



Threshold Effects on the Massless Neutrino in the Canonical Seesaw Mechanism

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Based on DZ, JHEP 10 (2024) 002

Synergies Towards the Future Standard Model

DESY Theory Workshop

25 September 2025

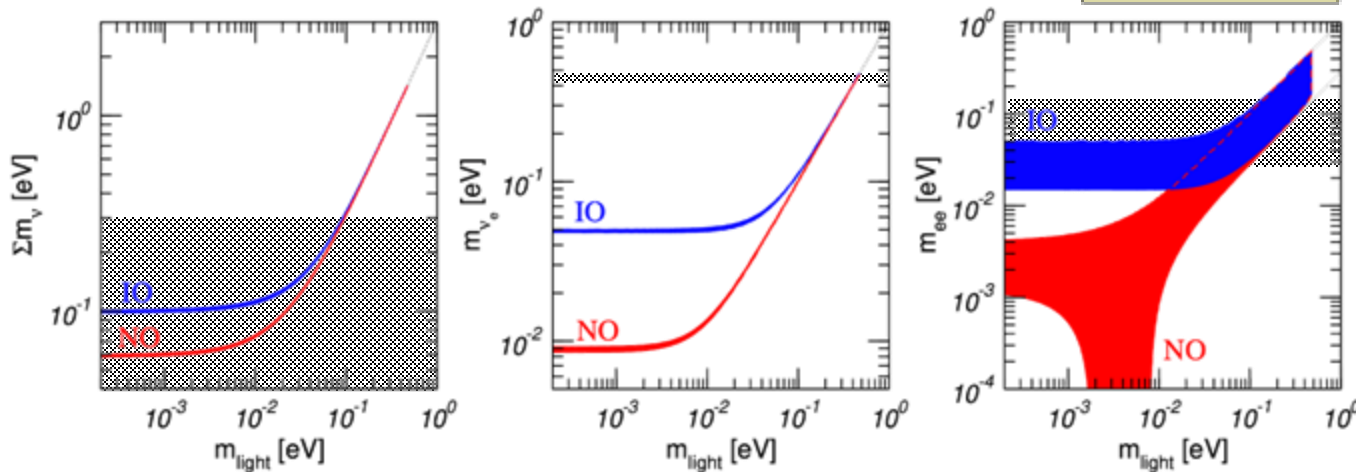
Status of Neutrino Masses

Oscillation

NuFIT 6.0 (2024)

	Normal Ordering ($\Delta\chi^2 = 0.6$)		Inverted Ordering (best fit)	
	bfp $\pm 1\sigma$	3σ range	bfp $\pm 1\sigma$	3σ range
$\frac{\Delta m_{21}^2}{10^{-6} \text{ eV}^2}$	$7.49^{+0.19}_{-0.19}$	$6.92 \rightarrow 8.05$	$7.49^{+0.19}_{-0.19}$	$6.92 \rightarrow 8.05$
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.513^{+0.021}_{-0.019}$	$+2.451 \rightarrow +2.578$	$-2.484^{+0.020}_{-0.020}$	$-2.547 \rightarrow -2.421$

Ivan Esteban et al., 2024



$$\sum m_\nu \lesssim 0.04 - 0.3 \text{ eV}$$

J. Q. Jiang et al., 2024

D. Naredo-Tuero et al., 2024

$$m_{\nu_e} < 0.45 \text{ eV (90\% CL)}$$

KATRIN collaboration, 2024

$$m_{ee} \lesssim 0.028 - 0.122 \text{ eV (90\% CL)}$$

KAMLAND-ZEN collaboration, 2024

$$m_{ee} \lesssim 0.079 - 0.180 \text{ eV (90\% CL)}$$

GERDA collaboration, 2020

Non-oscillation

No lower limit on the lightest neutrino mass allowing one massless neutrino

Neutrino Masses in Seesaw Mechanism

Can we have an exactly massless neutrino?

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Is the vanishing neutrino mass stable against quantum corrections if no extra symmetry protects it? If not, it provides theoretical lower limit

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The (minimal) type-I seesaw mechanism

P. Minkowski, 1977; T. Yanagida, 1979; M. Gell-Mann, P. Ramond, R. Slansky, 1979;
S. L. Glashow, 1980; R. N. Mohapatra, G. Senjanovic, 1980

Extending the SM with three (or two) right-handed neutrinos (RHNs):

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \overline{N_R} i \not{\partial} N_R - \left(\frac{1}{2} \overline{N_R^c} M_N N_R + \overline{\ell_L} Y_\nu \tilde{H} N_R + \text{h.c.} \right)$$



$$\langle H \rangle = v/\sqrt{2}$$

$$\mathcal{L}_\nu = -\frac{1}{2} \overline{\nu_L} M_\nu \nu_L^c + \text{h.c.} \quad \text{with} \quad M_\nu = -v^2 Y_\nu M_N^{-1} Y_\nu^T / 2$$

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One massless neutrino $\left\{ \begin{array}{l} \text{a) Two RHNs} \\ \text{b) Three RHNs and Rank-2 neutrino Yukawa matrix} \end{array} \right.$

Neutrino Masses in Seesaw Mechanism

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Is the vanishing neutrino mass stable against quantum corrections if no extra symmetry protects it? If not, it provides theoretical lower limit



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$$\langle H \rangle = v/\sqrt{2}$$

$$\mathcal{L}_\nu = -\frac{1}{2} \overline{\nu}_L M_\nu \nu_L^c + \text{h.c.} \quad \text{with} \quad M_\nu = -v^2 Y_\nu M_N^{-1} Y_\nu^T / 2$$

One massless neutrino $\left\{ \begin{array}{l} \text{a) Two RHNs} \\ \text{b) Three RHNs and Rank-2 neutrino Yukawa matrix} \end{array} \right.$

This vanishing neutrino mass is stable against **one-loop** Renormalization Group (RG) running effects **without threshold effects**

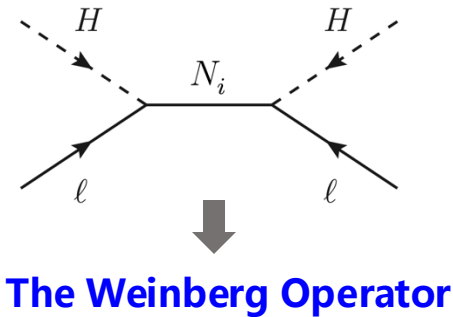
J. A. Casas, J. R. Espinosa, A. Ibarra, 2000; S. Antusch et al., 2001;2003;2005; J. W. Mei, Z. Z. Xing, 2004;
Z. Z. Xing, 2005; J. W. Mei, 2005; T. Ohlsson, H. Zhang, S. Zhou, 2013; T. Ohlsson, S. Zhou, 2014;...

But the case with hierarchical RHN masses is still unclear!!!

S. Antusch et al., 2002; N. J. Benoit et al, 2022

Threshold Effects in Seesaw Mechanism

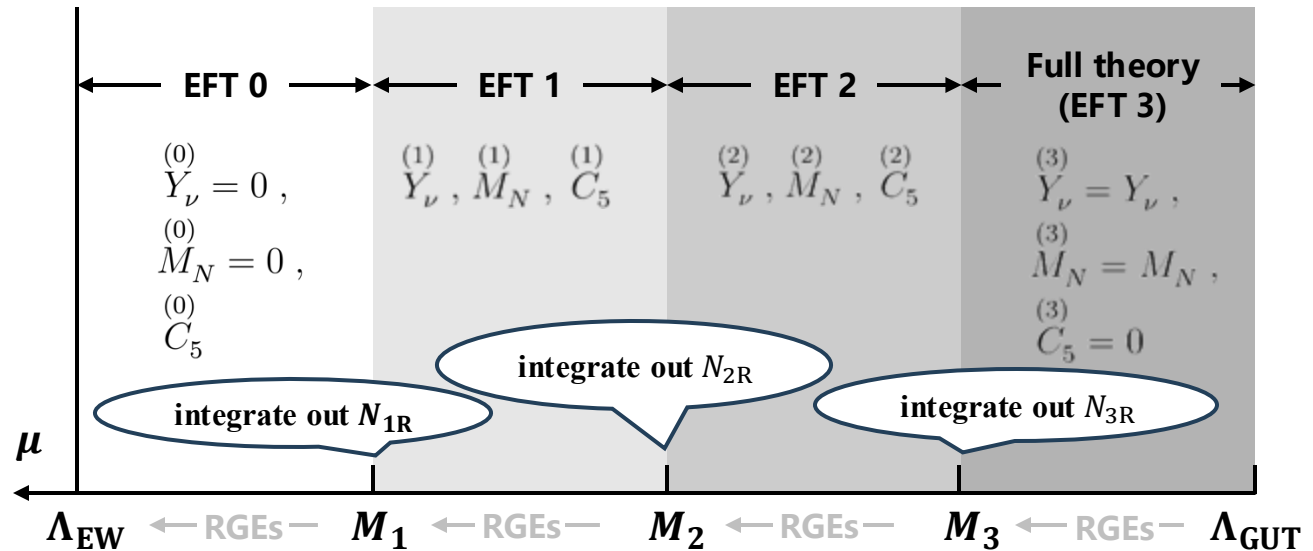
In the mass-independent renormalization scheme, e.g., the $\overline{\text{MS}}/\overline{\text{MS}}$ scheme:



$$\frac{1}{2} C_5^{\alpha\beta} \overline{\ell}_{\alpha L} \tilde{H} \tilde{H}^T \ell_{\beta L}^c$$

$$(C_5)_{\alpha\beta} = (Y_\nu)_{\alpha i} M_i^{-1} (Y_\nu)_{\beta i}$$

S. Weinberg, 1979



Matching conditions at each threshold scale $\mu = M_{n+1}$:

$$\triangleright (C_5)_{\alpha\beta} = (C_5)_{\alpha\beta}^{(n+1)} + (Y_\nu)_{\alpha(n+1)}^{(n+1)} M_{n+1}^{-1} (Y_\nu)_{\beta(n+1)}^{(n+1)}$$

$$\triangleright Y_\nu^{(n)} \text{ removing the last column of } Y_\nu^{(n+1)}$$

$$\triangleright M_N^{(n)} \text{ removing both the last column and the last row of } M_N^{(n+1)}$$

$$\begin{aligned} C_5^{(n_{\text{max}})} &= Y_\nu^{(0)} = M_N^{(0)} = 0 \\ Y_\nu^{(n_{\text{max}})} &= Y_\nu \quad M_N^{(n_{\text{max}})} = M_N \end{aligned}$$

One-loop RGEs and Threshold Effects

One-loop RGEs for non-degenerate seesaw scales:

S. Antusch et al., 2003; 2005

$$\mu \frac{dM_N^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[M_N^{(n)} Y_\nu^{(n)\dagger} Y_\nu^{(n)} + \left(Y_\nu^{(n)\dagger} Y_\nu^{(n)} \right)^T M_N^{(n)} \right] \quad T^{(n)} = \text{Tr} \left(Y_\nu^{(n)} Y_\nu^{(n)\dagger} + Y_l Y_l^\dagger + 3Y_u Y_u^\dagger + 3Y_d Y_d^\dagger \right)$$

$$\mu \frac{dY_\nu^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(T^{(n)} - \frac{3}{4}g_1^2 - \frac{9}{4}g_2^2 \right) Y_\nu^{(n)} + \frac{3}{2} \left(Y_\nu^{(n)} Y_\nu^{(n)\dagger} - Y_l Y_l^\dagger \right) Y_\nu^{(n)} \right]$$

$$\mu \frac{dC_5^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(4\lambda + 2T^{(n)} - 3g_2^2 \right) C_5^{(n)} + \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right) C_5^{(n)} + C_5^{(n)} \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right)^T \right]$$

The effect neutrino mass matrix:

$$M_\nu^{(n)} = -v^2 \kappa_\nu^{(n)} / 2$$

with $\kappa_\nu^{(n)} = C_5^{(n)} + \kappa_D^{(n)}$ and $\kappa_D^{(n)} = Y_\nu^{(n)} M_N^{-1} Y_\nu^{(n)T}$

$$\mu \frac{d\kappa_D^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_D^{(n)} + \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right) \kappa_D^{(n)} + \kappa_D^{(n)} \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right)^T \right]$$

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$$T^{(n)} = \text{Tr} \left(Y_\nu^{(n)} Y_\nu^{(n)\dagger} + Y_l Y_l^\dagger + 3Y_u Y_u^\dagger + 3Y_d Y_d^\dagger \right)$$

$$\mu \frac{dY_\nu^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(T^{(n)} - \frac{3}{4}g_1^2 - \frac{9}{4}g_2^2 \right) Y_\nu^{(n)} + \frac{3}{2} \left(Y_\nu^{(n)} Y_\nu^{(n)\dagger} - Y_l Y_l^\dagger \right) Y_\nu^{(n)} \right]$$

Additional terms

$$\mu \frac{dC_5^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(4\lambda + 2T^{(n)} - 3g_2^2 \right) C_5^{(n)} + \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right) C_5^{(n)} + C_5^{(n)} \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right)^T \right]$$

The effect neutrino mass matrix:

$$M_\nu^{(n)} = -v^2 \kappa_\nu^{(n)} / 2$$

with $\kappa_\nu^{(n)} = C_5^{(n)} + \kappa_D^{(n)}$ and $\kappa_D^{(n)} = Y_\nu^{(n)} M_N^{(n)-1} Y_\nu^{(n)T}$

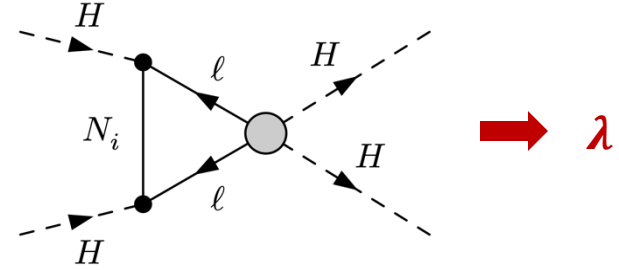
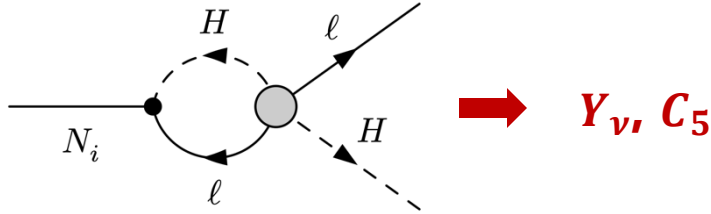
$$\mu \frac{d\kappa_D^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_D^{(n)} + \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right) \kappa_D^{(n)} + \kappa_D^{(n)} \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right)^T \right]$$

$$\mu \frac{d\kappa_\nu^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_\nu^{(n)} + \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right) \kappa_\nu^{(n)} + \kappa_\nu^{(n)} \left(\frac{1}{2} Y_\nu^{(n)} Y_\nu^{(n)\dagger} - \frac{3}{2} Y_l Y_l^\dagger \right)^T + \left(4\lambda + \frac{3}{2}g_1^2 + \frac{3}{2}g_2^2 \right) C_5^{(n)} \right]$$

Revisiting the One-loop RGEs

Overlooked UV divergent 1PI diagrams:

DZ, 2024



$$R_{\ell H N 1}^{\alpha \beta} = \overline{\ell}_{\alpha L} \tilde{H} \not{D} N_{\beta R}^c \quad \delta G_{\ell H N 1} = \frac{3i C_5 Y_\nu^*}{32\pi^2 \epsilon}$$

$$\delta G_{\ell H N 1}^{\alpha \beta} \mathcal{R}_{\ell H N 1}^{\alpha \beta} = -i (\delta G_{\ell H N 1} M_N)_{\alpha \beta} \overline{\ell}_{\alpha L} \tilde{H} N_{\beta R} - i \left(\delta G_{\ell H N 1} Y_\nu^T \right)_{\alpha \beta} \overline{\ell}_{\alpha L} \tilde{H} \tilde{H}^T \ell_{\beta L}^c$$

New contributions to the RGEs among seesaw scales:

$$\mu \frac{d\lambda}{d\mu} = \frac{1}{16\pi^2} \left[4\lambda \overset{(n)}{T} + \frac{3}{8} (g_1^2 + g_2^2)^2 + \frac{3}{4} g_2^4 - 3\lambda (g_1^2 + 3g_2^2) + 24\lambda^2 - 2\overset{(n)}{T}' \right. \\ \left. + 2\text{Tr} \left(\overset{(n)}{C}_5 \overset{(n)}{Y}_\nu^* \overset{(n)}{M}_N \overset{(n)}{Y}_\nu^\dagger + \overset{(n)}{Y}_\nu \overset{(n)}{M}_N^\dagger \overset{(n)}{Y}_\nu^T \overset{(n)}{C}_5^\dagger \right) \right],$$

$$\mu \frac{dY_\nu}{d\mu} = \frac{1}{16\pi^2} \left[\left(\overset{(n)}{T} - \frac{3}{4} g_1^2 - \frac{9}{4} g_2^2 \right) \overset{(n)}{Y}_\nu + \frac{3}{2} \left(\overset{(n)}{Y}_\nu \overset{(n)}{Y}_\nu^\dagger - Y_l Y_l^\dagger \right) \overset{(n)}{Y}_\nu - 3 \overset{(n)}{C}_5 \overset{(n)}{Y}_\nu^* \overset{(n)}{M}_N \right],$$

$$\mu \frac{dC_5}{d\mu} = \frac{1}{16\pi^2} \left[\left(4\lambda + 2\overset{(n)}{T} - 3g_2^2 \right) \overset{(n)}{C}_5 + \left(\frac{7}{2} \overset{(n)}{Y}_\nu \overset{(n)}{Y}_\nu^\dagger - \frac{3}{2} Y_l Y_l^\dagger \right) \overset{(n)}{C}_5 + \overset{(n)}{C}_5 \left(\frac{7}{2} \overset{(n)}{Y}_\nu \overset{(n)}{Y}_\nu^\dagger - \frac{3}{2} Y_l Y_l^\dagger \right)^T \right]$$

Revisiting the One-loop RGEs

$$\mu \frac{d\kappa_D}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_D + \left(\frac{1}{2}Y_\nu Y_\nu^\dagger - \frac{3}{2}Y_l Y_l^\dagger \right) \kappa_D + \kappa_D \left(\frac{1}{2}Y_\nu Y_\nu^\dagger - \frac{3}{2}Y_l Y_l^\dagger \right)^T \right. \\ \left. - 3Y_\nu Y_\nu^\dagger C_5 - 3C_5 \left(Y_\nu Y_\nu^\dagger \right)^T \right],$$

$$\begin{aligned} \kappa_\nu &= C_5 + \kappa_D \\ \kappa_D &= Y_\nu M_N^{-1} Y_\nu^T \end{aligned}$$

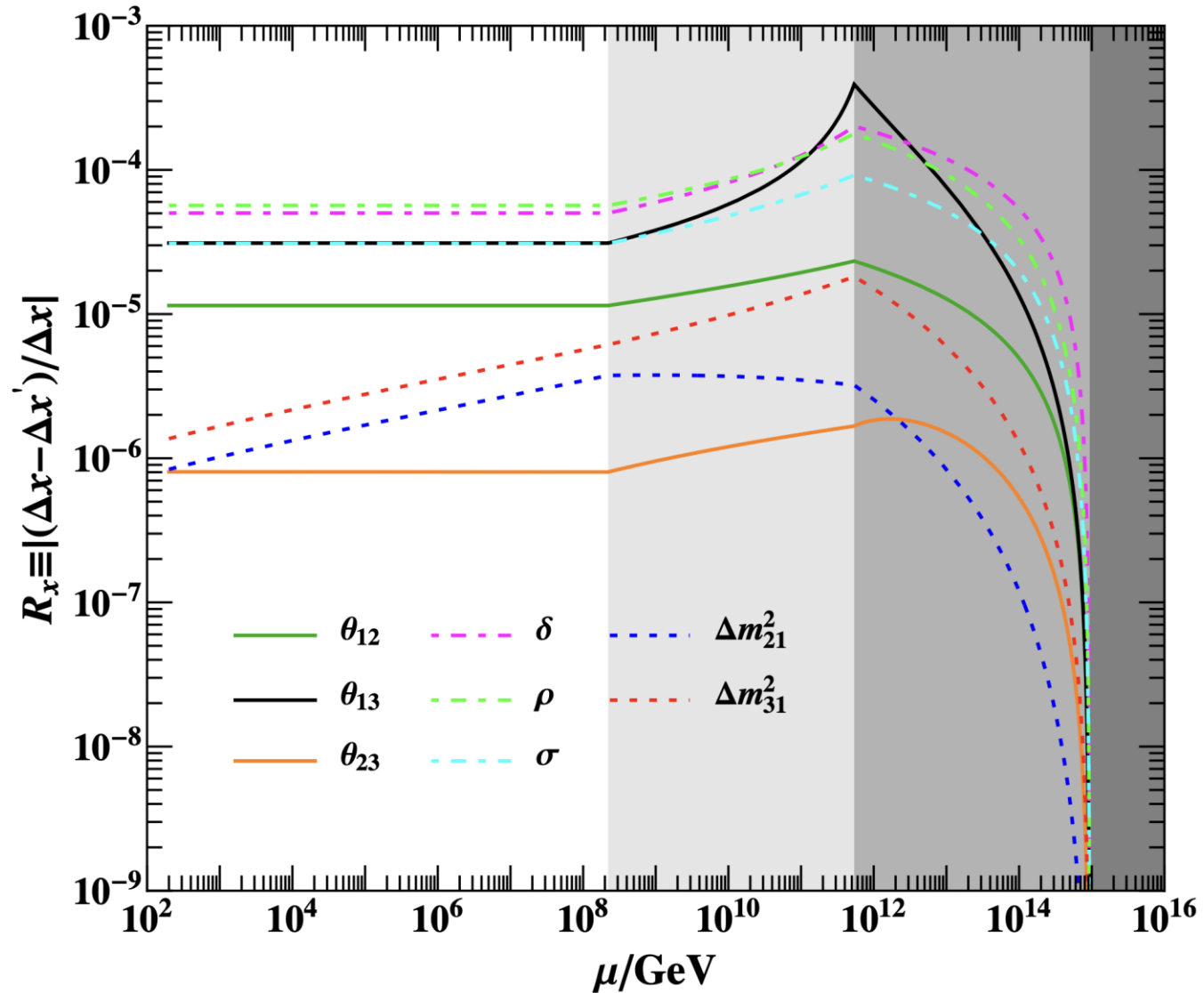
$$\mu \frac{dC_5}{d\mu} = \frac{1}{16\pi^2} \left[\left(4\lambda + 2T - 3g_2^2 \right) C_5 + \left(\frac{7}{2}Y_\nu Y_\nu^\dagger - \frac{3}{2}Y_l Y_l^\dagger \right) C_5 + C_5 \left(\frac{7}{2}Y_\nu Y_\nu^\dagger - \frac{3}{2}Y_l Y_l^\dagger \right)^T \right]$$

$$\mu \frac{d\kappa_\nu}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_\nu + \left(\frac{1}{2}Y_\nu Y_\nu^\dagger - \frac{3}{2}Y_l Y_l^\dagger \right) \kappa_\nu + \kappa_\nu \left(\frac{1}{2}Y_\nu Y_\nu^\dagger - \frac{3}{2}Y_l Y_l^\dagger \right)^T \right. \\ \left. + \left(4\lambda + \frac{3}{2}g_1^2 + \frac{3}{2}g_2^2 \right) C_5 \right].$$

- The counteraction of new contributions in κ_ν leaves κ_ν in the same form as before
- New contributions affect the running behaviors of λ , κ_D , and C_5
- The running behavior of κ_ν is indirectly influenced by those of λ , κ_D , and C_5

Revisiting the One-loop RGEs

Relative differences between the results with/without new contributions



Threshold Effects on the Massless Neutrino

$$\mu \frac{d\kappa_\nu^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_\nu^{(n)} + \left(\frac{1}{2}Y_\nu^{(n)}Y_\nu^{\dagger(n)} - \frac{3}{2}Y_l^{(n)}Y_l^{\dagger(n)} \right) \kappa_\nu^{(n)} + \kappa_\nu^{(n)} \left(\frac{1}{2}Y_\nu^{(n)}Y_\nu^{\dagger(n)} - \frac{3}{2}Y_l^{(n)}Y_l^{\dagger(n)} \right)^T \right. \\ \left. + \left(4\lambda + \frac{3}{2}g_1^2 + \frac{3}{2}g_2^2 \right) C_5^{(n)} \right] .$$

□ Above the highest or below the lowest seesaw scale:

$$\frac{d\kappa_\nu}{dt} = \alpha\kappa_\nu + \beta\kappa_\nu + \kappa_\nu\beta^T \quad U^\dagger\kappa_\nu U^* = \text{Diag}\{\chi_1, \chi_2, \chi_3\} \quad \tilde{\beta} = U^\dagger\beta U$$

$$\frac{d\chi_i}{dt} = \left(\alpha + 2\tilde{\beta}_{ii} \right) \chi_i \quad \chi_i(t) = \chi_i(t_0) \cdot \exp \left[\int_{t_0}^t \left(\alpha + 2\tilde{\beta}_{ii} \right) dt \right] \quad t \equiv \ln \mu$$

Neutrino masses are proportional to their initial values, therefore, they can not be generated if they are initially zero

□ Among seesaw scales (i.e., including threshold effects):

$$\frac{d\kappa_\nu}{dt} = \alpha\kappa_\nu + \beta\kappa_\nu + \kappa_\nu\beta^T + \alpha' C_5$$

It becomes much more complicated due to the coexistence of κ_ν and C_5

Threshold Effects on the Massless Neutrino

$$\mu \frac{d\kappa_\nu^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_\nu^{(n)} + \left(\frac{1}{2}Y_\nu^{(n)}Y_\nu^{\dagger(n)} - \frac{3}{2}Y_l Y_l^\dagger \right) \kappa_\nu^{(n)} + \kappa_\nu^{(n)} \left(\frac{1}{2}Y_\nu^{(n)}Y_\nu^{\dagger(n)} - \frac{3}{2}Y_l Y_l^\dagger \right)^T \right. \\ \left. + \left(4\lambda + \frac{3}{2}g_1^2 + \frac{3}{2}g_2^2 \right) C_5^{(n)} \right].$$

$$\text{Re}[\text{Det}(\kappa_\nu^{(n)})] = \chi_1^{(n)} \chi_2^{(n)} \chi_3^{(n)}$$



$$\frac{d\text{Det}(M)}{dx} = \text{Tr} \left(\text{Adj}(M) \frac{dM}{dx} \right)$$

$$M \text{Adj}(M) = \text{Adj}(M)M = \text{Det}(M) \mathbf{1}$$

$$\mu \frac{d\text{Det}(\kappa_\nu^{(n)})}{d\mu} = \frac{1}{16\pi^2} \left[\left(3\alpha_1^{(n)} + 2\beta^{(n)} \right) \text{Det}(\kappa_\nu^{(n)}) + \left(\alpha_2^{(n)} - \alpha_1^{(n)} \right) \text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)}) C_5^{(n)} \right) \right]$$

$$\alpha_1^{(n)} = 2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2, \quad \alpha_2^{(n)} = 4\lambda + 2T^{(n)} - 3g_2^2, \quad \beta^{(n)} = \frac{1}{2} \text{Tr} \left(Y_\nu^{(n)} Y_\nu^{\dagger(n)} - 3Y_l Y_l^\dagger \right)$$

Above the highest seesaw scale:

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n_{\max})}) C_5^{(n_{\max})} \right) = 0$$

Below the lowest seesaw scale:

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(0)}) C_5^{(0)} \right) = 3\text{Det}(\kappa_\nu^{(0)})$$

Among seesaw scales:

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)}) C_5^{(n)} \right) \quad ? \quad \text{Det}(\kappa_\nu^{(n)})$$

Threshold Effects on the Massless Neutrino

$$\mu \frac{d\kappa_\nu^{(n)}}{d\mu} = \frac{1}{16\pi^2} \left[\left(2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2 \right) \kappa_\nu^{(n)} + \left(\frac{1}{2}Y_\nu^{(n)}Y_\nu^{\dagger(n)} - \frac{3}{2}Y_l Y_l^\dagger \right) \kappa_\nu^{(n)} + \kappa_\nu^{(n)} \left(\frac{1}{2}Y_\nu^{(n)}Y_\nu^{\dagger(n)} - \frac{3}{2}Y_l Y_l^\dagger \right)^T \right. \\ \left. + \left(4\lambda + \frac{3}{2}g_1^2 + \frac{3}{2}g_2^2 \right) C_5^{(n)} \right].$$

$$\text{Re}[\text{Det}(\kappa_\nu^{(n)})] = \chi_1^{(n)} \chi_2^{(n)} \chi_3^{(n)}$$



$$\frac{d\text{Det}(M)}{dx} = \text{Tr} \left(\text{Adj}(M) \frac{dM}{dx} \right)$$

$$M \text{Adj}(M) = \text{Adj}(M)M = \text{Det}(M) \mathbf{1}$$

$$\mu \frac{d\text{Det}(\kappa_\nu^{(n)})}{d\mu} = \frac{1}{16\pi^2} \left[\left(3\alpha_1^{(n)} + 2\beta^{(n)} \right) \text{Det}(\kappa_\nu^{(n)}) + \left(\alpha_2^{(n)} - \alpha_1^{(n)} \right) \text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)}) C_5^{(n)} \right) \right]$$

$$\alpha_1^{(n)} = 2T^{(n)} - \frac{3}{2}g_1^2 - \frac{9}{2}g_2^2, \quad \alpha_2^{(n)} = 4\lambda + 2T^{(n)} - 3g_2^2, \quad \beta^{(n)} = \frac{1}{2} \text{Tr} \left(Y_\nu^{(n)} Y_\nu^{\dagger(n)} - 3Y_l Y_l^\dagger \right)$$

Above the highest seesaw scale:

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n_{\max})}) C_5^{(n_{\max})} \right) = 0$$

Below the lowest seesaw scale:

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(0)}) C_5^{(0)} \right) = 3\text{Det}(\kappa_\nu^{(0)})$$

**See backup or
DZ, 2024 for
more details**

Among seesaw scales:

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)}) C_5^{(n)} \right) = (3 - n) \text{Det}(\kappa_\nu^{(n)})$$

Threshold Effects on the Massless Neutrino

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)}) C_5^{(n)} \right) = (3 - n) \text{Det}(\kappa_\nu^{(n)}) \quad \rightarrow \quad \mu \frac{d\text{Det}(\kappa_\nu^{(n)})}{d\mu} = \frac{1}{16\pi^2} \left[3\alpha_1^{(n)} + (3 - n) \left(\alpha_2^{(n)} - \alpha_1^{(n)} \right) + 2\beta^{(n)} \right] \text{Det}(\kappa_\nu^{(n)})$$

Solution for the determinant:

$$\text{Det}(\kappa_\nu(t)) = I_{\text{Int}}(t) \cdot \text{Det}(\kappa_\nu(t_4))$$

$$I_{\text{Int}}(t) = \begin{cases} I_{\text{Int}}^{(3)}(t, t_4) , & t_3 < t \leq t_4 \\ I_{\text{Int}}^{(2)}(t, t_3) I_{\text{Int}}^{(3)}(t_3, t_4) , & t_2 < t \leq t_3 \\ I_{\text{Int}}^{(1)}(t, t_2) I_{\text{Int}}^{(2)}(t_2, t_3) I_{\text{Int}}^{(3)}(t_3, t_4) , & t_1 < t \leq t_2 \\ I_{\text{Int}}^{(0)}(t, t_1) I_{\text{Int}}^{(1)}(t_1, t_2) I_{\text{Int}}^{(2)}(t_2, t_3) I_{\text{Int}}^{(3)}(t_3, t_4) , & t \leq t_1 \end{cases}$$

$$t_{(n)} \equiv \ln(\mu_{(n)})$$

$$\mu_n = M_n$$

$$\mu_4 = \Lambda_{\text{GUT}}$$

$$I_{\text{Int}}^{(n)}(t, t_{n+1}) = \exp \left\{ \frac{1}{16\pi^2} \int_{t_{n+1}}^t \left[3\alpha_1^{(n)} + (3 - n) \left(\alpha_2^{(n)} - \alpha_1^{(n)} \right) + 2\beta^{(n)} \right] dt \right\}$$

**Massless neutrino
at the GUT scale**



**Zero determinant
at the GUT scale**



**Zero determinant
at the EW scale**

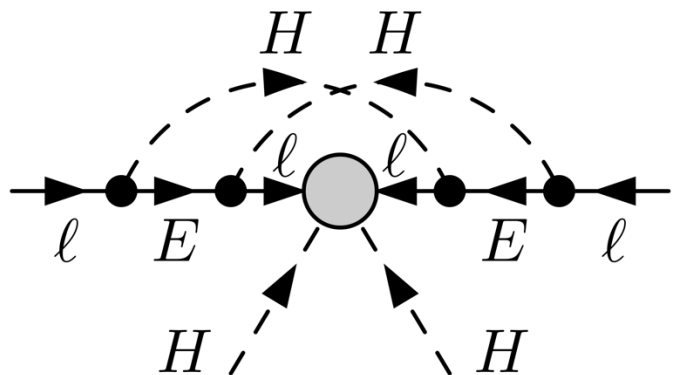


**Massless at
the EW scale**

The vanishing neutrino mass is stable against **one-loop** RG running effects
EVEN with **threshold effects**

Two-loop RG Running Effects

Rank-increase diagram for renormalization of the Weinberg operator:



$$\frac{dC_5}{dt} \ni \frac{1}{(16\pi^2)^2} \cdot 2Y_l Y_l^\dagger C_5 (Y_l Y_l^\dagger)^T$$

$$m_1(\Lambda_{\text{EW}}) \sim \frac{2y_\tau^4}{(16\pi^2)^2} \ln\left(\frac{\Lambda}{\Lambda_{\text{EW}}}\right) \cdot \sum_{i=2,3} m_i \text{Re}(U_{\tau 1}^* U_{\tau i})^2 \quad \textbf{(NMO)}$$

$$m_3(\Lambda_{\text{EW}}) \sim \frac{2y_\tau^4}{(16\pi^2)^2} \ln\left(\frac{\Lambda}{\Lambda_{\text{EW}}}\right) \cdot \sum_{i=1,2} m_i \text{Re}(U_{\tau 3}^* U_{\tau i})^2 \quad \textbf{(IMO)}$$

$m_{\text{lightest}} \sim 10^{-13} \text{ eV}$

S. Davidson, G. Isidori, A. Strumia, 2007; Z. Z. Xing, DZ, 2020; A. Ibarra, N. Leister, DZ, 2024

Conclusions

- We revisit the one-loop RGEs in the canonical seesaw mechanism among seesaw scales and obtain the missed terms
- We achieve the running behavior of the determinant of the effective neutrino mass matrix against the energy scale
- We strictly prove that if the lightest neutrino is initially massless, it remains massless at the one-loop level even threshold effects are taken into account.
- Nevertheless, two-loop RG running effects can generate a non-zero mass for the initially massless neutrino

THANKS FOR YOUR ATTENTION

Backup

For 3x3 matrix M , its adjugate can be written as

$$\text{Adj}(M) = \frac{1}{2} [(\text{Tr}(M))^2 - \text{Tr}(M^2)] \mathbf{1} - M\text{Tr}(M) + M^2$$



$$\begin{aligned} \text{Adj}(M + W) &= \text{Adj}(M) + \text{Adj}(W) - M\text{Tr}(W) - W\text{Tr}(M) + MW + WM \\ &\quad + [\text{Tr}(M)\text{Tr}(W) - \text{Tr}(MW)] \mathbf{1} . \end{aligned}$$



$$\text{Tr}(\text{Adj}(M + W)W) = 3\text{Det}(W) + 2\text{Tr}(\text{Adj}(W)M) + \text{Tr}(\text{Adj}(M)W)$$



$$\begin{aligned} \text{Tr}(\text{Adj}(\kappa_\nu^{(n)})C_5^{(n)}) &= 3\text{Det}(C_5^{(n)}) + 2\text{Tr}(\text{Adj}(C_5^{(n)})\kappa_D^{(n)}) + \text{Tr}(\text{Adj}(\kappa_D^{(n)})C_5^{(n)}) \\ \text{Tr}(\text{Adj}(\kappa_\nu^{(n)})\kappa_D^{(n)}) &= 3\text{Det}(\kappa_D^{(n)}) + 2\text{Tr}(\text{Adj}(\kappa_D^{(n)})C_5^{(n)}) + \text{Tr}(\text{Adj}(C_5^{(n)})\kappa_D^{(n)}) \end{aligned}$$



$$\text{Det}(\kappa_\nu^{(n)}) = \text{Det}(C_5^{(n)}) + \text{Det}(\kappa_D^{(n)}) + \text{Tr}(\text{Adj}(C_5^{(n)})\kappa_D^{(n)}) + \text{Tr}(\text{Adj}(\kappa_D^{(n)})C_5^{(n)})$$




$$\begin{aligned} \text{Tr}[\text{Adj}(\kappa_\nu^{(n)})C_5^{(n)}] &= \text{Det}(\kappa_\nu^{(n)}) + 2\text{Det}(C_5^{(n)}) - \text{Det}(\kappa_D^{(n)}) + \text{Tr}[\text{Adj}(C_5^{(n)})\kappa_D^{(n)}] \\ \text{Tr}[\text{Adj}(\kappa_\nu^{(n)})\kappa_D^{(n)}] &= 2\text{Det}(\kappa_\nu^{(n)}) + \text{Det}(C_5^{(n)}) - 2\text{Det}(\kappa_D^{(n)}) - \text{Tr}[\text{Adj}(\kappa_D^{(n)})C_5^{(n)}] \end{aligned}$$

Backup

Some essential points used for discussion:

- For 3-by-3 matrices, one can prove

$$\text{Tr}(\text{Adj}(M + W)W) = 3\text{Det}(W) + 2\text{Tr}(\text{Adj}(W)M) + \text{Tr}(\text{Adj}(M)W) \quad \text{Linear algebra relation}$$



$$\begin{aligned} \text{Tr} \left(\text{Adj}(\kappa_\nu)^{(n)} C_5^{(n)} \right) &= \text{Det}(\kappa_\nu)^{(n)} + 2\text{Det}(C_5)^{(n)} - \text{Det}(\kappa_D)^{(n)} + \text{Tr} \left(\text{Adj}(C_5)^{(n)} \kappa_D^{(n)} \right) \\ \text{Tr} \left(\text{Adj}(\kappa_\nu)^{(n)} C_5^{(n)} \right) &= 2\text{Det}(\kappa_\nu)^{(n)} + \text{Det}(C_5)^{(n)} - 2\text{Det}(\kappa_D)^{(n)} - \text{Tr} \left(\text{Adj}(\kappa_D)^{(n)} C_5^{(n)} \right) \end{aligned}$$

Independent of
the energy scale

- The differential equation of C_5 has the form $\mu \frac{dX}{d\mu} = \alpha X + \beta X + X\beta^T$, leading to that its rank will not be changed by running effects
- The rank of $\kappa_D^{(n)} = Y_\nu^{(n)} M_N^{-1} Y_\nu^{(n)T}$ is determined by those of Y_ν and M_N , i.e.,

$$\text{Rank}(\kappa_D^{(n)}) = \min\{\text{Rank}(Y_\nu^{(n)}), \text{Rank}(M_N^{(n)})\}$$
- If $\text{Rank}(M) \leq n - 2$ with n being the dimension of M , then $\text{Adj}(M) = 0$
- Subadditivity: $\text{Rank}(M+W) \leq \text{Rank}(M) + \text{Rank}(W)$

Backup

■ Two right-handed neutrinos: trivial

$$\begin{array}{l}
 \text{Rank}^{(2)}(Y_\nu(\mu)) = 2 \quad \Rightarrow \quad \text{Det}^{(2)}(\kappa_\nu(\mu)) = 0 \\
 \\
 \mathbf{M}_2 \quad (C_5)_{\alpha\beta}^{(1)} = (Y_\nu)_{\alpha 2}^{(2)} (Y_\nu)_{\beta 2}^{(2)} / M_2 \quad \Rightarrow \quad \text{Rank}^{(1)}(C_5(M_2)) = 1 \\
 \\
 \text{Rank}^{(1)}(C_5(\mu)) = 1 \quad \kappa_D^{(1)}(\mu) = Y_\nu^{(1)}(\mu) Y_\nu^{(1)\text{T}}(\mu) / M_1(\mu) \quad \Rightarrow \quad \text{Rank}^{(1)}(\kappa_D(\mu)) = 1 \\
 \\
 \text{Rank}^{(1)}(\kappa_\nu(\mu)) \leq \text{Rank}^{(1)}(C_5(\mu)) + \text{Rank}^{(1)}(\kappa_D(\mu)) = 2 \quad \Rightarrow \quad \text{Det}^{(1)}(\kappa_\nu(\mu)) = 0 \\
 \\
 \mathbf{M}_1 \quad \kappa_\nu^{(0)}(M_1) = \kappa_\nu^{(1)}(M_1) \quad \Rightarrow \quad \text{Det}^{(0)}(\kappa_\nu(M_1)) = 0 \\
 \\
 \text{Det}^{(0)}(\kappa_\nu(\mu)) = 0
 \end{array}$$

The determinant of the effect neutrino mass matrix remains vanishing against RG running effects



Massless neutrino

Backup

■ Three right-handed neutrinos:

M_3

$$\text{Rank}(\kappa_\nu^{(3)}(\mu)) = \text{Rank}(Y_\nu^{(3)}(\mu)) \leq 3$$

$$(C_5^{(2)})_{\alpha\beta} = (Y_\nu^{(3)})_{\alpha 3} (Y_\nu^{(3)})_{\beta 3} / M_3 \quad \Rightarrow \quad \text{Rank}(C_5^{(2)}(M_3)) = 1$$

$$\text{Rank}(C_5^{(2)}(\mu)) = \text{Rank}(C_5^{(2)}(M_3)) = 1 \quad \text{Rank}(\kappa_D^{(2)}(\mu)) = \text{Rank}(Y_\nu^{(2)}(\mu)) \leq 2$$

$$\text{Det}(C_5^{(2)}(\mu)) = \text{Det}(\kappa_D^{(2)}(\mu)) = 0 \quad \text{Adj}(C_5^{(2)}(\mu)) = 0$$

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(2)}(\mu)) C_5^{(2)}(\mu) \right) = \text{Det}(\kappa_\nu^{(2)}(\mu))$$

M_2

$$(C_5^{(1)})_{\alpha\beta} = (C_5^{(2)})_{\alpha\beta} + (Y_\nu^{(2)})_{\alpha 2} (Y_\nu^{(2)})_{\beta 2} / M_2 \quad \Rightarrow \quad \text{Rank}(C_5^{(1)}(M_2)) \leq \text{Rank}(C_5^{(2)}(M_2)) + \text{Rank}(\mathbf{y}_2^{(2)}(M_2) \mathbf{y}_2^{(2)\text{T}}(M_2) / M_2) = 2$$

$$\text{Rank}(C_5^{(1)}(\mu)) = \text{Rank}(C_5^{(1)}(M_2)) \leq 2 \quad \Rightarrow \quad \text{Det}(C_5^{(1)}(\mu)) = 0$$

$$\text{Rank}(\kappa_D^{(1)}(\mu)) = \text{Rank}(Y_\nu^{(1)}(\mu)) = 1 \quad \Rightarrow \quad \text{Det}(\kappa_D^{(1)}(\mu)) = \text{Adj}(\kappa_D^{(1)}(\mu)) = 0$$

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(1)}(\mu)) C_5^{(1)}(\mu) \right) = 2 \text{Det}(\kappa_\nu^{(1)}(\mu))$$

M_1

$$\kappa_\nu^{(0)}(M_1) = C_5^{(0)}(M_1) = \kappa_\nu^{(1)}(M_1)$$

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(0)}(\mu)) C_5^{(0)}(\mu) \right) = 3 \text{Det}(\kappa_\nu^{(0)}(\mu))$$

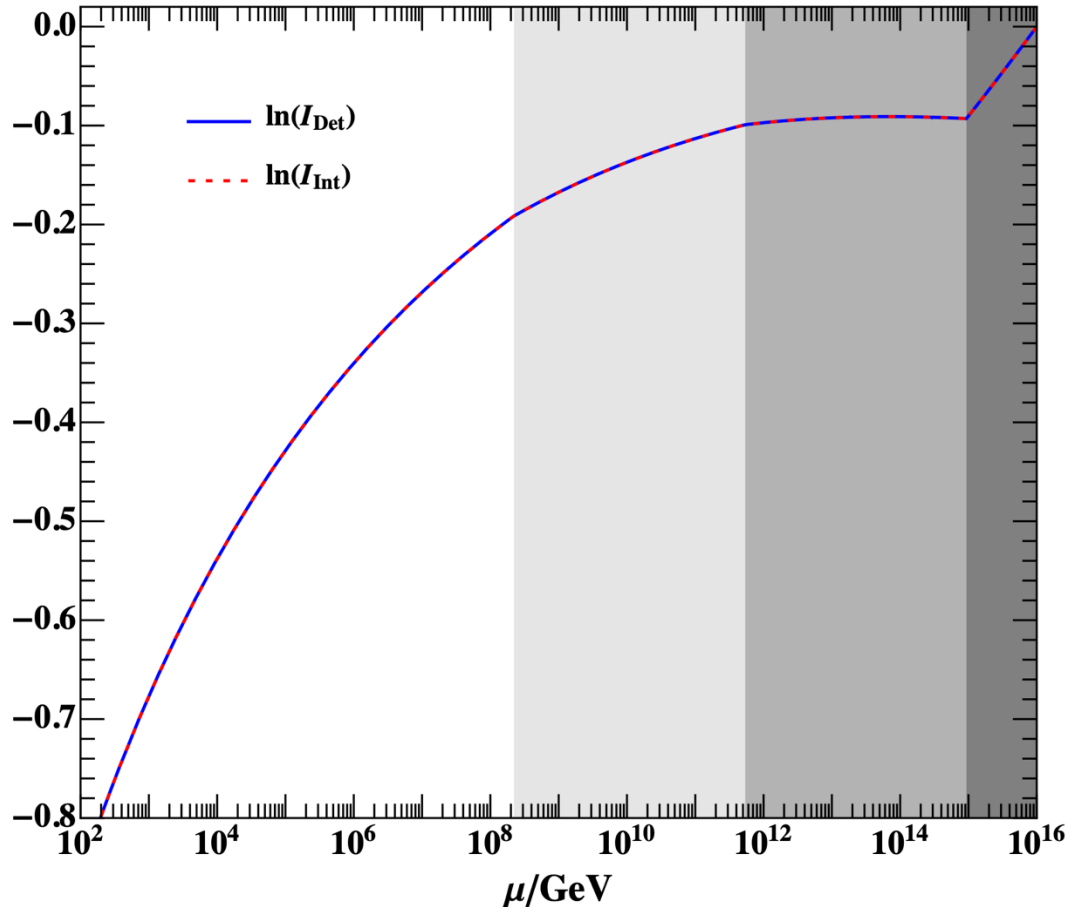
$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)} C_5^{(n)}) \right) = \text{Det}(\kappa_\nu^{(n)}) + 2 \text{Det}(C_5^{(n)}) - \text{Det}(\kappa_D^{(n)}) + \text{Tr} \left(\text{Adj}(C_5^{(n)}) \kappa_D^{(n)} \right)$$

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)} C_5^{(n)}) \right) = 2 \text{Det}(\kappa_\nu^{(n)}) + \text{Det}(C_5^{(n)}) - 2 \text{Det}(\kappa_D^{(n)}) - \text{Tr} \left(\text{Adj}(\kappa_D^{(n)}) C_5^{(n)} \right)$$

$$\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)} C_5^{(n)}) \right) = (3 - n) \text{Det}(\kappa_\nu^{(n)})$$

Backup

Running of the determinant against μ : $\text{Det}(\kappa_\nu(t)) = I_{\text{Int}}(t) \cdot \text{Det}(\kappa_\nu(t_4))$



$I_{\text{Det}}(t) = \text{Det}(\kappa_\nu(t))/\text{Det}(\kappa_\nu(t_4))$ is calculated directly from κ_ν

$I_{\text{Int}}(t)$ is calculated from its expression with Higgs, gauge and Yukawa couplings

Backup

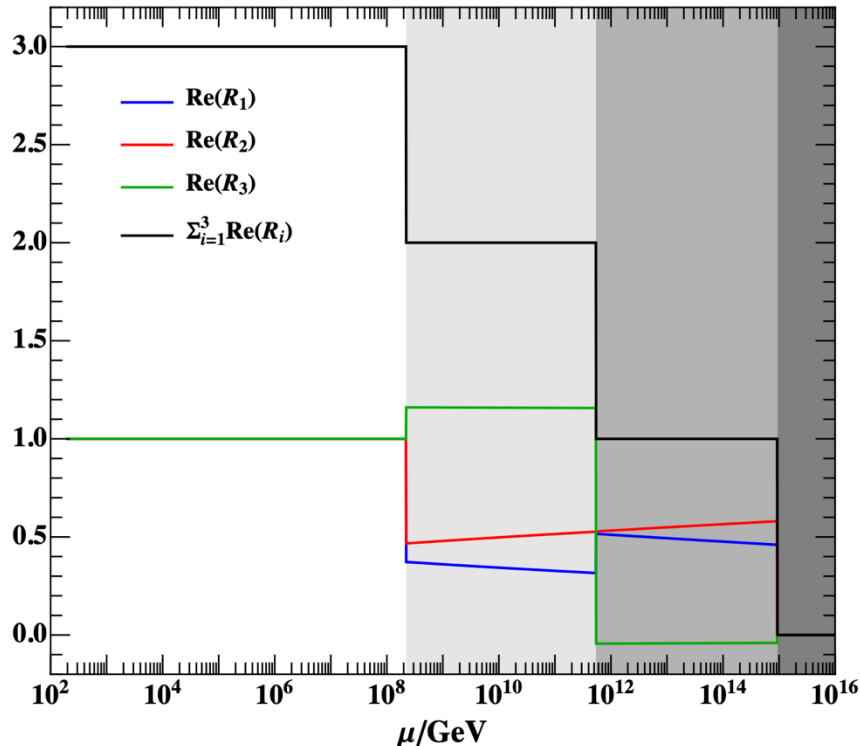
A little more discussion on $\text{Tr} \left(\text{Adj}(\kappa_\nu^{(n)}) C_5^{(n)} \right) = (3 - n) \text{Det}(\kappa_\nu^{(n)})$

$$(3 - n) \chi_1 \chi_2 \chi_3 = \chi_2 \chi_3 (\mathcal{C}_5)_{11} + \chi_1 \chi_3 (\mathcal{C}_5)_{22} + \chi_1 \chi_2 (\mathcal{C}_5)_{33}$$

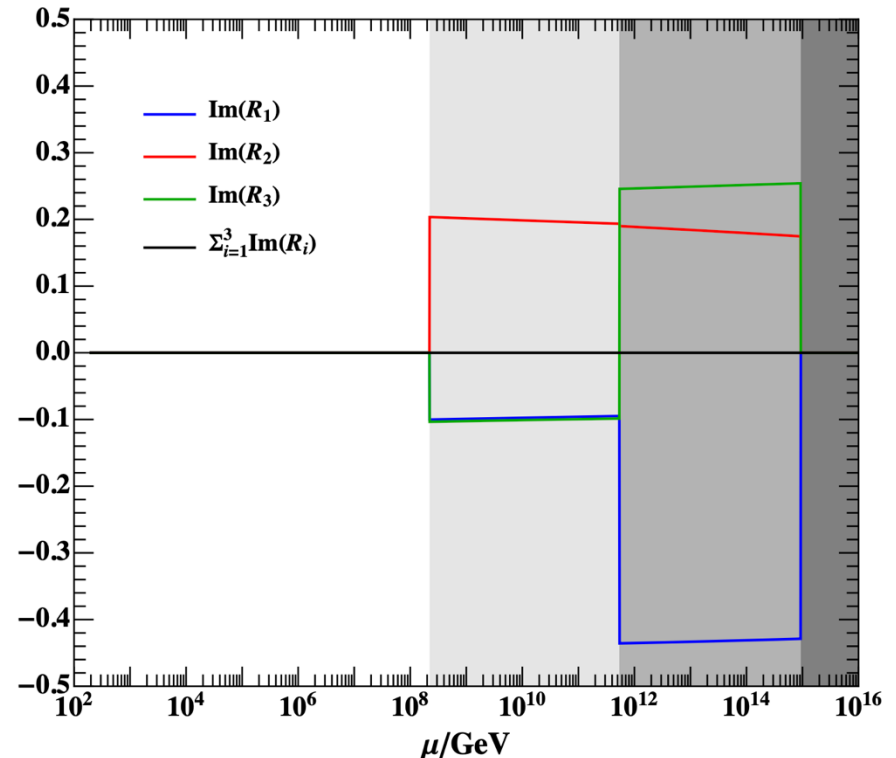
$$U_\nu^\dagger \kappa_\nu U_\nu^* = \text{Diag}\{\chi_1, \chi_2, \chi_3\}$$

$$\mathcal{C}_5 \equiv U_\nu^\dagger C_5 U_\nu^*$$

Running of $R_i \equiv (\mathcal{C}_5)_{ii} / \chi_i$ against the energy scale:



$$\sum_i \text{Re}[(\mathcal{C}_5)_{ii} / \chi_i] = 3 - n$$



$$\sum_i \text{Im}[(\mathcal{C}_5)_{ii} / \chi_i] = 0$$