

Renormalisation scheme dependence of the trilinear Higgs coupling in extended scalar sectors

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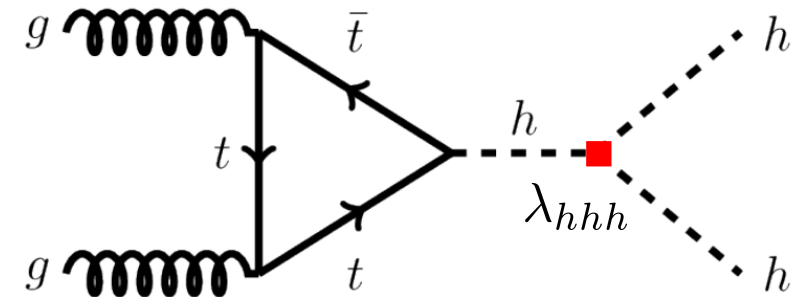
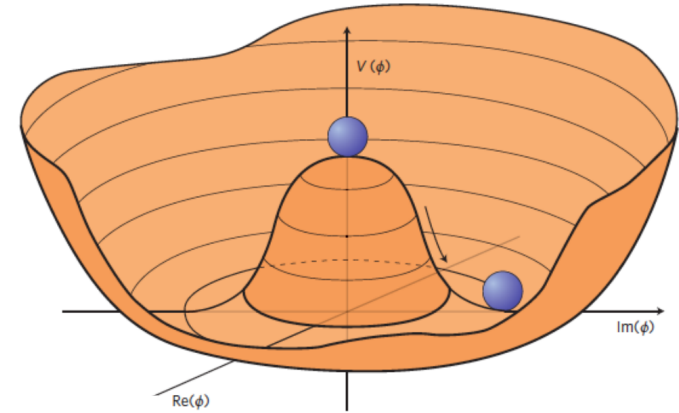
Motivation: Trilinear Coupling

[Bahl, Braathen, Gabelmann, Weiglein 2305.03015]

- Trilinear Higgs coupling determines the shape of the Higgs potential:
 - crucial for understanding dynamics of electroweak phase transitions

$$V_{\text{SM}} \supset \frac{m_h^2}{2} h^2 + \underbrace{\frac{3m_h^2}{v}}_{\lambda_{hhh}^{\text{SM},0}/6} h^3 + \frac{3m_h^2}{v^2} h^4$$

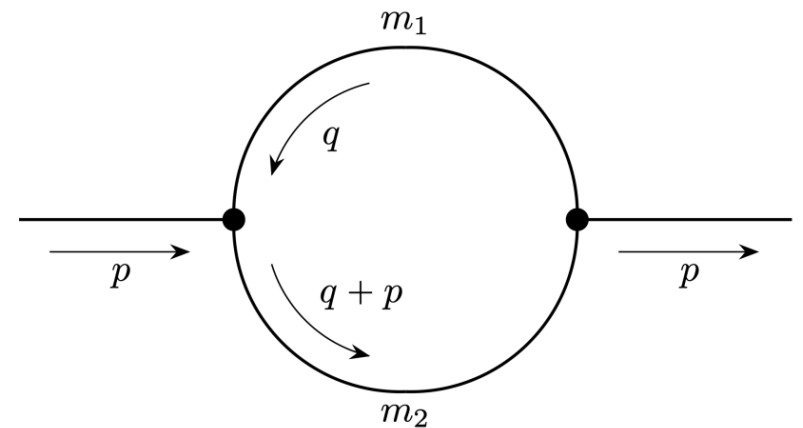
- sensitive to BSM physics
- exp. limit: $-1.2 < \frac{\lambda_{hhh}}{\lambda_{hhh}^{\text{SM},0}} < 7.2$ at 95% C.L.
 [ATLAS, 2406.09971]
 [CMS, 2407.13554]



Motivation: Precision Calculations

- reliable and numerically stable theoretical predictions of physical observables are imperative
 - important for comparison with experimental data
 - needed for consistent constraints on parameter space of BSM models

understanding of renormalisation scheme dependence needed



Singlet Extension of the Standard Model (SSM)

[Bahl, Braathen, Gabelmann, Weiglein 2305.03015]

- we consider the real singlet extension of the SM

$$V(\Phi, S) = \mu^2 |\Phi|^2 + \frac{\lambda}{2} |\Phi|^4 + \kappa_{SH} |\Phi|^2 S + \frac{\lambda_{SH}}{2} |\Phi|^2 S^2 + \frac{m_S^2}{2} S^2 + \frac{\kappa_S}{3} S^3 + \frac{\lambda_S}{2} S^4$$

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- after symmetry breaking: $\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2} G^+ \\ v+h+iG \end{pmatrix} \quad S = s + v_S$
- the fields mix to 2 CP-even eigenstates $h_{1,2}$
- the Lagrangian parameters are traded for $m_{h_1}, m_{h_2}, \alpha, v, v_S, \kappa_S, \kappa_{SH}$
- trilinear coupling at tree level: $\lambda_{h_1 h_1 h_1} = 3m_{h_1}^2 \left(\frac{\cos^3 \alpha}{v} + \frac{\sin^3 \alpha}{v_S} \right)$

5 new parameters

Choice of Renormalization scheme

OS renormalisation

- renormalised parameters correspond to measurable quantities
- OS condition: $m_0^2 = m_{\text{OS}}^2 + \delta m_{\text{OS}}^2$

$$\hat{\Sigma}(p^2 = m^2) = 0 \Rightarrow \delta m_{\text{OS}}^2 = \Sigma(p)|_{p^2=m^2}$$

renormalised
self energy

self energy

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$\overline{\text{MS}}$ renormalisation

- no conditions $\longrightarrow$ only UV divergent part and constants are subtracted

$$\Delta = \frac{2}{\epsilon} - \gamma_E + \ln 4\pi$$

- μ dependence $\longrightarrow$ running of parameters

$$m_B = m^{\overline{\text{MS}}}(\mu) + \delta m^{\overline{\text{MS}}}(\mu) = m^{\text{OS}} + \delta m^{\text{OS}}$$

Constructed renormalisation schemes

	m_{h_1}	m_{h_2}	α	κ_S	κ_{SH}	ν_S	Tads
OS	OS			MS		-	OS
MS	MS					-	OS
OSMSalpha	OS		MS			-	OS
OSkappas	OS					-	OS

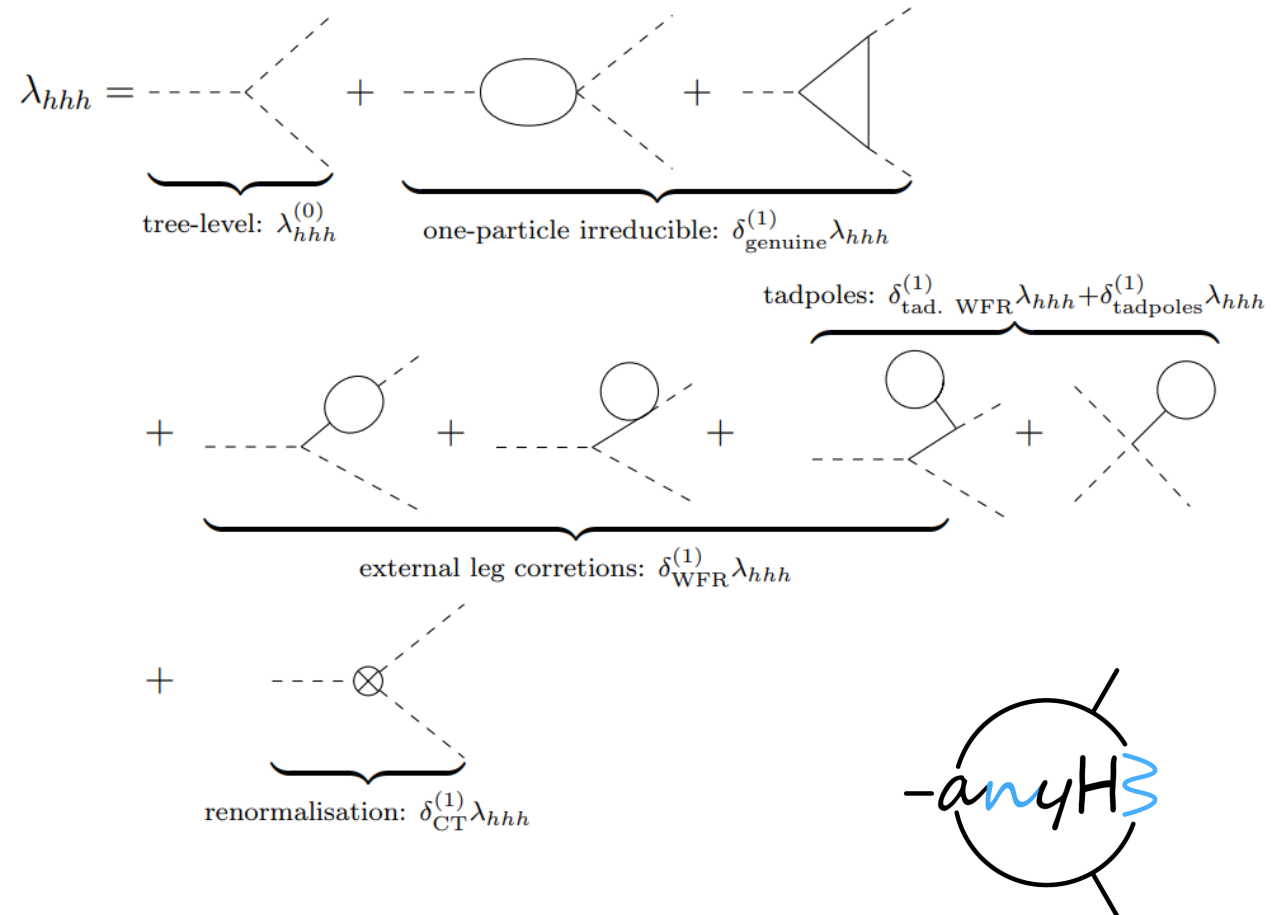
Computing the Trilinear Coupling: anyH3

[Bahl, Braathen, Gabelmann, Weiglein; 2305.03015]

- Python code anyH3 allows fast computation of λ_{hhh} at full one-loop level
- customize:
 - parameters/couplings
 - renormalisation
 - tadpole treatment
- applicable to all renormalisable models via UFO files

[Degrande et al.; 1108.2040]

[Darmé et al.; 2304.09883]



Analytic approximations

Applied simplifications:

1. set of particles: $h_1, h_2, G^\pm, G^0$

2. massless Goldstones

3. gaugeless limit: $m_W, m_Z \rightarrow 0$

4. α expansion

5. $m_{h_1} = 0 \rightarrow$ BSM contribution

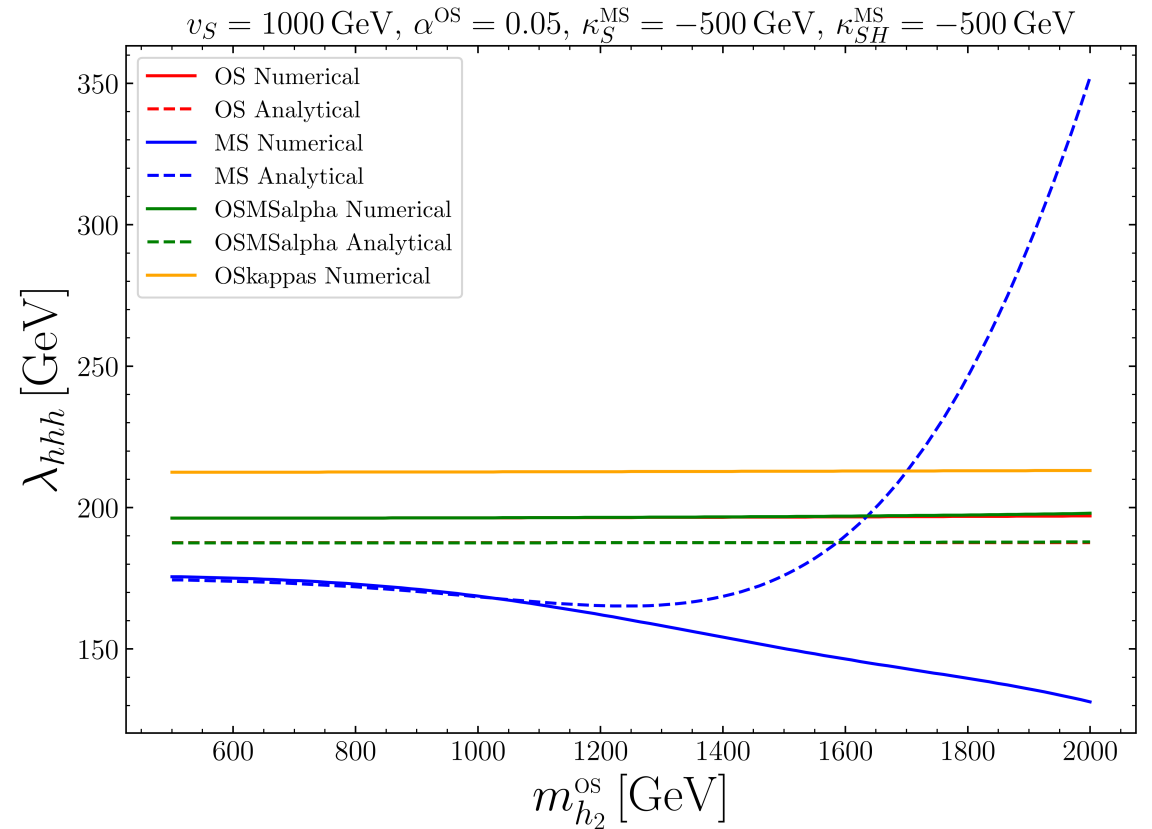
- Additional treelevel expression derived

$$\lambda_{hhh}^{OS,(1)} = \frac{\kappa_{SH}^2 v^2}{64 m_{h_2}^2 \pi^2 v_S^4} \left[6\alpha m_{h_2}^2 (-4 + \sqrt{3}\pi) v_S - 2\kappa_{SH} v_S v \right. \\ \left. - \alpha (-3 + \sqrt{3}\pi) (2\kappa_S v_S^2 - 3\kappa_{SH} v^2) \right]$$

$$\lambda_{hhh}^{MS,(1)} = \frac{1}{64 m_{h_2}^2 \pi^2 v_S^4 v^2} \left\{ 6\alpha m_{h_2}^6 v_S^3 + 6\kappa_{SH} m_{h_2}^4 v_S^3 v \right. \\ - 6\alpha \kappa_{SH}^2 v_S (-2m_{h_2}^2 + \kappa_S v_S) v^4 - 2\kappa_{SH}^3 v_S v^5 + 9\alpha \kappa_{SH}^3 v^6 \\ - 6 \left[\kappa_{SH} m_{h_2}^2 v_S^2 v (m_{h_2}^2 v_S - \kappa_{SH} v^2) \right. \\ + \alpha (m_{h_2}^6 v_S^3 + 2\kappa_{SH} m_{h_2}^2 v_S^2 (2m_{h_2}^2 - \kappa_S v_S) v^2 \\ \left. \left. + \kappa_{SH}^2 v_S (9m_{h_2}^2 - 2\kappa_S v_S) v^4 + 3\kappa_{SH}^3 v^6) \right] \cdot \ln \left(\frac{m_{h_2}^2}{\mu^2} \right) \right\}$$

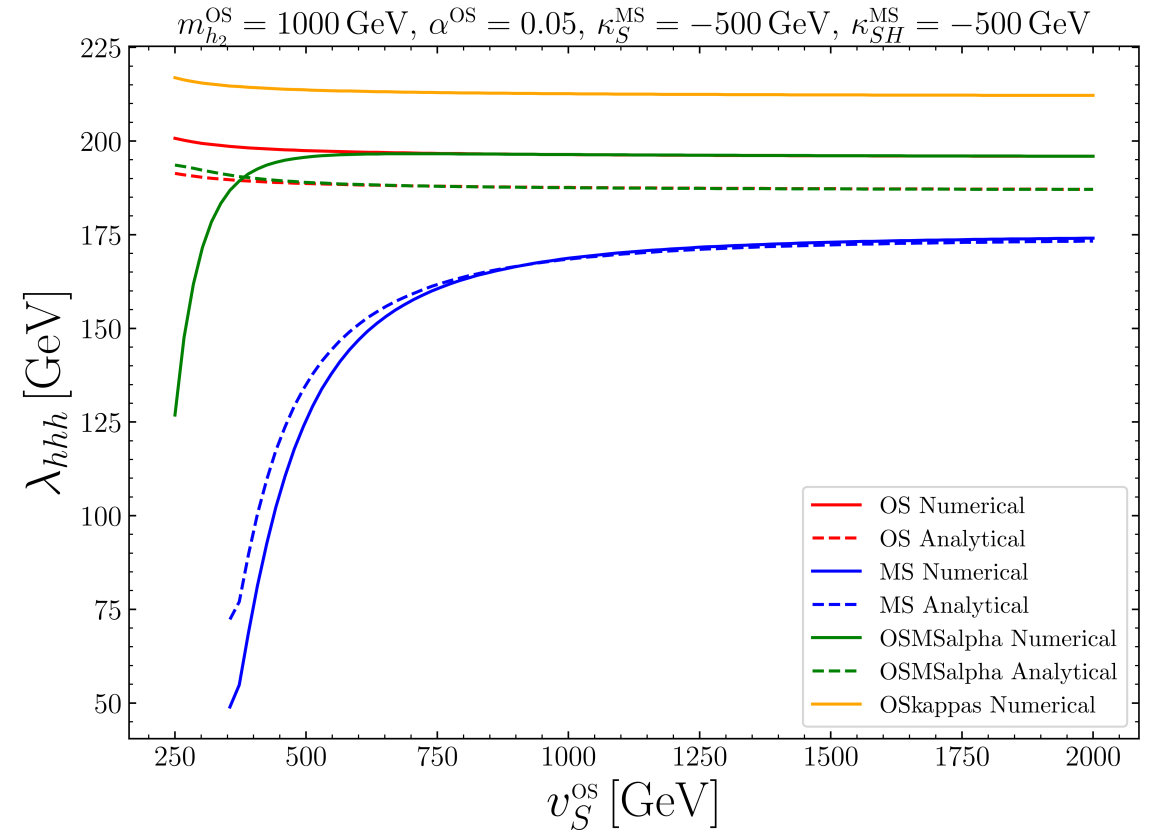
Scheme Dependence of the Trilinear Coupling

- different OS schemes show good
- constant shift between numerical and analytical calculation
- MS result differs for large BSM masses



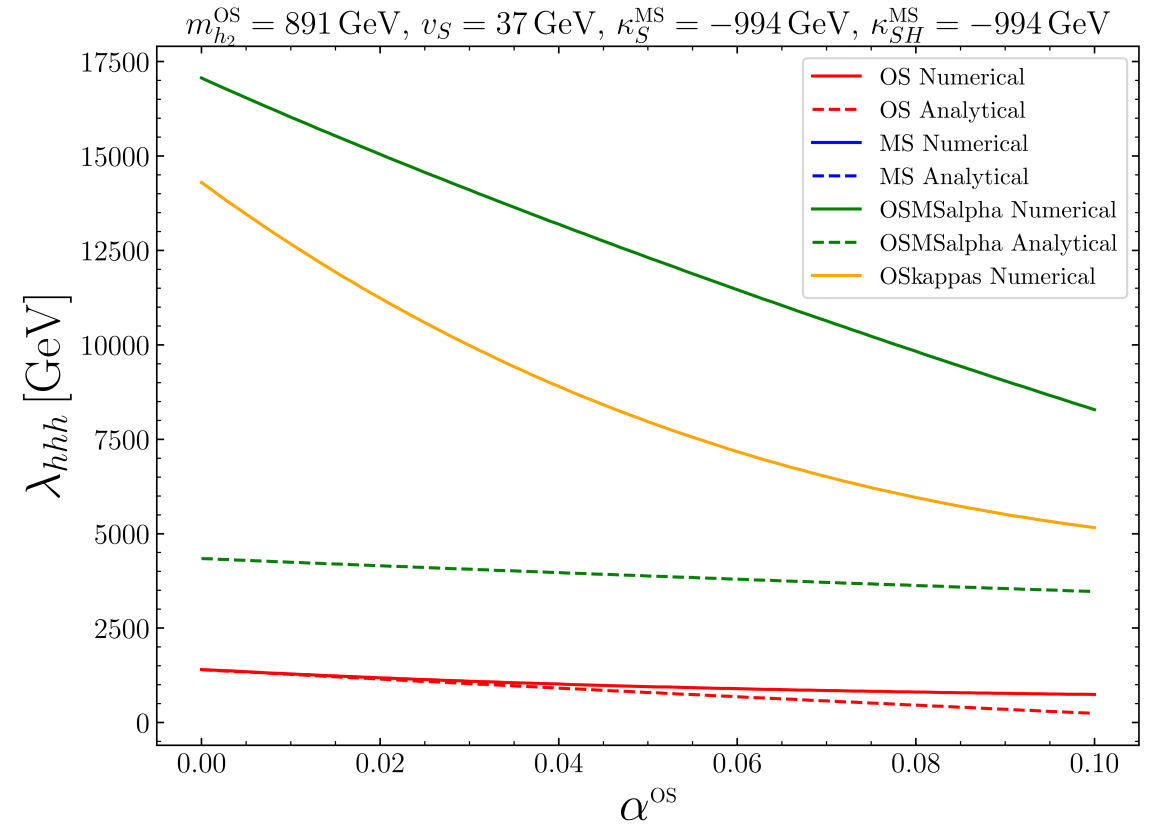
Scheme Dependence of the Trilinear Coupling

- high sensitivity to small singlets vevs
- $\overline{\text{MS}}$ conversion breaks down
- in general small deviations



Scheme Dependence of the Trilinear Coupling

- benchmark points with large corrections show a much greater uncertainty
- especially for small singlet vevs



Outlook

- Investigations of the renormalisation scheme dependence of the trilinear coupling in the SSM, derivation of analytic formulas
- Generally, larger differences between the schemes for large corrections to the coupling
- Extension to other BSM models e.g. 2HDM

Goal: derive a quantitative description of the uncertainty from the renormalisation scheme dependence → deliver a robust, consistent prediction for λ_{hhh}

Backup Slides

Tadpole Treatment

- Tadpole-free $\overline{\text{MS}}$ scheme:

$$T_h = t_h + \delta^{(1)} t_h = 0 \Rightarrow t_h = -\delta^{(1)} t_h = - \left. \frac{\partial V^{(1)}}{\partial h} \right|_{\text{min}}$$

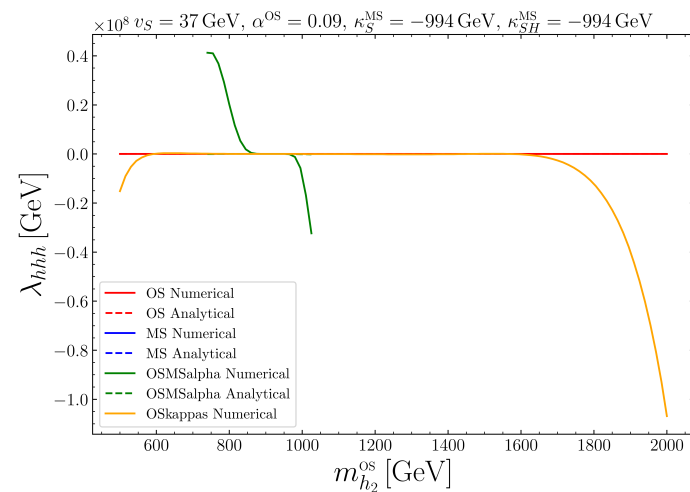
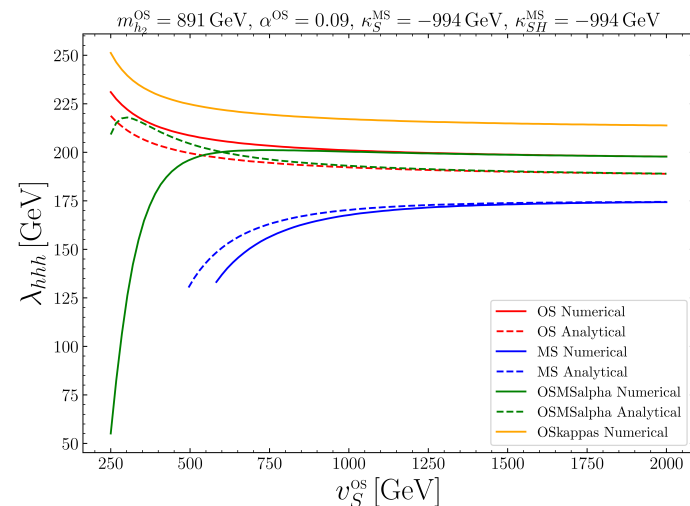
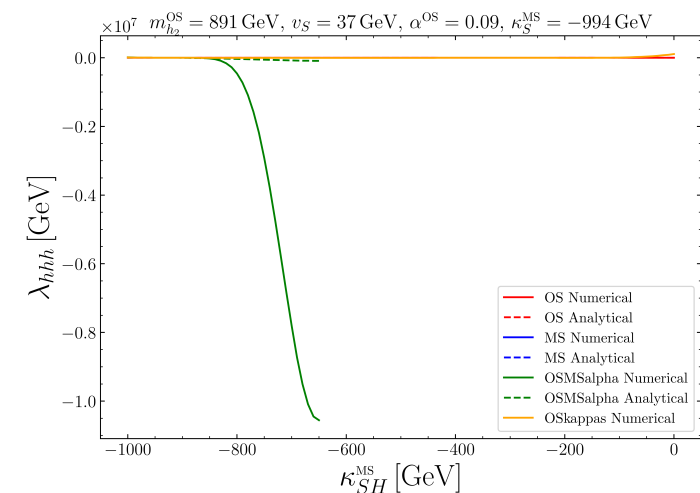
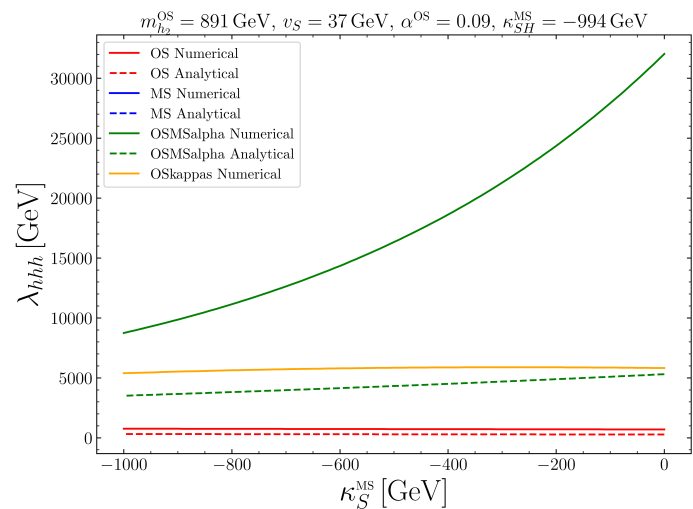
- OS tadpole renormalisation:

$$t_h = 0 \quad \delta_{\text{CT}}^{(1)} t_h = -\delta^{(1)} t_h$$

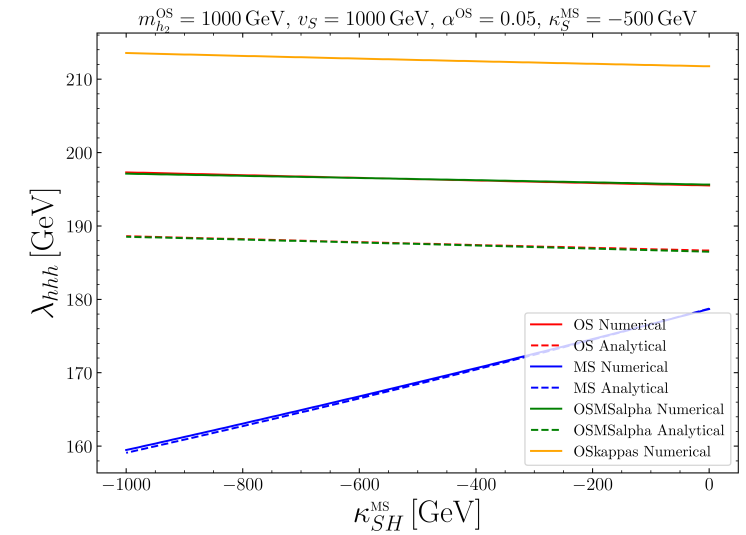
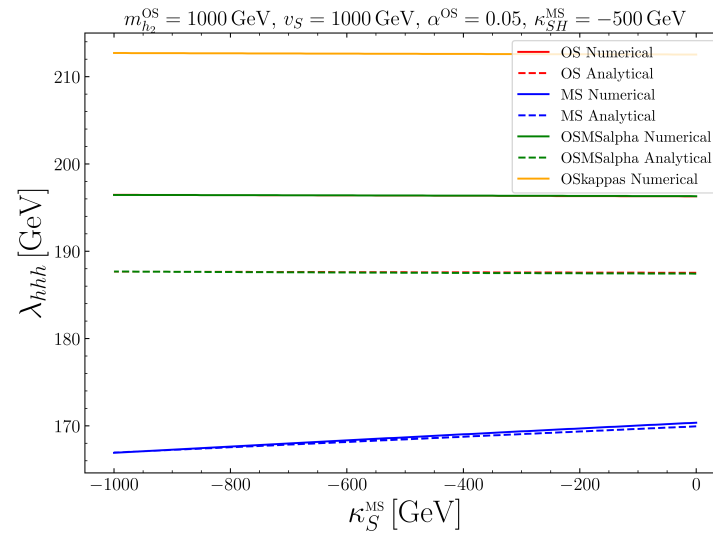
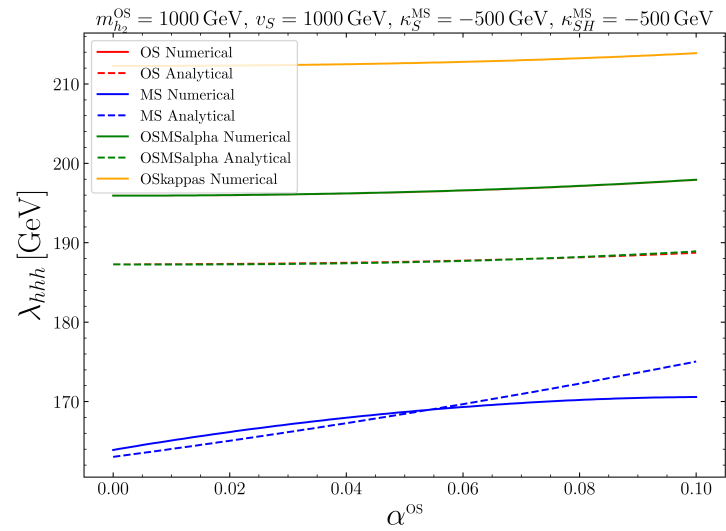
- $\overline{\text{MS}}$ tadpole renormalisation at the tree-level minimum (FJ):

$$t_h = 0 \text{ and } \overline{\text{MS}} \text{ renormalised tadpoles}$$

Benchmark Point



Test Point



Counter Terms

$$\delta\alpha = \frac{1}{2} \frac{\Sigma_{h_1 h_2}(p^2=m_{h_1}^2) + \Sigma_{h_1 h_2}(p^2=m_{h_2}^2)}{m_{h_2}^2 - m_{h_1}^2}$$

$$\delta\kappa_S = \frac{\frac{\partial\lambda_{HHH}^{(0)}}{\partial\kappa_{SH}} \left(\delta_{\text{gen+wfr}}^{(1)} \lambda_{hHH} + \sum_x \delta_x^{\text{CT}} \lambda_{hHH} \right) - \frac{\partial\lambda_{hHH}^{(0)}}{\partial\kappa_{SH}} \left(\delta_{\text{gen+wfr}}^{(1)} \lambda_{HHH} + \sum_x \delta_x^{\text{CT}} \lambda_{HHH} \right)}{\frac{\partial\lambda_{hHH}^{(0)}}{\partial\kappa_S} \frac{\partial\lambda_{HHH}^{(0)}}{\partial\kappa_{SH}} - \frac{\partial\lambda_{hHH}^{(0)}}{\partial\kappa_{SH}} \frac{\partial\lambda_{HHH}^{(0)}}{\partial\kappa_S}}.$$

$$\delta\kappa_{SH} = \frac{\frac{\partial\lambda_{hHH}^{(0)}}{\partial\kappa_S} \left(\delta_{\text{gen+wfr}}^{(1)} \lambda_{HHH} + \sum_x \delta_x^{\text{CT}} \lambda_{HHH} \right) - \frac{\partial\lambda_{HHH}^{(0)}}{\partial\kappa_S} \left(\delta_{\text{gen+wfr}}^{(1)} \lambda_{hHH} + \sum_x \delta_x^{\text{CT}} \lambda_{hHH} \right)}{\frac{\partial\lambda_{hHH}^{(0)}}{\partial\kappa_S} \frac{\partial\lambda_{HHH}^{(0)}}{\partial\kappa_{SH}} - \frac{\partial\lambda_{hHH}^{(0)}}{\partial\kappa_{SH}} \frac{\partial\lambda_{HHH}^{(0)}}{\partial\kappa_S}}$$