

Electroweak phase transition seeded by DFSZ axion string

Yu Hamada (DESY)

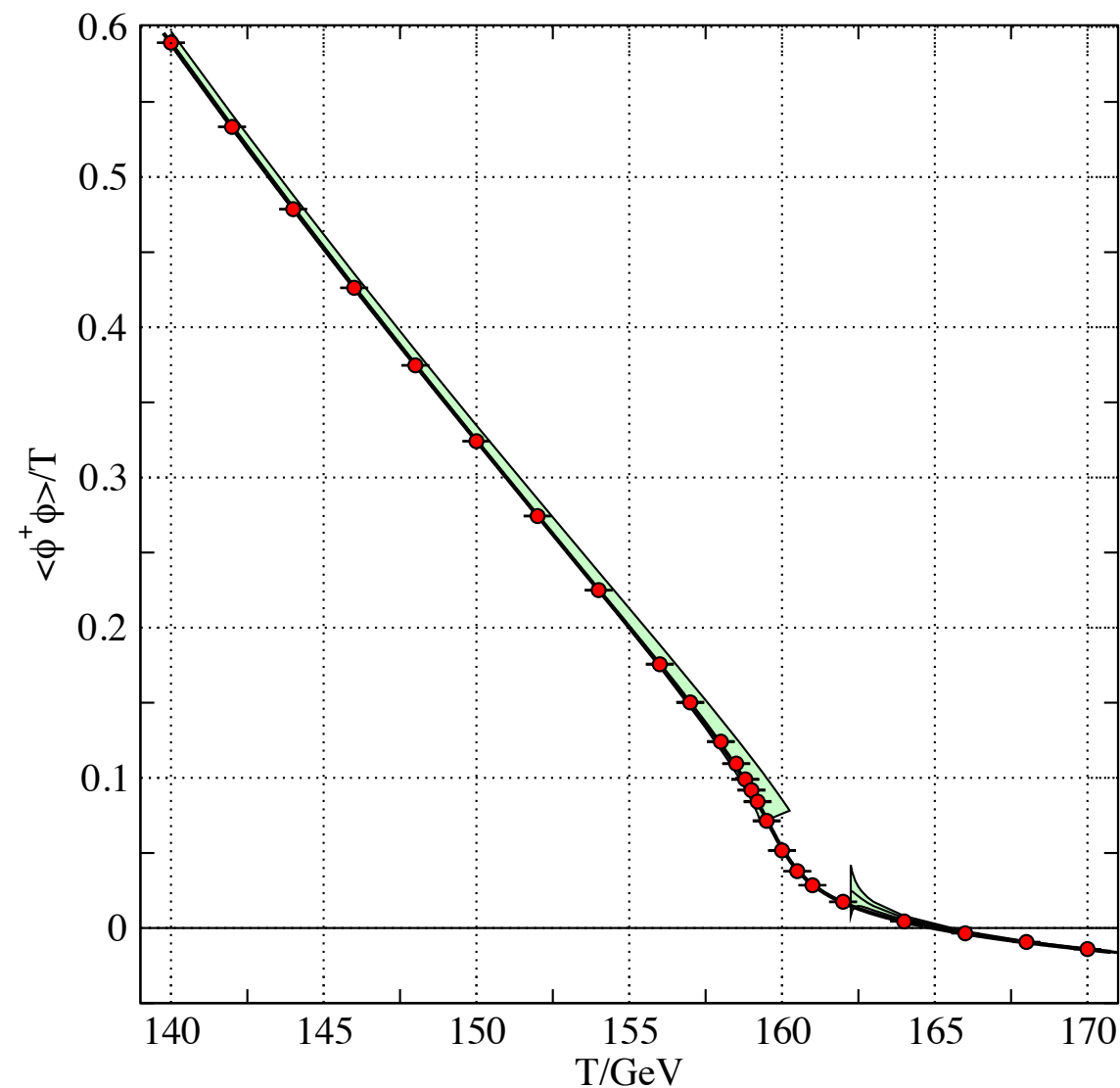
Work in progress w/ Simone Blasi (DESY)



Introduction

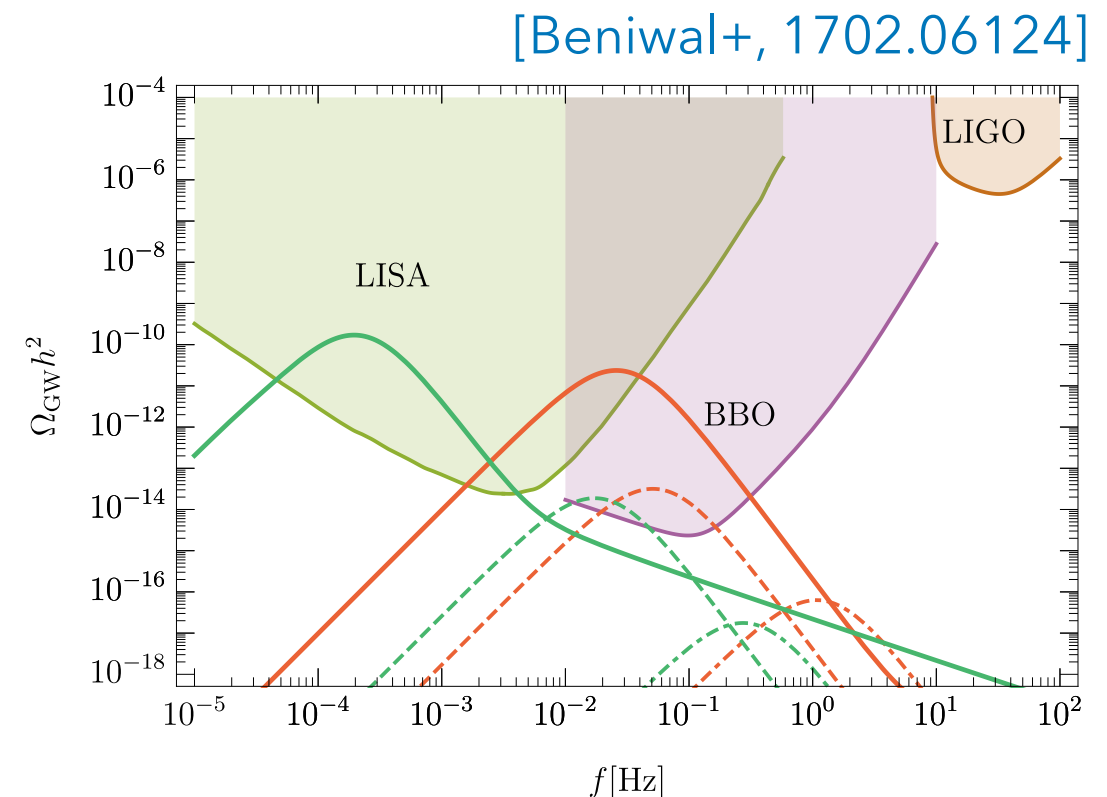
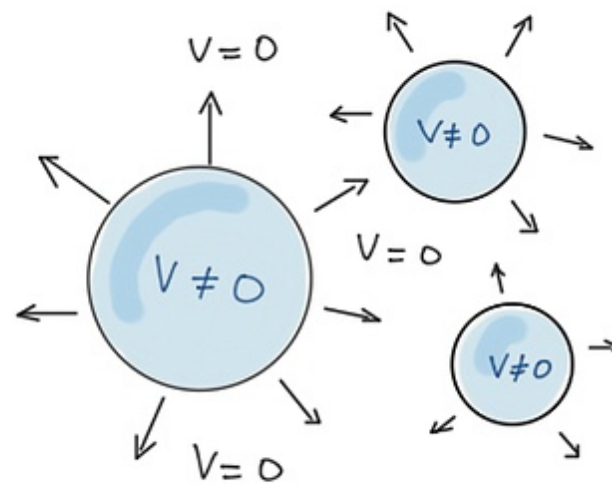
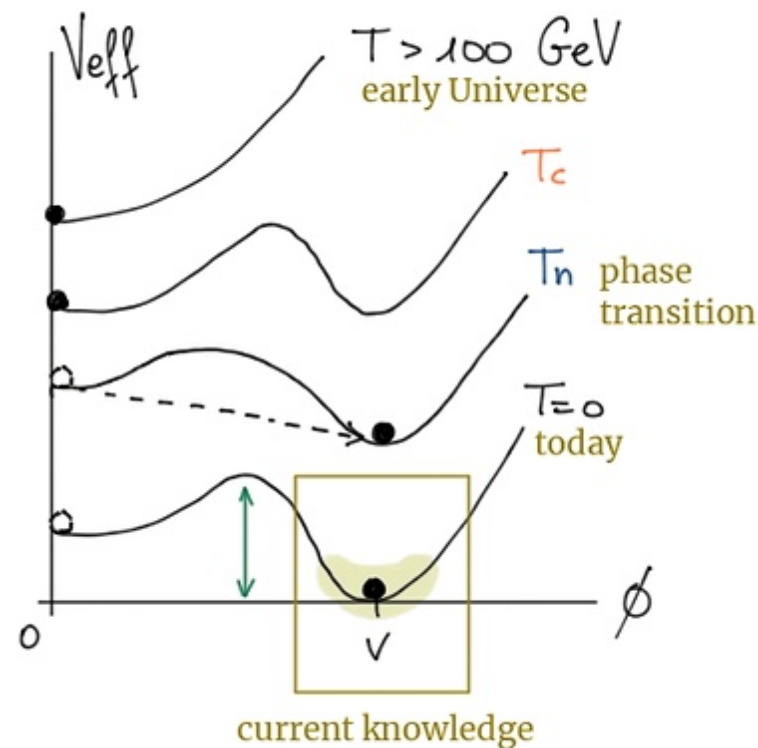
Electroweak Phase Transition

- SM predicts EW crossover, but we don't really know what happened at $T \sim v_{\text{EW}}$



1st-order EWPT in BSM

- A lot of motivation: EW Baryogenesis, Gravitational wave, etc



(Fig credit: Kateryna Radchenko)

- Bubble nucleation seeded by impurities** Cf) Simone's talk

ubiquitous in daily life, but still to be explored in cosmology

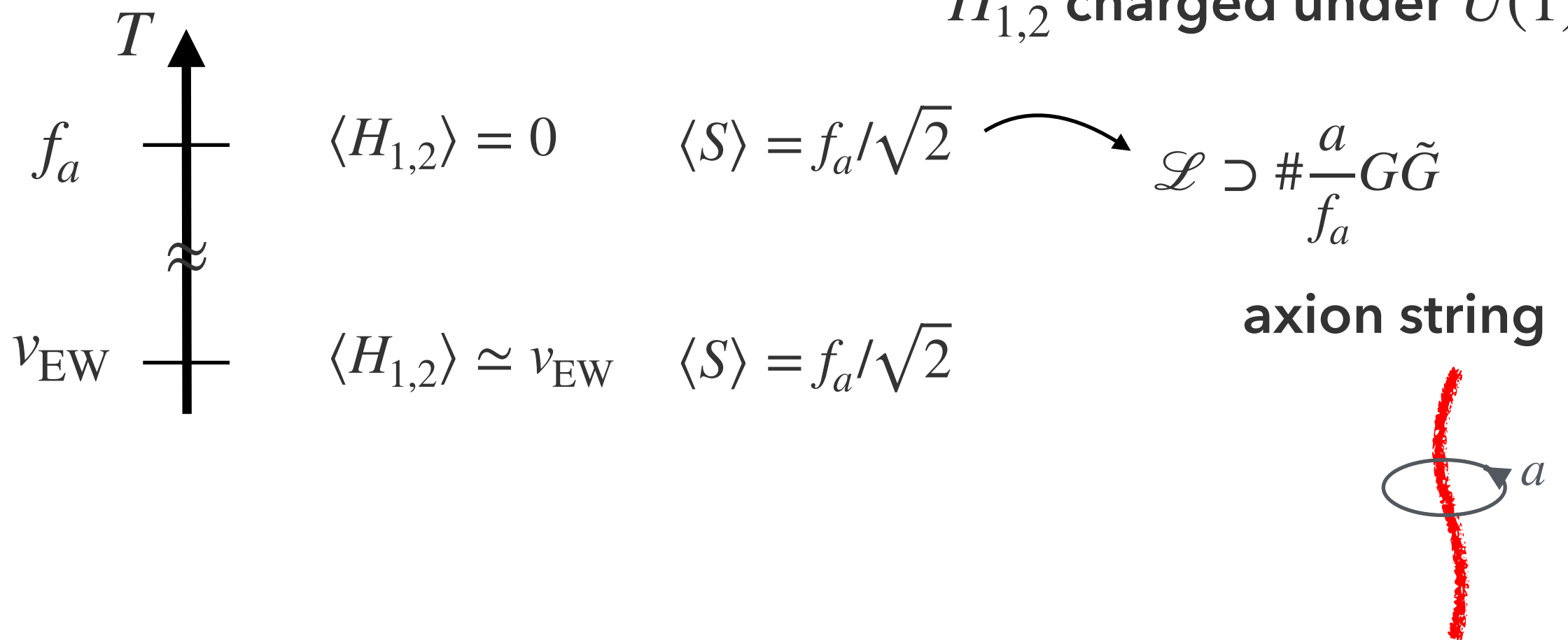
DFSZ axion model

[Zhitnitsky, '80] [Dine-Fischler-Srednicki, '83]

- 2HDM $H_{1,2}$ + Singlet PQ scalar S

$$\mathcal{L} = \mathcal{L}_{\text{kin.}} - V_S(|S|) - V_{\text{EW}}(H_1, H_2) - \left(\kappa S^2 H_1^\dagger H_2 + \text{h.c.} \right)$$

$H_{1,2}$ charged under $U(1)_{PQ}$



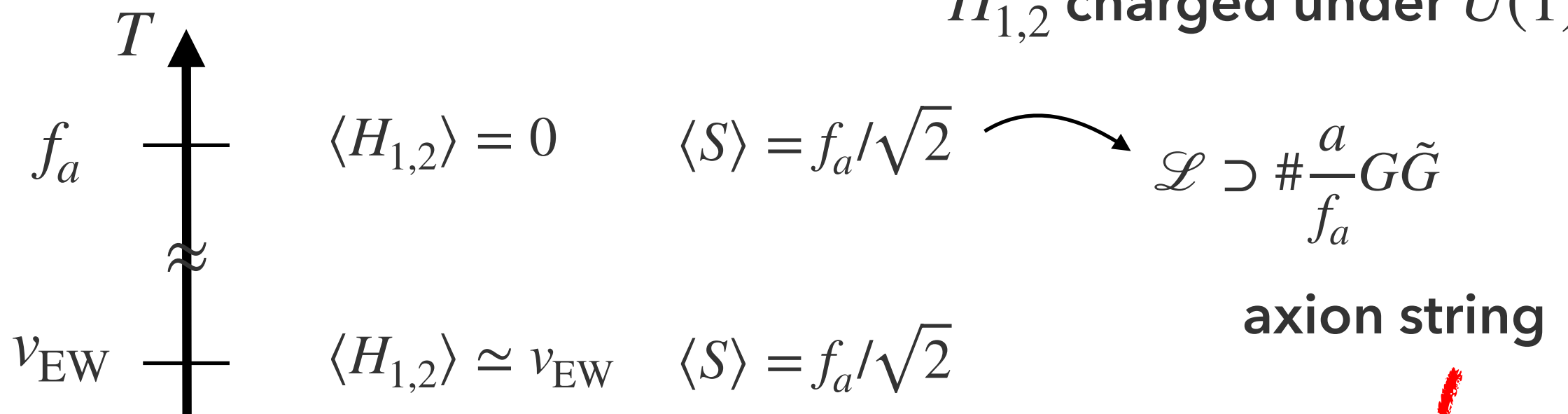
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$H_{1,2}$ charged under $U(1)_{PQ}$



- Nice setup to consider seeded EWPT
 - SSB of $U(1)_{PQ} \rightarrow$ impurity (axion string)
 - Two doublets $H_{1,2} \rightarrow$ can have 1st-order EWPT

Toy model inspired by DFSZ

DFSZ-like toy model

- Two complex scalar fields: φ (SM Higgs like), Φ (PQ-like)

$$\mathcal{L} = |\partial_\mu \varphi|^2 + |\partial_\mu \Phi|^2 - V_{PQ}(|\Phi|) - V(|\varphi|) + (\kappa \varphi^2 \Phi^2 + \text{h.c.})$$



- κ -term is DFSZ-ish operator

$$|\varphi|^2 |\Phi|^2 \text{ in KSVZ}$$

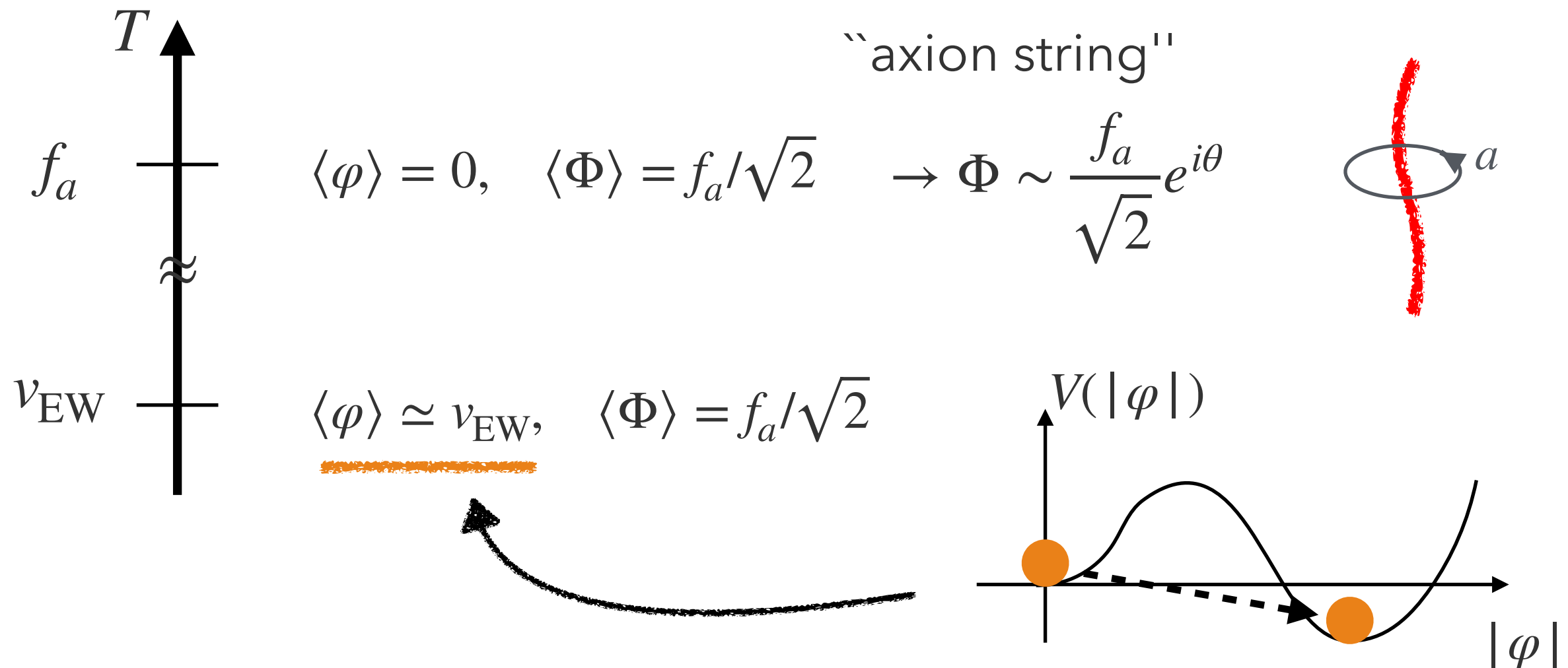
- The model has global $\underline{U(1)_{PQ}} \times \underline{\mathbb{Z}_2}$ sym.

$$\begin{cases} \Phi \rightarrow e^{i\xi} \Phi \\ \varphi \rightarrow e^{-i\xi} \varphi \end{cases} \quad \begin{cases} \Phi \rightarrow i\Phi \\ \varphi \rightarrow i\varphi \end{cases}$$

φ is involved for $U(1)_{PQ}$ (feature of DFSZ)

DFSZ-like toy model

- Heavier scalar Φ takes VEV earlier \rightarrow breaking of $U(1)_{PQ} \rightarrow$ axion!



How is this nucleation influenced by axion string?

Critical bubble in toy model

- The critical bubble in the background of axion string

$$\Phi \sim \frac{f_a}{\sqrt{2}} e^{i\theta}$$

← axion string

$$\varphi = \frac{1}{\sqrt{2}} h(x) e^{i\phi(x)}$$

← polar decomp. of Higgs

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 + \frac{h^2}{2}(\partial_\mu \phi)^2 - V(h) + \underline{2\kappa f_a^2 h^2 \cos 2(\theta + \phi)}$$

position-dependent cos potential



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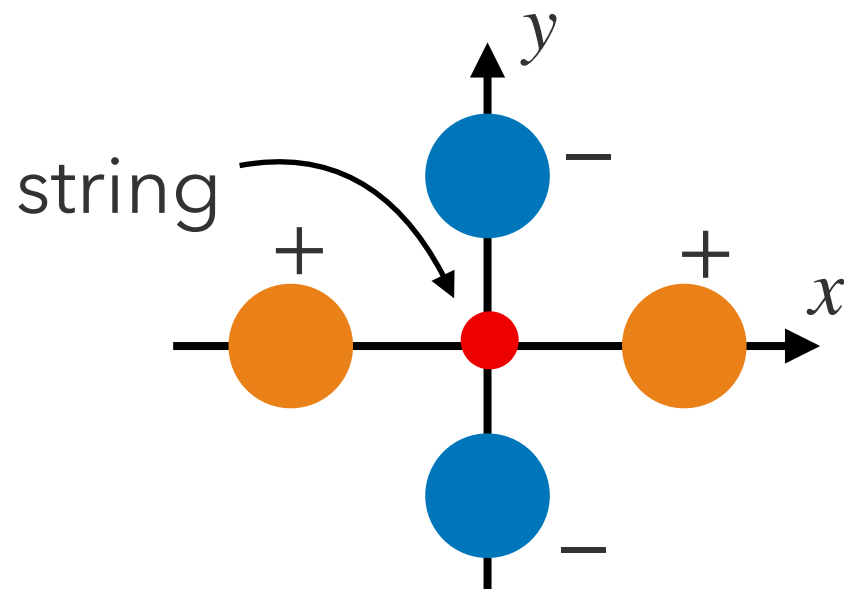
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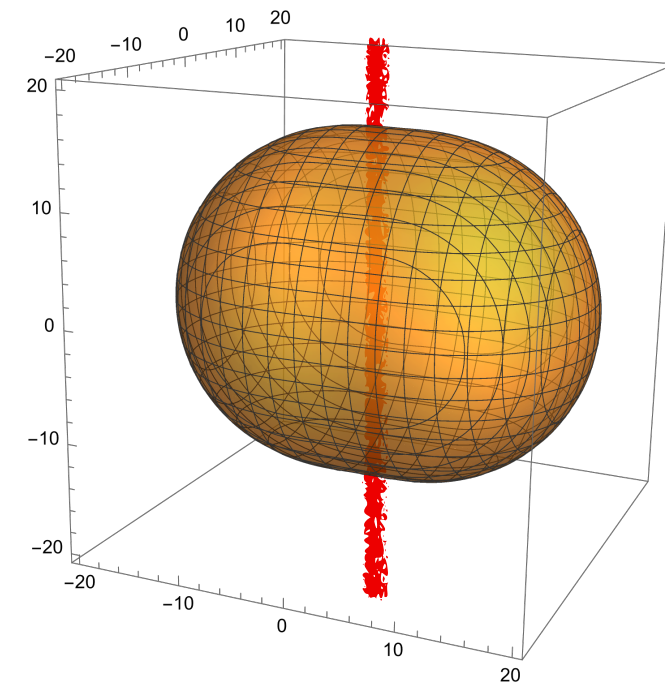
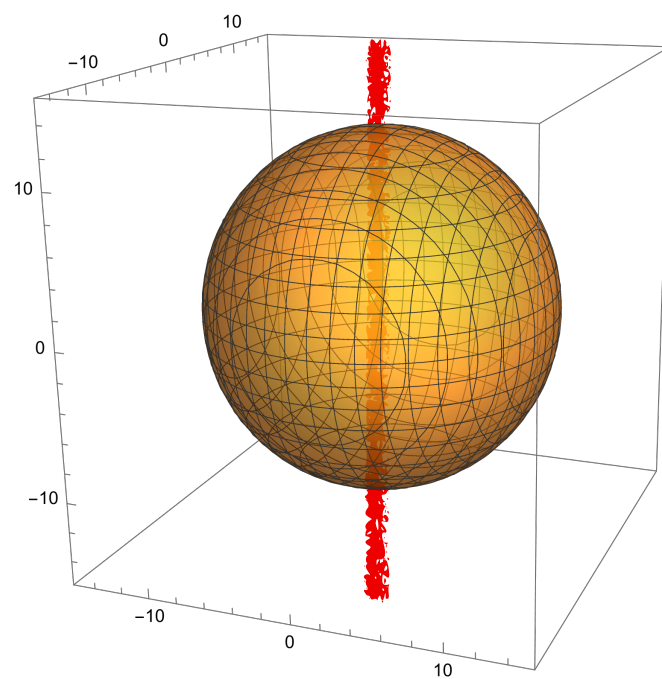
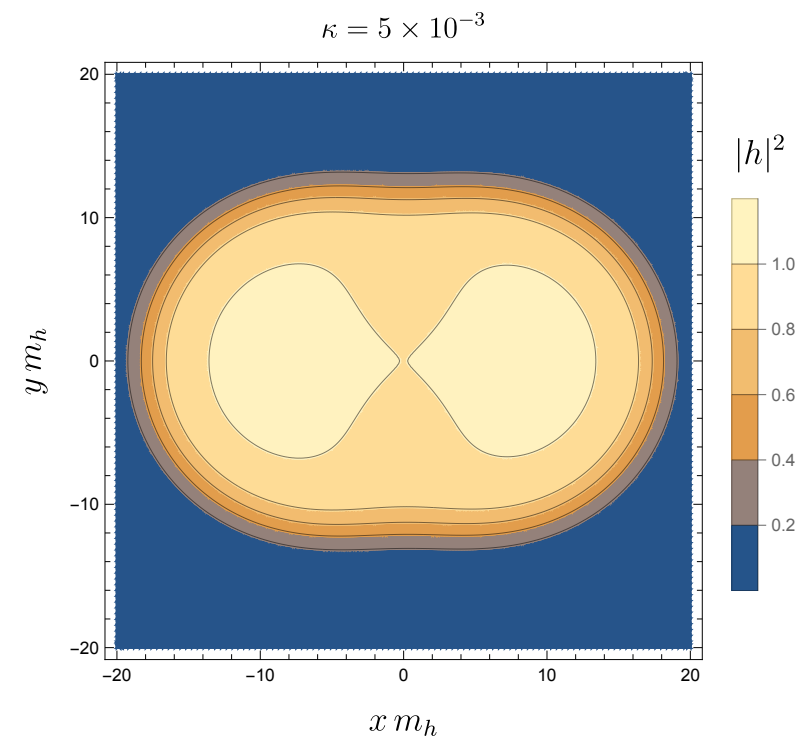
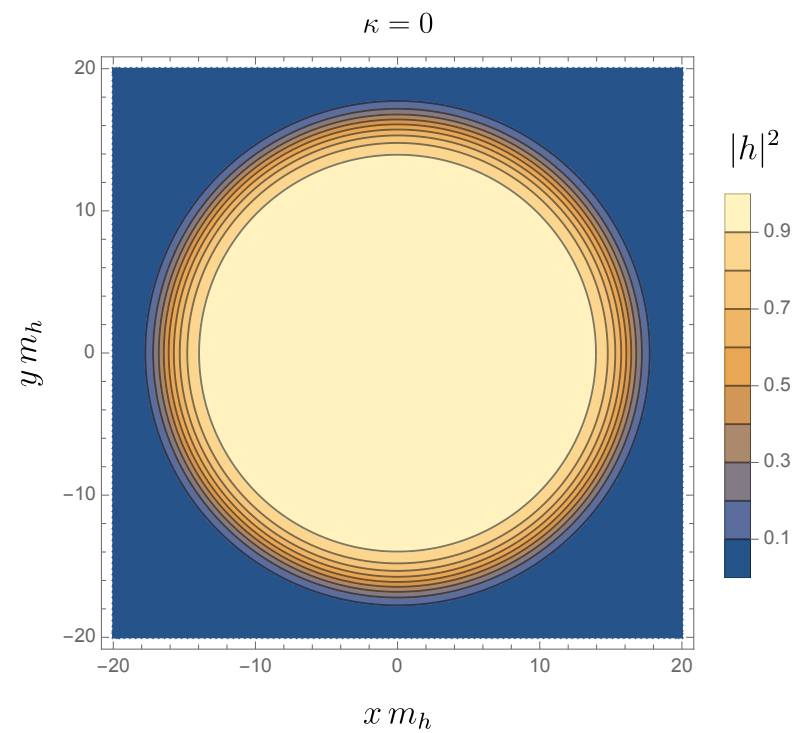
position-dependent cos potential



$$\cos 2\phi @ \theta = 0, \pi$$

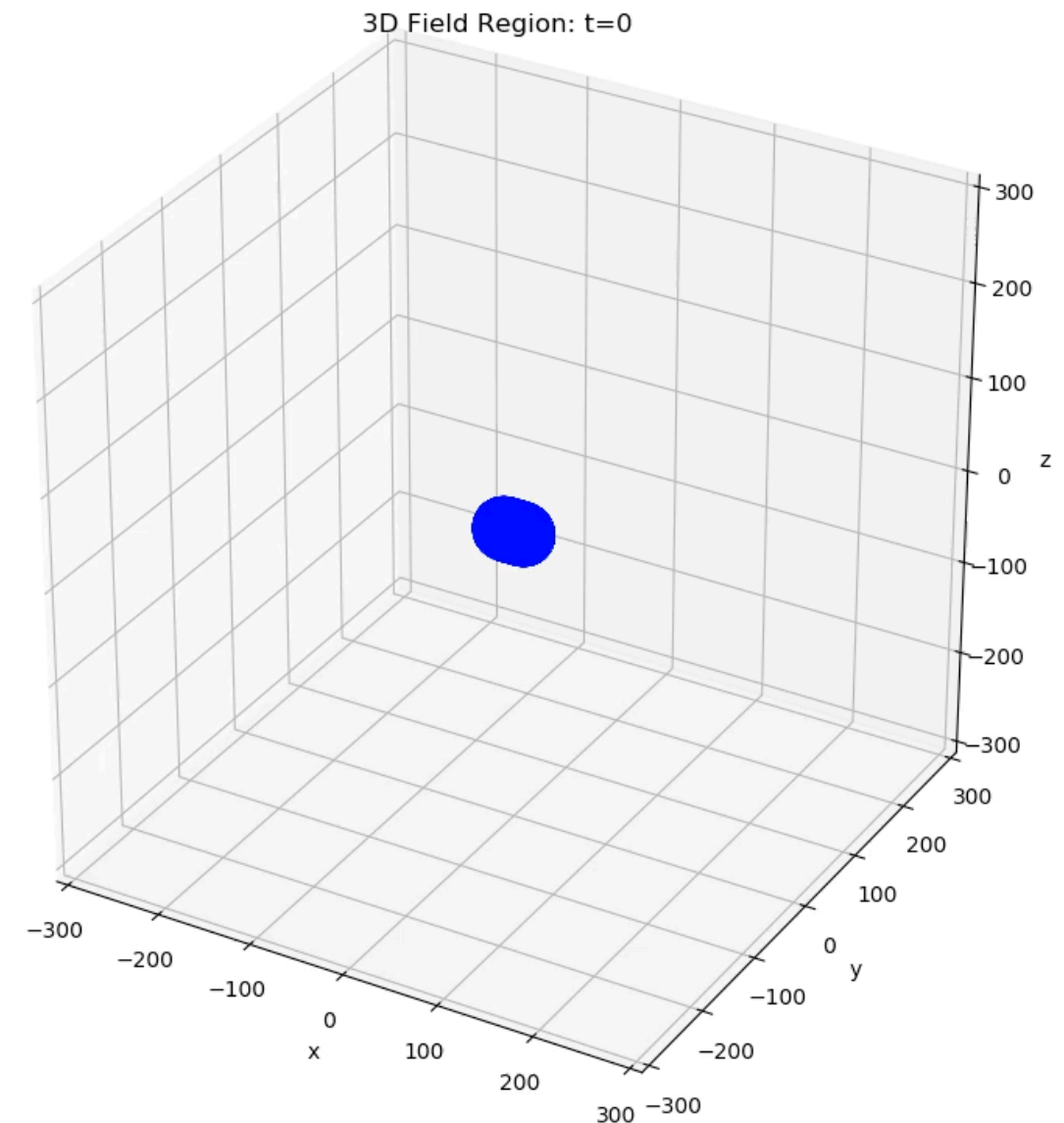
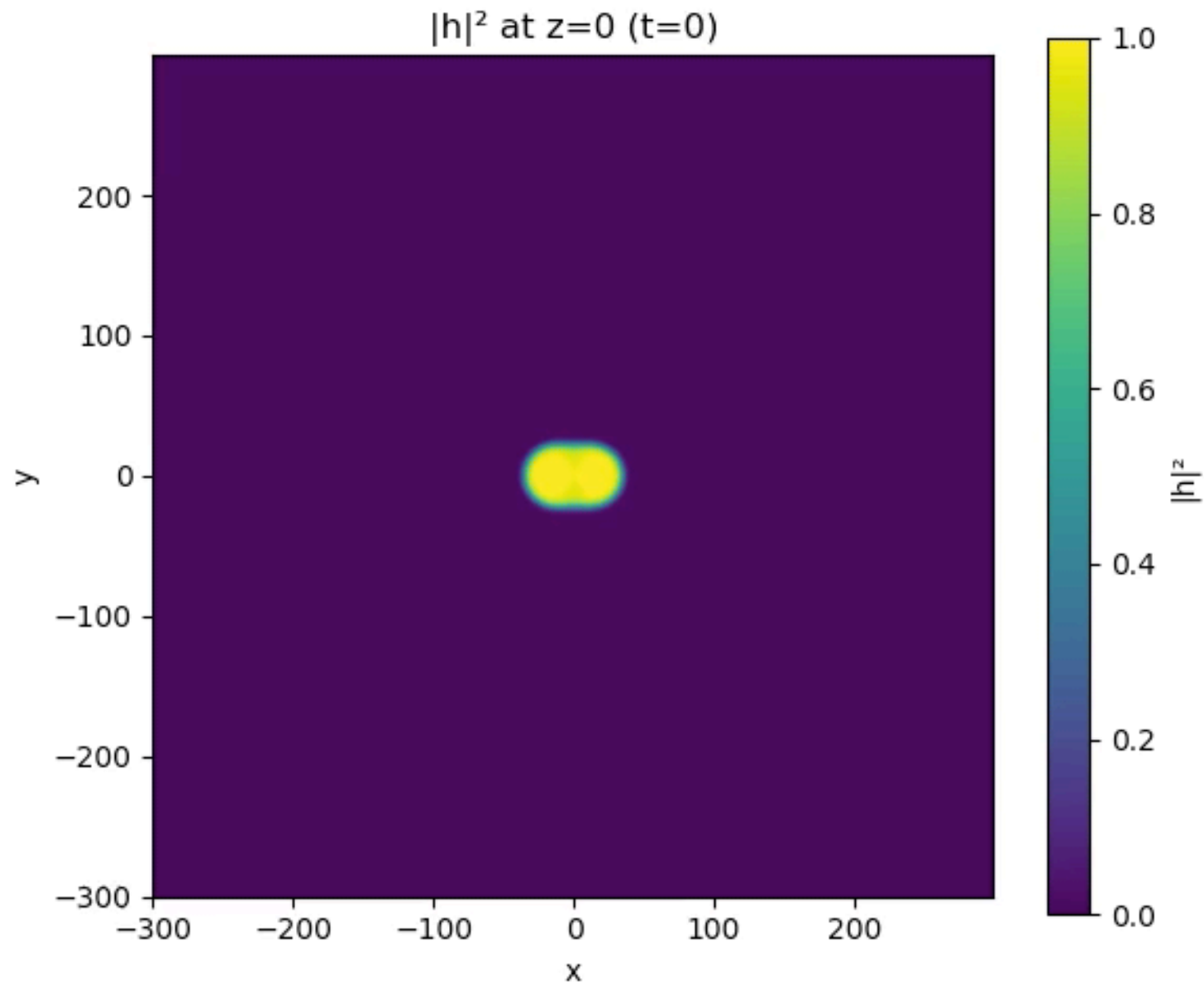
$$-\cos 2\phi @ \theta = \pi/2, 3\pi/2$$

Numerical results for critical bubble

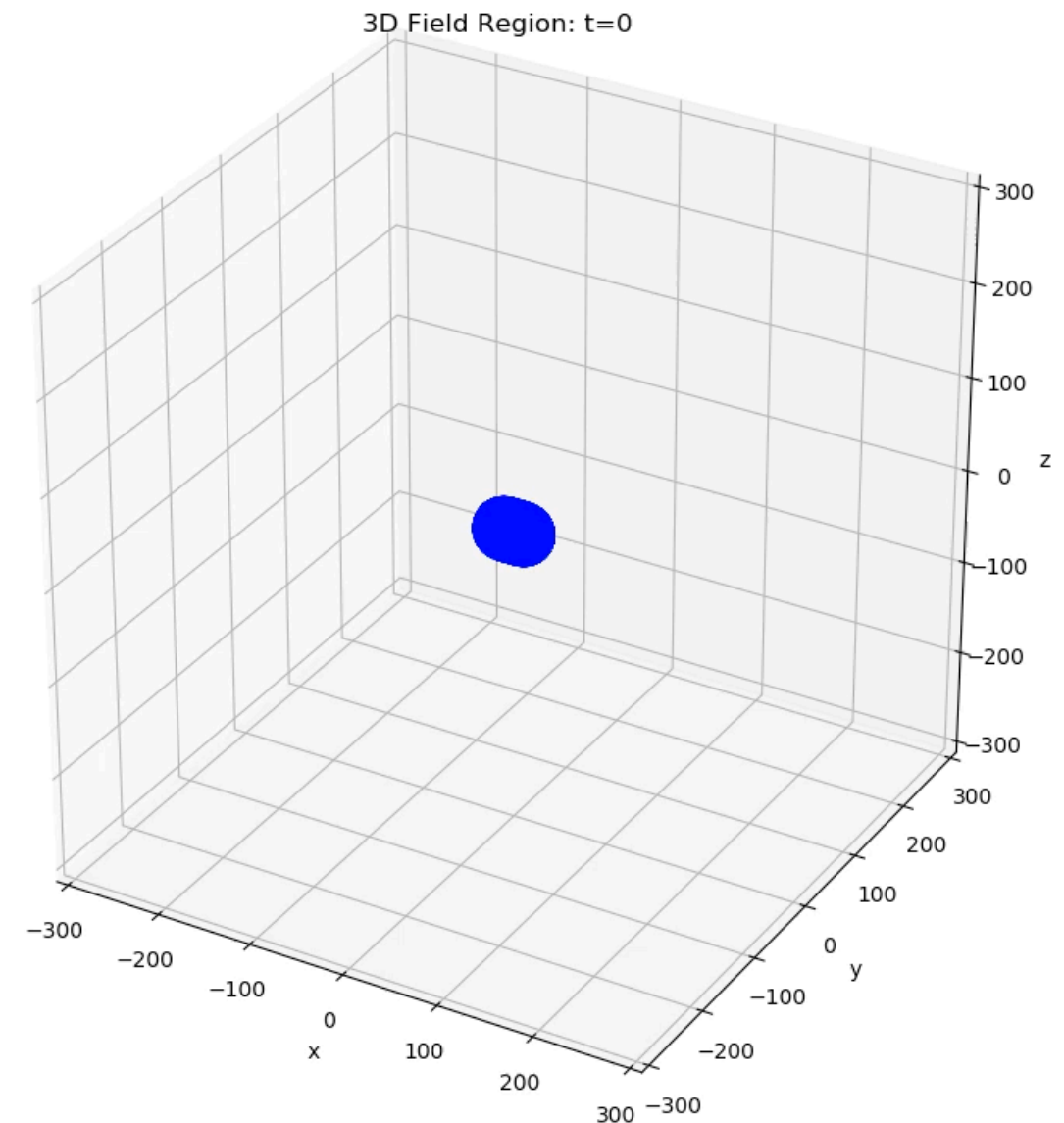
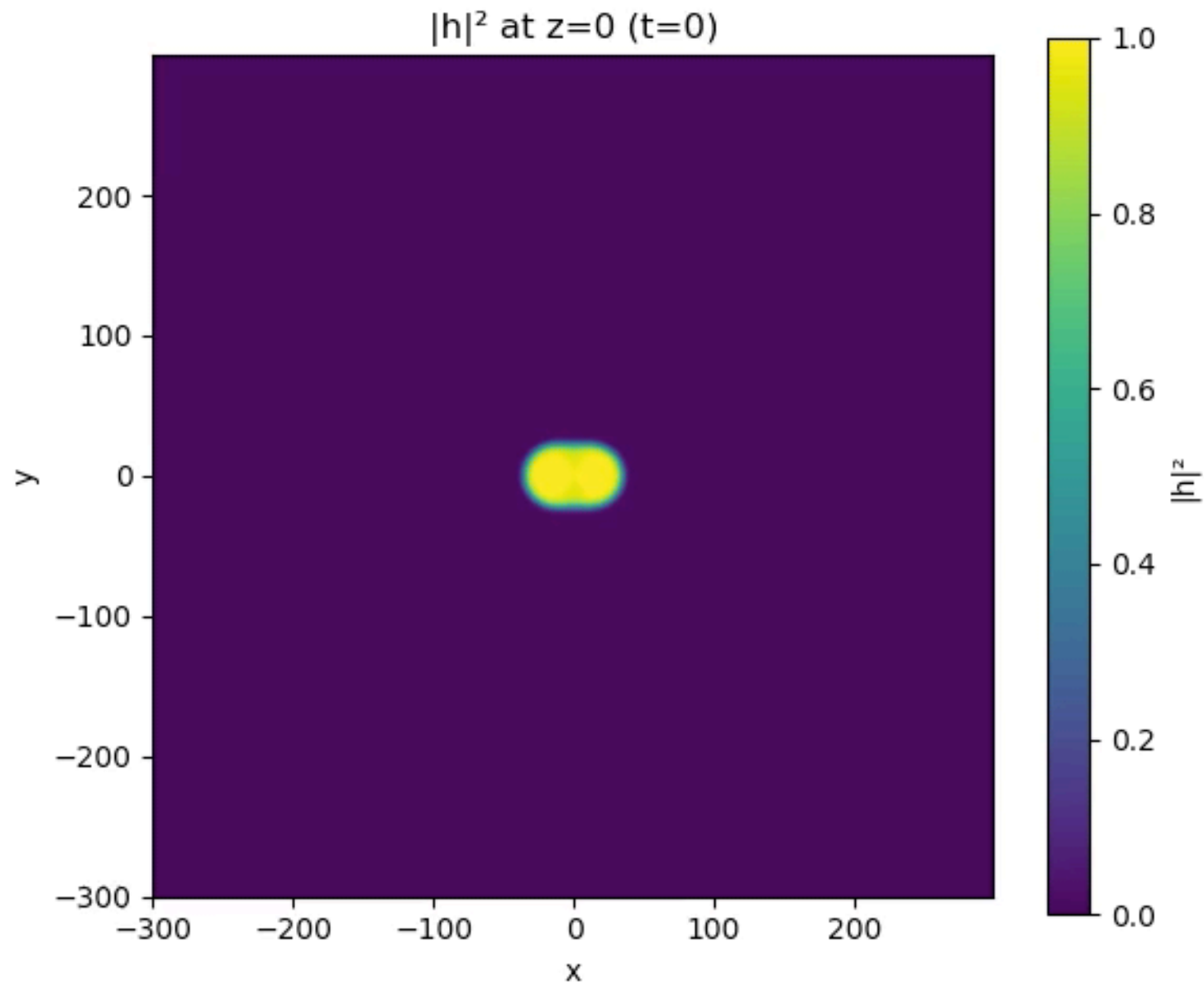


no $O(3)$ or $O(2)$ sym.

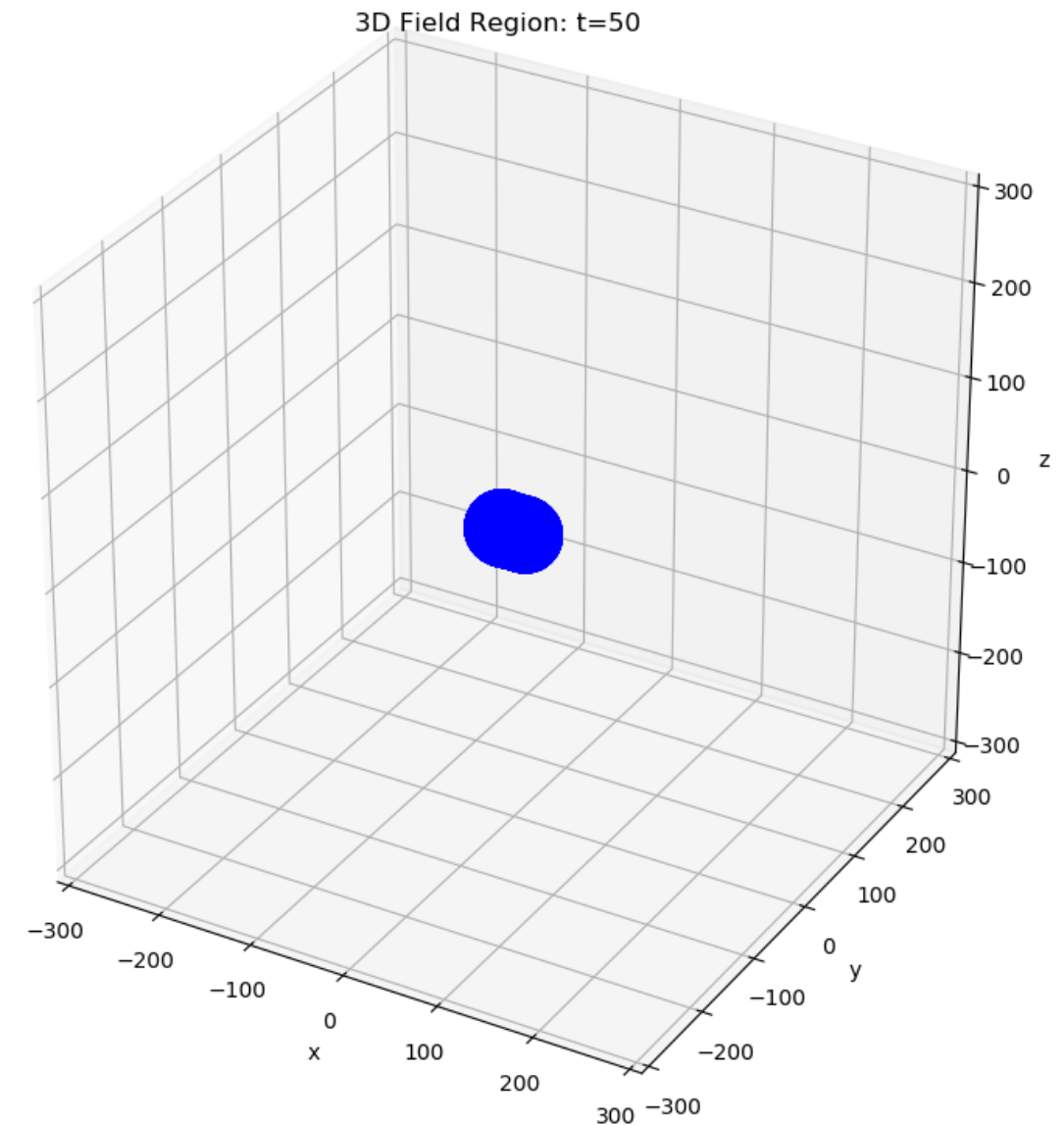
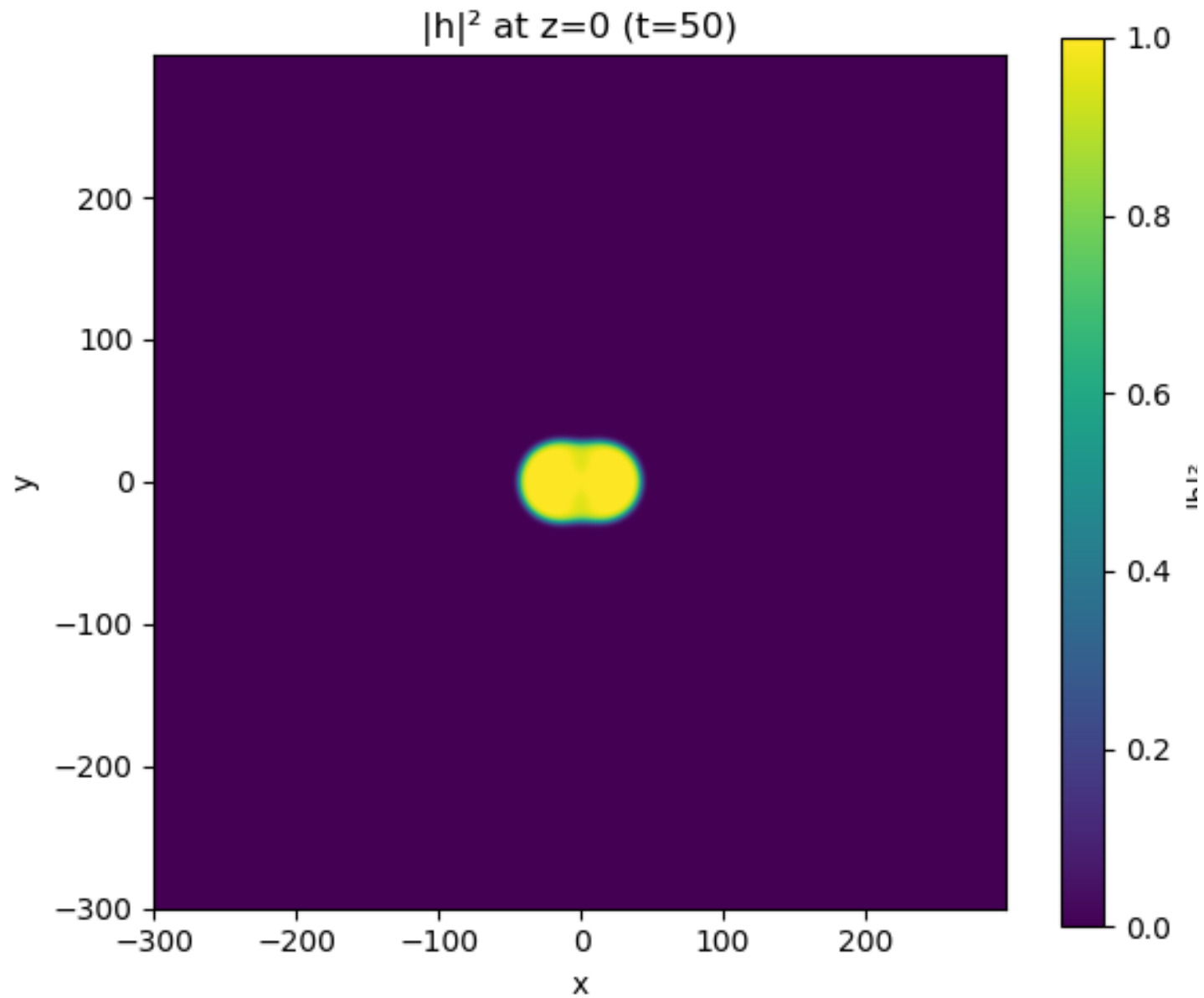
Expansion after nucleation



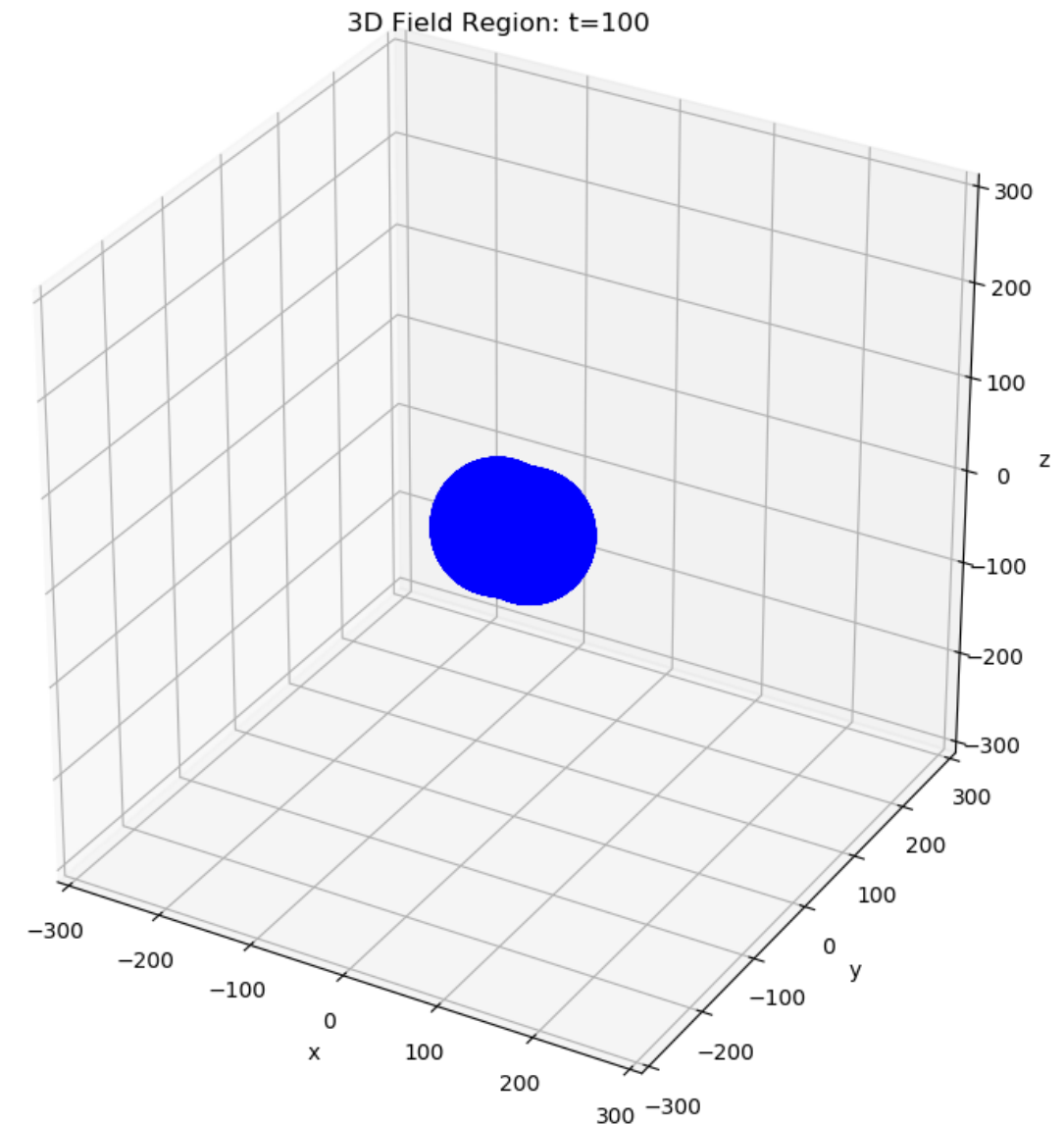
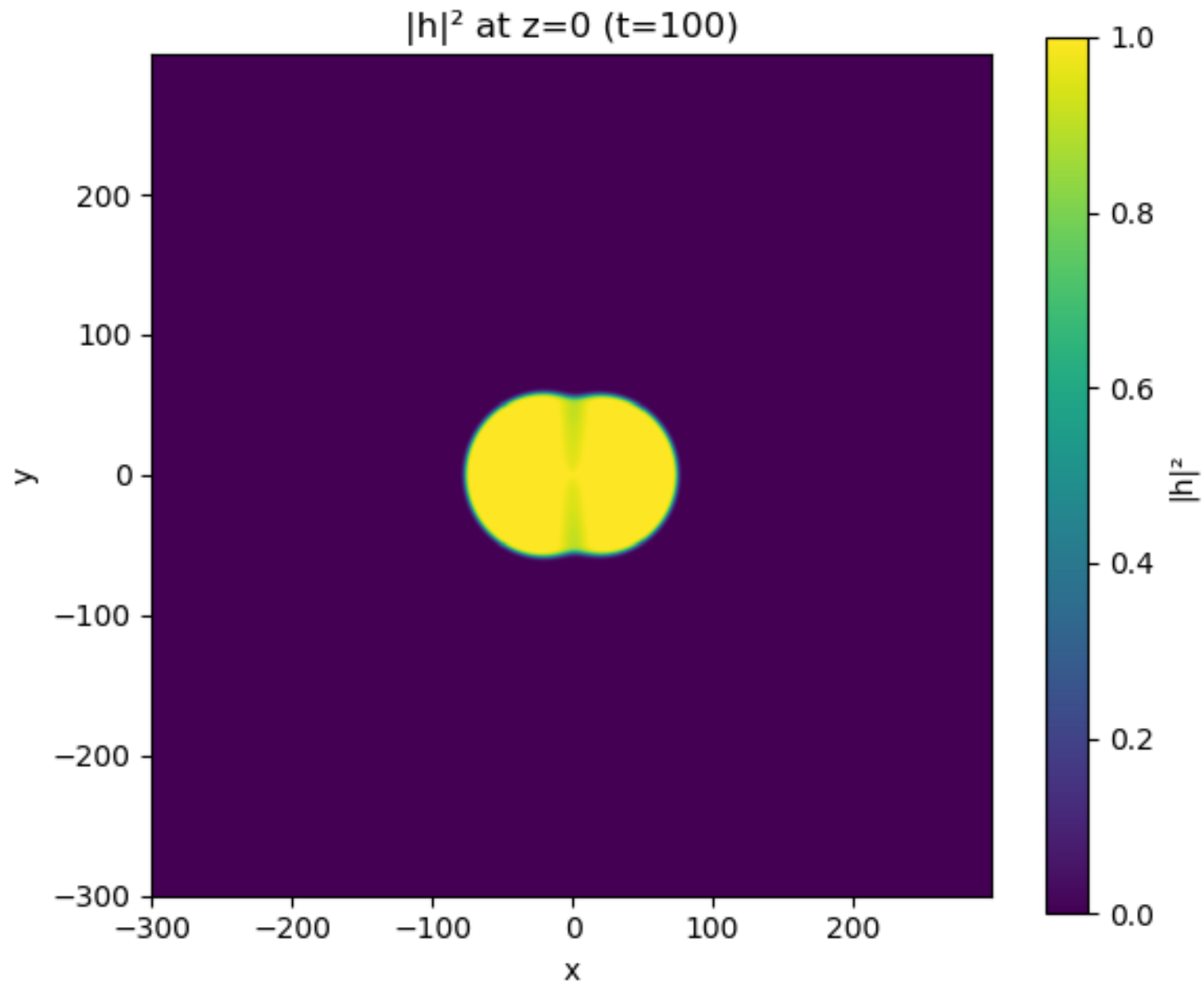
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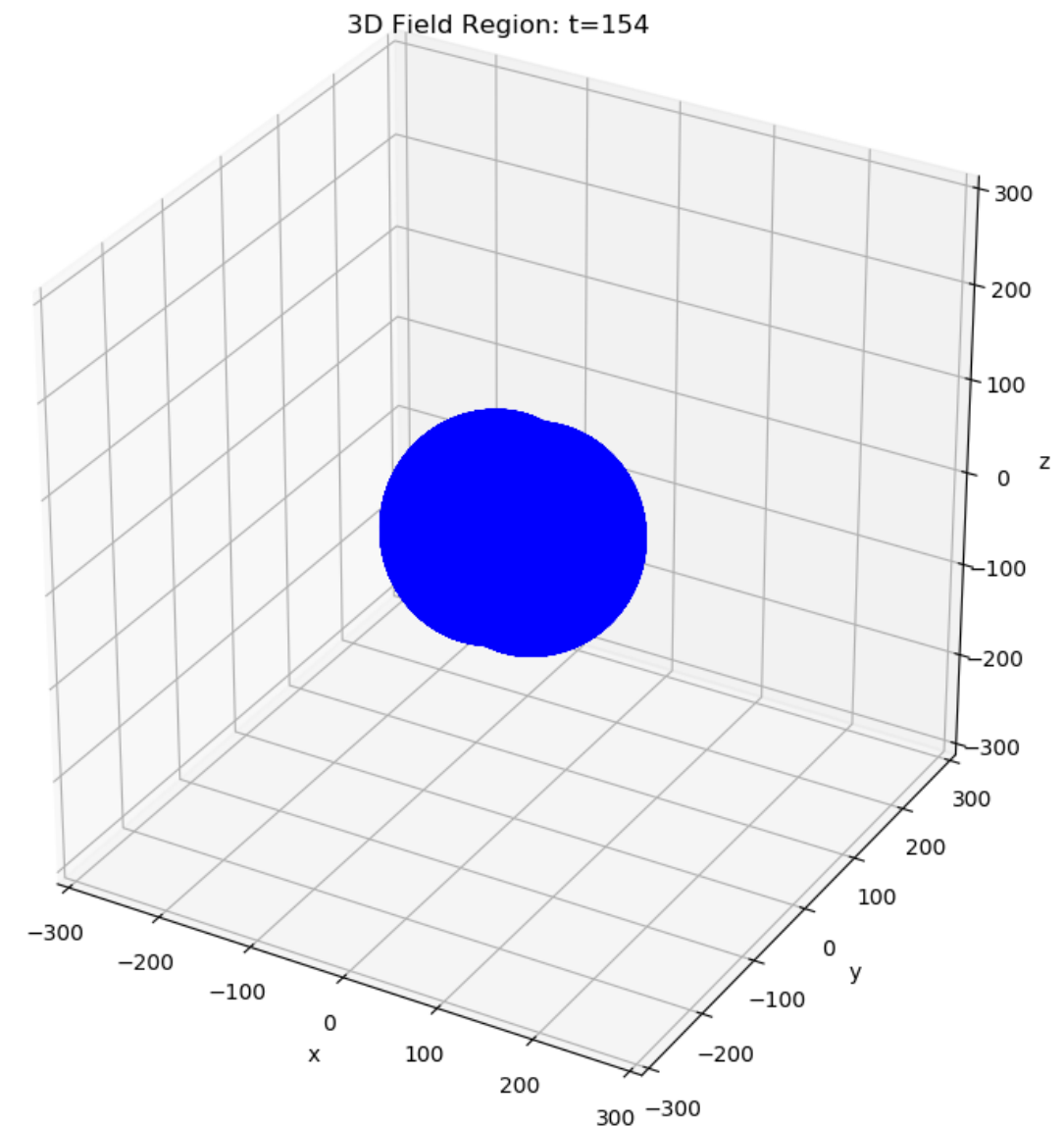
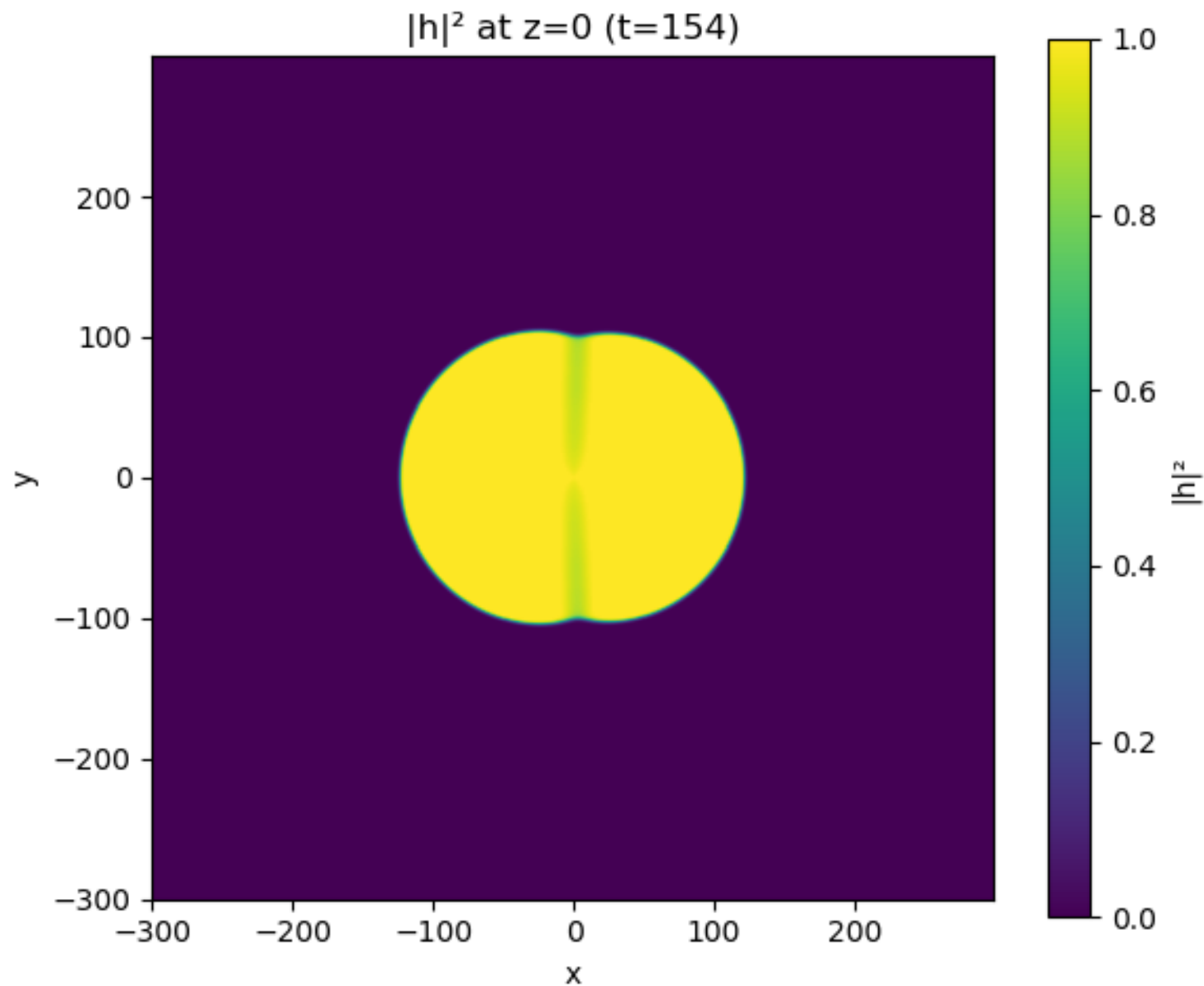
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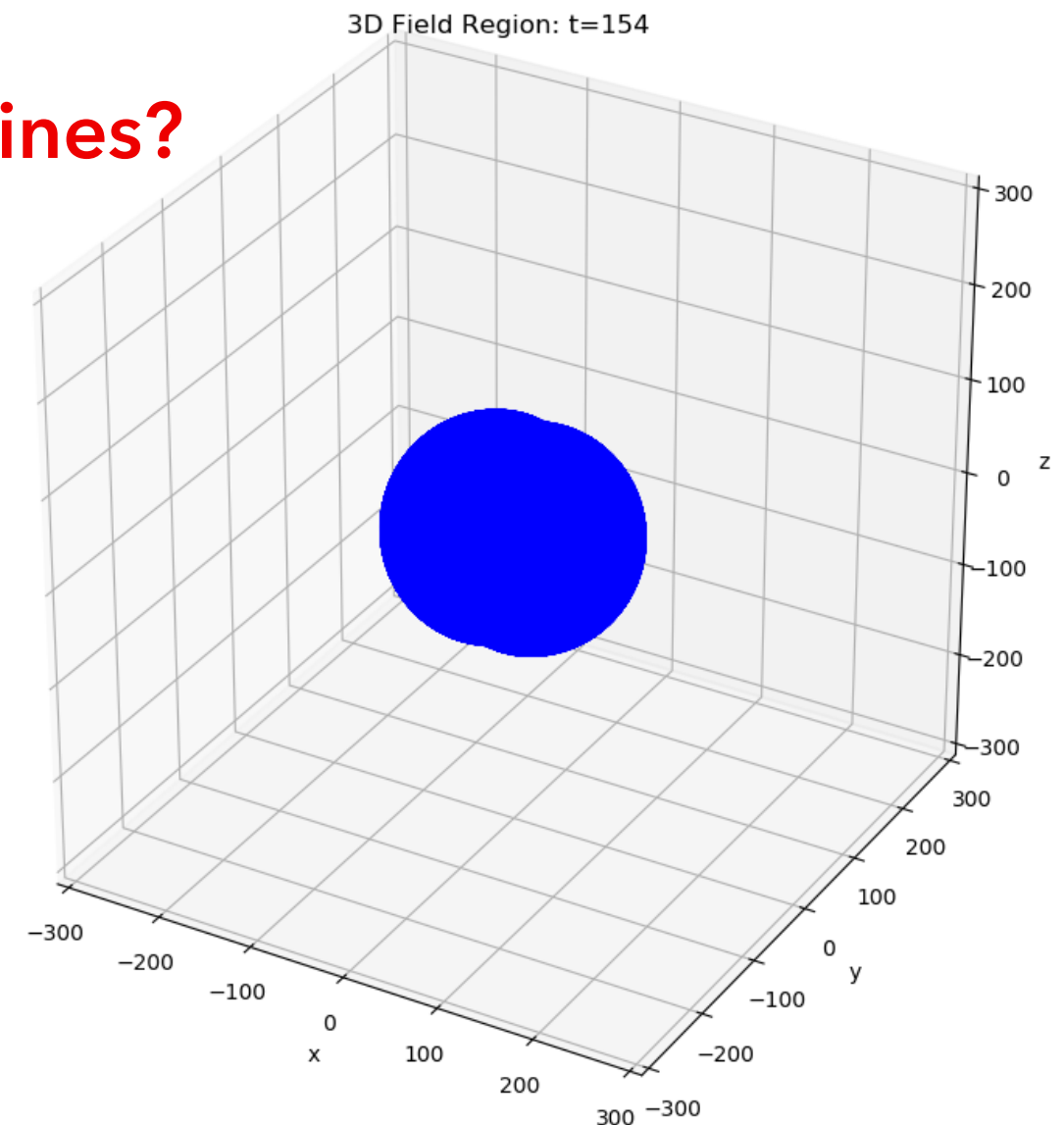
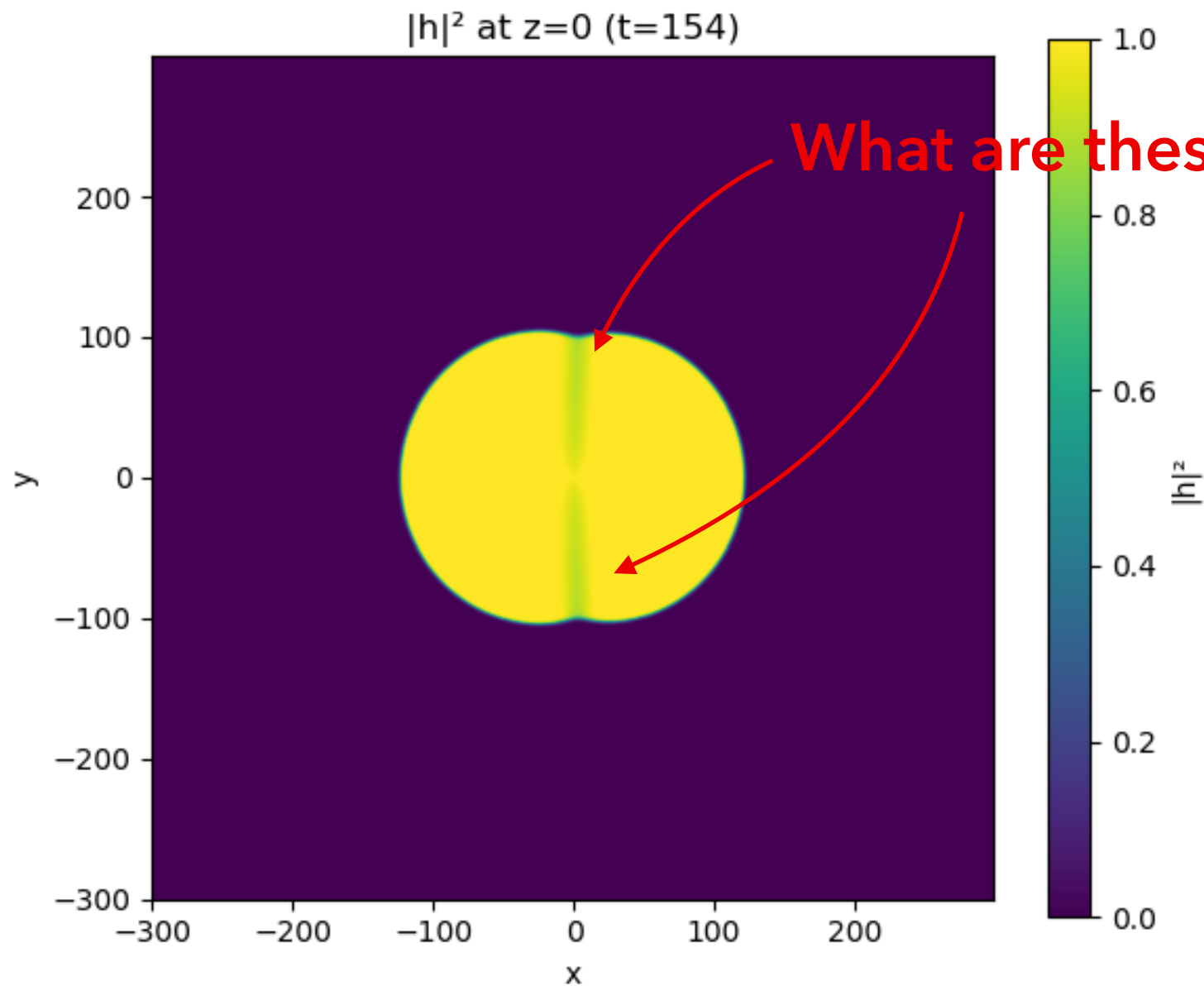
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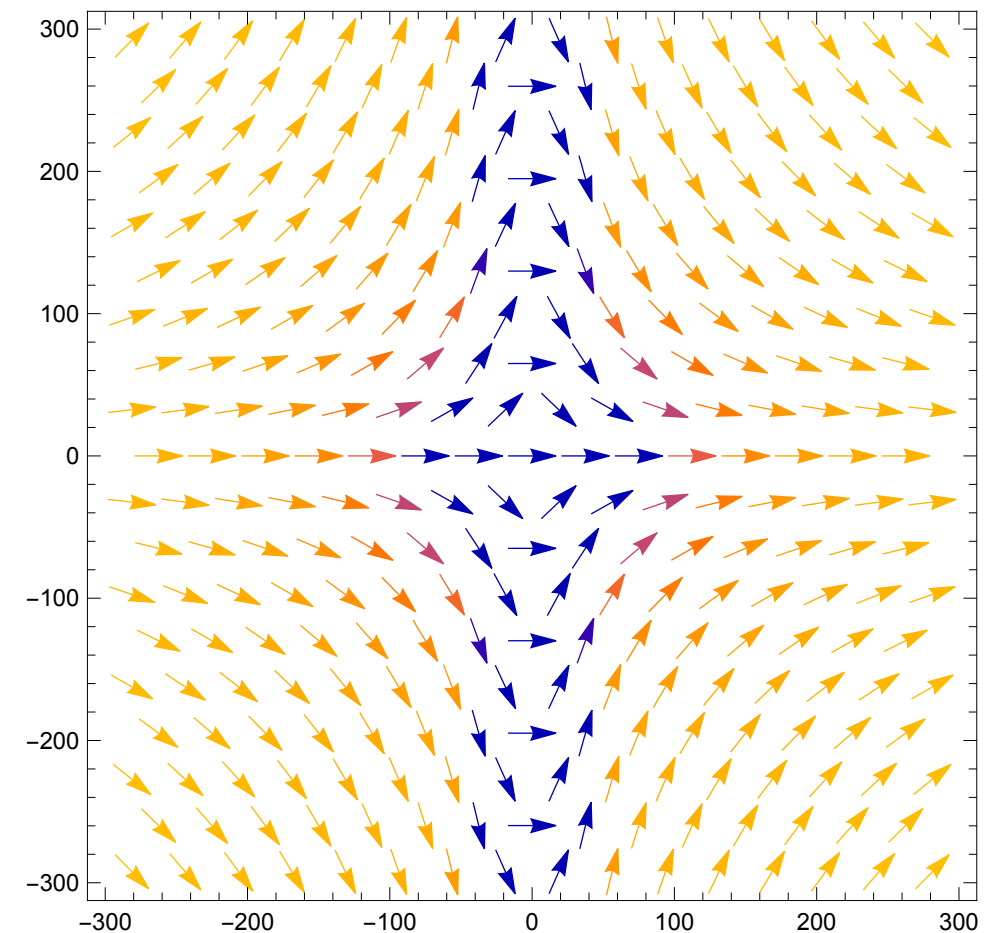
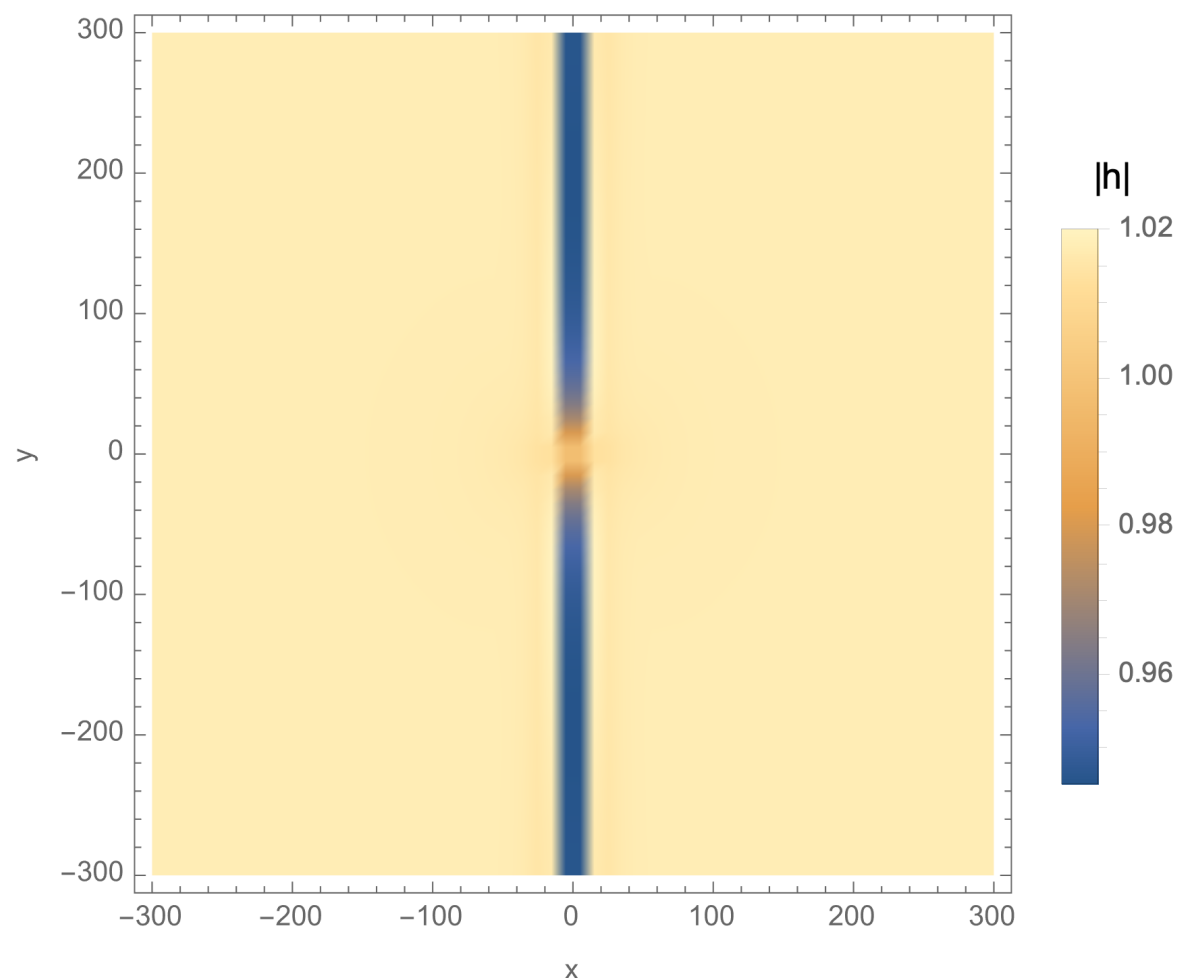


Domain wall

- DW exists in this model! → **New mechanism of DW production**

$$\mathcal{L} = |\partial_\mu \varphi|^2 + |\partial_\mu \Phi|^2 - V(|\varphi|) - V_{PQ}(|\Phi|) + (\kappa \varphi^2 \Phi^2 + \text{h.c.}) \quad U(1)_{PQ} \times \mathbb{Z}_2 \text{ sym.}$$

(Re φ , Im φ)

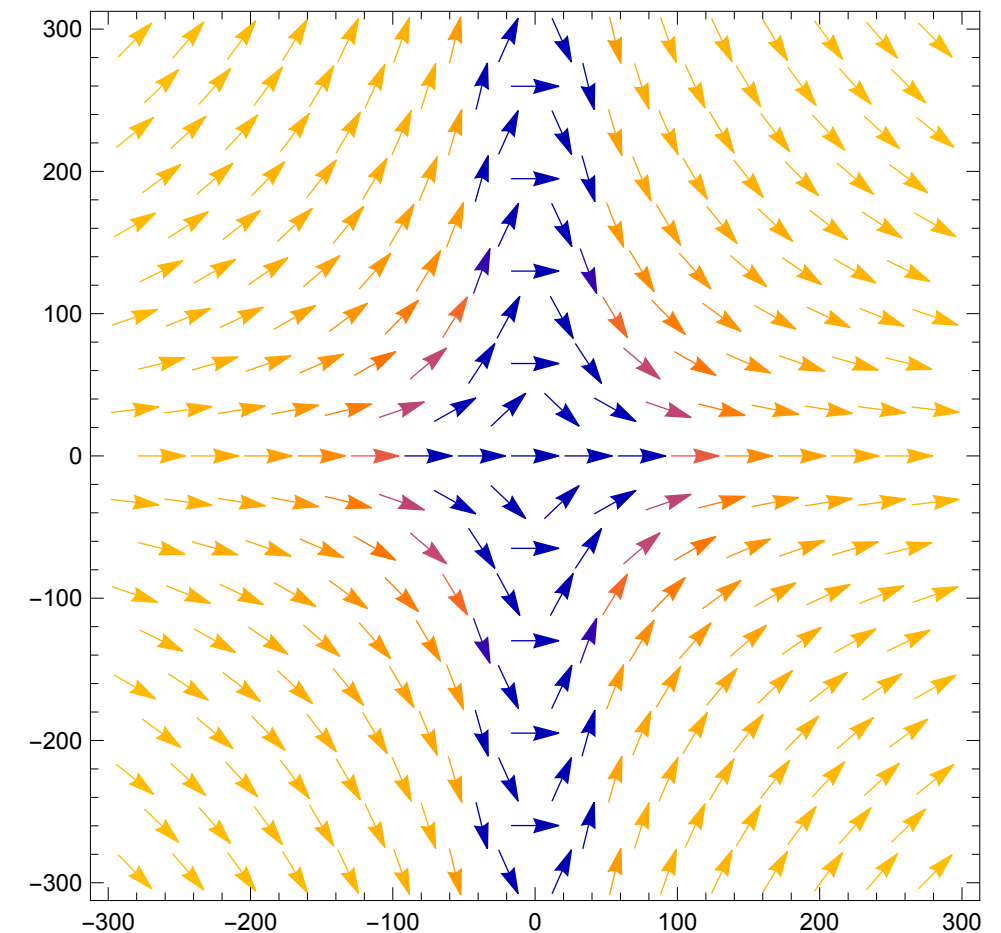
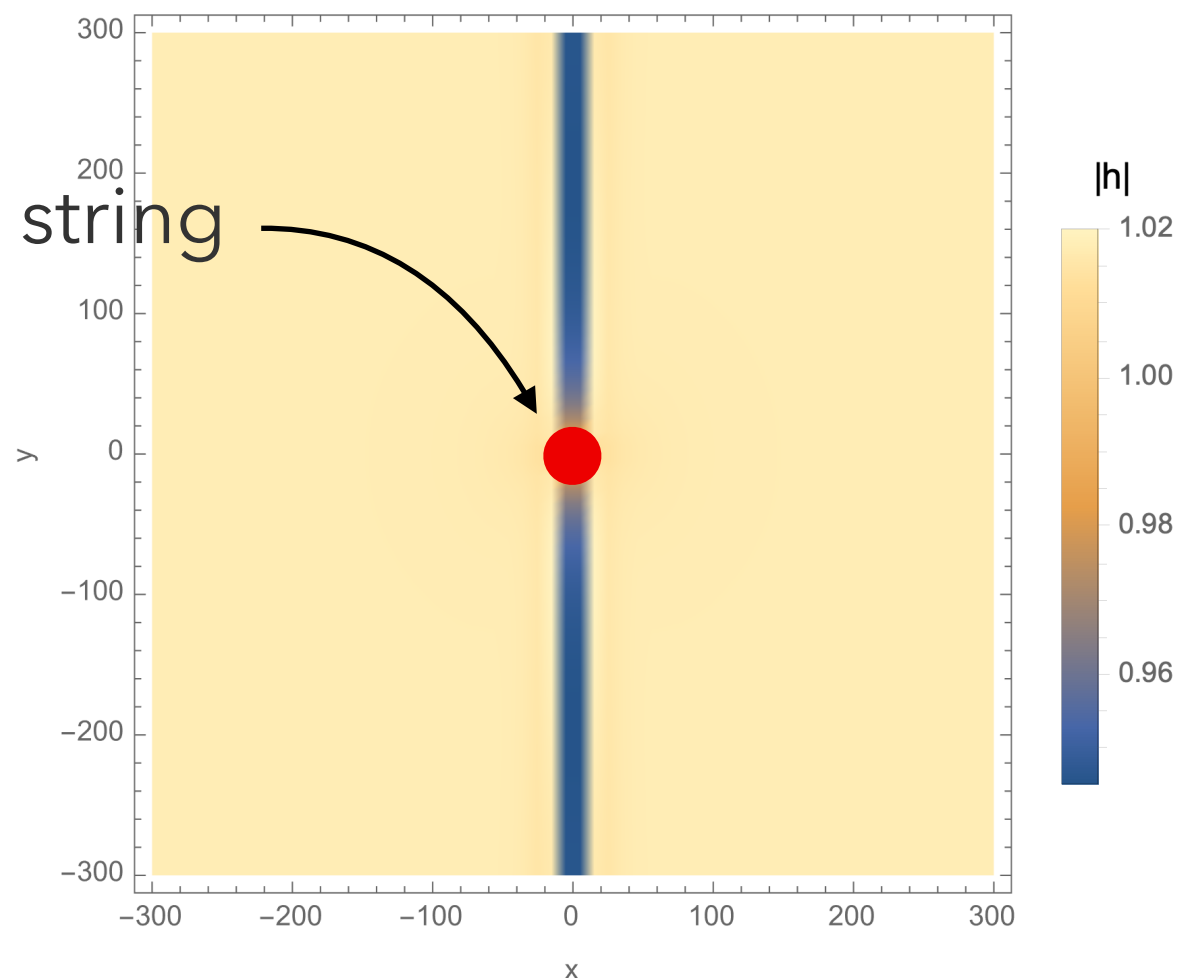


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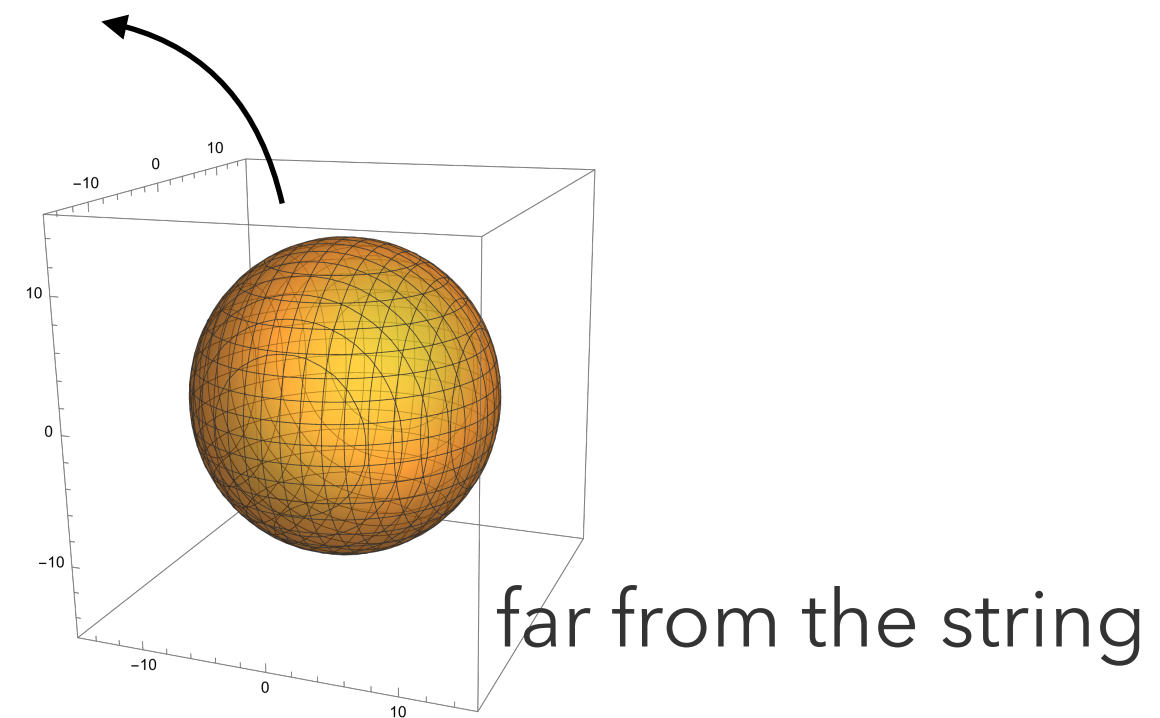
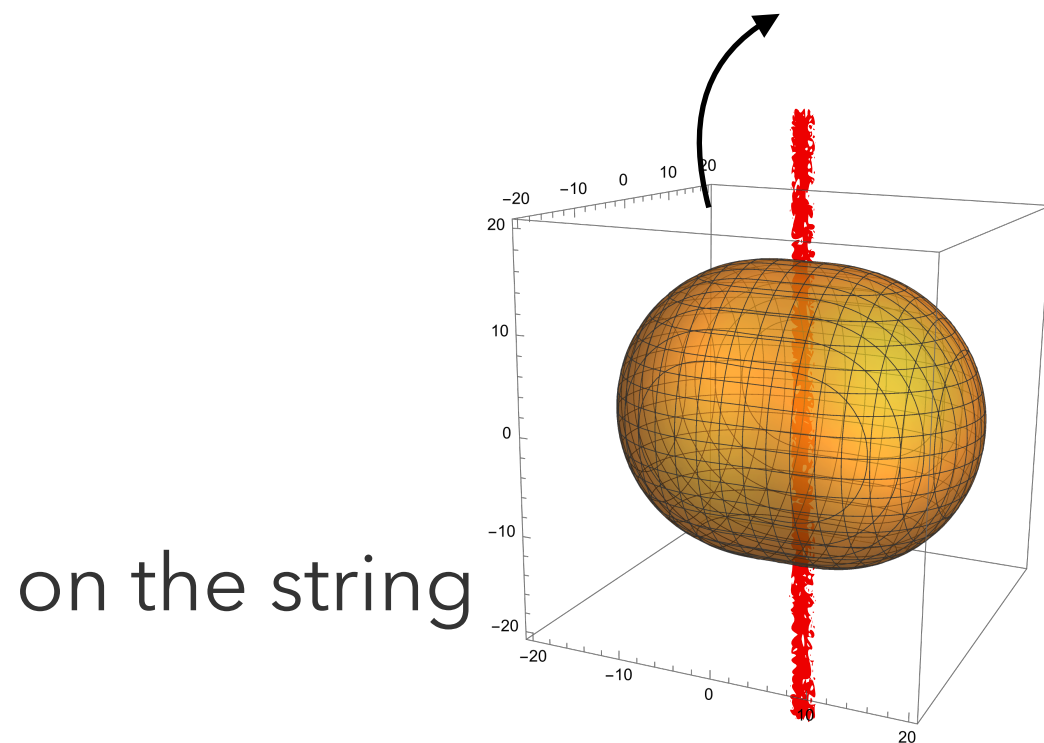
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Tunneling Probability

- This κ -term makes phenomenology much richer, but **doesn't enhance the seeded tunneling itself.** 😓

$$\because \Delta E_c \equiv E_c|_{\text{Seeded}} - E_c|_{\text{Hom}} \sim \kappa f_a^2 v_{\text{EW}}^2 R^3$$

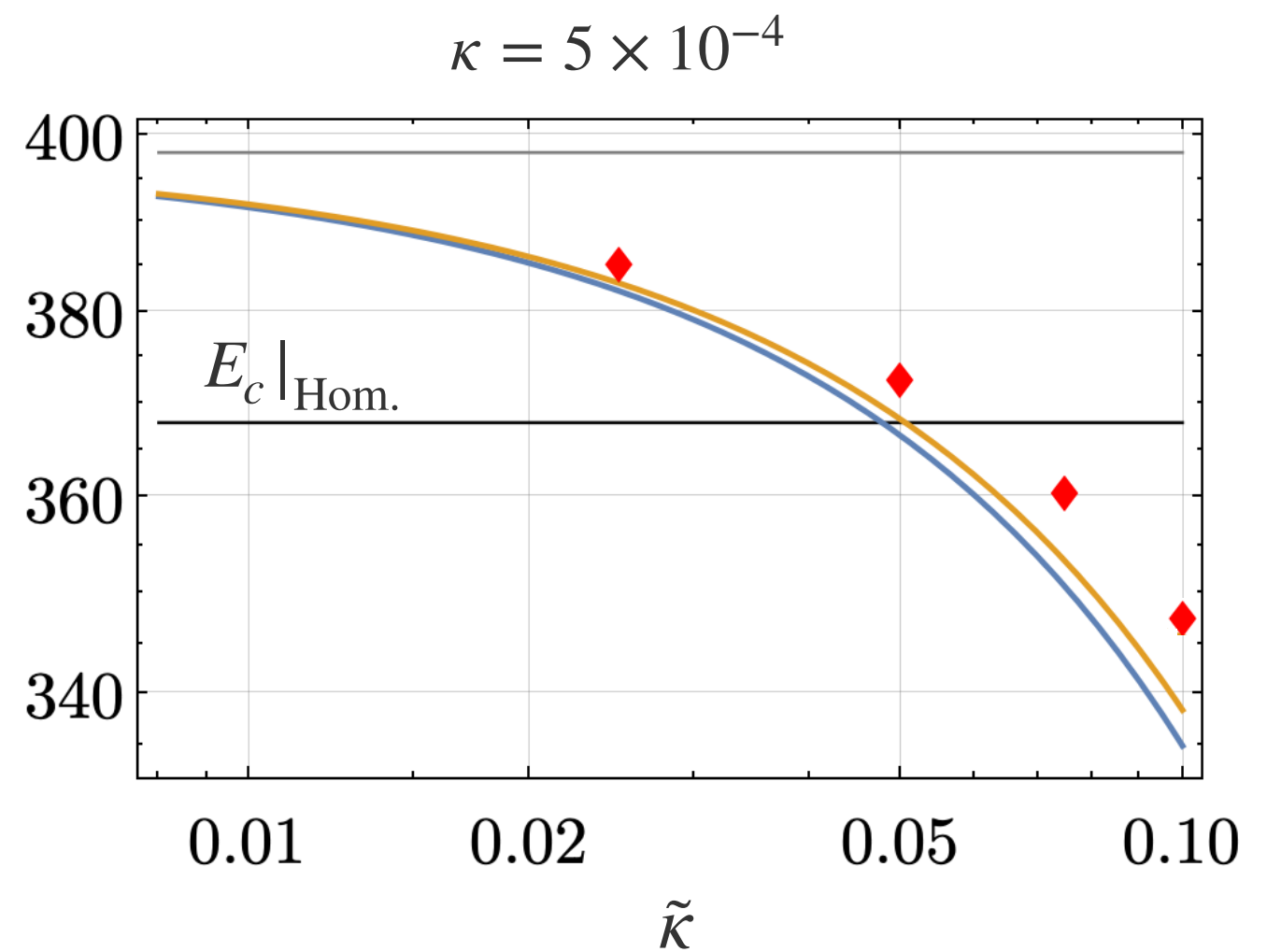
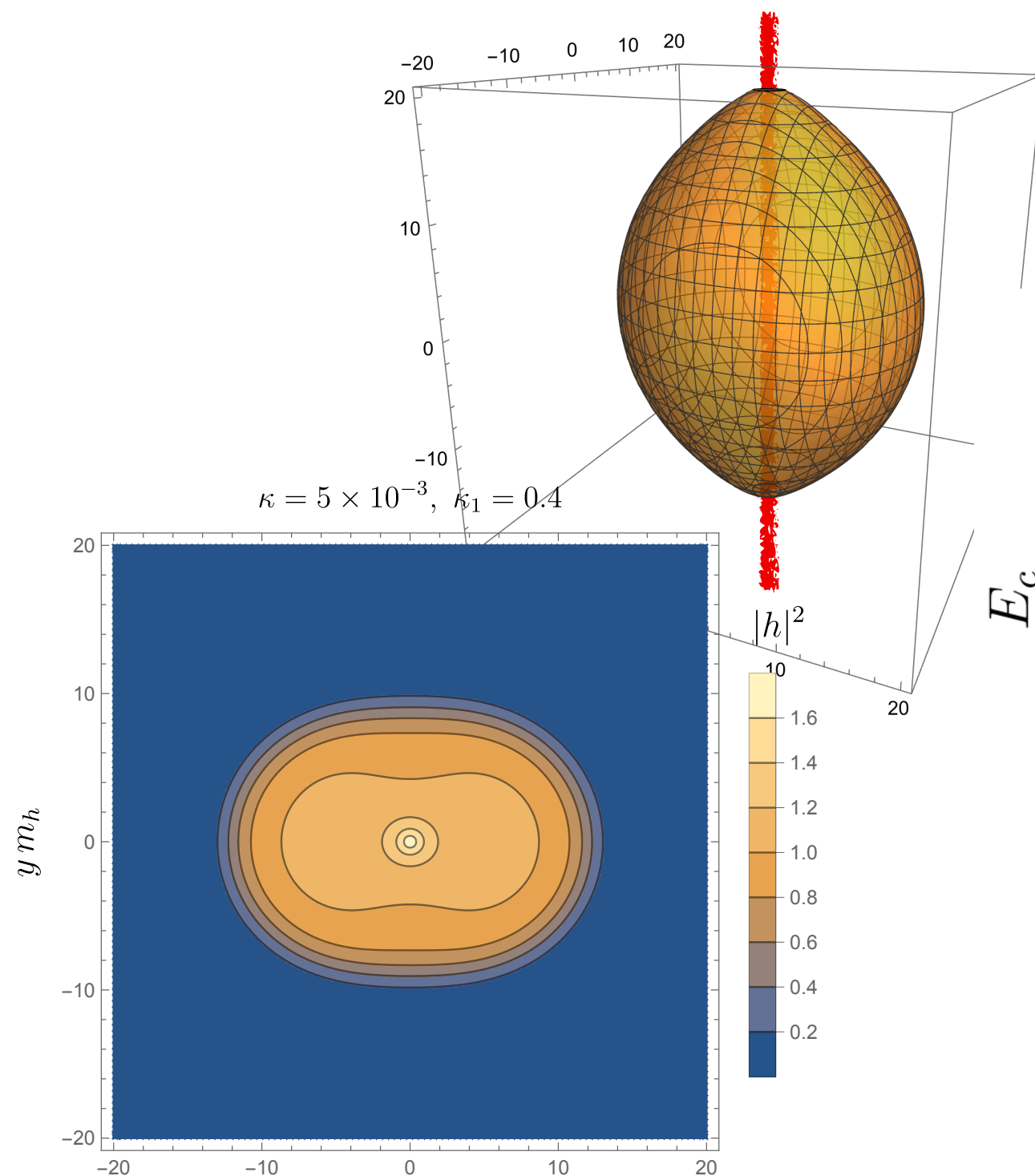


→ add $\tilde{\kappa} |\varphi|^2 |\Phi|^2$ like KSVZ setup [cf. Simone's talk]

[Blasi-Mariotti, 2405.08060]

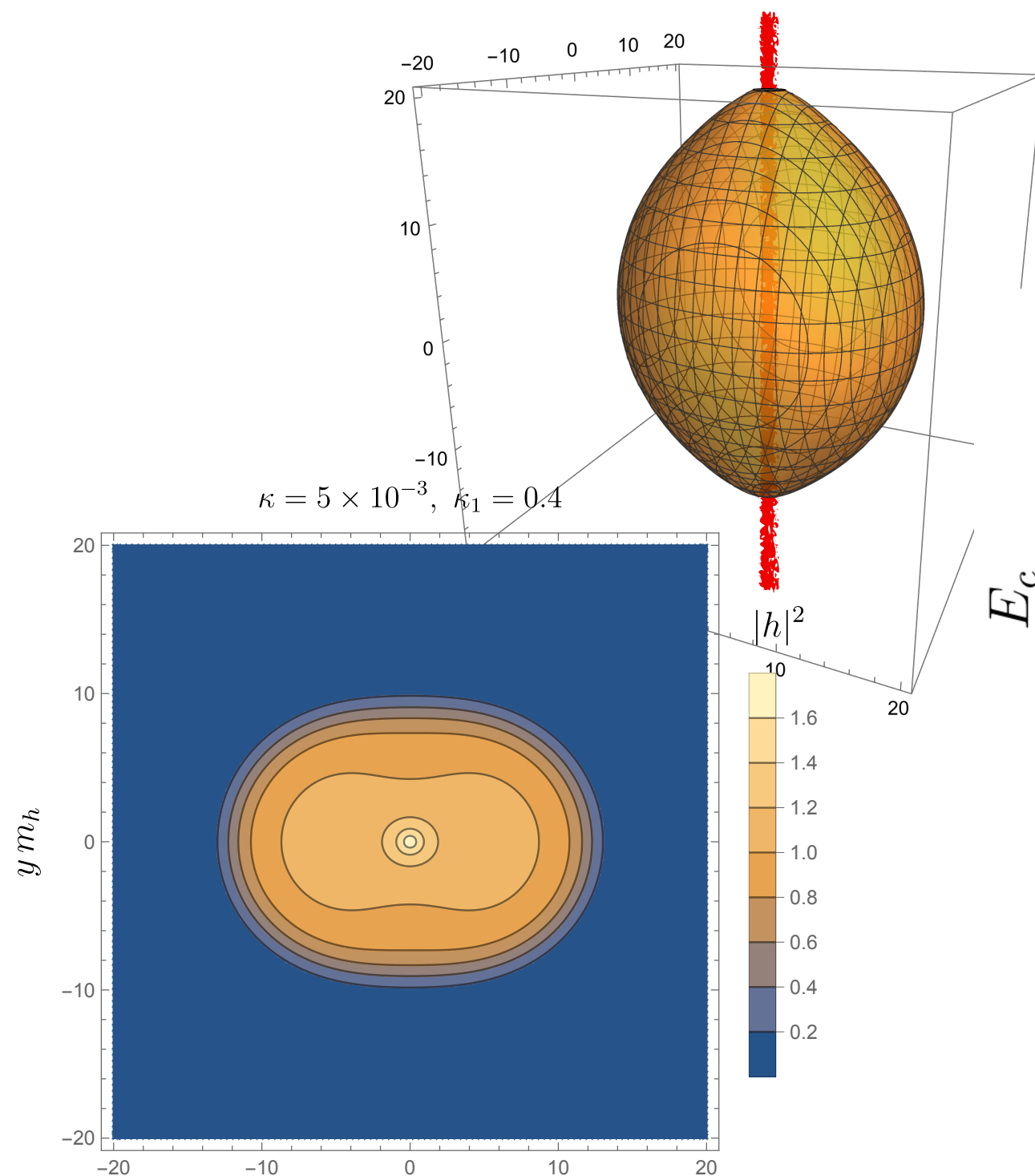
Critical bubble with KSVZ-like op.

- Comparison between seeded tunneling and normal one

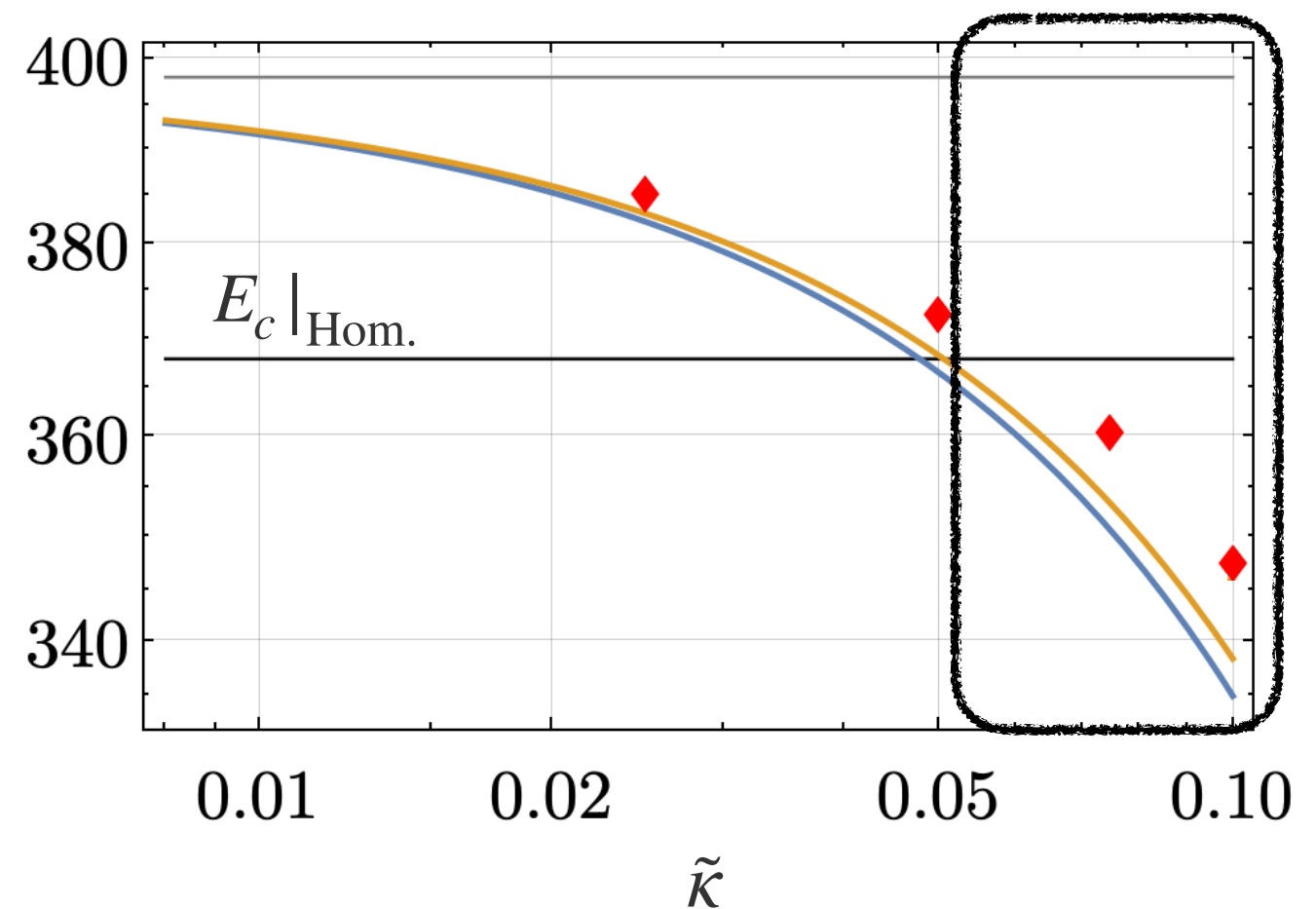


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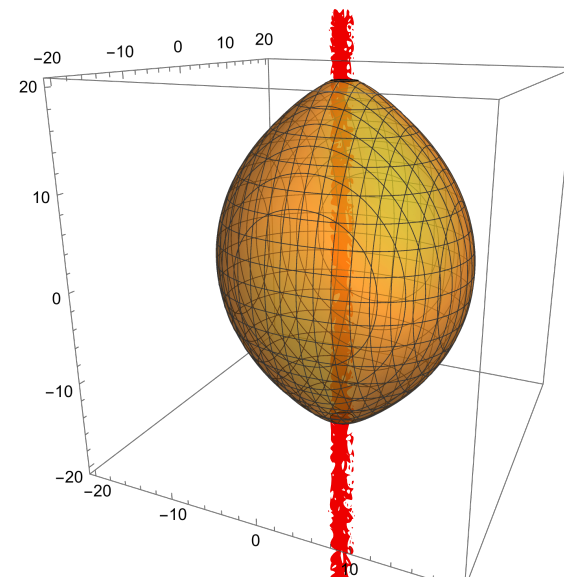
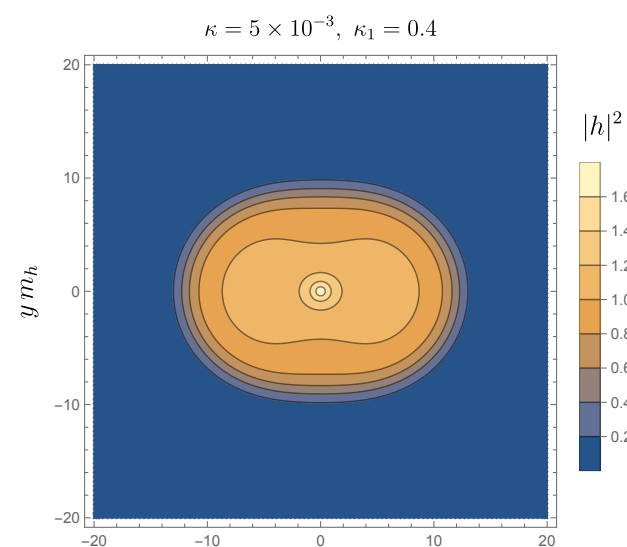


funny bubble is favored



Summary

- PT seeded by impurity is very interesting.
- A lot of future directions:
 - real DFSZ
 - funny shape of bubble \rightarrow enhancement of GW signal?
 - how affects EW baryogenesis?
 - What is consequence of DW production?



Backup

Perturbation theory

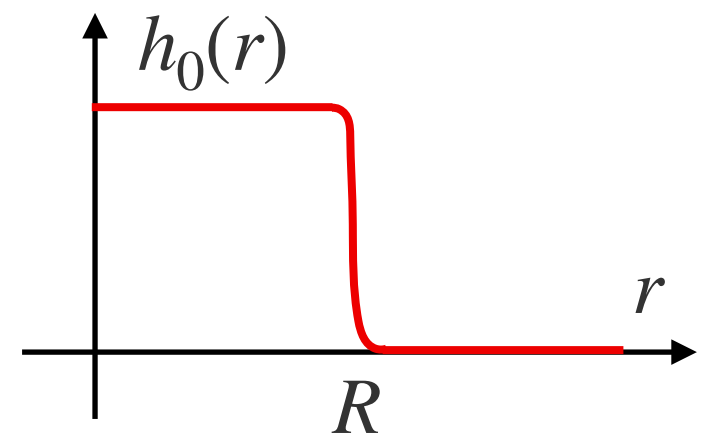
- It makes sense to consider perturbation w.r.t. κ :

higgs mode:

$$h(x) = h_0(x) + \delta h(x)$$

angular mode:

$$\phi(x) = \phi_0 + \delta\phi(x)$$



thin-wall as $\mathcal{O}(\kappa^0)$

Perturbation theory

- It makes sense to consider perturbation w.r.t. κ :

bounce solution w/ $\kappa = 0$

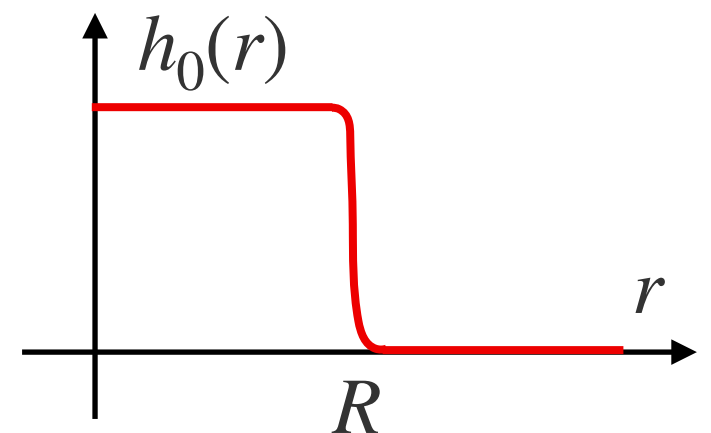
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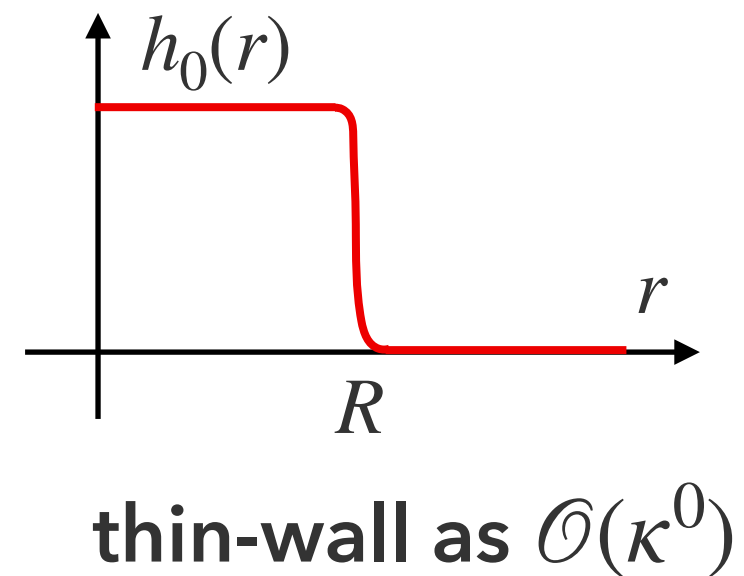


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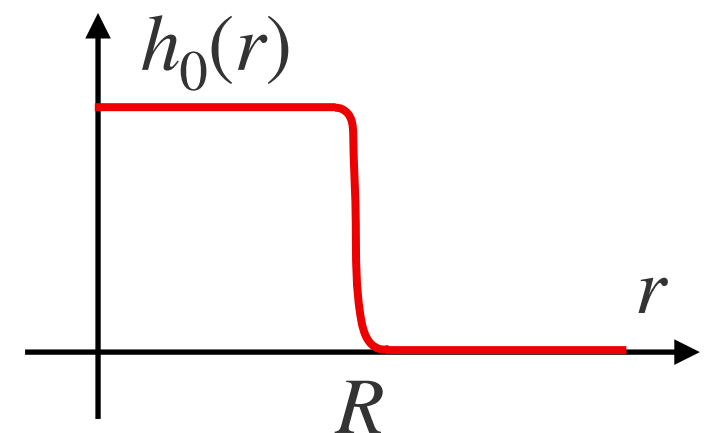
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thin-wall as $\mathcal{O}(\kappa^0)$

- EOM for $\mathcal{O}(\kappa^1)$ fluctuations:

higgs mode: $-\partial_i^2 \delta h + V''_{\text{EW}}(h_0) \delta h = \kappa h_0(r) f_a^2 \cos 2\theta$

angular mode: $-\partial_i [h_0^2(r) \partial_i \delta \phi] = -\kappa f_a^2 h_0^2(r) \sin 2\theta$

Linearized EOM for δh

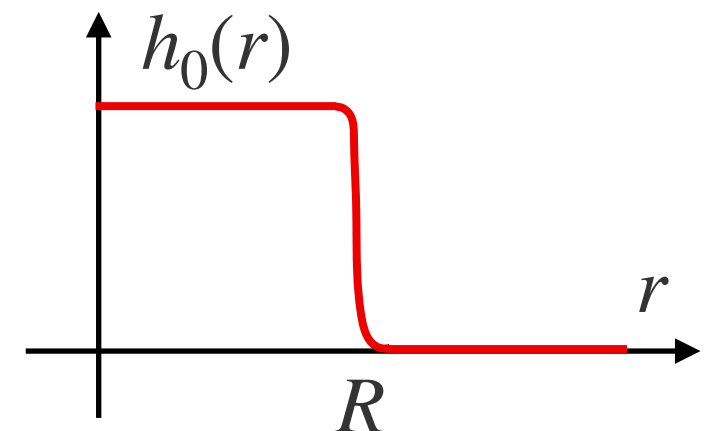
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$$\text{eigenvalue problem: } \hat{\mathcal{O}} f_\omega(x) = \omega^2 f_\omega(x)$$

- Spectral decomposition works well

$$\delta h = \sum_{\omega} \sum_{l,m} \frac{c_{\omega,l,m}}{\omega^2} Y_{lm}(\Omega) f_\omega(\rho, z)$$

$$\rightarrow c_{\omega,l,m} = \int d^3x S(x) Y_{l,m} f_\omega \propto \delta_{l2} \delta_{m2} \quad Y_{l,m}: \text{harmonic func}$$



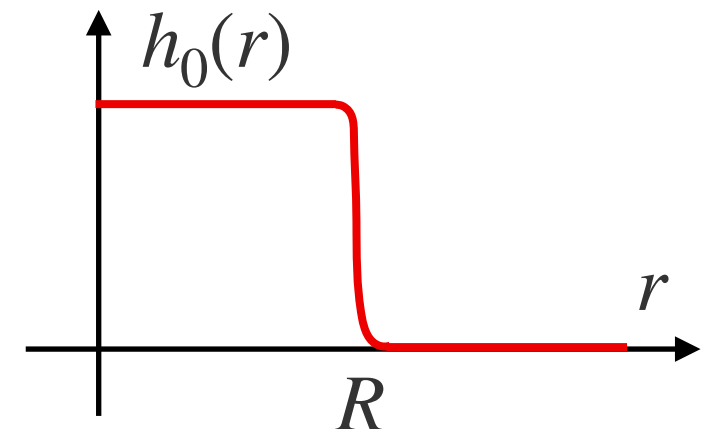
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- leading is quasi-zeromodes: translation of wall $\rightarrow \omega \sim 1/R$

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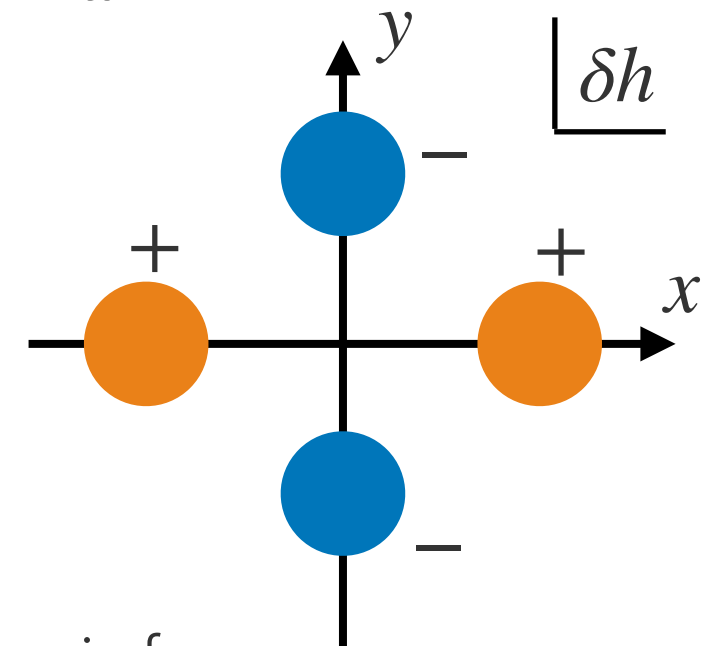
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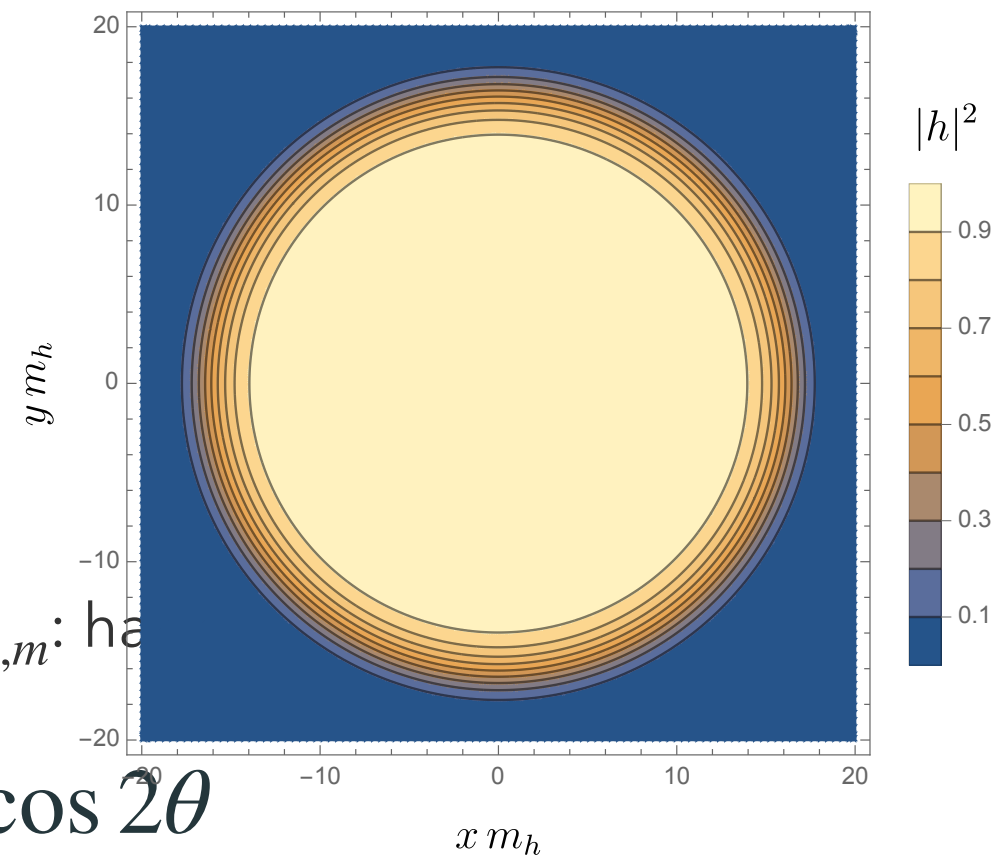
$$\text{eigenvalue problem: } \hat{\mathcal{O}} f_\omega(x) = \omega^2 f_\omega(x) \quad \kappa = 0$$

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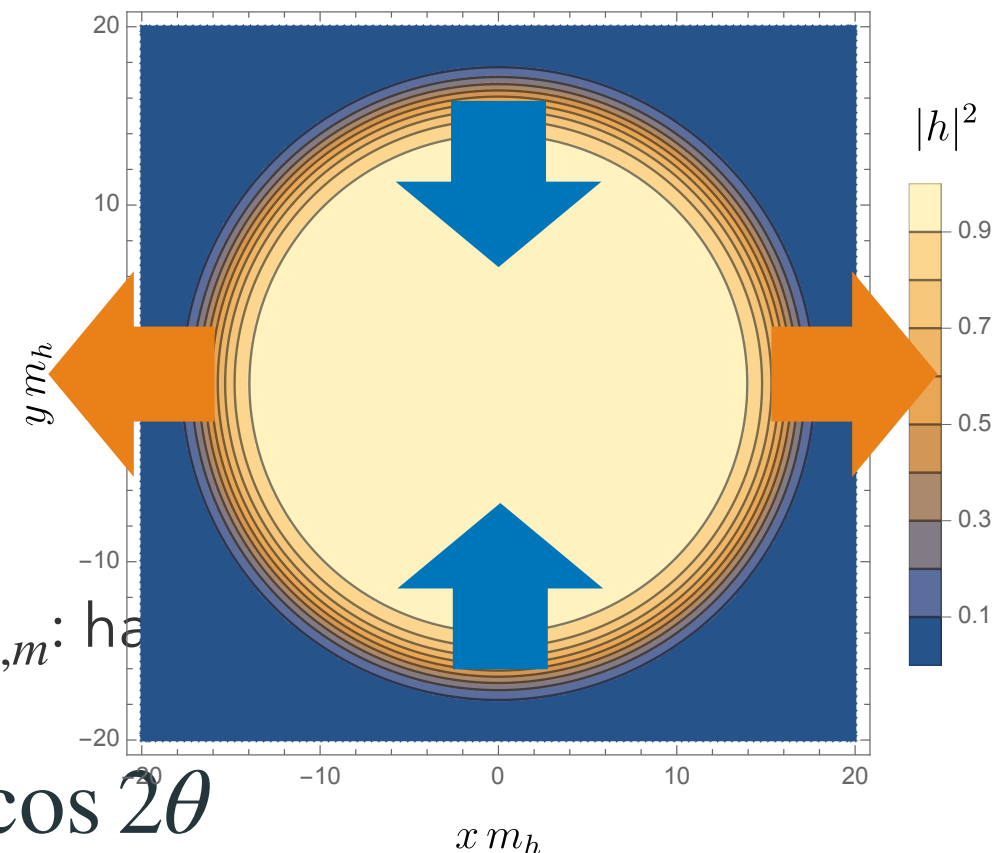
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- Green function method works well $Y_{l,m}(\Omega)$: harmonic func

$$\delta h = \sum_{\omega} \sum_{l,m} \frac{c_{\omega,l,m}}{\omega^2} Y_{lm}(\Omega) f_\omega(\rho, z) \rightarrow c_{\omega,l,m} = \int d^3x S(x) Y_{l,m} f_\omega$$

$$= \int d^3x' \sum_{\omega} \sum_{l,m} \frac{1}{\omega^2} Y_{lm}(\theta) f_\omega(x) \underline{Y_{l,m}^*(\theta') f_\omega^*(x') S(x')}$$

Green function

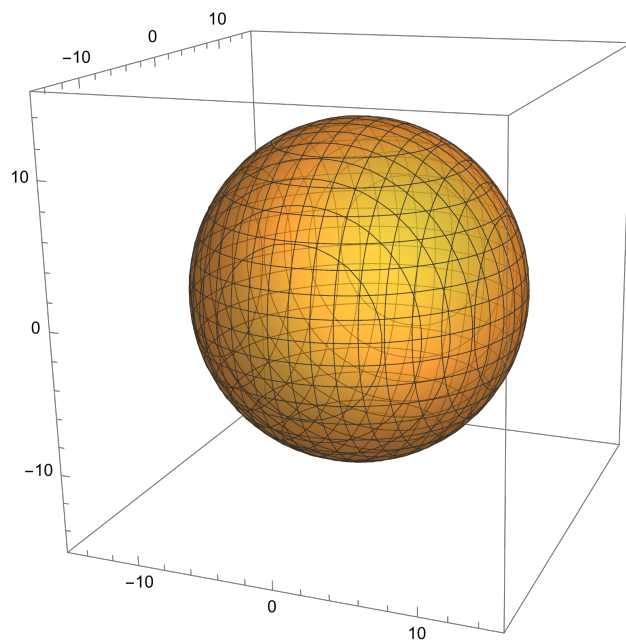
Perturbation theory

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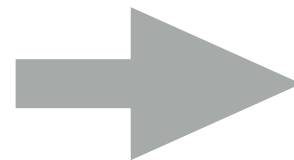
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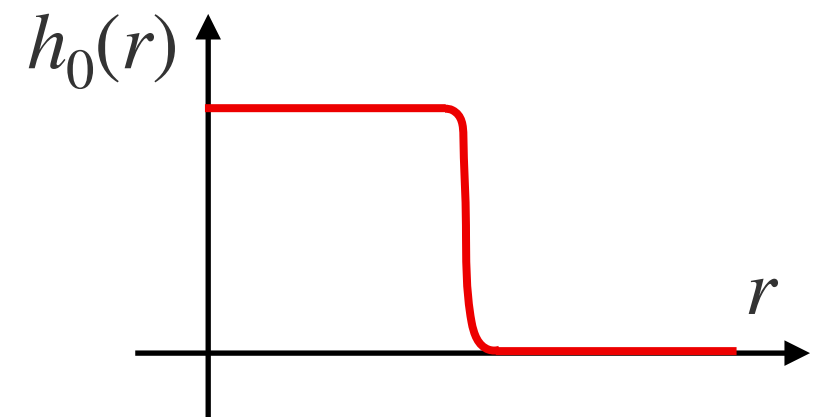


bounce up to $\mathcal{O}(\kappa^1)$

Linearized EOM

- EOM for $\mathcal{O}(\kappa^1)$ fluctuations:

$$\begin{aligned} -\partial_i^2 \delta h + V''_{\text{EW}}(h_0) \delta h &= \kappa h_0(r) f_a^2 \cos 2\theta \\ &\equiv \hat{\mathcal{O}} \delta h \end{aligned}$$



$$-\partial_i [h_0^2(r) \partial_i \delta \phi] = -\kappa f_a^2 h_0^2(r) \sin 2\theta$$

thin-wall as BG

- Eigenmode decomposition:

$$\delta h = \sum_{\omega} \sum_{l,m} \frac{c_{\omega,l,m}}{\omega^2} Y_{lm}(\theta) f_{\omega}(x)$$

coefficient of expansion
harmonic funcs

eigenvalue of "kernel": $\hat{\mathcal{O}} f_{\omega}(x) = \omega^2 f_{\omega}(x)$

solving Eq \Leftrightarrow determining $c_{\omega,l,m}$