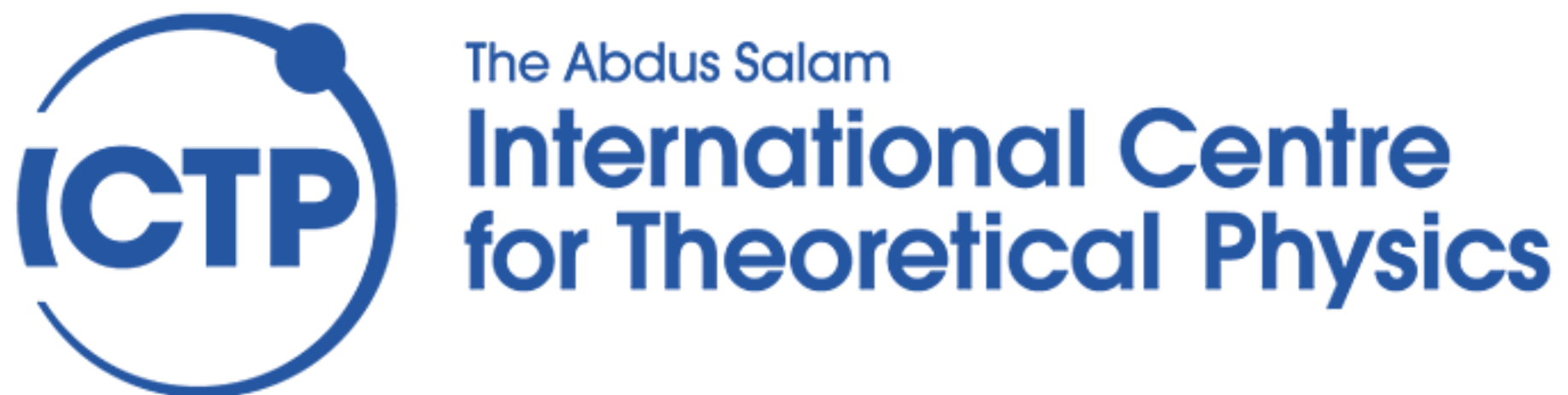


SymTFT for continuous spacetime symmetries.

Nicola Andrea Dondi (ICTP & INFN Trieste) - DESY Theory Workshop - 25/09/2025

Based on:

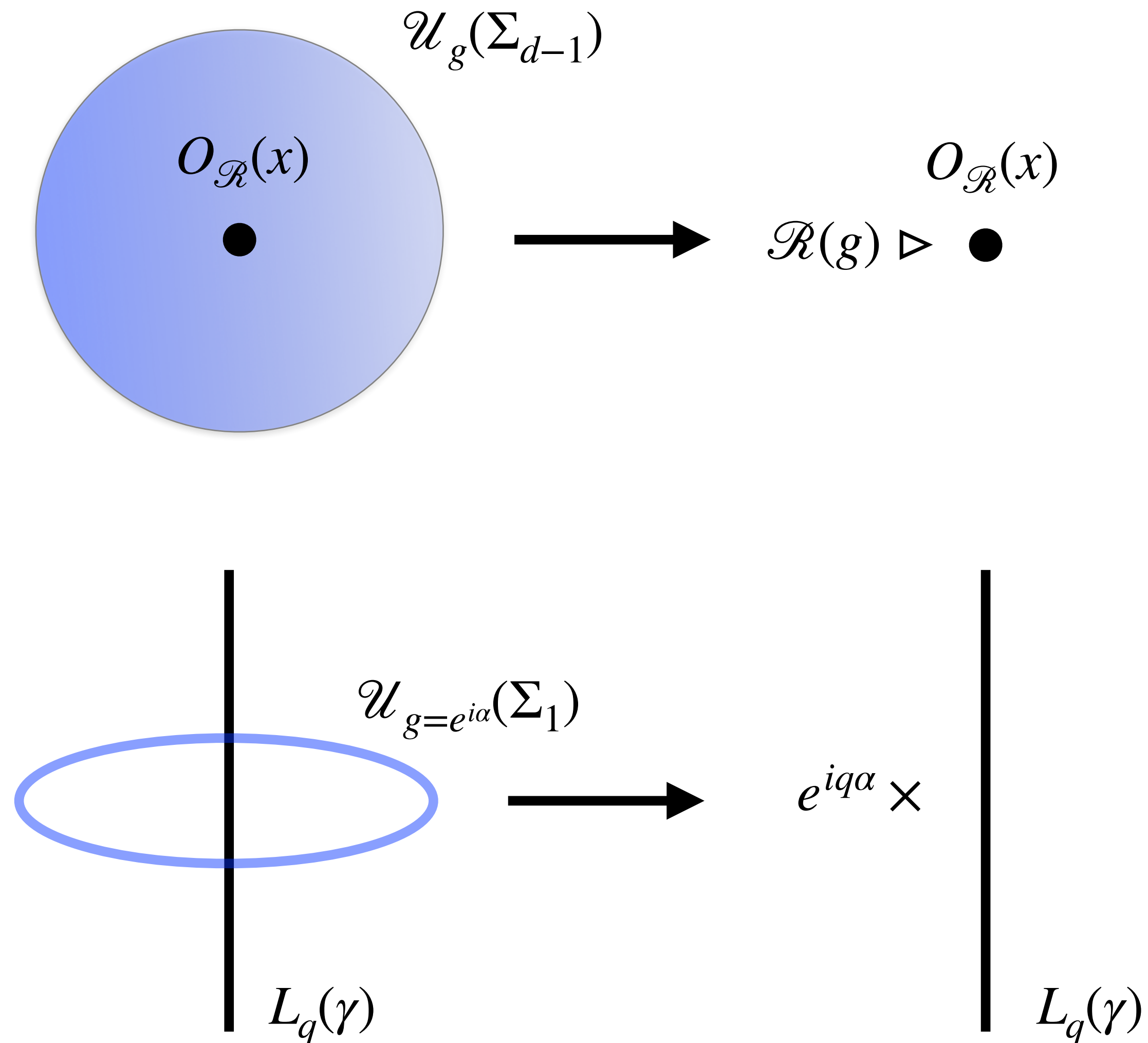
[2509.07965] with F. Apruzzi, I. García Etxebarria, H. T. Lam, S. Schäfer-Nameki



Symmetries and topological operators

[Gaiotto, Kapustin, Seiberg, Willet '14]

[Kapustin, Seiberg '14]



**Topological Defect
Operators (TDO)**

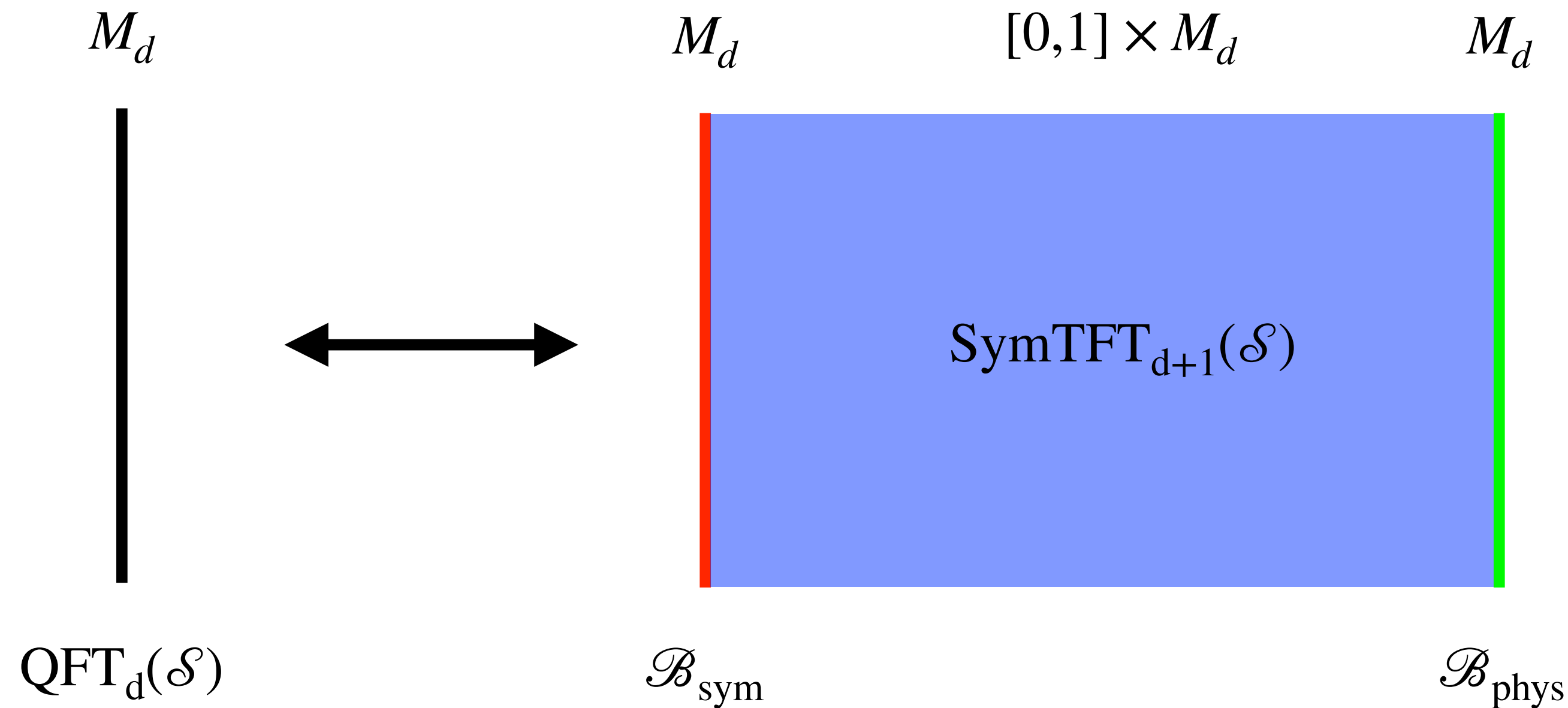


Symmetries

- Extension to non-invertible / categorical symmetries
- Can we incorporate spacetime symmetries?

[see also discussion in Harlow, Ooguri '19]

SymTFTs: overview



$\text{SymTFT}_{d+1}(\mathcal{S})$: TQFT flat-gauging $\mathcal{S}$ in one higher dimension.

- $\mathcal{B}_{\text{phys}}$: gapless / gapped boundary condition of $\text{SymTFT}_{d+1}(\mathcal{S})$ containing dynamical information on the original theory $\text{QFT}_d(\mathcal{S})$
- $\mathcal{B}_{\text{sym}}$: gapped boundary condition of $\text{SymTFT}_{d+1}(\mathcal{S})$.
- Bulk topological operator corresponds to symmetry operators / charged objects depending on their endability.

Recent progress

- Finite group symmetries: symmetry theory is Dijkgraaf-Witten-type.

[Freed, Moore, Teleman 22], [Apruzzi, Bonetti, García Etxebarria, Hosseini, Schäfer-Nameki '21], [Gaiotto, Kulp '20]...

- Continuous group symmetries: typically BF-type theories.

[Brennan, Sun 24], [Bonetti, Del Zotto, Minasian '24], [Apruzzi, Bedogna, Dondi '24], [Antinucci, Benini '24]...

Useful for studying various symmetry properties in a theory-independent ways:

- Encode different symmetry global variants

[Apruzzi, Bonetti, García Etxebarria, Hosseini, Schäfer-Nameki '21]...

- Classify anomalies and higher-charges.

[Bhardwaj, Schafer-Nameki 23], [Kaidi, Nardoni, Zafrir, Zheng '23]...

- Construct gapped and gapless phases

[Bhardwaj, Bottini, Pajer, Schafer-Nameki, 24], [Bonetti, Del Zotto, Minasian '25]...

Non-abelian BF theories: Generalities

$$S = \frac{i}{2\pi} \int_{M_{d+1}} \langle B_{d-1}, F \rangle, \quad F = dA + A \wedge A.$$

$$\text{EOM: } d_A B_{d-1} = F = 0.$$

- $A : G$ -connection.
- $B_{d-1} : \mathfrak{g}$ -valued $(d-1)$ -form.
- $\langle \dots, \dots \rangle : \text{Adjoint-invariant bilinear on the algebra.}$

This theory is topological and has an extended gauge symmetry:

$$\delta_{(\epsilon, \sigma)} A = d_A \epsilon$$

$$\delta_{(\epsilon, \sigma)} B_{d-1} = [B_{d-1}, \epsilon] + d_A \sigma_{d-2}$$

$$\mathcal{L}_\xi A = \iota_\xi(F) + d_A(\iota_\xi A)$$

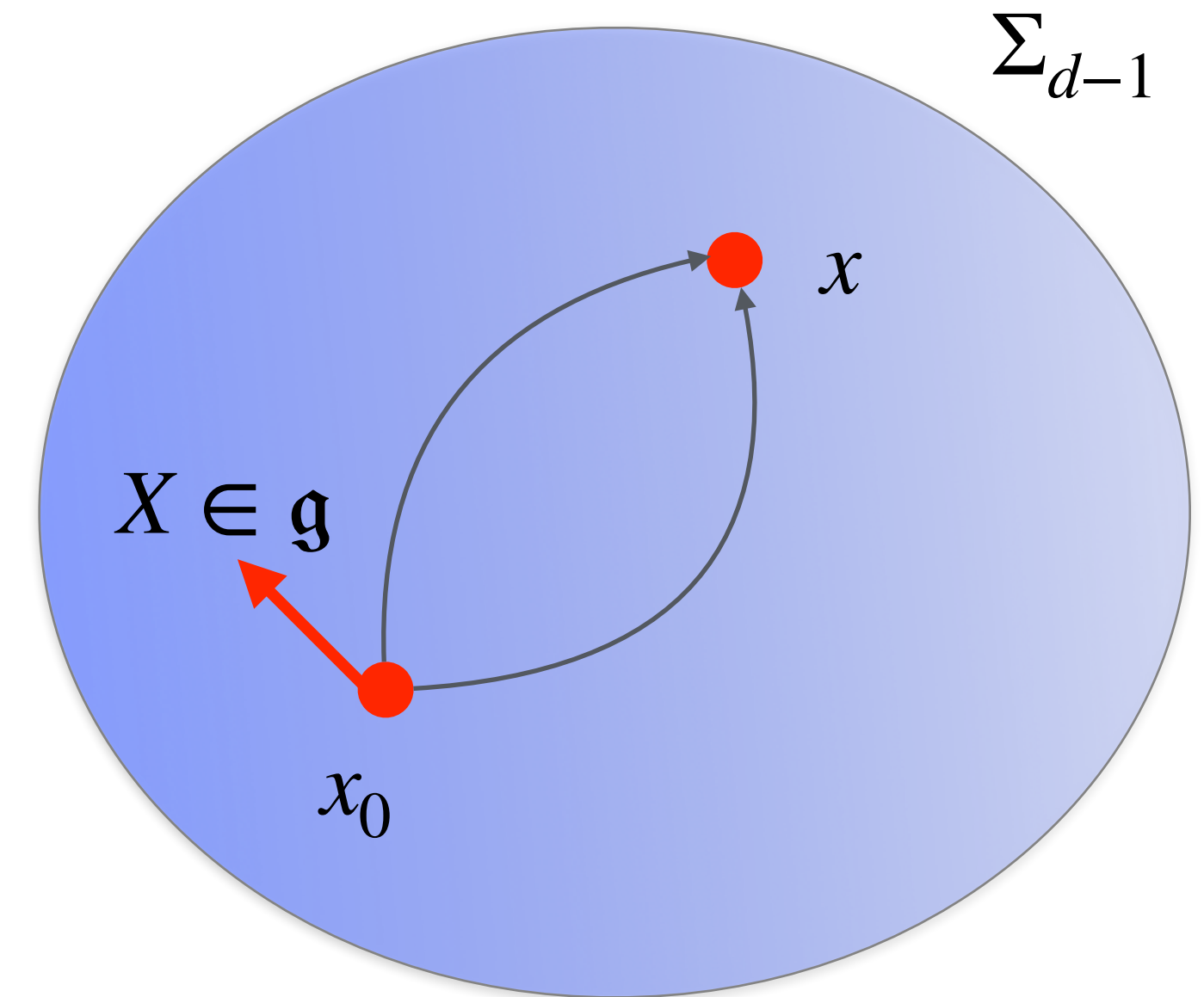
$$\mathcal{L}_\xi B_{d-1} = \iota_\xi(d_A B_{d-1}) + [\iota_\xi A, B_{d-1}] + d_A(\iota_\xi B_{d-1})$$

On flat connections $\text{Diff} \simeq \text{Gauge}$

Non-abelian BF theories: TDOs

Spectra of topological defects:

- $\mathcal{U}_{\mathcal{R}}(\Sigma_1) = \text{tr}_{\mathcal{R}} \mathcal{P} \exp \left\{ i \int_{\Sigma_1} A \right\}$
- $\mathcal{U}_{[e^X]}(\Sigma_{d-1}) = \int_{[X]} dX \exp \left\{ i \int_{\Sigma_{d-1}} \langle \bar{\alpha}_0(x), B_{d-1} \rangle \right\}$
 $\bar{\alpha}_0 \in \Omega^0(\Sigma_{d-1}, \mathfrak{g}), \quad d_A \bar{\alpha}_0 = 0, \quad \bar{\alpha}_0(x_0) = X \in \mathfrak{g}$



Specific cases of a class of defects labeled by $([g], \mathcal{R}^{H_g})$ with H_g centralizer of $g \in G$.

A SymTFT for conformal symmetry

Proposal: BF theory for gauge group $SO(d + 1, 1)$. Gauge fields on boundaries acquire geometric interpretations

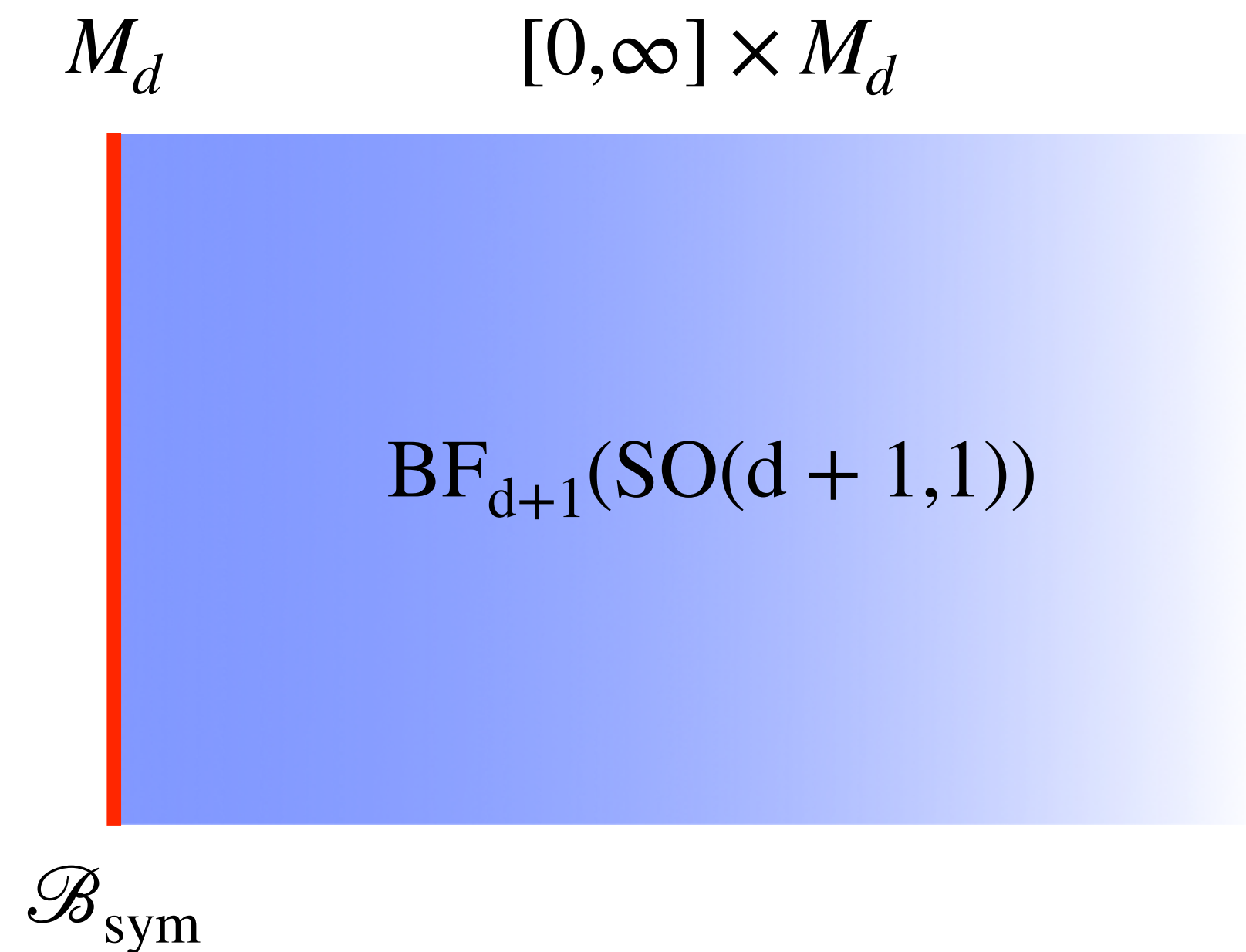
$$A|_{\mathcal{B}_{\text{sym}}} \equiv \mathcal{A} = e^a P_a + \frac{1}{2} \omega^{ab} L_{ab} + f^a K_a + \sigma D$$

In this talk, we take $\mathcal{B}_{\text{sym}}$ to be a *flat space*

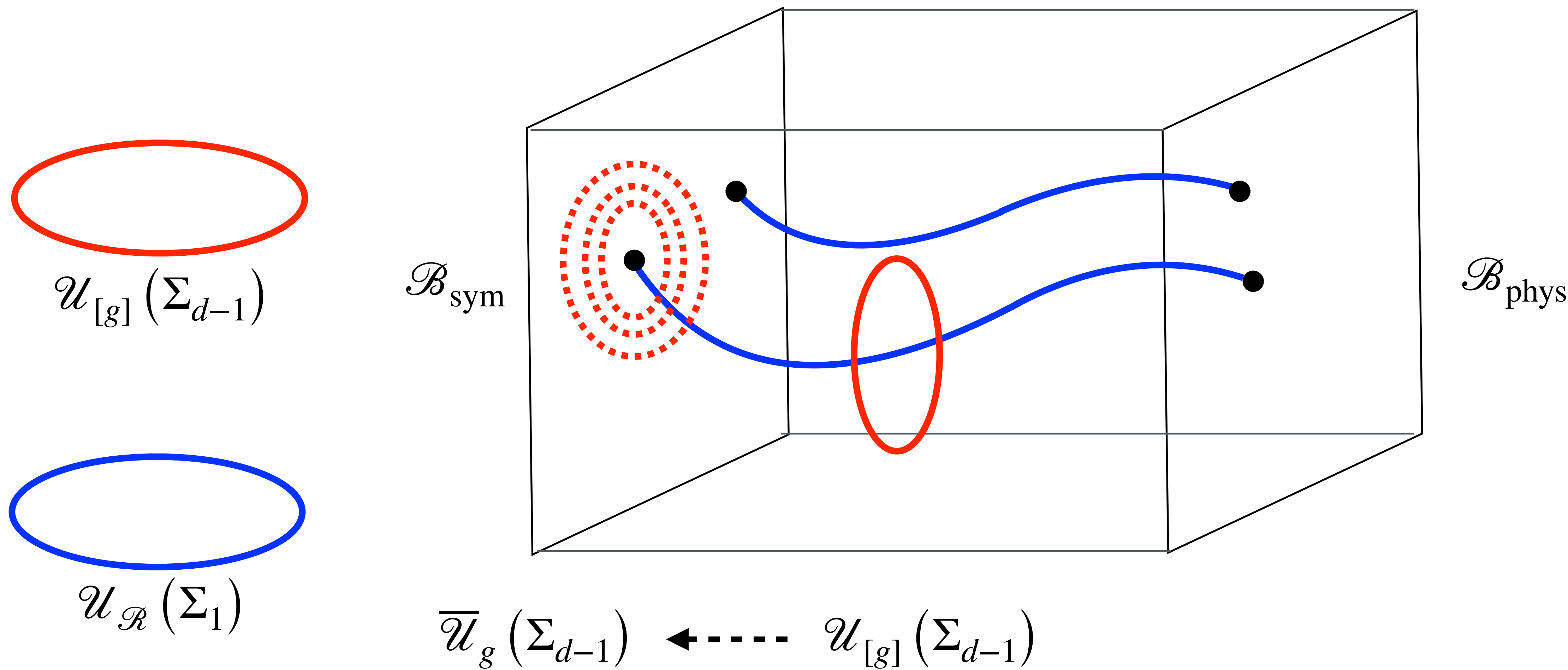
Dirichlet boundary condition:

$$S_\partial = \frac{i}{2\pi} \int_{\mathcal{B}_{\text{sym}}} \langle A - \mathcal{A}, B_{d-1} \rangle, \quad \mathcal{F} = 0$$

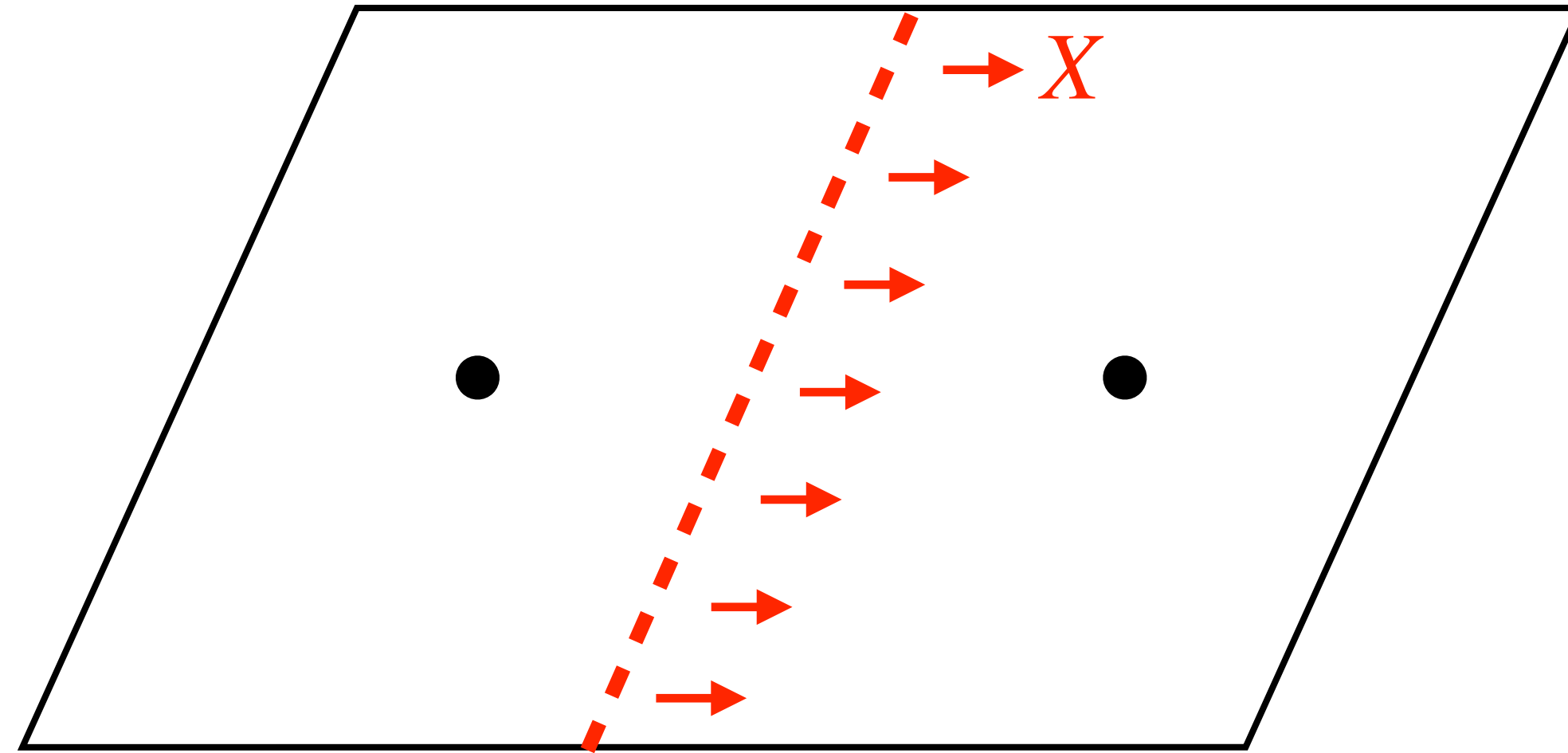
$$\mathcal{A} = \bar{e}^a \otimes P_a, \quad \bar{e}^a_\mu \in SO(d)$$



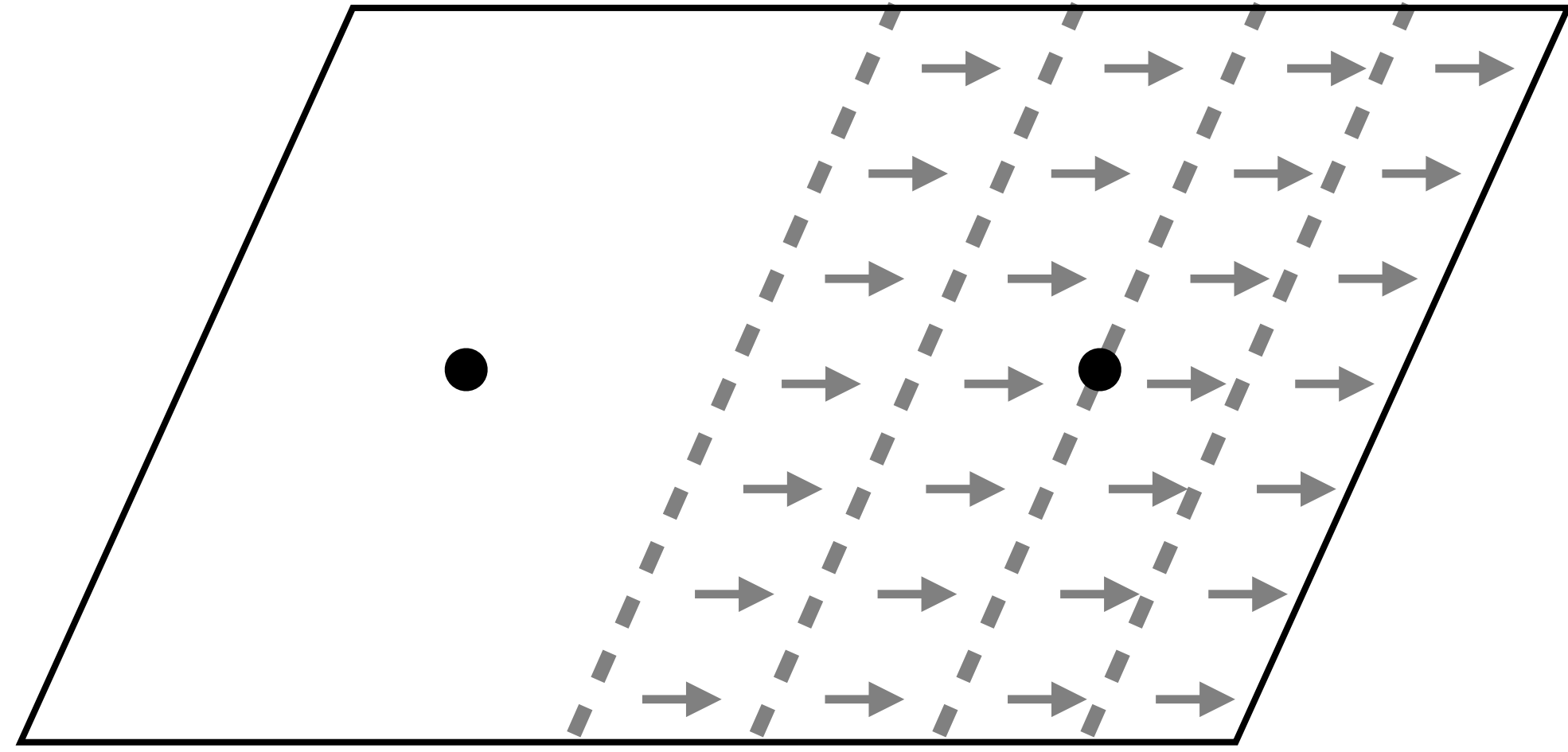
Linking action for TDOs



How to move a point



$$\mathcal{B}_{\text{sym}} : \mathcal{A} = \bar{e}^a \otimes P_a$$



$$\mathcal{B}_{\text{sym}} : \mathcal{A}^{(X)} = [\bar{e}^a + X \delta^{(1)}(\Sigma_{d-1})] \otimes P_a$$

The new background $\mathcal{A}^{(X)}$ is also flat $\implies$ we can find a gauge transformation back to $\mathcal{A}$, or equivalently, a diffeomorphism $\xi(X)$!

$$\langle O(x) \bar{\mathcal{U}}_X(\Sigma_{d-1}) O(y) \rangle_{\mathcal{A}} = \langle O(x) O(y) \rangle_{\mathcal{A}^{(X)}} = \langle O(x) e^{\mathcal{L}_{\xi(X)}} O(y) \rangle$$

Incorporating anomalies

In dimension $d + 1 = 2n + 1$, anomalies are incorporated by decorating the SymTFT with anomaly theories:

$$S^{(d=2)} = \frac{ik}{2\pi} \int_{I \times M_2} \left\{ \langle A, F \rangle - \frac{1}{3} \langle A, A^2 \rangle \right\} \quad S^{(d=4)} = \frac{ik}{6(2\pi)^2} \int_{I \times M_4} \left\{ \langle A, F, F \rangle - \frac{1}{2} \langle A, A^2, F \rangle + \frac{1}{10} \langle A, A^2, A^2 \rangle \right\}$$

$$\langle \dots \rangle : \text{Sym}^{n+1}(\mathfrak{g}) \rightarrow \mathbb{R}/\mathbb{C}, \quad \langle gX_1g^{-1}, \dots, gX_{n+1}g^{-1} \rangle = \langle X_1, \dots, X_{n+1} \rangle$$

For $SO(2n + 1, 1)$ with generators M_{ab} there can be more than one multilinear products:

- For $n = 2\ell + 1$, $\ell \in \mathbb{Z}_{\geq 0}$ there are two: symmetries version of $\sim (Tr[M_{ab}M^{ab}])^{\frac{n+1}{2}}$
and $\sim \epsilon_{a_1b_1\dots a_{n+1}b_{n+1}} M^{a_1b_1} \dots M^{a_{n+1}b_{n+1}}$ (gravitational anomaly + conformal type-A)
- For $n = 2\ell$, $\ell \in \mathbb{Z}_{\geq 0}$ there is only the ϵ -like (only conformal type-A).

Future directions

- Interplay between spacetime and internal (categorical) symmetries:
 - Extend the classification of gapless spacetime symmetry breaking phases.
 - Mixed gravitational - higher form anomalies? [Nicolis, Penco, Piazza, Rattazzi '15]
- Can we realise spacetime supersymmetries?
[Grassi, Penati, '25], [Ambrosino, Luo, Wang, Zhang '24], [Wang, Zhang '23]
- Are type-B conformal anomalies realised by bulk topological couplings?
[Imbimbo, Porro, '25], [Imbimbo, Rovere, Warman, '23]

Thank you for your attention!