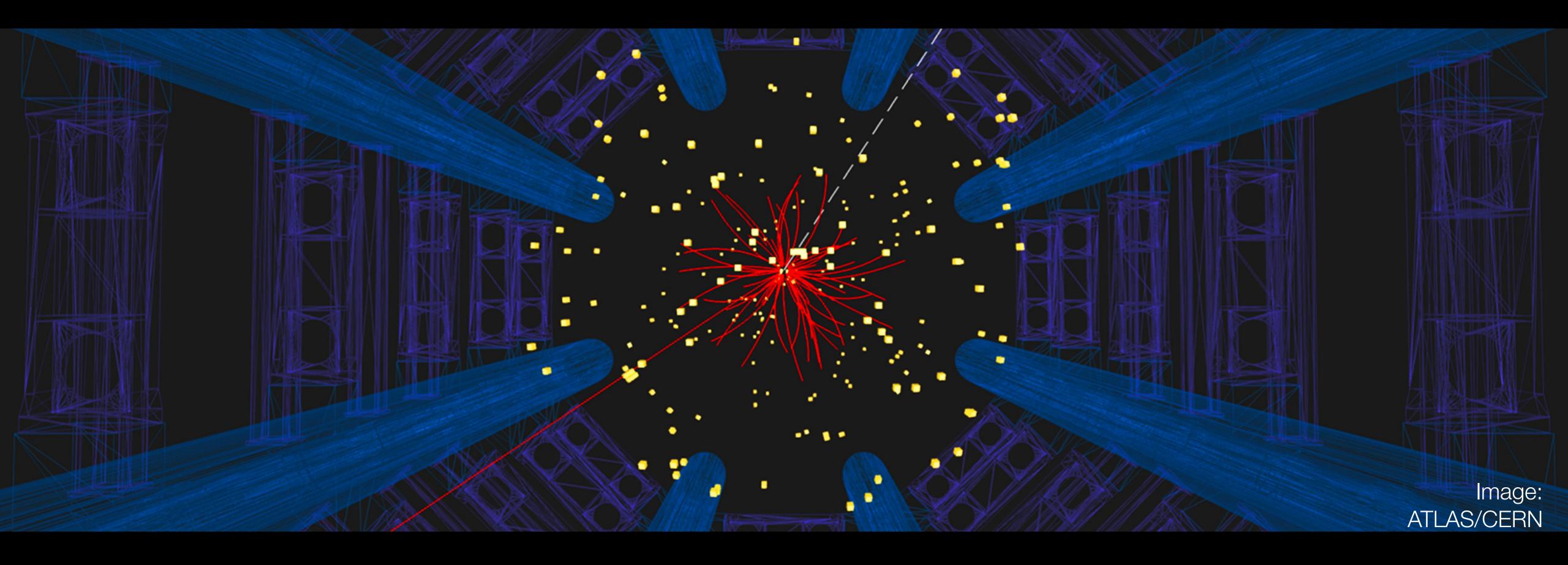
LHC Physics - BSM searches & beyond



Lydia Beresford

DESY Summer Student Lectures 05.08.25



Physics Goals of the LHC

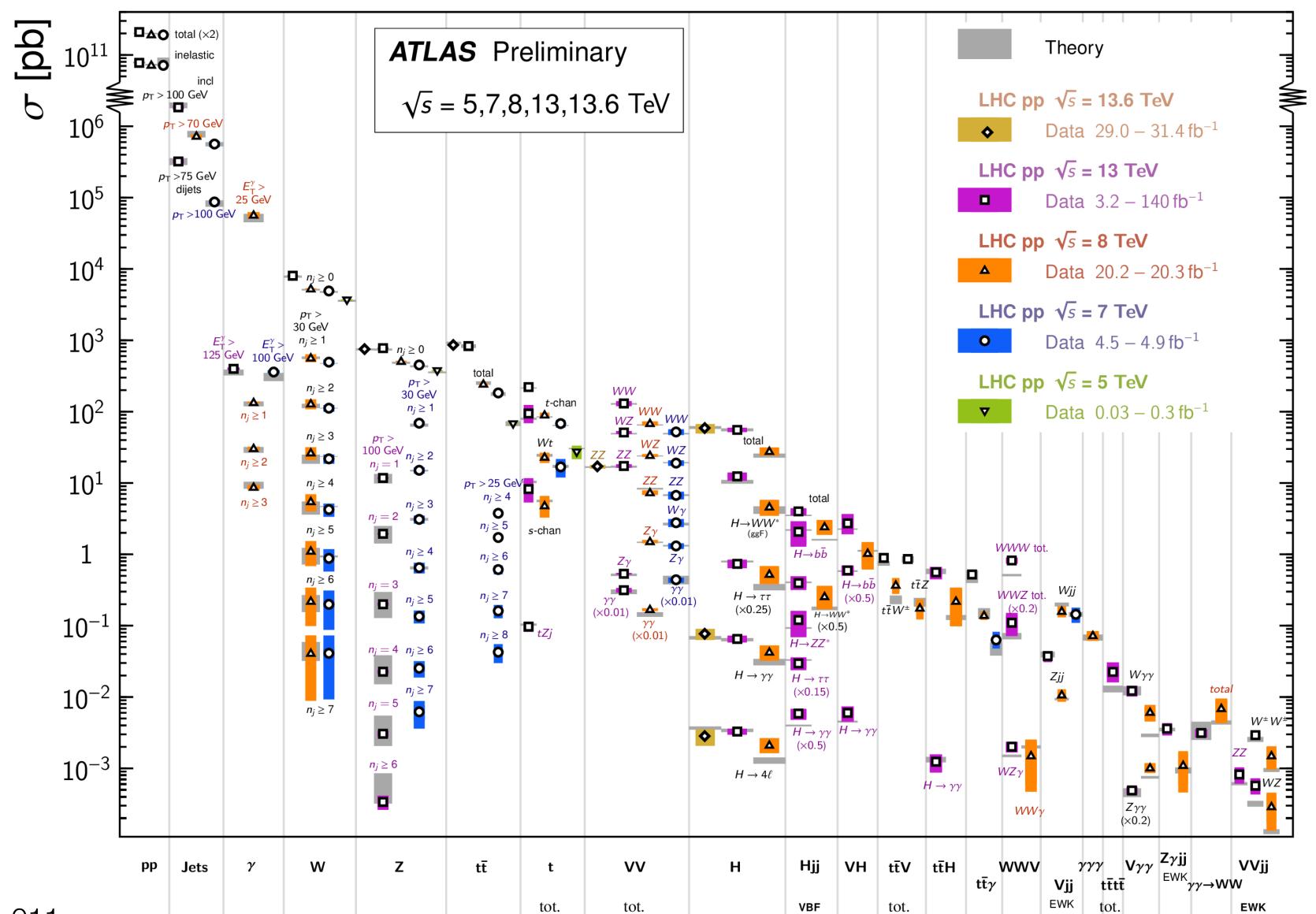


Search for New Physics

Standard Model success

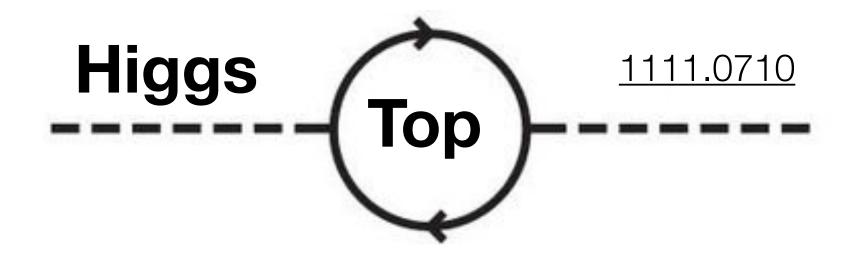
Standard Model Production Cross Section Measurements

Status: October 2023



Shortcomings of the Standard Model

- No explanation for Dark Matter or Dark Energy
- Gravity is not included in the Standard Model
- Doesn't explain why the Higgs is so light (hierarchy problem)



 Does not explain matter-antimatter asymmetry in the universe

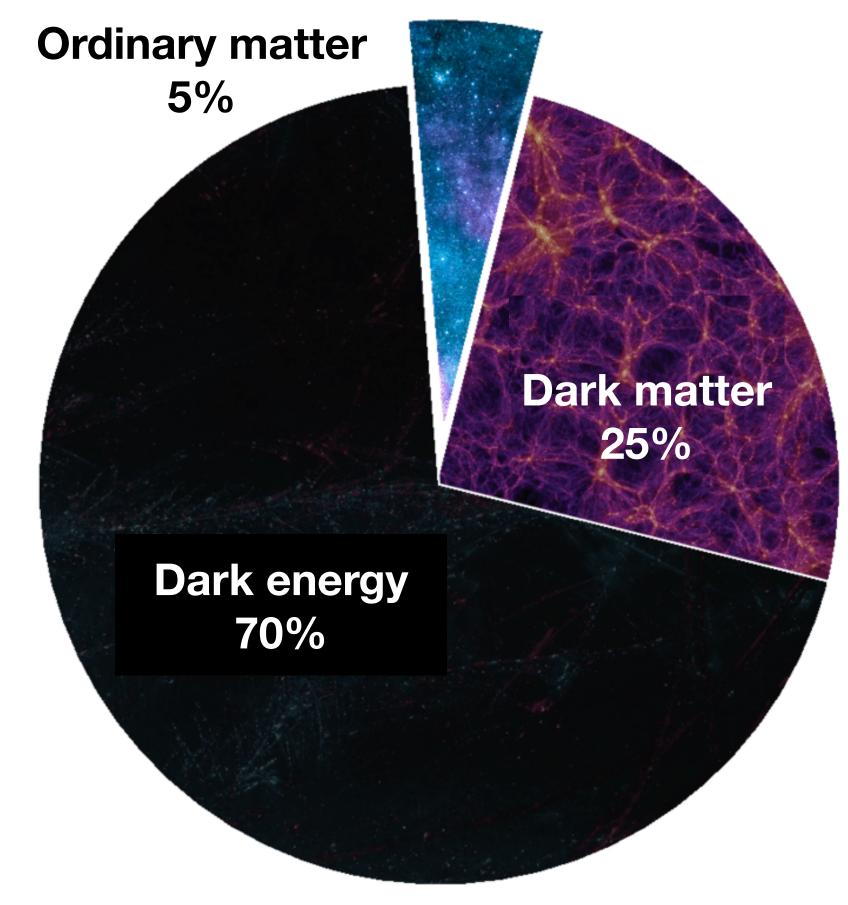
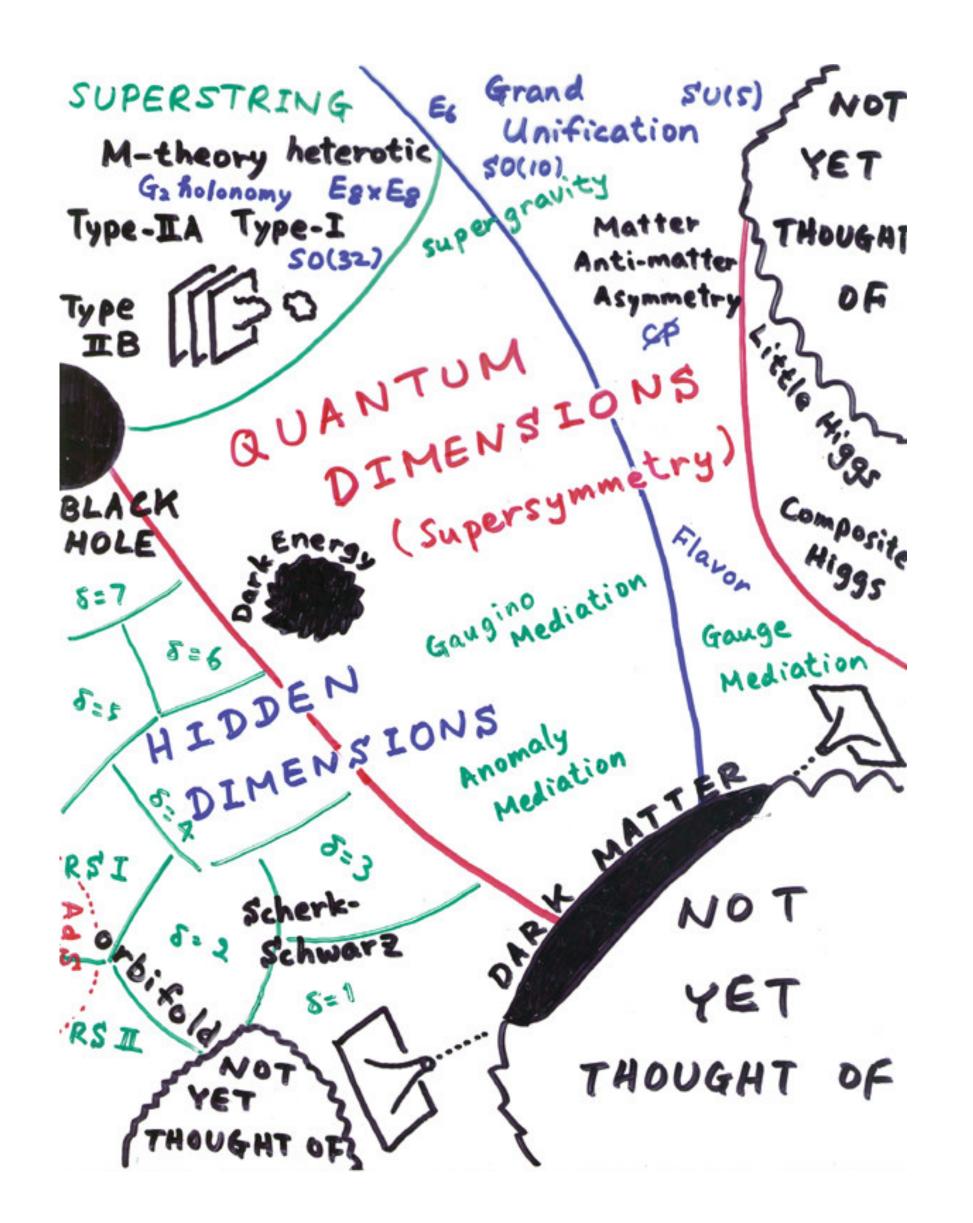


Image: adapted from Florian Wolz

Energy density of the universe

Landscape of proposed New Physics scenarios is vast!



Two ways of approaching this

Model driven

Start from a specific theory prediction

→ design and optimise for that specific signal

Signature driven

Look for deviation from the SM anywhere

→ look at specific final states (dijet, high MET)

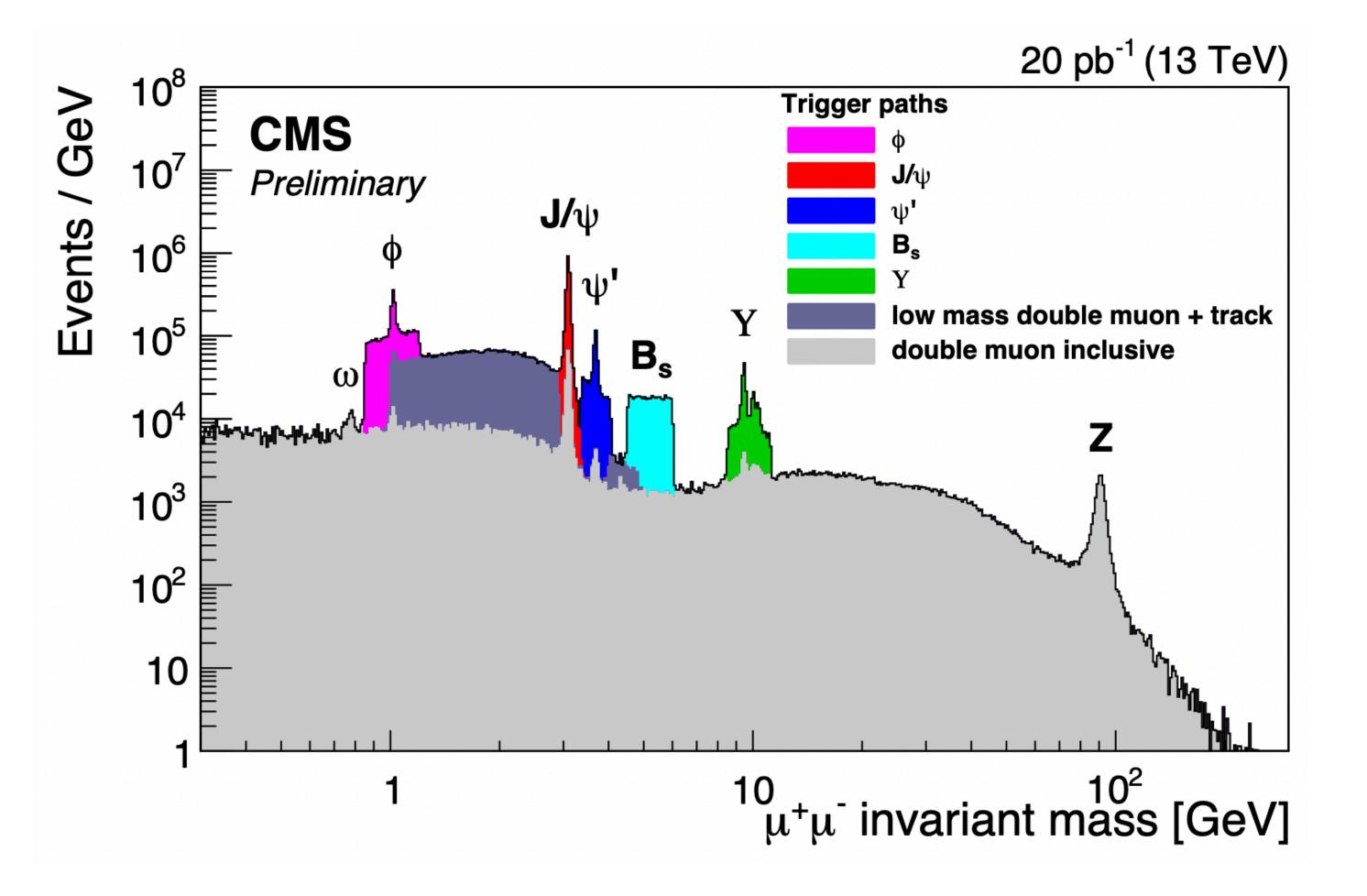
Both strategies are followed at the LHC

Crucial: excellent understanding of Standard Model backgrounds is needed!

Resonance searches

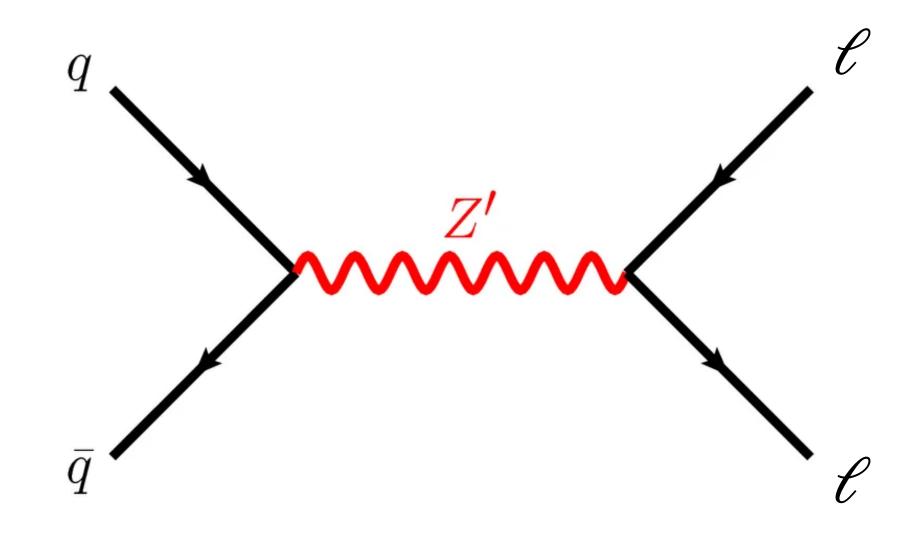
Many models predict the existence of additional, so-far undiscovered particles

→ Would likely find them through their decay products

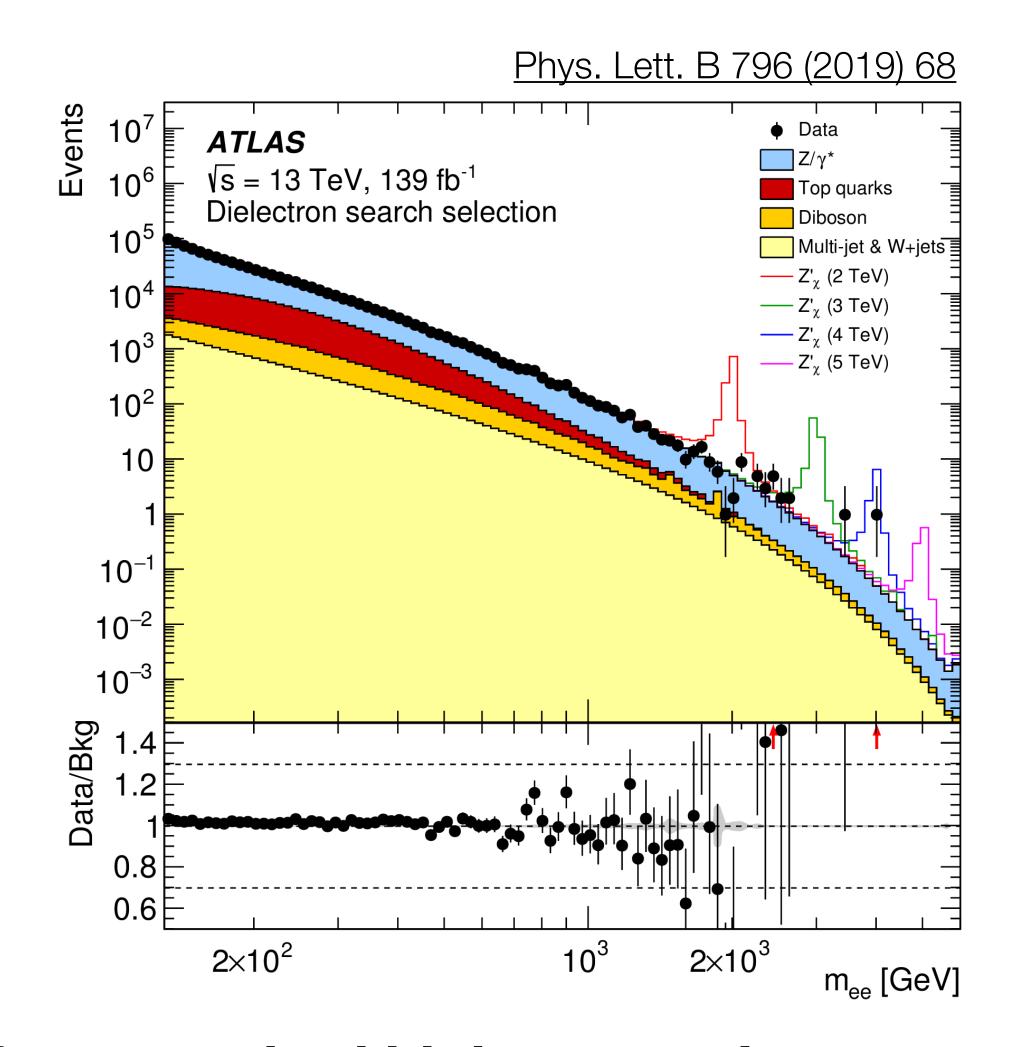


Dilepton resonances

Heavy partners of the known bosons e.g. $Z' \to \ell \ell$ Occur in e.g. grand unified theories (GUT)



Now probing masses up to 4 TeV

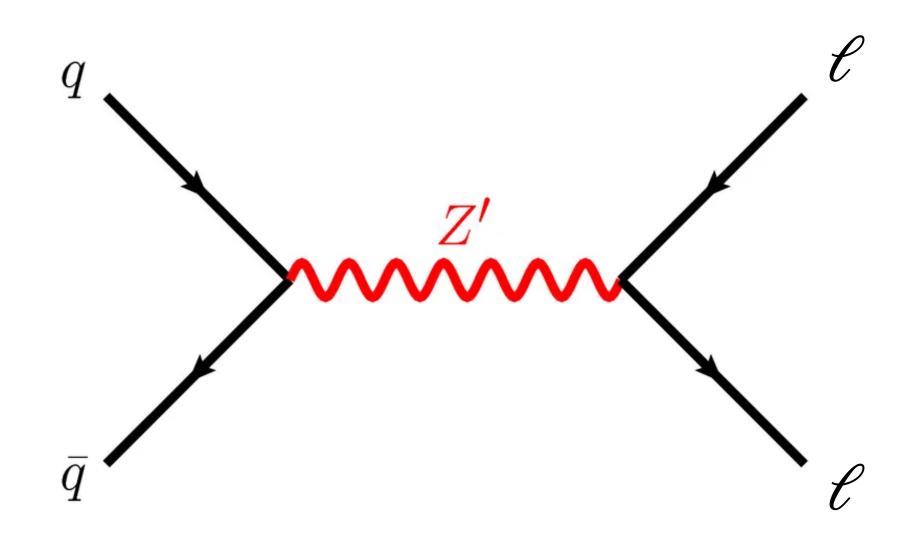


LHC strength: high CM energy → sensitivity to so far un-probed high masses!

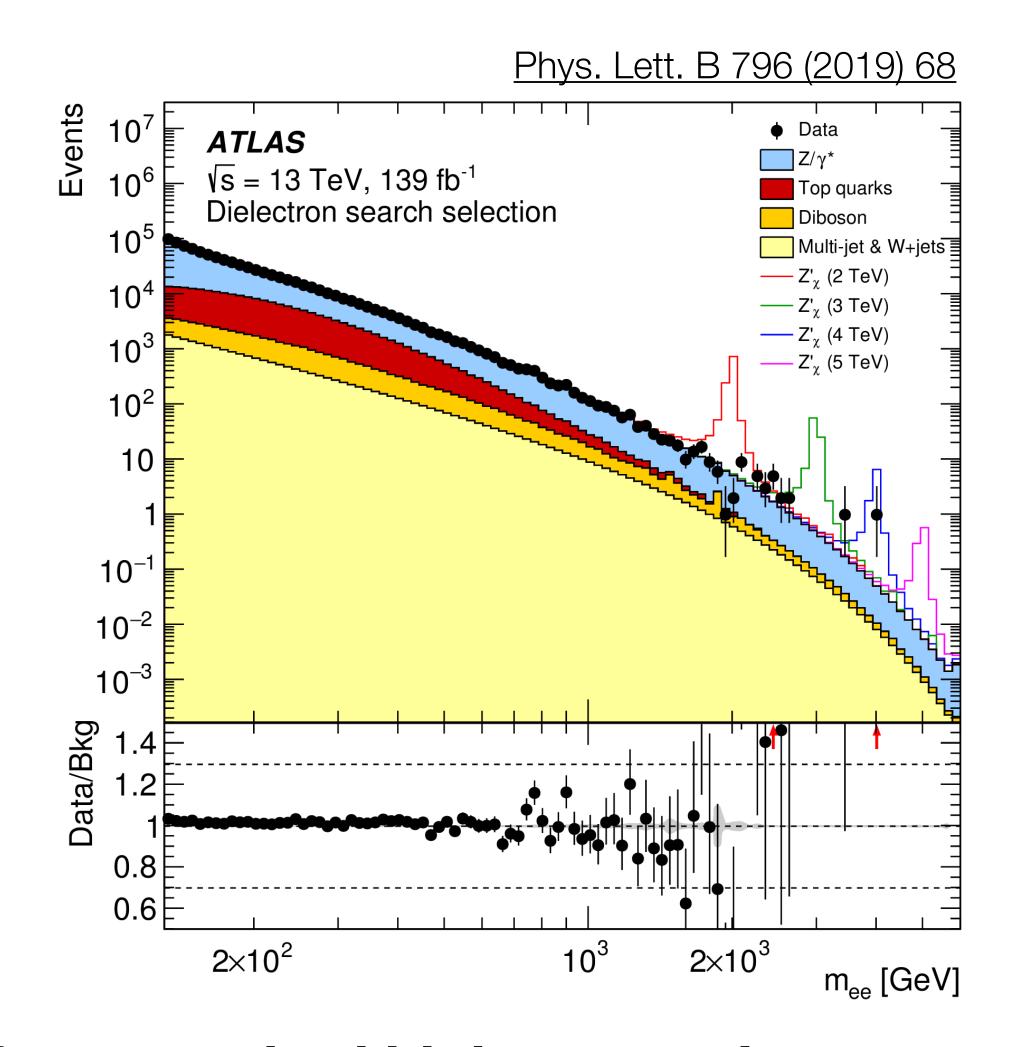
Dilepton resonances

Quiz: what is a grand unified theory?

Heavy partners of the known bosons e.g. $Z' \to \ell \ell$ Occur in e.g. grand unified theories (GUT)



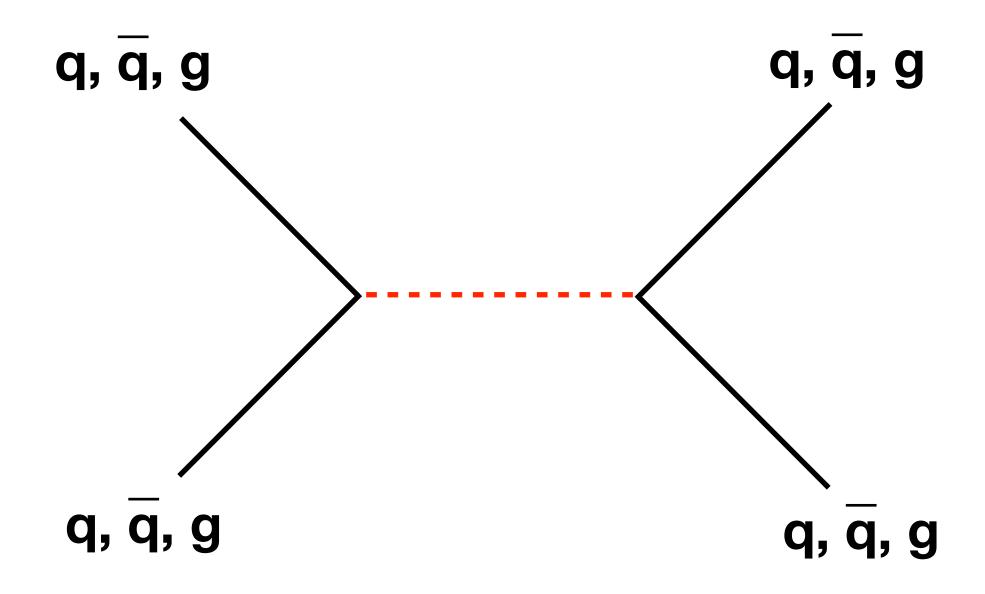
Now probing masses up to 4 TeV

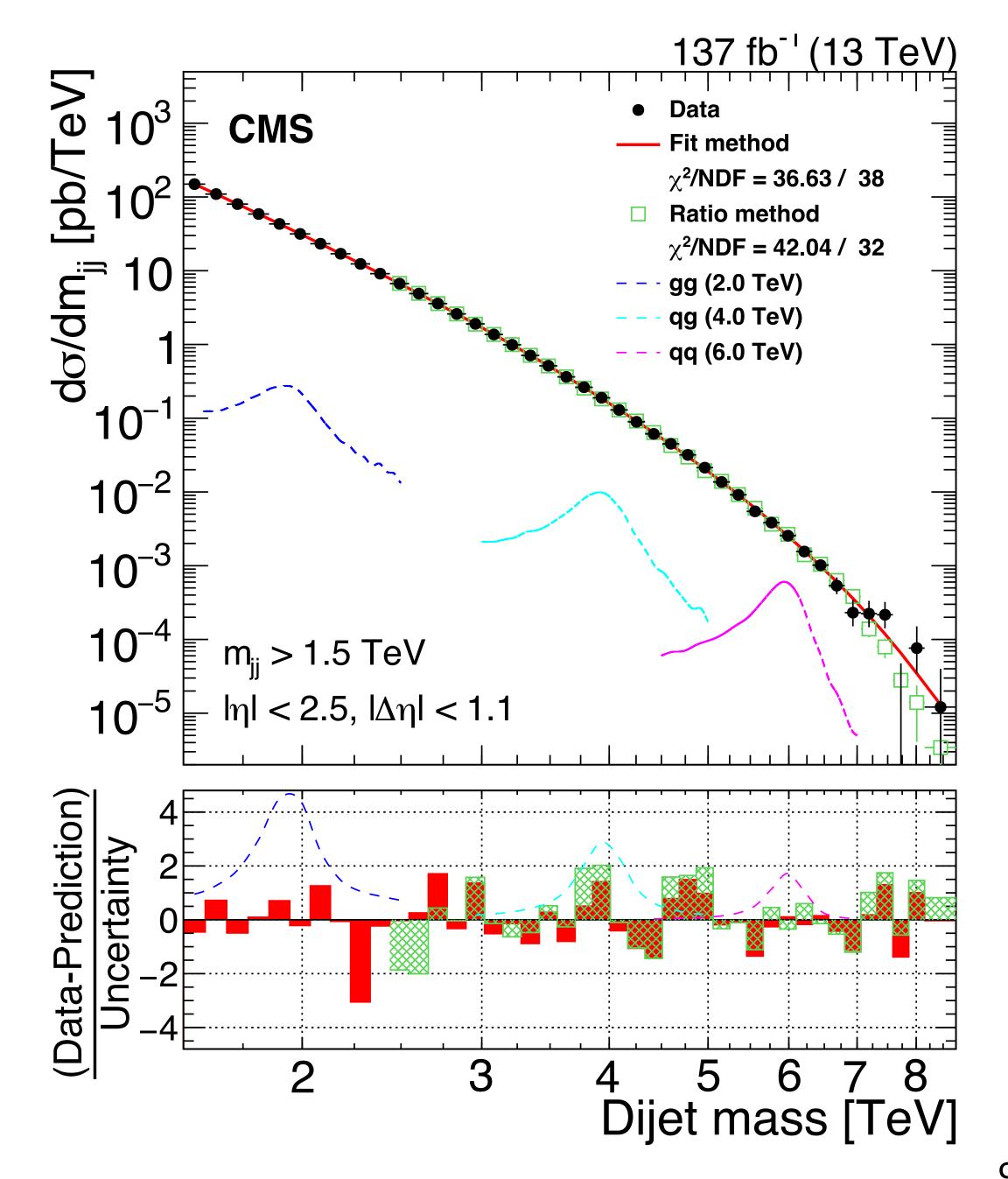


LHC strength: high CM energy → sensitivity to so far un-probed high masses!

Dijet resonances

Several new physics models predict heavy resonances that decay into dijets (qq, qg or gg)



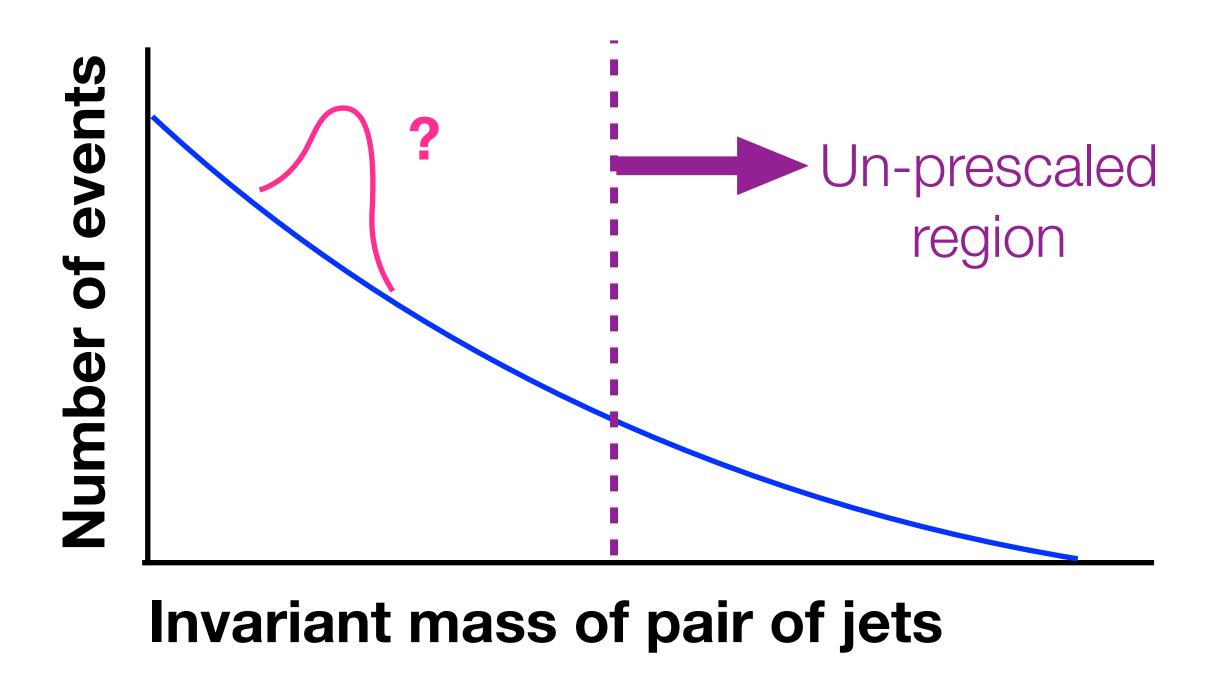


What about lower masses?

Experimental Challenge

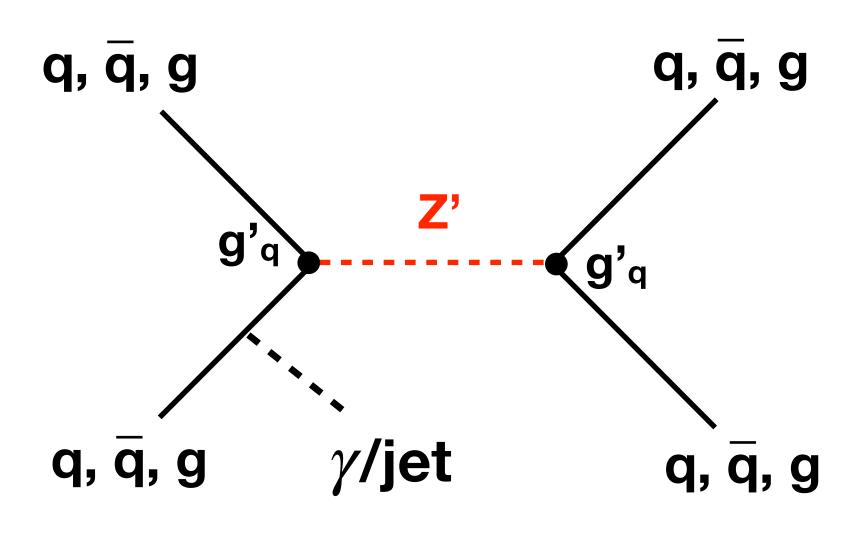
LHC produces huge amount of jet events at low mass

- → Single jet triggers highly prescaled at low masses
- → New techniques needed

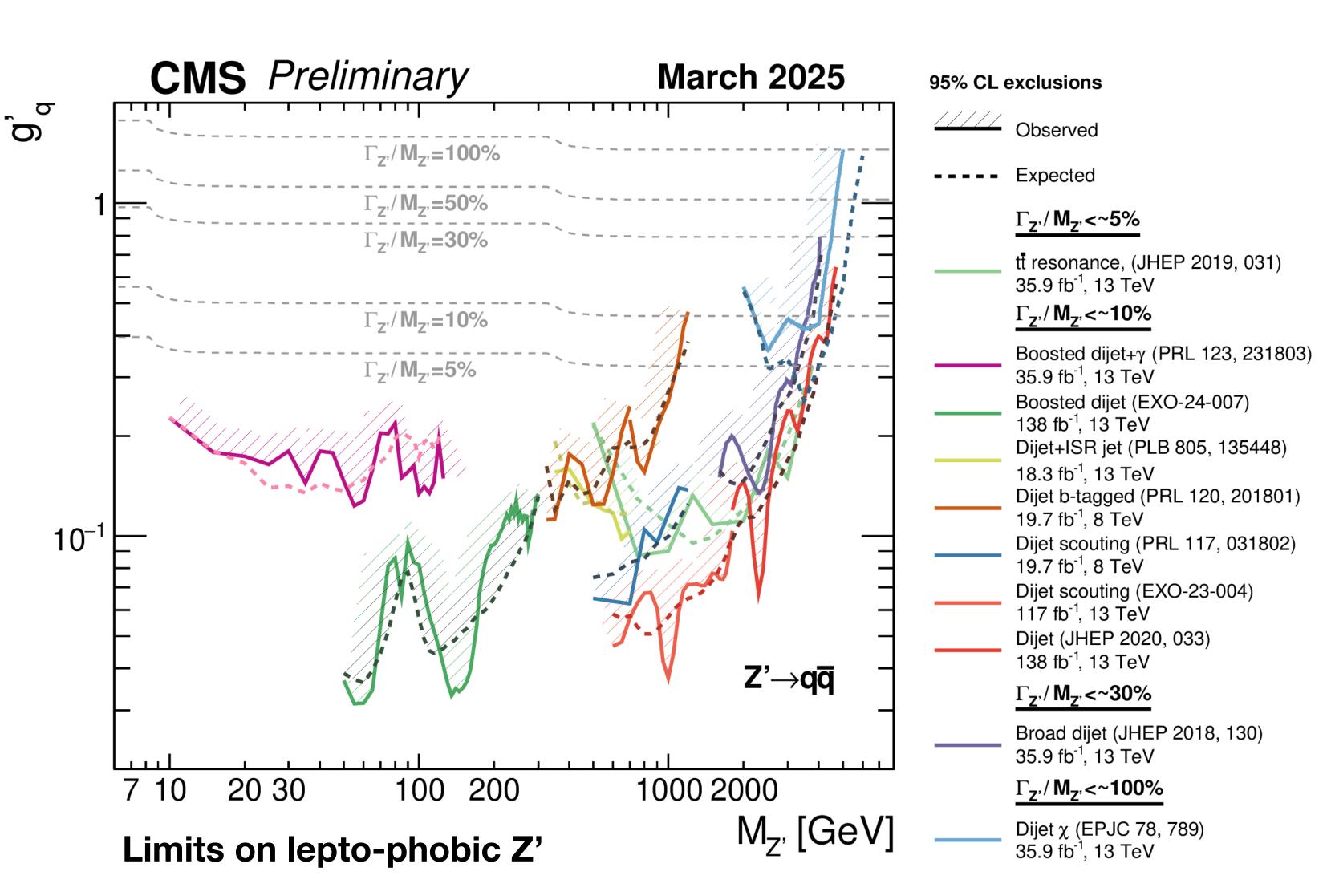


What about lower masses?

Trigger on radiated photon or jet

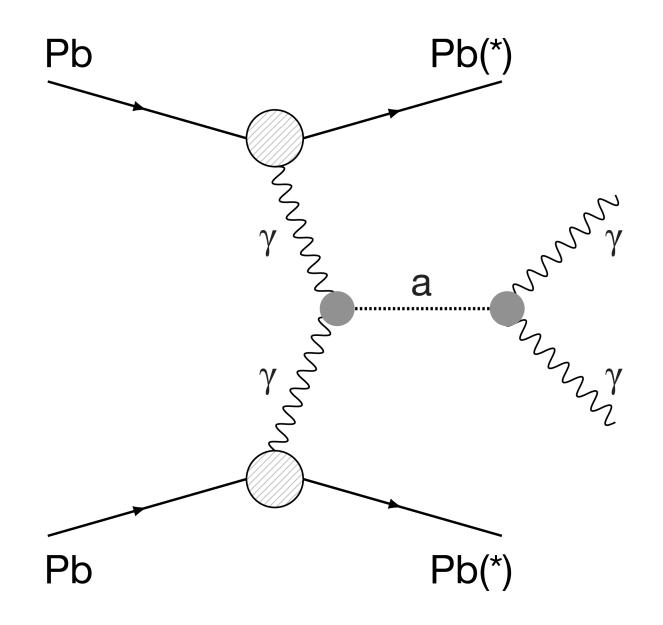


Reconstruct resonance as a single large radius jet & use sub-structure

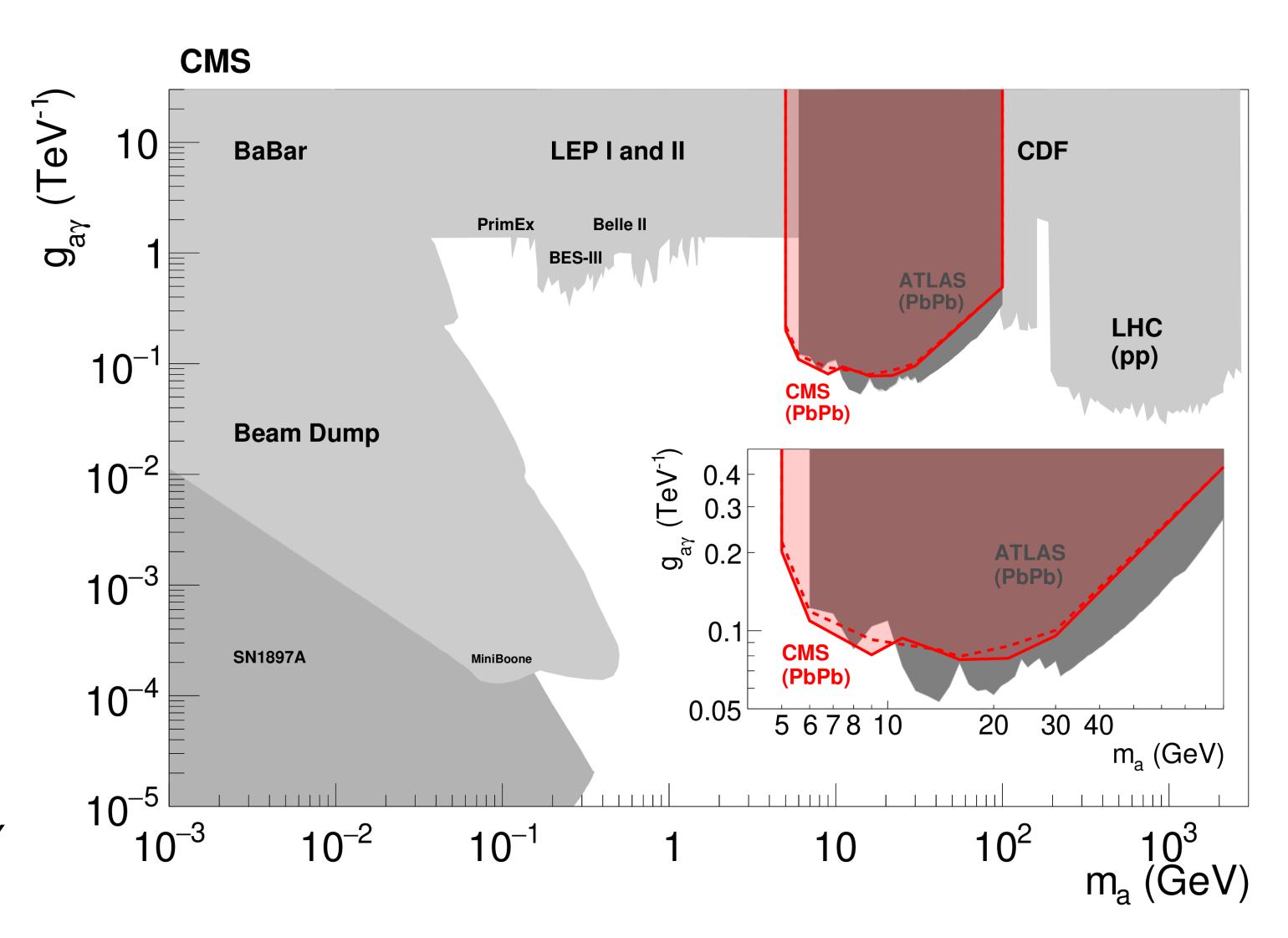


What if the new resonance only couples to photons?

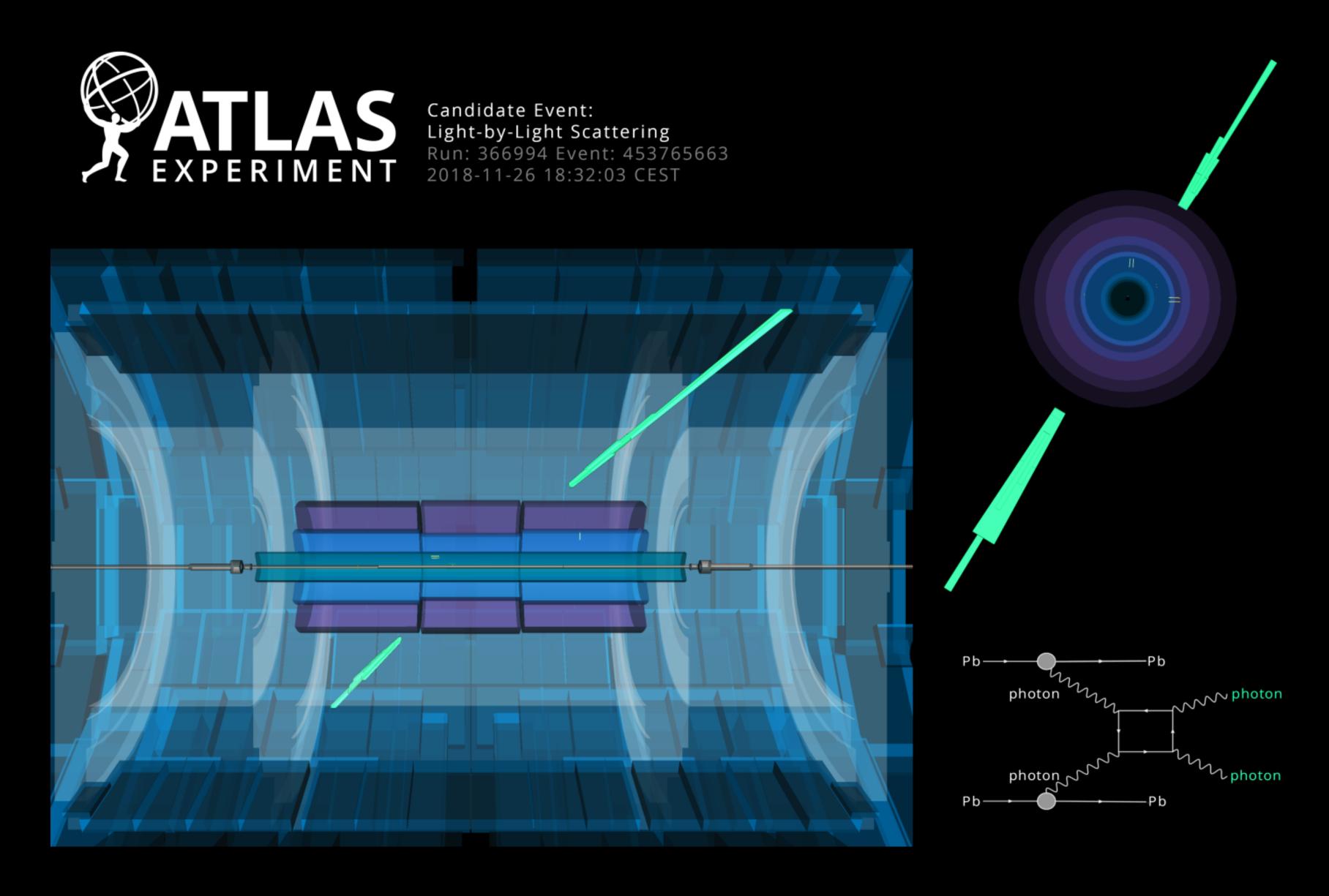
Axion-Like-Particle (ALP)



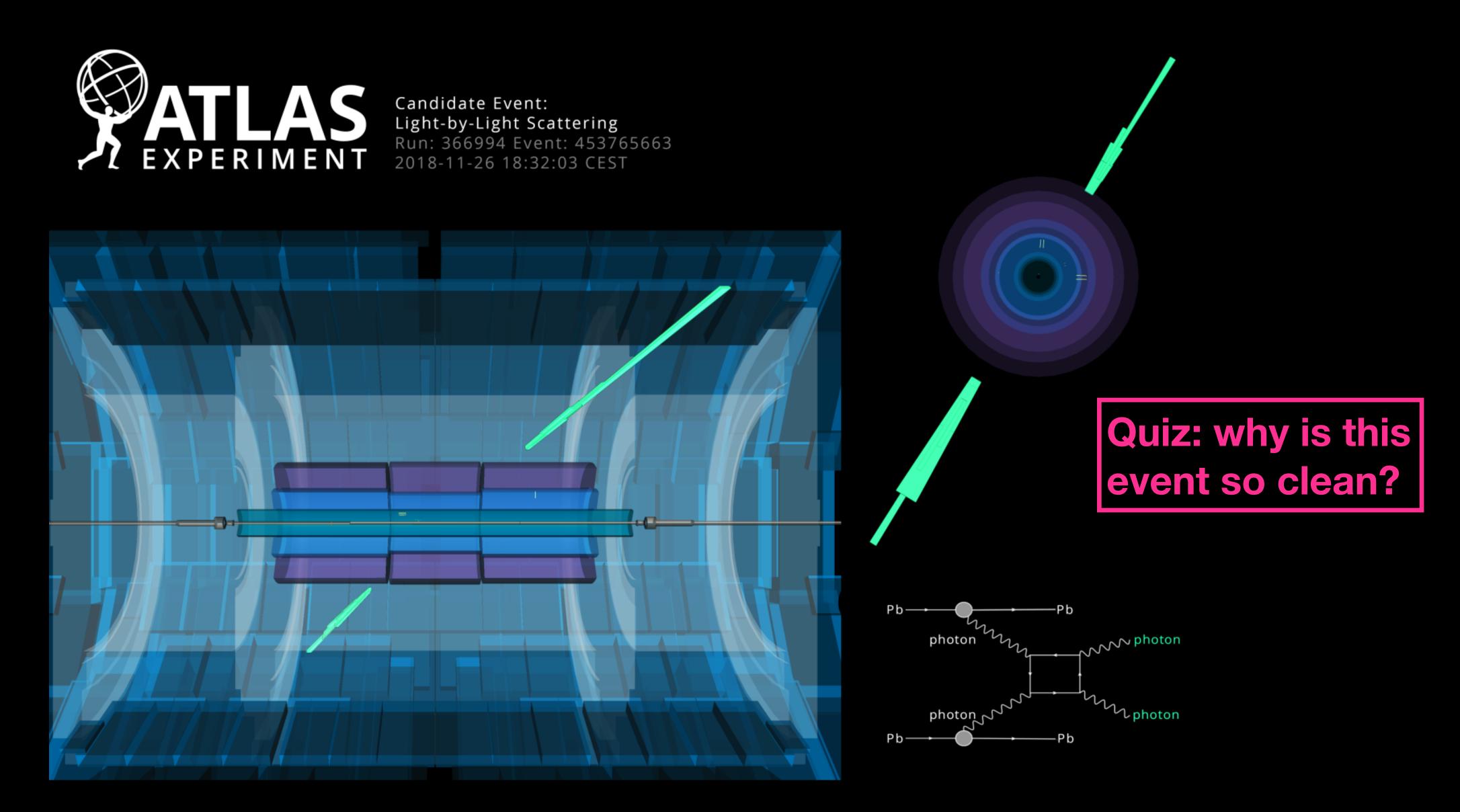
Look for narrow resonance in $m_{\gamma\gamma}$



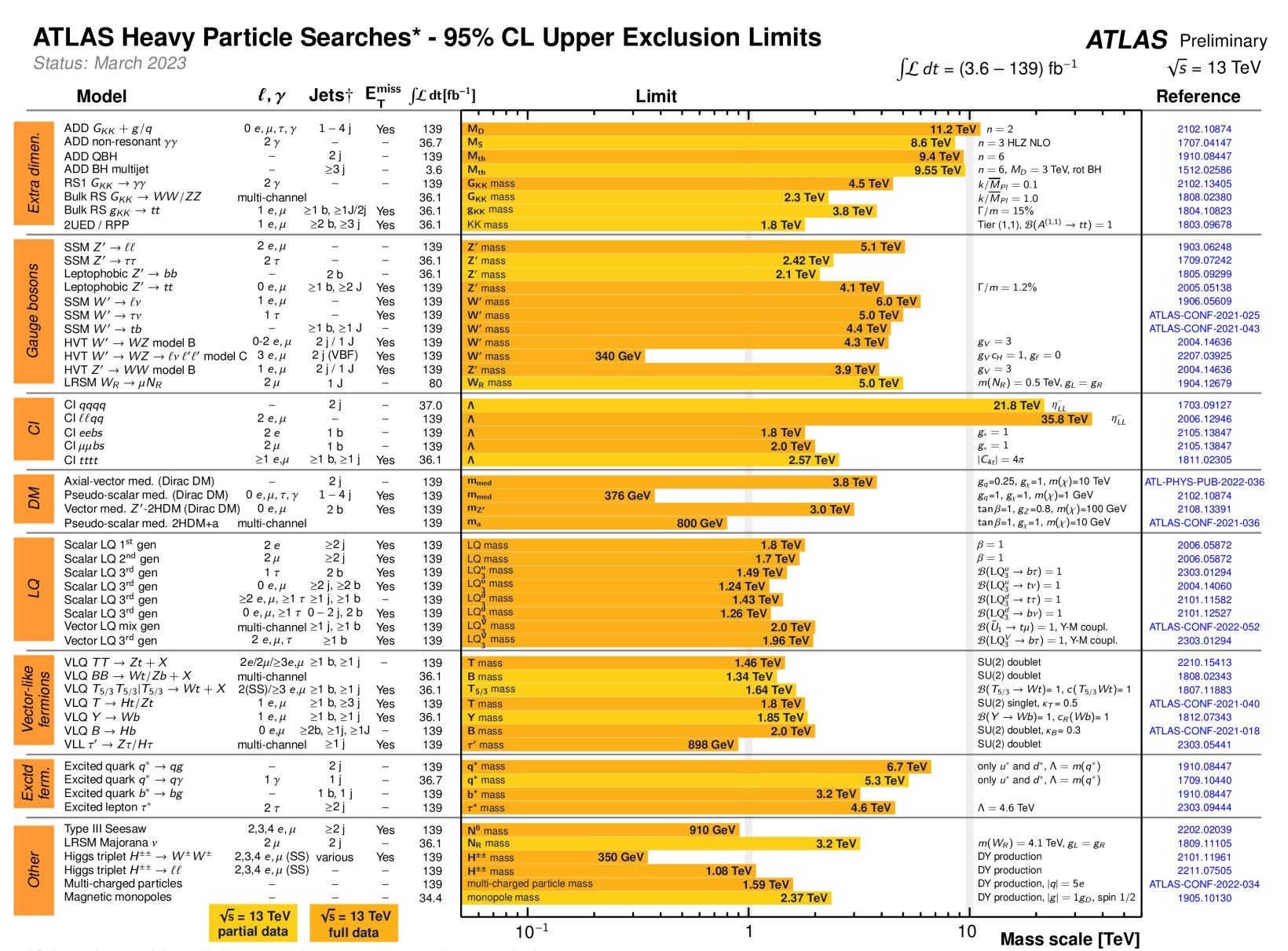
Pb+Pb Light-by-light scattering candidate event



Pb+Pb Light-by-light scattering candidate event



Exotics search summary plot - a subset

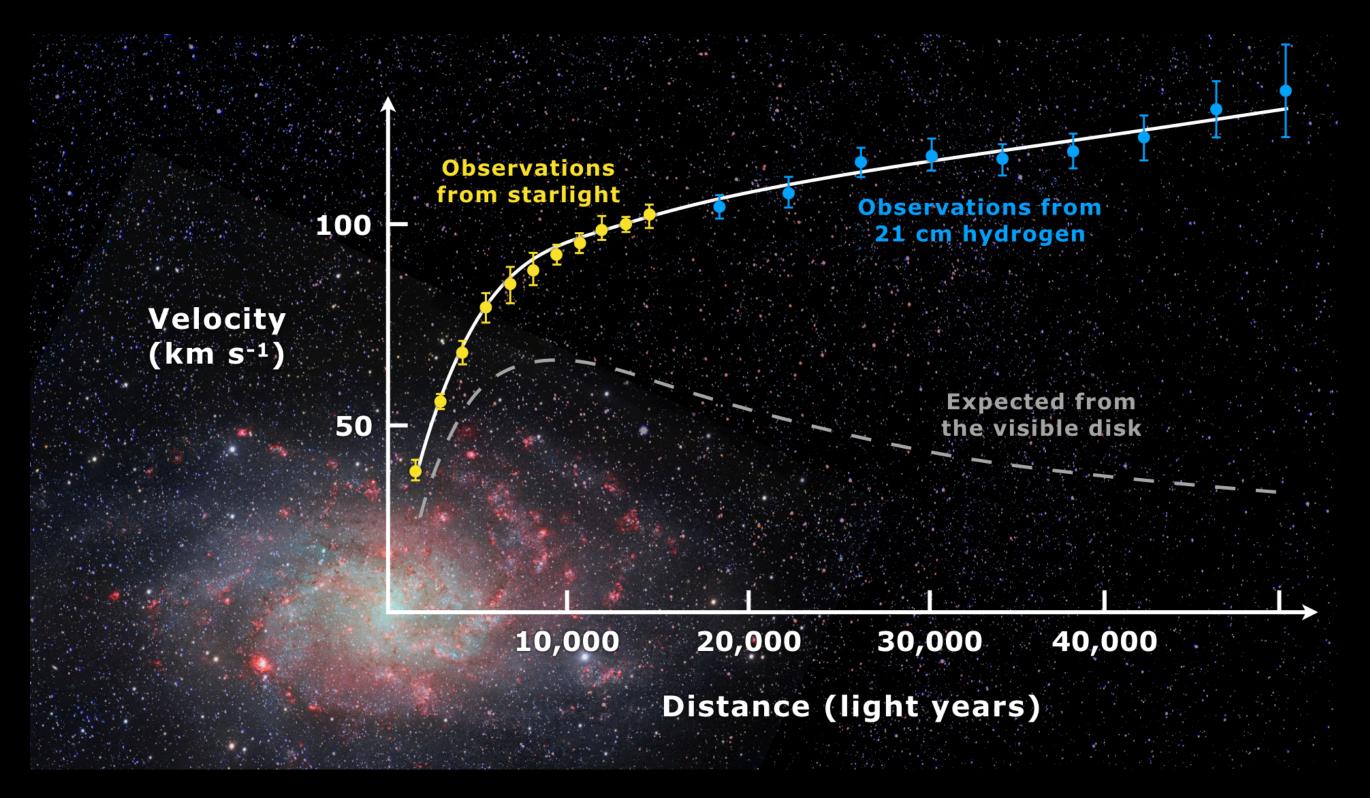


^{*}Only a selection of the available mass limits on new states or phenomena is shown.

[†]Small-radius (large-radius) jets are denoted by the letter i (J).

Dark Matter

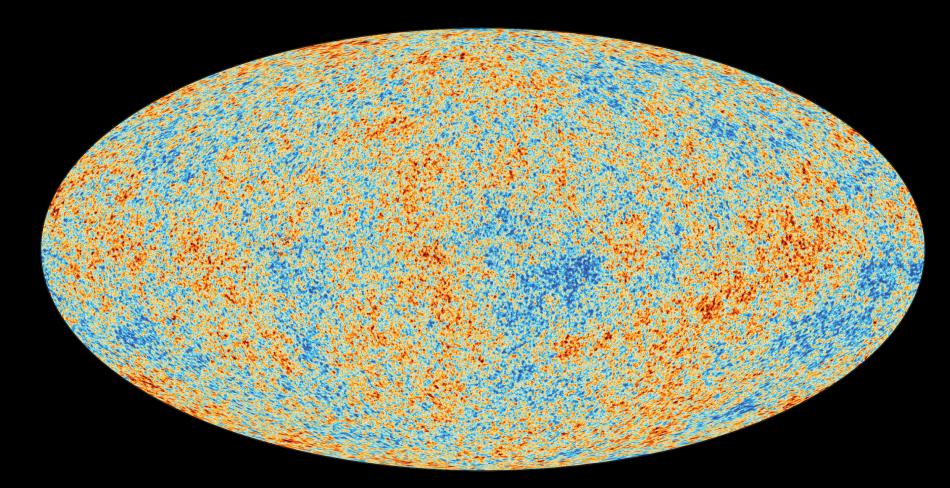
Astrophysical observations give strong evidence for a new kind of matter that interacts gravitationally



Rotational curves of galaxies



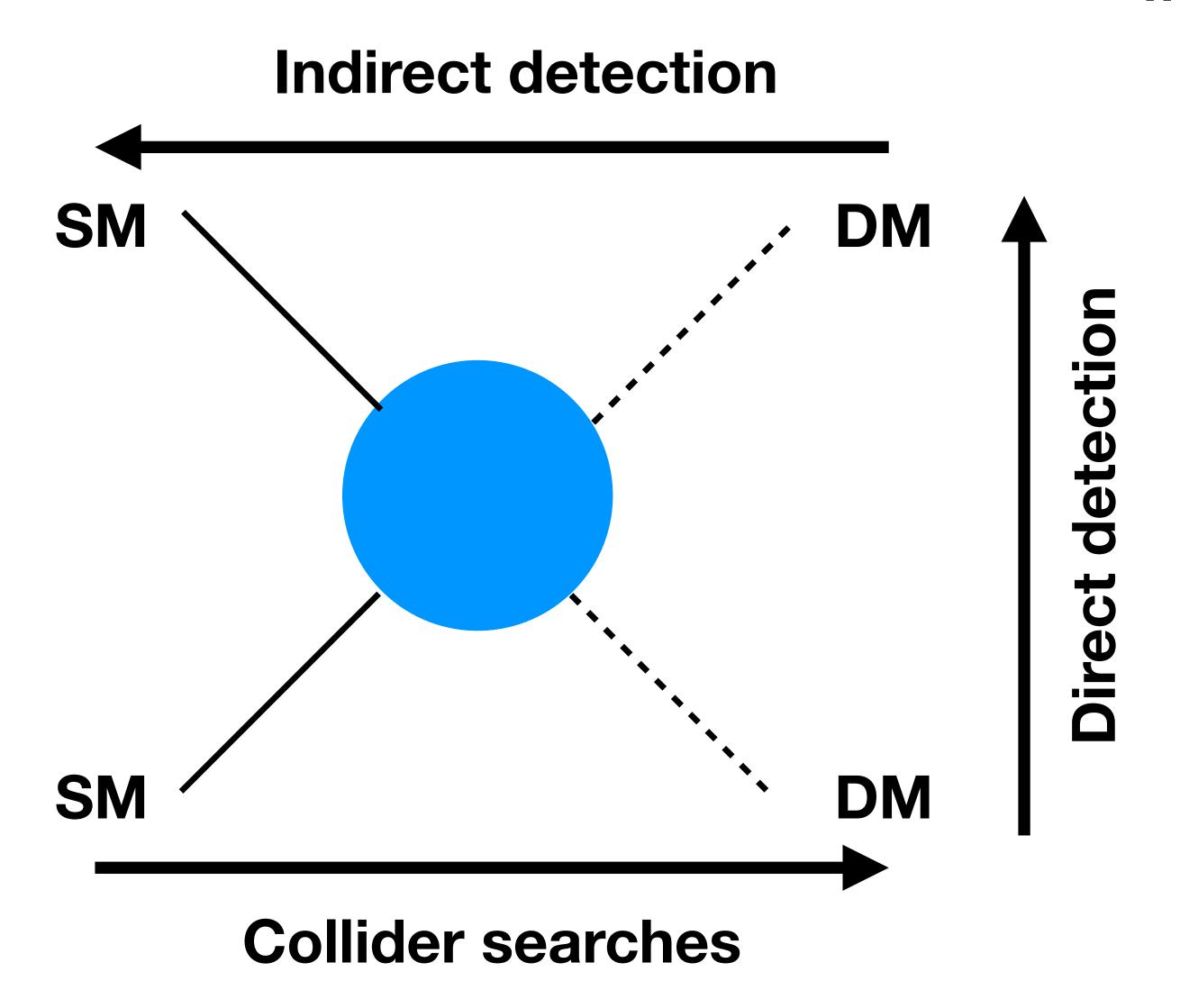
Gravitational lensing



Cosmic microwave background

Searching for Dark Matter

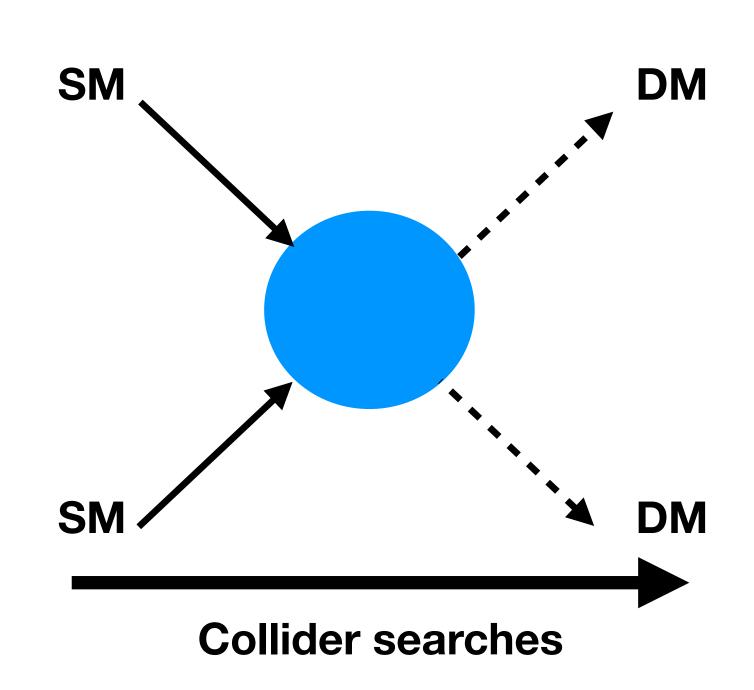
See lectures by Axel Lindner and Aaron Dean Spector for more on Dark Matter

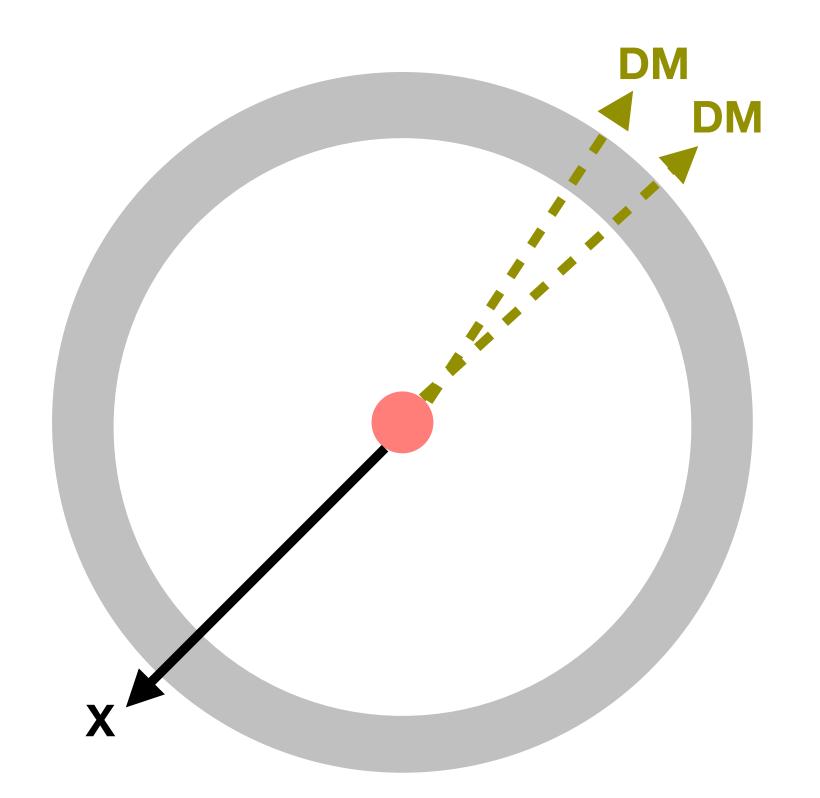


Dark Matter searches at colliders

DM could be produced in proton-proton collisions

- DM does not interact with the detector
- Infer due to momentum imbalance when produced in association with particle "X"
 - → mono-X (can also be several particles)



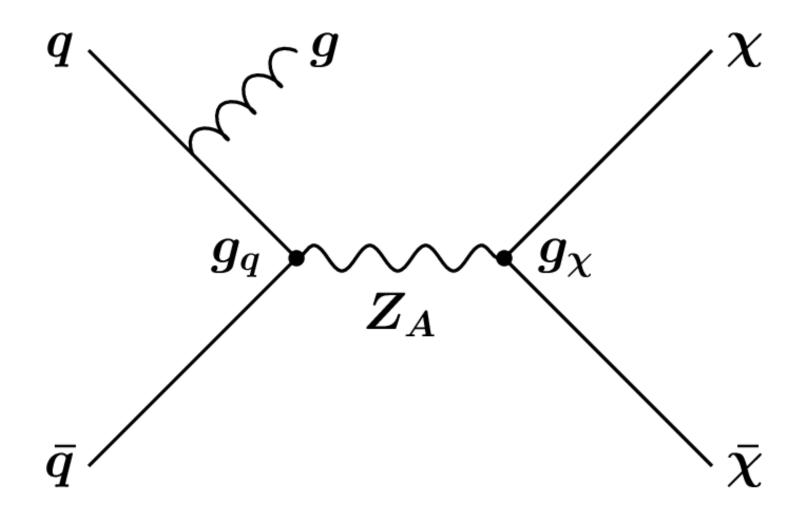


Mono-X signature

ATLAS and CMS cannot detect what they cannot see

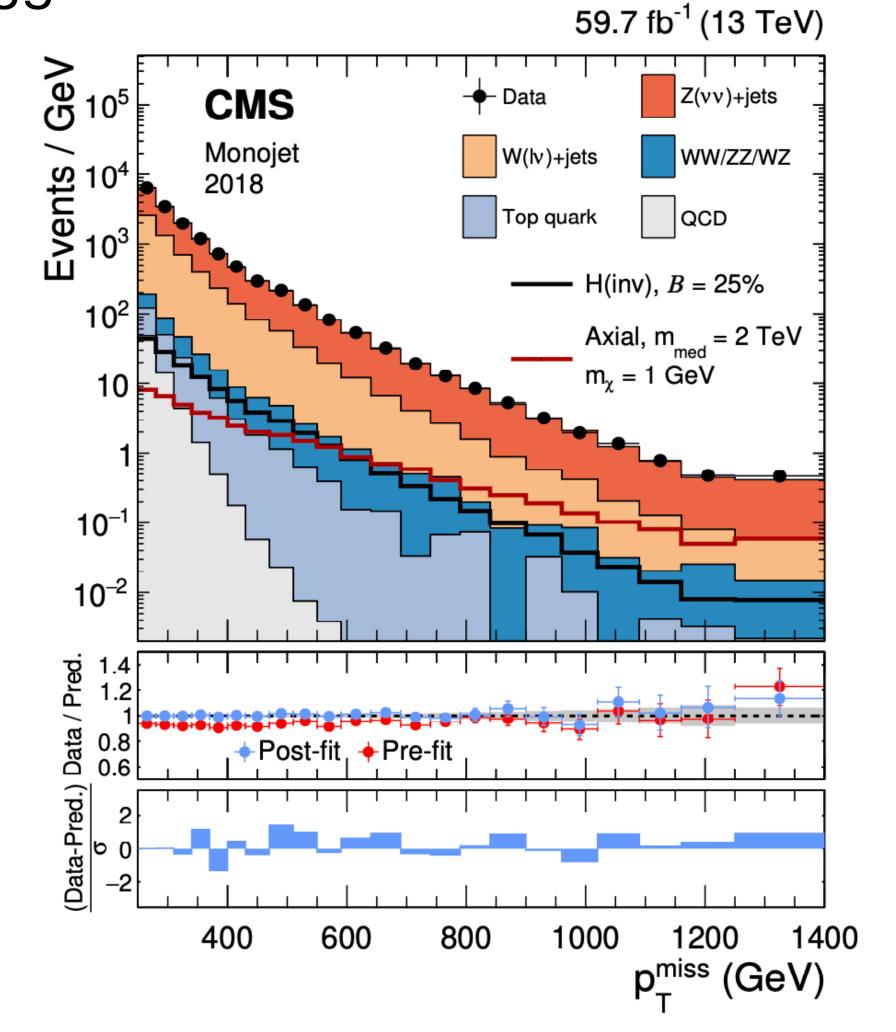
→ If DM is produced at the LHC need some X to trigger on

Simplified model



ISR jet with axial-vector Z (spin 1, parity even)

Could also be: photon, W, Z

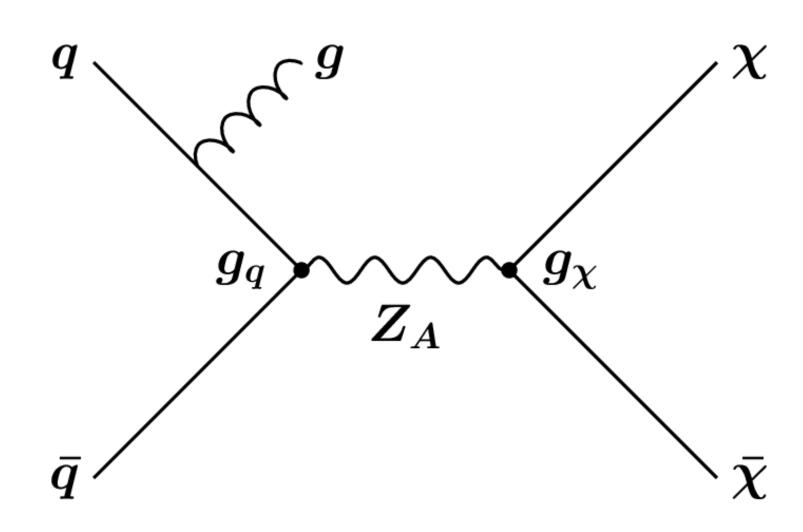


Mono-X signature

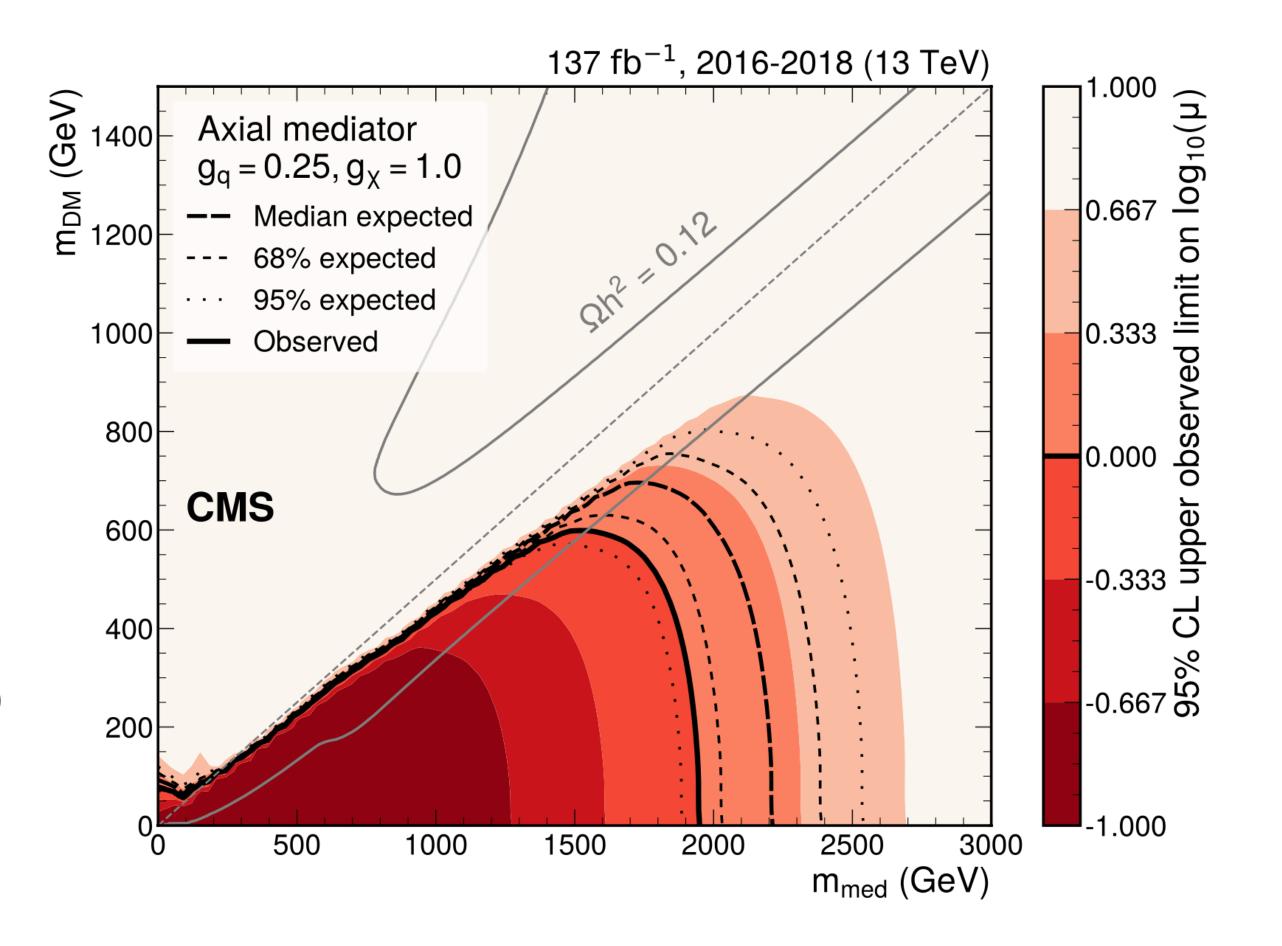
ATLAS and CMS cannot detect what they cannot see

→ If DM is produced at the LHC need some X to trigger on

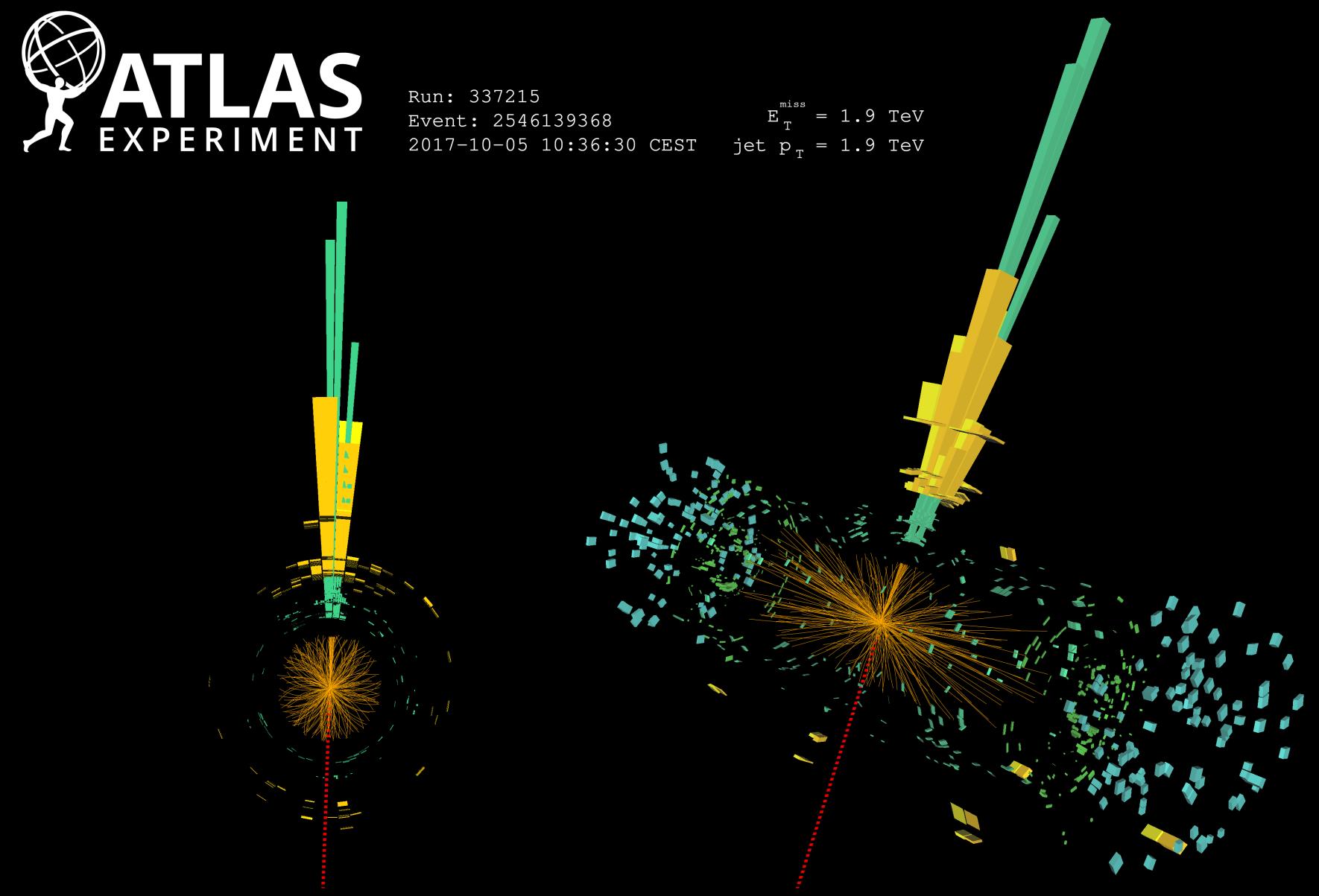
Simplified model



ISR jet with axial-vector Z (spin 1, parity even)
Could also be: photon, W, Z



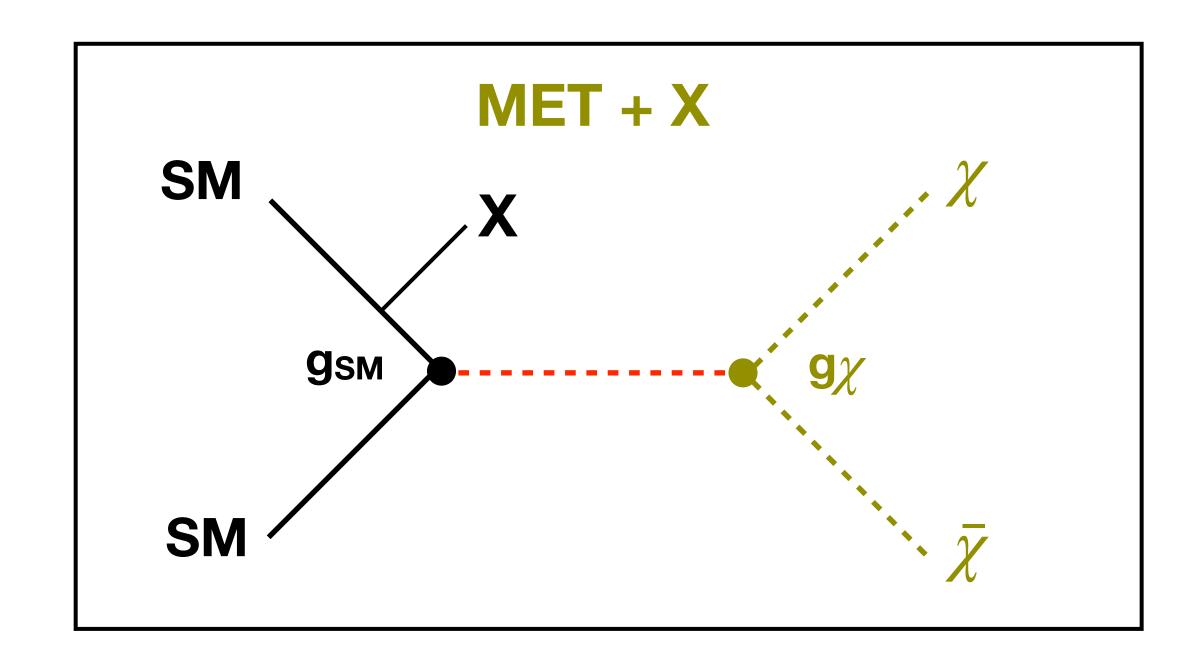
A spectacular mono-jet event

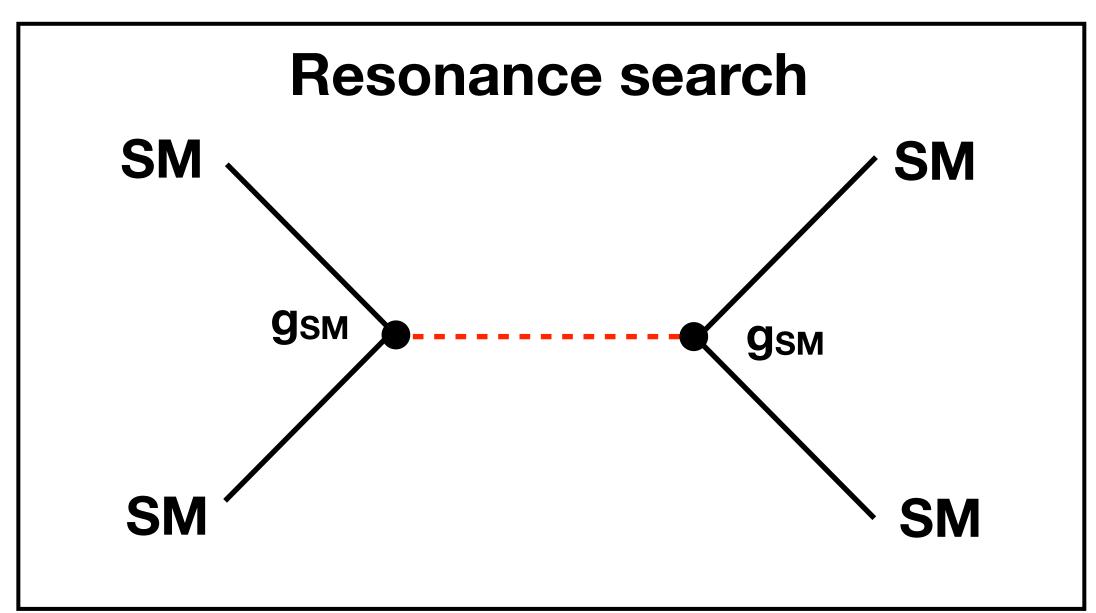


Phys. Rev. D 103 (2021) 112006

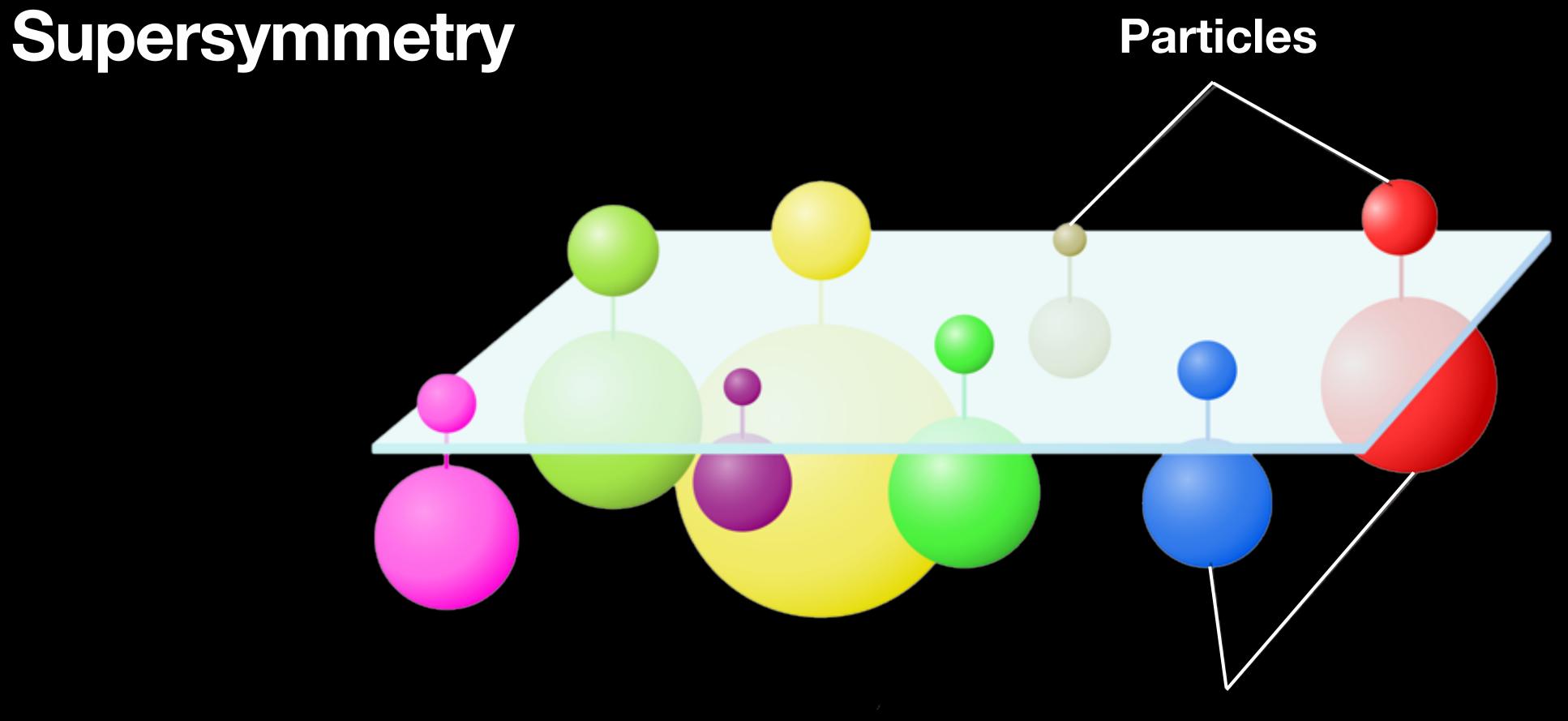
Searching for the mediator

Can also look for dark matter mediators at the LHC!





Need a model to combine the searches which increases sensitivity



Supersymmetric particles

Every SM particle has a SUSY partner (sparticles) with spin difference of 1/2 i.e. Fermions ↔ Bosons Broken symmetry → different masses

Supersymmetry

R-parity: (B=baryon number, L=lepton number, s=spin)

$$R = (-1)^{(2s+3B+L)} = \begin{cases} +1 \text{ for SM particles} \\ -1 \text{ for SUSY particles} \end{cases}$$

R-parity conservation (RPC)

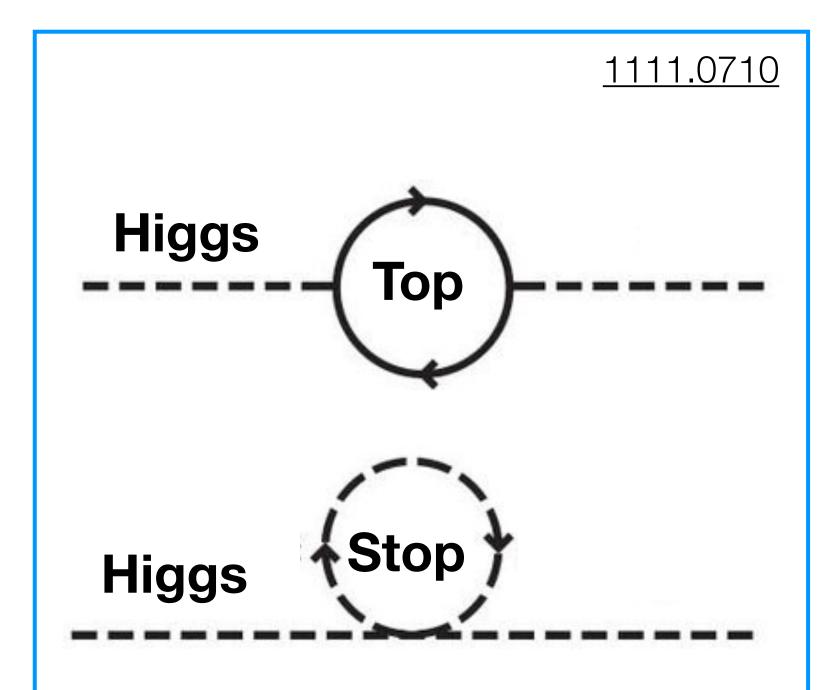
originally introduced for proton stability

- Always pairs of sparticles
- Lightest supersymmetric particle (LSP)
 is stable and escapes detection
- Final state decay has at least one LSP

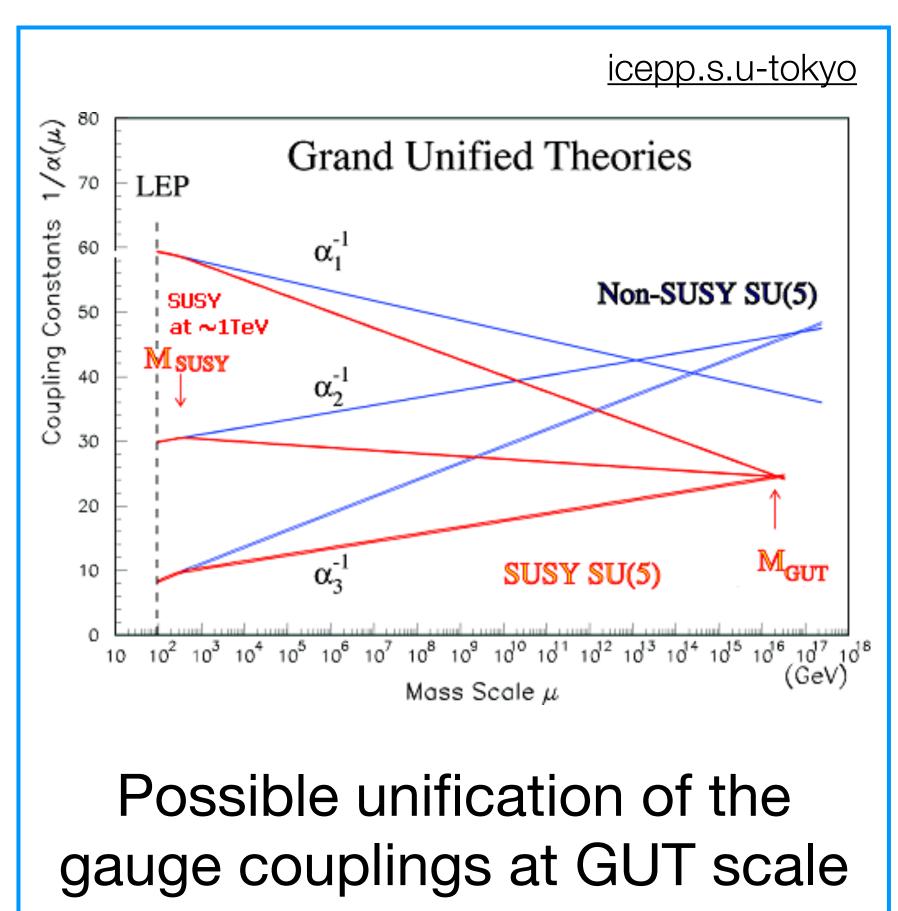
R-parity violation (RPV)

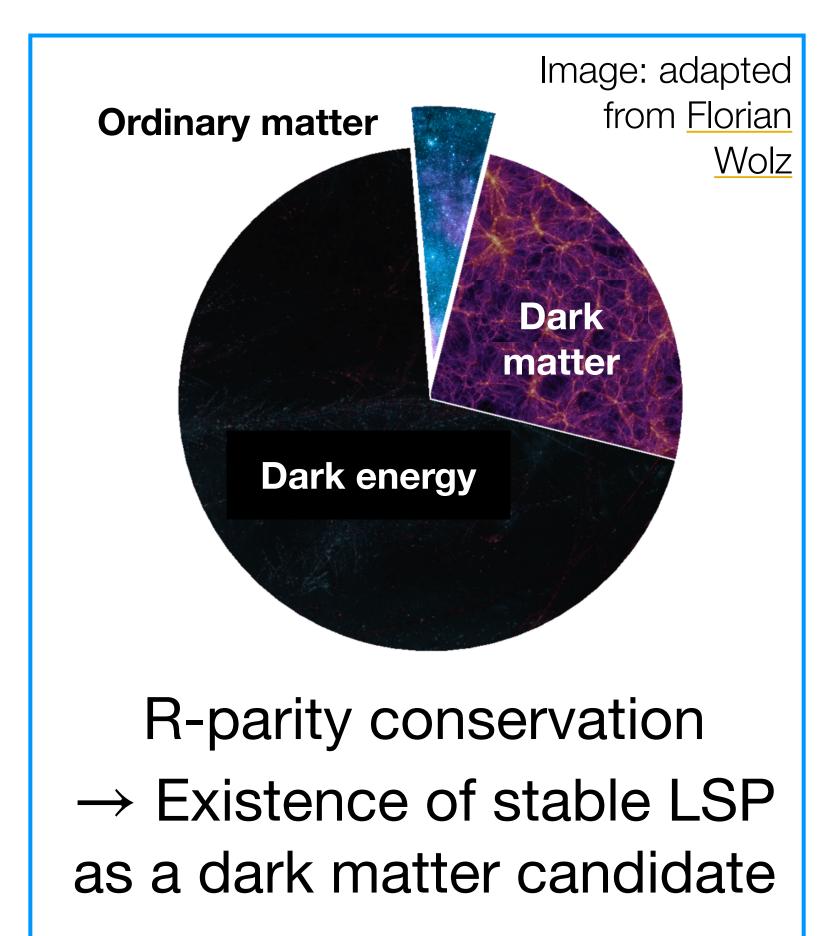
- Either lepton or baryon number violation
- Sparticles can decay exclusively to SM particles
- Low missing energy in the final state

Supersymmetry



Sparticle loops cancel corrections to Higgs mass



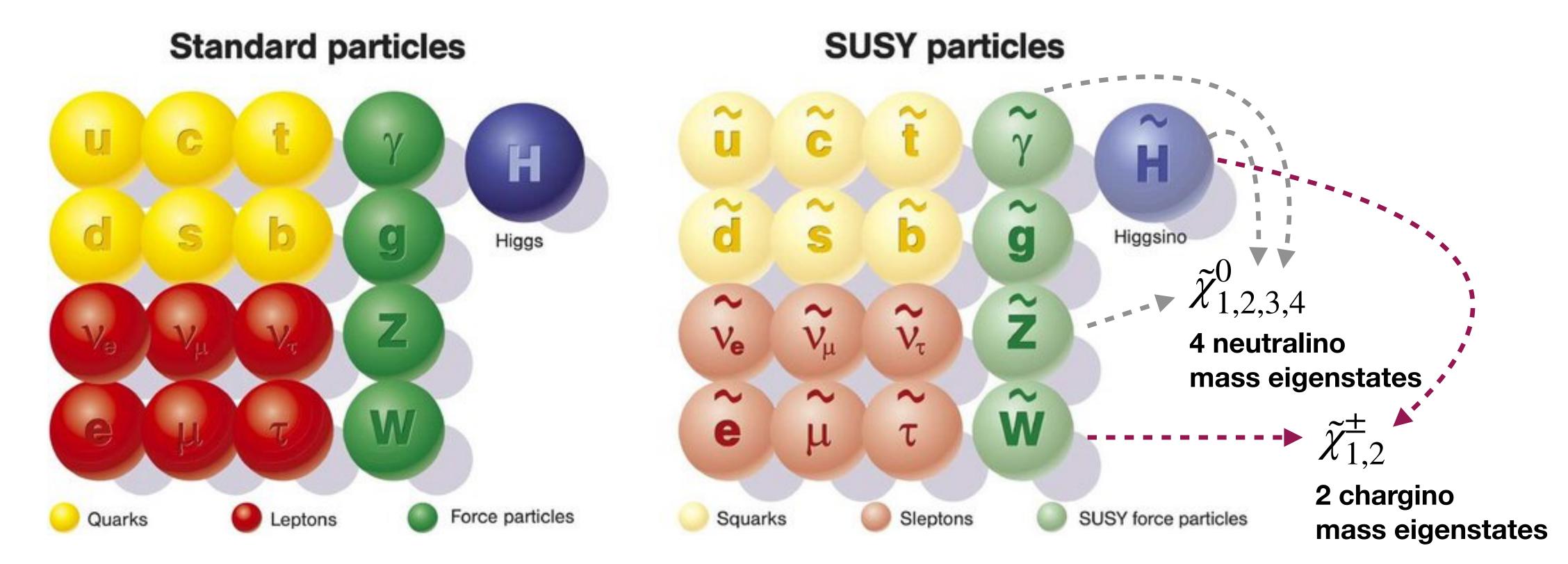


Supersymmetry

SUSY predicts a plethora of new particles

Potential parameter space is huge

e.g. MSSM: ~100, pMSSM: 19



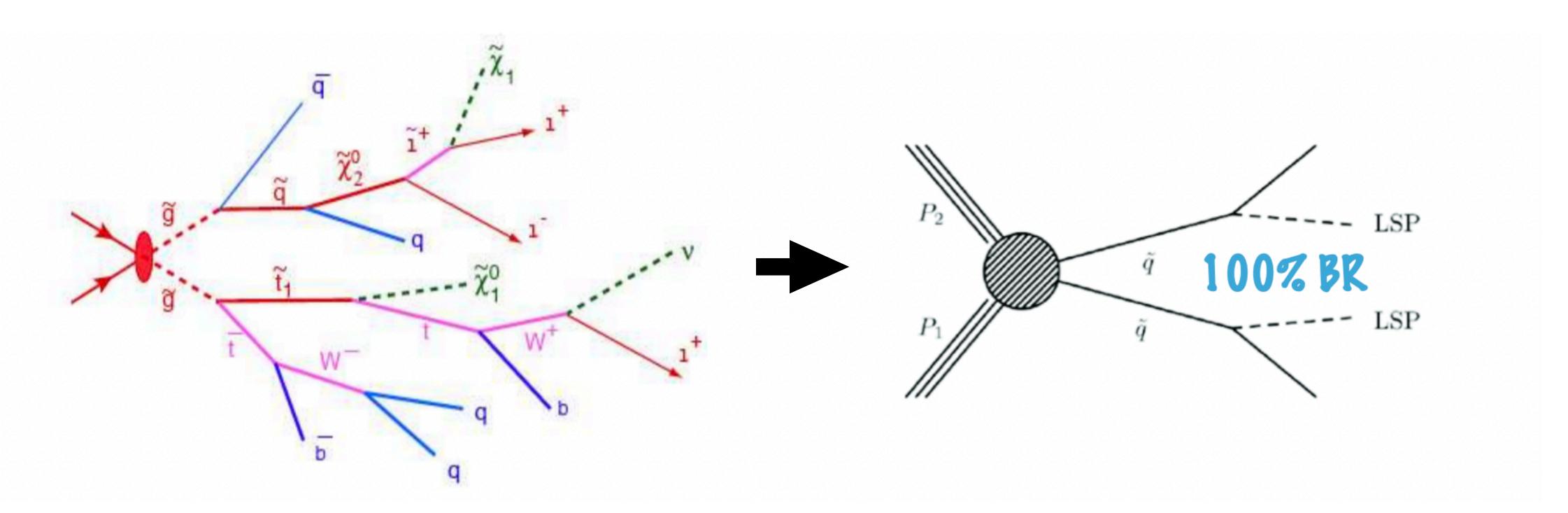
MSSM = Minimal Supersymmetric Standard Model

Simplified models in SUSY searches

Often we simplify further e.g.

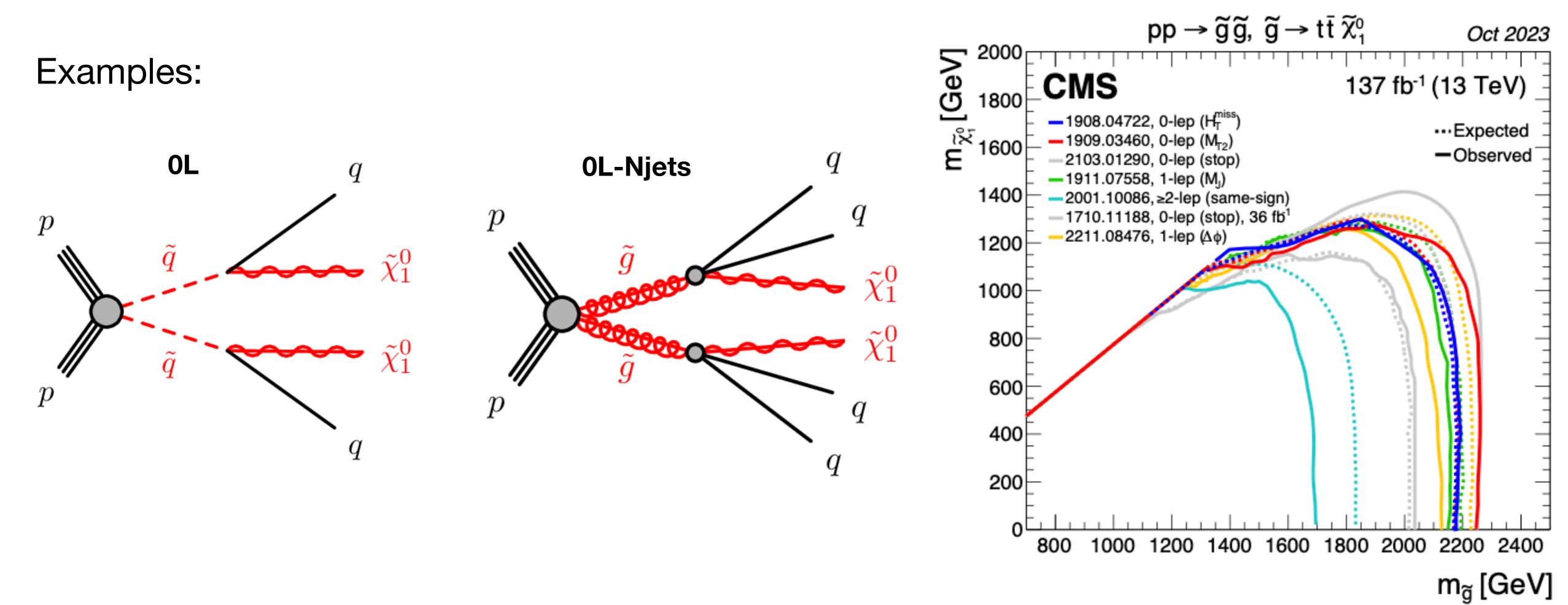
Assume masses of most SUSY particles are so high they do not affect the decay chains Make assumptions about mass degeneracy

Be careful when interpreting limits, always check assumptions



Squark and Gluino searches

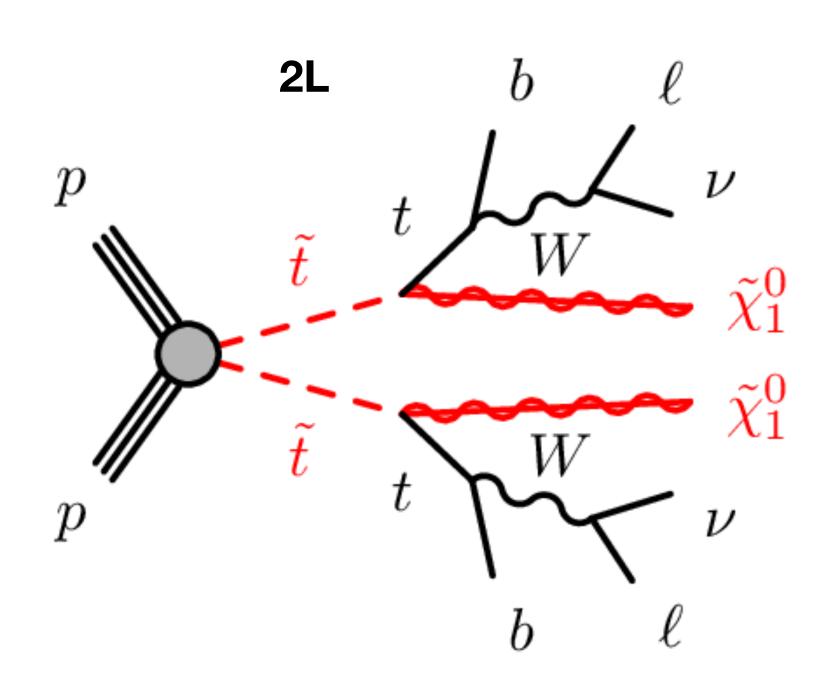
Assuming pair production of squarks and gluinos, decays to LSP + various SM quarks (light and heavy flavour) \rightarrow Highest production cross sections

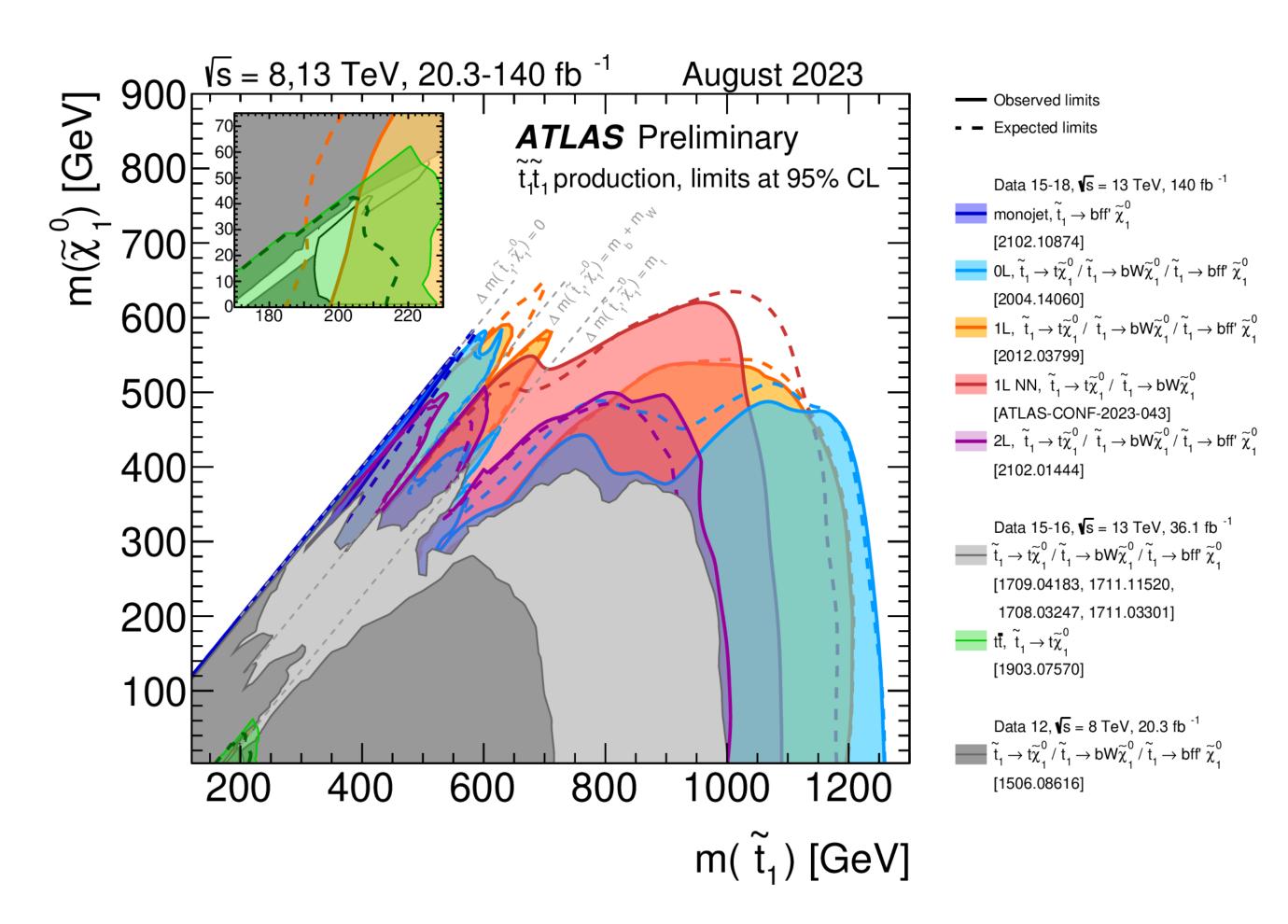


Stop searches

Stops plan important role in stabilising the Higgs mass (cancel top contribution) 0, 1 and 2 lepton channels \rightarrow remember $t\bar{t}$ decay modes + extra MET

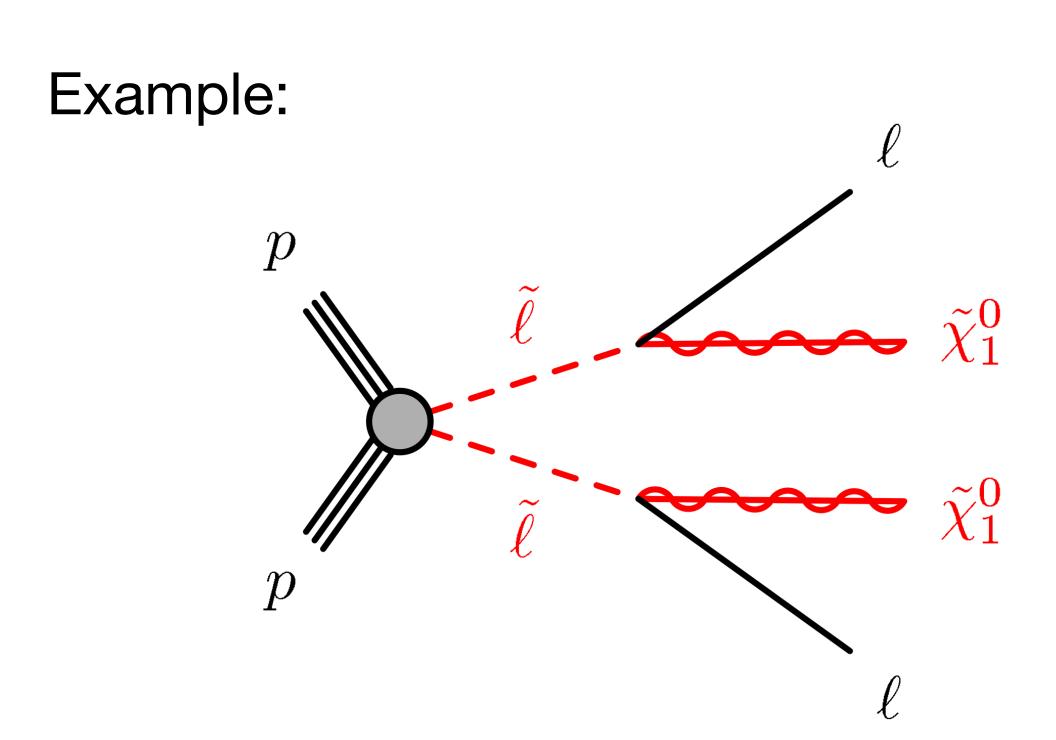
Example:

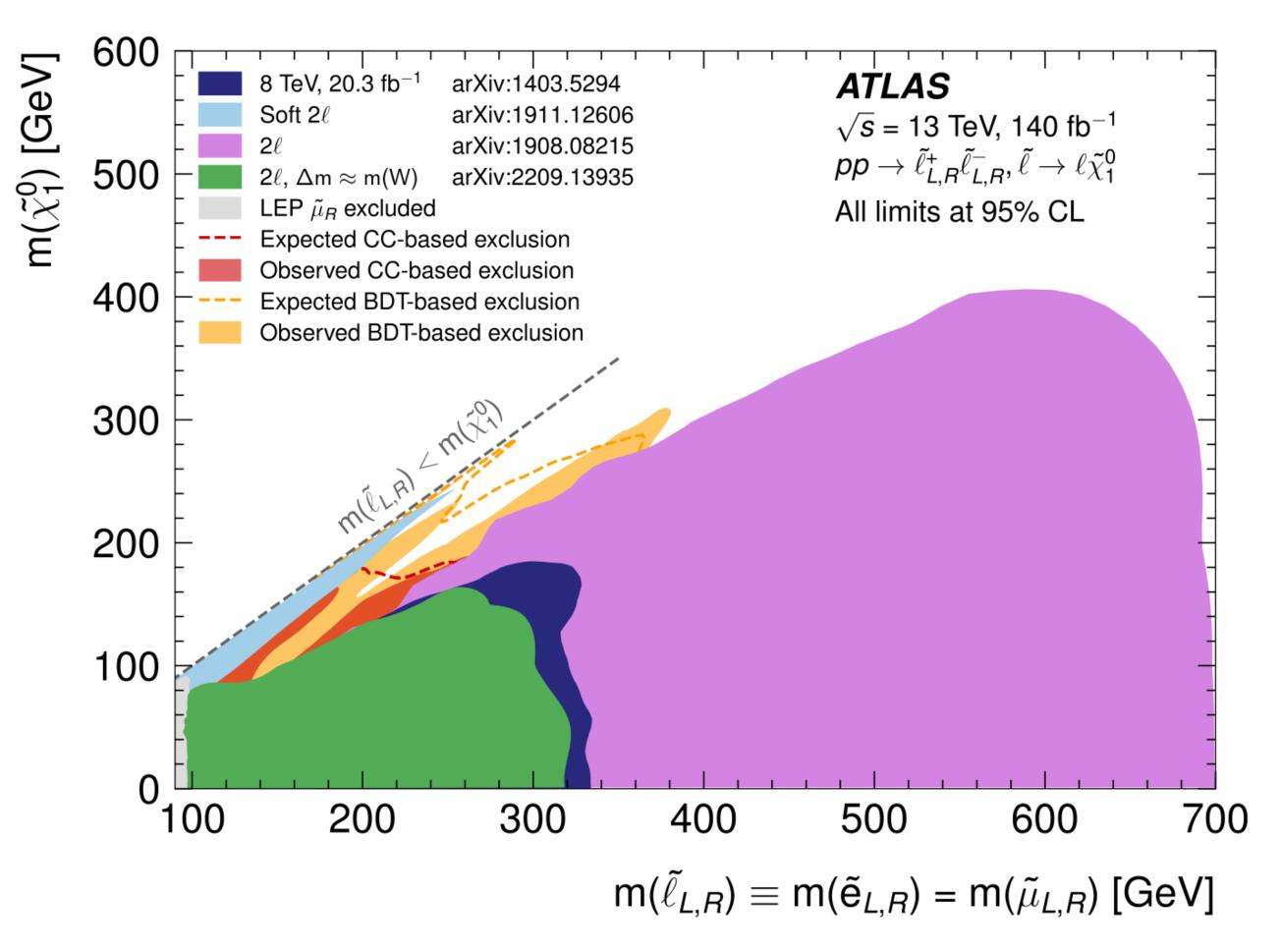




Electroweak searches

Coloured sparticles could be heavier than EW sparticles Production cross section much lower for charginos & sleptons than squarks & gluino

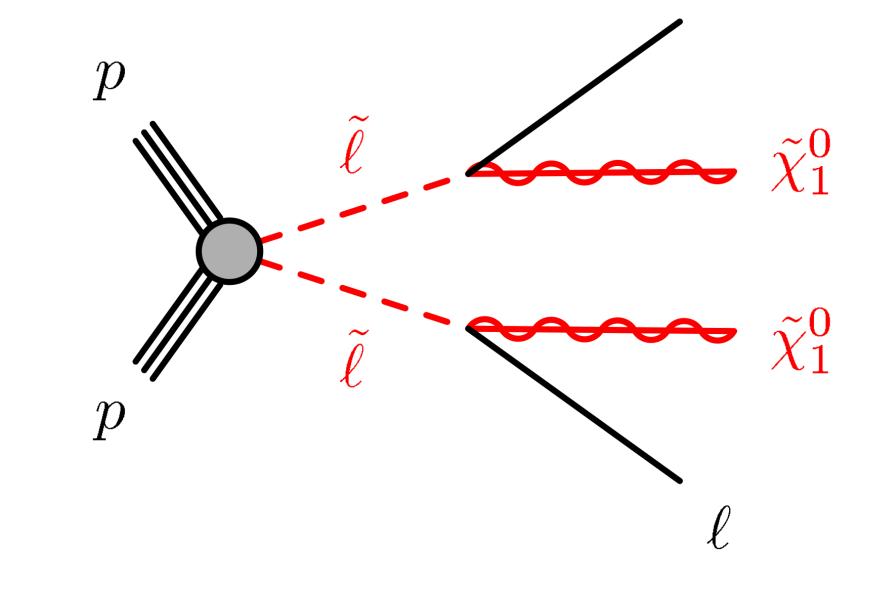




Slepton limits pushing 700 GeV in simplified models but gaps remain

From low to high-hanging fruits

- Strong
- Electroweak
- Compressed → Small mass differences
- · R-parity violating (low MET)
- Unconventional signatures



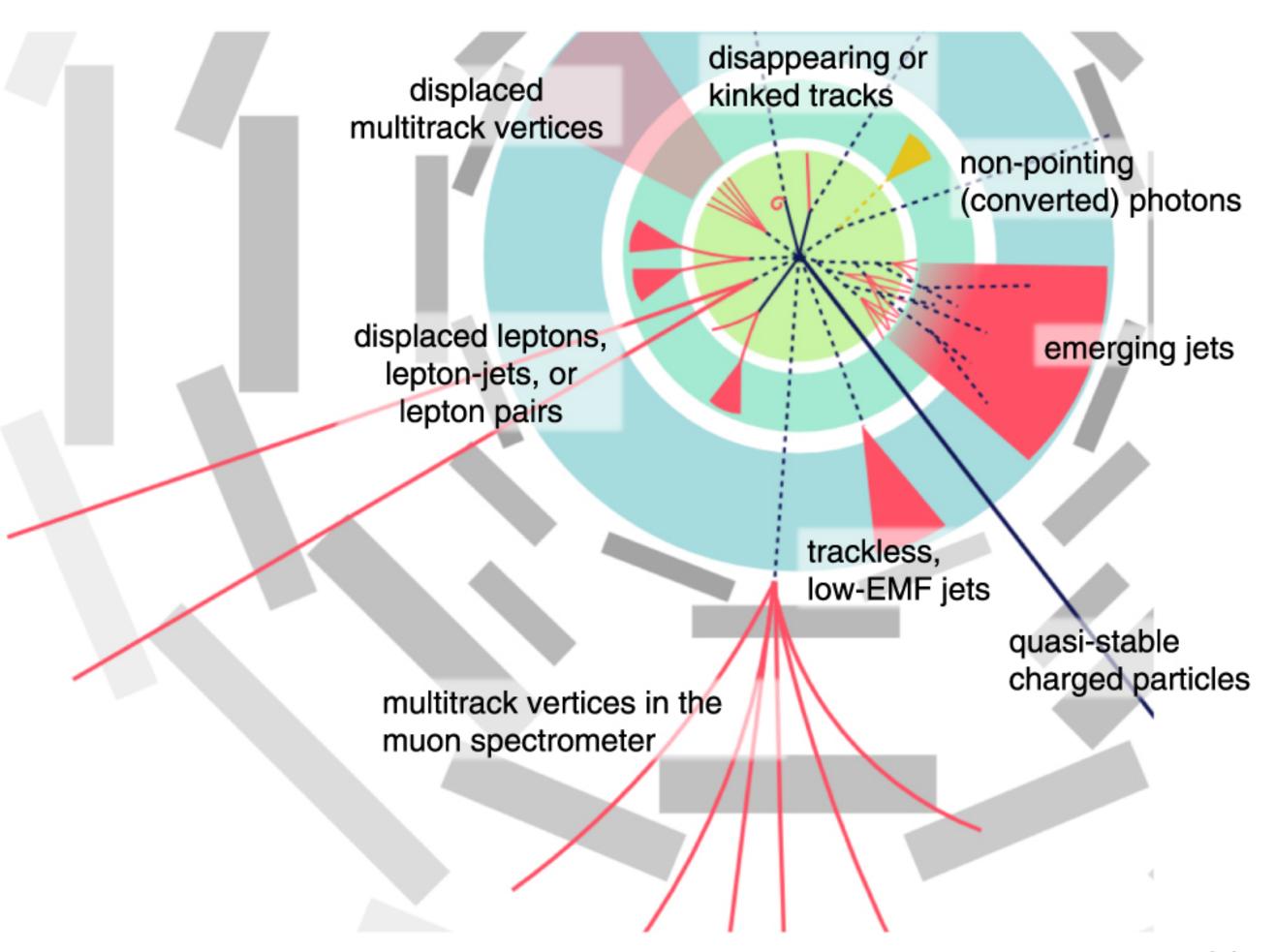


Unconventional signatures

So far the models were BSM but the signatures contained "known" SM particles

→ What if the decay products are "BSM"?

- Often "long lived" new particles
- Bkgs very different from other searches
 - Data driven approaches to estimate them
- Often dedicated data reconstruction algorithms needed



Where do we stand after Run 2?

 10^{-1}

ATLAS Preliminary ATLAS SUSY Searches* - 95% CL Lower Limits July 2024 \sqrt{s} = 13 TeV Signature Model $\int \mathcal{L} dt \, [fb^{-1}]$ Mass limit Reference $E_T^{ m miss}$ $E_T^{ m miss}$ $m(\tilde{\chi}_1^0) < 400 \text{ GeV}$ 2010.14293 $\tilde{q}\tilde{q}, \, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ 1-3 jets 140 \tilde{q} [8× Degen.] 0.9 2102.10874 $m(\tilde{q})-m(\tilde{\chi}_1^0)=5 \text{ GeV}$ $E_T^{ m miss}$ 140 2-6 jets $\tilde{g}\tilde{g}, \, \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$ $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ 2010.14293 1.15-1.95 $m(\tilde{\chi}_1^0)=1000 \text{ GeV}$ 2010.14293 2-6 jets 2.2 $m(\tilde{\chi}_1^0)$ <600 GeV $\tilde{g}\tilde{g}, \, \tilde{g} \rightarrow q\bar{q}W\tilde{\chi}_1^0$ 140 2101.01629 $ee, \mu\mu$ 2 jets 140 2.2 $m(\tilde{\chi}_1^0)$ <700 GeV 2204.13072 $\tilde{g}\tilde{g}, \, \tilde{g} \rightarrow q\bar{q}(\ell\ell)\tilde{\chi}_1^0$ 7-11 jets 140 1.97 $m(\tilde{\chi}_{1}^{0}) < 600 \,\text{GeV}$ 2008.06032 $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$ SS e, μ 140 1.15 2307.01094 $m(\tilde{g})-m(\tilde{\chi}_1^0)=200 \text{ GeV}$ $\tilde{g}\tilde{g}, \; \tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$ 3 *b* $m(\tilde{\chi}_1^0)$ <500 GeV 2211.08028 SS e, μ 140 1.25 6 jets 1909.08457 $m(\tilde{g})-m(\tilde{\chi}_1^0)=300 \text{ GeV}$ $\tilde{b}_1 \tilde{b}_1$ 1.255 $0e, \mu$ 2 *b* 140 $m(\tilde{\chi}_1^0)$ <400 GeV 2101.12527 10 GeV< Δ m($\tilde{b}_1\tilde{\chi}_1^0$)<20 GeV 2101.12527 $\tilde{b}_1 \tilde{b}_1, \, \tilde{b}_1 {\rightarrow} b \tilde{\chi}_2^0 \rightarrow b h \tilde{\chi}_1^0$ 6 *b* 0.23-1.35 $\Delta m(\tilde{\chi}_{2}^{0}, \tilde{\chi}_{1}^{0}) = 130 \text{ GeV}, m(\tilde{\chi}_{1}^{0}) = 100 \text{ GeV}$ Forbidden 1908.03122 140 0.13-0.85 2 b 2τ 2103.08189 $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130 \text{ GeV}, m(\tilde{\chi}_1^0) = 0 \text{ GeV}$ $\tilde{t}_1 \tilde{t}_1, \, \tilde{t}_1 \rightarrow t \tilde{\chi}_1^0$ ≥ 1 jet 140 1.25 2004.14060, 2012.03799 $m(\tilde{\chi}_1^0)=1 \text{ GeV}$ $\tilde{t}_1\tilde{t}_1, \, \tilde{t}_1 {\rightarrow} Wb\tilde{\chi}_1^0$ 140 Forbidden 1.05 $m(\tilde{\chi}_1^0)=500 \text{ GeV}$ 2012.03799, 2401.13430 $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b \nu, \tilde{\tau}_1 \rightarrow \tau \tilde{G}$ 1-2 auForbidden $m(\tilde{\tau}_1)=800 \text{ GeV}$ 2108.07665 $\tilde{t}_1 \tilde{t}_1, \, \tilde{t}_1 \rightarrow c \tilde{\chi}_1^0 / \, \tilde{c} \tilde{c}, \, \tilde{c} \rightarrow c \tilde{\chi}_1^0$ 36.1 $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ 1805.01649 140 0.55 2102.10874 $m(\tilde{t}_1,\tilde{c})-m(\tilde{\chi}_1^0)=5$ GeV $\tilde{t}_1\tilde{t}_1, \, \tilde{t}_1 {\rightarrow} t\tilde{\chi}_2^0, \, \tilde{\chi}_2^0 {\rightarrow} Z/h\tilde{\chi}_1^0$ 1-2 e, μ 1-4 *b* 140 0.067-1.18 $m(\tilde{\chi}_2^0)=500 \text{ GeV}$ 2006.05880 $\tilde{t}_2\tilde{t}_2, \, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ $3e,\mu$ Forbidden 0.86 $m(\tilde{\chi}_{1}^{0})=360 \text{ GeV}, m(\tilde{t}_{1})-m(\tilde{\chi}_{1}^{0})=40 \text{ GeV}$ 2006.05880 $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via WZMultiple ℓ/jets $m(\tilde{\chi}_1^0)=0$, wino-bino 2106.01676, 2108.07586 $ee, \mu\mu$ ≥ 1 jet 140 0.205 $m(\tilde{\chi}_1^{\pm})-m(\tilde{\chi}_1^{0})=5$ GeV, wino-bino 1911.12606 $E_T^{ m miss}$ $\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp}$ via WW140 0.42 $2e,\mu$ $m(\tilde{\chi}_1^0)=0$, wino-bino 1908.08215 $\tilde{\chi}_1^{\pm}/\tilde{\chi}_2^0$ Multiple ℓ/jets $\tilde{\chi}_1^{\pm} \tilde{\chi}_2^0$ via Wh140 Forbidden $m(\tilde{\chi}_1^0)=70$ GeV. wino-bino 2004.10894, 2108.07586 $ilde{\chi}_1^{\pm} ilde{\chi}_1^{\mp}$ via $ilde{\ell}_L/ ilde{\nu}$ 140 $2e,\mu$ 1.0 $m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^{\pm}) + m(\tilde{\chi}_1^{0}))$ 1908.08215 $\tilde{\tau}\tilde{\tau},\,\tilde{\tau}{ ightarrow} au ilde{\chi}_1^0$ 2τ 140 0.5 $m(\tilde{\chi}_1^0)=0$ 2402.00603 0 jets ≥ 1 jet $\tilde{\ell}_{L,R}\tilde{\ell}_{L,R},\,\tilde{\ell}{ ightarrow}\ell\tilde{\chi}_1^0$ $2e,\mu$ 140 $m(\tilde{\chi}_1^0)=0$ 1908.08215 140 0.26 $m(\tilde{\ell})-m(\tilde{\chi}_1^0)=10 \text{ GeV}$ 1911.12606 $\begin{array}{ccc} \geq 3 \ b & E_{T}^{\text{miss}} \\ \text{0 jets} & E_{T}^{\text{miss}} \\ \geq 2 \ \text{large jets} & E_{T}^{\text{miss}} \end{array}$ $\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$ 2401.14922 0.55 140 $BR(\tilde{\chi}_1^0 \to Z\tilde{G})=1$ 2103.11684 140 0.45-0.93 $BR(\tilde{\chi}_1^0 \to Z\tilde{G}) = \tilde{\chi}_1^0$ 2108.07586 $\mathsf{BR}(\tilde{\chi}_1^0 \to Z\tilde{G}) = \mathsf{BR}(\tilde{\chi}_1^0 \to h\tilde{G}) = 0.5$ 140 0.77 2204.13072 0.66 Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^{\pm}$ Disapp. trk 2201.02472 0.21 2201.02472 Pure higgsino Stable \tilde{g} R-hadron pixel dE/dx 140 2.05 2205.06013 pixel dE/dx 140 \tilde{g} [$\tau(\tilde{g})$ =10 ns] Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq \tilde{\chi}_1^0$ $m(\tilde{\chi}_1^0)=100 \text{ GeV}$ 2205.06013 $\tilde{\ell}\tilde{\ell},\,\tilde{\ell}{ ightarrow}\ell\tilde{G}$ 140 Displ. lep 0.74 $\tau(\tilde{\ell}) = 0.1 \text{ ns}$ ATLAS-CONF-2024-011 0.36 0.36 $\tau(\tilde{\ell}) = 0.1 \text{ ns}$ ATLAS-CONF-2024-011 pixel dE/dx $\tau(\tilde{\ell}) = 10 \text{ ns}$ 2205.06013 $\tilde{\chi}_{1}^{\pm}\tilde{\chi}_{1}^{\mp}/\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{\pm} \rightarrow Z\ell \rightarrow \ell\ell\ell$ $3e,\mu$ 1.05 Pure Wino 2011.10543 0 jets 0.95 1.55 $\tilde{\chi}_1^{\pm} \tilde{\chi}_1^{\mp} / \tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\ell\ell\nu\nu$ $m(\tilde{\chi}_1^0)=200 \text{ GeV}$ 2103.11684 $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_{1}^{0}, \tilde{\chi}_{1}^{0} \rightarrow qqq$ 140 2.34 Large $\lambda_{112}^{\prime\prime}$ 2401.16333 $\tilde{t}\tilde{t}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$ 36.1 0.55 1.05 ATLAS-CONF-2018-003 $m(\tilde{\chi}_1^0)$ =200 GeV, bino-like $\tilde{t}\tilde{t}, \tilde{t} \rightarrow b\tilde{\chi}_1^{\pm}, \tilde{\chi}_1^{\pm} \rightarrow bbs$ 0.95 $\geq 4b$ 140 Forbidden 2010.01015 $m(\tilde{\chi}_1^{\pm})=500 \text{ GeV}$ $\tilde{t}_1\tilde{t}_1, \, \tilde{t}_1 {\rightarrow} bs$ 2 jets + 2 b0.61 1710.07171 $\tilde{t}_1\tilde{t}_1, \, \tilde{t}_1 \rightarrow q \tilde{t}_1$ 140 136 0.4-1.85 BR($\tilde{t}_1 \rightarrow be/b\mu$)>20% 2406.18367 f_1 [1e-10< λ'_{22k} <1e-8, 3e-10< λ'_{22k} <3e-9] $BR(\tilde{t}_1 \rightarrow q\mu) = 100\%$, $\cos\theta_t = 1$ 2003.11956 $\tilde{\chi}_{1}^{\pm}/\tilde{\chi}_{2}^{0}/\tilde{\chi}_{1}^{0}, \, \tilde{\chi}_{1,2}^{0} \rightarrow tbs, \, \tilde{\chi}_{1}^{+} \rightarrow bbs$ 1-2 e, μ ≥6 jets 140 0.2-0.32 Pure higgsino 2106.09609

Mass scale [TeV]

Remember simplified models are only part of the story

Always check assumptions in the model and how it fits with the bigger picture

^{*}Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

Searches for new physics at the LHC

So far no smoking gun for signs of new physics

Here and there local 2-3 sigma excesses → Important to follow up

However, luminosity will now double only every couple of years

New Ideas?

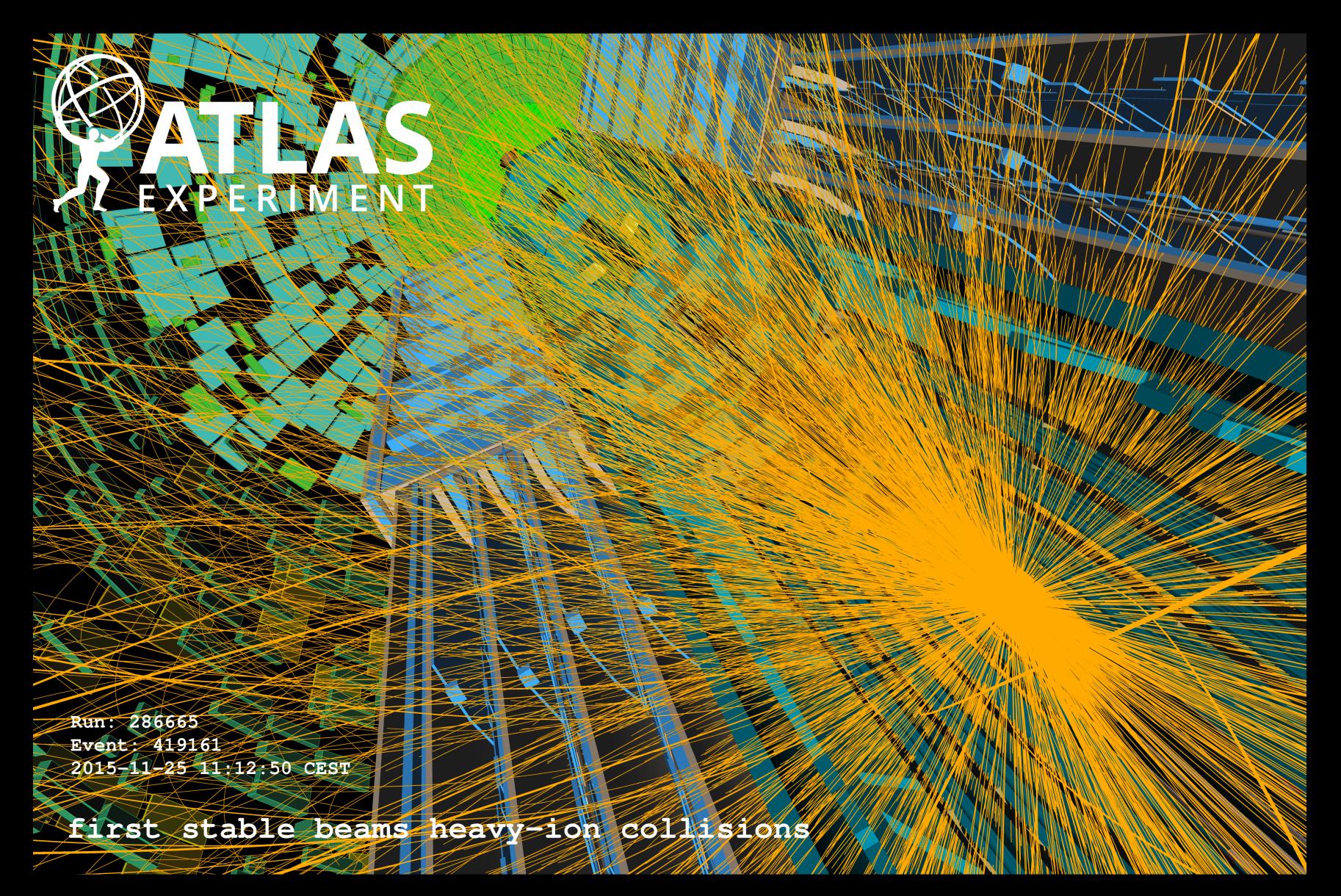
- Make sure we didn't miss something \rightarrow what we didn't trigger on is lost forever
- Maybe something super long lived is produced in ATLAS/CMS but can only be detected with other experiments \rightarrow e.g. FASER?

Physics Goals of the LHC



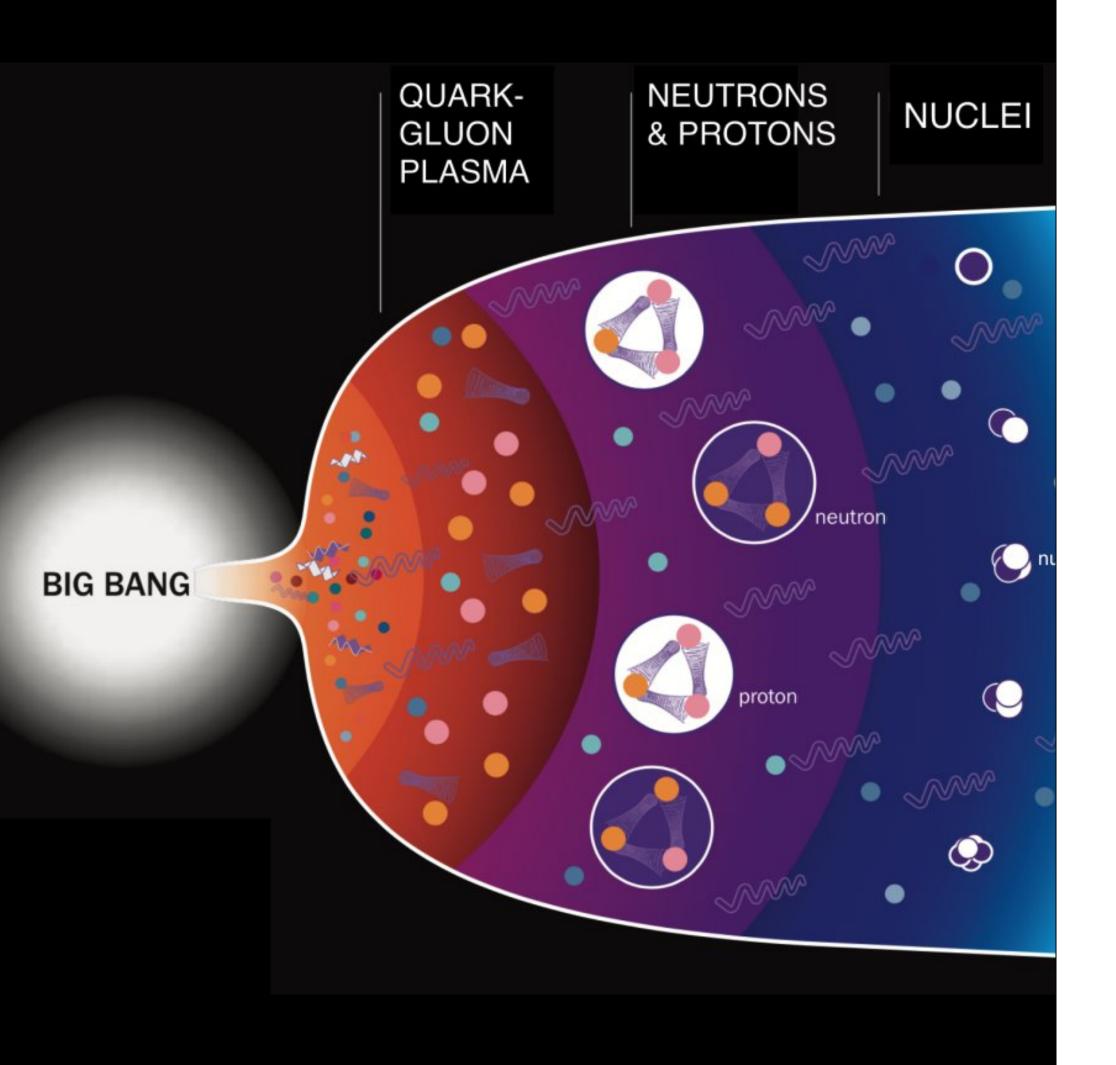
Study the Quark-Gluon Plasma

LHC Pb+Pb collision



ATLAS heavy ion event display

Quark Gluon Plasma (QGP)



Our early universe was too hot & dense for the strong force to bind quarks and gluons together

→ Quark Gluon Plasma (QGP)
Hot dense medium with free colour charges

We can create it in the lab in LHC heavy ion collisions

We study how high momentum partons strongly interact in QGP & lose energy before hadronising

→ Jet quenching

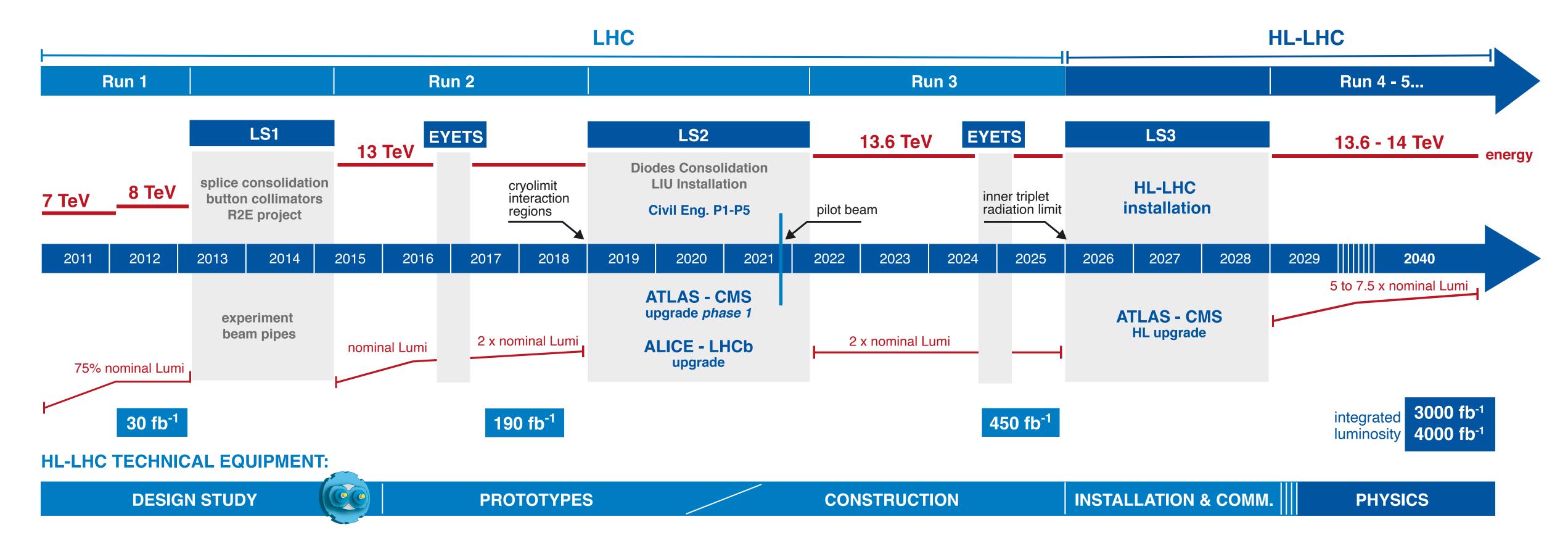
Binding potential of quarkonia e.g. $\Upsilon(b\bar{b})$ can be modified in QGP (colour screening by QGP constituents)

→ Suppression of Quarkonia

Where do we go next?







38

HL-LHC

Major detector upgrades underway

Expected pile-up of ~200 poses immense challenge

Precision measurements of the Higgs boson

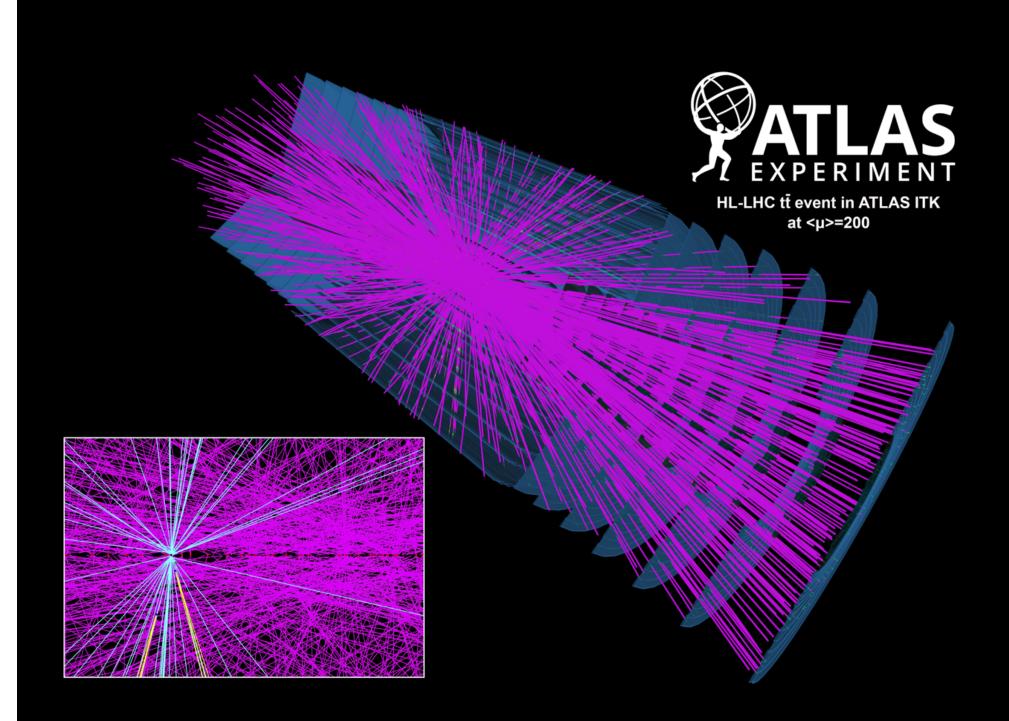
• Establish Higgs self-coupling at the 5σ level

Push the boundaries of SM precision measurements

Could find deviations that could hint at new physics

Pursue dedicated searches for new physics

Your new idea!



Physics Goals of the LHC



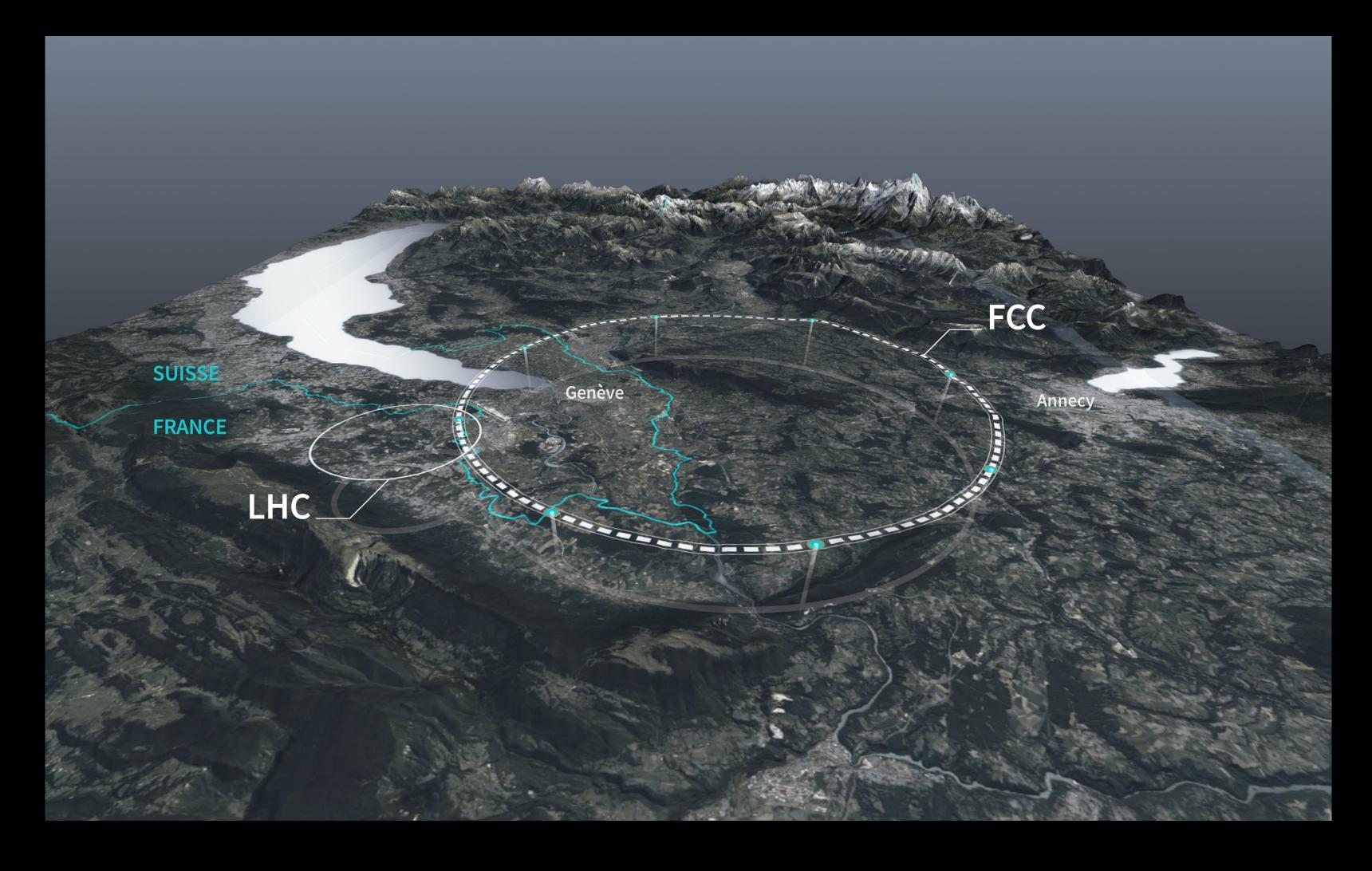
Measure the Standard Model

Search for the Higgs Boson

Search for New Physics

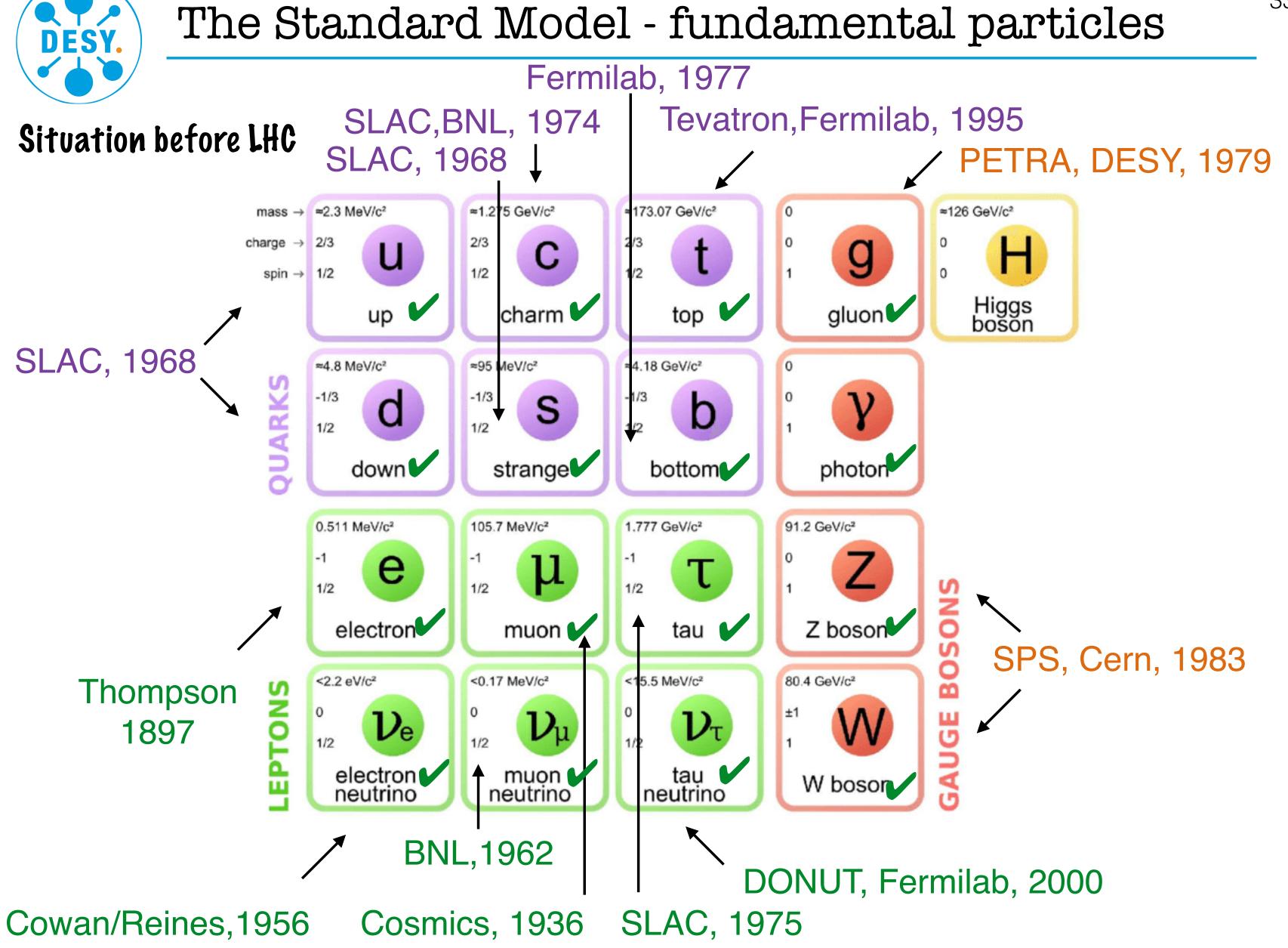
Study the Quark-Gluon Plasma

Outlook

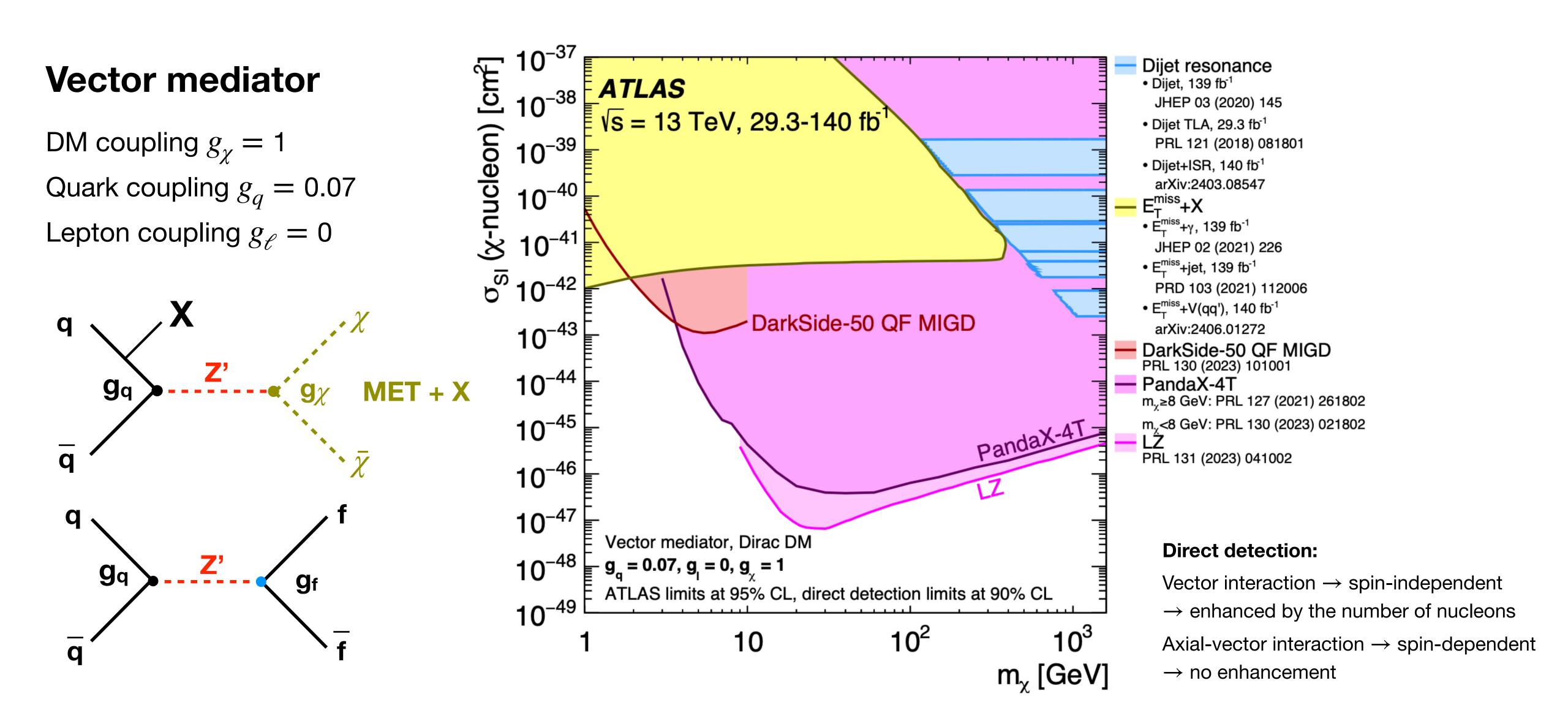


See lectures by Karsten Buesser for more on Future Colliders

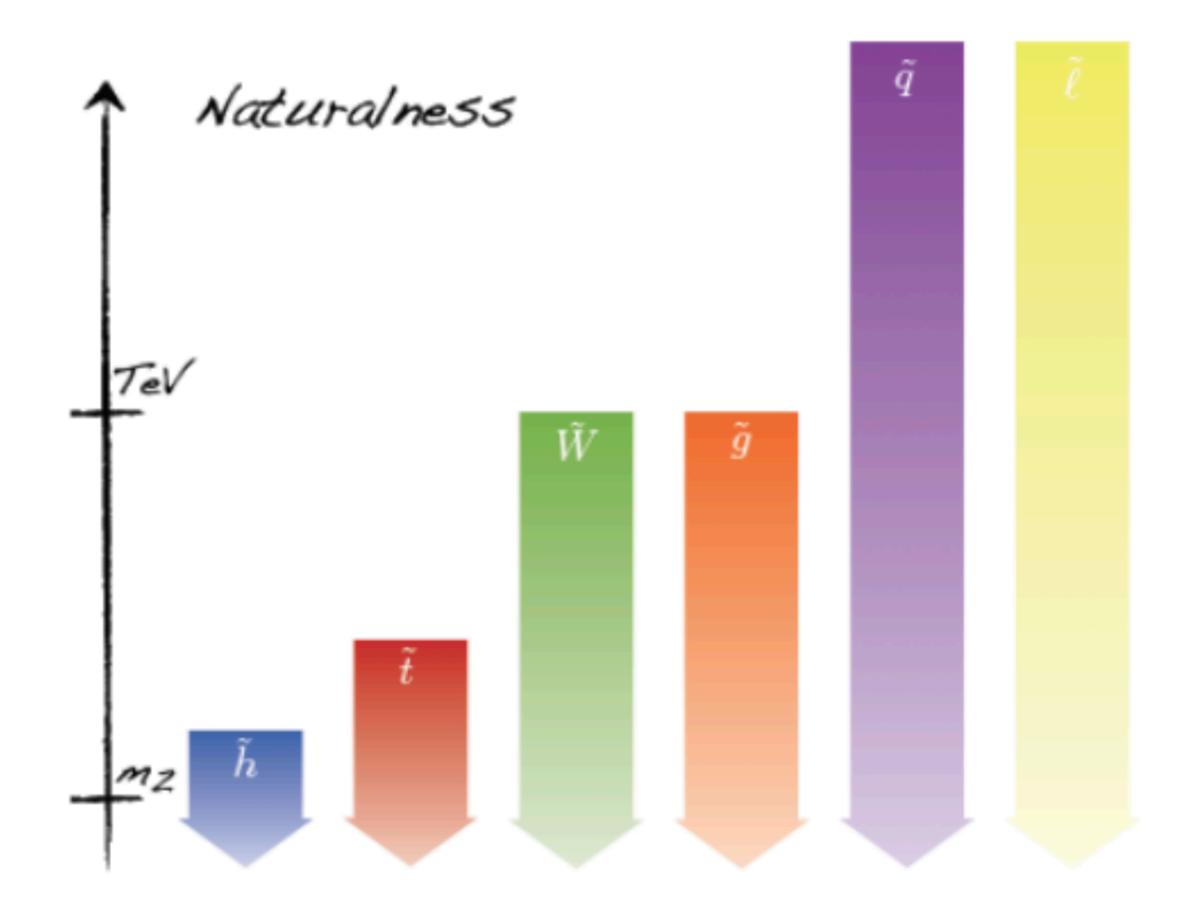
https://fcc.web.cern.ch



Comparing LHC results with direct detection



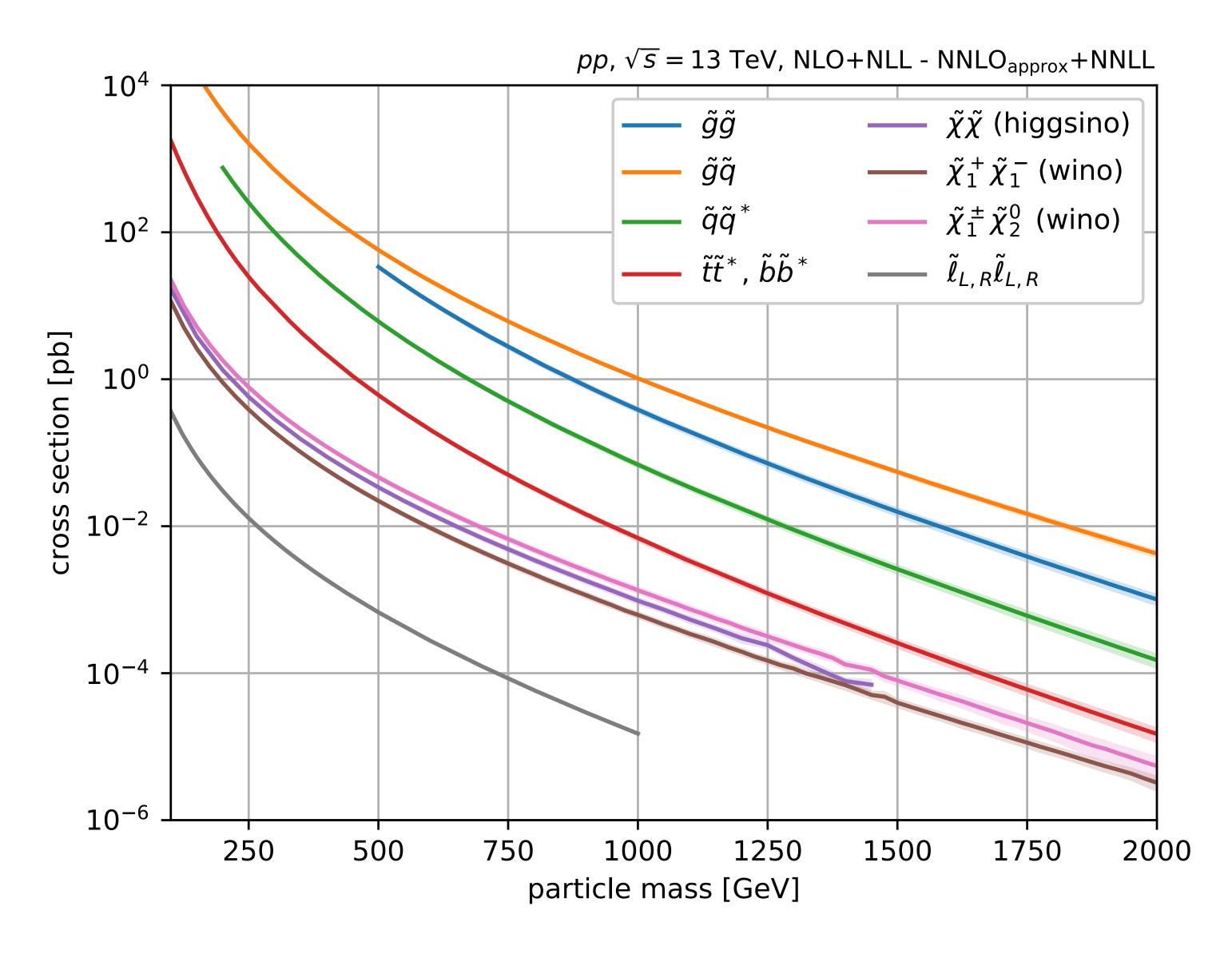
Can only compare results for a given model and set of parameters



SUSY mass scales motivated by electroweak naturalness

N. Craigh, arXiv:1309.0528

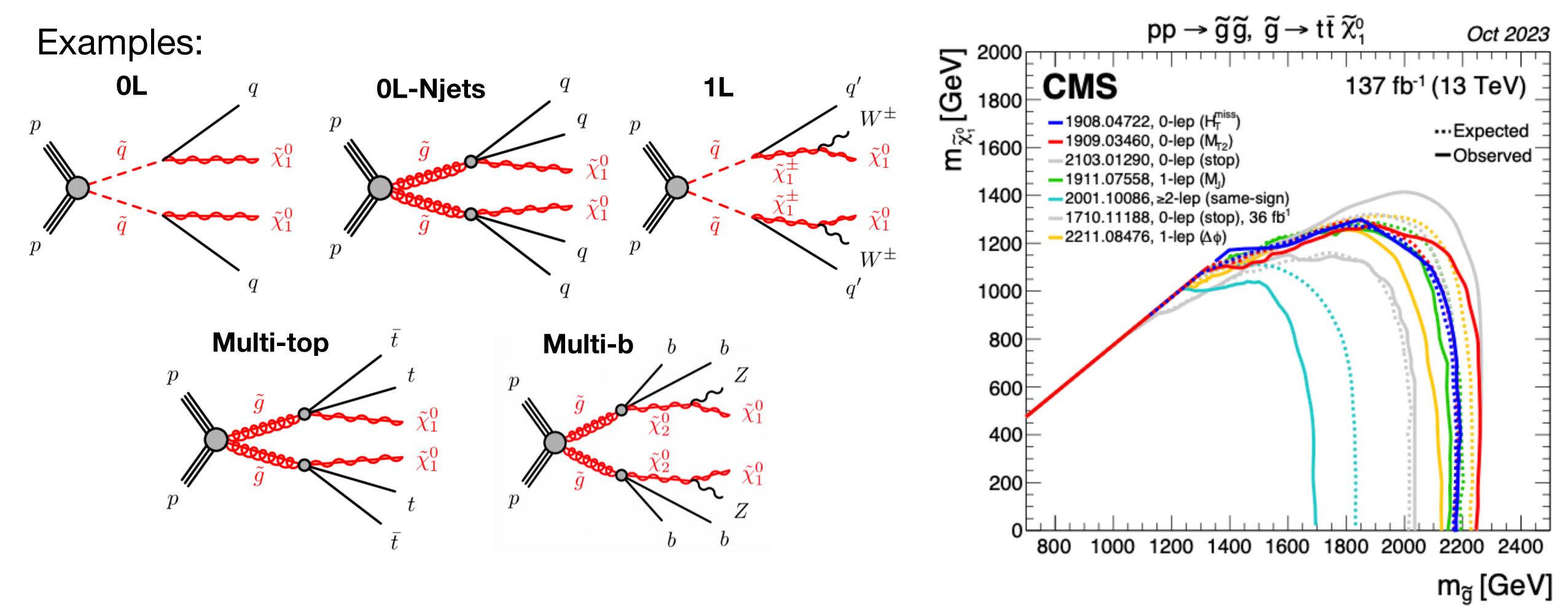
SUSY cross sections



45

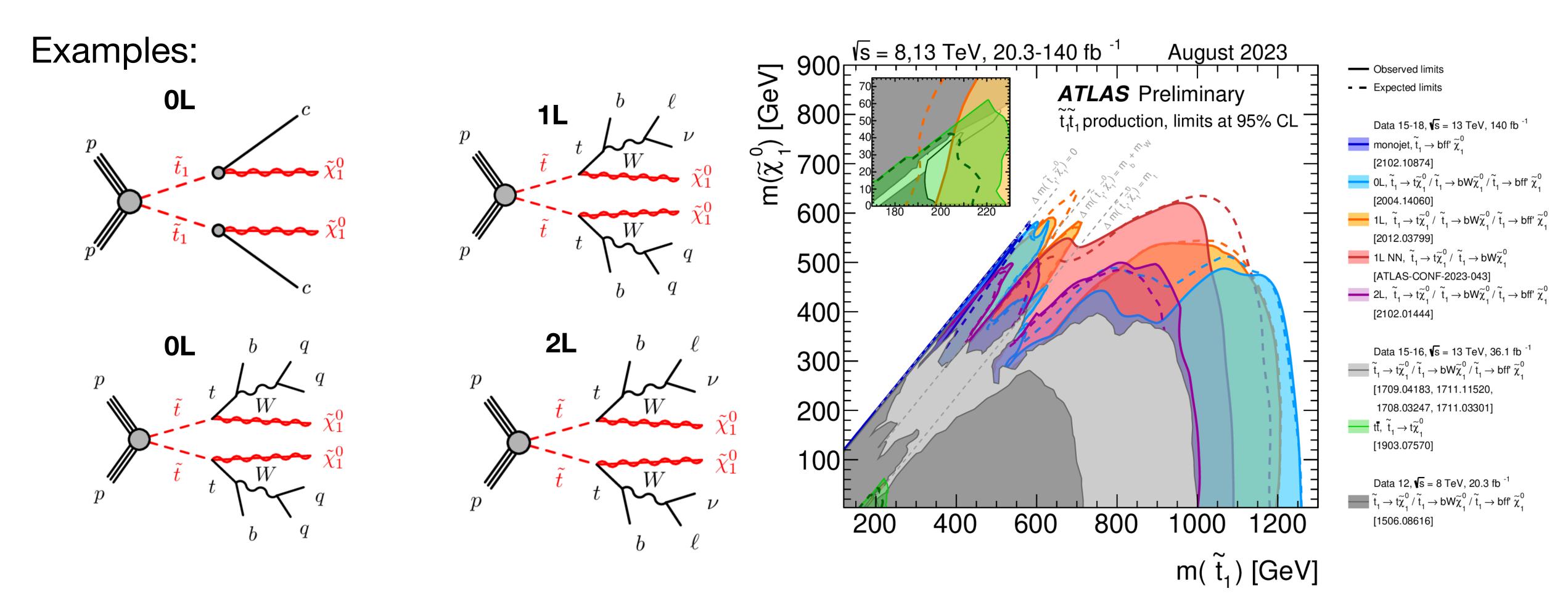
Squark and Gluino searches

Assuming pair production of squarks and gluinos, decays to LSP + various SM quarks (light and heavy flavour) \rightarrow Highest production cross sections

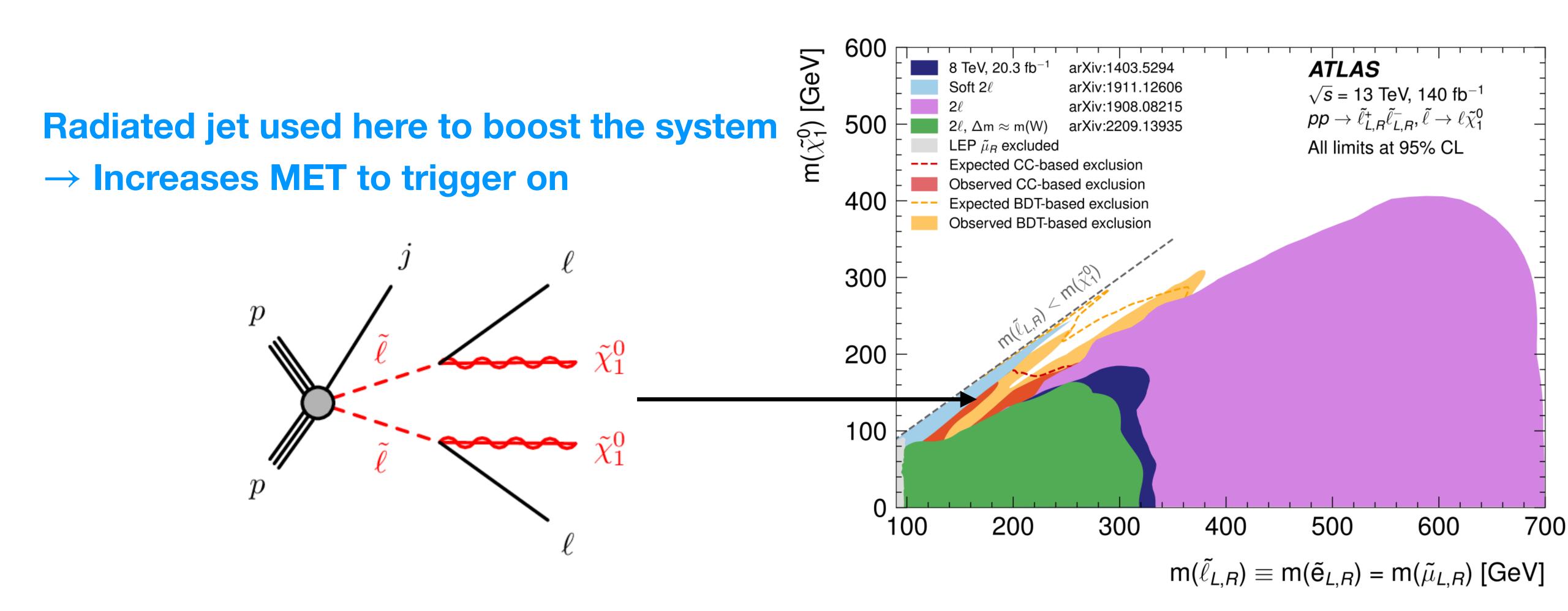


Stop searches

Stops plan important role in stabilising the Higgs mass (cancel top contribution) 0, 1 and 2 lepton channels \rightarrow remember $t\bar{t}$ decay modes + extra MET



Electroweak searches

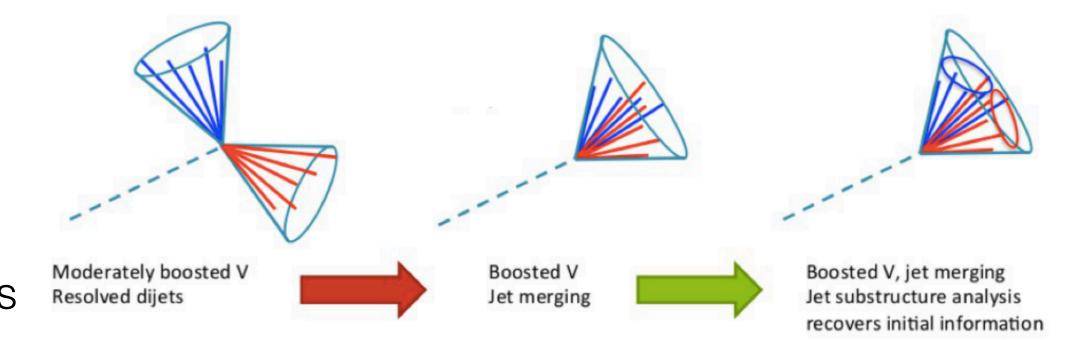




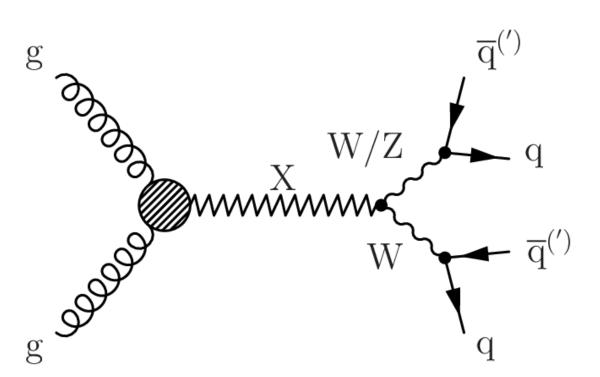
Advanced jet reconstruction - substructure

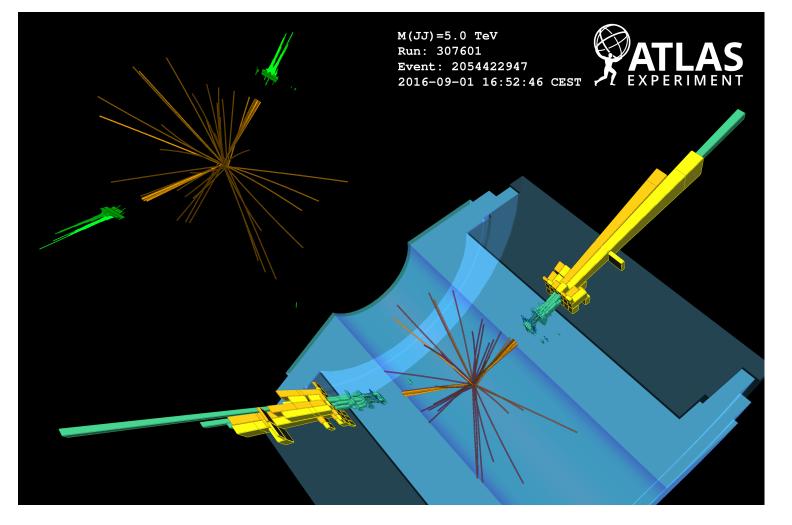
At very high momenta, jets (p.ex. from boson decays) can merge to form one big jet

 analyzing the jet substructure allows to reduce background



- example: high mass resonance search with diboson final state





Diboson candidate event

WIMP miracle

In the early universe dark matter and SM particles were in thermal equilibrium

- Constant production and annihilation of dark matter $\chi\chi \iff ff$
- Freeze out
 - Universe cools
 - WIMP mass too high for production $\chi\chi \to ff$
 - Universe expands, particles don't meet anymore
- Can obtain relic density from Boltzmann equation
- Compare with observation (CMB)

Corresponds to ~100 GeV particle interacting with the weak force

