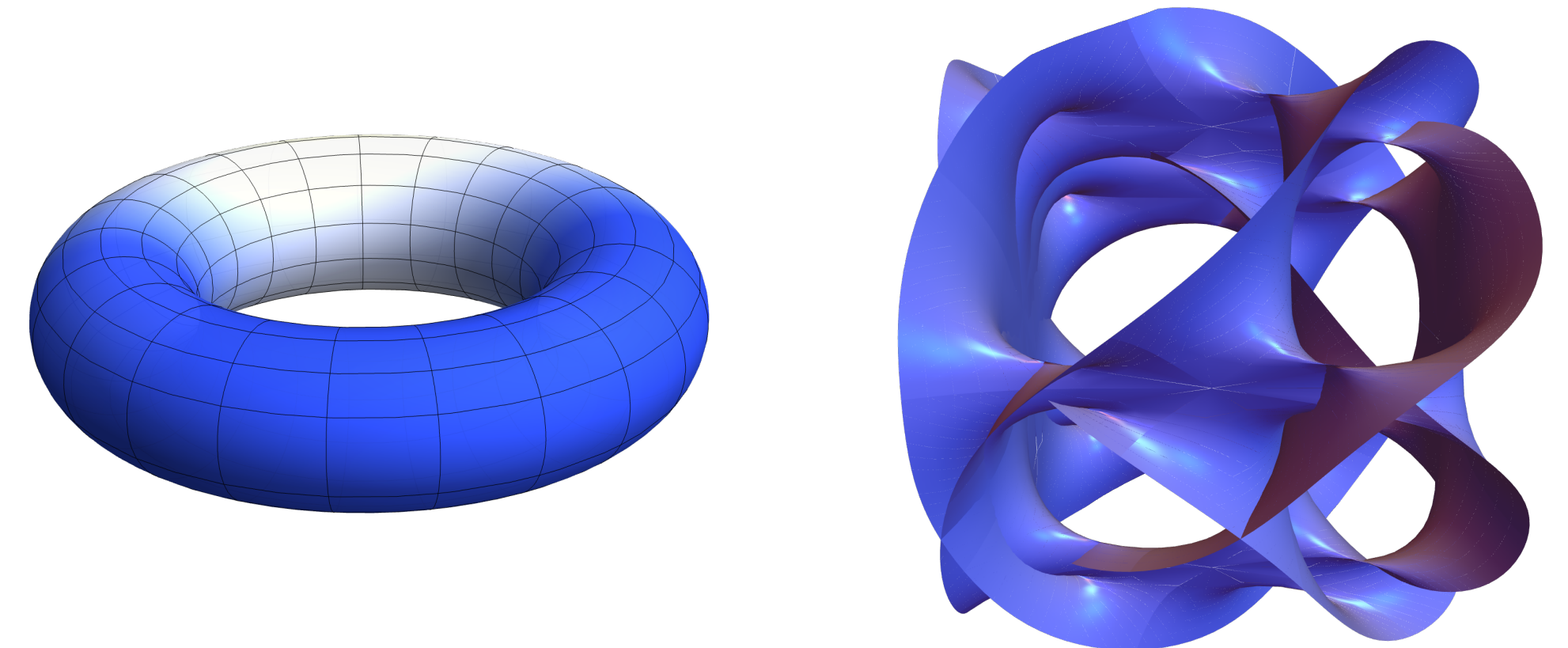


# Canonical Feynman integrals beyond polylogarithms

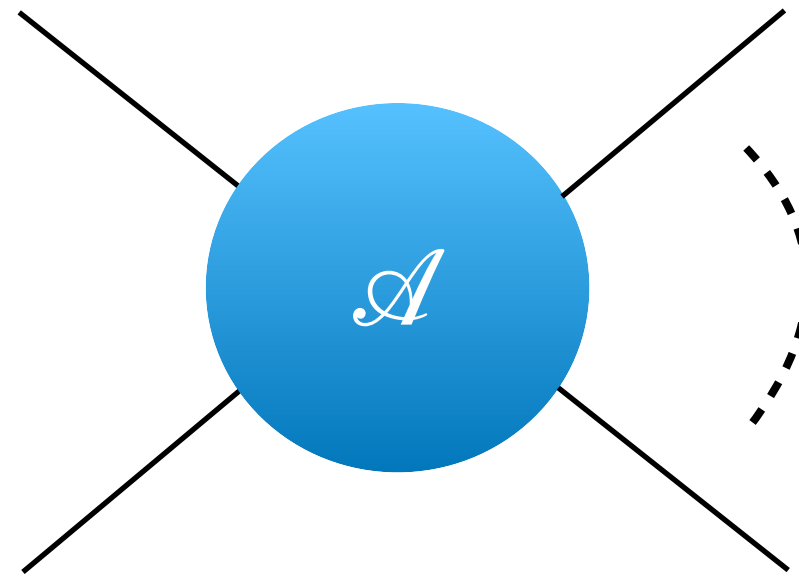
with **Claude Duhr, Christoph Nega, Benjamin Sauer,  
Lorenzo Tancredi, Fabian J. Wagner**

Based on: *JHEP* 06 (2025) 128, arXiv 2503.20655



# Motivation

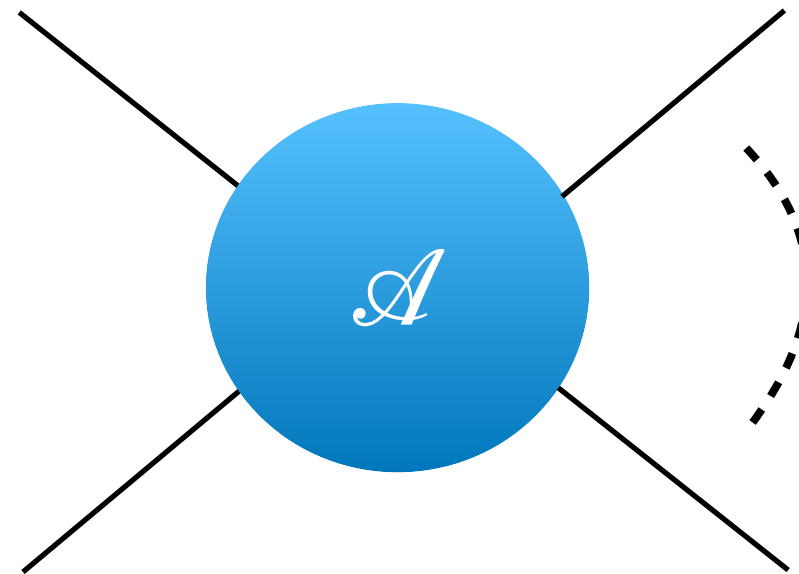
# Motivation





# Motivation

## Singularities

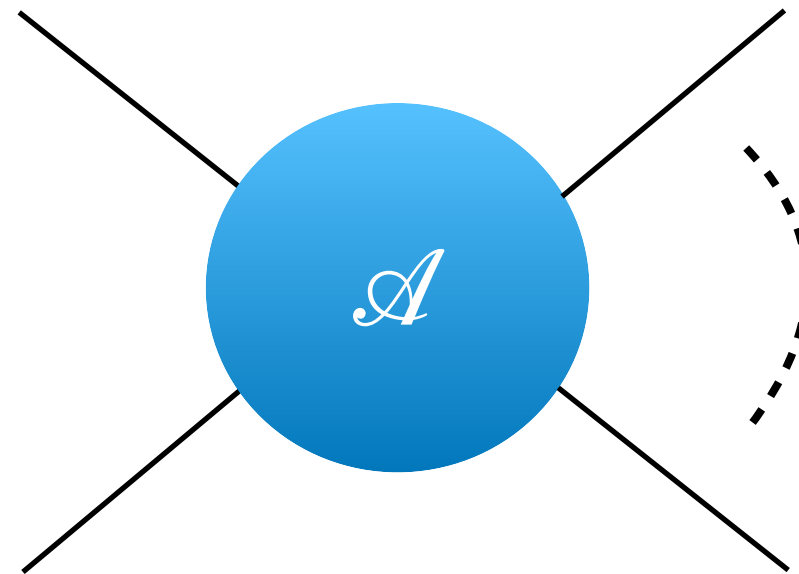




# Motivation

## Singularities

poles

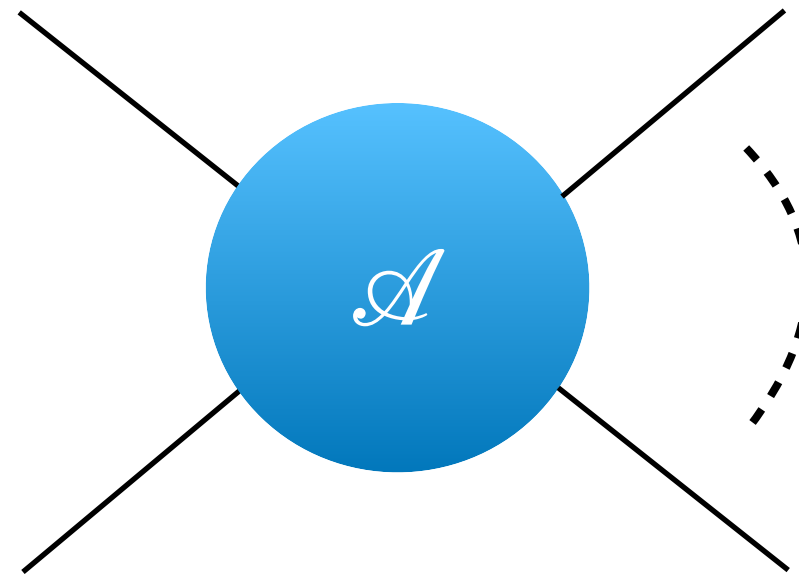


# Motivation

## Singularities

poles

*single*-particle  
states go on-shell

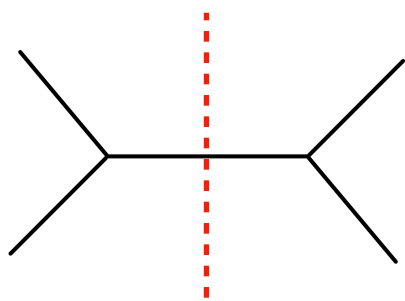


# Motivation

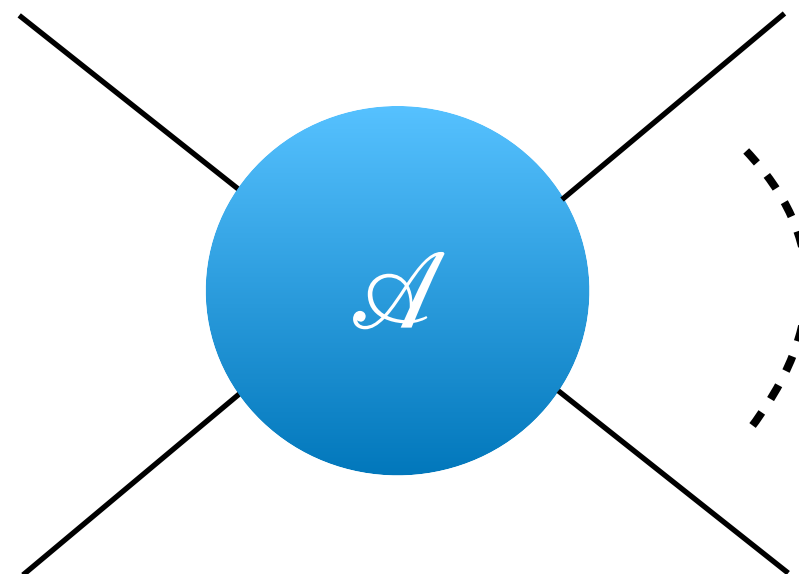
## Singularities

poles

*single*-particle  
states go on-shell



$$\sim \frac{1}{s}$$



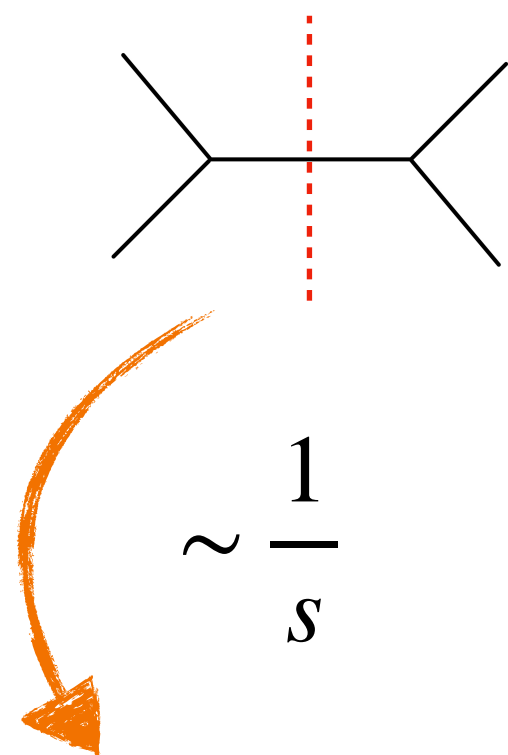


# Motivation

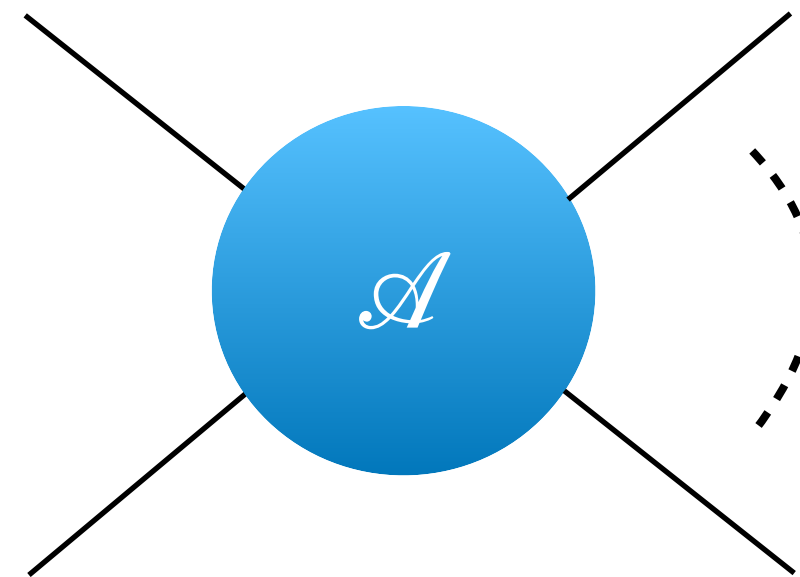
## Singularities

poles

*single-particle*  
states go on-shell



“cutting” the propagator ~  
the particle goes on shell



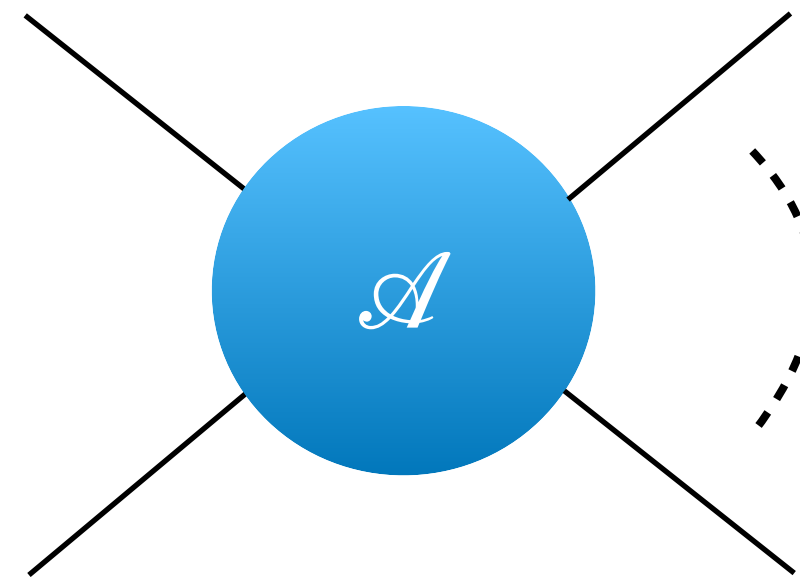
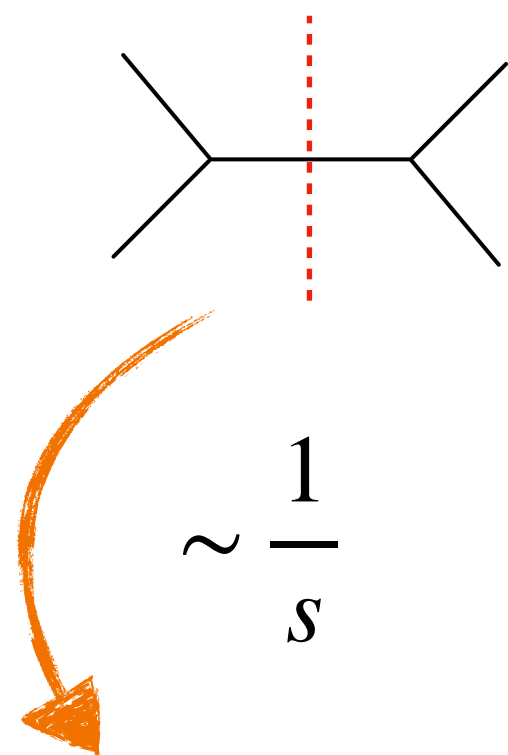
# Motivation

## Singularities

poles

branch-cuts

*single-particle*  
states go on-shell



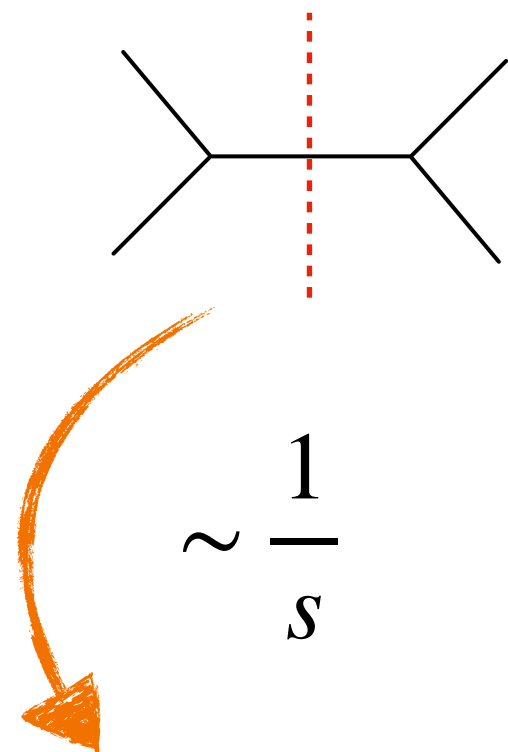
“cutting” the propagator  $\sim$   
the particle goes on shell

# Motivation

## Singularities

poles

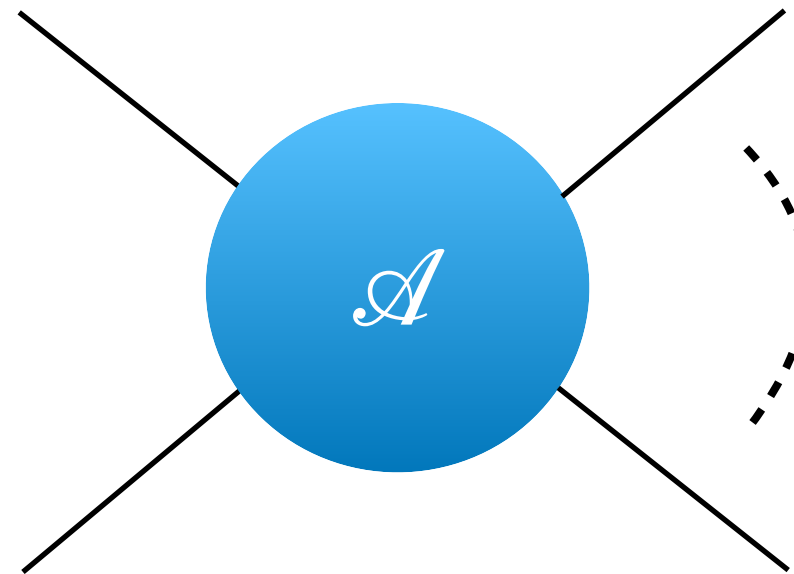
*single*-particle  
states go on-shell



$$\sim \frac{1}{s}$$

branch-cuts

*multi*-particle  
states go on-shell



“cutting” the propagator  $\sim$   
the particle goes on shell

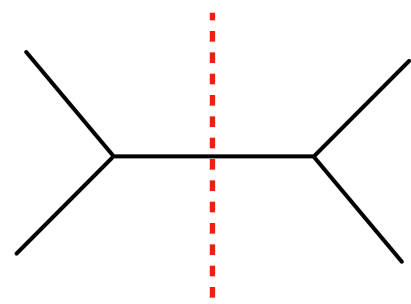


# Motivation

## Singularities

poles

*single*-particle  
states go on-shell

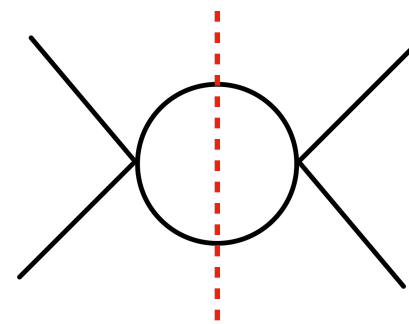


$\sim \frac{1}{s}$

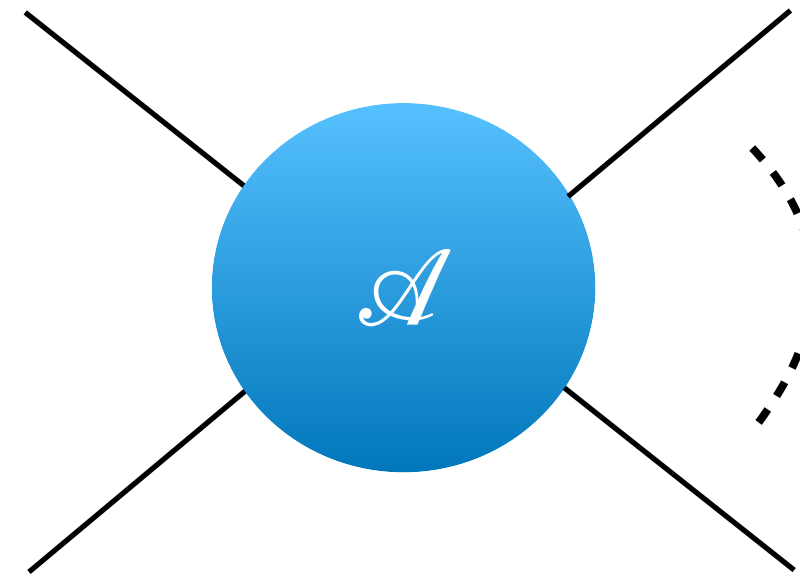
An orange curved arrow pointing from the pole singularity diagram towards the text "cutting the propagator".

branch-cuts

*multi*-particle  
states go on-shell



$\sim \sqrt{s}, \log(s)$



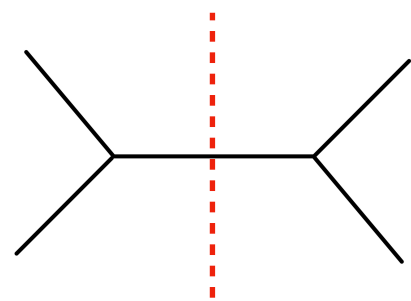
“cutting” the propagator ~  
the particle goes on shell

# Motivation

## Singularities

poles

*single*-particle  
states go on-shell

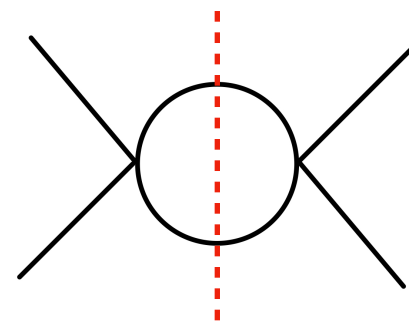


$\sim \frac{1}{s}$

An orange curved arrow pointing from the pole diagram towards the text "cutting the propagator".

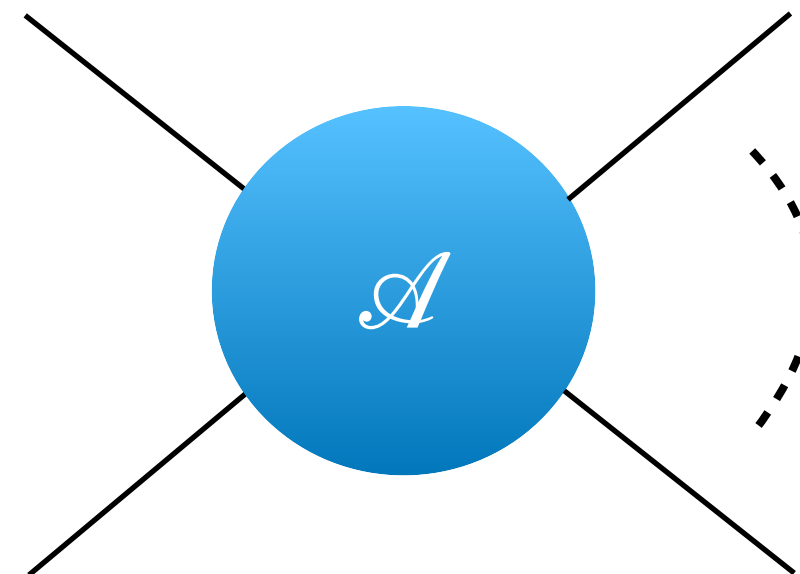
branch-cuts

*multi*-particle  
states go on-shell



$\sim \sqrt{s}, \log(s)$

## Computation



“cutting” the propagator ~  
the particle goes on shell



# Motivation

## Singularities

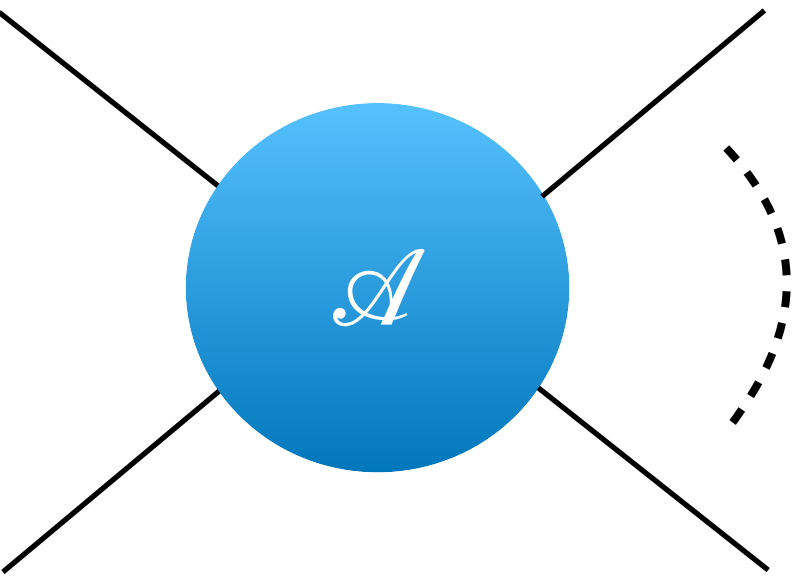
## Computation

poles

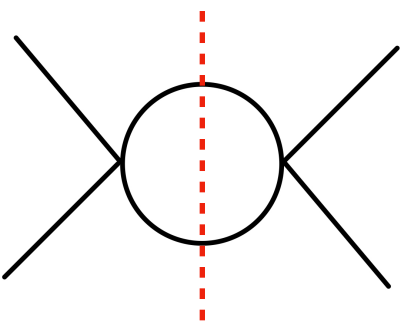
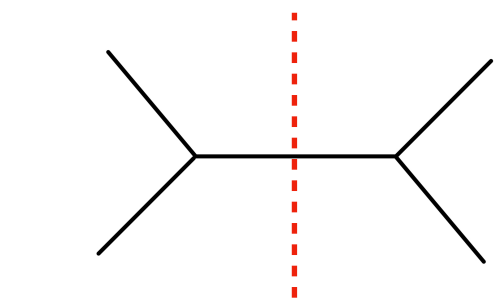
branch-cuts

*single*-particle  
states go on-shell

*multi*-particle  
states go on-shell



$\mathcal{A} \sim$



$\sim \frac{1}{s}$

$\sim \sqrt{s}, \log(s)$

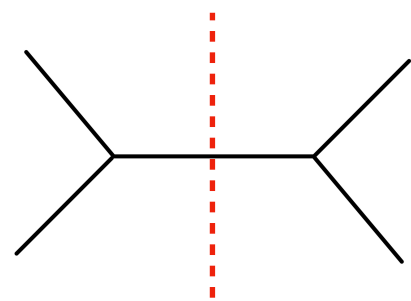
“cutting” the propagator  $\sim$   
the particle goes on shell

# Motivation

## Singularities

poles

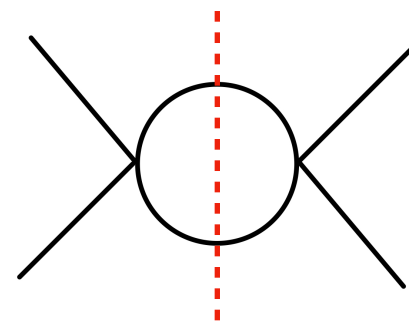
*single*-particle  
states go on-shell



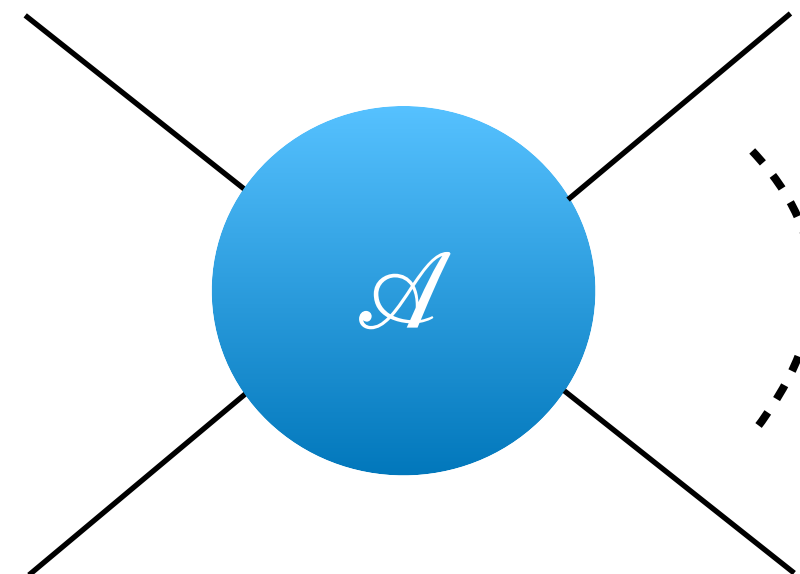
$$\sim \frac{1}{s}$$

branch-cuts

*multi*-particle  
states go on-shell



$$\sim \sqrt{s}, \log(s)$$



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

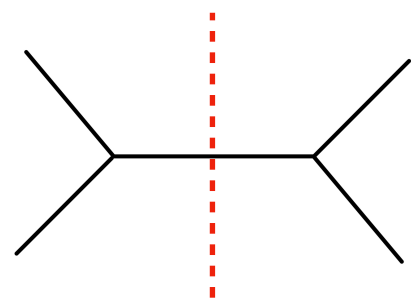
“cutting” the propagator ~  
the particle goes on shell

# Motivation

## Singularities

poles

*single-particle*  
states go on-shell

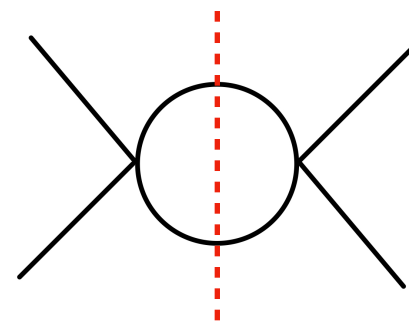


$\sim \frac{1}{s}$

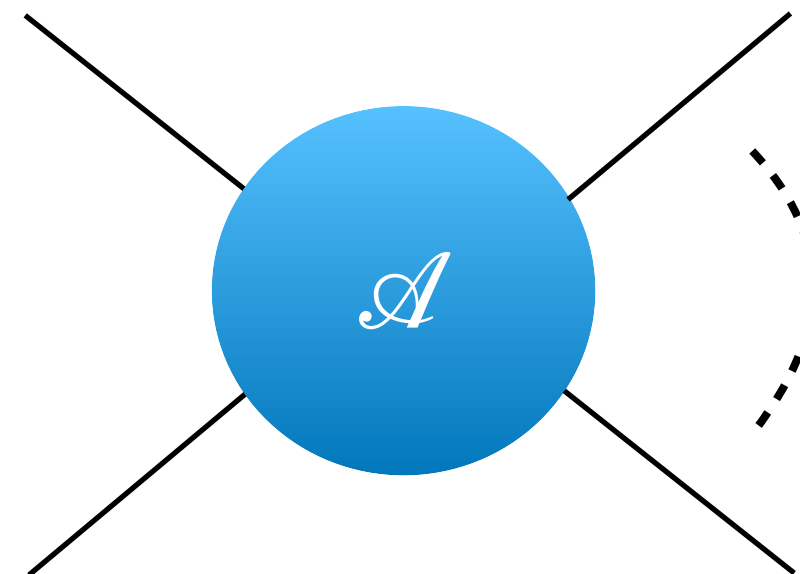
An orange curved arrow points from the propagator line in the diagram above to the text below.

branch-cuts

*multi-particle*  
states go on-shell



$\sim \sqrt{s}, \log(s)$



## Computation

$\mathcal{A} \sim T^{\mu\nu}$

tensor  
structure

“cutting” the propagator  $\sim$   
the particle goes on shell

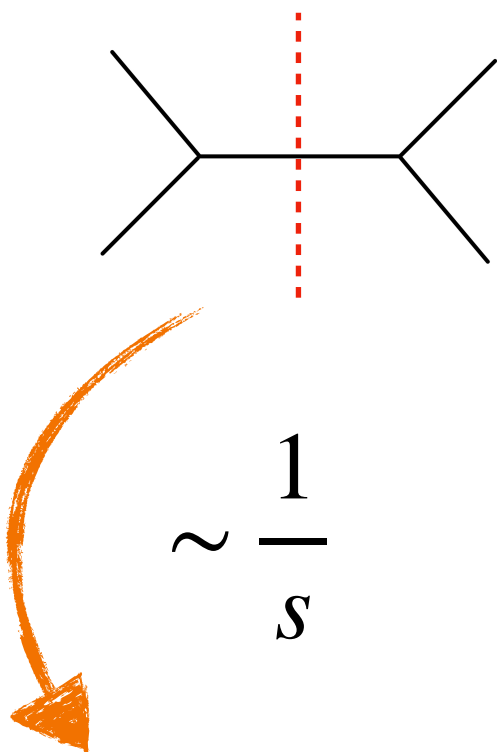


# Motivation

## Singularities

poles

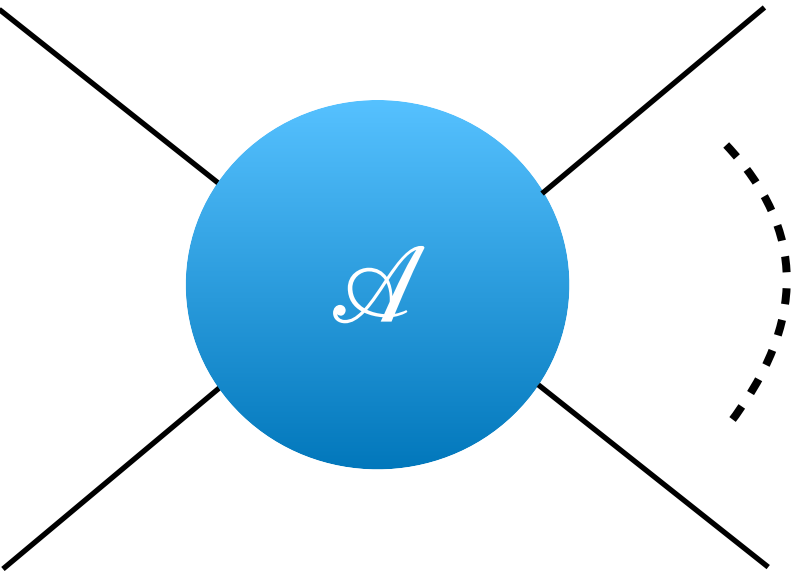
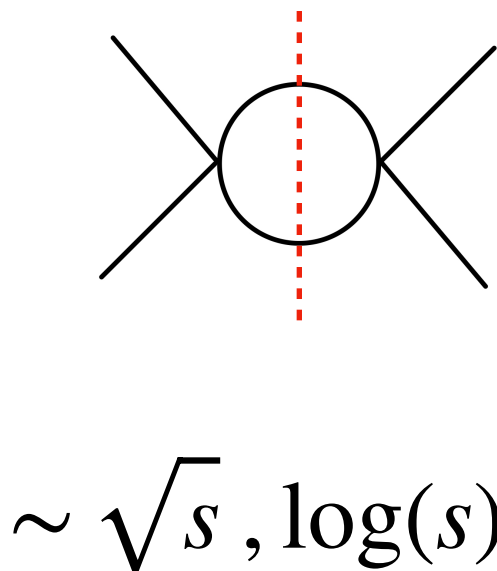
single-particle  
states go on-shell



“cutting” the propagator ~  
the particle goes on shell

branch-cuts

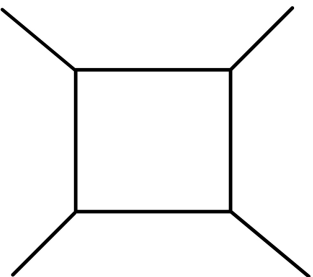
multi-particle  
states go on-shell



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure

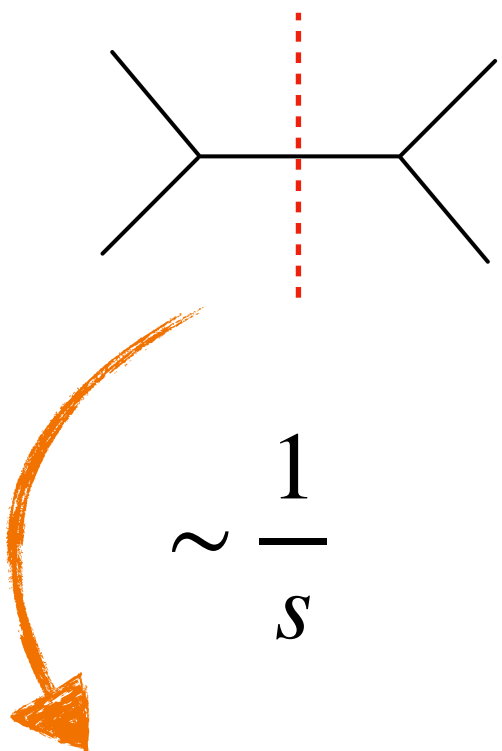


# Motivation

## Singularities

poles

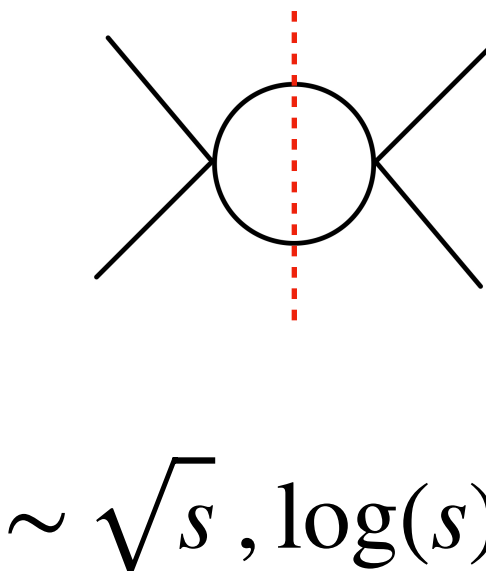
single-particle  
states go on-shell



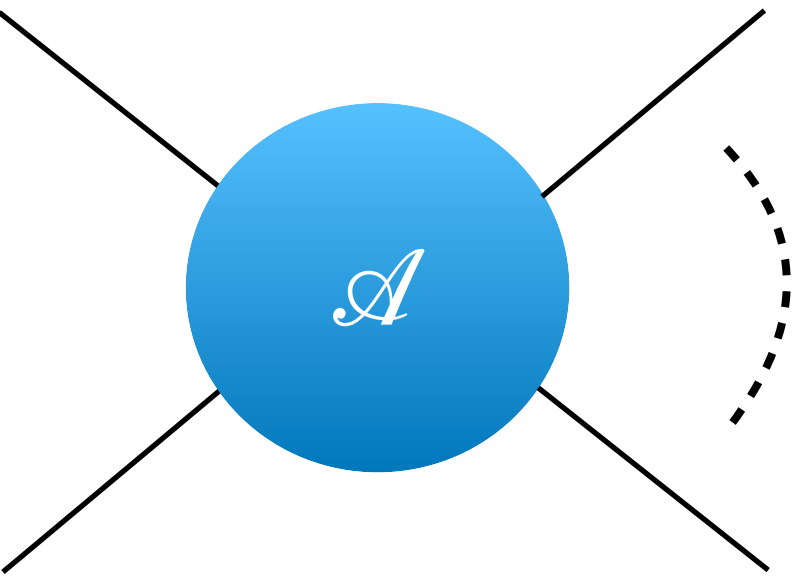
$$\sim \frac{1}{s}$$

branch-cuts

multi-particle  
states go on-shell



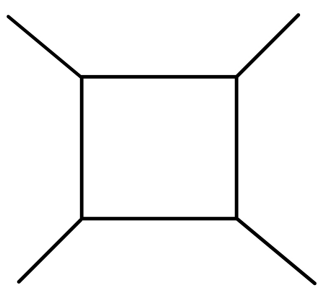
$$\sim \sqrt{s}, \log(s)$$



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

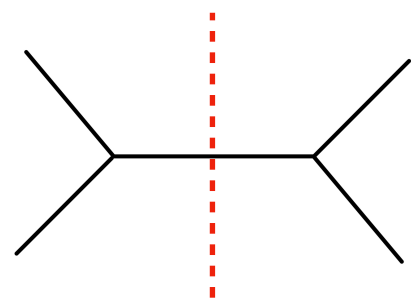
“cutting” the propagator  $\sim$   
the particle goes on shell

# Motivation

## Singularities

poles

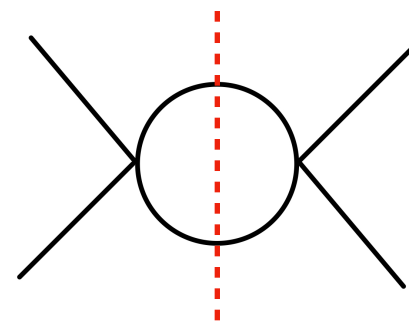
single-particle  
states go on-shell



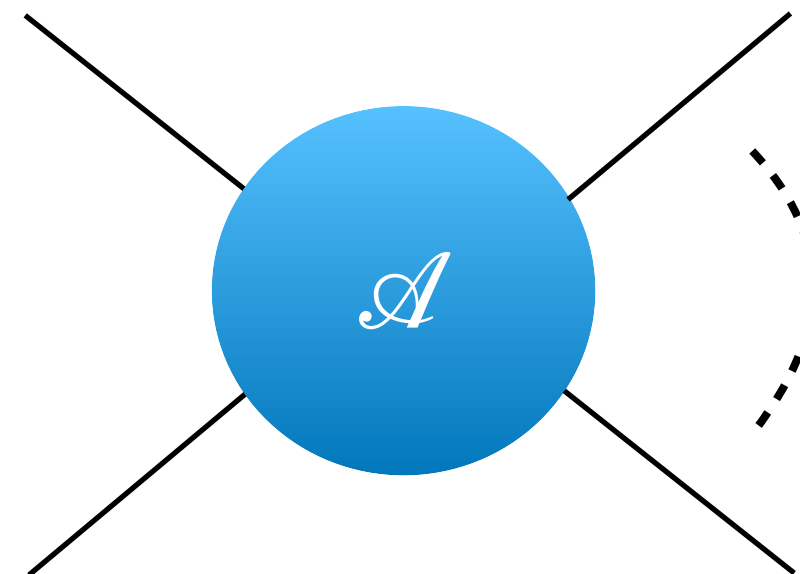
$$\sim \frac{1}{s}$$

branch-cuts

multi-particle  
states go on-shell



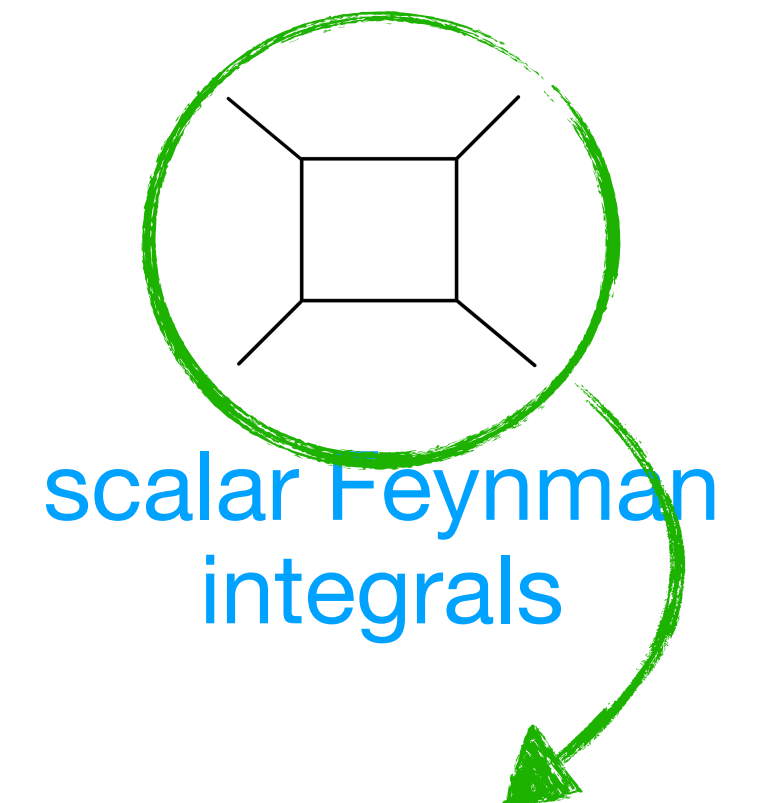
$$\sim \sqrt{s}, \log(s)$$



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

“cutting” the propagator  $\sim$   
the particle goes on shell

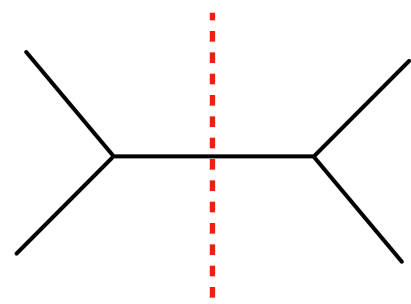


# Motivation

## Singularities

poles

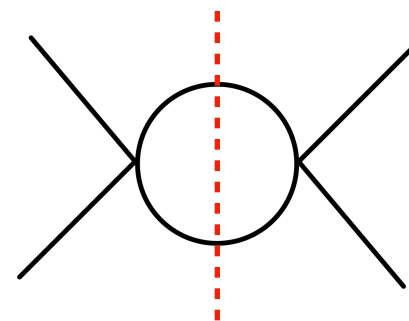
single-particle  
states go on-shell



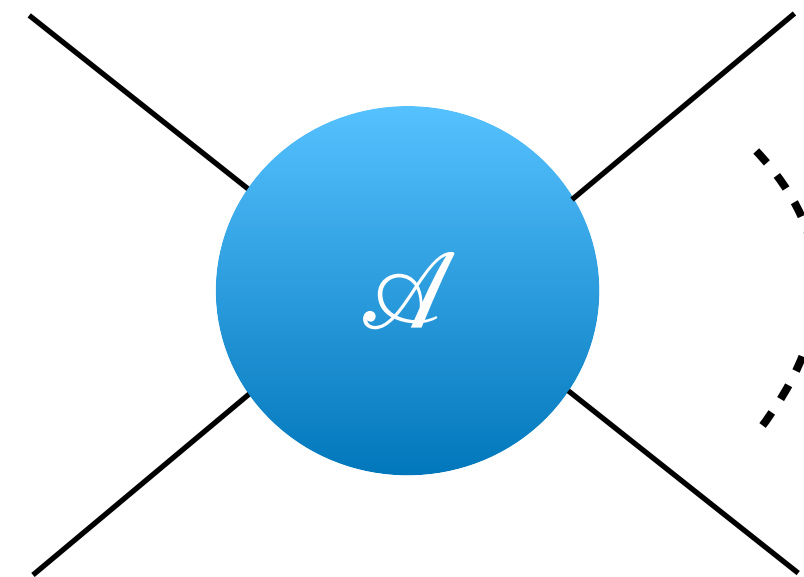
$$\sim \frac{1}{s}$$

branch-cuts

multi-particle  
states go on-shell



$$\sim \sqrt{s}, \log(s)$$



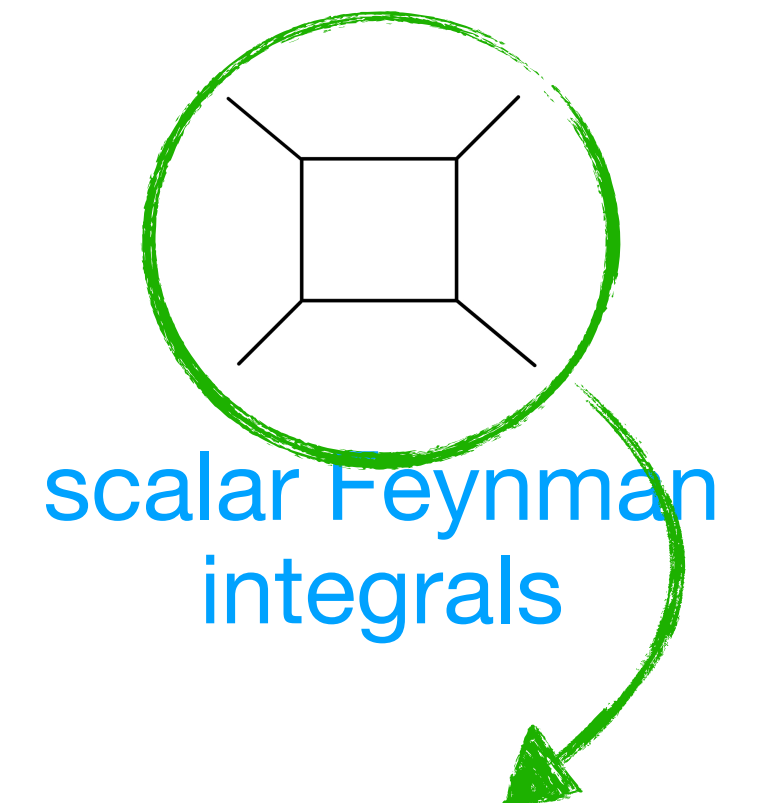
one-loop



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

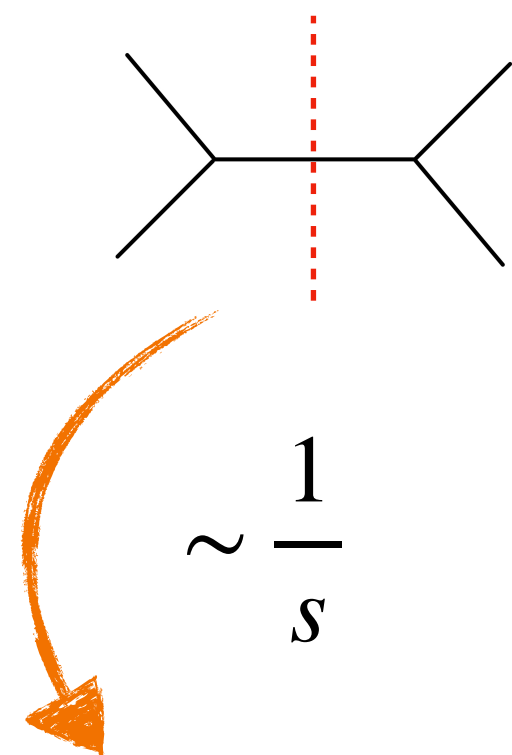
“cutting” the propagator  $\sim$   
the particle goes on shell

# Motivation

## Singularities

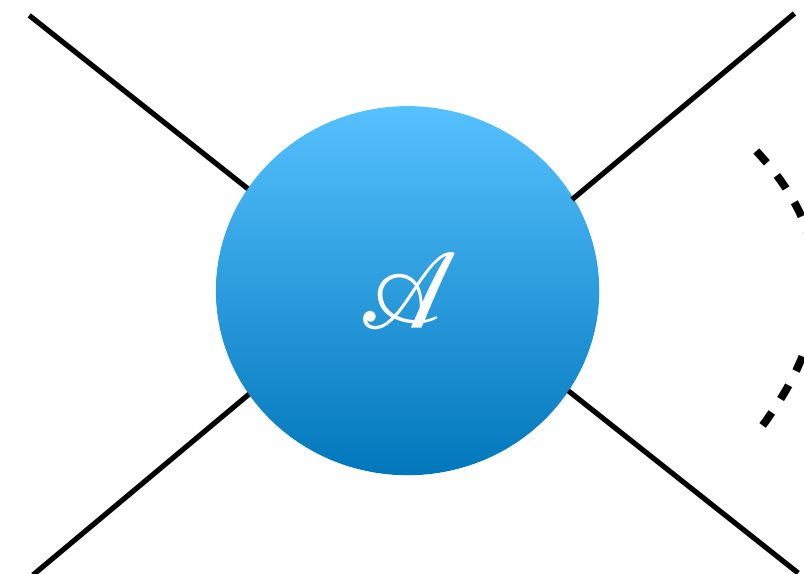
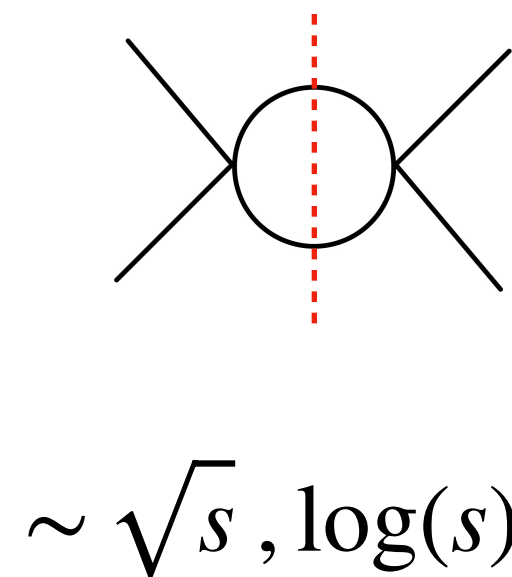
poles

single-particle  
states go on-shell



branch-cuts

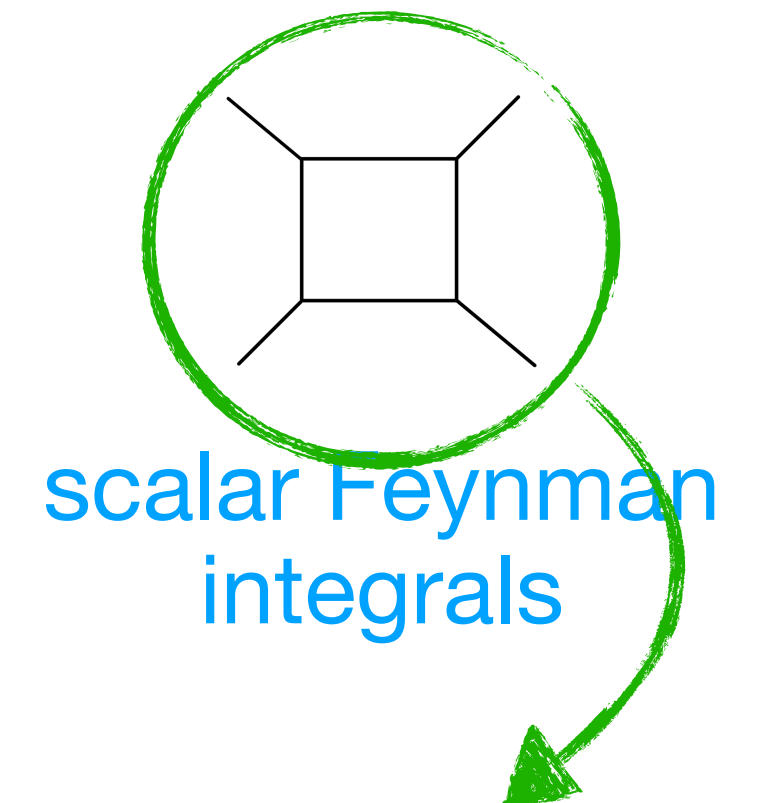
multi-particle  
states go on-shell



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



one-loop

$$\sum_i R_i(s_{ij}) \int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

“cutting” the propagator  $\sim$   
the particle goes on shell

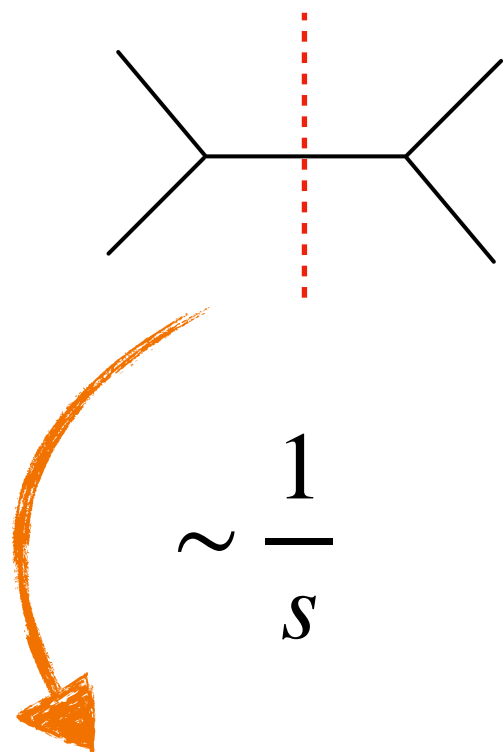


# Motivation

## Singularities

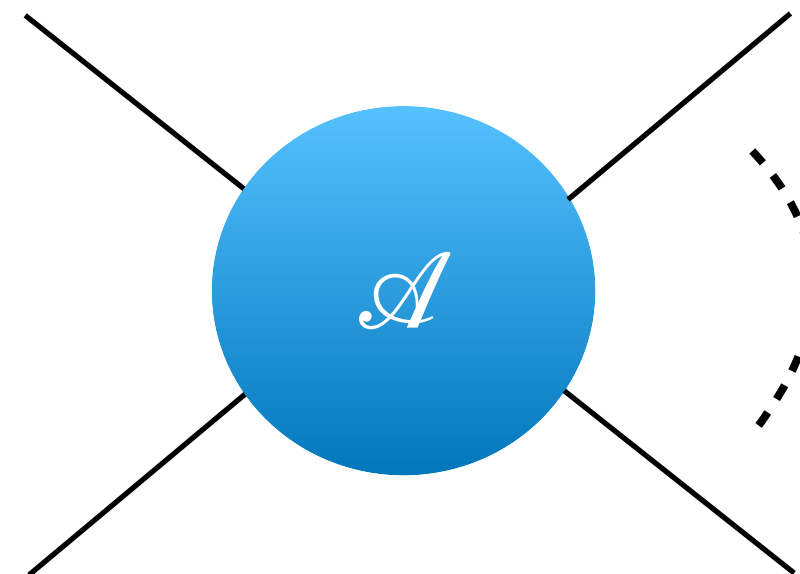
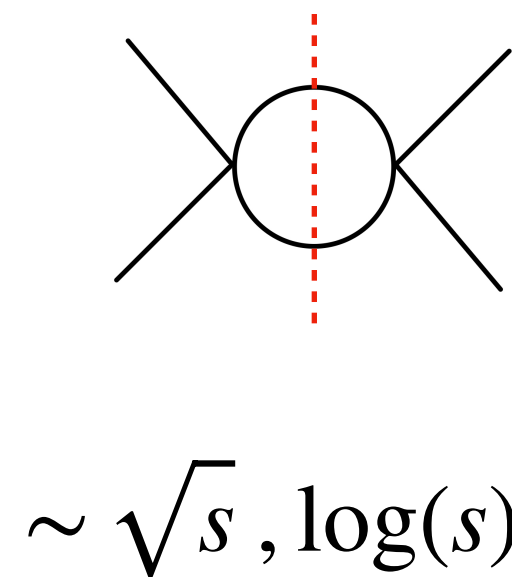
poles

single-particle  
states go on-shell



branch-cuts

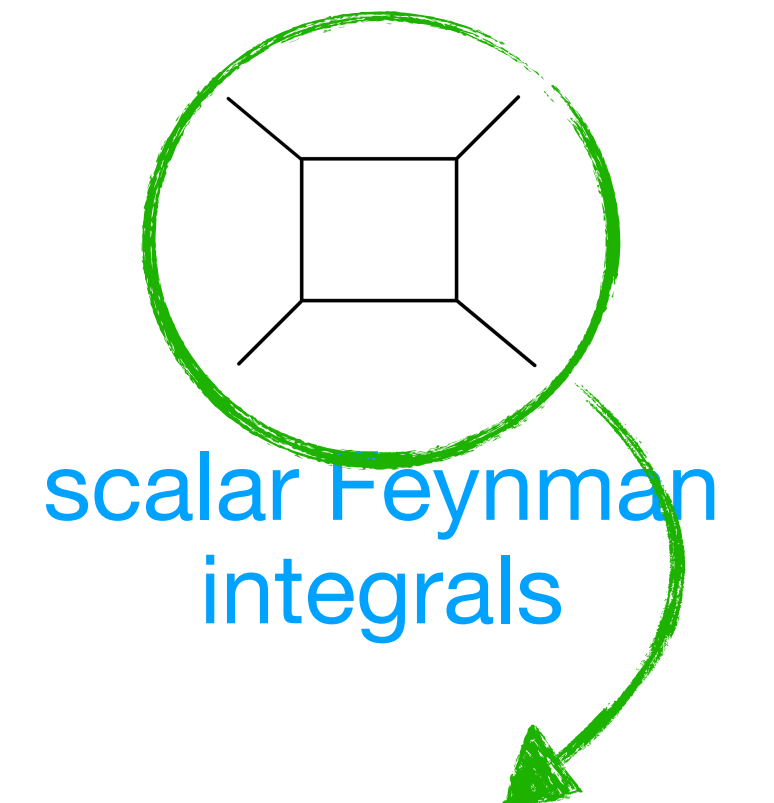
multi-particle  
states go on-shell



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



one-loop

$$\sum_i R_i(s_{ij}) \int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

algebraic  
functions

poles

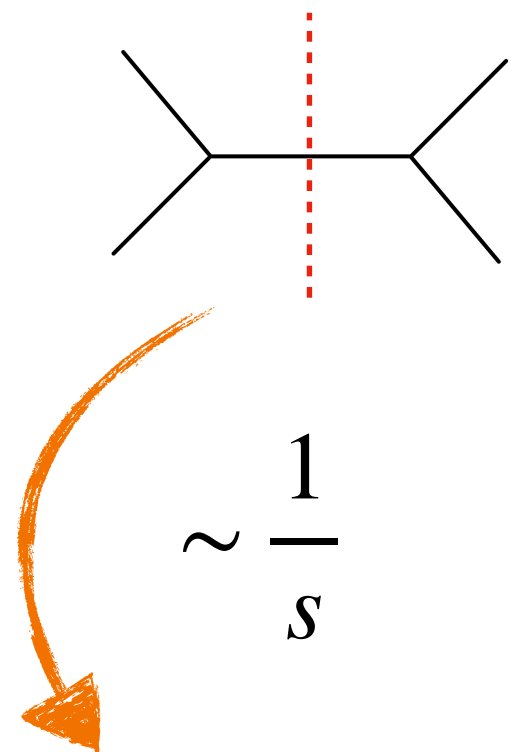
“cutting” the propagator  $\sim$   
the particle goes on shell

# Motivation

## Singularities

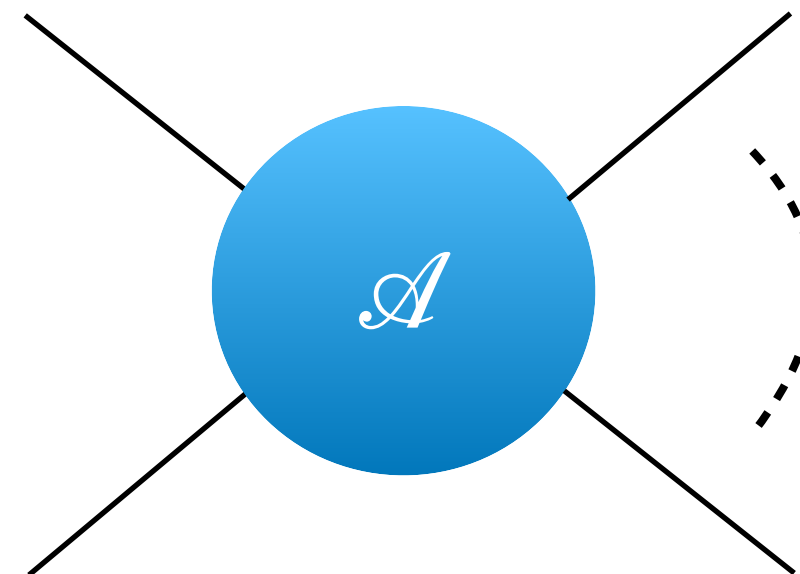
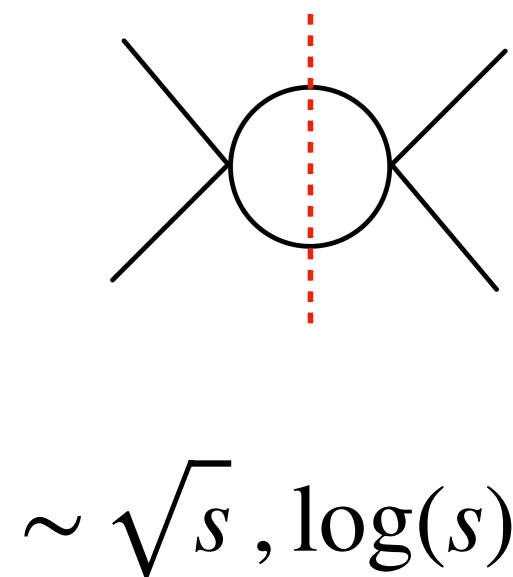
poles

single-particle  
states go on-shell



branch-cuts

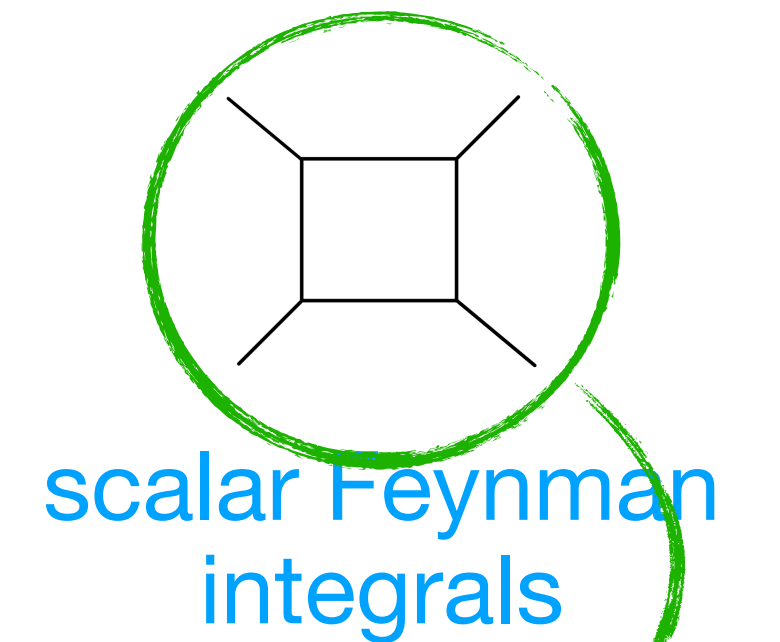
multi-particle  
states go on-shell



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



one-loop



$$\sum_i R_i(s_{ij})$$

algebraic  
functions

poles

$$\int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

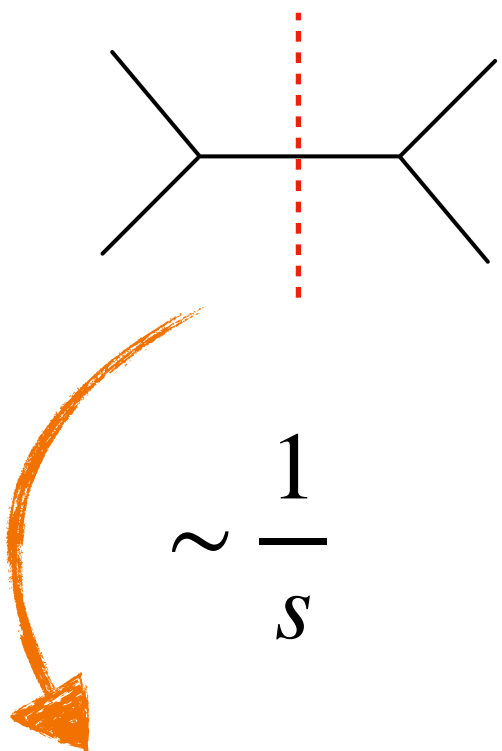
“cutting” the propagator  $\sim$   
the particle goes on shell

# Motivation

## Singularities

poles

single-particle  
states go on-shell

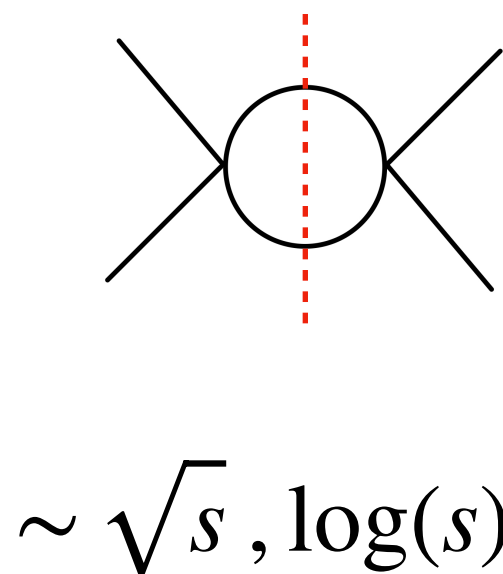


$$\sim \frac{1}{s}$$

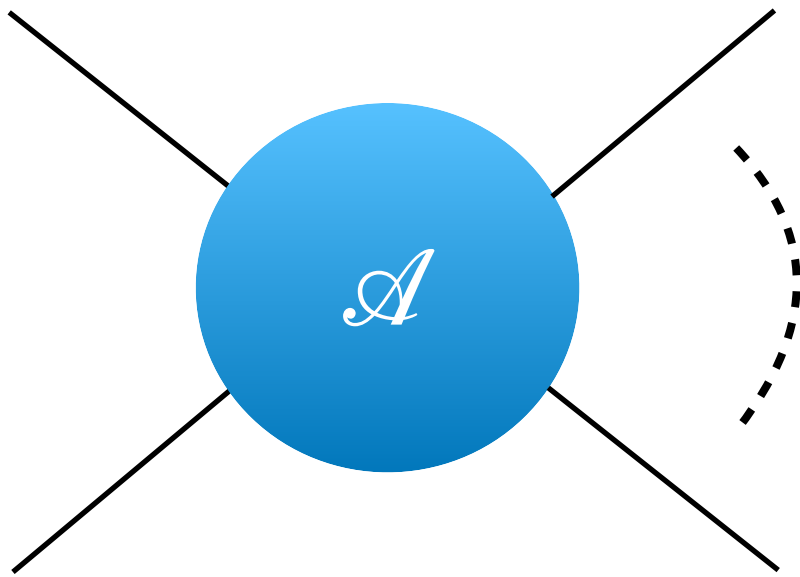
“cutting” the propagator ~  
the particle goes on shell

branch-cuts

multi-particle  
states go on-shell



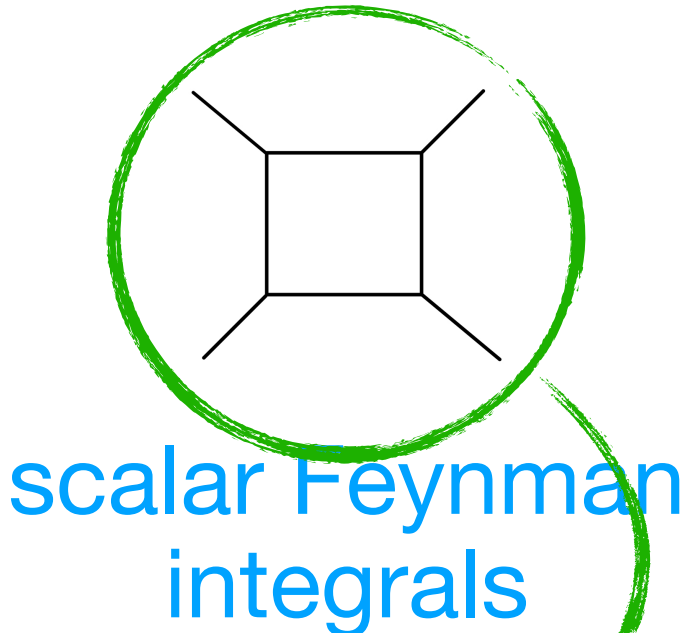
$$\sim \sqrt{s}, \log(s)$$



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

one-loop



$$\sum_i R_i(s_{ij})$$

algebraic  
functions

poles

$$\int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

special functions

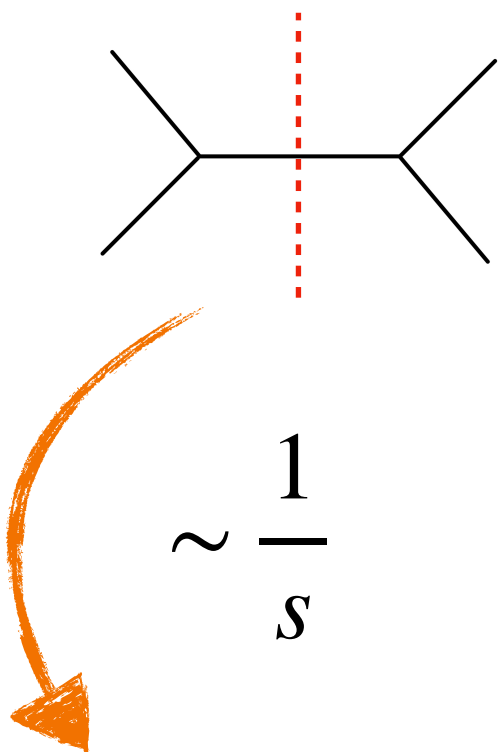


# Motivation

## Singularities

poles

single-particle  
states go on-shell

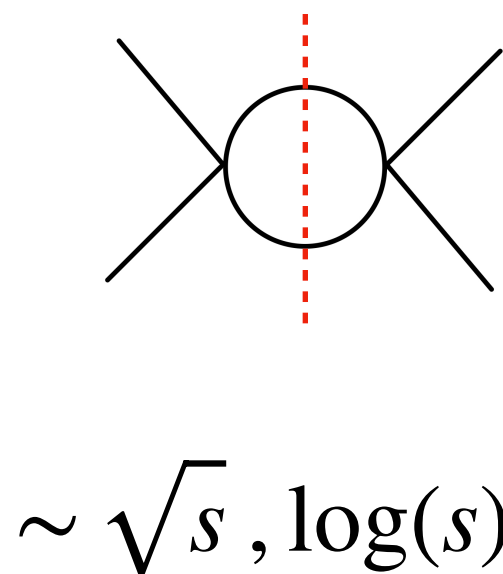


$$\sim \frac{1}{s}$$

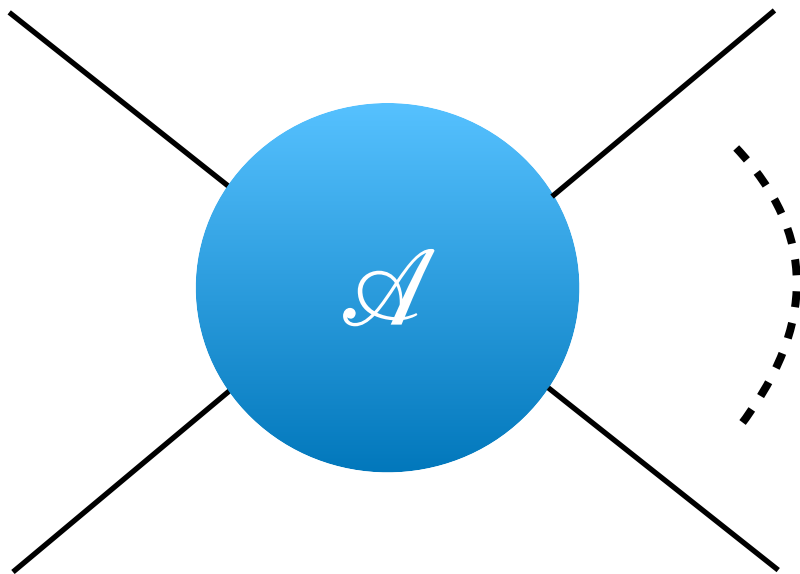
“cutting” the propagator ~  
the particle goes on shell

branch-cuts

multi-particle  
states go on-shell



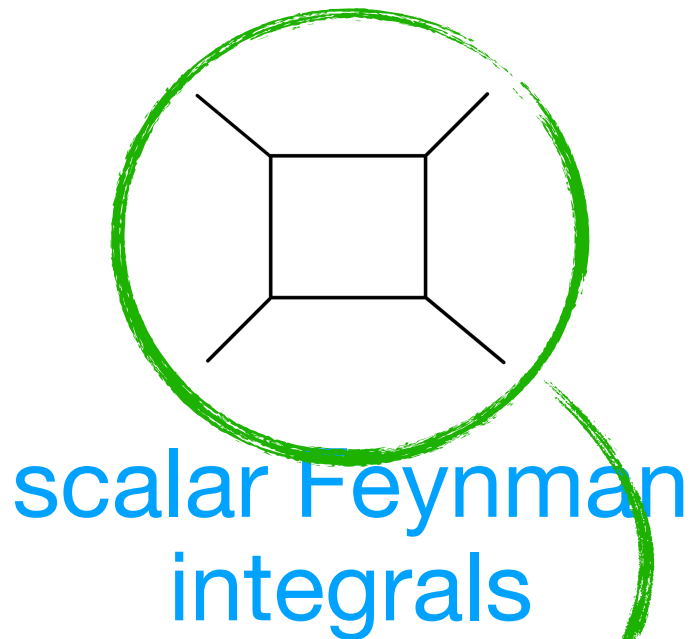
$$\sim \sqrt{s}, \log(s)$$



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

one-loop



$$\sum_i R_i(s_{ij})$$

algebraic  
functions

poles

$$\int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

special functions

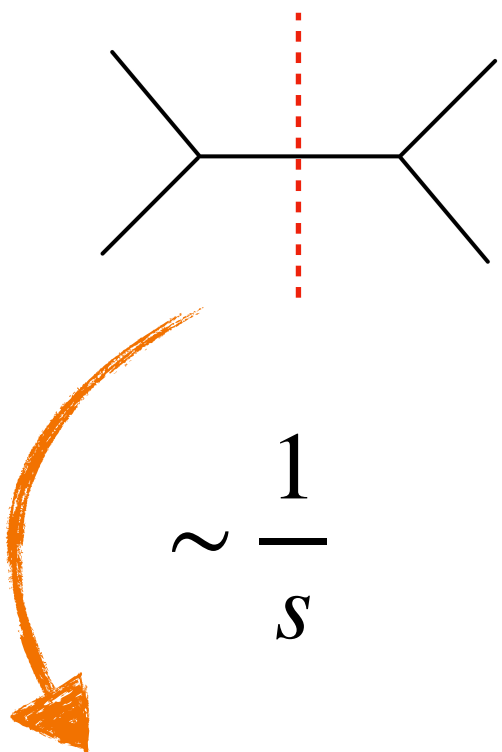
transcendental:  
from the integration

# Motivation

## Singularities

poles

single-particle  
states go on-shell

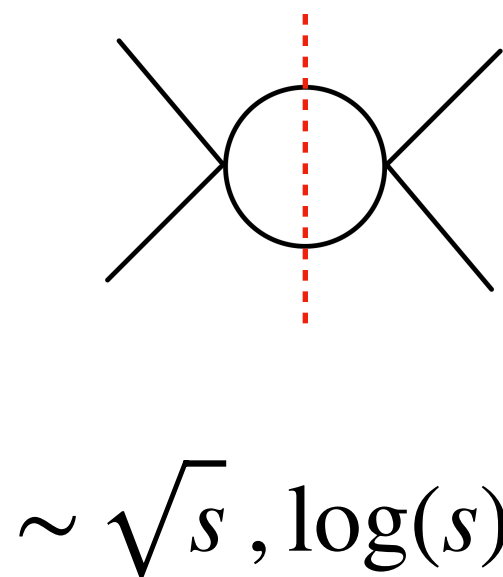


$$\sim \frac{1}{s}$$

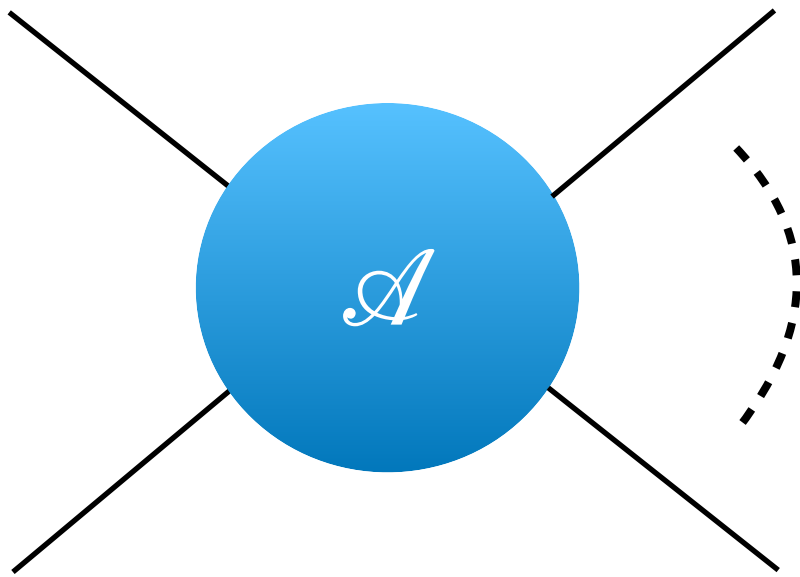
“cutting” the propagator ~  
the particle goes on shell

branch-cuts

multi-particle  
states go on-shell



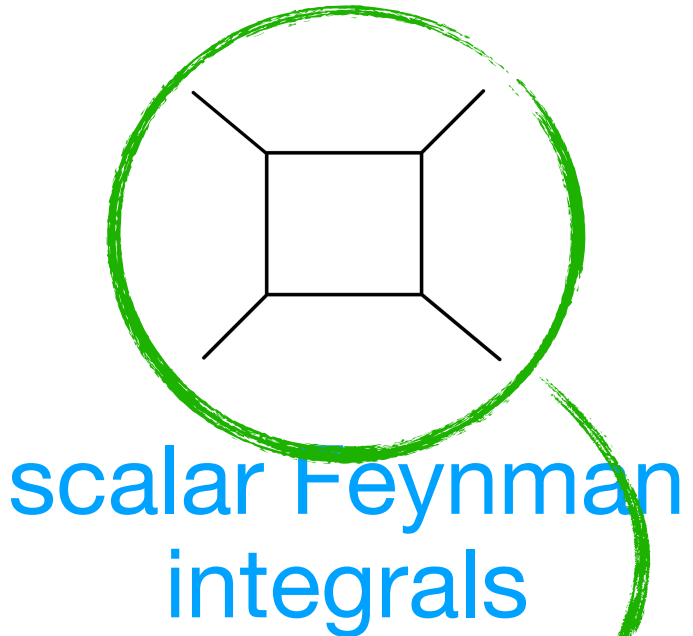
$$\sim \sqrt{s}, \log(s)$$



## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

one-loop

$$\sum_i R_i(s_{ij}) \int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

algebraic  
functions

poles

special functions

transcendental:  
from the integration

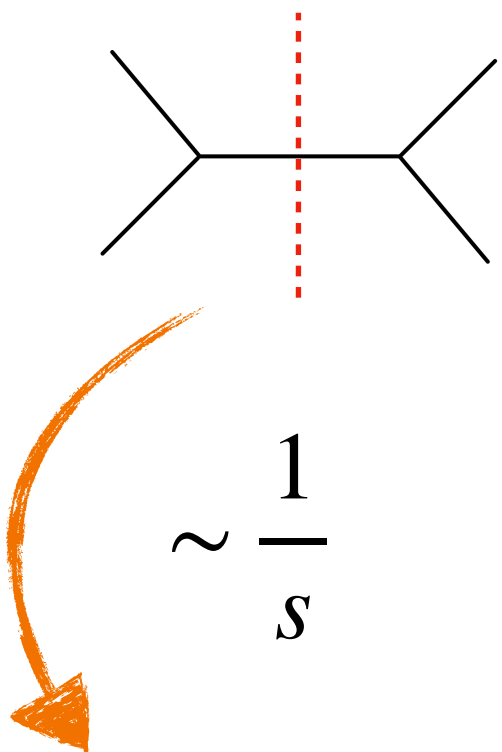
evident with *pure* integrals  
[Arkani et al., '10]

# Motivation

## Singularities

poles

single-particle  
states go on-shell

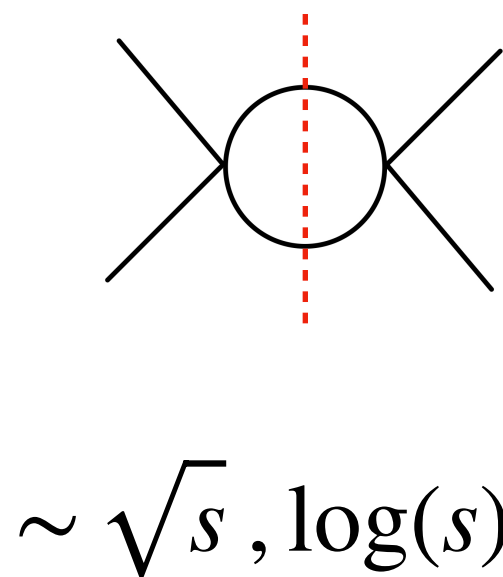


$$\sim \frac{1}{s}$$

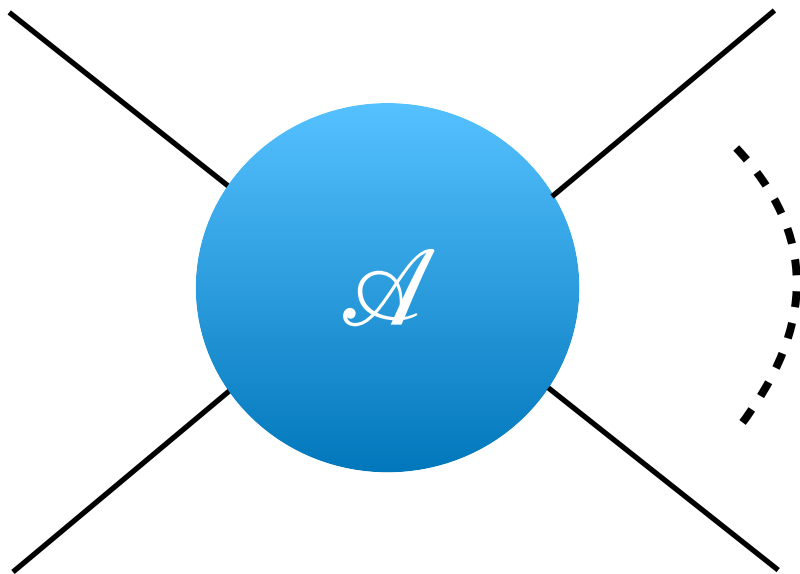
“cutting” the propagator ~  
the particle goes on shell

branch-cuts

multi-particle  
states go on-shell



$$\sim \sqrt{s}, \log(s)$$

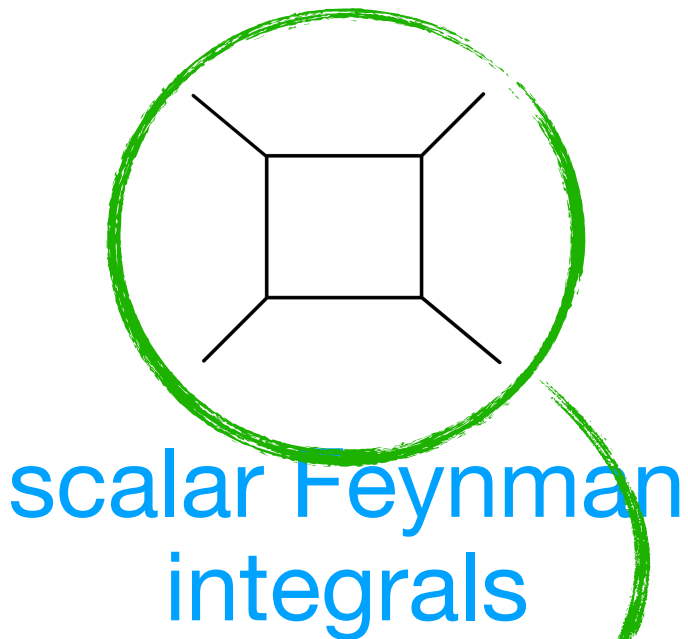


What stays beyond?

## Computation

$$\mathcal{A} \sim T^{\mu\nu}$$

tensor  
structure



scalar Feynman  
integrals

one-loop



$$\sum_i R_i(s_{ij})$$

algebraic  
functions

poles

$$\int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

special functions

transcendental:  
from the integration

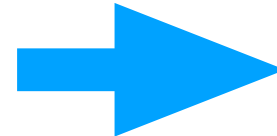
evident with *pure* integrals  
[Arkani et al., '10]



# First of all... when are there only (poly)logs?

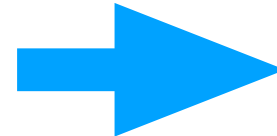


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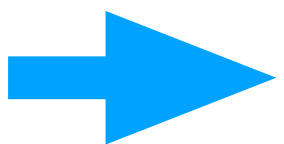


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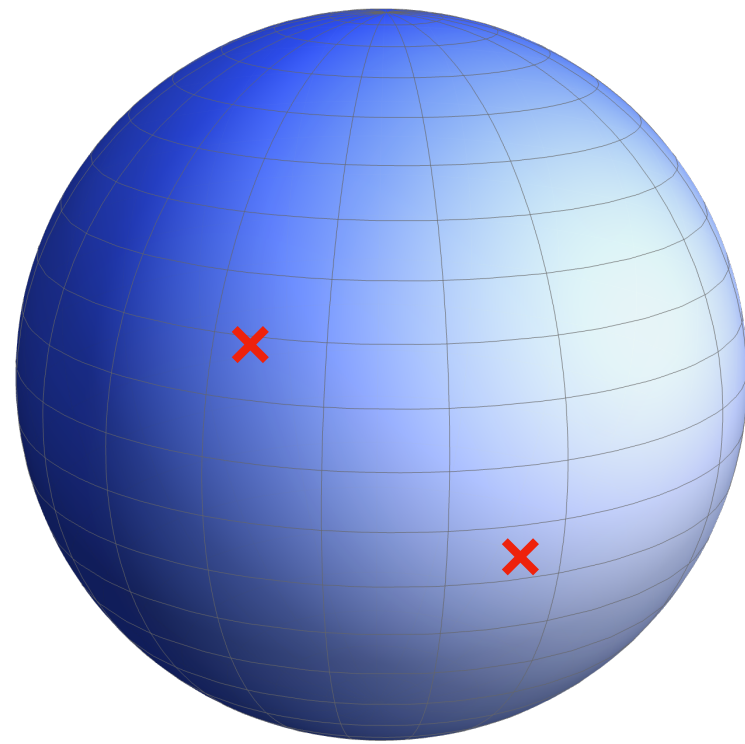
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The *integrand* lives on the Riemann sphere with some marked points

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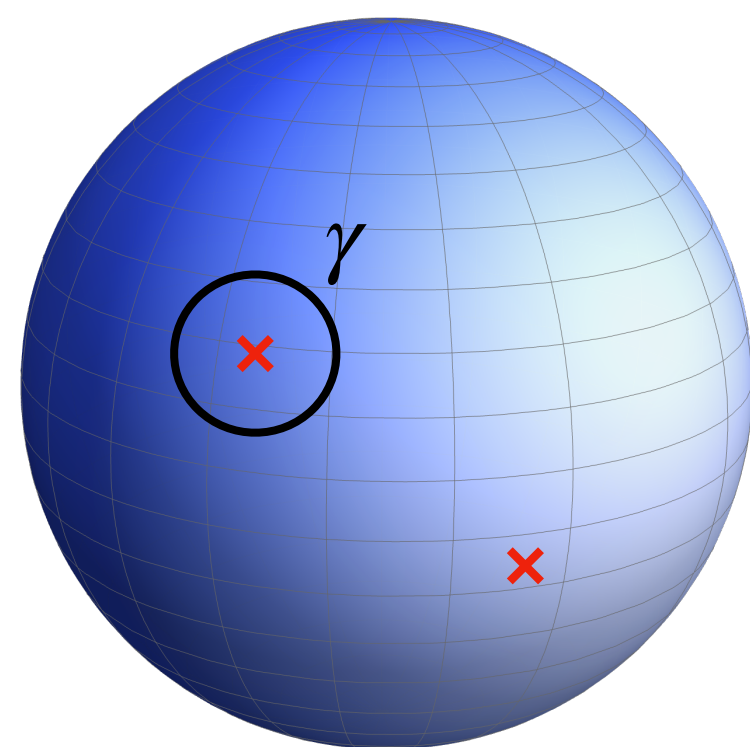
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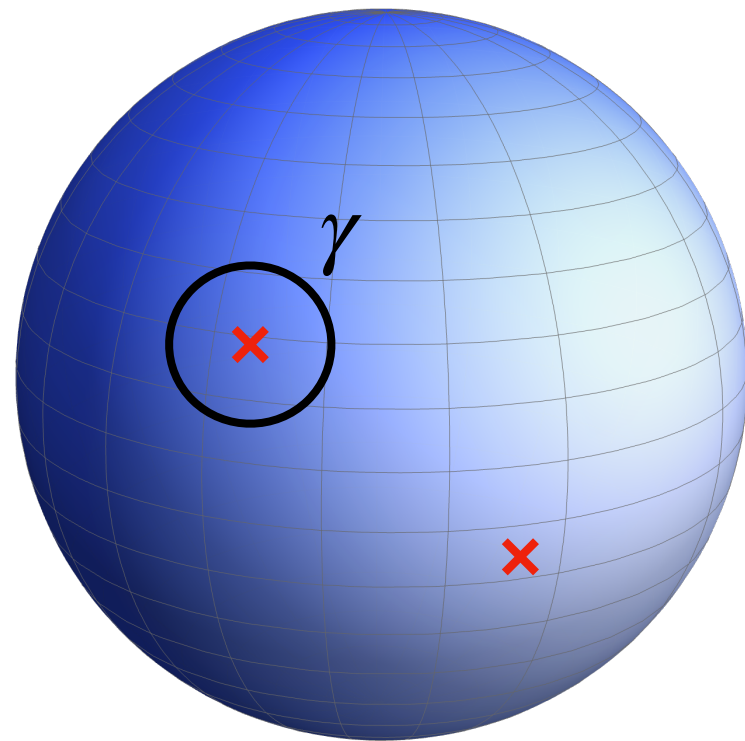
Independent forms:  $\int \frac{dt}{t-a}$



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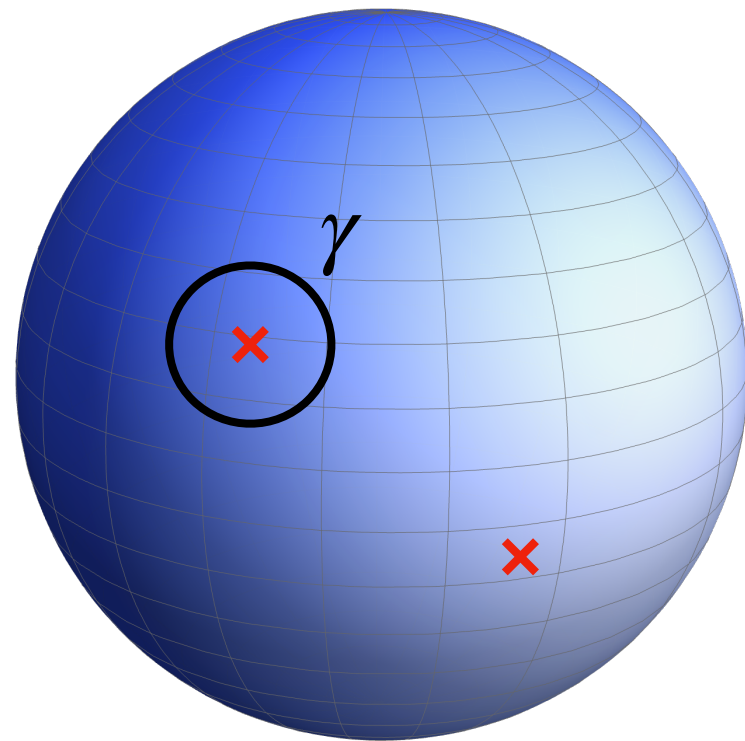


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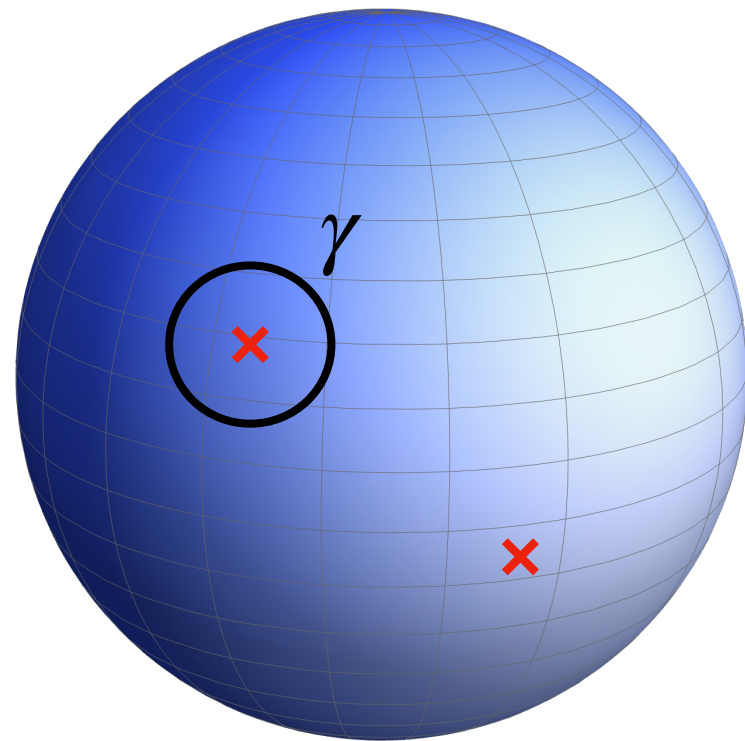
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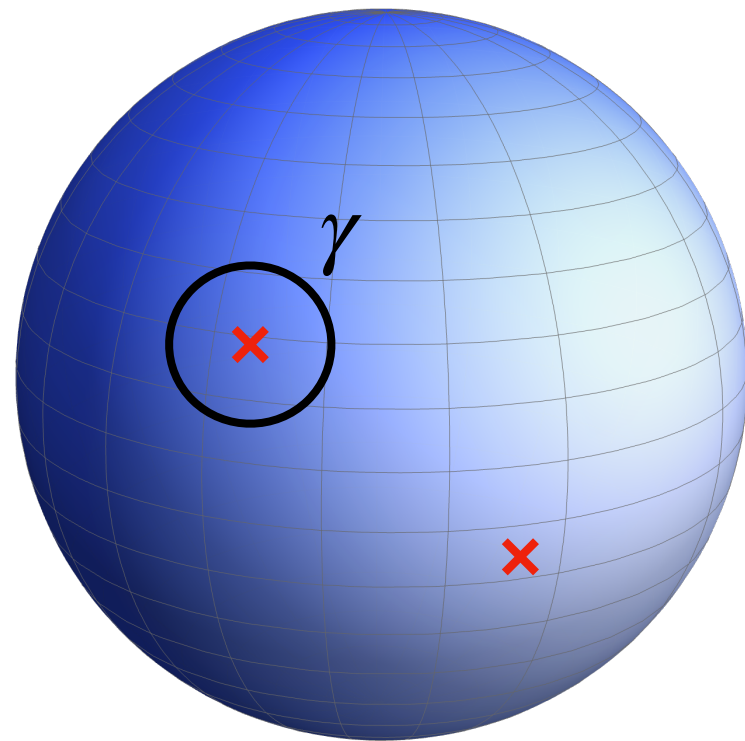
dim-reg  
parameter



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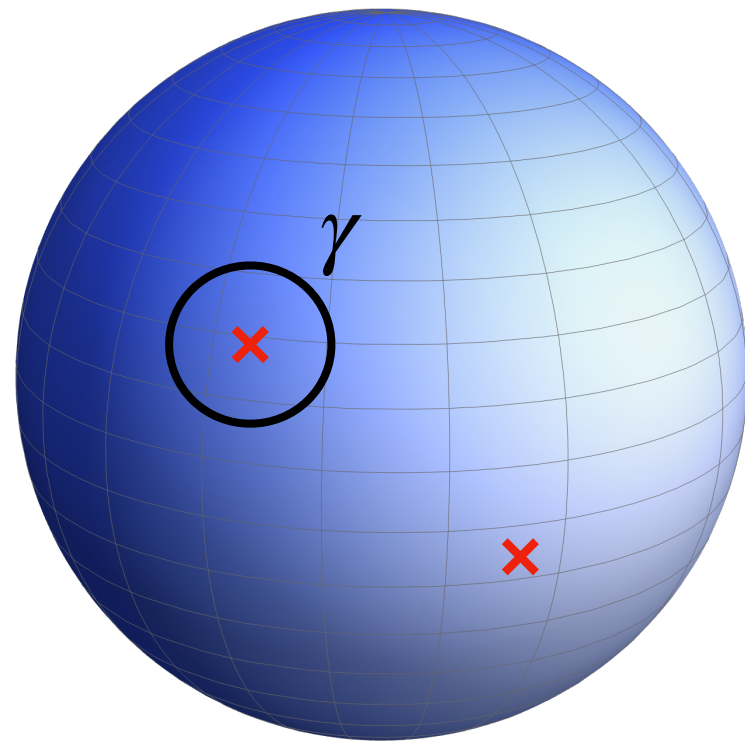
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[Henn, '13]

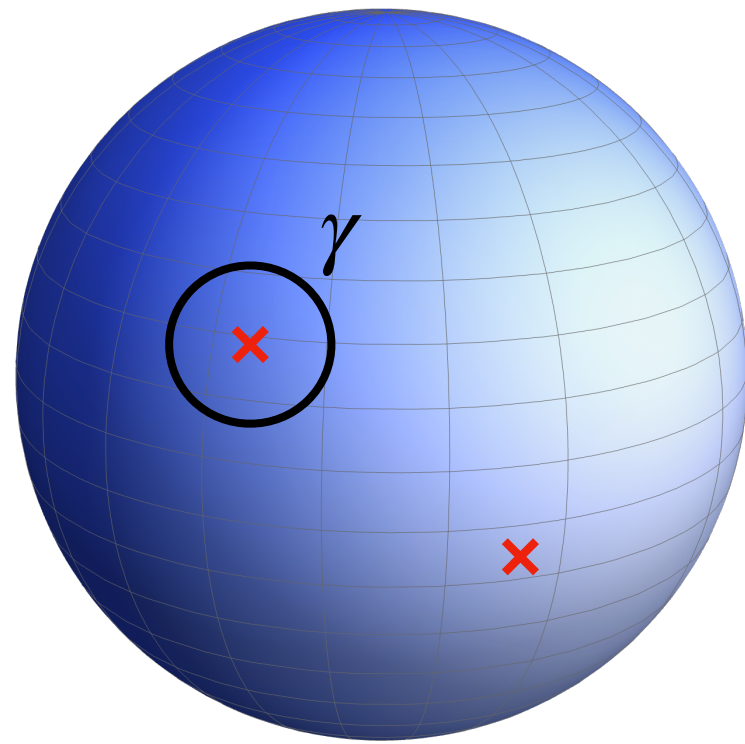




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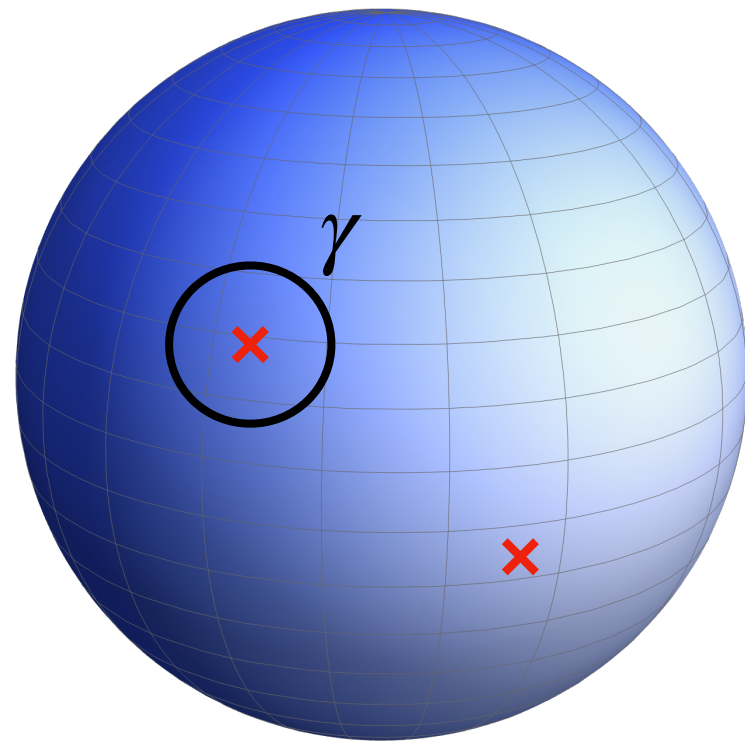
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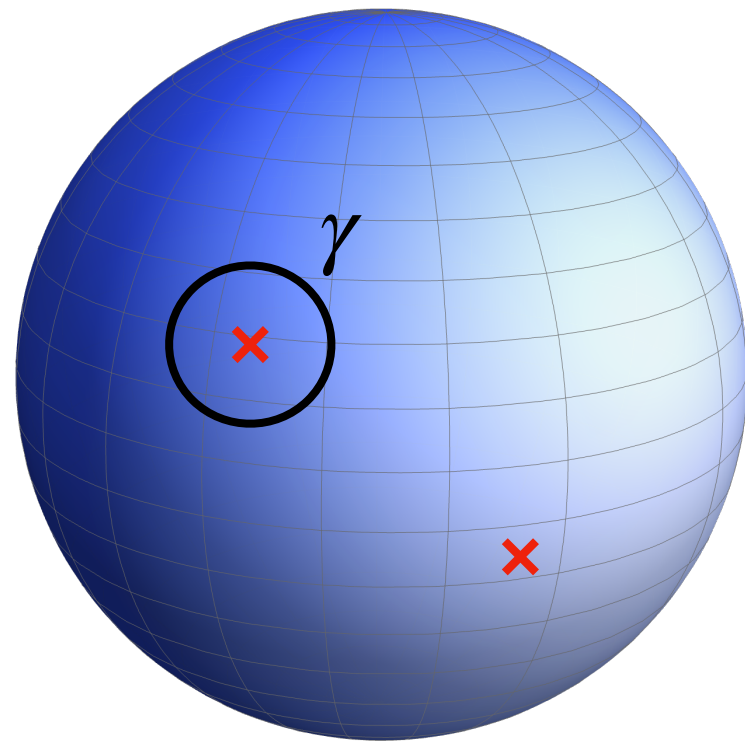
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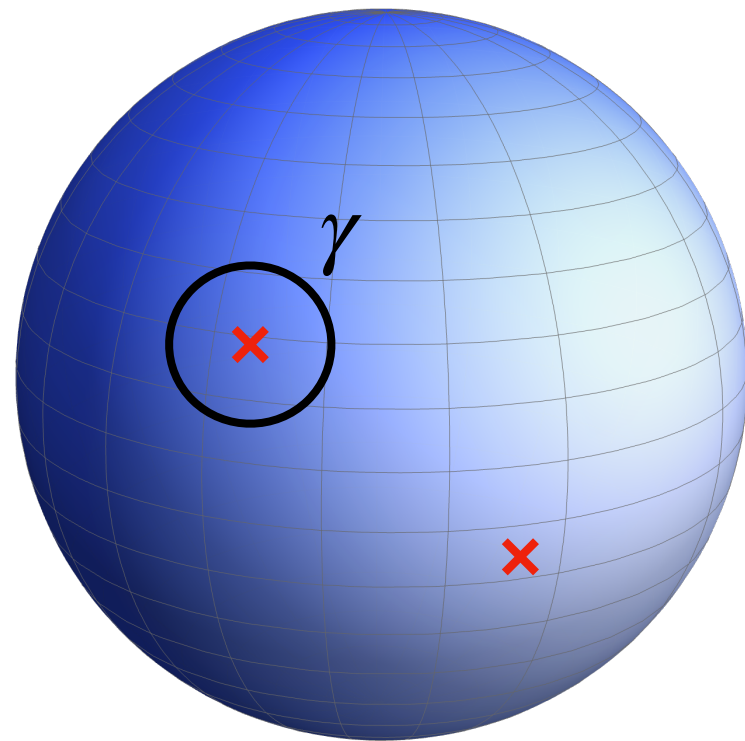
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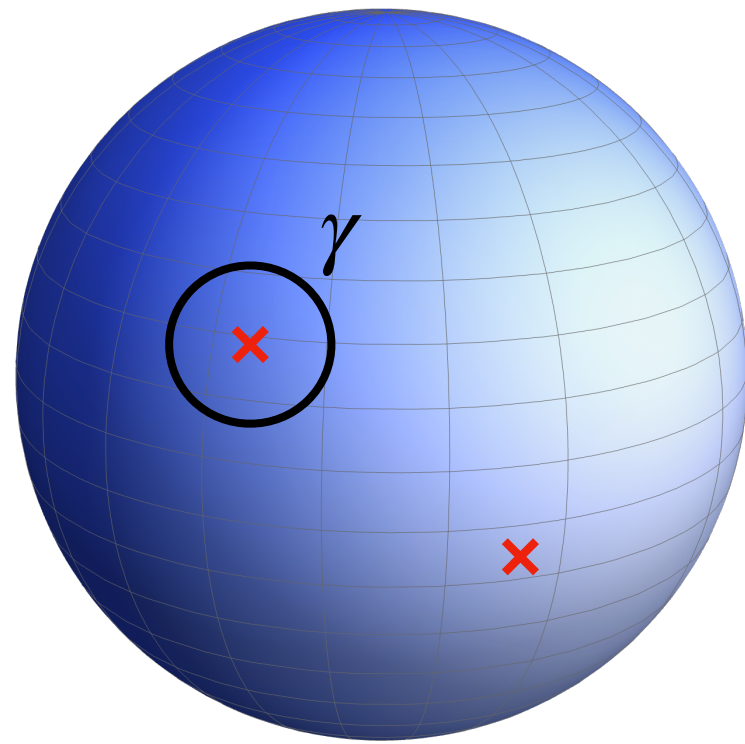
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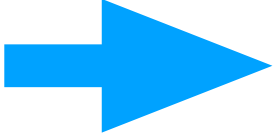
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uniform transcendental weight  
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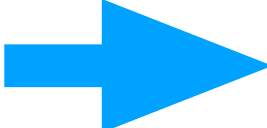
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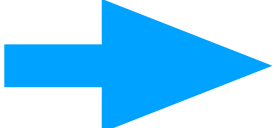
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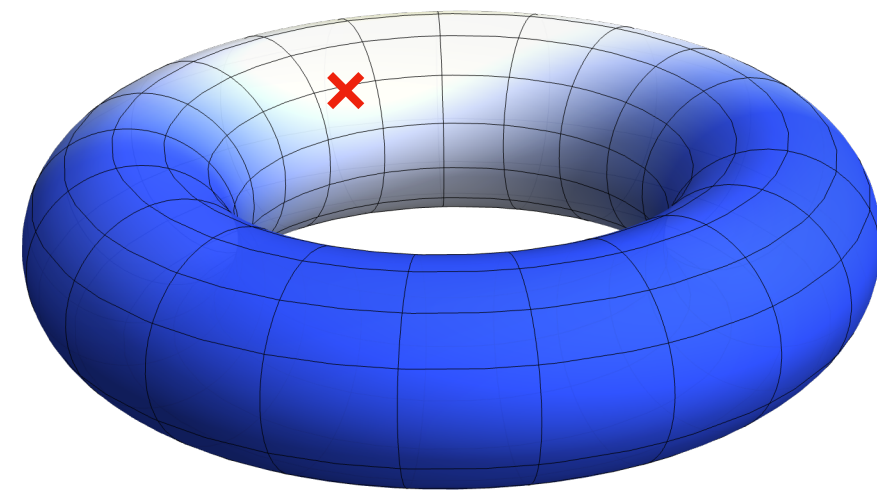
simplest generalisation: **elliptic curves**



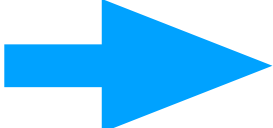
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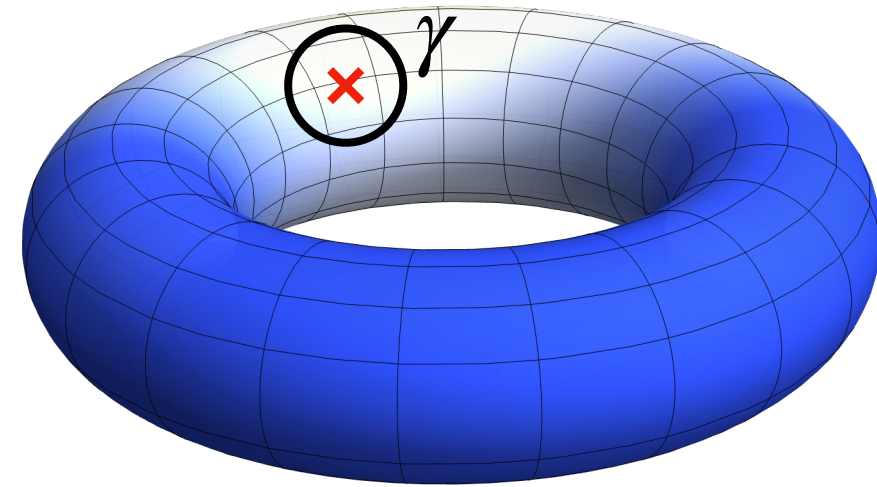
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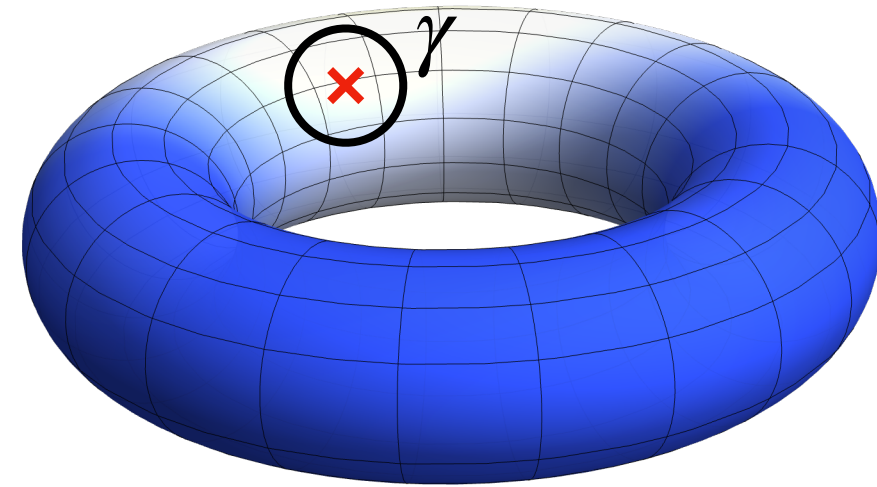
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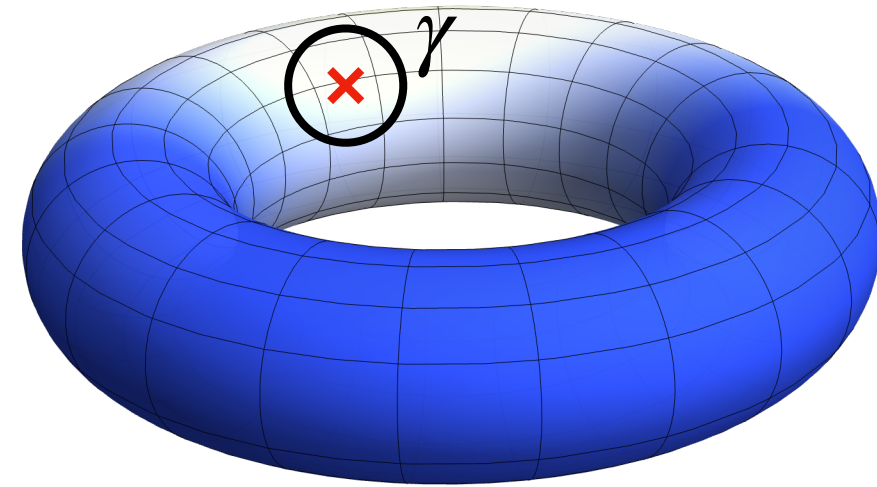


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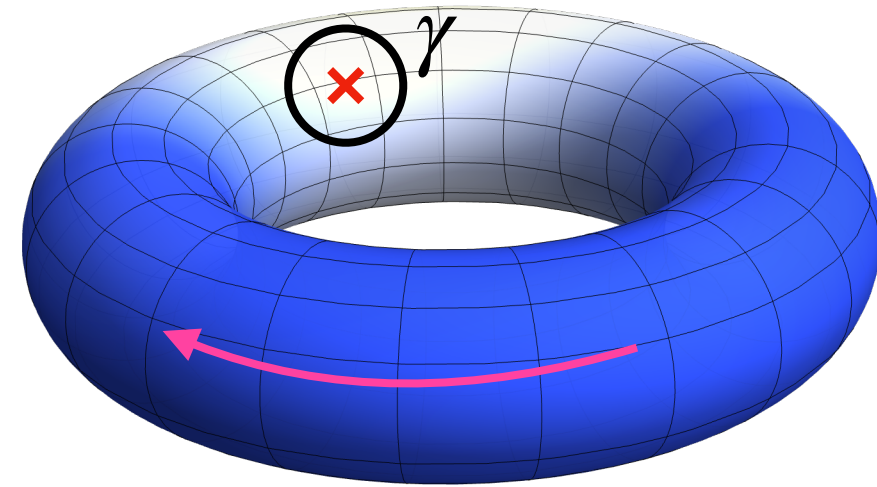
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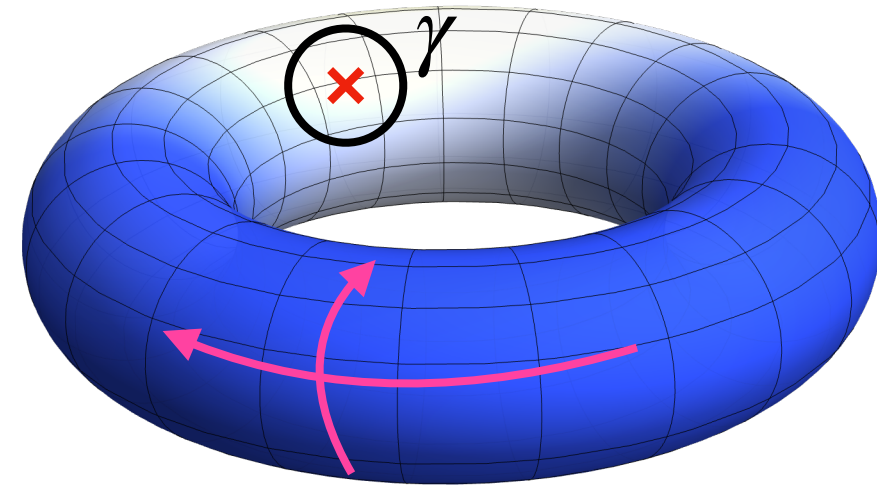
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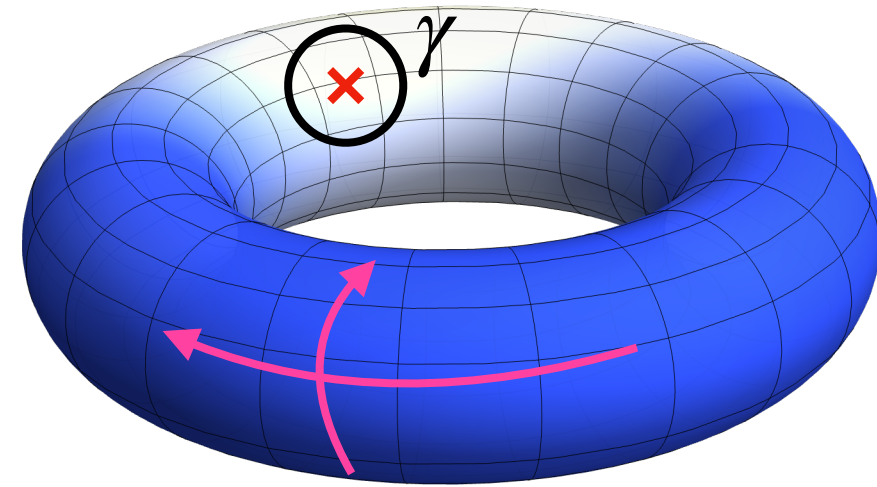


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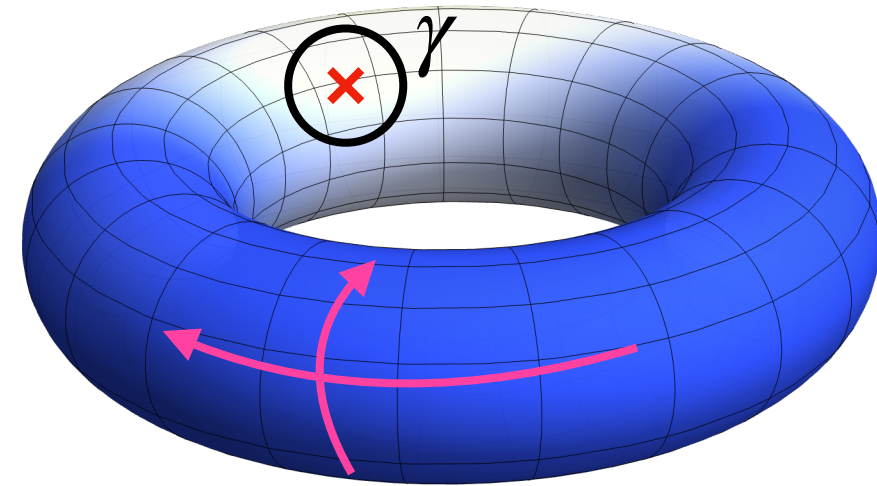


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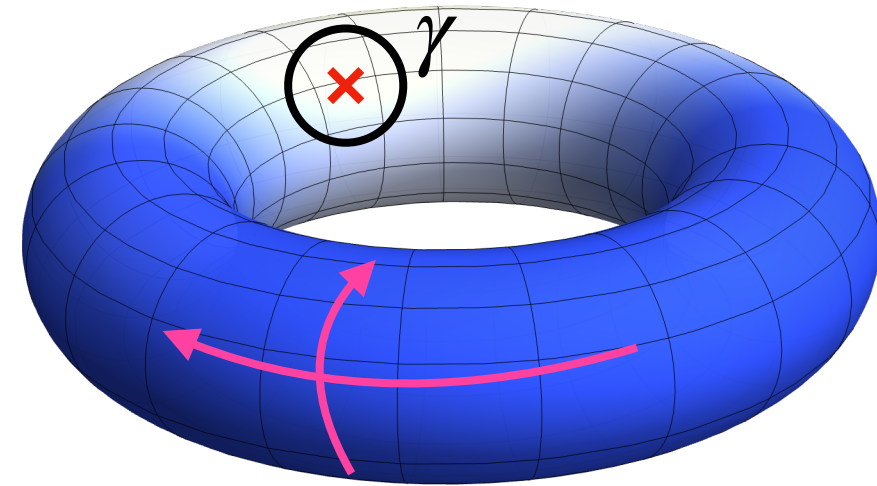
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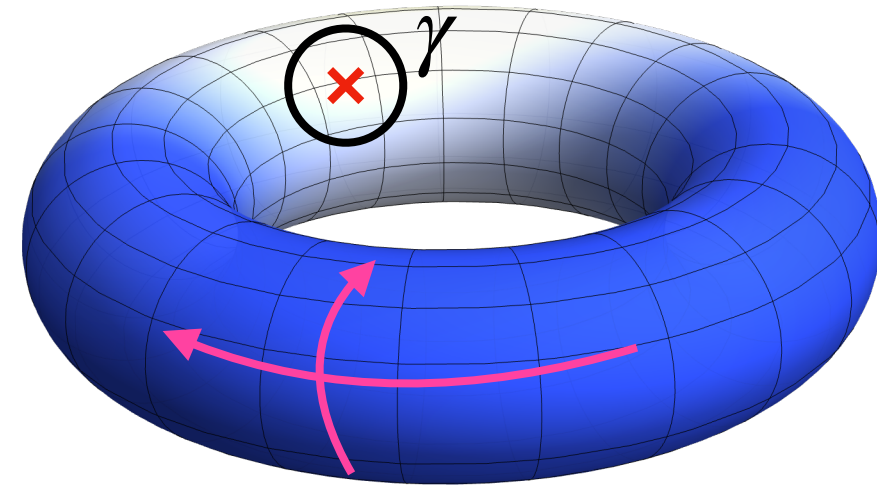
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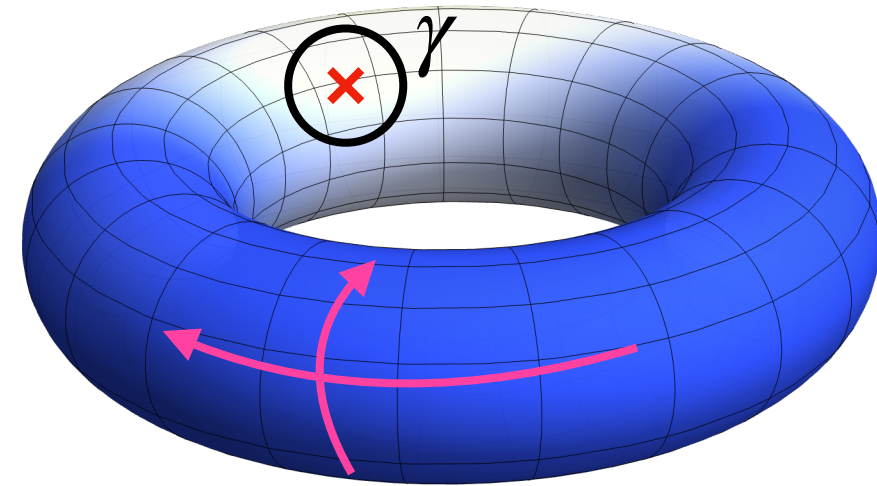


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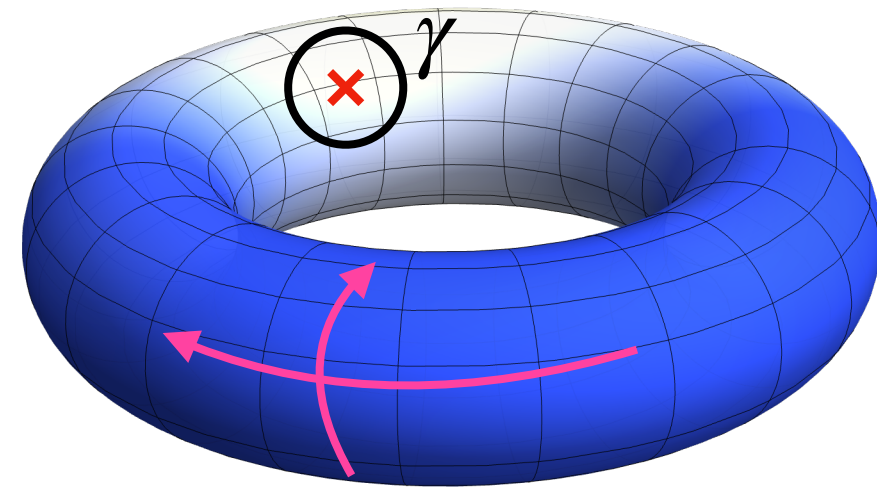
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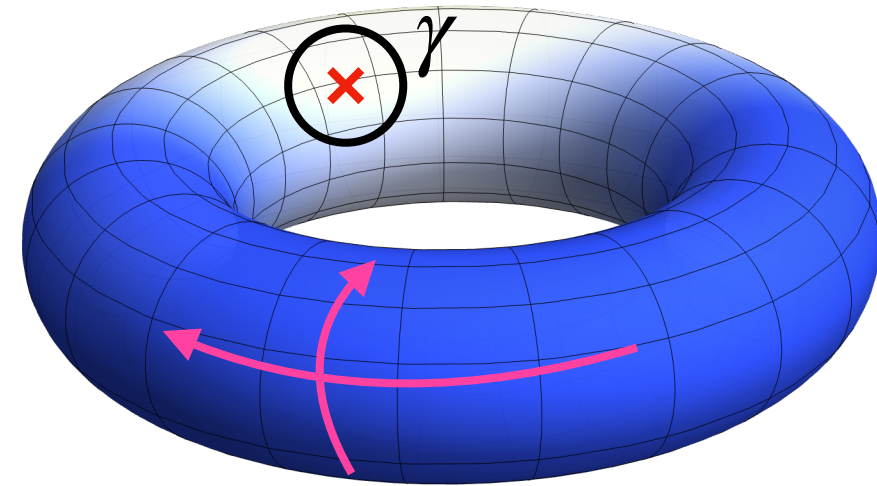


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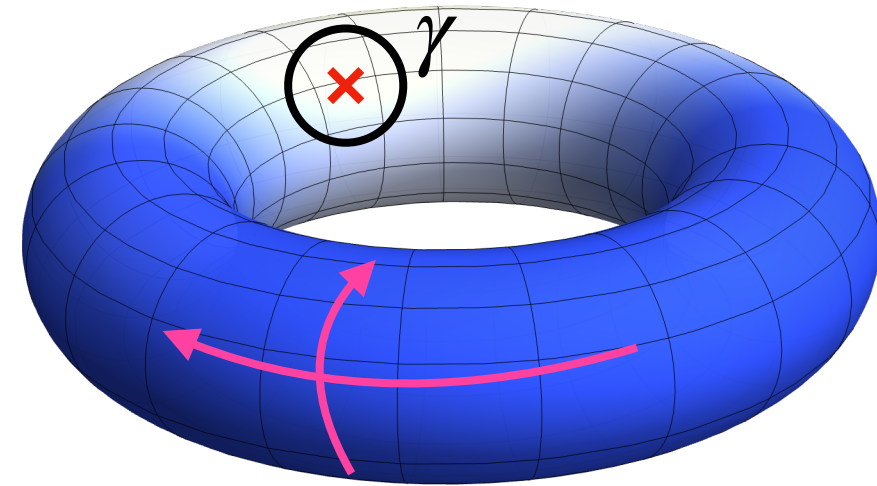
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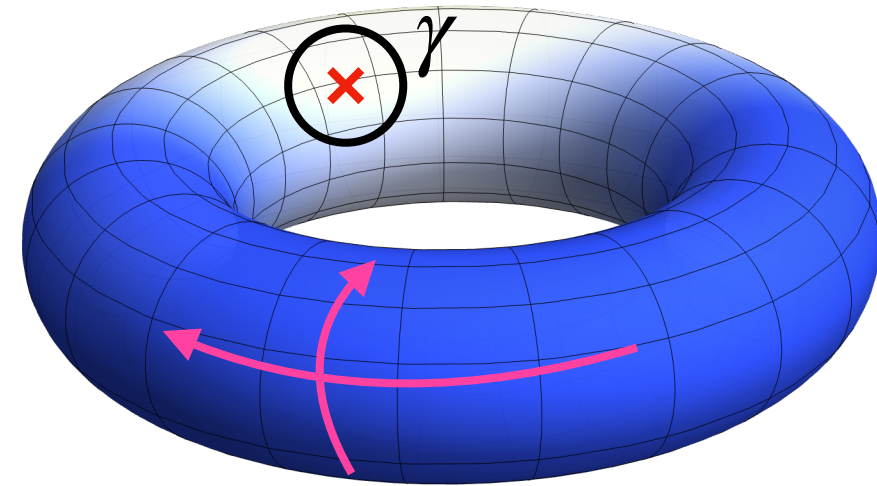


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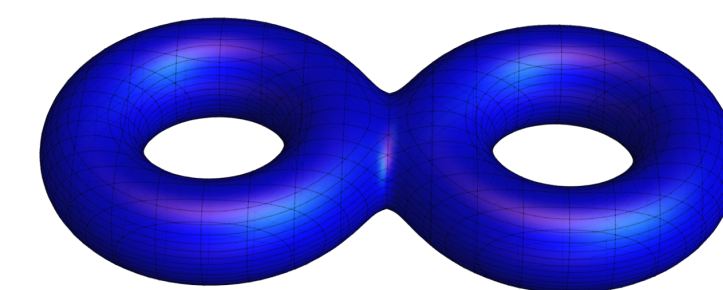
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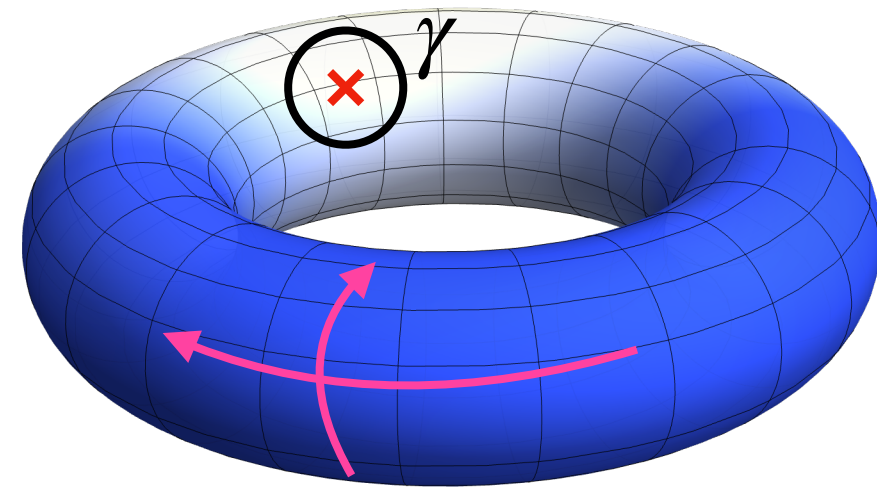


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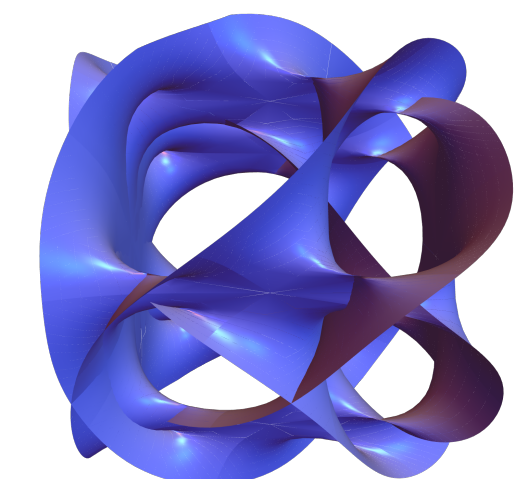
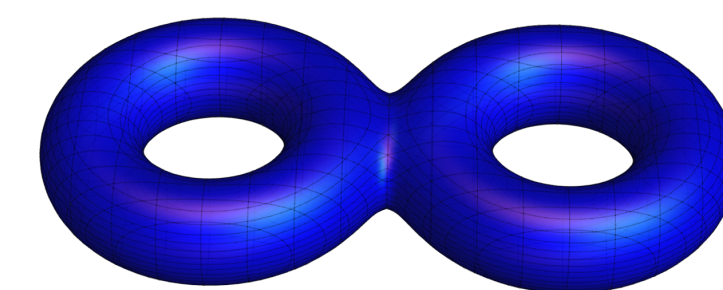
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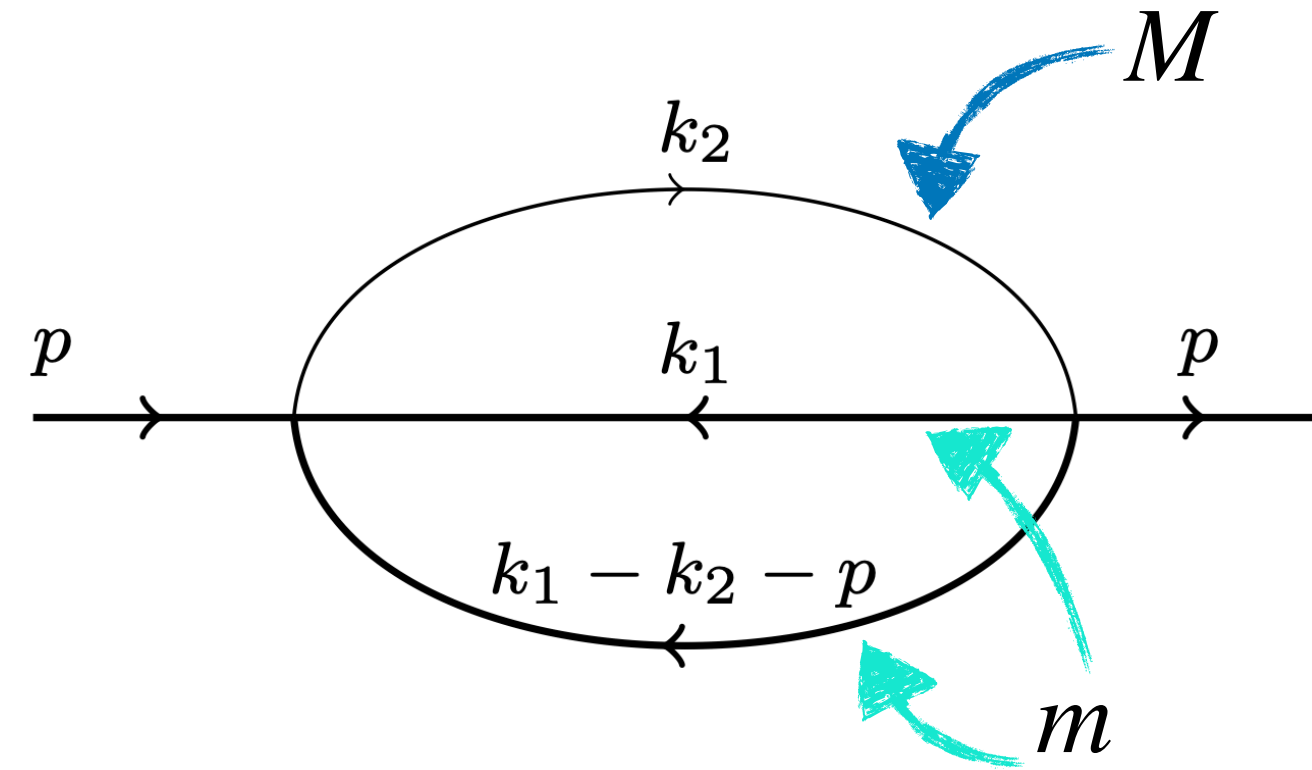
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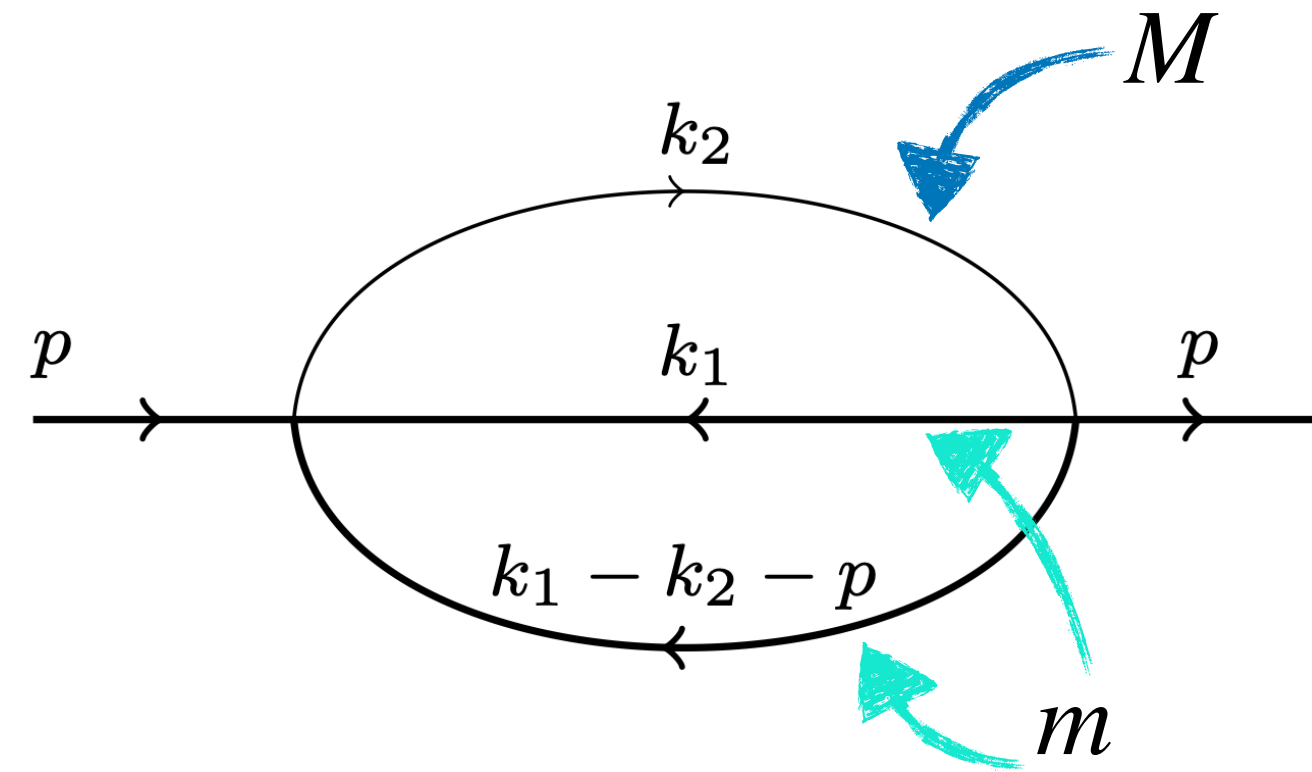
# Running example: two-loop sunrise



$$I_{\nu_1, \dots, \nu_5}(\mathbf{s}; d) = \int \left( \prod_{j=1}^2 \frac{d^d k_j}{i\pi^{d/2}} \right) \frac{(k_1 \cdot p)^{-\nu_4} (k_2 \cdot p)^{-\nu_5}}{(k_1^2 - m^2)^{\nu_1} (k_2^2 - M^2)^{\nu_2} ((k_1 - k_2 - p)^2 - m^2)^{\nu_3}}$$



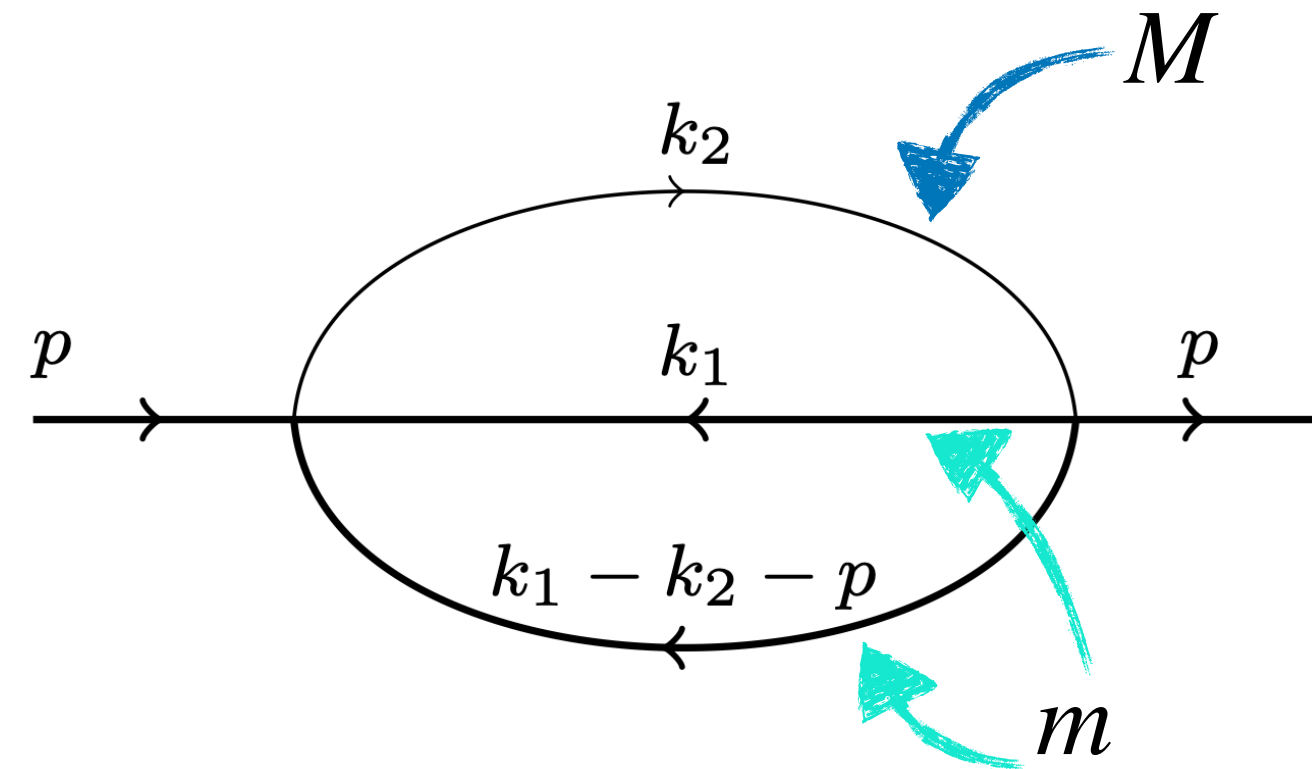
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Two cases:

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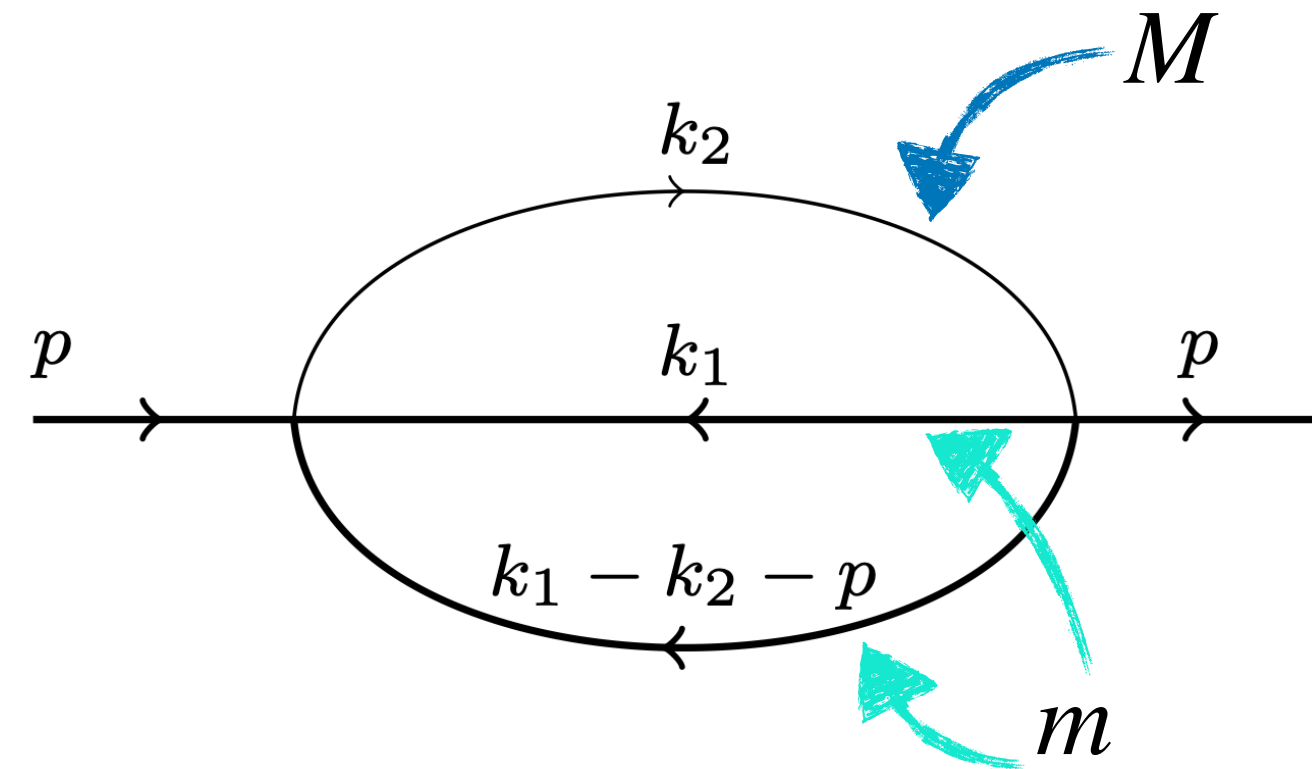


Two cases:

- $M = 0$  polylogarithmic

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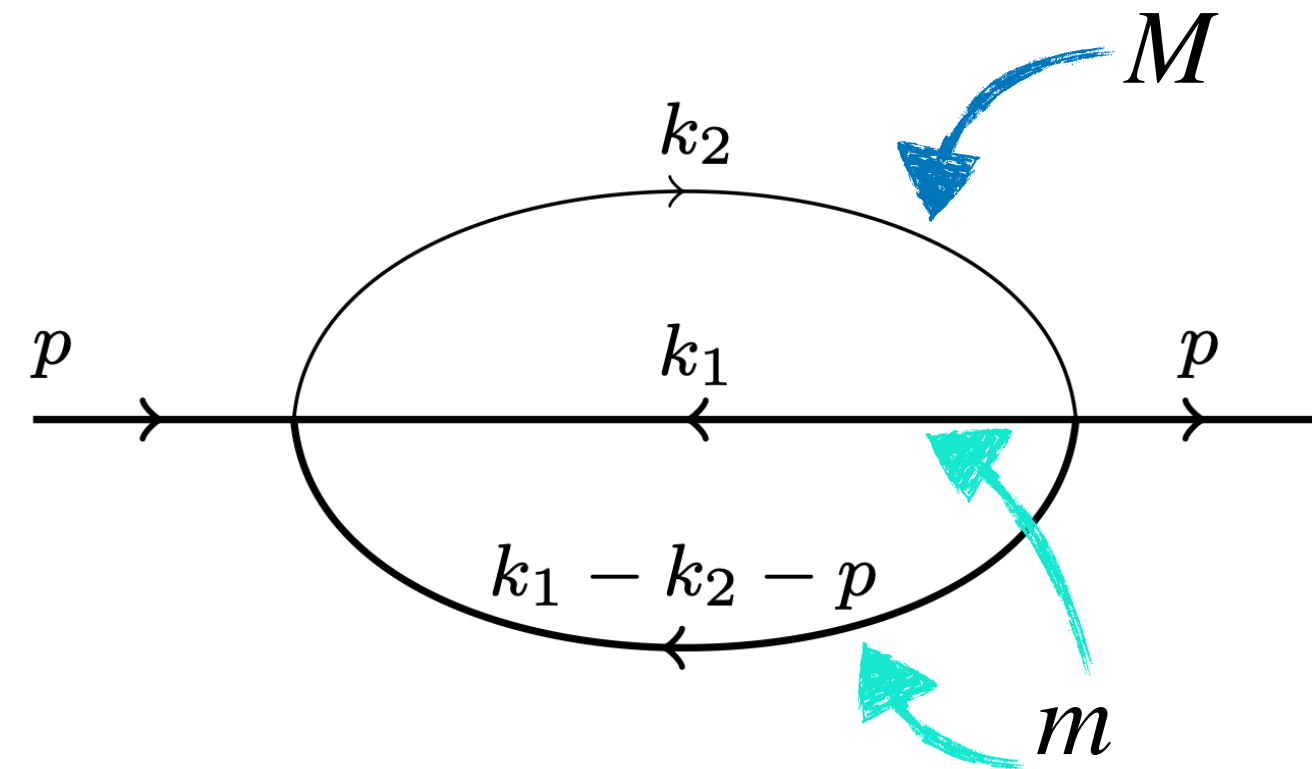
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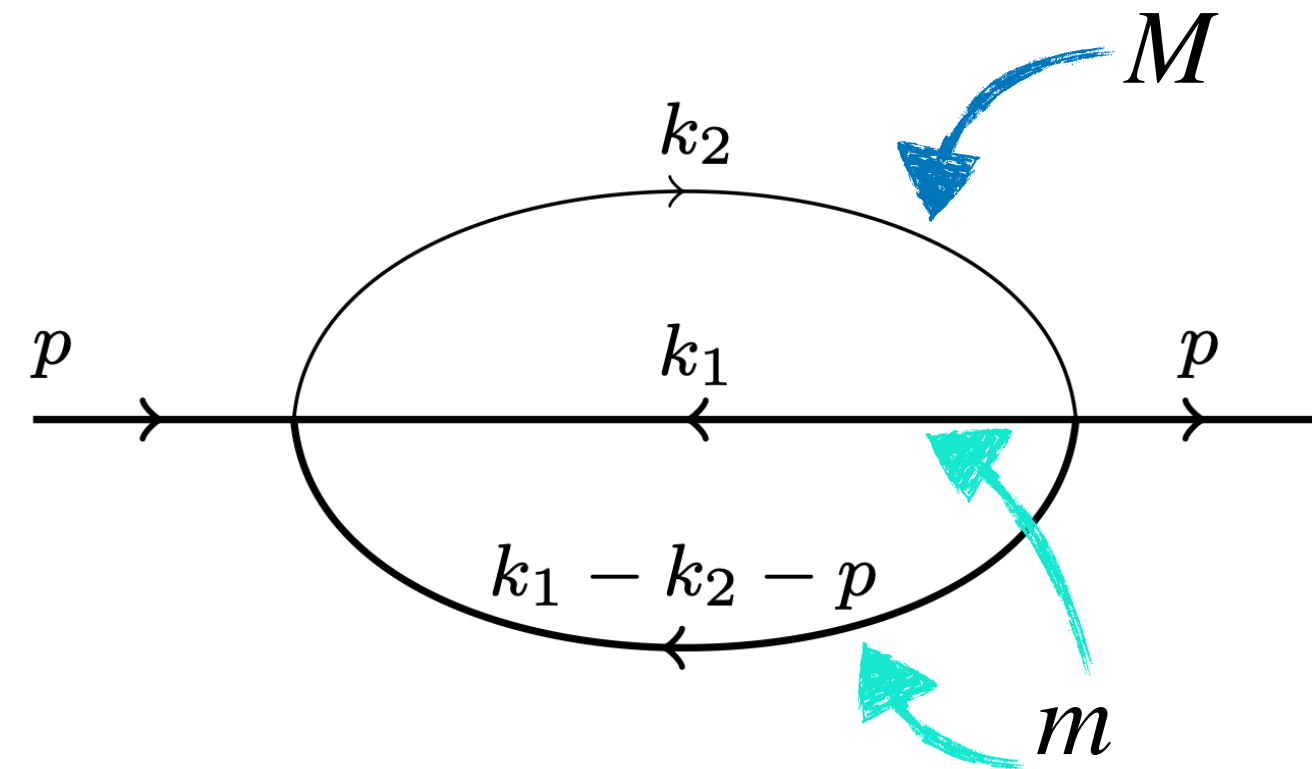
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# Integrand analysis

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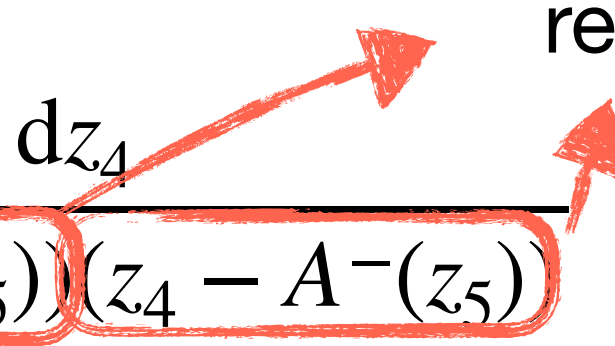
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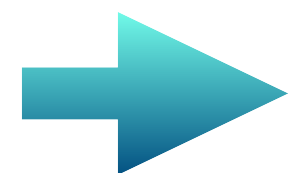
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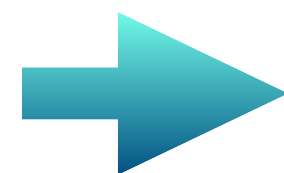
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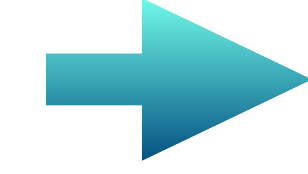
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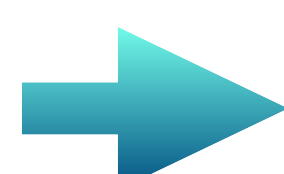
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$\sqrt{P_4(z_5)}$ : elliptic curve!



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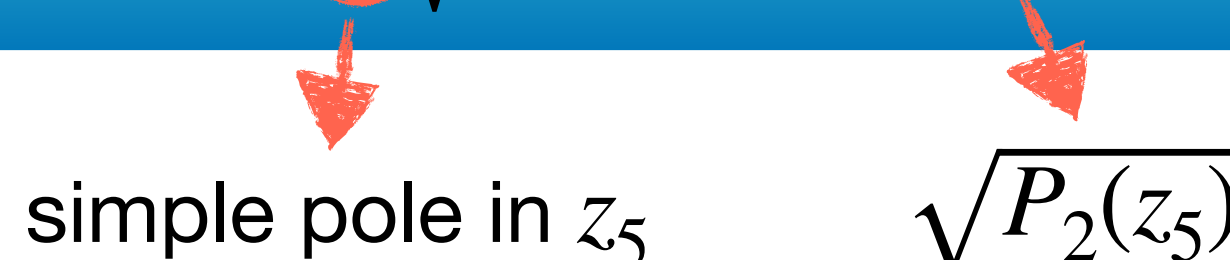
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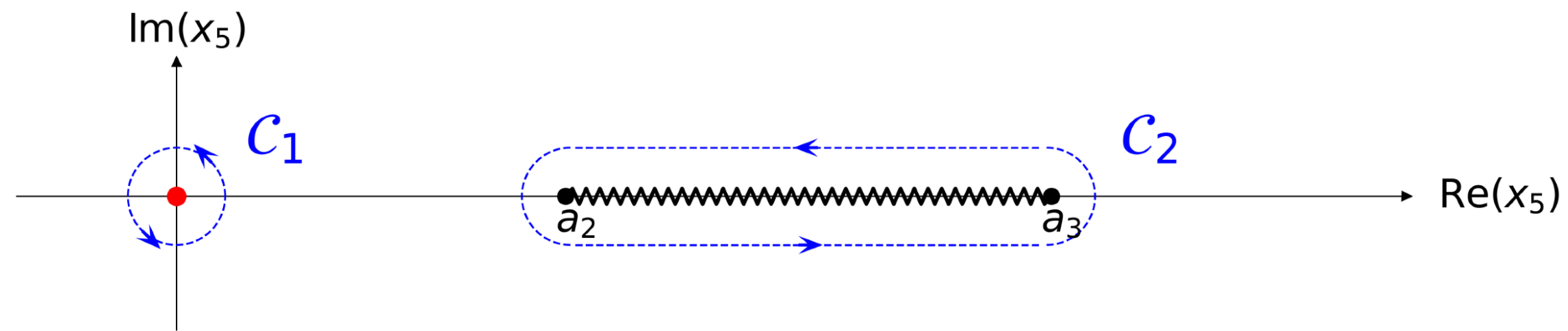
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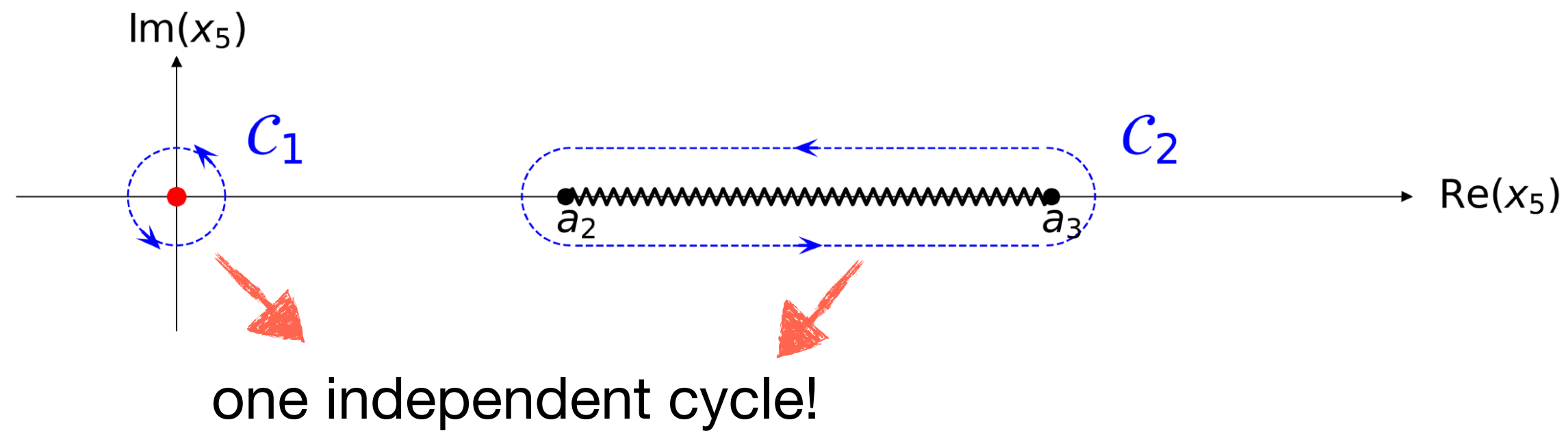


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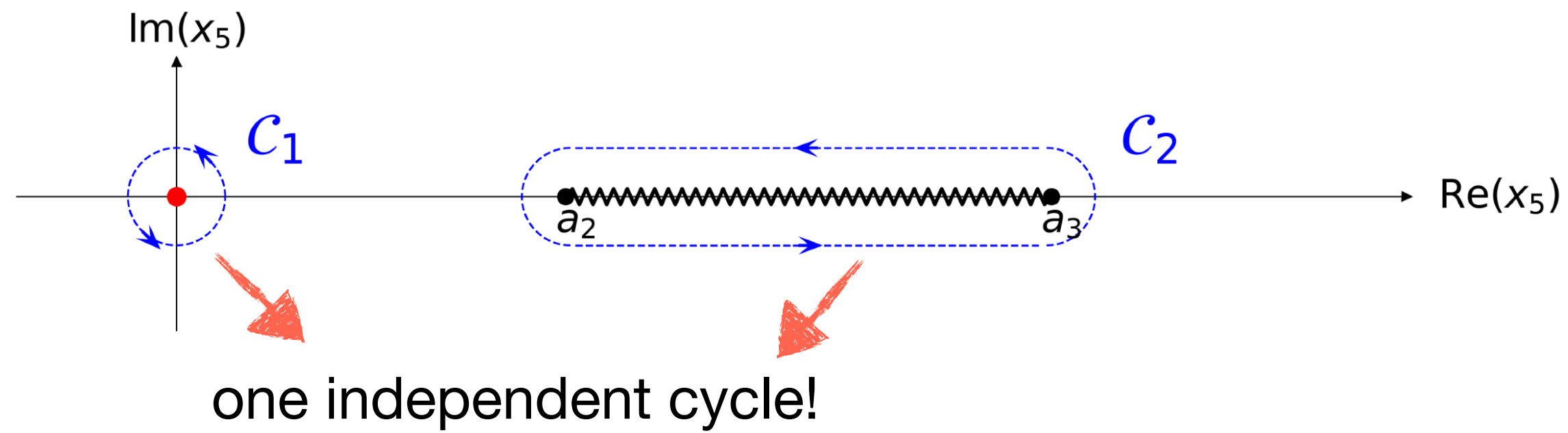


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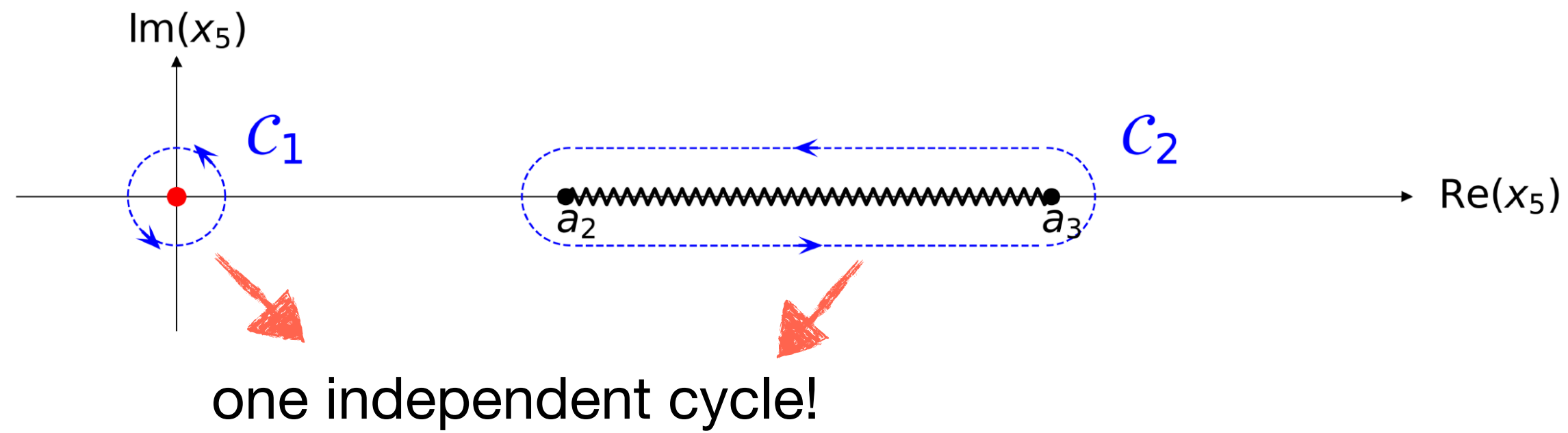
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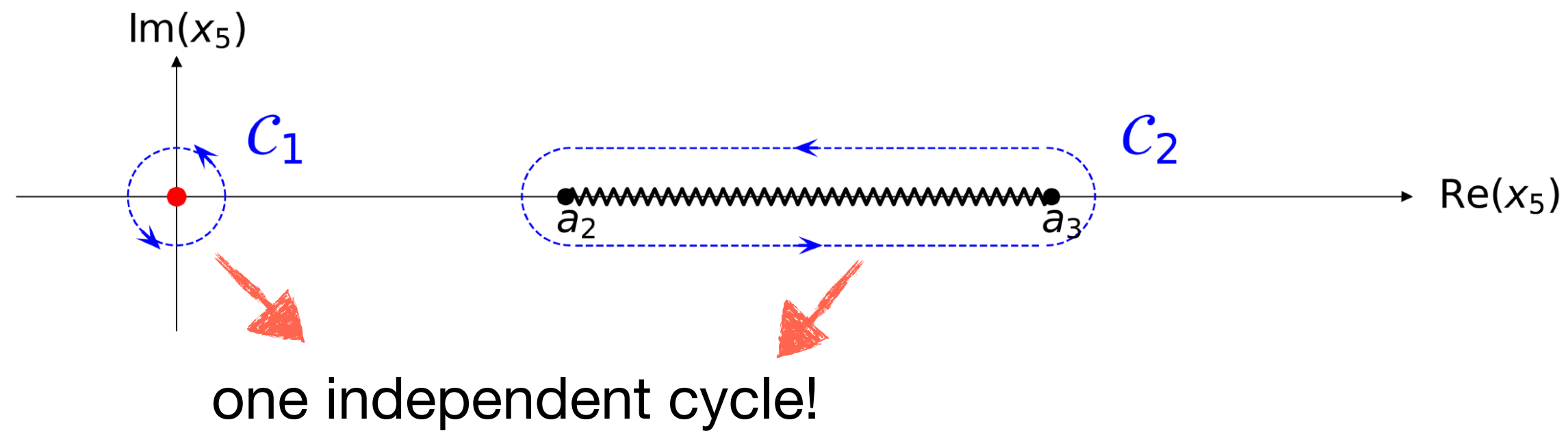


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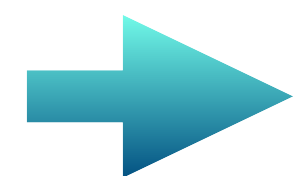
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$$J_1 = \sqrt{s(s - 4m^2)} I_{1,1,1,0,0}$$



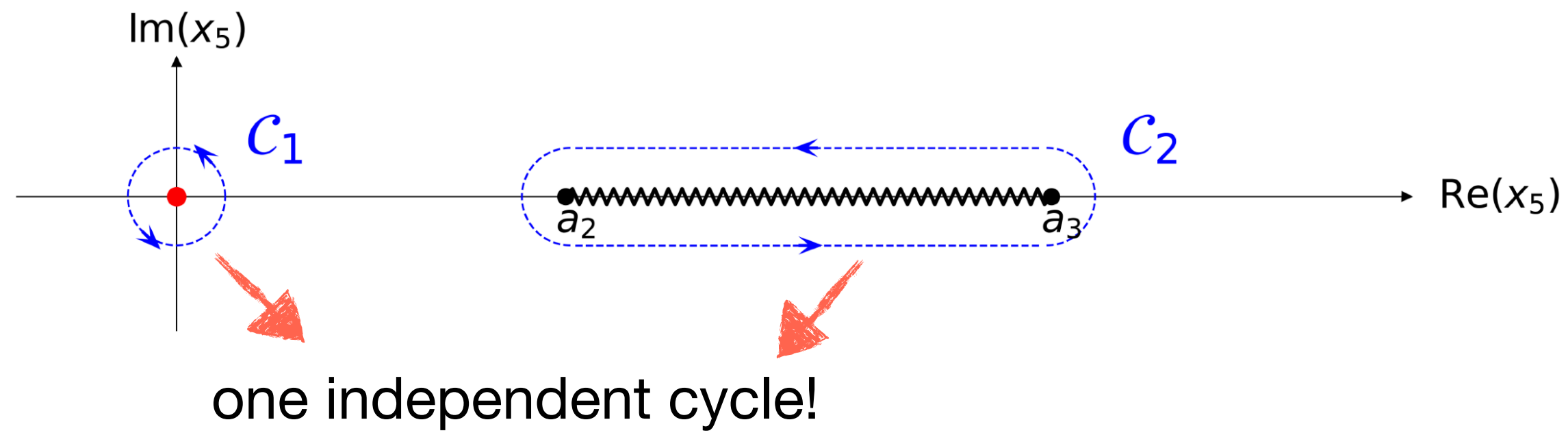
# Case 1: $M = 0$

$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

simple pole in  $z_5$ 
 $\sqrt{P_2(z_5)}$

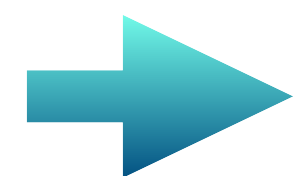
•  $\nu_5 = 0$

•  $\nu_5 = -1$



$$\int_{C_1} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \int_{C_2} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \frac{1}{\sqrt{s(s - 4m^2)}}$$

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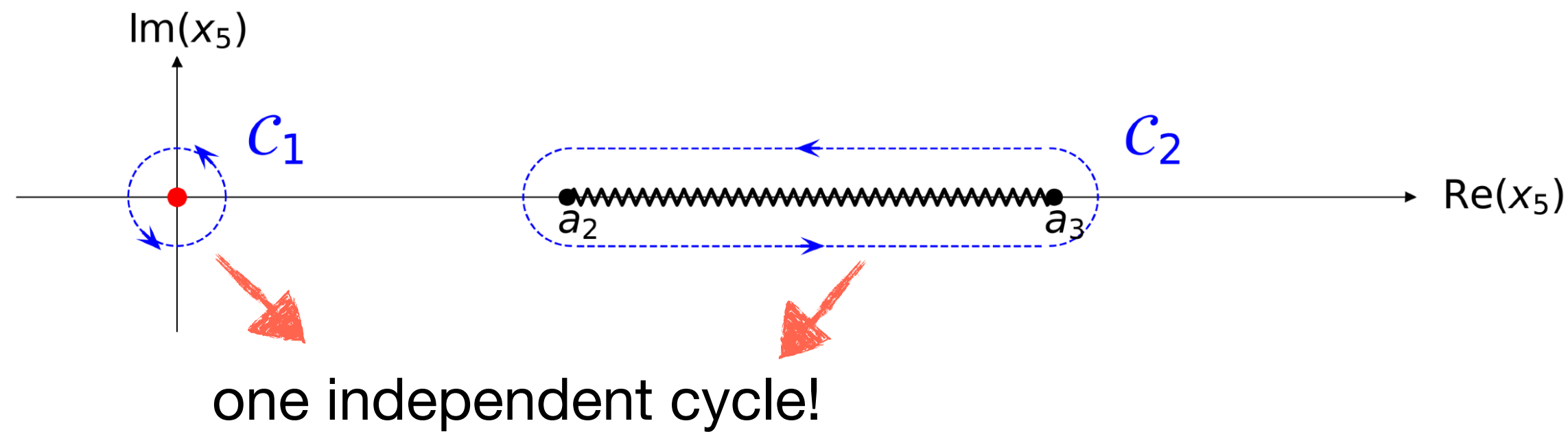
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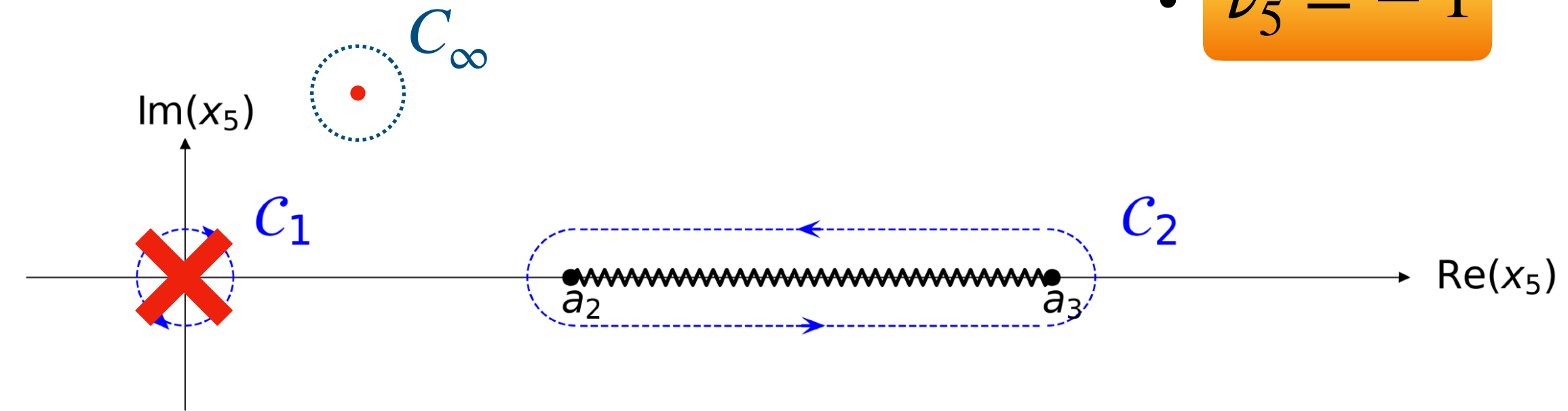
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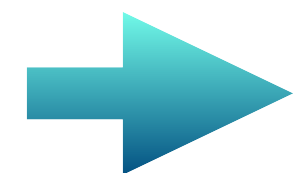


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$$\int_{C_1} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \int_{C_2} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \frac{1}{\sqrt{s(s - 4m^2)}}$$

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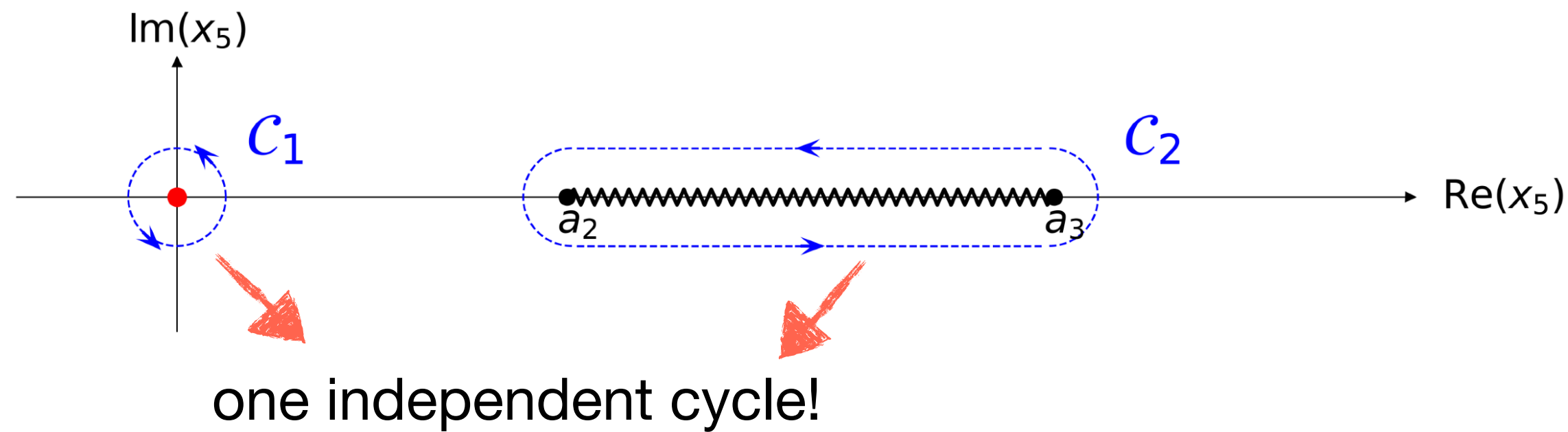
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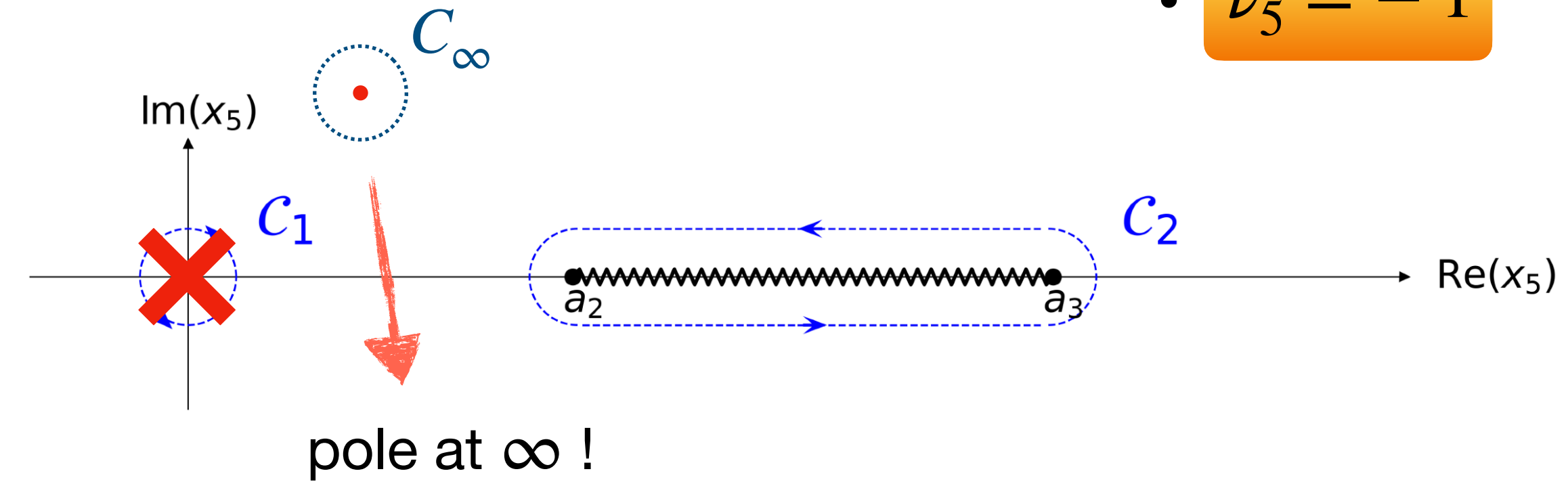
$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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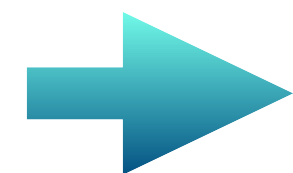


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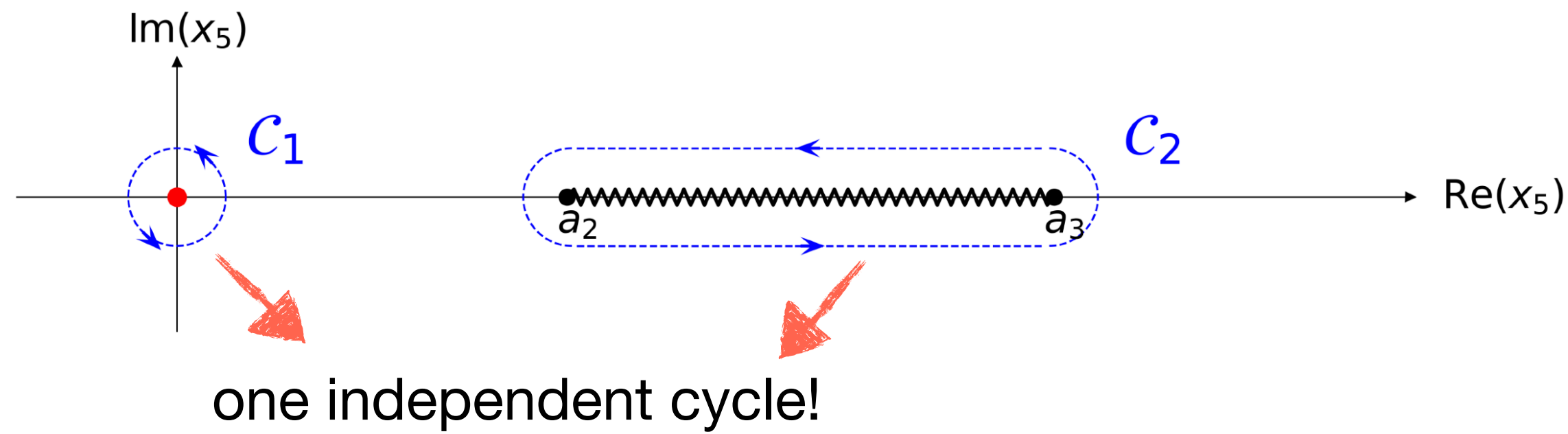
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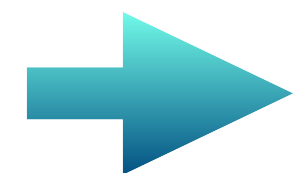
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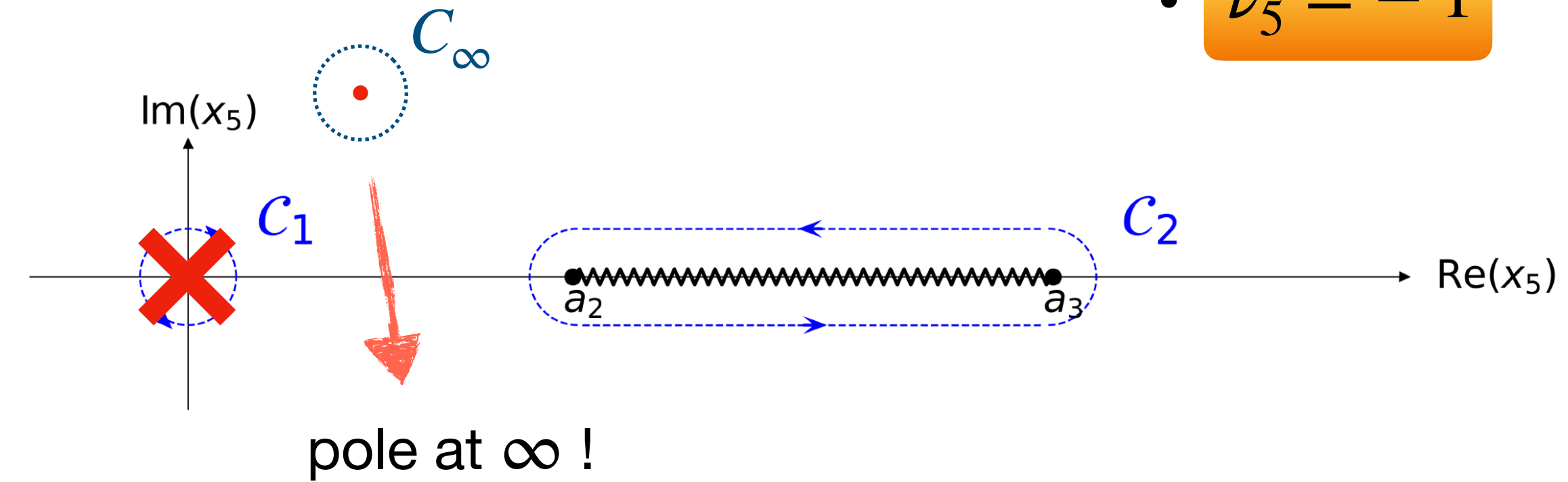
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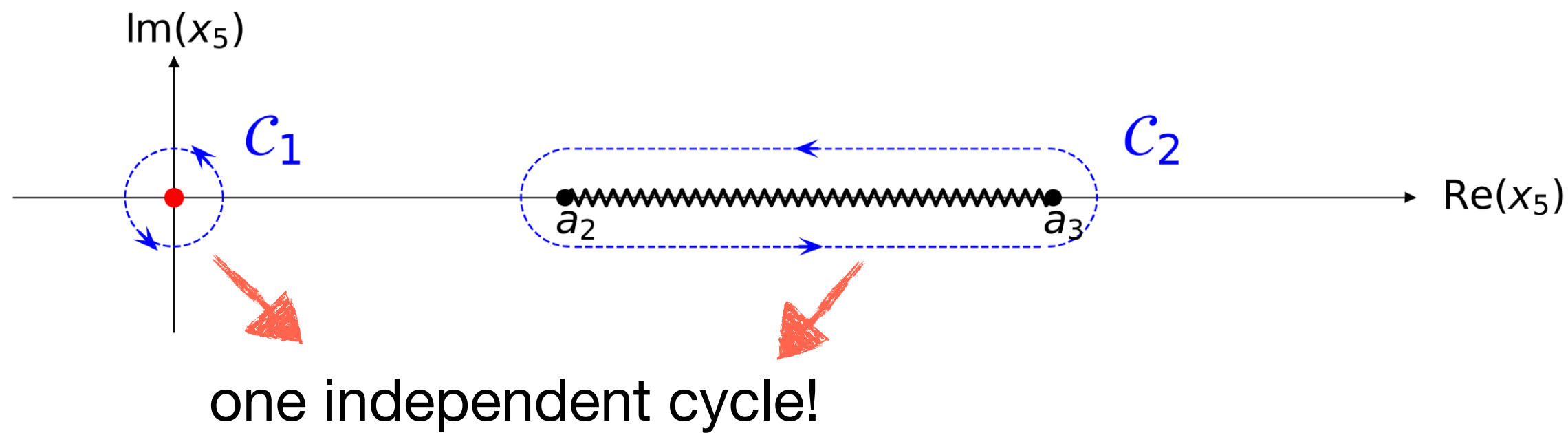
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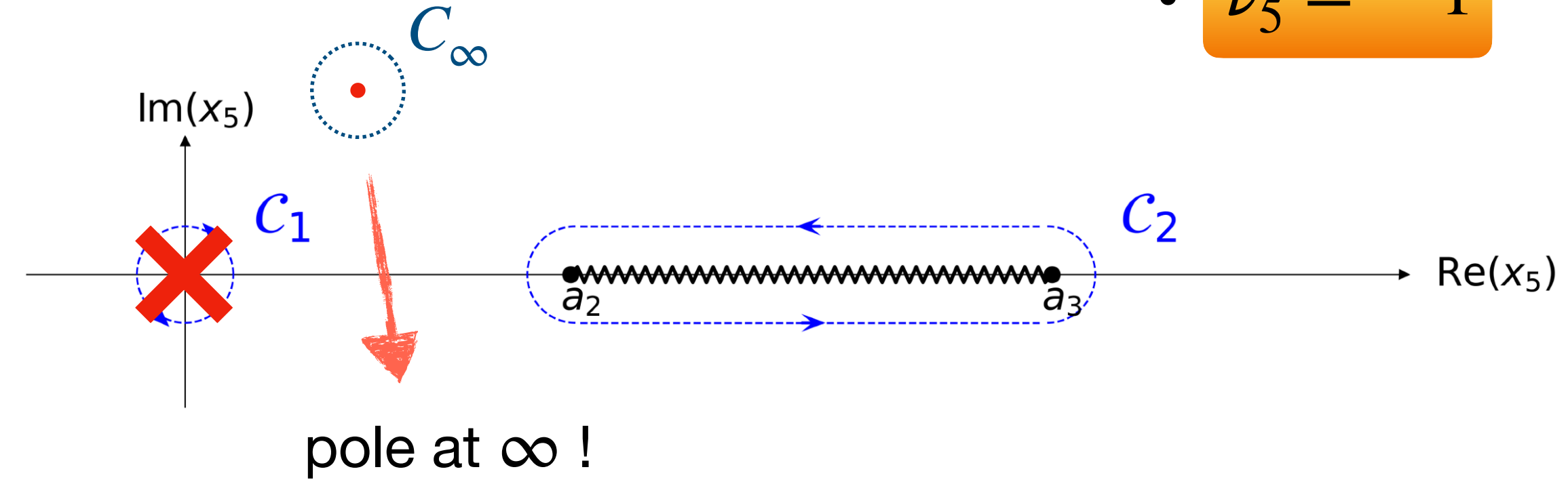
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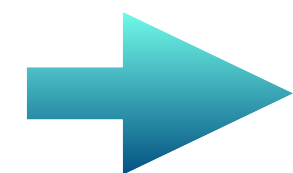
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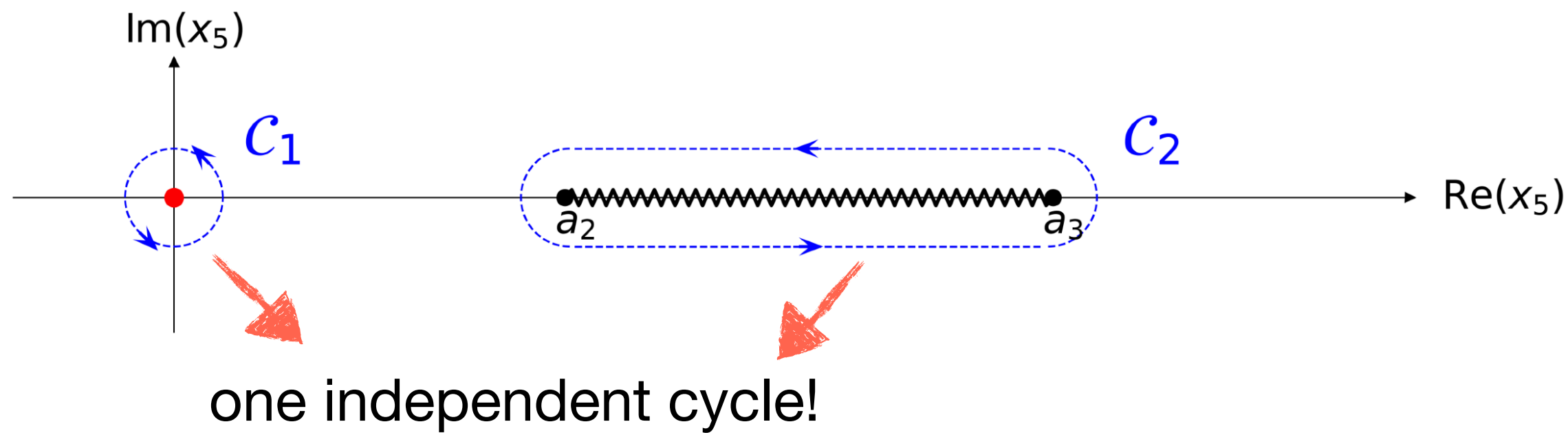
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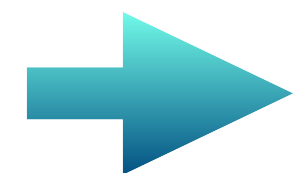
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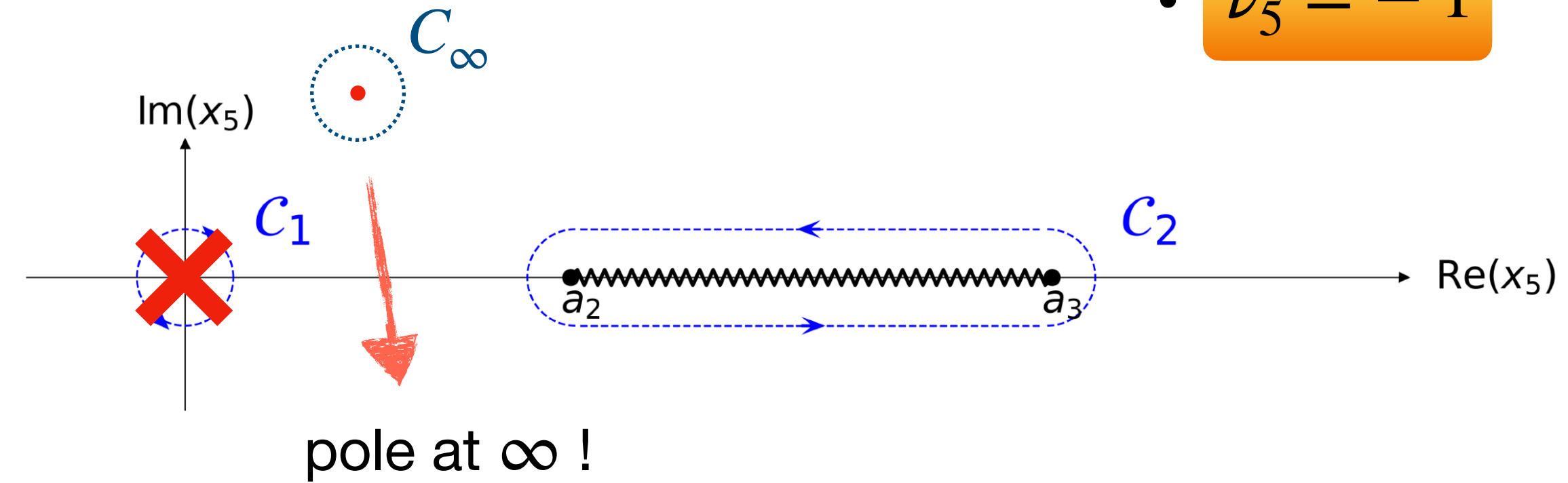
$$\int_{C_1} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \int_{C_2} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \frac{1}{\sqrt{s(s - 4m^2)}}$$

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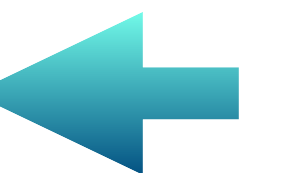
•  $\nu_5 = -1$



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$$J_2 = I_{1,1,1,0,-1}$$



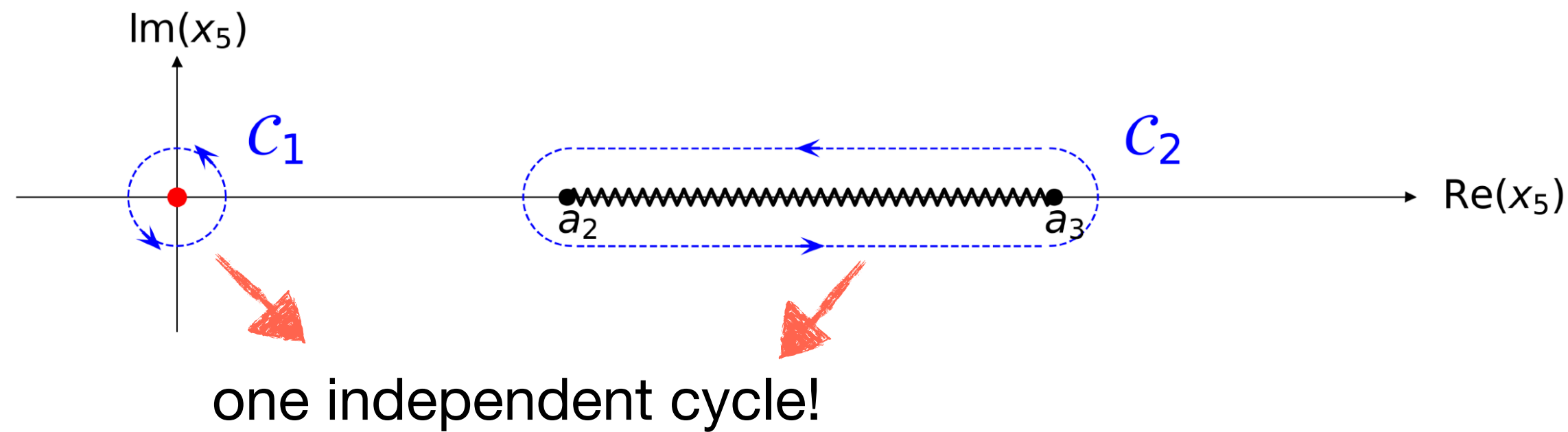
# Case 1: $M = 0$

$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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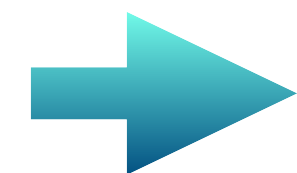
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$$\int_{C_1} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \int_{C_2} \frac{dz_5}{z_5 \sqrt{(2z_5 + s)(4m^2 - s - 2z_5)}} \propto \frac{1}{\sqrt{s(s - 4m^2)}}$$

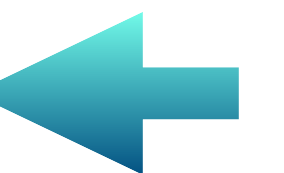
$$I_{1,1,1,0,0} \big|_{z_1=z_2=z_3=0} \sim \frac{1}{\sqrt{s(s - 4m^2)}} \int d \log g_1 \int d \log f$$



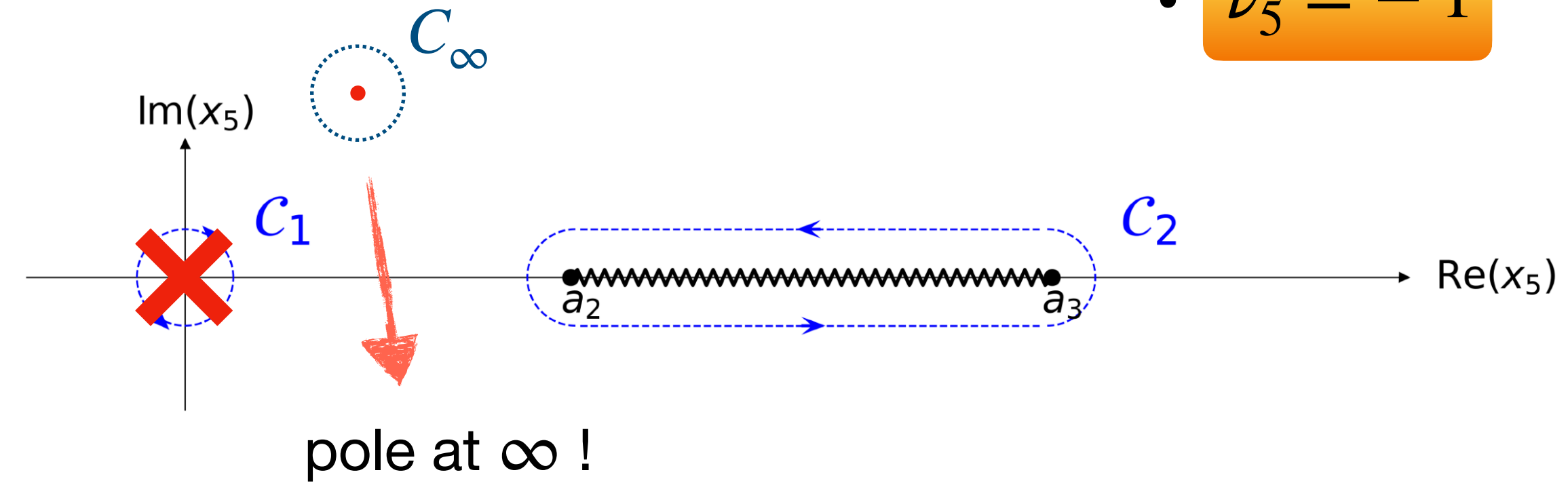
$$J_1 = \sqrt{s(s - 4m^2)} I_{1,1,1,0,0}$$

$$\sum_i R_i(s_{ij}) \int_{\gamma} d \log f_n \wedge \dots \wedge d \log f_1$$

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## Case 2: $M \neq 0$

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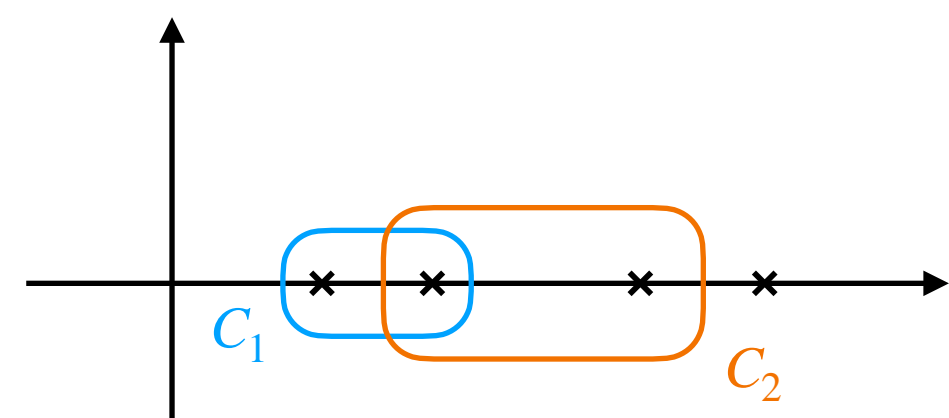
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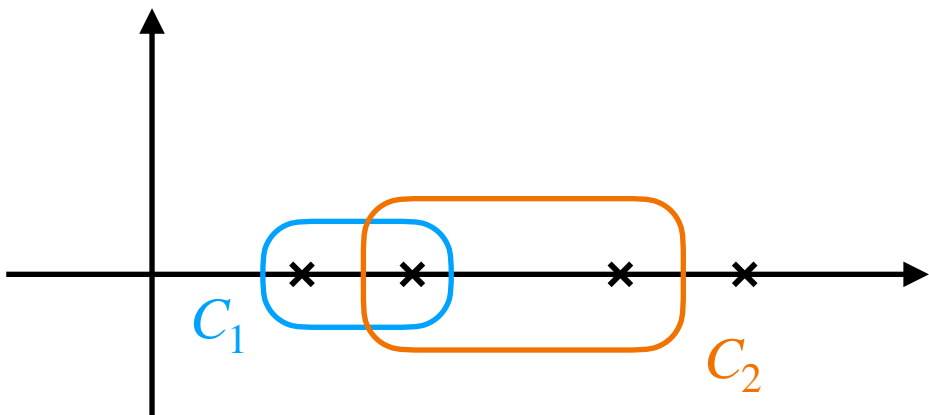
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# Case 2: $M \neq 0$



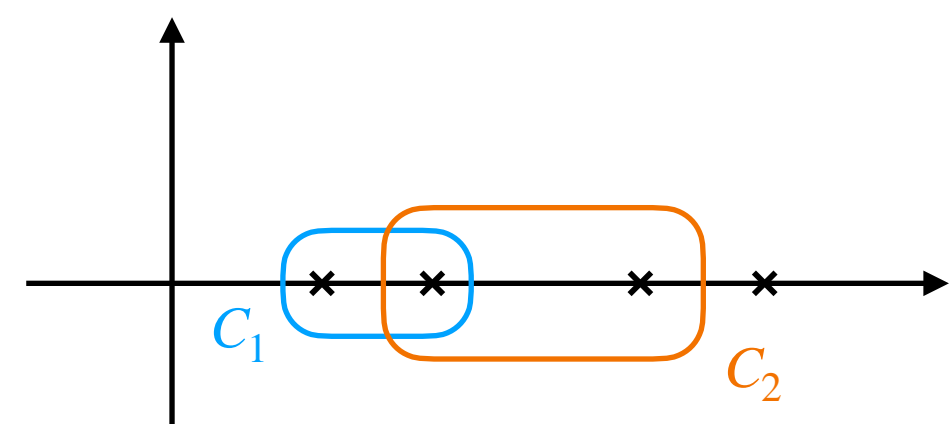
- $\nu_5 = 0$

$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

$\leftarrow$ 

$:= y$

# Case 2: $M \neq 0$



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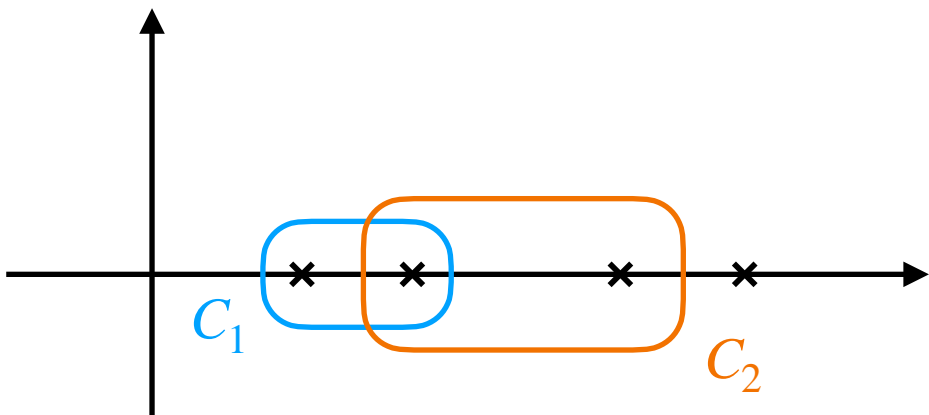
$$I_{1,1,1,0,0}|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5}{y} \int d\log f(z_4, z_5, \mathbf{s})$$

$$I_{1,1,1,0,\nu_5}|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d\log f(z_4, z_5, \mathbf{s})$$

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# Case 2: $M \neq 0$



- $\nu_5 = 0$  1st kind integral

$$I_{1,1,1,0,0} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5}{y} \int d \log f(z_4, z_5, s)$$

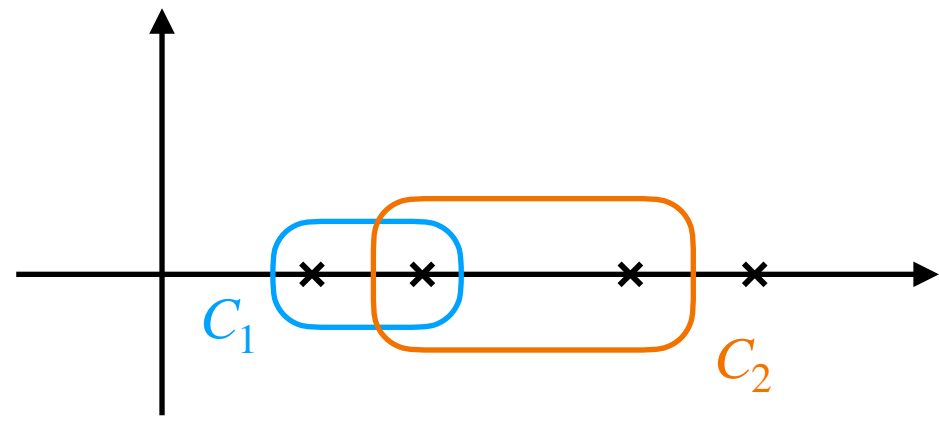
$$\int_{C_1} \frac{dz_5}{y} \propto \psi_0$$

$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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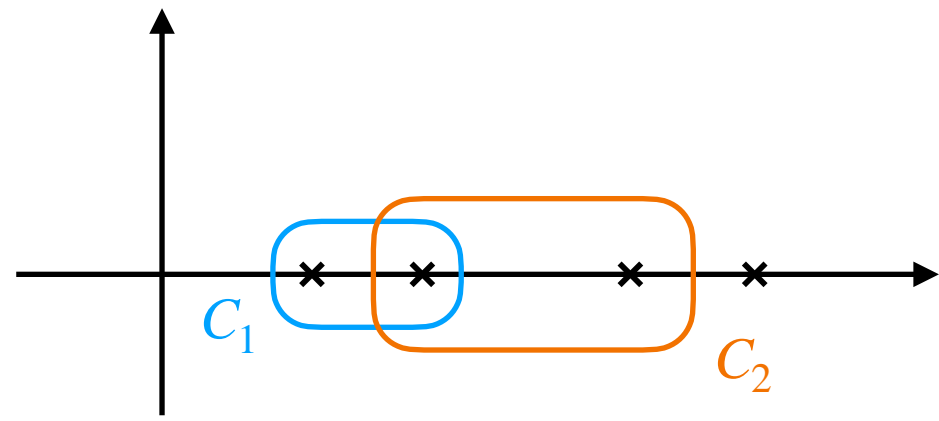
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holomorphic period of  
the elliptic curve!

$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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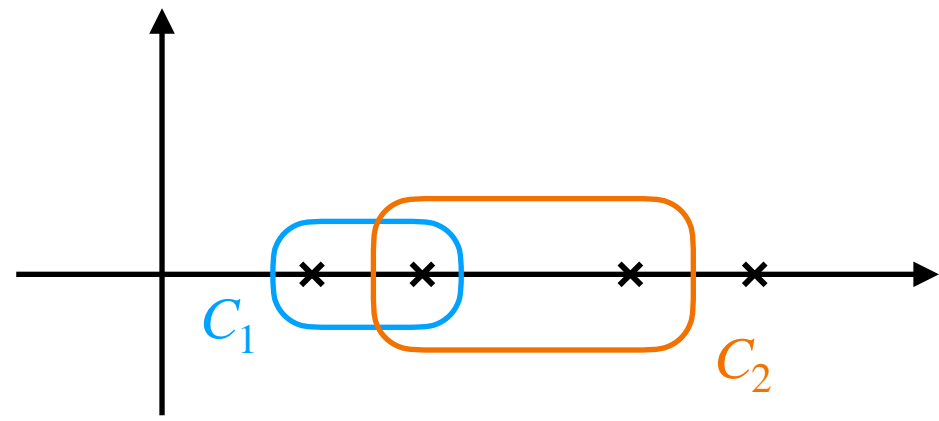
holomorphic period of the elliptic curve!

$$J_1 = \frac{1}{\psi_0} I_{1,1,1,0,0}$$

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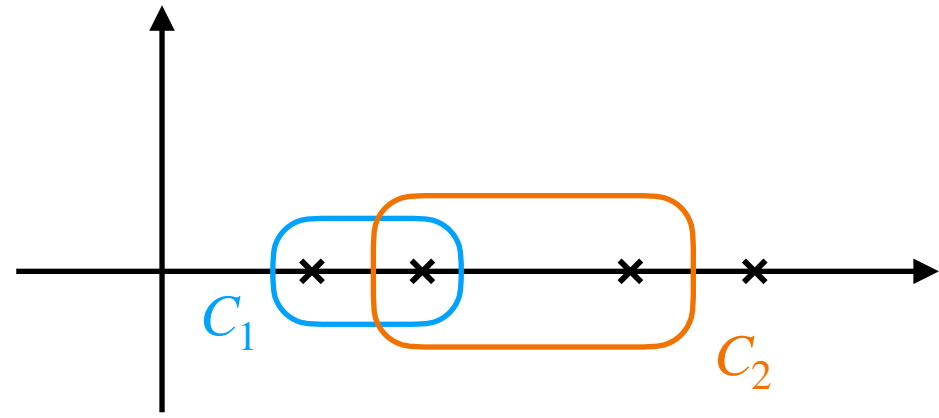
$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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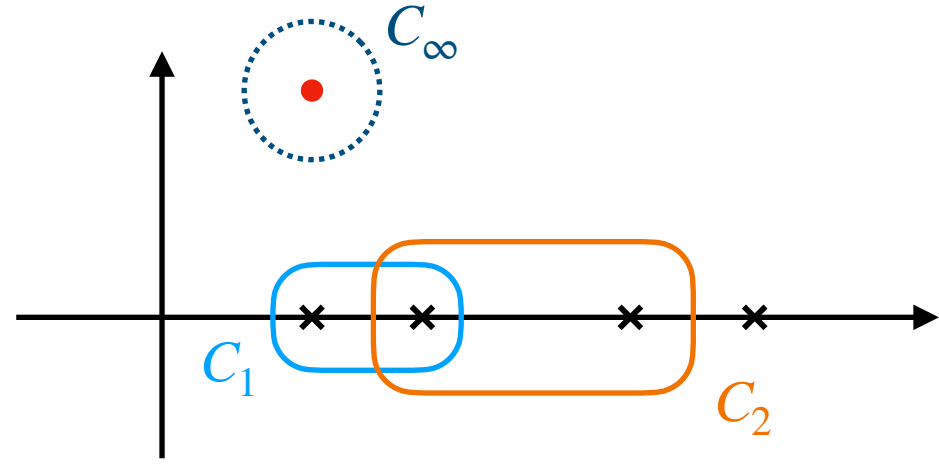
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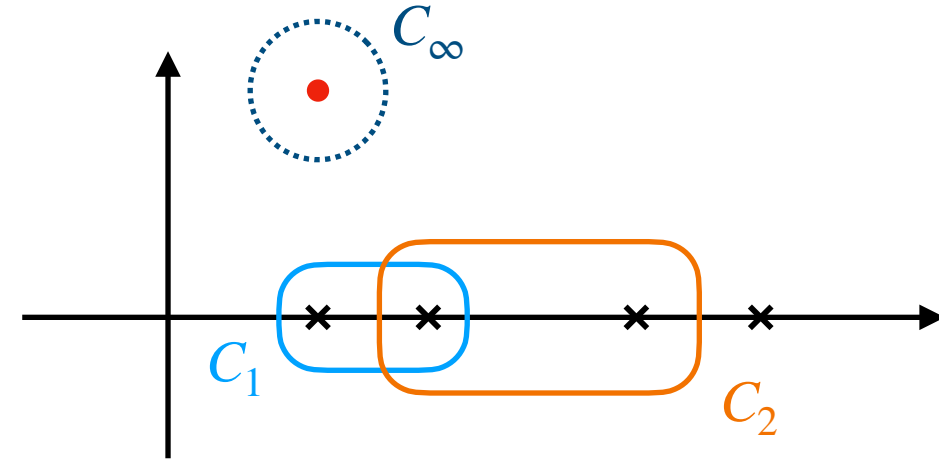
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$$I_{1,1,1,0,-1} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5}{y} \int d \log f(z_4, z_5, \mathbf{s})$$

## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

$:= y$

- $\nu_5 = 0$  1st kind integral

$$I_{1,1,1,0,0} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5}{y} \int d \log f(z_4, z_5, s)$$

$$\int_{C_1} \frac{dz_5}{y} \propto \psi_0 \quad \text{holomorphic period of the elliptic curve!}$$

$$J_1 = \frac{1}{\psi_0} I_{1,1,1,0,0}$$

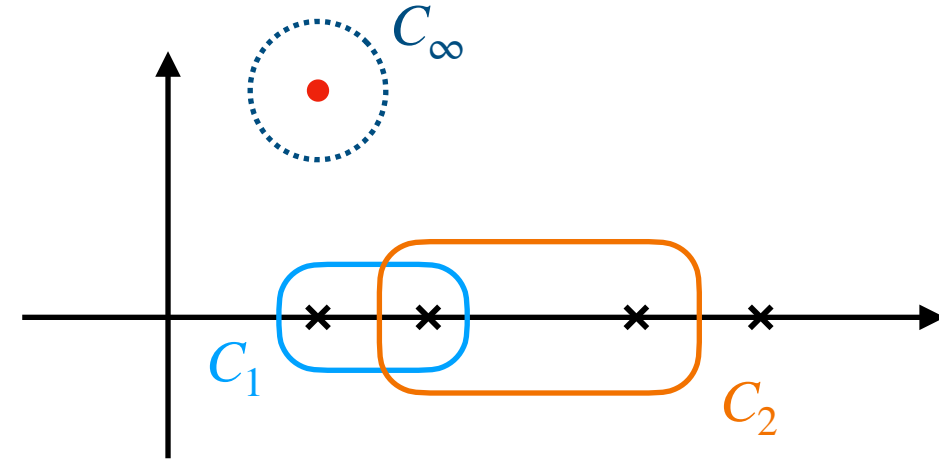
- $\nu_5 = -1$  3rd kind integral

$$I_{1,1,1,0,-1} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5}{y} \int d \log f(z_4, z_5, s)$$

$$\int_{C_\infty} \frac{dz_5 z_5}{y} \propto 1 \quad \text{simple pole at } \infty$$



## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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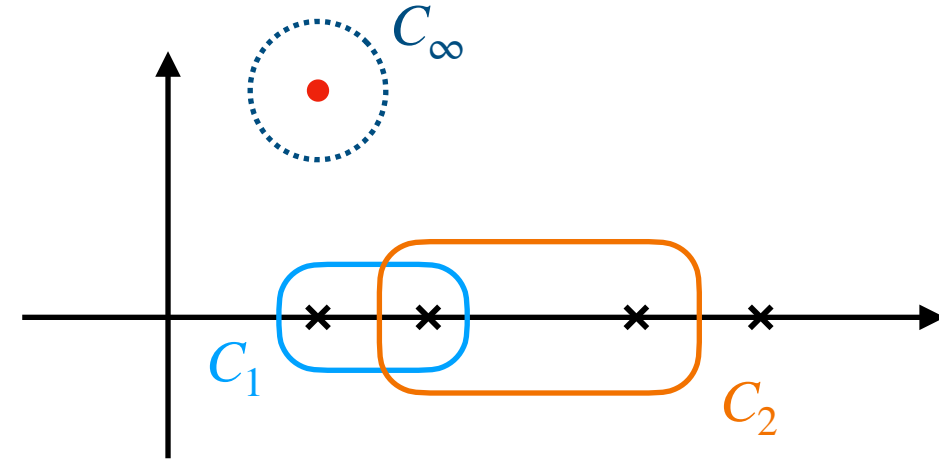
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$$J_2 = I_{1,1,1,0,-1}$$



## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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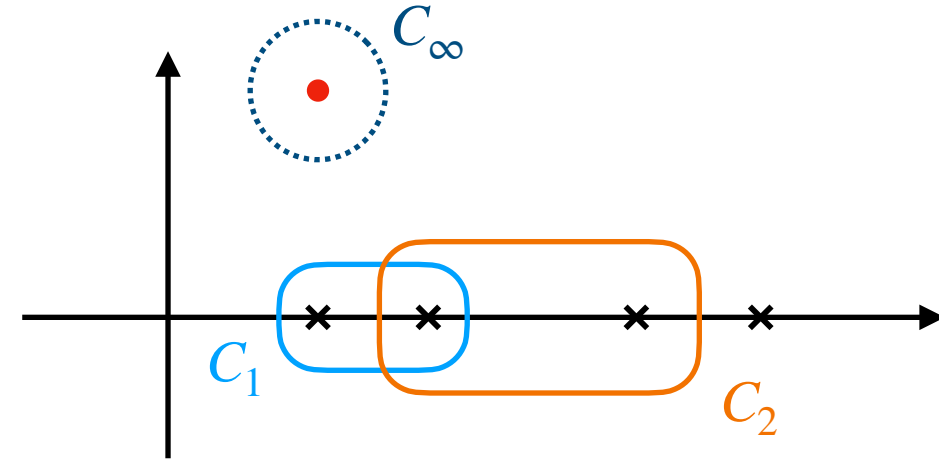
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- $\nu_5 = -2$

## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

$$:= y$$

- $\nu_5 = 0$  1st kind integral

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holomorphic period of the elliptic curve!

$$J_1 = \frac{1}{\psi_0} I_{1,1,1,0,0}$$

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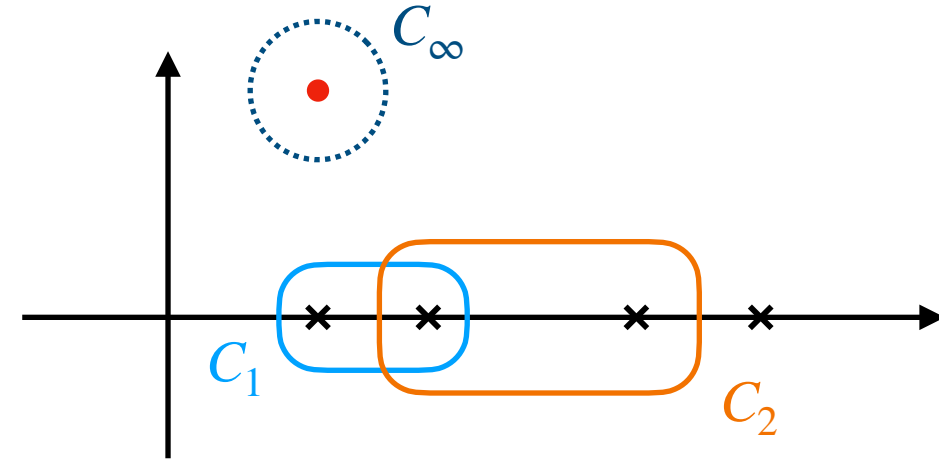
simple pole at  $\infty$

$$J_2 = I_{1,1,1,0,-1}$$

- $\nu_5 = -2$

$$I_{1,1,1,0,-2} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^2}{y} \int d \log f(z_4, z_5, s)$$

## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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simple pole at  $\infty$

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- $\nu_5 = -2$  2nd kind integral

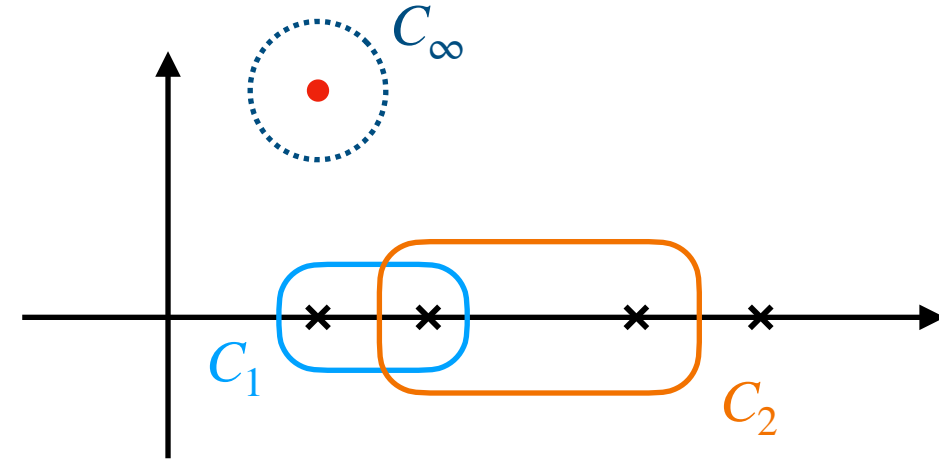
$$I_{1,1,1,0,-2} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^2}{y} \int d \log f(z_4, z_5, s)$$

$$\int \frac{dz_5 z_5^2}{y}$$

double pole at  $\infty$  (no residue)



## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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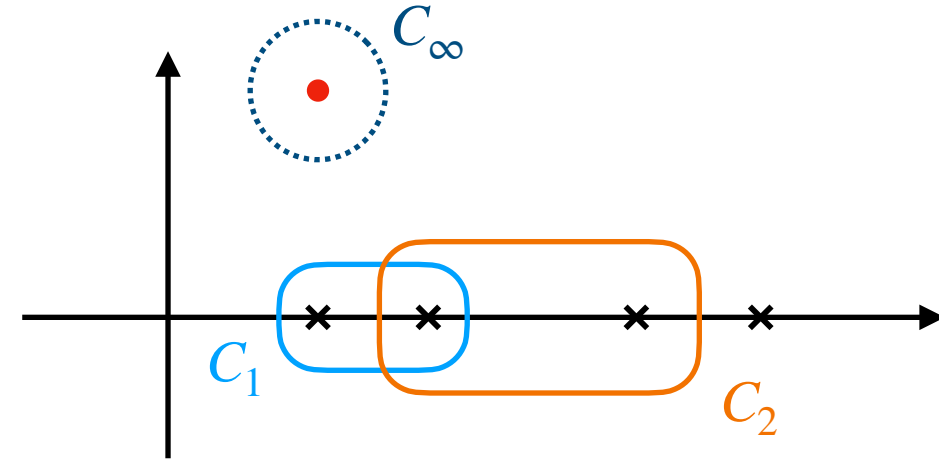
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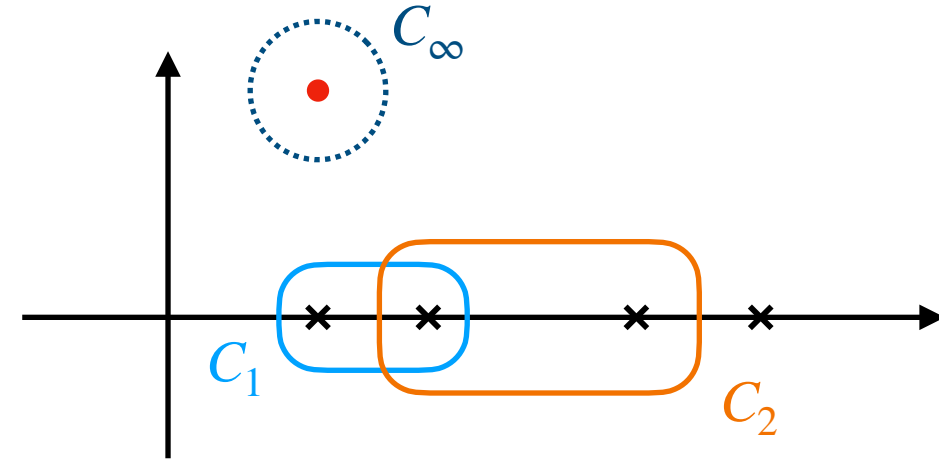
$$\int \frac{dz_5 z_5^2}{y}$$

double pole at  $\infty$  (no residue)



**no log singularities?!**

## Case 2: $M \neq 0$



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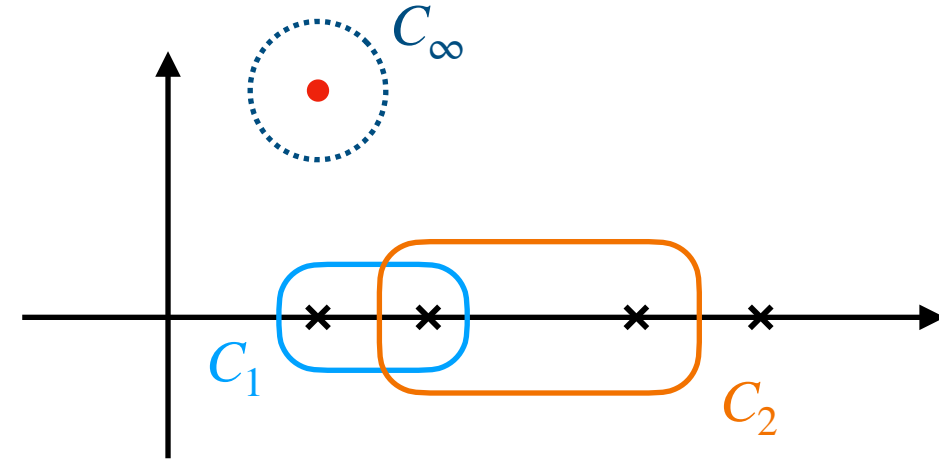
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**no log singularities?!**

[Görge, Nega, Tancredi, Wagner '23]  
[Duhr, **SM**, Nega, Sauer, Tancredi, Wagner '25]



## Case 2: $M \neq 0$



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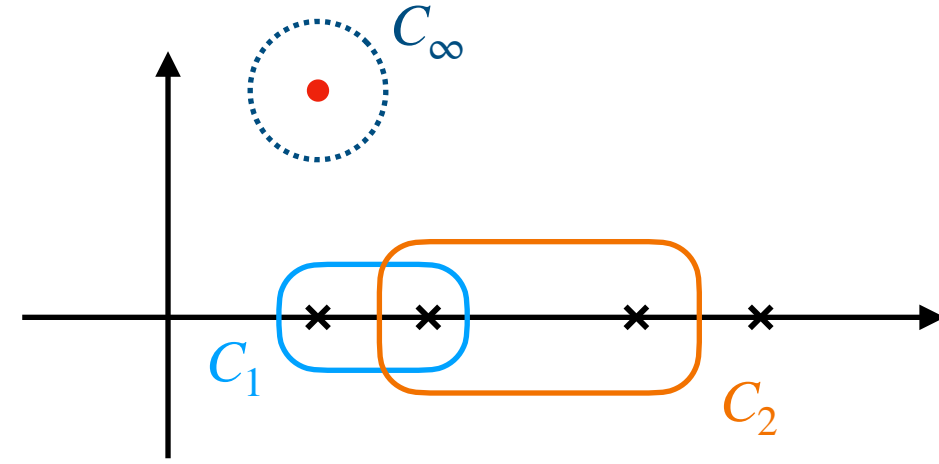
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[Görge, Nega, Tancredi, Wagner '23]  
[Duhr, **SM**, Nega, Sauer, Tancredi, Wagner '25]

$$\tilde{I}_3 = \partial I_{1,1,1,0,0}$$



## Case 2: $M \neq 0$



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simple pole at  $\infty$

$$J_2 = I_{1,1,1,0,-1}$$

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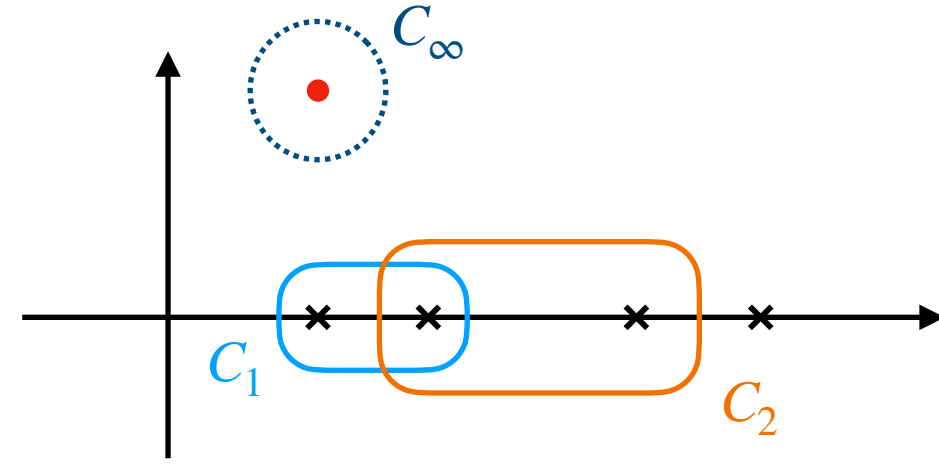
[Görge, Nega, Tancredi, Wagner '23]

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- $\varepsilon$ -factorised differential equations:  $\partial_{s_{ij}} \mathbf{J} = \varepsilon \mathbf{A}(s_{ij}) \mathbf{J}$

## Case 2: $M \neq 0$



$$I_{1,1,1,0,\nu_5} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^{-\nu_5}}{\sqrt{(2z_5 + s + M^2)(M^2 s - z_5^2)(4m^2 - M^2 - s - 2z_5)}} \int d \log f(z_4, z_5, s)$$

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$$I_{1,1,1,0,0} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5}{y} \int d \log f(z_4, z_5, s)$$

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holomorphic period of the elliptic curve!

$$J_1 = \frac{1}{\psi_0} I_{1,1,1,0,0}$$

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simple pole at  $\infty$

$$J_2 = I_{1,1,1,0,-1}$$

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$$I_{1,1,1,0,-2} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5 z_5^2}{y} \int d \log f(z_4, z_5, s)$$

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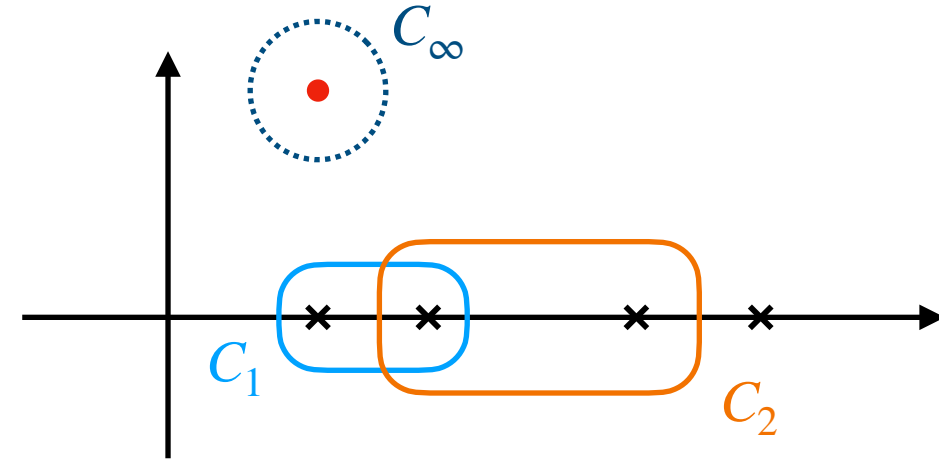
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## Case 2: $M \neq 0$



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$$:= y$$

- $\nu_5 = 0$  1st kind integral

$$I_{1,1,1,0,0} \big|_{z_1=z_2=z_3=0} \sim \int \frac{dz_5}{y} \int d \log f(z_4, z_5, s)$$

$$\int_{C_1} \frac{dz_5}{y} \propto \psi_0$$

holomorphic period of the elliptic curve!

$$J_1 = \frac{1}{\psi_0} I_{1,1,1,0,0}$$

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simple pole at  $\infty$

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