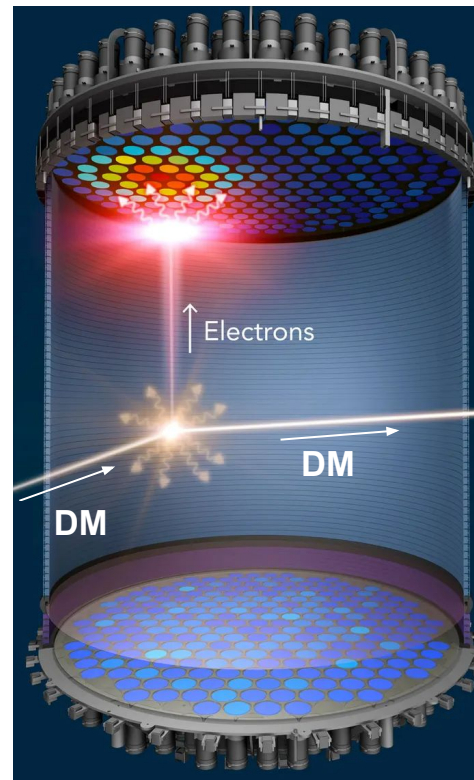
Image credit bottom:

https://th-thumbnailer.cdn-si-edu.com/3Qxw0k7nJ0cC7BUwbFSP5zHO6w8=/fit-in/1600x0/https%3A%2F%2Ftf-cmsv2-smithsonianmag-media.s3.amazonaws.com%2Ffiler%2F42%2Fd5%2F42d5a303-b006-4972-9efa-92797a25ba9c%2F31667821088_f762d2c200_o.jpg



Cryogenic detectors:

- At very low T (mK)
- Scattering DM increases T
- Bkg. ionises material
→ separate signal
- Materials: Ge, Si



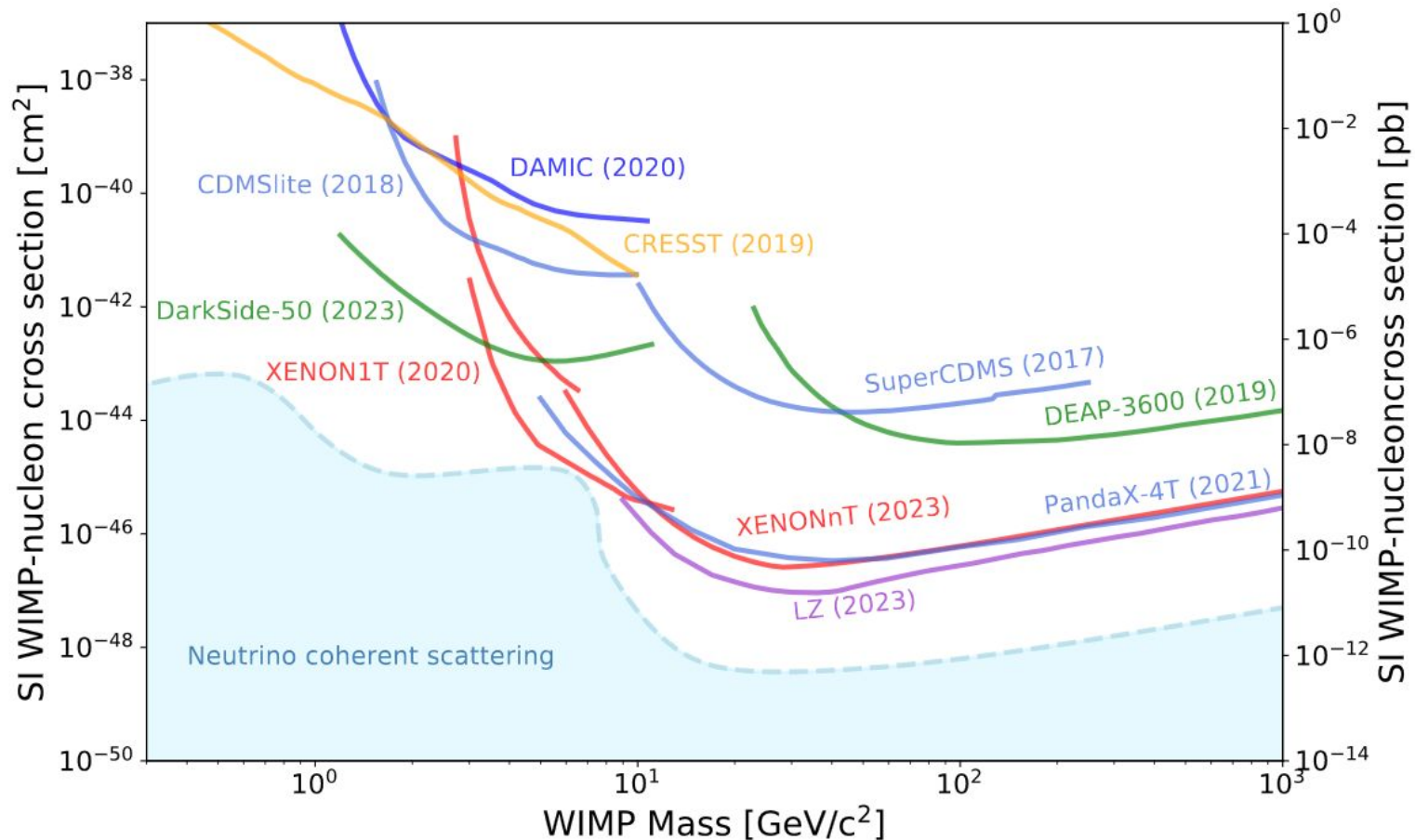
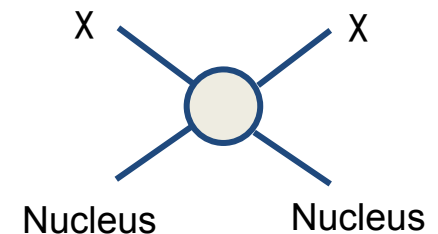
Noble liquid detectors:

- E.g. liquid Ar, Xe
- Phonons from DM → photons + ionisation
- Example XENON experiment
- Detect photons with PMTs at the side
- Detect ionisation by E-field + scintillation signal at the top
- Ratio of photons/ionisation different for DM / bkg.

DD experiments – results

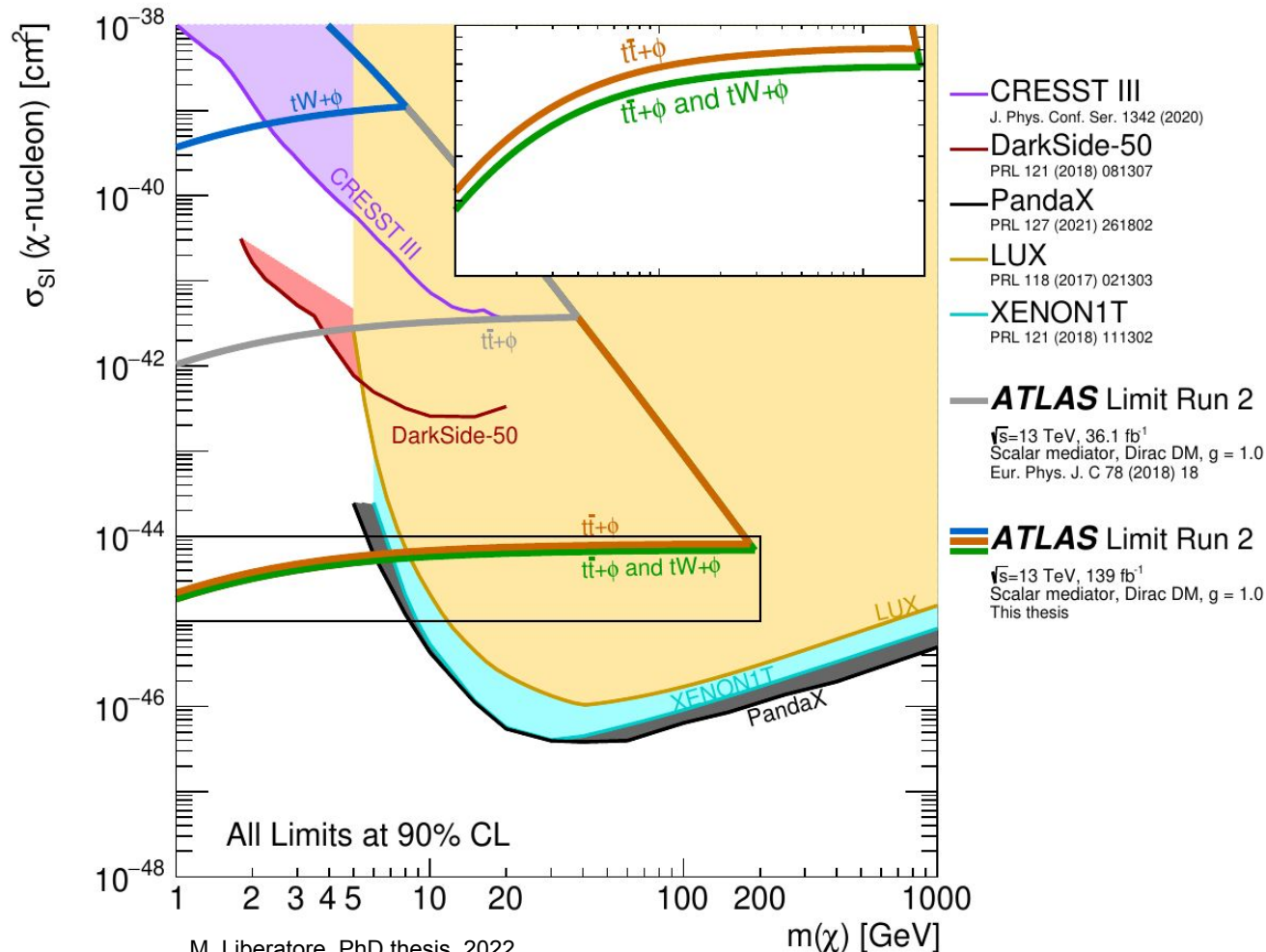
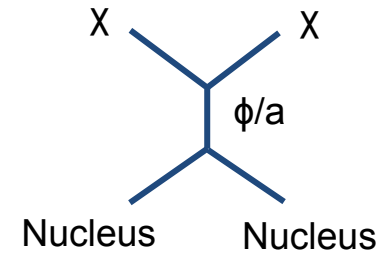
[rpp2024-rev-dark-matter](#)

- Direct detection experiment results interpreted in terms of Effective Field Theories, no assumption on interaction mechanism of DM and nucleus
- The neutrino floor is not that far!



DD experiments – comparison to collider

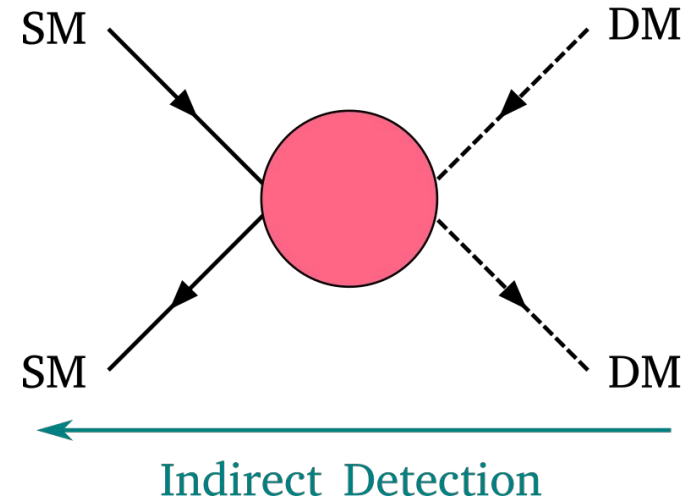
- Comparing to collider: need to use a “simplified DM model”
→ additional mediator ϕ/a
- Collider searches dominate at low masses!



Indirect detection Dark Matter searches

[rpp2022-rev-dark-matter](#)

- In areas of high mass density: Dark Matter will annihilate → indirect detection
- Annihilation yields signals such as photon pairs, neutrinos, baryons
- High mass density areas: sun, galaxy centers, ...
- Production rate of IDD events depends on:
 - Annihilation rate / cross-section
 - Dark Matter density
 - The number of final state particles

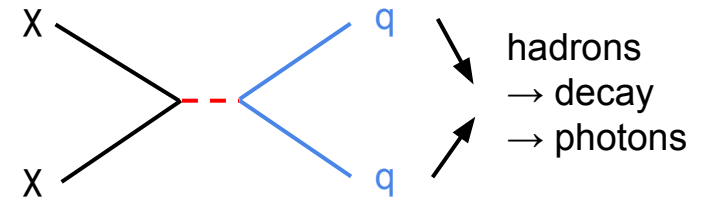


https://www.youcanseethemilkyway.com/wp-content/uploads/2023/01/Massive_Black_Hole_at_the_Center_of_the_Milky_Way.jpg

Experimental techniques of indirect detection

• Photons

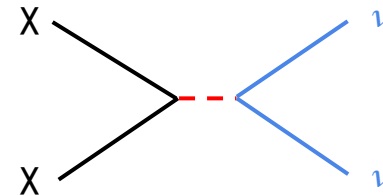
- produced e.g. when $DM + DM \rightarrow$ quarks
- search for high energy photons (“gamma rays”) e.g. from the galactic center
- experiments: e.g. FERMI-LAT



https://www.svmmetrymagazine.org/sites/default/files/legacy/images/200912/article05_image01.jpg

• Neutrinos

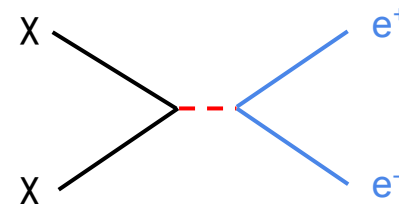
- produced e.g. in the sun
- search for high energy neutrinos from sun
- experiments: e.g. IceCube, Kamiokande



<https://cdn.mos.cms.futurecdn.net/CJ4pXsqqF4VwYdRXPyT3tK.jpg>

• Antiparticles

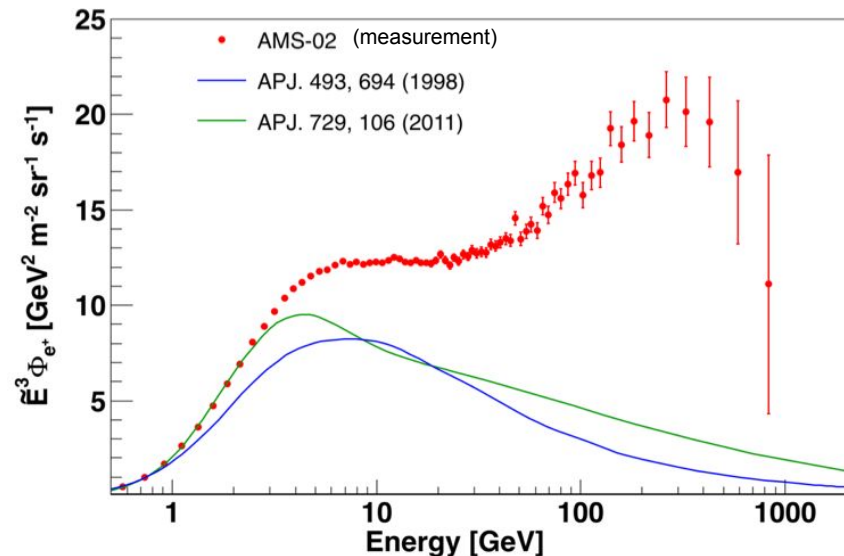
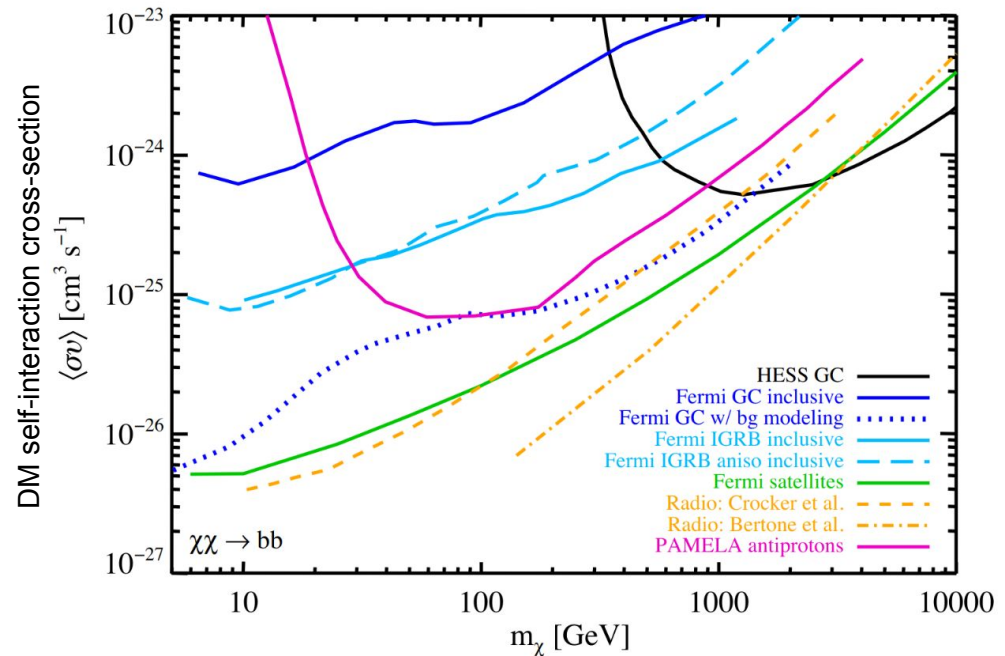
- produced e.g. if $DM + DM \rightarrow e^+e^-$
- charged particles deflected in universe → generally search for antiparticles
- experiments: e.g. HEAT, AMS, HESS



Indirect detection: results

[arxiv:1604.00014](https://arxiv.org/abs/1604.00014)

- Covered mass range comparable to DD
- Exciting result: excess of positrons observed by multiple experiments (AMS, FERMI, PAMELA)
- Yet unclear if it is due to Dark Matter, astrophysical origin investigated

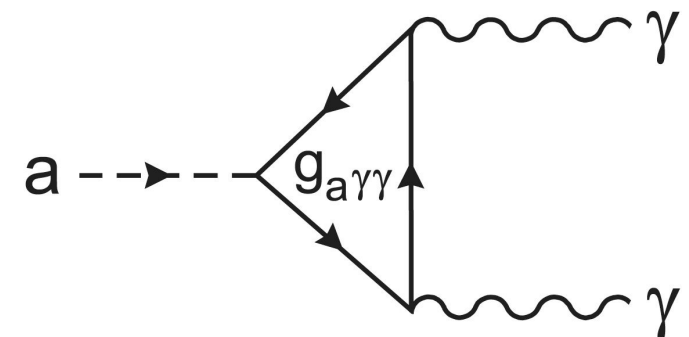
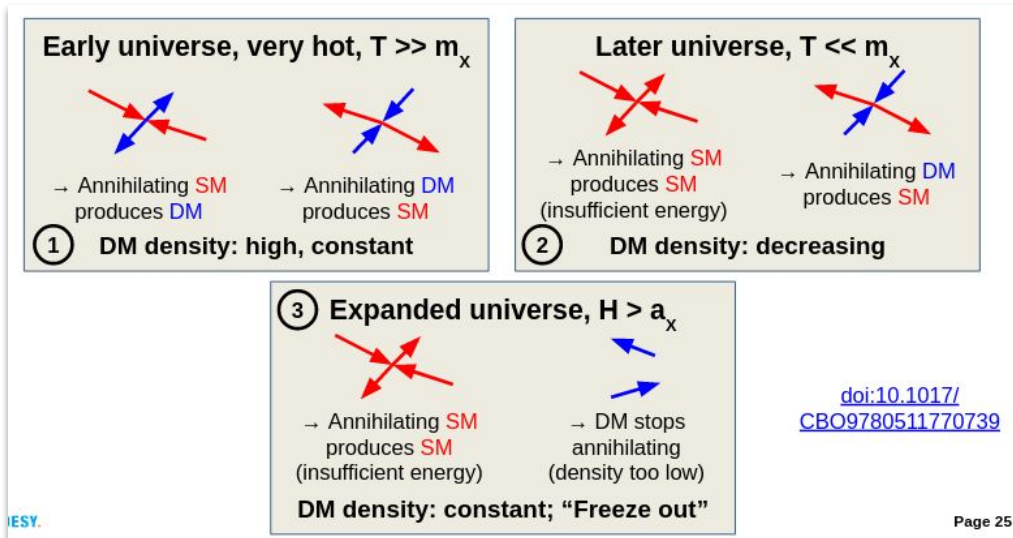
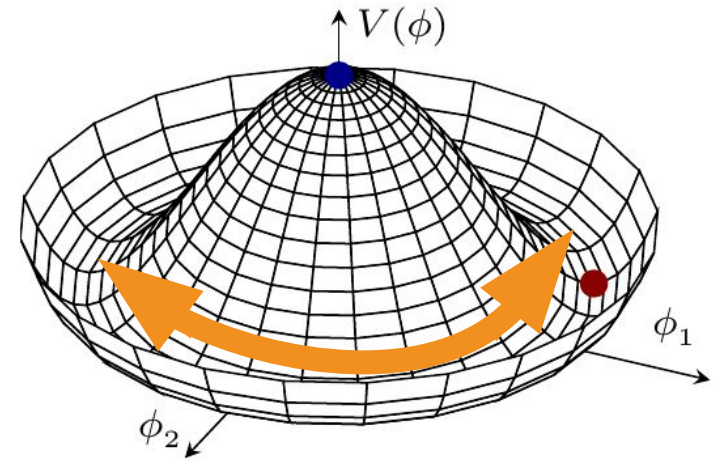


https://ams02.space/sites/default/files/inline-images/%20PhysRevLett.122.041102.Fig5_.png

Axions

- Searches so far considered WIMP DM
→ mass & x-sec. similar to weak bosons
- Also covered Axion Dark Matter last time
→ much lighter DM candidates
→ origin from solution of strong CP probl.
- Interact with SM matter by 2-photon vertex

(Lecture 1)

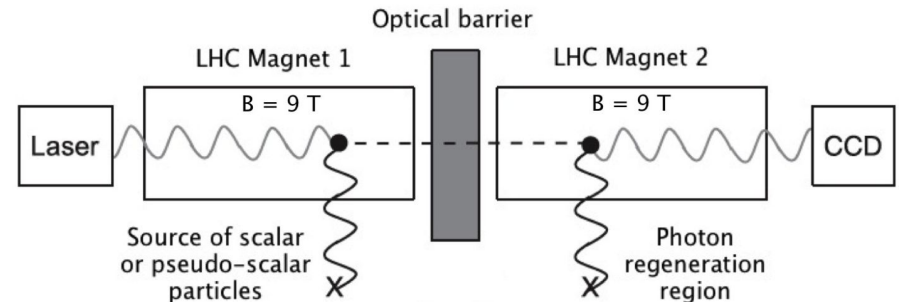


Axions search experiments

- Idea of Axion experiments:
 - can interact with Axions by exposing them to strong magnetic fields (=high photon flux)
 - can produce Axions by using strong lasers and strong magnetic fields
- ALPs experiment: shine laser into a strong magnetic field
 - photons converted to Axions
 - photons stopped by wall, Axions pass through
 - on the other side: another strong magnet → Axions converted to photons



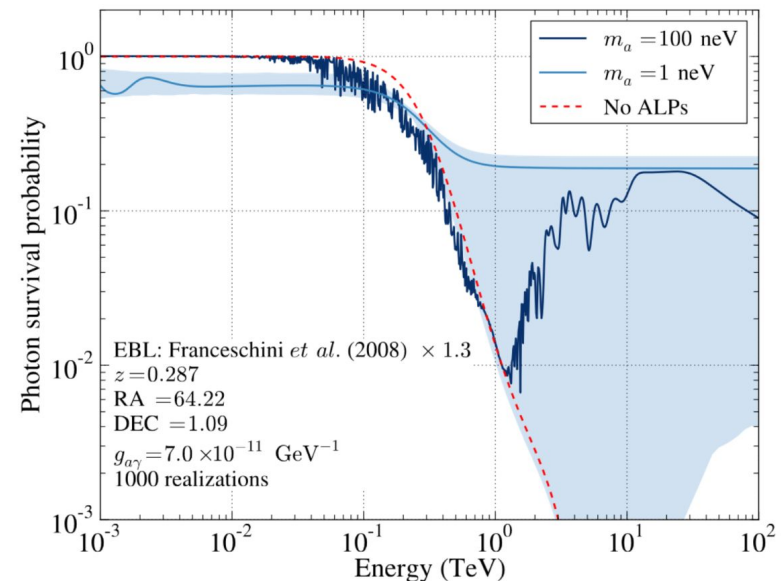
https://particle-physics.desy.de/sites/site_particle-physics/content/e221990/e222445/e228223/e228225/e228229/ALPS_ger.jpg



[arxiv:1410.2566](https://arxiv.org/abs/1410.2566)

Searching for Axions from & in the universe

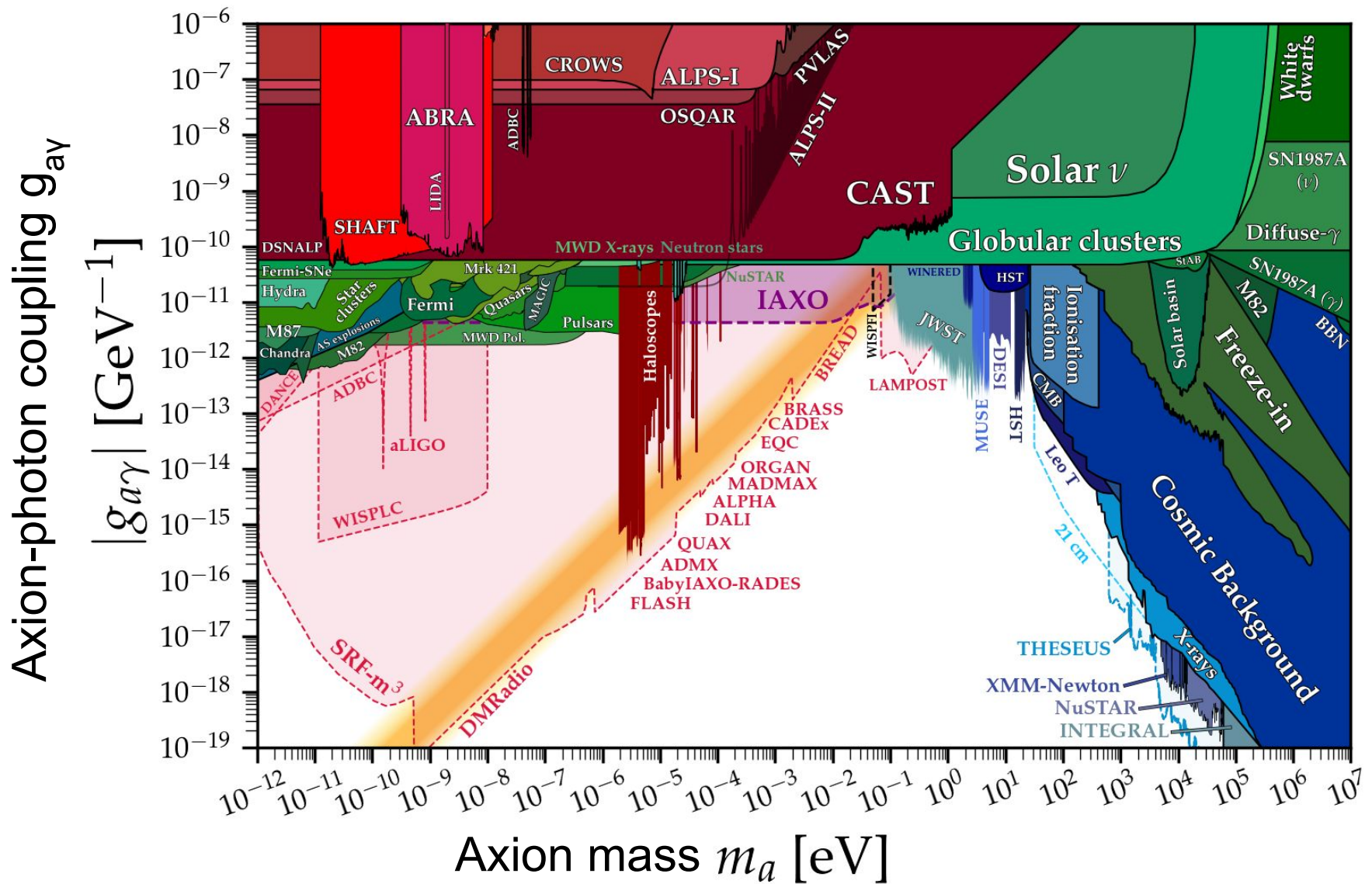
- CAST / (Baby-)IAXO experiments: search for Axions produced in sun
 - use a large magnet to point at the sun: Axions converted to photons
 - measure these photons
- The magnetic fields converting (photons $\rightarrow$ Axions) & (Axions $\rightarrow$ photons) can also be of astrophysical origin
 - expect to measure more high energy photons if there are Axions



M. Meyer, The Opacity of the Universe for High and Very High Energy γ -Rays,
<https://inspirehep.net/literature/1254304>

Axions: exclusion limits

Ciaran O'Hare, AxionLimits,
doi:10.5281/zenodo.3932430, [Link](#)



The magnetic moment of the muon

[rpp2022-rev-g-2-muon-anom-mag-moment](#)

- Recall: circular current = magnetic field → magnetic dipole moment
- Similarly: charged particles on circular orbits & spinning around themselves exhibit magnetic dipole moment
- Magnetic dipole moment basically indicates how strongly a particle is affected by a magnetic field

$$\mathbf{m}_S = - \frac{g_S \mu_B \mathbf{S}}{\hbar}, \quad \mu_B = \frac{e\hbar}{2m_e}$$

Annotations:
 - Red arrow pointing to g_S : ??
 - Blue arrow pointing to $\mathbf{S}$: spin-vector
 - Blue arrow pointing to $\hbar$: constant

- For the muon: $g_\mu = 2$ (SM LO),
loop effects → $g_\mu > 2$

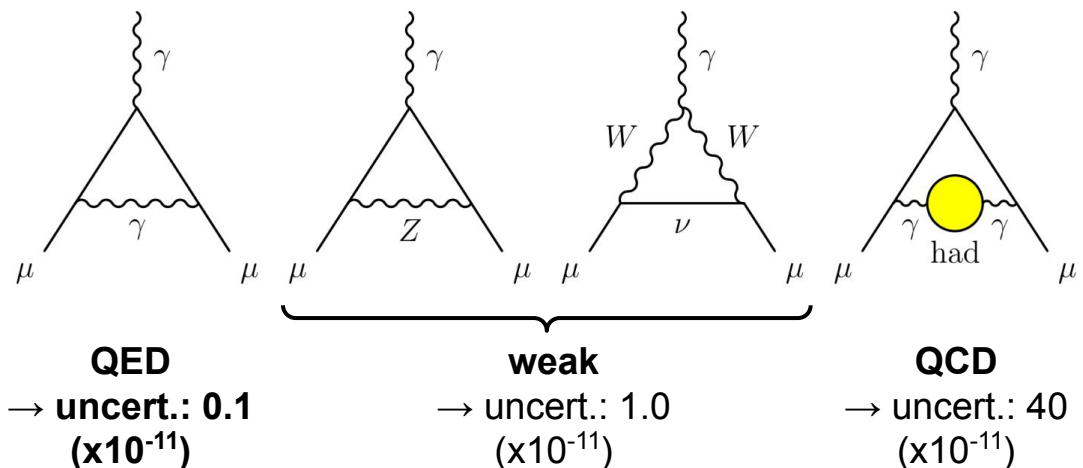
- Measure the difference:

$$a_\mu = (g_\mu - 2) / 2$$

→ **stress-test**

Standard Model

(the loops it predicts)

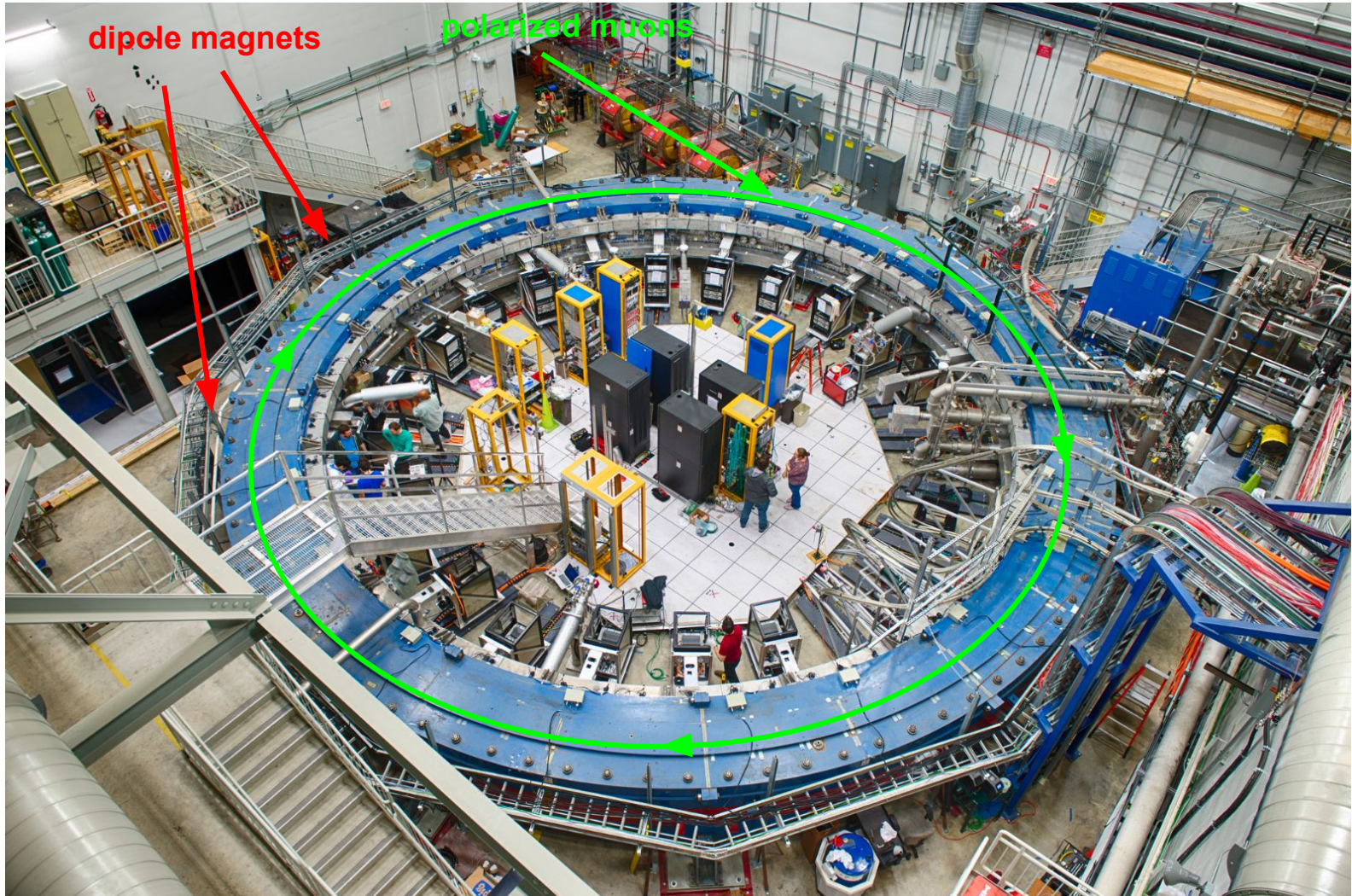


<https://www.particlebites.com/?p=8972>

How to measure g-2

- Produce muons, store ring via dipole magnets

$$\mathbf{F} = q \mathbf{v} \times \mathbf{B}$$



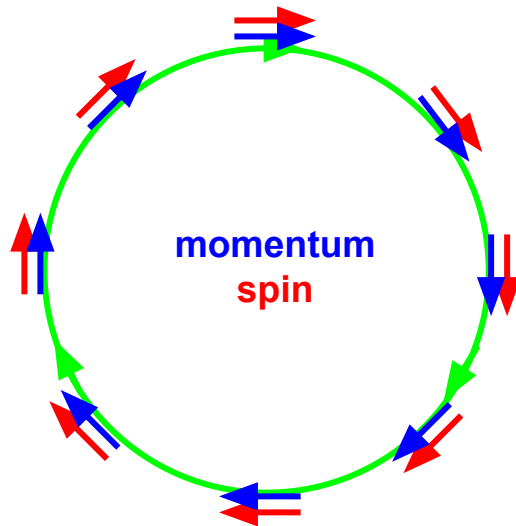
https://upload.wikimedia.org/wikipedia/commons/7/7c/Fermilab_g-2_%28E989%29_ring.jpg

How to measure g-2 (2)

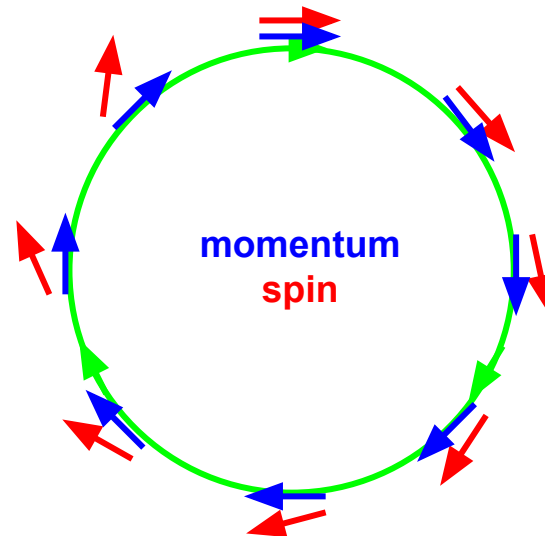
- Two frequencies:
 - rotation frequency of muons (cyclotron frequency) ω_c
 - precision frequency due to spin ω_s
- Their difference depends on g-2!!!
→ measure ω_a & $\mathbf{B} \rightarrow g-2 (*)!!$

$$\omega_a \equiv \omega_s - \omega_c = -a_\mu \frac{eB}{m_\mu}$$

*in reality you measure ω_P - the Larmor frequency of the free proton - instead of $\mathbf{B}$



$g = 2$

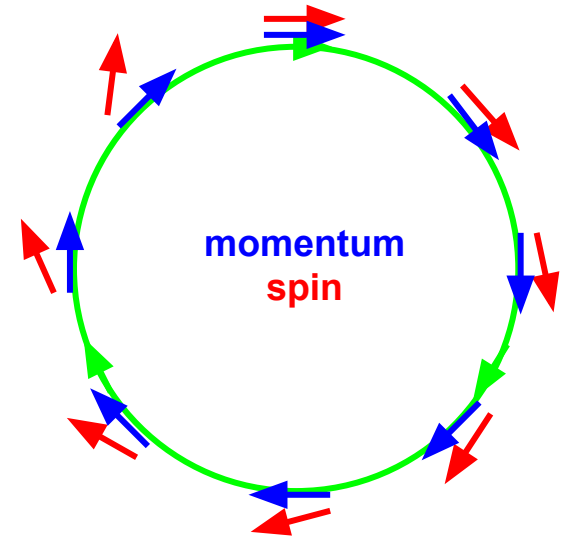


$g \neq 2$

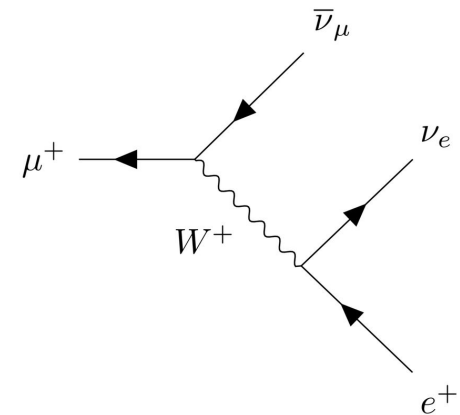
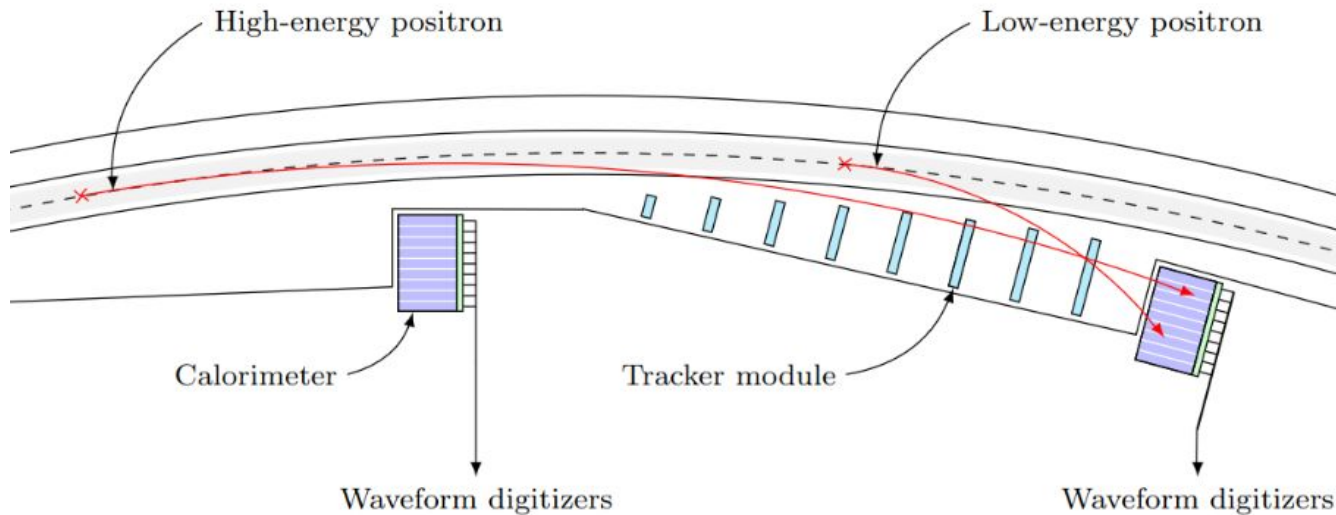
<https://muse.inf.infn.it/wp-content/uploads/2019/09/lusiani-patras-may19-v2.pdf>

How to measure g-2 (3)

- Muon is not stable and decays
- Measure energy of decay positrons
→ Most energetic if muon momentum & spin align!
- **If filter out most energetic positrons, see modulation with frequency ω_a !!!**

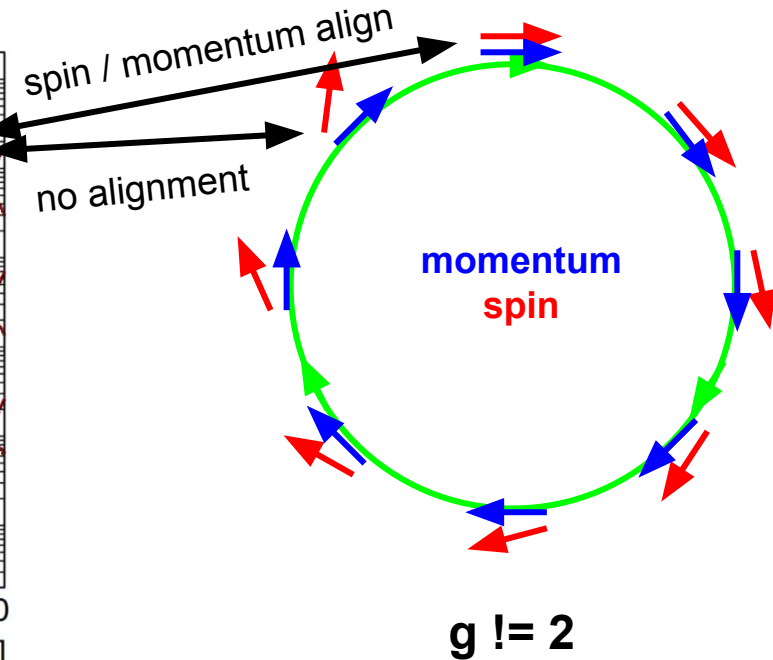
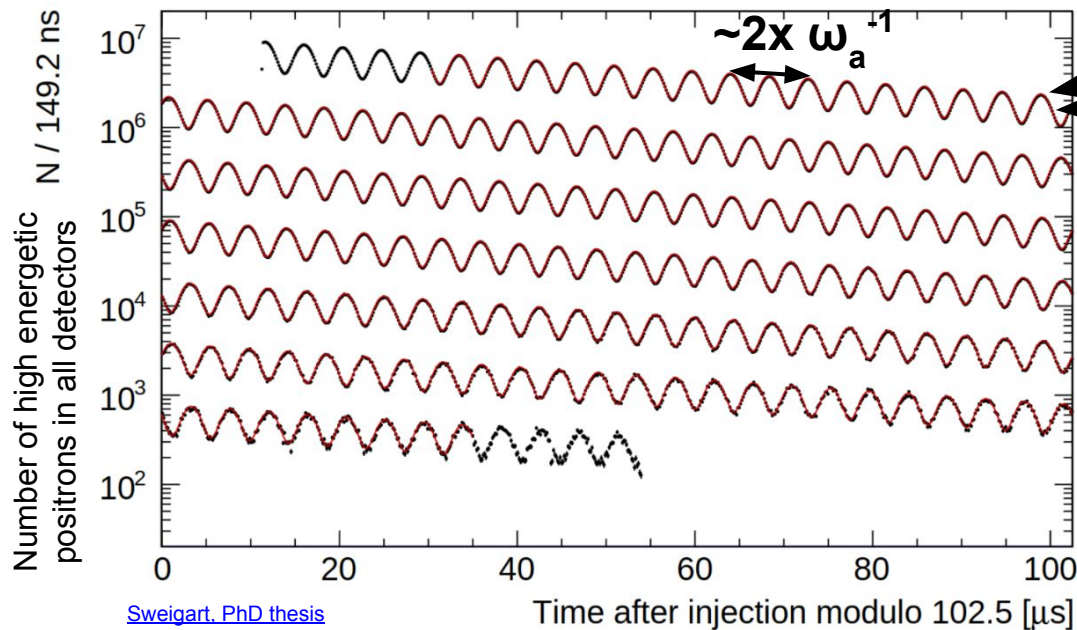


$g \neq 2$



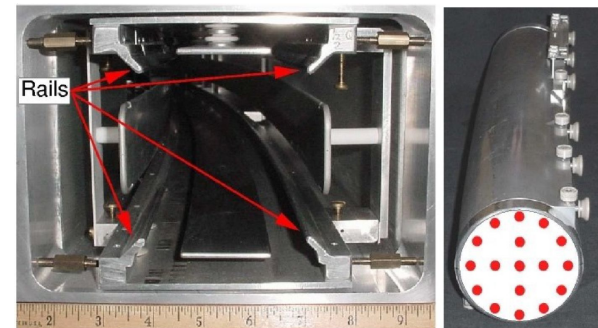
How to measure g-2 (4)

<https://www.particlebites.com/?p=8999>



- If filter out most energetic positrons, see modulation with frequency ω_a !!!

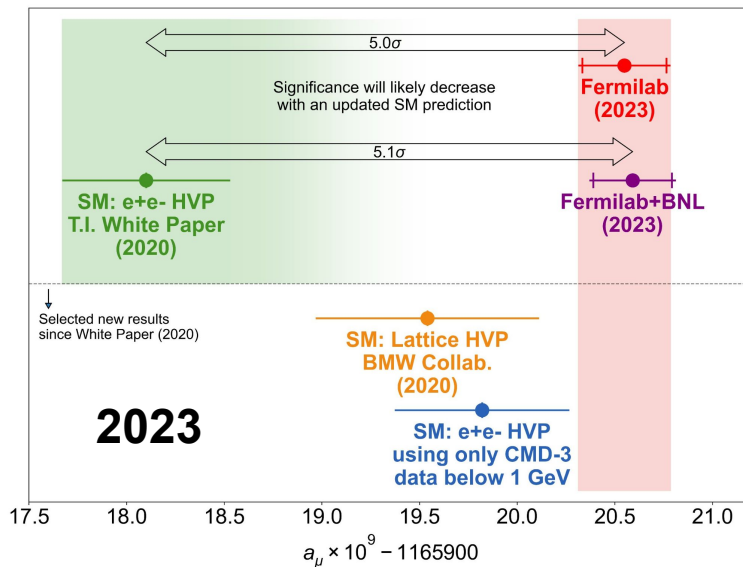
- Besides ω_a , must measure **B**!
- Measure magnetic field with special “trolley” measuring the field very precisely every three days & many other tools



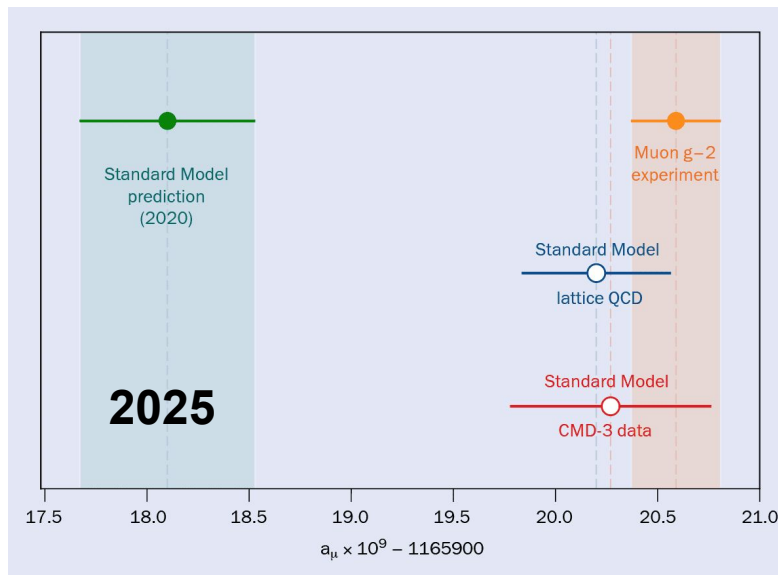
https://conference.ippp.dur.ac.uk/event/999/contributions/5220/attachments/4218/4977/Schreckenberger_Planck2021.pdf

Results of g-2

<https://pbs.twimg.com/media/F3LmADlWcAAIAxb?format=jpg&name=4096x4096>



<https://physicsworld.com/a/the-muons-magnetic-moment-exposes-a-huge-hole-in-the-standard-model-unless-it-doesnt/>



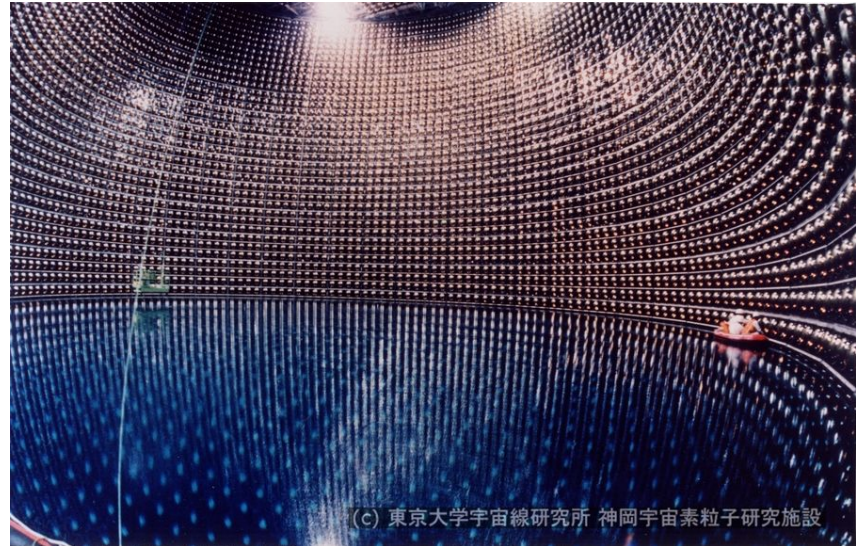
- g-2 measured extremely precisely to ~10 digits!
- In 2023: 5.1 σ discrepancy of theory pred. and meas.
 - theory pred. use data-driven approach
 - less discrepancy from lattice QCD
- In 2025: new measurement as input to theory pred.
 - good agreement in g-2
 - but meas. very different from previous ones (!?)
 - lattice result reproduced

g-2 seems resolved for now.

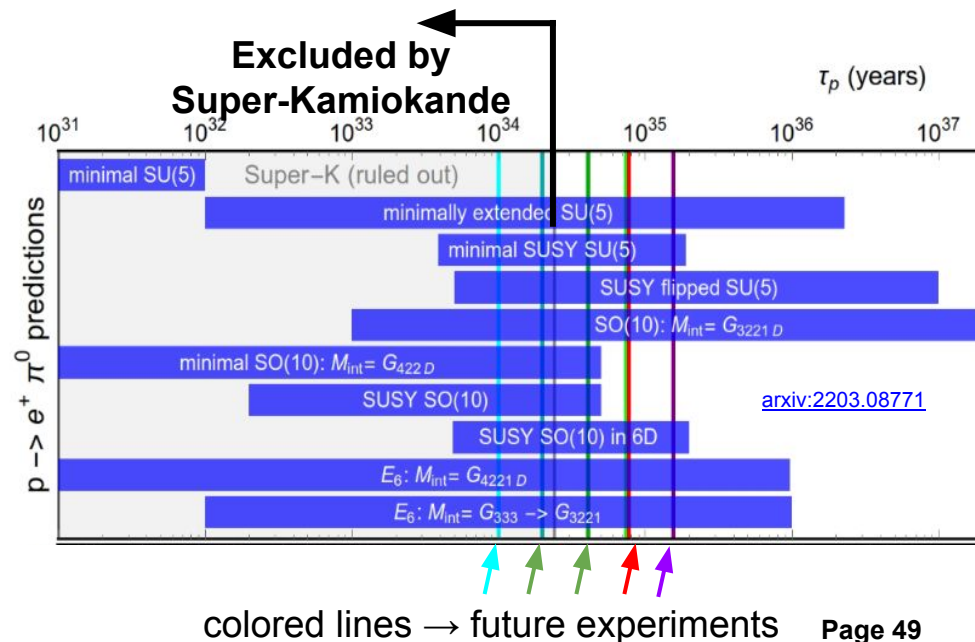
Proton stability

[arxiv:2010.16098](https://arxiv.org/abs/2010.16098)

- Proton is stable in SM (Baryon number conservation)
- Grand Unified Theories can predict that the proton is not stable
- “Not stable” $\rightarrow$ lifetimes $> 10^{30}$ y
- Search for proton decay e.g. with Super-Kamiokande
 - 27.5 T of Water $\rightarrow$ many protons
 - Search for $p \rightarrow \ell^+ \pi^0$
 - Proton lifetime $> 10^{34}$ y
- Multiple models excluded, but still room for SUSY ;)



<https://wechselljetzt.de/uploads/Super-K2-7d5e2bccc48802c113419e3c24de77c03662e919.jpg>



Summary

- Many, many searches / experiments targeting BSM physics
- Constraints on many model parameters, yet not clear what BSM physics is
- Many exciting experiments/results coming up investigating BSM physics further, e.g.
 - XENON n-ton → in operation
 - ALPS II → in operation
 - Hyper Kamiokande
 - HL-LHC
 - BabyIAXO
 - ...and many more!
- **What will we find?**



<https://thumbs.dreamstime.com/b/lama-closeup-funny-expression-face-mouth-open-appears-to-be-talking-humorous-96704071.jpg>

Thank you

Backup slides

SUSY-2018-08

A 3.5-inch hard disk drive (HDD) with its top cover removed, revealing the internal components. The drive is black and rectangular. The top cover is hinged on the left side and is lifted, showing the internal platters and the actuator arm. The platters are silver and circular, mounted on a central spindle. The actuator arm is yellow and white, with a read/write head at the end. The drive is shown from a top-down perspective, highlighting the internal mechanical parts.

ATLAS
 $\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$
 2-body selection
 $\text{SR}_{\text{DF}}^{2\text{-body}}$

Events / Bin

Legend:

- Data
- $t\bar{t}$
- $t\bar{t}Z$
- Diboson
- ▨ Standard Model
- ▨ Wt
- FNP
- Others

$\cdots t\bar{t}+\phi: m(\phi, Z) = (150, 1) \text{ GeV}$
 $\cdots t\bar{t}+a: m(a, \gamma) = (150, 1) \text{ GeV}$

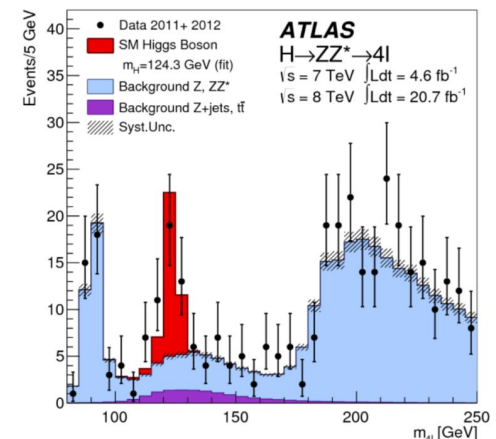
stop proxy invariant mass $m_{T2} [\text{GeV}]$

Data / SM

How compatible is the data with the Standard Model prediction? ([Data Analysis lecture](#)) $\ell = -2 \ln \left(\frac{\mathcal{L}(\text{data}, H_1)}{\mathcal{L}(\text{data}, H_0)} \right)$

Figure 1 is a plot showing the ratio of data to the Standard Model (Data / SM) versus the stop proxy invariant mass m_{T2} [GeV]. The y-axis ranges from 0 to 2, and the x-axis ranges from 100 to 280 GeV. Data points with error bars are shown, along with a horizontal line at Data / SM = 1. Two shaded regions indicate the fit range for the BSM model. Blue arrows point from the text "How much BSM could fit here? (none)" to the data points at higher m_{T2} values.

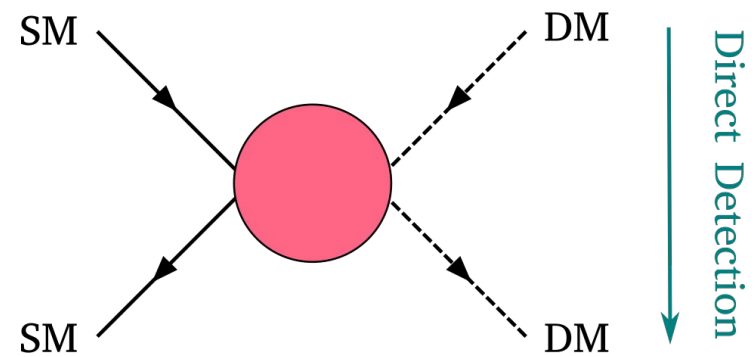
https://media.istockphoto.com/id/133838482/photo/foto/champaign.jpg?s=170667a&w=0&k=20&c=e7OMitD6AU5BnKQqnhD4GV70YpIZ_QiCGJVJhvp6Z8=



Direct detection Dark Matter searches

[1903.03026](#)

- Assume: permanent flow of DM particles through planet earth
- DM scatters with SM particles → **measurable!**
- DM scatters with nuclei as no EM int.
- No. scattered DM particles depends:
 - Density of DM → from rotation curve: $\rho_0 = 0.3 \text{ GeV}/c^2/\text{cm}^3$ ($\pm 50\%$)
 - DM velocity w.r.t. planet earth, $\sim 220 \text{ km/s}$ with some variation
 - DM-nuclei cross-section
 - Mass of DM & interacting nuclei
 - Spin of DM particle and sensitivity of target material



There is an ideal ratio of m_χ / m_N !

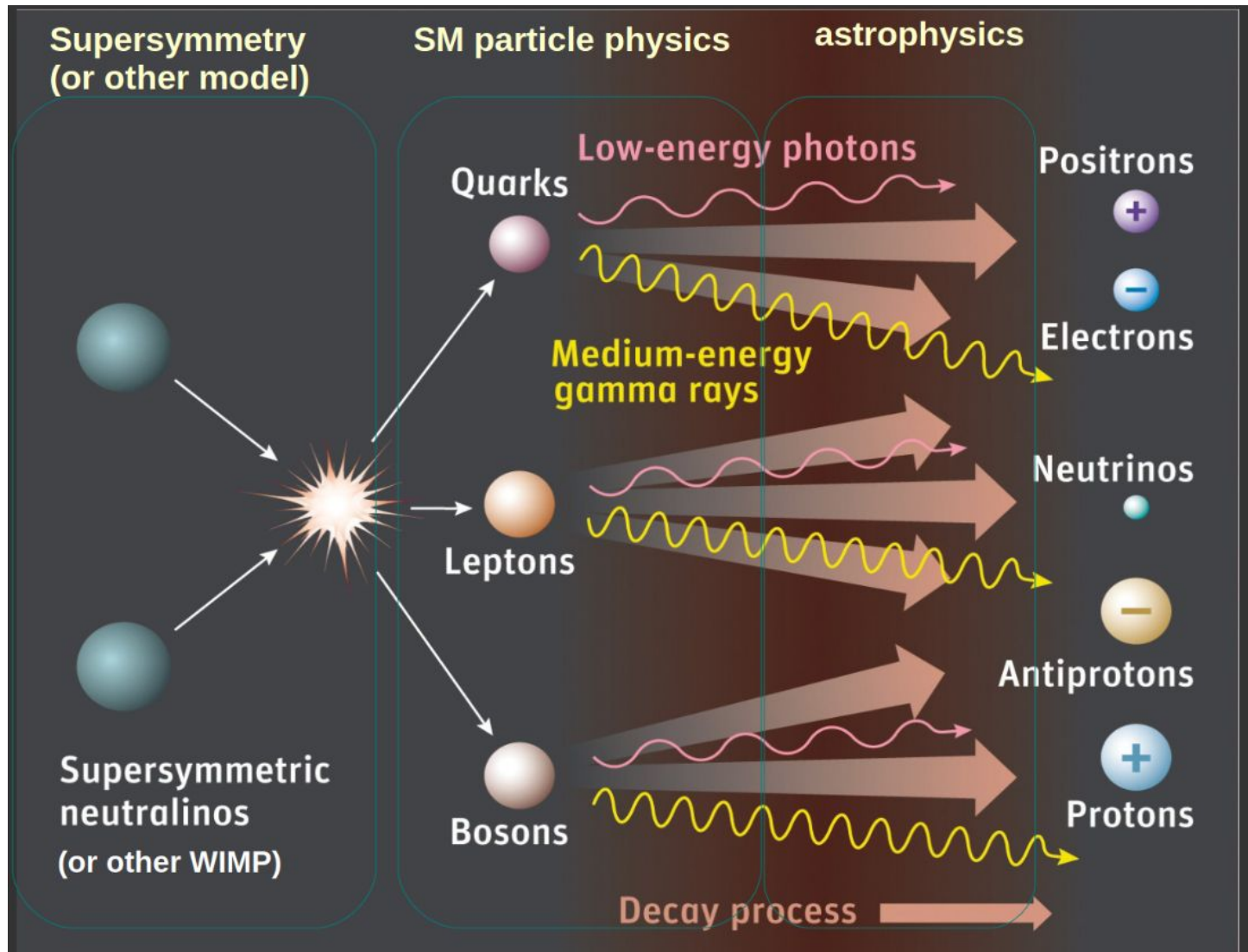
Recoil energy
→ **higher = less**

$$\frac{dR}{dE_{\text{nr}}} \propto \exp \left(-\frac{E_{\text{nr}}}{E_0} \frac{4m_\chi m_N}{(m_\chi + m_N)^2} \right)$$

Events per recoil energy

Most probably kinetic energy of DM (depends on velocity)
→ **higher = more**

Indirect detection all in one slide



https://www.mpi-hd.mpg.de/lin/events/isapp2011/pages/lectures/de_los_Heros.pdf

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